Tracking aerosols and SO$_2$ clouds from the Raikoke eruption: 3D view from satellite observations

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Abstract. The 21 June 2019 eruption of the Raikoke volcano (Kuril Islands, Russia; 48°N, 153°E) produced significant amounts of volcanic aerosols (sulfate and ash) and sulfur dioxide (SO$_2$) gas that penetrated into the lower stratosphere. The dispersed SO$_2$ and sulfate aerosols in the stratosphere were still detectable by multiple satellite sensors for many months after the eruption. For this study of SO$_2$ and aerosol clouds we use data obtained from two of the Ozone Mapping and Profiler Suite sensors on the Suomi National Polar-orbiting Partnership satellite: total column SO$_2$ from the Nadir Mapper and aerosol extinction profiles from the Limb Profiler as well as other satellite data sets. We evaluated the limb viewing geometry effect (the “arch effect”) in the retrieval of the LP standard aerosol extinction product at 674 nm. It was shown that the amount of SO$_2$ decreases with a characteristic period of 8–18 d and the peak of stratospheric aerosol optical depth recorded at a wavelength of 674 nm lags the initial peak of SO$_2$ mass by 1.5 months. Using satellite observations and a trajectory model, we examined the dynamics of an unusual atmospheric feature that was observed, a stratospheric coherent circular cloud of SO$_2$ and aerosol from 18 July to 22 September 2019.

1 Introduction

An eruption of the Raikoke volcano (see Fig. 1; 0.55 km summit altitude; uninhabited island at 48.29°N, 153.4°E) occurred on 21 June 2019 at 18:00 UT. The eruption was so strong that a cloud of ash and volcanic gases was ejected to a height of 17–19 km (Gorkavyi et al., 2020), putting a significant part of the Raikoke volcanic plume into the stratosphere and above the heights of meteorological water and ice clouds and the tropopause, which is ∼10–12 km for such northern latitudes (see Fig. 1). The Microwave Limb Sounder (MLS) satellite data for 23–27 June 2019 show that the observed parts of the dispersed SO$_2$ cloud had heights between 11 and 18 km, with a peak concentration at 14 km. The maximum volcanic cloud top height increased by more than 6 km within 4 d after the eruption due to aerosol–radiation interaction (Muser et al., 2020).

The study of dispersed volcanic clouds and the transformation of their SO$_2$ into sulfuric acid aerosols is interesting for several reasons. Firstly, clouds of ash and sulfate aerosol are aviation hazards (Carn et al., 2009; Krotkov et al., 2014). Secondly, the lifetime of volcanic SO$_2$ depends on its injection height, SO$_2$ uptake on ash, and the concentrations of oxidants that vary with season and location (Guo et al., 2004; Krotkov et al., 2010; Beirle et al., 2014; Zhu et al., 2020). SO$_2$ from volcanic eruptions has a longer lifetime if it is injected into the stratosphere (Carn et al., 2016). SO$_2$ in the stratosphere is converted to sulfate aerosol, which...
interacts with the long-lived Junge layer – the naturally occurring background stratospheric aerosol layer. Furthermore, the dynamics of the Junge layer itself is not well understood. In addition, stratospheric aerosols are an important climate forcing factor, because aerosols modify both the short-wave and longwave radiation in the atmosphere and reaching the Earth’s surface (Toohey et al., 2019; von Savigny et al., 2020). Thus, each case of volcanic injection of large amounts of SO$_2$ into the stratosphere is of great interest to climate scientists (Robock 2000; Foster et al., 2008). Our ability to study volcanic SO$_2$ has evolved in recent decades along with satellite remote sensing technology (Krueger 1983; Krueger et al., 2000, 2008; Bovensmann et al., 1999; Guo et al., 2004; Carn et al., 2003, 2007, 2008, 2009, 2017; Clerbaux et al., 2009; Clarisse et al., 2010, 2013; Penning de Vries et al., 2014; Sandvik et al., 2019; Hedelt et al., 2019; Theys et al., 2019; Fisher et al., 2019). The spectral data obtained from satellites have made it possible to analyze in detail transient volcanic SO$_2$ clouds, such as of the eruptions of Pinatubo (Bluth et al., 1992; Guo et al., 2004; Fisher et al., 2019), El Chichón (Krueger 1983; Krueger et al., 2008), and Kasatochi (Bourassa et al., 2010; Prata et al., 2010; Krotkov et al., 2010; Clarisse et al., 2011).

