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Lidar-Radiometer Inversion Code (LIRIC) for the retrieval of vertical aerosol properties from combined lidar/radiometer data: development and distribution in EARLINET

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LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



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LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Abstract

This paper presents a detailed description of LIRIC (Lidar-Radiometer Inversion Code) algorithm for simultaneous processing of coincident lidar and radiometric (sun photometric) observations for the retrieval of the aerosol concentration vertical profiles.

As the lidar/radiometric input data we use measurements from European Aerosol Research Lidar Network (EARLINET) lidars and collocated sun-photometers of Aerosol Robotic Network (AERONET). The LIRIC data processing provides sequential inversion of the combined lidar and radiometric data by the estimations of column-integrated aerosol parameters from radiometric measurements followed by the retrieval of height-dependent concentrations of fine and coarse aerosols from lidar signals using integrated column characteristics of aerosol layer as a priori constraints. The use of polarized lidar observations allows us to discriminate between spherical and non-spherical particles of the coarse aerosol mode.

The LIRIC software package was implemented and tested at a number of EARLINET stations. Intercomparison of the LIRIC-based aerosol retrievals was performed for the observations by seven EARLINET lidars in Leipzig, Germany on 25 May 2009. We found close agreement between the aerosol parameters derived from different lidars that supports high robustness of the LIRIC algorithm. The sensitivity of the retrieval results to the possible reduction of the available observation data is also discussed.

1 Introduction

The aerosol impact on the radiation balance of the atmosphere is an important climate forming factor. In addition, aerosol particles are among the unhealthiest pollutants of the air. It is dramatized by rapid propagation of pollutants in the atmosphere that expands local ecocatastrophes to a global scale. Therefore, the monitoring of the aerosol evolution and transport in the atmosphere is an obligatory prerequisite for predicting climatic and ecological changes.

LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Sun-radiometer and lidar networks contribute to aerosol remote sensing. The global Aerosol Robotic Network (AERONET) of ground-based sun/sky-scanning radiometers (e.g. Holben et al., 1998) provides reliable data on columnar aerosol properties from more than 200 world-wide distributed sites. The results of AERONET observations are the aerosol optical thickness (AOT) obtained from direct sun observations and additional microphysical and optical properties of aerosol particles (single scattering albedo, volume distribution of aerosol particles, complex refractive index, fraction of spherical particles, etc.) derived by the inversion of direct and scattered radiation measurements (Dubovik and King, 2000; Dubovik et al., 2002, 2004). The regional radiometer network SKYNET was established in the South Eastern Asian regions (Takamura et al., 2004) and it employs its own equipment and processing procedure (Hashimoto et al., 2012).

The lidar measurements are basically used to provide information on the vertical variability of the aerosol characteristics. Currently, lidar networks, such as the European Aerosol Research Lidar Network (EARLINET) (Bösenberg et al., 2000; Pappalardo et al., 2014), the micro-pulse lidars network (MPL-Net) (Welton et al., 2002), the Asian Dust Network (AD-Net) (Murayama et al., 2001), the lidar network in Former SU countries CIS-LiNet (Chaikovsky et al., 2005), the northeast American CREST Lidar Network (CLN) (Hoff et al., 2009), and the Latin-American Lidar Network LALINET (Antuña et al., 2012), monitor aerosol vertical distributions in the atmosphere over the vast regions of the Earth. The Global Atmosphere Watch (GAW) Aerosol Lidar Observation Network (GALION), also known as the “network of networks” (e.g. Bösenberg and Hoff, 2007), was established under the aegis of GAW to coordinate lidar activity all over the world. The outcome of the lidar observations are presented in the lidar network databases as vertical profiles of aerosol backscatter and extinction coefficients.

Aerosol columnar properties from AERONET and pointed above vertical profiles of aerosol parameters from lidar networks are complementary pieces of information characterizing aerosol properties. Nowadays, lidars and sun/sky-scanning radiometers are among the basic tools in comprehensive experiments aimed at studying the transfor-

LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



mation and transport of smoke (e.g. Lund Myhre et al., 2007; McKendry et al., 2011; Colarco et al., 2004), dust (e.g. Ansmann et al., 2009; McKendry et al., 2007; Müller et al., 2003; Papayannis et al., 2008) and volcanic ash (e.g. Ansmann et al., 2010, 2011, 2012; Papayannis et al., 2012; Gasteiger et al., 2010). A number of SKYNET sites (Takamura et al., 2004) and most of the EARLINET stations are equipped with lidar and radiometer instruments. Further enhancement of the aerosol characterization is expected from the synergy of co-located radiometer and lidar observations. Namely, the coordination of measurement procedures of the two systems and the derivation of aerosol parameters from combined measurements results in advanced characterization of the aerosol layer with a superior performance compared to the aerosol information that would have been obtained from independent processing of lidar and radiometer data.

The idea of combined lidar and radiometer sounding (LRS) for retrieving vertical distributions of aerosol characteristics was first proposed by Chaikovsky et al. (2002), and it gave rise to the development of the lidar-radiometer synergetic algorithm, later named LIRIC (Lidar-Radiometer Inversion Code). LIRIC is based on processing co-located lidar and radiometer measurements by using a two-step *sequential inversion*. First, the radiometer data was processed in tune to the standard AERONET inversion algorithm. Then, first-step results were used as a priori constrains on aerosol properties for lidar data processing.

First application of LIRIC technique to the actual data processing was presented by Chaikovsky et al. (2004a). In that study, the technique was adapted to the EARLINET-AERONET stations in Minsk (Belarus) and Belsk (Poland) (e.g. Chaikovsky et al., 2004b, 2010a). Results of the LRS observations were of interest for studying of long range aerosol transport in the East European region (Kabashnikov et al., 2010; Chaikovsky et al., 2010b; Papayannis et al., 2014).

Another algorithm for data processing in combined lidar-and-radiometer experiments exploits the decomposition of the AERONET column-integrated aerosol size distribu-

tion into log-normal modes and selection of some of these modes for characterization of aerosol layers using measured lidar data (Cuesta et al., 2008).

Besides, the single-wavelength POLIPHON technique was alternatively developed (e.g. Tesche et al., 2009; Ansmann et al., 2012). This technique retrieves particle volume concentration profiles of aerosol separately for fine and coarse fractions. The algorithm relies on the measured profiles of the particle linear depolarization ratio and lidar ratio and it does not require the assumption of a specific particle shape. Columnar concentrations of aerosol modes retrieved by AERONET are used in POLIPHON as additional input data. The algorithm POLIPHON is designed for the data processing in lidar sounding of the aerosol layers with coarse non-spherical particles (dust, volcano ash).

In the recent years, the LRS technique has been implemented within the advanced research network ACTRIS in the frame of EU 7th Framework Programme project. To date, a number of joint EARLINET/AERONET stations have implemented regular atmospheric observations using LIRIC for processing combined sun-radiometer and lidar-measured data (e.g. Chaikovsky et al., 2012; Papayannis et al., 2014; Tsekeri et al., 2013). The aerosol model and mathematical basis of the LIRIC algorithm became the prerequisite for further development of algorithms for parallel processing of combined lidar-radiometer measurements (Lopatin et al., 2013) and the results of ground/satellite bond experiments (Dubovik et al., 2014). At the same time, a comprehensive description of the LIRIC algorithm has not been yet documented in detail.

This paper describes the basic physical and mathematical aspects of LIRIC algorithm with all necessary equations, thus filling up this gap. The appendices contain the details of the inversion scheme and can be useful for advanced users to modify and improve this code.

LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



2 The algorithm concept and structure

The aerosol retrievals from combined lidar and radiometer measurements belong to a class of “ill-posed” inverse problems that, in particular, is characterized by non-unique and highly unstable solutions arising even under small measurement or simulation errors. In practice, solution of the “ill-posed” problems requires to introduce a priori information (e.g. Turchin et al., 1971; Tikhonov et al., 1977; Twomey, 1977; Tarantola, 1987; Rodgers, 2000). LIRIC algorithm was designed on the basis of multi-term LSM (Least Square Method) (Dubovik, 2004). This method was implemented in AERONET data processing (Dubovik and King, 2000) and then it was refined in the retrieval algorithms for the data processing of the combined optical measurements (e.g. Dubovik, 2004; Dubovik et al., 2011, 2014; Lopatin et al., 2013).

The inversion algorithm LIRIC can be divided into three key procedures (e.g. Tarantola, 1987): (i) parameterization of the object under study (i.e., development of the aerosol layer model), (ii) forward modeling, i.e. derivation of the equations that relate observed signals with specified parameters of the aerosol model, and (iii) inverse modeling or retrieval of the target parameters of the aerosol model that minimize discrepancies between the measured and the calculated input signals.

2.1 Combined lidar/radiometer experiment and aerosol model

The lidar/radiometer input data assumed to come from measurements of EARLINET lidars (e.g. Matthias et al., 2004; Freudenthaler et al., 2010) and spectral-scanning sun-radiometers of AERONET (Holben et al., 1998). The majority of EARLINET lidar stations provides day-time measurements of elastic backscatter signals at three wavelengths (355, 532 and 1064 nm) and two cross/parallel polarization components of the signal at a single wavelength. Additional information on aerosol parameters is expected to come from day-time Raman lidar measurements.

Radiometric data includes results of direct-Sun and almucantar (scanning) measurements (Holben et al., 1998; Dubovik and King, 2000). Direct-Sun measurements are

LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

[Title Page](#)

[Abstract](#)

[Introduction](#)

[Conclusions](#)

[References](#)

[Tables](#)

[Figures](#)

[◀](#)

[▶](#)

[◀](#)

[▶](#)

[Back](#)

[Close](#)

[Full Screen / Esc](#)

[Printer-friendly Version](#)

[Interactive Discussion](#)



LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



carried out in 15 min intervals. Almost clear-sky measurements are required to obtain almucantar data and about 2–6 successful measurements are made during the daytime under favorable meteorological conditions at EARLINET/AERONET stations. Under these circumstances time synchronization of lidar and radiometric observations usually means nearly simultaneous measurements within the same 1 h interval.

These radiometric measurements enable the retrieval of the aerosol properties over the entire atmospheric column. Thus, except for volcanic events, the maximum lidar sounding height, h_{\max} , can be limited to the tropopause level because the stratospheric aerosol layer does not significantly contribute to columnar aerosol optical parameters. In contrast, aerosols in the lower troposphere are key contributors to the observed columnar characteristics. Consequently, it is desirable to perform the lidar sounding from the lowest possible altitude. Likewise, the contribution of the bottom layer (which is not observed by lidar) to the columnar optical parameters must be small enough to be modeled by a homogeneous layer with the same aerosol parameters as at the lowest level of lidar sounding. In practice, the lower sounding limit for most of the lidar measurements in EARLINET is about 200 m or more that can be too high especially for low boundary layers in winter seasons. It should be decreased in winter to compensate reduction of the boundary layer height. Therefore, lidar measurements in the lower layer have to be carried out by a co-operative receiving system with smaller objective and larger field of view or by sounding the atmosphere along a slant trajectory.