The 2019 eruption of Raikoke was well observed by a number of satellite sensors including the Ozone Monitoring Instrument (OMI), the Tropospheric Monitoring Instrument (TROPOMI), the Ozone Mapping and Profiler Suite (OMPS) Nadir Mapper (NM) and Limb Profiler (LP), and the Cloud–Aerosol Lidar with Orthogonal Polarization (CALIOP). Studying the nadir passive data of OMPS NM and TROPOMI, our attention was drawn to an interesting phenomenon, the presence of compact long-lived stratospheric SO$_2$ clouds, hereafter referred to as coherent circular clouds or CCC, each about 300 km in diameter. The first was noted on 29 June 2019 centered at 51$^\circ$ N, 157$^\circ$ W (near Alaska), and the second was observed starting on 18 July 2019 near Kamchatka. The first of these moved towards the North Pole while the second moved south to $\sim$20–30$^\circ$ N, where it remained for more than 2 months, having made almost three complete revolutions around the Earth.

A dense cloud of SO$_2$ and aerosol that formed after the eruption of the Raikoke volcano spread over a week at latitudes above 40$^\circ$ N (see TROPOMI SO$_2$ map on 29 June 2019 in Fig. 2). Atmospheric currents stretched the cloud into long jets, twisted them with spirals, and even formed CCCs, as shown in Fig. 2 in the region near the coordinates 52$^\circ$ N, 156$^\circ$ W. To analyze in detail the evolution and dynamics of the Raikoke volcanic cloud, we examine (1) the amount and height of SO$_2$ and aerosol emitted during the eruption using satellite nadir and limb data; (2) the mutual evolution of SO$_2$ and aerosol for 100 d, starting from the moment of the eruption with limb data using an adjustment correcting for limb viewing geometry effects; and (3) the propagation of the CCCs with a case study in July 2019 using a trajectory model.

2 Data and methods

In this section, we detail the instruments and methods used to (1) examine SO$_2$ and aerosol distributions obtained from
We use SO\(_2\) total column data obtained from two satellite spectrometers: (1) TROPOMI (see Fig. 2; Theys et al., 2017, 2019) on the Sentinel-5 Precursor satellite and (2) OMPS Nadir Mapper (NM) on the Suomi National Polar-orbiting Partnership (SNPP) satellite (Li et al., 2013, 2015; Zhang et al., 2017). Both instruments observe the Earth’s backscattered radiance and solar irradiance at hyperspectral ultraviolet (UV) wavelengths with two-dimensional charge-coupled devices (CCDs), measuring in the spectral domain in one dimension and in the spatial domain (across a satellite track) in the other dimension. The satellite motion provides measurements along the satellite track. The spatial resolution of TROPOMI (3.5 km by 7 km) is much finer than that from OMPS NM measurements made in the nominal mode (50 km by 50 km), and the spectral resolution is higher for TROPOMI as well (∼0.5 nm as compared with ∼1 nm for OMPS NM). The satellite swaths are wide (2700 km for TROPOMI and 2800 km for OMPS NM), providing nearly global daily coverage. Column SO\(_2\) in Dobson units (1 DU = 2.69 × 10\(^{16}\) molecules cm\(^{-2}\)) is retrieved from hyperspectral solar backscatter at UV wavelengths (312–390 nm). TROPOMI uses a spectral fitting algorithm based on differential optical absorption spectroscopy (DOAS) (Theys et al., 2017, 2019). OMPS uses a spectral fitting approach with a principal component analysis (PCA) scheme (Li et al., 2013, 2017). For large volcanic SO\(_2\) signals like Raikoke, comparisons between TROPOMI and SNPP/OMPS so far (for several eruptions) show little bias, with the total SO\(_2\) mass estimates from the two normally agreeing to within 5%–10% (with the exception of the very early stages of large eruptions, where the density of SO\(_2\) and/or volcanic ash is too high to be fully accounted for in operational algorithms). For retrieval noise on a pixel-to-pixel basis, SNPP/OMPS SO\(_2\) (for stratospheric clouds) is less than 0.1 DU. TROPOMI’s noise on a pixel-by-pixel base is several times greater, but once TROPOMI pixels are averaged to OMPS footprints, the noise is reduced by ∼30%. We also used the Ozone Monitoring Instrument (OMI) on NASA’s Aura satellite as an additional data source.

### 2.2 Aerosol profile data

We use backscatter data from CALIOP (Fairlie et al., 2014), the OMPS Limb Profiler (LP) (Loughman et al., 2018; Chen et al., 2018), and the Stratospheric Aerosol and Gas Experiment instrument aboard the International Space Station (SAGE III/ISS). The OMPS LP V1.5 aerosol retrieval algorithm is described Sects. 2 and 3 of Chen et al. (2018). The height of an aerosol cloud can be estimated with both OMPS LP and CALIOP. We use the CALIOP lidar images for aerosol, specifically 532 nm total attenuated backscatter signal with a spatial resolution of 40 km along track and...
a vertical resolution of 120 m (altitude < 20 km) and 360 m (altitude > 20 km).

OMPS LP hyperspectral measurements capture aerosol data with a sampling of ~ 1 km in altitude. The actual cross-track field of view (FOV) of each OMPS LP slit is ~ 4 km (Dittman et al., 2002). OMPS LP views the atmosphere in a backward direction along the orbit track with three vertical slits, one (central) aligned with the orbit track and the other two (left and right) separated by ±4.25° horizontally. Here, we use aerosol data from version 1.5 (Chen et al., 2018).