The choice of the optical aerosol model is a key step of the retrieval algorithm. The optical model should be constructed following the principle of parsimony or “Occam’s razor”: the number of aerosol parameters has to be minimal but complete in order to provide unbiased retrieval from available measurements.

In this work, we use the AERONET model approach to characterize the aerosol layer of the atmosphere (Dubovik and King, 2000): aerosols are modeled by several modes with a certain aerosol particle size distribution, wherein each mode is a mixture of homogeneous spherical particles and randomly-oriented spheroids (Dubovik et al., 2002, 2006). The distribution of the spheroid aspect ratio is fixed. The number of

aerosol modes, K depends on specification of the lidar data. If we use only total (scalar) backscatter lidar measurements, the aerosol model includes fine and coarse modes ($K = 2$). There is boundary size between fine and coarse fractions in the algorithm, which is determined as the value in 0.194–0.576 μm range that corresponds to a minimum of the column particle volume size distribution, $dV(r)/d\ln r$. If measurements of cross and parallel co-polarized components are available, spherical and non-spherical particles of the coarse mode are considered as two different fractions ($K = 3$).

Thus, two sets of parameters characterize the aerosol layer:

- i. A number of columnar aerosol parameters retrieved from radiometer measurements (Dubovik et al., 2000a, 2002, 2006). This set of parameters is formed by:
 - (1) the total content of each aerosol mode, (i.e. columnar volume concentrations),

$$\hat{C}_k^V = \int_{r_{\min,k}}^{r_{\max,k}} \frac{dV_k(r)}{d\ln r} d\ln r, \quad (1)$$

where $r_{\min,k}$ and $r_{\max,k}$ is the minimum and the maximum radius of the k th aerosol mode ($k = 1, \dots, K$), respectively; (2) the particle volume size distribution $dV_k(r)/d\ln r$ for each aerosol mode, (3) complex refractive indices at the wavelength λ , $m(\lambda) = n(\lambda) + ik(\lambda)$, (4) the “sphericity”, ζ_{sph} (the ratio of spherical particle’s volume to the total volume), (5) aerosol optical thickness (AOT) of the k th aerosol mode, $\hat{E}_k(\lambda_j)$, (6) the single scattering albedo for the k th aerosol mode, $\varpi_k(\lambda)$, (7) the elements of the backscattering matrix, $\mathbf{P}_{x,x}^k(\lambda, 180^\circ)$, and (8) coefficients \mathbf{a}_k and \mathbf{b}_k , which determine optical extinction and backscatter characteristics of aerosol particles for the k -aerosol mode (see Sect. 3.1). Parameters $m(\lambda)$ and ζ_{sph} are assumed the same for particles of all sizes. Definitions and detailed description of the columnar aerosol parameters are available at the AERONET information system; cloud screening and quality control algorithms were described by Holben et al. (2006).

LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Discussion Paper | Discussion Paper | Discussion Paper | Discussion Paper | Discussion Paper

Title Page	
Abstract	Introduction
Conclusions	References
Tables	Figures
◀	▶
◀	▶
Back	Close
Full Screen / Esc	
Printer-friendly Version	
Interactive Discussion	



LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



ii. The height, h , distributions of particle volume concentrations (PVC) for each of aerosol mode, $c_k(h)$, which define the vertical variability of the aerosol features.

A lack of lidar data to resolve height-variation of aerosol microstructure motivates the assumption of altitude-independent microphysical parameters of the aerosol modes.

2.2 Algorithm's structure

Two options of the retrieval procedure for the processing LRS data have been developed:

1. First one deals with sequential inversion of lidar and radiometer data. It is carried out by preliminary calculation of the i-type parameters defined in Sect. 2.1 from radiometric measurements by using the AERONET inversion algorithm (Dubovik and King, 2000), followed by subsequent inversion of the ii-type parameters by using lidar data with columnar characteristics of aerosol layer passed as a priori data (Chaikovsky et al., 2012);
2. Second option suggests parallel inversion approach for retrieving optimal i-and-ii-type parameters of the aerosol model by using a joint inversion procedure from combined lidar and radiometer data.

While the sequential algorithm could be considered as an unsophisticated inversion procedure to combine lidar and AERONET data, the parallel inversion method leads, in principle, to more effective estimation of aerosol parameters because it allows simultaneously retrieved columnar aerosol parameters to be specified in accordance with the additional lidar data. Currently, the parallel inversion algorithm for a two-component aerosol model is implemented in GARRLIC (Lopatin et al., 2013). Similar aerosol mode concentration profiles and residual discrepancies between measured and calculated input signals are obtained from both retrieval procedures when processing experimental data (Lopatin et al., 2013).

Advantages of the “parallel inversion approach” are expected for more involved measurements, such as in the unified algorithm GRASP (Generalized Retrieval of Aerosol and Surface Properties), which aimed at characterizing atmospheric properties from remote ground and satellite observations (Dubovik et al., 2014).

LIRIC algorithm described below was created on the base of the sequential inversion approach. Figure 1 shows the structure of the algorithm.

The algorithm is divided into several rather independent modules to provide flexibility of the software package. Module 1 (preprocessing of lidar data) creates a set of smoothed and normalized lidar signals, L^* , covariance matrix, Ω_L , and setting parameters (type of lidar measurement, sounding wavelength, geographical coordinates of lidar station and date of measurement, etc.) for modeling aerosol and molecular layers. Module 2 (recalculation of radiometer data) estimates i-type columnar parameters of the aerosol model for lidar sounding wavelengths. Level 1.5 or Level 2.0 AERONET data are acceptable as input data in LIRIC. Initial profiles of the aerosol-mode concentrations, $c_k^0(h)$, as well as molecular (Rayleigh) extinction, $\sigma_r(\lambda, h)$, and molecular backscatter coefficients, $\beta_r(\lambda, h)$, are generated by module 3 (atmospheric model). Module 4 (forward model) calculates arrays of lidar signals, $L_j(c_k^{m-1}(h))$, and columnar volume concentrations, $\hat{C}_k^{V,m-1}$, given aerosol concentration profiles, $c_k^{m-1}(h)$, in iterative inversion procedure, where “ m ” stands for the m th retrieval iteration and “ j ” is the number of receiving channel. Module 5 (numerical inversion) is responsible for fitting aerosol-mode concentration profiles for the retrieved aerosol model, $c_k^{m-1}(h)$, given measured data and a priori information. Inversion parameters, constraints on the smoothness characteristics, and error signals for the sensitivity test are passed to the algorithm by module 6 (inversion settings&error modeling).

LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



3 Forward modeling of LRS experiment

Range-corrected normalized lidar signals and columnar-aerosol parameters retrieved from radiometer measurements are the input data to the LRS processing procedure (see Fig. 1). Below we define a set of basic equations that are needed for the forward modeling of the measured quantities as well as to estimate the error-covariance matrix.

3.1 Basic lidar equations

The multichannel lidar carries out J “different” lidar measurements ($j \in 1, \dots, J$) that yields a set of lidar signal records, P_j^* , $j \in 1, \dots, J$. The term “different” means that different kinds of lidar measurements are performed, such of total intensity, cross- and parallel-polarized signal components at different wavelengths. Here we consider that each “different” lidar measurement is provided by a specific j th channel. Parameter J stands for the number of lidar channels irrespective of the actual implementation of the lidar system.

Range-corrected normalized lidar signals are calculated at the preprocessing stage of the inversion procedure (Module 1 in Fig. 1):

$$L_j^*(h) = \frac{S_j^*(\lambda_j, h)}{\hat{S}_j^*(\lambda, h_{\text{ref}})} \exp(-2\tau_r(\lambda_j, h, h_{\text{ref}})), \quad (2)$$

where $S_j^*(\lambda_j, h) = P_j^*(\lambda_j, h)h^2$; $\hat{S}_j^*(\lambda_j, h_{\text{ref}})$ is the value of $S_j^*(\lambda_j, h)$ at the reference point, h_{ref} is usually defined in the end of the sensing range; $\tau_r(\lambda_j, h, h_{\text{ref}})$ is the molecular optical thickness related to the range of (h, h_{ref}) , λ_j is the wavelength, and h is the height. The set of lidar signals, $L_j^*(h)$, constitutes the input lidar vector, L^* .

The lidar system provides measurements from the lowest to the highest altitude levels specified by h_{min} and h_{max} , respectively. Currently, it is assumed that the radiometer is co-located at a height of $h_0 < h_{\text{min}}$, so columnar aerosol optical properties of the layer $h_0 < h < h_{\text{min}}$ are to be taken into consideration. If there is no information on the aerosol

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



parameters in the surface layer, this layer is assumed to be homogeneous. Under this assumption, scattering parameters for the altitude range $h_0 < h < h_{\min}$ of the lidar vector L^* are set equal to the values at h_{\min} .

The relationship between the measured lidar signals $L^*(\lambda)$, and the aerosol mode concentration, $c_k(h)$, can be written as follows:

$$L^* = L(\lambda, c_k(h), \mathbf{a}_k, \mathbf{b}_k) + \mathbf{\Delta}_L, \quad (3)$$

where $\mathbf{\Delta}_L$ is the vector of measurement uncertainties. Here, an asterisk (*) denotes “measured” and no-asterisk denotes “model estimated”.

Since function $L(\dots)$ in Eq. (3) depends on the type of lidar measurement, it is expedient to introduce special parameter, $\rho_j \in 1, u, U$ that indicates the type of measurement associated to the j -channel of the lidar. In our case, $\rho_j \in 1, 2, 3$, indicates total intensity, cross-polarized, and parallel-polarized measurements, correspondingly.