Limb measurements of the atmosphere that view an altitude \( H \) at the tangent point also view altitudes above \( H \) in the foreground and background of the line of sight. However, the converse situation is also true: an object at a fixed altitude \( H \) will appear to be at a lower altitude \( h \) if it is located closer or farther than the tangent point. This is shown schematically in Fig. 3, where a single cloud is located at true altitude \( H \) and five successive LP measurement events (each event generates a new profile) are illustrated. For event C, the cloud location is at the tangent point, and the apparent altitude \( h \) equals the true altitude \( H \). For events A and B, the cloud position is closer to the LP instrument than the tangent point, and \( h \) is less than \( H \). For events D and E, the cloud position is farther from the LP instrument than the tangent point, and again \( h < H \).

Plotting successive altitude profiles of LP (e.g., extinction coefficient) as function of tangential latitude when a vertically limited feature (such as a cloud) is present gives the impression of an “arch” in the data. The arch effect is observed when the length of the visible part of the cloud is less than ~ 1100 km (at a cloud height of 25 km). Figure 4 shows cloud \( F_0G_0 \), 1 km thick and 226 km long, centered above tangent point \( T \). Due to the curvature of the globe, such a cloud has an observed thickness of 2 km (see Fig. 4). If we take a cloud 226 km long and with a real thickness of \( N \) km, then the observed cloud thickness will be \( N + 1 \) km. Thus, the real average height of a thin (1–2 km) cloud is underestimated by 0.5–1 km even under the most optimal observation conditions. Consider a cloud FG (see Fig. 4), the center of which is displaced from the tangent point by 273 km (or by \( \phi \sim 2.5° \)). The real height of the FG cloud is 24–25 km, but its observed height varies from 13 to 22 km. If we consider the \( F_0G \) cloud with a length of 499 km, then its real height above the Earth’s surface will be 24–25 km, and the observed height is 13–25 km. Let us take into account that the Limb Profiler assigns the latitude of the tangent point to any extended cloud. Therefore, a single cloud shown in Fig. 3 in five different observed positions, instead of one real geographic latitude, receives several observed latitudes, which creates an arch effect. Let the region \( F_0G_0 \) be a gap in a continuous cloud. Then this gap, together with the arch effect, will lead to a decrease in the maximum observed height of the cloud layer by 1 km (see Fig. 4).

Figure 5, showing an OMPS LP extinction profile track across the volcanic plume, illustrates the arch effect. The maximum values (yellow) appear to vary in altitude by 2–3 km due to the projection effect shown in Fig. 3. Note that this effect will be present for high, thin clouds such as cirrus clouds as well as for aerosol plumes. If we believe that these lower-altitude values do not represent a true aerosol signal, we need to apply a correction in order to accurately determine overall aerosol loading.

The relationship between true cloud height \( H \) and apparent cloud height \( h \) in Fig. 3 is given by a simple function (DeLand and Gorkavyi, 2021):

\[
\cos \varphi = \frac{R + h}{R + H},
\]

where \( R \) is the radius of the Earth and \( \varphi \) is the angular displacement. Table 1 gives some specific examples of the an-
Figure 5. (a) The extinction coefficient profiles (674 nm) from OMPS LP for the orbit 40 636 (central slit, same orbit as in Fig. 12, which shows the profiles for events 102, 105, and 108) for 31 August 2019. To accurately calculate the extinction coefficient it is necessary to take into account the arch effect, which is clearly visible (OMPS LP data). (b) (Left) An arch that appears during the limb observation of an aerosol cloud 2 km thick ($H = 23–25$ km) and $2^\circ$ long. Each frame corresponds to a satellite orbital shift by 1$^\circ$. Radiance from aerosol is proportional to the number of observed particles, taking into account the distance to them (the left part of the arch is slightly brighter than the right, because on frames 17–21 the cloud was closer to the satellite than on frames 23–27). The radiance units are arbitrary. (b) (Right) The red line is the profile of a cloud of 400 particles, which is observed on frame 22 (middle of the arch in Fig. 5a. The black line is the number of visible particles, summed over all frames, that is, over the entire arch.

Table 1. The observed height $h$ and angular displacement $\varphi$ for the cloud with true height $H$.