The lidar functions, $L_{j,\rho_j}(\dots)$, for the ρ_j -type measurements are defined by the following equations:

Equation for the total backscatter signal:

$$L_{j,1}(\lambda_j, c_k(h)) = \frac{\beta_{a,1}(\lambda_j, h) + \beta_r(\lambda_j, h)}{R_{j,1}(\lambda_j, h_{\text{ref}})\beta_r(\lambda_j, h_{\text{ref}})} \exp\left(-2 \int_{h_{\text{ref}}}^h \sigma_a(\lambda_j, h) dh\right), \quad (4)$$

where

$$R_{j,1}(\lambda_j, h) = \frac{\beta_{a,1}(\lambda_j, h) + \beta_r(\lambda_j, h)}{\beta_r(\lambda_j, h)}. \quad (5)$$

Equation for the parallel-polarized signal component:

$$L_{j,3}(\lambda_j, c_k(h)) = \frac{\beta_{a,3}(\lambda_j, h) + \frac{1}{1+\chi}\beta_r(\lambda_j, h)}{\frac{1}{1+\chi}\beta_r(\lambda_j, h_{\text{ref}})R_{j,3}(\lambda_j, h_{\text{ref}})} \exp\left(-2 \int_{h_{\text{ref}}}^h \sigma_a(\lambda_j, h) dh\right), \quad (6)$$

where

$$R_{j,3}(\lambda_j, h_i) = \frac{\beta_{a,3}(\lambda_j, h_i) + \frac{1}{1+\chi}\beta_r(\lambda_j, h_i)}{\frac{1}{1+\chi}\beta_r(\lambda_j, h_i)}. \quad (7)$$

Equation for the cross-polarized signal component:

$$L_{j,2}(\lambda_j, c_k(h)) = \frac{(\beta_{a,2}(\lambda_j, h) + \mu\beta_{a,3}(\lambda_j, h) + \frac{\chi+\mu}{\chi+1}\beta_r(\lambda_j, h))}{\frac{\chi+\mu}{\chi+1}\beta_r(\lambda_j, h)R^{\text{eff}}(\lambda_j, h_{\text{ref}})} \exp\left(-2 \int_{h_{\text{ref}}}^h \sigma_a(\lambda_j, h)dh\right) \quad (8)$$

5 where

$$R^{\text{eff}}(\lambda_j, h) = \left(\frac{\chi}{\chi + \mu} \frac{\beta_{a,2}(\lambda_j, h) + \frac{\chi}{\chi+1}\beta_r(\lambda_j, h)}{\frac{\chi}{\chi+1}\beta_r(\lambda_j, h)} + \frac{\mu}{\chi + \mu} \frac{\beta_{a,3}(\lambda_j, h) + \frac{1}{1+\chi}\beta_r(\lambda_j, h)}{\frac{1}{1+\chi}\beta_r(\lambda_j, h)} \right). \quad (9)$$

In Eqs. (4)–(9), $\beta_{a,1}$, $\beta_{a,3}$, and $\beta_{a,2}$ denote the aerosol backscatter coefficient and its parallel- and cross-polarized components, respectively; $\sigma_a(\lambda_j, h)$ is the aerosol extinction coefficient; $\chi(\lambda_j) = \frac{\beta_{r,2}(\lambda_j)}{\beta_{r,3}(\lambda_j)}$ is the ratio of cross- and parallel-polarized components of the molecular backscatter coefficient.

Different cross-talk factors contribute to the spurious signal in the cross polarized receiving channel. These factors include the residual of cross polarized component of the laser beam, non-ideal adjustment of the polarization planes between transmitter/receiver channels and depolarization by optical elements. Equations (6) and (8) allow for these cross-talk effects in a similar manner to (Chaikovsky, 1990; Biele et al., 2000). Thus, parameter μ in Eqs. (8)–(9) represents the leakage of the parallel component of the sounding beam into the cross polarized lidar receiving channel.

The aerosol extinction and backscatter coefficients in the Eqs. (3)–(9) are expressed as a function of the parameters of the aerosol modes:

$$\sigma_a(\lambda_j, h) = \sum_k c_k(h) a_k(\lambda_j), \quad (10)$$

$$\beta_{a,1}(\lambda_j, h) = \sum_k c_k(h) b_{k,1}(\lambda_j), \quad (11)$$

$$5 \quad \beta_{a,2}(\lambda_j, h) = \sum_k c_k(h) b_{k,2}(\lambda_j), \quad (12)$$

and

$$\beta_{a,3}(\lambda_j, h) = \sum_k c_k(h) b_{k,3}(\lambda_j). \quad (13)$$

The coefficients $a_k(\lambda_j)$ and $b_{k,x}(\lambda_j)$, pointed out in Sect. 2.1, are determined by columnar optical parameters of aerosol modes:

$$10 \quad a_k(\lambda_j) = \frac{\hat{E}_k(\lambda_j)}{\hat{C}_k^V}, \quad (14)$$

$$b_{k,1}(\lambda_j) = \frac{1}{4\pi} \varpi_k(\lambda_j) a_k(\lambda_j) P_{1,1}^k(\lambda_j, 180^\circ), \quad (15)$$

$$b_{k,3}(\lambda_j) = \frac{1}{4\pi} \varpi_k(\lambda_j) a_k(\lambda_j) \frac{P_{1,1}^k(\lambda_j, 180^\circ) + P_{2,2}^k(\lambda_j, 180^\circ)}{2}, \quad (16)$$

$$b_{k,2}(\lambda_j) = \frac{1}{4\pi} \varpi_k(\lambda_j) a_k(\lambda_j) \frac{P_{1,1}^k(\lambda_j, 180^\circ) - P_{2,2}^k(\lambda_j, 180^\circ)}{2}. \quad (17)$$

LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



3.2 Forward model of radiometer data

In accordance with the multi-term LSM approach (Dubovik, 2004), the columnar concentrations of aerosol modes, \hat{C}_k^V , obtained from radiometer measurements are formally considered in LIRIC as a result of additional independent measurements.

The equation for the vector, \hat{C}^{*V} , which is defined as the “measured” columnar volume concentrations of the aerosol modes given vector of aerosol modes concentration, $c(h_i)$, $i \in 1, I$ can be written in the following form

$$\hat{C}^{*V} = \mathbf{H}\mathbf{c} + \mathbf{\Delta}_V \quad (18)$$

The k th component of the vector \hat{C}^{*V} is defined by the equation

$$C_k^{*V}(c_k(h_i)) = \sum_{i=1}^I c_k(h_i) \Delta h_i. \quad (19)$$

The structure of the vectors \hat{C}^{*V} , \mathbf{c} and matrix \mathbf{H} is considered in Appendix C.

4 Numerical inversion

Statistical regularization technique (e.g. Turchin et al., 1971; Tarantola, 1987; Rodgers, 2000) considers errors, $\mathbf{\Delta}_L$ and $\mathbf{\Delta}_C$, in Eqs. (3) and (18) as random variables. Under the additional assumption that errors have independent normal distributions, the multidimensional conditional probability density function (PDF) (or “likelihood function”) is defined by (Chaikovsky et al., 2004a)

$$F(L^*, \hat{C}^{*V} | \mathbf{c}) \sim \exp \left[\left(-\frac{1}{2} \left((L^* - L(\mathbf{c}))^T \mathbf{\Omega}_L^{-1} (L^* - L(\mathbf{c})) + (\hat{C}^{*V} - \mathbf{H}\mathbf{c})^T \mathbf{\Omega}_V^{-1} (\hat{C}^{*V} - \mathbf{H}\mathbf{c}) \right) \right) \right] \quad (20)$$

12774

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Here, $F(L^*, \hat{C}^{*V} | c)$ is the PDF of measurement vectors L^* and \hat{C}^{*V} , $L(c)$ is the vector function in Eq. (3), H is the matrix in Eq. (18), c is the target retrieval vector of aerosol modes concentration, and Ω_L and Ω_V are the covariance matrices of error vectors Δ_L and Δ_V , respectively.

An extensively used tool for the regularization of an “ill-posed” problem is the application of a priori constraint on the smoothness of retrieved characteristics. LIRIC restricts the norms of the second differences of functions $c_k(h_j)$. Following the statistical regularization approach (Turchin et al., 1971) we included a priori probability function,

$$F_{\text{apr}}(c) \sim \exp\left(-\frac{1}{2} \left(c^T \Omega_S c\right)\right) \quad (21)$$

into the retrieval procedure as the additional constraint. Here, $\Omega_S = S_2^T Q_2^{-1} S_2$ is smoothing matrix, S_2 is the matrix of the second-order differences, and Q_2 is diagonal weighting matrix (Twomey, 1977; Dubovik et al., 2011).

The Bayes’ strategy (Turchin et al., 1971; Tarantola, 1987; Rodgers, 2000) for solving an “ill-posed” problem combined with multi-term LSM technique (Dubovik, 2004; Dubovik et al., 2011) defines the solution \hat{c} in accordance with the maximum a posteriori rule

$$\hat{c} = \arg \min_c \{\Psi(c)\},$$

where the objective or cost function, $\Psi(c)$, has the following multi-term representation (Dubovik, 2004; Dubovik et al., 2011)

$$\Psi(c) = (L^* - L(c))^T \Omega_L^{-1} (L^* - L(c)) + (\hat{C}^{*V} - Hc)^T \Omega_V^{-1} (\hat{C}^{*V} - Hc) + (c^T S_2^T Q_2^{-1} S_2 c) \quad (22)$$

We assume that the errors Δ_L in Eq. (3) and Δ_V in Eq. (18) are uncorrelated. In this case the non-zero diagonal elements of the covariance matrices Ω_L and Ω_V are the variances of the elements of the vectors Δ_L and Δ_V , respectively.

Since the minimization procedure does not prescribe a residual value for $\Psi(\mathbf{c})$, it is convenient to reformulate weight matrices as follows (Dubovik, 2004):

$$\check{\check{\Omega}}_L = \frac{1}{\varepsilon_L^2} \Omega_L; \quad \check{\check{\Omega}}_V = \frac{1}{\varepsilon_V^2} \Omega_V; \quad \check{\check{\Omega}}_S = \frac{1}{\varepsilon_S^2} \Omega_S, \quad (23)$$

where ε_L^2 , ε_V^2 , and ε_S^2 are the first elements of the corresponding covariance matrices.

After substitution of the covariance matrices expressed through the weight matrices into Eq. (22) and multiplication it by ε_L^2 , the $\Psi(\mathbf{c})$ takes the form of the sum of three components:

$$\check{\check{\Psi}}(L^*, \hat{\mathbf{C}}_V, \mathbf{c}) = \check{\check{\Psi}}_L(L^*, \mathbf{c}) + \gamma_V \check{\check{\Psi}}_V(\hat{\mathbf{C}}^{*V}, \mathbf{c}) + \gamma_S \check{\check{\Psi}}_S(\mathbf{c}), \quad (24)$$

where

$$\check{\check{\Psi}}_L(L^*, \mathbf{c}) = (L^* - L(\mathbf{c}))^T \check{\check{\Omega}}_L^{-1} (L^* - L(\mathbf{c})), \quad (25)$$

is related to “lidar-measured” data, Eq. (3),

$$\check{\check{\Psi}}_V(\hat{\mathbf{C}}^{*V}, \mathbf{c}) = (\hat{\mathbf{C}}^{*V} - \mathbf{H}\mathbf{c})^T \check{\check{\Omega}}_V^{-1} (\hat{\mathbf{C}}^{*V} - \mathbf{H}\mathbf{c}), \quad (26)$$

is related to radiometer-measured data, Eq. (18),

$$\check{\check{\Psi}}_S(\mathbf{c}) = (\mathbf{c}^T \mathbf{S}^T \check{\check{Q}}_2^{-1} \mathbf{S}\mathbf{c}), \quad (27)$$

is related to a priori information, Eq. (21),

$$\gamma_V = \frac{\varepsilon_L^2}{\varepsilon_V^2}; \quad \gamma_S = \frac{\varepsilon_L^2}{\varepsilon_2^2}. \quad (28)$$

The coefficients γ_V and γ_S are so-called *Lagrange* multipliers that determine the weight of different contributors from each source of information (i.e., “measurements”

LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



and “a priori” contribution) to the retrieval solution relative to the contribution of the first data source (since $\gamma_L = 1$). Equations (22) and (24) are equivalent; however Eq. (24) is more convenient for the analysis of the relative contribution from different data source.