<table>
<thead>
<tr>
<th>$h \backslash H$</th>
<th>$H = 10$ km</th>
<th>15 km</th>
<th>20 km</th>
<th>25 km</th>
<th>30 km</th>
</tr>
</thead>
<tbody>
<tr>
<td>$h = 0$ km</td>
<td>$\varphi = 3.21^\circ$</td>
<td>3.93$^\circ$</td>
<td>4.53$^\circ$</td>
<td>5.07$^\circ$</td>
<td>5.55$^\circ$</td>
</tr>
<tr>
<td>5 km</td>
<td>2.27$^\circ$</td>
<td>3.21$^\circ$</td>
<td>3.93$^\circ$</td>
<td>4.53$^\circ$</td>
<td>5.07$^\circ$</td>
</tr>
<tr>
<td>10 km</td>
<td>–</td>
<td>2.27$^\circ$</td>
<td>3.21$^\circ$</td>
<td>3.92$^\circ$</td>
<td>4.53$^\circ$</td>
</tr>
<tr>
<td>15 km</td>
<td>–</td>
<td>–</td>
<td>2.27$^\circ$</td>
<td>3.20$^\circ$</td>
<td>3.92$^\circ$</td>
</tr>
<tr>
<td>20 km</td>
<td>–</td>
<td>–</td>
<td>–</td>
<td>2.27$^\circ$</td>
<td>3.20$^\circ$</td>
</tr>
<tr>
<td>25 km</td>
<td>–</td>
<td>–</td>
<td>–</td>
<td>–</td>
<td>2.26$^\circ$</td>
</tr>
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gular displacement predicted for specific combinations of $H$ and $h$. This displacement is approximately equal to a shift in latitude for a sun-synchronous satellite such as SNPP (inclination = 98.9$^\circ$) outside the polar regions. From this table, we see that for a true cloud at $H = 25$ km, the apparent cloud signal at $h = 15$ km would be displaced by $\sim 3.2^\circ$ both earlier and later along the orbit. This total separation of $\sim 6.4^\circ$ is consistent with the arch width at 15 km shown in Fig. 4. Note also that because of differences in overall path length, we expect events A and B in Fig. 3 (left branch of arch in Fig. 4) to have a stronger signal than the corresponding events D and E at the same apparent altitude. We can therefore use Eq. (1) to calculate and apply a correction for

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determining the magnitude and position of an aerosol cloud (we believe that all parts of the arch below the real height are artifacts, so the value of the extinction coefficient for them should be equal to zero).

The arch in Fig. 5a is formed by several measurements of one CCC, which has a true height of 23–25 km (see the analysis of this CCC in Sect. 3.2). Parts of the arch (between events 102 and 108), indicated with arrows in Fig. 5a below 23 km, are artifacts of multiple registrations of the same object and should be removed when calculating the total amount of aerosol at given latitudes or heights. The low stratospheric aerosol layer at an altitude of 18–20 km decreasing with latitude to 15–16 km (blue color) is the background layer. It merges in frames 101–109 with an apparent volcanic cloud arch. Note that the CCC has been the southernmost part of the Raikoke plume since late July 2019 (see Sect. 3.3). The tropical cloud at latitude ~ 8° N extends up to 16 km altitude.

The reality of the discussed arch effect is confirmed by simple modeling. Arch model in Fig. 5b was obtained by direct modeling of a cloud of 400 particles located at the nodes of a uniform grid (20 particles are distributed at a length of 2°, and 20 rows of particles are uniformly distributed along the radius in the range of 23–25 km). The model does not use radiative transfer models. In Fig. 5b (left) only the distance between the particles and the satellite is taken into account, while Fig. 5b (right) shows the simply visible (for the limb sensor) number of particles at a given altitude $h$ (in 1 km step). The red line corresponds to a one-time observation of a cloud located at the tangent point at an altitude of $h = 22–25$ km (an increase in the apparent thickness of the cloud by 1 km is associated with the curvature of the Earth, which is why the cloud itself turns out to be curved – see Figs. 3 and 4). That is, it is the most realistic observation of the cloud at its optimal location. The black line shows the sum of the cloud particles observed at different times. As a result of this summation, the number of particles visible at a given observed height $h$ (which differs from the real constant cloud height $H = 23–25$ km) turns out to be overestimated. Therefore, when analyzing the picture in Fig. 5, we must remember that it is composed of frames received at different times, so the one cloud will be registered many times. The same effect of multiple registrations will be observed for clouds of any complexity and configuration, including a uniform aerosol layer, because any cloud can be divided into a large number of elementary pieces, similar to a simple compact cloud considered in Fig. 5b. Obviously, the arch effect for a spherically uniform aerosol layer should be fully compensated for by the 1D retrieval model. But the further the real system is from the spherical symmetry, the more difficult it will be to take into account the arch effect.

The considered model of the arch effect does not depend on the specific model of radiation transfer and the methods of retrieval of the spatial distribution of aerosol. Therefore, for specific limb sensors (OMPS-LP, SAGE III, OSIRIS), it is necessary to evaluate how accurately the available retrieval packages handle compact clouds. This is especially true for 1D retrieval methods, which assume spherical symmetry of the atmosphere and which are used to obtain aerosol extinction in the OMPS/LP. To estimate the possible retrieval uncertainty due to the arch effect, we apply a simple compensation method for one specific case of the Raikoke aerosol cloud (see next section). This compensation method assumes that the considered aerosol clouds form compact clusters or a highly heterogeneous aerosol layer and that the arch effect was not taken into account in retrieval. Thus, this example should be regarded as a maximal estimation for the arch effect. Where this assumption is not valid, our correction will be overestimated, as, for example, happened with the correction of the background aerosol value (see Fig. 6), which was observed before the Raikoke eruption.