If $\gamma_V, \gamma_S \rightarrow 0$, we return to a non-regularized solution for vector c that is based solely on measured lidar data with the minimum discrepancy between measured and calculated input signals. This solution, however, could be non-physical, multivalued, and unstable. The possible solution space should be restricted by increasing the *Lagrange* multipliers despite the fact that it results in increasing of discrepancy between measured and model signal. The algorithms to determine the *Lagrange* multipliers by finding a reasonable compromise between the solution quality and the closeness of the measured and model signals are described in Hansen, 2001; Vogel, 2002; and Doicu et al., 2010.

The final step of the retrieval procedure is calculation of the concentration profiles $c_k(h_i)$ for each aerosol mode. Initial approximations $c_k^0(h_i)$ are set and stepwise improved to provide the minimum of the objective function (Eq. 25). Increments are calculated by means of the Levenberg–Marquardt method (Levenberg, 1944; Marquardt, 1963).

The analytical expressions of the terms of Eq. (25), the covariance matrices, as well as the details of the inversion procedure, are described in Appendices A, B and C.

5 Program package for processing combined lidar and radiometer data

Figure 2 shows the structure of the software package that implements the LIRIC algorithm. A set of specific programs are joined in three sub-packages.

The sub-package LiOpt implements module (3) of the LIRIC algorithm (Fig. 1), which provides preprocessing of the AERONET retrieval products. Program AERLID recalculates the columnar optical characteristics for the lidar sounding wavelengths, including the elements of the scattering matrices for the spherical and non-spherical particles

LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



as well as for fine and coarse aerosol modes. Then, this code writes data down to the Radiometer Database.

The preprocessing of lidar data is carried out by the SignalSuite sub-package. It contains several programs. Among them:

- ULIS is an operational program that provides measurement procedures and record of raw lidar data to Microsoft ACCESS database (DB Lidar raw);
- nc2mdb is a program to convert EARLINET standard raw-lidar nc-files into mdb-files to process by LIRIC;
- program Synthesizer averages the series of lidar signals, converts the profiles to the optimal altitude scale and, then, “glues” signals and provides the “dead-time” correction of photo-counting lidar signals;
- program Tropoexport calculates a normalized smoothed lidar signal and its variance, and generates molecular and aerosol atmospheric models; this program aims at implementing modules 2 and 3 of the algorithm.

Finally, the main sub-package ProfileRetriever implements the LIRIC inversion procedure. Program ConcentRetriver retrieves profiles $c_V^{k,m}(h)$ of the aerosol mode concentrations and writes data down to Access database, DB-processed. Module Inversion setting&Errors modeling generates a set of noise-corrupted input data files by adding “white noise” and amplitude distortions to the initial lidar signals and perturbing aerosol model parameters retrieved from radiometer measurements in order to provide the error sensitivity analysis. Real measurement conditions, technical features of the lidar system and the accuracy of columnar aerosol parameters retrieved from the radiometer measurements (Dubovik et al., 2000a, b) are taken into account in setting parameters of the module. Program OutputViewer allows viewing the output data and their conversion from mdb-files into other formats.

LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page	
Abstract	Introduction
Conclusions	References
Tables	Figures
◀	▶
◀	▶
Back	Close
Full Screen / Esc	
Printer-friendly Version	
Interactive Discussion	



6 Verification of operability and sensitivity tests

The LRS technique applies the aerosol model to the retrieval scheme that was basically designed in AERONET. This model assumes that aerosol consists of fine and coarse modes and that both are mixtures of spherical particles and randomly-oriented homogeneous spheroids. The advanced T matrix code (Mishchenko et al., 2000, 2002) provides computation of scattering matrices of the aerosol particles. Thus, any optical characteristic of the aerosol layer can be calculated using data of the LRS experiment.

The applicability analysis of the AERONET spheroid model to aerosol particles is beyond the scope of this paper. We only note that this model was validated by the comparison of calculated optical parameters and laboratory measurements of light scattering matrices for mineral dust particles (Volten et al., 2001). Incorporation of the spheroid model into AERONET operational retrieval code has significantly improved AERONET products when evaluating parameters of coarse non-spherical particles (Cattrell et al., 2005; Dubovik et al., 2006). This model has also been incorporated when processing data from ground-based polarimetric measurements (e.g. Li et al., 2009), lidar sounding data (e.g. Veselovskii et al., 2010; David et al., 2013; Müller et al., 2013), and satellite-base observations (e.g. Levy et al., 2007a, b; Dubovik et al., 2011; Schuster et al., 2012).

6.1 Verification of LIRIC program package: EARLI09 intercomparison experiment

EARLI09 intercomparison experiment was held in May 2009 at Leibniz Institute for Tropospheric Research in Leipzig, Germany (Wandinger et al., 2015). This campaign provided excellent opportunity to validate the LRS technique for network measurements. The results of the LIRIC data processing for simultaneous measurements by seven lidars of different scientific teams on 25 May 2009 in Leipzig were compared.

LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Total optical depth distribution (Fig. 3a) and back-trajectories analysis (Fig. 3b) indicates that LRS measurements were carried out during the Saharan dust event in the Leipzig region and the dust was transported in the layer above 2 km.

Figures 4 and 5 show PVC profiles, $c_k(h)$, retrieved from lidar data of the different EARLINET teams combined with the same AERONET information, as well as their root mean square deviations and relative deviations for the two types of input data set, namely, with and without depolarization measurements.

It is evident from Figs. 4a and 5a that $c_k(h)$ profiles show close agreement in structure and magnitude over the troposphere except for the lower layer. The relative deviations increase only when values of the aerosol concentration become negligible.

We explain the discrepancy between $c_k(h)$ profiles in the near-surface atmosphere by the uncertainty in geometrical overlap factors and the differences in lower-boundary heights of the considered lidar systems. Also some differences in the retrieved concentration profiles $c_k(h)$ are due to measurement errors, uncertainties in aerosol modeling as well as specificities on the inverse operator (see Appendix C).

The potential errors in the PVC profiles for the specific combined lidar/radiometer experiment were estimated by using the Errors modeling module of the LIRIC package (Fig. 2). Figures 6 and 7 illustrate the sensitivity of the retrieved aerosol concentration profiles to the errors of the lidar measurements. The original lidar signals were taken as they measured by München lidar (curves 4 in Fig. 6) and have been perturbed by adding “white noise” with different root-mean-square deviations (rms-deviations), α_j , and have been distorted by multiplying them by the coefficient,

$$k_j(h_i) = 1 + \frac{\Delta_j}{100} \frac{h_{\text{ref}} - h_i}{h_{\text{ref}}}, \quad (29)$$

where percentage parameter Δ_j determines the amount of non-linearity.

In response, the program module generated twelve disturbed lidar signal sets that allowed us to estimate the impact of measurement errors. As an illustration, Figs. 6 and 7 simulate higher errors than typical ones in most EARLINET lidars. Four realizations of

the disturbed signals are shown in Fig. 6. Coefficient $k_j(h_j)$ increases/decreases from referent to start point that results in divergence of the lidar signals in Fig. 6.

PVC profiles, $c_k(h)$, corresponding to the lidar signals in Fig. 6 and their rms-deviations calculated for full ensembles of input data are shown in Fig. 7. Changes in the PVC profiles of the dominant coarse non-spherical mode are shown by the Fig. 7 to be minor (Fig. 7c). Although profiles $c_k(h)$ of fine and coarse spherical particles (Fig. 7a and b) are not very stable; they qualitatively retain similarity with the initial distributions.

Figure 8 illustrates the effect of uncertainties in columnar aerosol parameters retrieved from radiometer data. Variations of the columnar aerosol characteristics lead to changes in coefficients a and b of lidar-related Eqs. (14)–(17) (Sect. 3.1). Statistical characteristics of aerosol concentration profiles retrieved with relative deviation of the parameter $\vartheta_{k,p}^j$ (effective lidar ratio of the aerosol fraction, see Appendix B) in the range $\pm 20\%$ are presented in Fig. 8. Relative deviation of aerosol concentration profile becomes significant only for small values of the concentration.

6.2 Dependence of retrieved aerosol concentration profiles on the content of the input data set

Three types of data set related to different sources of information compose the LIRIC input data-file: three or four measured lidar signals, column-aerosol parameters from radiometer measurements, and a priori smoothness constraints. Two- or three-mode aerosol models are used according to the type of the measured lidar signals. Formally, we deal with redundant input information and, hence, the number of input data set can be decreased. Consequently, the significance of the different information components in retrieval procedure is of interest as well as variations of the retrieved profiles, $c_k(h)$, in the absence of some input data.

As pointed in Sect. 4, the objective functions of LIRIC regularization algorithm (Eq. 22), consists of a set of terms, which implement contribution of different types of input data into the retrieval process. Setting the variance of the specific kind of mea-

LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



surement to a large value implies neglecting of the correspondent term in the objective function (Eq. 22) and elimination of this part of the input data in estimation of the final aerosol parameters. Program package implements this option and makes allowing one to analyze the contribution of different measured data in the processing procedure of specific experiment.

Below we shortly examine sensitivity of the retrieved profiles, $c_k(h)$, to the input data selection for the case of combined lidar/radiometer sounding of the atmospheric aerosol during the last period of Eyjafjallajökull volcano ash transport to the European area in Lille, France, on the 19 May 2010. Air mass back-trajectories (Fig. 9) forecasted the possibility of appearance of volcanic ash in the layer between 1300 and 2500 m. The structure of the retrieved profiles, $c_k(h)$, shown in Fig. 10a agrees well with the forecast. Deviations $\delta(c_k(h_i))$ associated to the profiles $c_k(h)$ have been calculated by an “error modelling” procedure similar to the one described in Sect. 6.1.

A mixture of spherical and non-spherical particles constitutes the aerosol layer at the height of about 2000 m. The profile of particle depolarization ratio at 532 nm and its deviation have been calculated from the retrieved aerosol mode concentrations, $c_k(h)$. The profiles are shown in Fig. 10b, curves D(3) and rms_dev(3). The results of the direct calculation of depolarization ratio and their deviations from lidar measurements are presented by curves D(2) and rms_dev(2). Profiles D(2) and D(3) show rather close agreement in magnitude and vertical structure that could confirm the efficiency of the aerosol modeling used in this study.

The curves in Fig. 11 show the deviations in the retrieved concentration profiles, $c_k(h)$, after elimination one of the lidar signals or columnar volume concentrations of aerosol modes, C^V , from the input data set. As can be seen from Fig. 11, the concentration profile of the fine-particle mode undergoes minor changes upon elimination of a single lidar signal or columnar volume concentrations. This implies that our experiment well-defined with respect to the fine-mode concentration. On the other hand, concentrations of coarse modes are sensitive to input information. Thus, lidar data at 1064 nm wavelength plays a crucial role in the retrieval of the coarse spherical mode.

LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



In the same manner, lidar depolarization measurement is the key factor in the retrieval of the coarse spheroid particle mode. Evaluations of columnar volume concentrations from radiometer measurement are necessary for all cases.

Figure 12a shows concentration profiles, $c_k(h)$, which were retrieved for two- and three-mode aerosol models and characterized the aerosol layer in the same LRS experiment. The fine-mode concentration profiles for two aerosol models are practically coincident. Profiles $c_k(h)$ of coarse modes for two-mode aerosol model, coarse(3), and the sum of two coarse components for three-mode aerosol model, coarse(4), are similar in shape but quantitatively are a bit different. The column concentrations of the course (3) and (4) modes are equal.

The curves in Fig. 12b and c show the deviations of the concentration profiles, $c_k(h)$, for the two-mode aerosol model after reduction of the input data set. Deviations of $c_k(h)$ profiles are rather similar to those for the three-mode aerosol model in Fig. 11. Deviations of fine-mode concentration profile are small, even if any single sub-set of input data is eliminated. Coarse-mode concentration profiles preserve original forms when one of the lidar signals at the 355 or 532 nm wavelength is excluded from the processing procedure.

Generally, for measurement conditions that characterize the experiment under discussion, two-wavelength lidar sounding (at 355 and 1064 or at 532 and 1064 nm) combined with radiometer measurement provides retrieving concentration profiles of fine and coarse aerosol modes for two-mode aerosol model.

7 Discussion and conclusions

The active process of dissemination of the LIRIC in EARLINET started in 2012. Nowadays, 11 EARLINET teams participate in implementation of LRS technique (see Fig. 13).

Since that time experimental works on the validation of the LIRIC product have been carried out. For example, evaluation of the uncertainties in the retrieved aerosol pa-

LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



rameters for different aerosol types, aerosol loads, and overlap characteristics of the lidar systems was made in Granados-Muñoz et al., 2014. Comparisons of aerosol optical and microphysical parameters retrieved from LIRIC against profiles of aerosol backscatter coefficients and depolarization ratios directly calculated from lidar-data measurements (e.g. Tsekeri et al., 2012, 2013; Wagner et al., 2013; Kokkalis et al., 2013; Granados-Muñoz et al., 2014) as well as against modeled or airborne in situ measured profiles of aerosol mode concentrations (e.g. Kokkalis et al., 2012, 2013; Nemuc et al., 2013) have shown reasonable agreement for different types of aerosols.

The LIRIC concentration profiles of aerosol fractions during dust and volcano ash events have been compared with those for spherical and non-spherical particles derived from polarization measurements using the POLIPHON technique (e.g. Wagner et al., 2013; Nemuc et al., 2013; Papayannis et al., 2014). In spite of the noticeable difference between the aerosol models and independent processing algorithms, the retrieved aerosol concentration profiles have proved to be similar. This is quite natural because both approaches use the depolarization of backscatter signal to distinguish between spherical and non-spherical particles.

The number of aerosol studies using LIRIC algorithm increases. These studies focus on the dynamics of aerosol microstructure during transport of air masses polluted by dust (e.g. Chaikovsky et al., 2010b; Tsekeri et al., 2013; Binietoglou et al., 2015; Granados-Muñoz et al., 2015), fire smoke (e.g. Chaikovsky et al., 2004, 2010b; Pietruczuk and Chaikovsky, 2007), and volcano ash (Kokkalis et al., 2013). The list of lidar teams that take advantage of the LIRIC is still expanding.

In its turn, the experience of EARLINET teams in data processing of LRS experiments allows one to outline some common features of LIRIC solutions:

- i. The assumption that microphysical parameters for the different aerosol modes are altitude-independent can be inappropriate for some cases. On the other hand, lidar signals calculated using profiles, $c_k^m(h_i)$, in the LIRIC fitting procedure agree well with the measured profiles (within the measurement errors). This is the evidence that there is not much information on the profile shape of microphysical

We will use the utility function

$$\delta_{p_j, u}^j = \begin{cases} 1 & \dots \text{if} \dots p_j = u \\ 0 & \dots \text{if} \dots p_j \neq u \end{cases}, \quad (\text{A1})$$

along with the following definitions of combinations of aerosol and molecular optical parameters in Eqs. (4)–(9):

$$\begin{aligned} \beta_a^{\text{ef}}(\lambda_j, p_j, h) &= \left(\beta_{a, p_j}(\lambda_j, h) + \delta_{p_j, 2}^j \mu \beta_{a, 3}(\lambda_j, h) \right) \\ &= \left(\sum_k c_k(h) b_{k, p_j}(\lambda_j) + \delta_{p_j, 2}^j \mu \sum_k c_k(h) b_{k, p_j}(\lambda_j) \right) \end{aligned} \quad (\text{A2})$$

$$\beta_r^{\text{ef}}(\lambda_j, p_j, h) = \left(\delta_{p_j, 2}^j (p_j) \left(\frac{\mu - 1}{\chi + 1} \right) + \frac{1}{1 + \delta_{p_j, 3}^j \chi} \right) \beta_r(\lambda_j, h) \quad (\text{A3})$$

$$\beta^{\text{ef}}(\lambda_j, p_j, h) = \beta_a^{\text{ef}}(\lambda_j, p_j, h) + \beta_r^{\text{ef}}(\lambda_j, p_j, h) \quad (\text{A4})$$

$$\hat{R}_j^{\text{ef}}(\lambda_j, p_j, h) = \frac{\beta_a^{\text{ef}}(\lambda_j, p_j, h) + \beta_r^{\text{ef}}(\lambda_j, p_j, h)}{\beta_r^{\text{ef}}(\lambda_j, p_j, h_i)} \quad (\text{A5})$$

$$\tau_a(\lambda_j, h, h_{\text{ref}}) = \int_h^{h_{\text{ref}}} \sigma_a(\lambda_j, h) dh \quad (\text{A6})$$

This permits Eqs. (4), (6) and (7) to be written in general form:

$$L_j(p_j, \lambda_j, h) = \frac{\beta^{\text{ef}}(\lambda_j, p_j, h) \exp(2\tau_a(\lambda_j, h, h_{\text{ref}}))}{\beta_r^{\text{ef}}(\lambda_j, p_j, h_{\text{ref}}) \hat{R}_j^{\text{ef}}(\lambda_j, p_j, h_{\text{ref}})}. \quad (\text{A7})$$

Therefore, the related to the lidar objective function, $\check{\Psi}_L(L^*, \mathbf{c})$, (Eq. 24), is given by the equation:

$$\check{\Psi}_L(L^*, \mathbf{c}) = \sum_j \sum_i \frac{\Delta h_i}{\check{\Omega}_{L_j}(i, i)} \left(L_{j,i}^* - \frac{\left(\sum_k c_k(h_i) b_{k,p_j}(\lambda_j) + \delta_{p,2}^j \mu \sum_k c_k(h_i) b_{k,p_j}(\lambda_j) \right)}{\beta_r^{\text{ef}}(\lambda_j, p_j, h_{\text{ref}}) \hat{R}_j^{\text{ef}}(\lambda_j, p_j, h_{\text{ref}})} \right)^2 \times \exp \left(2 \sum_k \sum_i c_k(h_i) a_k(\lambda_j) \Delta h_i \right)$$

$i \in 1, \dots, l.$

(A8)

Equation (26), $\check{\Psi}_V(\hat{\mathbf{C}}^{*V}, \mathbf{c})$, which brings radiometer data into the processing procedure can be expressed as follows:

$$\check{\Psi}_V(\hat{\mathbf{C}}^{*V}, \mathbf{c}) = \sum_k \frac{1}{\check{\Omega}_V(k, k)} \left(\hat{\mathbf{C}}^{*V} - \sum_i c_k(h_i) |\Delta h_i| \right)^2.$$

(A9)

Calculation of the “smoothness” part of the objective function is described in details in Dubovik, 2004 and Dubovik et al., 2011.

Appendix B: Evaluation of covariance matrix

The covariance matrixes, $\mathbf{\Omega}_L$, $\mathbf{\Omega}_V$, and $\mathbf{\Omega}_2$, defined in Sect. 4 characterize uncertainties of the complex input vector, $(L^*, \hat{\mathbf{C}}^{*V}, \hat{\mathbf{O}})$, where \mathbf{O}^* is “zero” vector that is defined to formalize a priori smoothness restrictions on concentration profiles (e.g. Dubovik, 2004). These matrices determine the “weights” of different parts of input information through the minimization procedure of the objective function (Eq. 22).

In our case the measure of the smoothness for concentration profiles, $c_k(h_j)$, should be chosen as a priori evaluated parameters. Aerosol columnar volume concentrations, \hat{C}^{*V} , and variances, $\Omega_V(k, k)$, are the parts of input radiometer data. Thus, only evaluation of covariance matrix, Ω_L , is to be done.

The assumption of independent normal distribution for variations of “lidar” vector, L^* , at different heights implies the diagonal covariance matrix. The non-zero diagonal elements, $\Omega_{L_j}(h_j, h_j)$, of the covariance matrix are the variances of differences between the components, $L_{j,i}^*$, of the lidar vector and the appropriate modeled function, $L_j(p_j, \lambda_j, h_j)$, in Eq. (A7).