The principle of our posterior algorithm for estimating the arch effect is as follows: the data in the uppermost pixel of the arch are considered true. For a given pixel, artifact extinctions are calculated for pixels in the side branches of the arch. These artifact extinctions are subtracted from the initial arch extinction. The procedure is repeated for the second-highest pixel in the already modified arch image. This continues until the side branches of the arch completely disappear. The top pixels with the corrected extinction are summed into the cloud with the corrected data. We consider this procedure of correction only as an estimate, which shows the possible significance of the arch effect.

The arch effect is a specific example of the effects of inhomogeneity along the line of sight that is an issue for all types of limb sounders. One way to account for such effects is to use a two-dimensional (2D) radiative transfer model (RTM) that is able to account for such effects along with multiple
observations in a tomographic retrieval (e.g., Livesey et al., 2006; Zawada et al., 2018; Loughman et al., 2018). Instead, we have developed an a posteriori adjustment method that is effective in correcting for the arch-type effects that appear with isolated features within the lines of sight for a series of observations.

The arch effect characteristic of the limb observations should be taken into account when calculating the optical thickness of aerosol clouds, determining a characteristic cloud height, or using calibration heights of ~45 km. Calibration heights should be free of aerosol, and they are indeed free of aerosol but can be contaminated with polar stratospheric and mesospheric clouds (PMCs) and other stratospheric aerosols. For example, PMCs located at an altitude of 80–85 km, due to the arch effect, can be projected onto the calibration heights. Such contamination of the calibration heights distorts the whole picture of aerosol distribution with height.

SAGE III/ISS makes direct measurements of aerosol extinction through the attenuation of the solar beam during sunrise and sunset. The vertical resolution of these measurements is about 1/2 km. Because they are occultation measurements, SAGE observations are less frequent than OMPS-LP.

Using the color ratio (extinction at 512 nm to extinction at 1022 nm) along with the extinction at 1022 nm, clouds can be removed from the SAGE observations to produce average aerosol extinctions (Thomason and Vernier, 2013).

### 2.3 The NASA ftraj trajectory model

The “ftraj” trajectory model from NASA’s Goddard Space Flight Center Atmospheric Chemistry and Dynamics Laboratory uses a fourth-order Runge–Kutta integration scheme to track parcels isentropically, with optional diabatic adjustments (Schoeberl and Sparling, 1995). The model is driven with winds at 1/4° horizontal resolution and spaced every 6 h from the Goddard Earth Observing System (GEOS) forward-processing system produced by the NASA Global Modeling and Assimilation Office (GMAO). Each model run was initialized with 3000 parcels distributed at random horizontal and vertical positions within a cylinder of radius 150 km centered about the location of the observed cloud and stretching between the lowest and highest potential temperature (Θ) values of the observed cloud. Those Θ values were derived by interpolating the GEOS fields to the parcels’ initial longitude, latitude, altitude, and time. Both the starting and ending cloud observations extended through a range of altitudes and hence Θ values.

### 3 Results

The sensitivity of the satellite data (OMPS LP, CALIOP) we examined is such that the aerosol cloud from the Raikoke volcano was observable for many months following the eruption. Thus, a synergistic study of Raikoke volcanic emissions using various satellite-based instruments provides a good opportunity to study the dynamics of the volcanic cloud long-term dispersion. For quantitative analysis, we will use 2D SO$_2$ vertical column density (VCD) maps obtained by OMPS NM, as well as vertical profiles of aerosol extinction that are obtained from the OMPS LP.

#### 3.1 SO$_2$ and aerosol evolution

The retrieved Raikoke SO$_2$ mass increased for 2 d after the eruption and then exponentially decayed with an e-folding timescale of 8–18 d. The amount of stratospheric SO$_2$ decreases due to photochemical conversion to sulfuric acid through gas-phase reaction with the hydroxyl radical, OH.
Sulfuric acid nucleates new particles and condenses onto pre-existing particles to form long-lived stratospheric sulfate aerosol. Figure 6 shows the change in the total SO$_2$ mass (kt) retrieved by the OMPS NM and the altitude-averaged (13–18 km) aerosol extinction coefficient retrieved with the OMPS LP in the latitude range 45–65° N after the Raikoke eruption. The aerosol extinction data in Fig. 6 are shown before and after removing the arch effect. These estimates show the potential significance of the arch effect. As we discussed earlier, the arch effect estimation assumes high cloud heterogeneity and should be considered the maximum estimation, as indicated by the arrow bars on Fig. 6. For more accurate calculations, the arch effect should be investigated using a 2D RTM.

Figure 6 shows that retrieved SO$_2$ mass increases for 2 d after the eruption during rapid dispersion of the fresh opaque volcanic cloud. Following this initial dispersion, the SO$_2$ mass decreases exponentially due to chemical conversion to sulfate. Aerosol extinction has opposite behavior; it increases by approximately a factor of 4 as the concentration of SO$_2$ decreases, reaching a maximum ∼ 50 d after the eruption. After that, the aerosol extinction starts to decrease due to gravitational sedimentation and other processes, but very slowly (from OMPS LP and NM data). Adjustments for the arch effect do not change the overall temporal pattern.

As can be seen from Fig. 7, the total SO$_2$ mass decreases exponentially with a variable (8–18 d) timescale. Krotkov et al. (2010) derived a volcanic SO$_2$ mass decay of ∼ 9 d following the 8 August 2008 Kasatochi eruption (30–90° N) using the Aura OMI. The overall apparent $e$-folding time for Raikoke is ∼ 8–10 d (Fig. 7, solid black lines), in good agreement with the Kasatochi $e$-folding time. However, we see two different regimes in Raikoke SO$_2$ clouds with significantly different estimated $e$-folding time, as discussed below.

Figure 7a shows that during the first 20 d after the eruption, the SO$_2$ mass decreases with a longer apparent $e$-folding lifetime of ∼ 19 d, compared to the later period (days 21 to 60), when the decrease in SO$_2$ mass accelerates with the apparent $e$-folding lifetime ∼ 8 d. The initial estimate of the $e$-folding lifetime of ∼ 19 d, derived from the earlier stage after the eruption, probably represents the true timescale for chemical conversion of SO$_2$ into sulfate aerosol on altitudes > 20 km.
Figure 9. (a) Daily location of the maximum altitude reached by the Raikoke aerosol plume as detected by OMPS LP, colored by the day number since 1 June and plotted every other day until day 150. (b) Daily zonal mean aerosol extinction profiles at 675 nm (km$^{-1}$) between 20° S–90° N (which is close to hemispheric coverage) measured by OMPS LP from June 2019 to June 2020 and smoothed spatially using a five-point boxcar averaging. Only profiles measured above tropopause +1 km are used. The white line is the aerosol plume maximum altitude in kilometers (km). (b) (Top) Stratospheric aerosol optical depth ($\times 10^3$, sAOD) at 675 nm, derived by integrating aerosol extinction profiles above the tropopause to 35 km.

(Carn et al., 2016). Following the drop of the SO$_2$ column amounts, the SO$_2$ detection limit of the OMPS instrument becomes a more important factor, leading to faster apparent decay of the total observed mass of SO$_2$. As the plume spreads out to larger areas, and more and more pixels with SO$_2$ fall below the detection limit of the OMPS NM sensor (OMPS in the stratosphere can typically detect 0.2–0.3 DU of SO$_2$), the apparent decay rate becomes larger, likely reflecting the combined effects of the chemical SO$_2$ loss and the diminished OMPS sensitivity to weaker SO$_2$ signals. Figure 7b shows a similar pattern for change in SO$_2$ mass within the CCC. For the period between 28 and 45 d after the eruption, the $e$-folding lifetime is $\sim 19$ d. This suggests that the time after the eruption does not play a big role here – only the compactness of the CCC with a high level of SO$_2$ is essential. As in the case of Fig. 7a, the second phase of the evolution of the CCC is characterized by a faster apparent (the $e$-folding lifetime $\sim 8$ d) rate of decrease in the recorded SO$_2$.

3.2 Stratospheric aerosol from limb observations

Figure 8 shows the time–latitude evolution of the zonal average aerosol extinction at different heights (without adjustment for the arch effect). Most of the aerosol is transported poleward at 14.5 km (Fig. 8a), and the effects of the eruption lasted for almost a year. In addition, Figs. 8b and c also show the aerosol transport to subtropics and tropics at higher altitudes. Increased aerosol loading in the lower stratosphere can also be attributed to two pyrocumulonimbus (pyroCb) events that took place before and after the eruption, Alberta fires (18 June) and Siberian fires (2 July) (Kloss et al., 2021). OMPS LP detected both plumes in the stratosphere at 12–13 km, although it became difficult to separate them from the Raikoke plume once it spread around the Northern Hemisphere. The zonal mean aerosol extinction profile between 20 and 90° N shows the vertical transport of the Raikoke plume to higher altitudes and its persistence in the lower stratosphere (Fig. 9). Figure 9a shows that the maximum altitudes of the plume are around 25 km, when the plume penetrates the tropics. The plume altitude is derived using the OMPS LP cloud al-
The ratio of the zonal mean, daily average aerosol extinction for the indicated dates after the Raikoke eruption to the averaged aerosol extinction immediately prior to the eruption shows the evolution of the volcanic aerosol in the background aerosol layer (13–20 km) and the occurrence of the CCC near 25 km from OMPS LP data.