Given Eqs. (2), (3) and (A1)–(A8), the elements of vector, Δ_{L_j} , are defined by:

$$\begin{aligned} \Delta_{L_j}(h_j) &= L_{j,i}^* - L_j(p_j, \lambda_j, h_j) \\ &= \frac{S^*(\lambda_j, h_j)}{\hat{S}^*(\lambda_j, h_{\text{ref}})} \exp(-\tau_r(\lambda_j, h_j, h_{\text{ref}})) - \frac{\beta^{\text{ef}}(\lambda_j, p_j, h_j) \exp(2\tau_a(\lambda_j, h_j, h_{\text{ref}}))}{\beta_r^{\text{ef}}(\lambda_j, p_j, h_{\text{ref}}) \hat{R}_j^{\text{ef}}(\lambda_j, p_j, h_{\text{ref}})} \end{aligned} \quad (\text{B1})$$

Using the finite differences technique (e.g. Russell et al., 1979) one can expand $\Delta_{L_j}(h_j)$ in *Taylor* series, and then neglect all the terms of the second or higher order. As a result, variation $\delta(\Delta_{L_j}(h_j))$ can be expressed as a function of variations related with the input parameters, $\delta(S^*(\lambda_j, h_j))$, $\delta(\beta^{\text{ef}}(\lambda_j, p_j, h_j))$, $\delta(\tau_a(\lambda_j, h_j, h_{\text{ref}}))$, and $\delta(\tau_r(\lambda_j, h_j, h_{\text{ref}}))$:

$$\begin{aligned} \delta(\Delta_{L_j}(h_j)) &= -2L_{j,i}^* \delta(\tau_r(\lambda_j, h_j, h_{\text{ref}})) + L_{j,i}^* \frac{\delta(S^*(\lambda_j, h_j))}{S^*(\lambda_j, h_j)} \\ &+ \frac{\beta^{\text{ef}}(\lambda_j, p_j, h_j) \exp(2\tau_a(\lambda_j, h_j, h_{\text{ref}}))}{\beta_r^{\text{ef}}(\lambda_j, p_j, h_{\text{ref}}) \hat{R}_j^{\text{ef}}(\lambda_j, p_j, h_{\text{ref}})} \frac{\delta(\beta^{\text{ef}}(\lambda_j, p_j, h_j))}{\beta^{\text{ef}}(\lambda_j, p_j, h_j)} \end{aligned} \quad (\text{B2})$$

LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



$$\begin{aligned}
 & -2 \frac{\beta^{\text{ef}}(\lambda_j, \rho_j, h_j) \exp(2\tau_a(\lambda_j, h_j, h_{\text{ref}}))}{\beta_r^{\text{ef}}(\lambda_j, \rho_j, h_{\text{ref}}) \hat{R}_j^{\text{ef}}(\lambda_j, \rho_j, h_{\text{ref}})} \delta(\tau_a(\lambda_j, h_j, h_{\text{ref}})) \\
 & \approx L_{j,i}^* \left(\frac{\delta(S^{*j}(h_j))}{S^{*j}(h_j)} - \frac{\delta(\beta^{\text{ef}}(\lambda_j, \rho_j, h_j))}{\beta^{\text{ef}}(\lambda_j, \rho_j, h_j)} - 2\delta(\tau_r(\lambda_j, h_j, h_{\text{ref}})) - 2\delta(\tau_a(\lambda_j, h_j, h_{\text{ref}})) \right).
 \end{aligned}$$

Under the assumption of independent variations of different parameters, the variance $\Omega_L(h_i, h_n)$ is expressed as follows

$$\begin{aligned}
 \Omega_L(h_i, h_i) &= \langle \delta(\Delta_{L_j}(h_i)) \delta(\Delta_{L_j}(h_i)) \rangle \\
 &= L_{j,n}^{*2} \left(\frac{\delta^2(P_{j,i}^*)}{(P_{j,i}^*)^2} + \frac{\delta^2(\beta^{\text{ef}}(\lambda_j, \rho_j, h_j))}{(\beta^{\text{ef}}(\lambda_j, \rho_j, h_j))^2} + 4\delta^2(\tau_r(\lambda_j, h_j, h_{\text{ref}})) + 4\delta^2(\tau_a(\lambda_j, h_j, h_{\text{ref}})) \right),
 \end{aligned} \tag{B3}$$

where $\langle \dots \rangle$ denotes ensemble averaging over measurement realizations, and $P_{j,i}^* = P_j^*(h_i)$.

The terms in the large round parentheses in Eq.(B3) determine contributions of measurement errors and uncertainties of a priori defined optical characteristics. We aim at approximate estimation of $\Omega_{L_j}(h_i, h_i)$ at the preprocessing stage without involving of retrieved parameters. This feedback-free approach greatly simplifies the structure of the inversion algorithm.

Uncertainties of the optical parameters

The term $\delta^2(\beta^{\text{ef}}) / (\beta^{\text{ef}})^2$ in Eq. (B3) is the relative variance of the total backscatter coefficient. It can be transformed into the sum of relative variances of aerosol and

molecular backscatter coefficients:

$$\frac{\delta^2 \left(\beta^{\text{ef}}(\lambda_j, p_j, h_i) \right)}{\left(\beta^{\text{ef}}(\lambda_j, p_j, h_i) \right)^2} = \frac{\delta^2 \left(\beta_r^{\text{ef}}(\lambda_j, p_j, h_i) \right)}{\left(\beta_r^{\text{ef}}(\lambda_j, p_j, h_i) \right)^2} \frac{1}{\left(\hat{R}_j^{\text{ef}}(\lambda_j, p_j, h_i) \right)^2} + \frac{\delta^2 \left(\beta_a^{\text{ef}}(\lambda_j, p_j, h_i) \right) \left(\hat{R}_j^{\text{ef}}(\lambda_j, p_j, h_i) - 1 \right)^2}{\left(\beta_a^{\text{ef}}(\lambda_j, p_j, h_i) \right)^2 \left(\hat{R}_j^{\text{ef}}(\lambda_j, p_j, h_i) \right)^2} \quad (\text{B4})$$

The International Standard Atmosphere ISO 2533 and seasonal latitudinal changed model CIRA (Committee on Space Research (COSPAR), 2012; Fleming et al., 1988), as well as measurements by radiosondes are applied in LIRIC for the calculation of molecular optical parameters. The relative variance of calculated molecular backscatter coefficient

$$\alpha_1^2 = \frac{\delta^2 \left(\beta_r^{\text{ef}}(\lambda_j, p_j, h_i) \right)}{\left(\beta_r^{\text{ef}}(\lambda_j, p_j, h_i) \right)^2} \quad (\text{B5})$$

is assumed to be a constant and its value can be reduced to $\alpha_1 = 0.01$ (e.g. Russell et al., 1979) if data of coordinated radiosonde measurements is available.

The aerosol backscatter coefficients, $\beta_a^{\text{ef}}(\lambda_j, p_j, h_i)$, are estimated by using Eqs. (10)–(17). Uncertainties of $\beta_a^{\text{ef}}(\lambda_j, p_j, h_i)$ basically follow from estimation errors of the coefficient $b(v, j, p_j)$ in Eqs. (15)–(17) that can be written by equation:

$$b(j, p_j, k) = \frac{1}{\vartheta_{k,p}^j} \frac{E_k^*(\lambda_j)}{\hat{C}_k^V} \quad (\text{B6})$$

LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



where

$$\frac{1}{\vartheta_{k,p}^j} = \frac{1}{4\pi} \omega_k(\lambda_j) A_{k,p}^j, \quad (\text{B7})$$

$$A_{k,p}^j = \begin{cases} P_{1,1}^V(\lambda_j, \gamma = 180^\circ) & \text{if } \rho_j = 1 \\ \frac{P_{1,1}^k(\lambda_j, \gamma = 180^\circ) - P_{2,2}^k(\lambda_j, \gamma = 180^\circ)}{2} & \rho_j = 2 \\ \frac{P_{1,1}^k(\lambda_j, \gamma = 180^\circ) + P_{2,2}^k(\lambda_j, \gamma = 180^\circ)}{2} & \rho_j = 3 \end{cases} \quad (\text{B8})$$

Parameter $\vartheta_{k,p}^j = \sigma_a^k(\lambda_j, h_j) / \beta_a^{\text{ef},k}(\lambda_j, h_j)$ in Eq. (B8) is the extinction-to-backscatter ratio or “lidar ratio” of the k -aerosol mode.

Parameters $\vartheta_{k,p}^j$ are retrieved from the data of radiometric direct Sun and almucantar measurements that are usually performed with the maximum scattering angle less than 150° . The range of the scattering angles decreases with decreasing the sun zenith angle. Retrieval of optical parameters in the backscatter direction, in a certain sense, is an extrapolation procedure out of the measured range with possible increasing of estimation uncertainties. One assumes that the errors of the estimation of $\vartheta_{k,p}^j$ are the main reason of the incorrect calculation of backscatter coefficients $\beta_a^{\text{ef}}(\lambda_j, \rho_j, h_j)$ and introduces parameter α_2 for characterization of the standard deviation of coefficients $1/\vartheta_{k,p}^j$ in LRS measurements.

Thus, Eq. (B4) is transformed to

$$\frac{\delta^2 \left(\beta_a^{\text{ef}}(\lambda_j, \rho_j, h_j) \right)}{\left(\beta_a^{\text{ef}}(\lambda_j, \rho_j, h_j) \right)^2} = \frac{\alpha_1^2}{\left(R_j^{\text{ef}}(\lambda_j, \rho_j, h_j) \right)^2} + \alpha_2^2 \frac{\left(R_j^{\text{ef}}(\lambda_j, \rho_j, h_j) - 1 \right)^2}{\left(R_j^{\text{ef}}(\lambda_j, \rho_j, h_j) \right)^2}. \quad (\text{B9})$$

The backscatter ratio $R_j^{\text{ef}}(\lambda_j, \rho_j, h_j)$ in Eq. (B9) under assumption $\mu = 0$ is approximately calculated at the pre-processing stage using the Klett algorithm (Klett, 1981).

LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Basically, the variance of aerosol optical thickness, $\delta^2(\tau_a(\lambda_j, h_i, h_{\text{ref}}))$, arises from altitude variations of aerosol modes that are not assumed by the aerosol model. Relative error of $\tau_a(\lambda_j, h_i, h_{\text{ref}})$ is zero at the reference point and is equal to α_3^2 at the start point h_1 , where α_3 is close to the error of AOT calculation from radiometer measurements.

Thus, the following approximation is used in the LIRIC algorithm:

$$\delta^2(\tau_a(\lambda_j, h_i, h_{\text{ref}})) = \alpha_3^2 \tau_a^2(\lambda_j, h_i, h_{\text{ref}}). \quad (\text{B10})$$

Term $\delta^2(\tau_r(\lambda_j, h_i, h_{\text{ref}}))$ in Eq. (B3) denotes the variance of molecular optical thickness of the atmospheric layer (h_n, h_{ref}). Only long scale or systematic deviations of molecular density contribute to the variance $\delta^2(\tau_r(\lambda_j, h_i, h_N))$. Similar to Eq. (B10),

$$\delta^2(\tau_r(\lambda_j, h_i, h_{\text{ref}})) = \alpha_4^2 \tau_r^2(\lambda_j, h_i, h_{\text{ref}}). \quad (\text{B11})$$

Measurement errors

Optical signals, detected by the lidar data acquisition system consists of backscatter $P_{j,i}^*$ and background B_j^* components. A suitable algorithm for estimating the measurement errors is described by Slesar' et al. (2013) and Slesar' et al. (2015). Regardless of the type of the photo-receiving sensor, three factors determine the measurement errors:

- non-linearity of the recording channel, which consist of nonlinearity of the photodetector and electronic units;
- “nonsynchronous” noise (non-correlated with the sounding pulse);
- “synchronized” noise (correlated with the sounding pulse).

Non-linearity of a receiving channel basically originates from saturation of an output signal at high incident light because of photo-sensor or electronic unit limitations. Likewise, deviations of an amplifier gain cause linear distortions of the detecting signal within the working range of photo-receiving module.

LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Basic difference between two types of noise is that “nonsynchronous” noise can be reduced by accumulation of input signals or by decreasing frequency bandwidth of the receiving channel, while this method is ineffective for “synchronized” noise. The main type of the “nonsynchronous” noise is the Schottky noise. “Synchronous” noise is basically caused by the interference of the electrical impulses from the laser power supply, synchronous with the sounding optical pulse. It is predominantly a low-frequency noise, and acceptable limitation of the frequency band of the photo-receiving channel does not lead to its decline.

We assume that the accumulation of the receiving lidar signal with A sounding pulses and averaging of the lidar signal over $2M + 1$ bins are carried out at the measurement and pre-processing stages.

Summing up the contributions of the noise components, one can write the following expression for the variances of the receiving analog and photon-counting signals (Slesar et al., 2013, 2015):

– for analog channel:

$$\frac{\delta^2(P_{j,i}^*)}{(P_{j,i}^*)^2} = \omega_j^2 \frac{(P_{j,n}^* + B_j^*)^2}{(P_{j,n}^*)^2} + \frac{(G_j^*)^2 + q_j^2 (P_{j,n}^* + B_j^*)}{A(2M + 1)(P_{j,n}^*)^2} + \frac{(U_j^*)^2}{(P_{j,n}^*)^2}, \quad (\text{B12})$$

where ω is the coefficient of nonlinearity, G_j^* is the amplitude of electrical noise, q_j^2 is the coefficient characterizing the power of the Schottky noise, U_j^* is the amplitude of “synchronized” noise;

– for counting channel:

$$\frac{\delta^2(N_{j,i}^*)}{(N_{j,i}^*)^2} = \omega_j^2 \frac{(N_{j,i}^* + N_{j,B})^2}{(N_{j,i}^*)^2} + \frac{(N_{j,G}^*)^2 + N_{j,i}^* + N_{j,B}}{A(2M + 1)(N_{j,i}^*)^2} + \frac{(N_{j,U}^*)^2}{(N_{j,i}^*)^2}, \quad (\text{B13})$$

where $N_{j,i}^*$ is the detected lidar signal, $N_{j,B}$ is the background signal, $N_{j,G}$ is the external “nonsynchronous” noise, and $N_{j,U}^*$ is “synchronized” noise.

Parameters ω_j , G_j^* , $N_{j,G}^*$, q_j , U_j^* , $N_{j,G}^*$, $N_{j,U}^*$ for specific photo-receiving module can be evaluated on a dedicated test bench by means of special calibration procedures (Slesar et al., 2013, 2015).

Appendix C: Details of inversion procedure

One can understand intuitively that optical parameters of aerosol modes, which constitute the aerosol model (see Sect.1.1), should be different to allow retrieving aerosol mode concentrations by means of algorithm described in Sect. 4. More corrected definition of this requirement is deficiency of linear relation between the sets of coefficients $\{\mathbf{a}_k, \mathbf{b}_k\}$ which define optical characteristics of k th aerosol mode. This conclusion results from the linear approximation of the Eq. (3). It means that we seek the solution, $\mathbf{c}(h_i)$, only from data of multi-wavelength lidar sounding. The linear least squares solution of Eq. (2) can be written as

$$\mathbf{c} = \left(\mathbf{K}_L^T \mathbf{\Omega}_L^{-1} \mathbf{K}_L \right)^{-1} \mathbf{K}_L^T \mathbf{\Omega}_L^{-1} \mathbf{L}^*, \quad (\text{C1})$$

where \mathbf{K}_L is the Jacobi matrix of the first partial derivatives $\{\mathbf{K}_L\}_{x,y} = \delta L_x / \delta c_y \Big|_c$. The following definitions are used in Eq. (C1) for measured vector, $\hat{\mathbf{L}}^*$ and state vector, \mathbf{c} ,

with dimensions $J \times 1$ and $K \times 1$, correspondently:

$$L_{j,i}^* = \begin{cases} L_1^*(h_1) \\ L_2^*(h_1) \\ \dots \\ L_j^*(h_i) \\ \dots \\ L_J^*(h_j) \end{cases}, \quad c_{k,i} = \begin{cases} c_1(h_1) \\ c_2(h_1) \\ \dots \\ c_k(h_i) \\ \dots \\ c_K(h_j) \end{cases}. \quad (C2)$$

The formula (C1) is valid if $\det(\mathbf{U}_L = \mathbf{K}_L^T \mathbf{\Omega}_L^{-1} \mathbf{K}_L) \neq 0$.

We use additional requirements that optical thickness of the aerosol layer is small, and the variances of the measured errors, $(\varepsilon_{L,1}^2, \varepsilon_{L,i}^2, \varepsilon_{L,l}^2)$, do not depend on h_i . So the matrix $\mathbf{U}_L = \mathbf{K}_L^T \mathbf{\Omega}_L^{-1} \mathbf{K}_L$ with dimensions $K \times K$ takes the block-diagonal form

$$U_L = \begin{bmatrix} \frac{1}{\varepsilon_1^2} \mathbf{U}_{k,k^*} & \dots & & & \\ & & \dots & & \\ & & & \frac{1}{\varepsilon_i^2} \mathbf{U}_{k,k^*} & \dots \\ & & & & \dots \\ & & & & \dots & \frac{1}{\varepsilon_l^2} \mathbf{U}_{k,k^*} \end{bmatrix}, \quad (C3)$$

where matrix \mathbf{U}_{k,k^*} , ($k \in 1, \dots, K$), does not depend on the superscript i . For 3-mode aerosol model ($K = 3$) and 4-channel lidar measurements ($J = 4$) matrix \mathbf{U}_{k,k^*} can be

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

written

$$\mathbf{U}_{k,k^*} = \begin{bmatrix} \sum_j b^2[j, p_j, 1] & \sum_j b[j, p_j, 1]b[j, p_j, 2] & \sum_j b[j, p_j, 1]b[j, p_j, 3] \\ \sum_j b[j, p_j, 1]b[j, p_j, 2] & \sum_j b^2[j, p_j, 2] & \sum_j b[j, p_j, 2]b[j, p_j, 3] \\ \sum_j b[j, p_j, 1]b[j, p_j, 3] & \sum_j b[j, p_j, 2]b[j, p_j, 3] & \sum_j b^2[j, p_j, 3] \end{bmatrix}. \quad (C4)$$

Thus, results of the retrieval depend on the specifics of matrix \mathbf{U}_{k,k^*} . The well-conditioned matrix \mathbf{U}_{k,k^*} provides suitable solution of Eq. (3). On the analogy with Veselovsky et al. (2005) the eigenvalue decomposition technique has been used to evaluate the “condition number” of matrix \mathbf{U}_{k,k^*} .

$$\text{Cond} \{ \mathbf{U}_{k,k^*} \} = |\psi_{\max}| / |\psi_{\min}| \quad (C5)$$

where ψ_{\max} and ψ_{\min} are the maximum and minimum eigenvalues of matrix $\mathbf{U}_{k,k}$, respectively. Parameter $\sqrt{\text{Cond} \{ \mathbf{U}_{k,k^*} \}}$ is a coefficient of increasing relative error of $c_{k,i}$ as compared to the relative error of $L_{n,j}^*$ estimation (Trefethen and Bau, 1997).

The data of radiometric measurements in Minsk during 2002–2010 were used to calculate the parameters $\text{Cond} \{ \mathbf{U}_{k,k^*} \}$ for the aerosol models with 2 and 3 aerosol fractions (3 and 4 measuring channels, correspondently). The Cumulative Distribution Functions (CDF) of parameter $\text{Cond} \{ \mathbf{U}_{k,k^*} \}$ is shown in Fig. C1.

Matrix $\mathbf{U}_{k,k}$ is sufficiently well conditional for the two-fraction aerosol model, and solution (Eq. C1) is applicable for the calculation of aerosol mode concentrations. In the case of the three-fractional aerosol model, parameters $\text{Cond} \{ \mathbf{U}_{k,k^*} \}$ increase approximately by 10, and the matrix \mathbf{U}_{k,k^*} becomes ill-conditioned. In such case we have to involve the Eq. (18) in retrieving procedure, i.e. to use information on the vector $\hat{\mathbf{C}}^{*V}$

from radiometric measurements. With our definitions, matrix \mathbf{H} in Eq. (18) is written as

$$\mathbf{H}_{K \times IK} = \begin{bmatrix} \Delta h_1 & \dots & 0 & \Delta h_2 & \dots & 0 & \dots & \Delta h_l & \dots & 0 \\ 0 & \Delta h_1 & 0 & 0 & \dots & 0 & \dots & 0 & \dots & 0 \\ \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\ 0 & \dots & \Delta h_1 & 0 & \dots & \Delta h_2 & \dots & 0 & \dots & \Delta h_l \end{bmatrix} \quad (\text{C6})$$

Finally, a priori smoothness restrictions are used as the additional factor for regularizing “ill-posed” problem solution.

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**LIRIC for the retrieval
of vertical aerosol
properties from
combined
lidar/radiometer data**

A. Chaikovsky et al.

Title Page	
Abstract	Introduction
Conclusions	References
Tables	Figures
⏪	⏩
◀	▶
Back	Close
Full Screen / Esc	
Printer-friendly Version	
Interactive Discussion	



LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Ansmann, A., Tesche, M., Groß, S., Freudenthaler, V., Seifert, P., Hiebsch, A., Schmidt, J., Wandinger, U., Mattis, I., Müller, D., and Wiegner, M.: The 16 April 2010 major volcanic ash plume over central Europe: EARLINET lidar and AERONET photometer observations at Leipzig and Munich, Germany, *Geophys. Res. Lett.*, 37, L13810, doi:10.1029/2010GL043809, 2010.