Figure 10 shows the altitude–latitude profile of a daily zonal mean ratio of the aerosol extinction coefficient for 1 and 21 July and 1 and 30 September 2019 to a similar daily zonal mean for a quiet period (16–20 June 2019) before the eruption. The means were computed without taking into account the arch effect (see DeLand and Gorkavyi, 2021, and Sect. 2.2). Therefore, the lower parts of the layers shown may have an overestimated extinction coefficient. In addition to increasing the density of the background aerosol layer (the contribution of seasonal changes is possible here), an interesting CCC appears at an altitude of about 25 km that will be discussed in Sect. 3.3.

The SAGE 1022 nm extinctions for the month of September 2019 are shown in Fig. 11 along with the profile lati-
Figure 12. The aerosol formed as a result of the Raikoke eruption is clearly visible in the extinction coefficient profiles obtained by the OMPS LP aerosol extinction vertical profile months after the eruption.

Figure 13 shows satellite observations of the CCC from 18 July to 22 September 2019. During this period, the height of the CCC has increased significantly. There are more points by the end of the CCC observation period, because the cloud spreads out and is more often observed by CALIOP.

Figure 17 shows the results of backward trajectory modeling of the CCC for the 1-month period with August 18 through 19 September 2019. The basic idea is to use the trajectories to see if this long-lived CCC can be correctly transported using assimilated meteorological observations.

We have two hypotheses to evaluate. The first hypothesis ("the great red spot" hypothesis) is that the observed CCC is contained within a gyre (i.e., a whirlpool) that existed from around the time of the eruption on 21 June to 1 September or even later, advected around the globe by the large-scale circulation. The second hypothesis (the "dead fish" hypothesis) is that the CCC was simply advected around by the large-scale
winds without being sheared apart because the environment of the summer stratosphere is fairly quiescent.

We inspected the meteorological fields from NASA’s Global Modeling and Assimilation Office (GMAO) Modern-Era Retrospective analysis for Research and Applications, version 2 (MERRA-2) reanalysis (Gelaro et al., 2017) and from the Global Data Assimilation System (GDAS) by the National Centers for Environmental Prediction (NCEP) at the location and altitude of the observed volcanic CCC on 21 and 24 July; the winds show no discernible gyre. This is to be expected, as the satellite data that go into these assimilation products are unlikely to pick up features on sufficiently small scales that they would show up clearly and prominently in the winds. Therefore, it is not possible to test the great red spot hypothesis using a trajectory model with the wind products available to us.

We can examine the dead fish hypothesis. If a trajectory model run shows a cloud of parcels maintaining its integrity over a long period of time, then this hypothesis becomes plausible. There are difficulties, however. First, the winds used to drive the trajectory model have spatial and temporal resolutions that are coarse enough that trajectories cannot be modeled with perfect accuracy. One can expect that a cloud of parcels will lose many of its parcels along the way. It would be unrealistic to expect to duplicate the volcanic cloud’s integrity and position as a function of time. Instead, the test is whether the simulated cloud of parcels can maintain its existence over the course of a month or more during this time period. Second, a more serious problem is that of altitude. We can expect a volcanic CCC to experience self-lofting, and the CCC indeed is observed to increase in altitude with time. However, the winds vary with altitude, and thus the parcel trajectories will depend critically on getting their altitude changes right.

A carbonaceous aerosol plume associated with wild fires in British Columbia in August 2017 reached the stratosphere a few days following the initial injection and resulting from self-lofting triggered by solar heating (Torres et al., 2020; see similar effect for Australian fires – Khaykin et al., 2020). The case of a volcanic CCC, which we are considering, is interesting in that the effect of an increase in height from 19 to 26 km is observed for a cloud presumably consist-
Figure 14. OMPS SO$_2$ STL columns plotted using a linear scale for 30 July 2019. The data show the westward drift of the CCC with easterly stratospheric winds at $\sim 14$–$15^\circ$ d$^{-1}$, or $\sim 1400$ km d$^{-1}$, $\sim 60$ km h$^{-1}$.

Figure 15. OMPS SO$_2$ STL columns plotted using a linear scale for 4 August 2019.
The CALIOP total attenuated backscatter data (km$^{-1}$ sr$^{-1}$) for 18 August and 1 September 2019, where the CCC is visible at the indicated latitudes and heights.

The observed path of the CCC as inferred from OMI/OMPS/TROPOMI and CALIOP observations. The individual observations are shown by dots color-coded by month.

The thermodynamic processes within the cloud are not well known, so rather than attempting to model the cloud’s self-lofting and the overall vertical advection, we impose an approximation to the cloud’s vertical motion. We determined also confirmed by independent satellite data from the Atmospheric Chemistry Experiment–Fourier transform spectrometer (ACE-FTS) (Cameron et al., 2021).

In this case, the main increase in the height of the CCC occurred during its movement at approximately the same latitude. Early on (19 July) it is at 80 hPa (450 K), but by 1 September it had risen to 30 hPa (590 K). This is a diabatic self-lofting heating rate of about 3 K d$^{-1}$. Note that this CCC, increasing its height, simultaneously contained a detectable amount of SO$_2$, which is also confirmed by independent satellite data from the Atmospheric Chemistry Experiment–Fourier transform spectrometer (ACE-FTS) (Cameron et al., 2021).
Figure 18. Backward trajectory modeling. Panel (a) shows several discrete time snapshots from the Goddard “ftraj” back-trajectory modeling for 3000 parcels within a 150 km radius of the observed CCC position initialized on 18 August 2019. (b) The trajectory model run was initialized with the observed cloud position on 19 September 2019 and then run backwards to 18 August 2019. Panel (c) shows the same model results as for panel (a) but using a different visualization. Here the parcels are binned and counted in longitude and time, regardless of altitude or latitude. (d) Similar to panel (c), except that the model run is initialized on 19 September 2019 and run backwards to 18 August 2019.

Figure 18a superimposes several discrete time snapshots from the trajectory model output, starting with the observed cloud position on 18 August 2019. The modeled positions are generally close to the observed, except that the model fails to pick up the effects of the synoptic weather feature beginning around 23 July. Figure 18b shows the output from another model run, initialized with the observed CCC position on 19 September 2019 and then run backwards to 18 August 2019. Figure 18c and d show the same model results but using a Hovmöller-type plot.

The CCC appears not to be a whirlpool but rather a CCC with easterly flow and drifting (dead fish model). Neither the turbulent mixing nor gravity waves penetrate well into the summer stratosphere to mix features, so they tend to move intact if injected at the right time.

4 Conclusion

We studied the process of converting SO$_2$ to aerosol for the 21 June 2019 eruption of the Raikoke volcano (Kuril Islands, Russia; 48°N, 153°E). The peak sulfate aerosol extinction is 1.5 months behind the time of the release of SO$_2$ in the Raikoke eruption. We see two different regimes in Raikoke SO$_2$ clouds with significantly different $e$-folding times. Initially, the SO$_2$ mass decreases with a longer apparent $e$-
folding lifetime of $\sim 19$ d, compared to the later period when the decrease in SO$_2$ mass accelerates with the apparent $e$-folding lifetime $\sim 8$ d. The $e$-folding lifetime of $\sim 19$ d, derived from the earlier stage after the eruption, probably reflects the true conversion of SO$_2$ into sulfate aerosol. After that, the detection limit of the OMPS instrument becomes a more important factor, leading to faster apparent decay of the total observed SO$_2$.

We evaluated the influence of the arch effect on the calculation of the extinction coefficient of a finite volcanic cloud using a homogeneous spherical shell atmosphere assumption adopted for processing the data of the OMPS LimbProfiler. We have shown that this effect is significant; therefore, it should be investigated by more accurately accounting for inhomogeneous atmospheric composition along the line of sight within the framework of a 2D model of radiative transfer.

We examined the unusual coherent circular cloud (CCC) of SO$_2$ and aerosol, which was observed for more than 2 months (from 18 July to 22 September). The CCC was embedded in the summer easterly flow in the stratosphere and demonstrated an unprecedented structural cohesion along with a diabatic self-lofting heating rate of about 3 K d$^{-1}$. For 2 months, the CCC circled the globe almost three times at the latitude of $30^\circ$ N (during September, the CCC shifted south to $15$–$20^\circ$ N) and increased its height from 19 to 26 km.

Data availability. The SO$_2$ data for OMPS, OMI, data, and TROPOMI data are available at https://so2.gsfc.nasa.gov/ (Krotkov, 2021) and the Goddard Earth Sciences Data and Information Services Center (GES DISC) at https://doi.org/10.5067/A9O02ZH0J94R (Yang, 2017). CALIOP data are available at https://www-calipso.larc.nasa.gov/products/ (Vaughan et al., 2004). The last version of the OMPS-NPP LP L2 Aerosol Extinction Vertical Profile swath multi-wavelength daily slit collection data are accessible from the Goddard Earth Sciences Data and Information Services Center (GES DISC) at https://snpp-omps.gesdisc.eosdis.nasa.gov/data/SNPP_OMPS_Level2/OMPS_NPP_LP_L2_AER_DAILY.2/doc/README.OMPS_NPP_LP_L2_AER_DAILY.pdf (Johnson et al., 2020) and https://doi.org/10.5067/CX2B9NW6FI27 (Taha, 2020). The SAGE III/ISS data were obtained from the NASA Langley Research Center Atmospheric Science Data Center at https://asdc.larc.nasa.gov/project/SAGEIII-ISS (see Alfred et al., 2007).

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Competing interests. The contact author has declared that neither they nor their co-authors have any competing interests.

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