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LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

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LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

[Title Page](#)[Abstract](#)[Introduction](#)[Conclusions](#)[References](#)[Tables](#)[Figures](#)[◀](#)[▶](#)[◀](#)[▶](#)[Back](#)[Close](#)[Full Screen / Esc](#)[Printer-friendly Version](#)[Interactive Discussion](#)

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LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



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**LIRIC for the retrieval
of vertical aerosol
properties from
combined
lidar/radiometer data**

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



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**LIRIC for the retrieval
of vertical aerosol
properties from
combined
lidar/radiometer data**

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

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**LIRIC for the retrieval
of vertical aerosol
properties from
combined
lidar/radiometer data**

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

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**LIRIC for the retrieval
of vertical aerosol
properties from
combined
lidar/radiometer data**

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



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LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



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LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



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AMTD

8, 12759–12822, 2015

LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

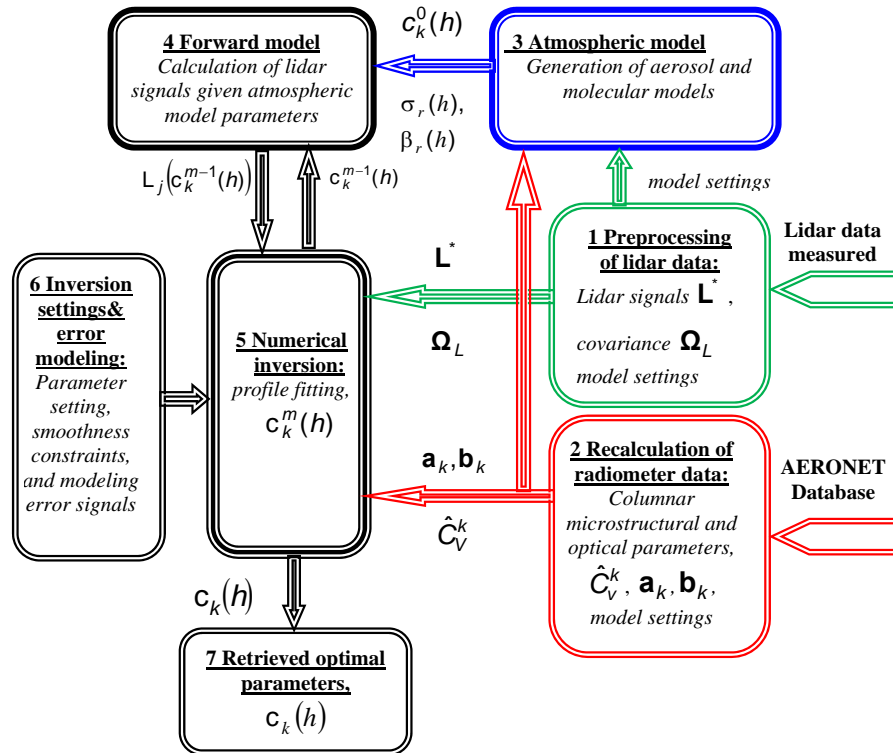


Figure 1. Flowchart of LIRIC algorithm. Details are in Sect. 2.2.

Title Page

Abstract Introduction

Conclusions References

Tables Figures

◀ ▶

◀ ▶

Back Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

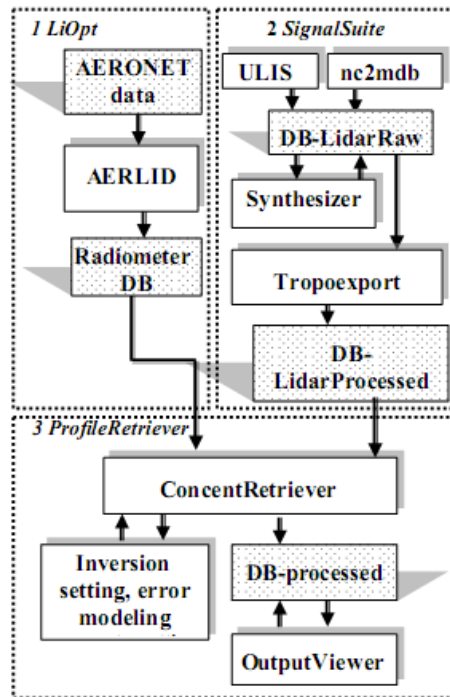


Figure 2. Flowchart of the program package.

Title Page	
Abstract	Introduction
Conclusions	References
Tables	Figures
◀	▶
◀	▶
Back	Close
Full Screen / Esc	
Printer-friendly Version	
Interactive Discussion	



LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

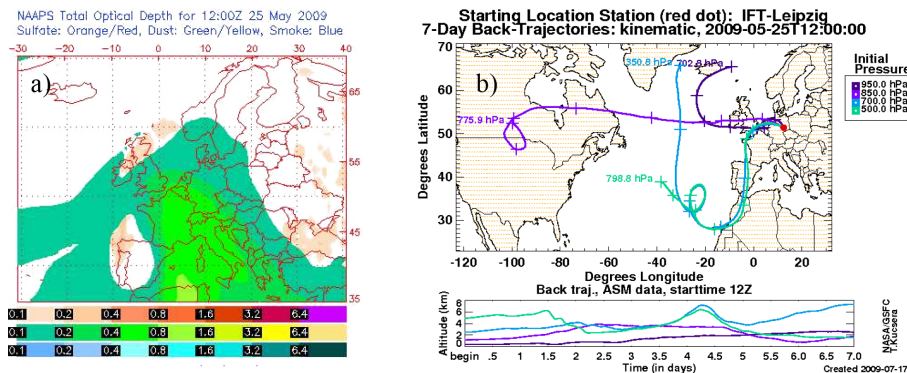


Figure 3. EARLI09 intercomparison experiment: **(a)** NAAPS Total Optical Depth forecast, 25 May 2009 at 12:00 UTC; **(b)** 7 day backtrajectories ending over Leipzig, Germany at 12:00 UTC on 25 May 2009.

LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

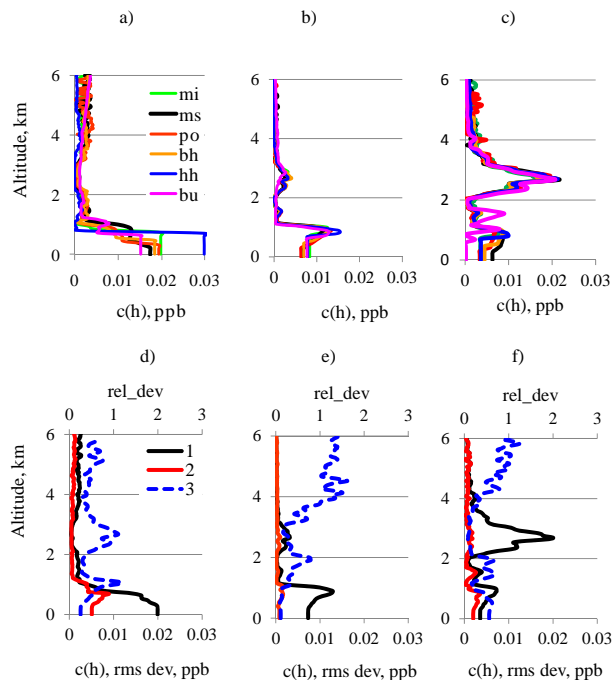


Figure 4. PVC profiles, $c_k(h)$, and estimated deviations retrieved from data of EARLI09 inter-comparison campaign, 10:20–11:40 UTC, 25 May 2009, Leipzig, Germany, measured in Leipzig by six EARLINET lidars: mi – Minsk, ms – München, po – Potenza, bh – Bilthoven, hh – Hamburg, bu – Bucharest; **(a, d)** – fine, **(b, e)** – coarse spherical; **(c, f)** – coarse non-spherical; 1 – average PVC profile, 2 – rms-deviation (rms_dev), 3 – relative deviation (rel_dev). Measured data from four lidar channels (355, 532-parallel, 532-cross, 1064 nm) and three-mode aerosol model were used.

LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

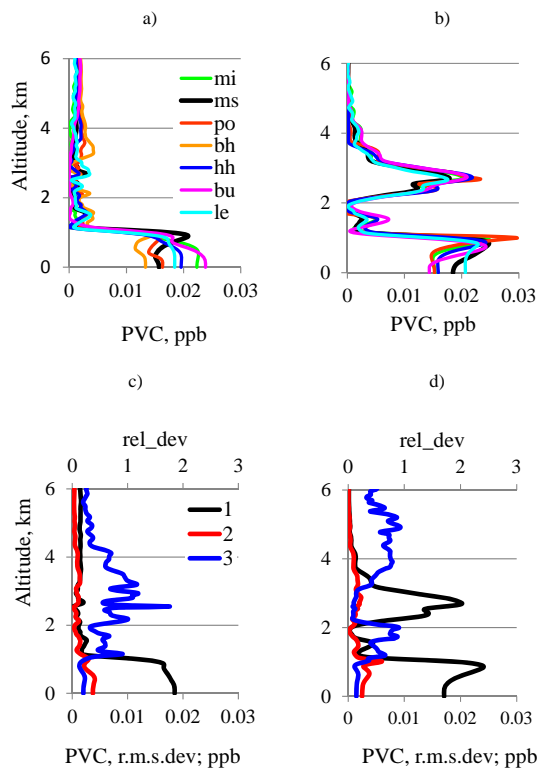


Figure 5. Same as Fig. 4 except for data from three lidar channels (355, 532 – intensity/parallel polarized component, and 1064 nm) and two mode aerosol model were used. Label “le” stands for lidar “PollyXT” of TROPOS, Leipzig.

LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

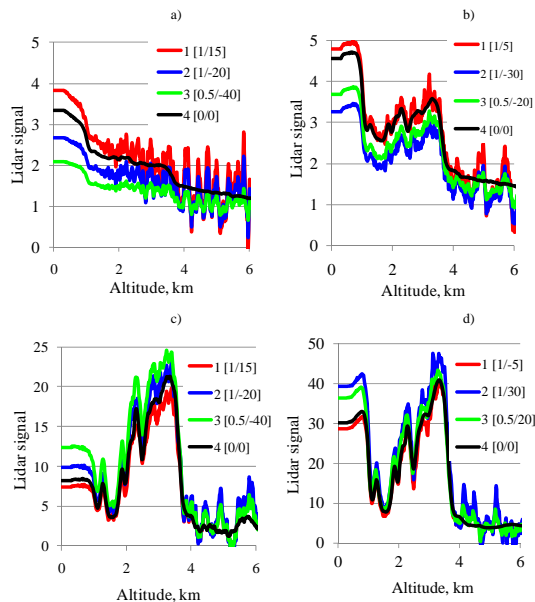


Figure 6. Range-corrected normalized lidar signals, L^* , corrupted with noise and amplitude distortions. Original data are provided by the München lidar team in the frame of EARLI09 intercomparison campaign, 14:30–15:30 UTC, 25 May 2009, Leipzig, Germany: **(a)** – 355 nm, **(b)** 1064 nm, **(c)** – 532 nm, parallel polarized, **(d)** – 532 nm, cross polarized; 4 – original signal, $1 \div 3$ – corrupted signals. In square brackets distortion parameters α_j/Δ_j are given.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

**LIRIC for the retrieval
of vertical aerosol
properties from
combined
lidar/radiometer data**

A. Chaikovsky et al.

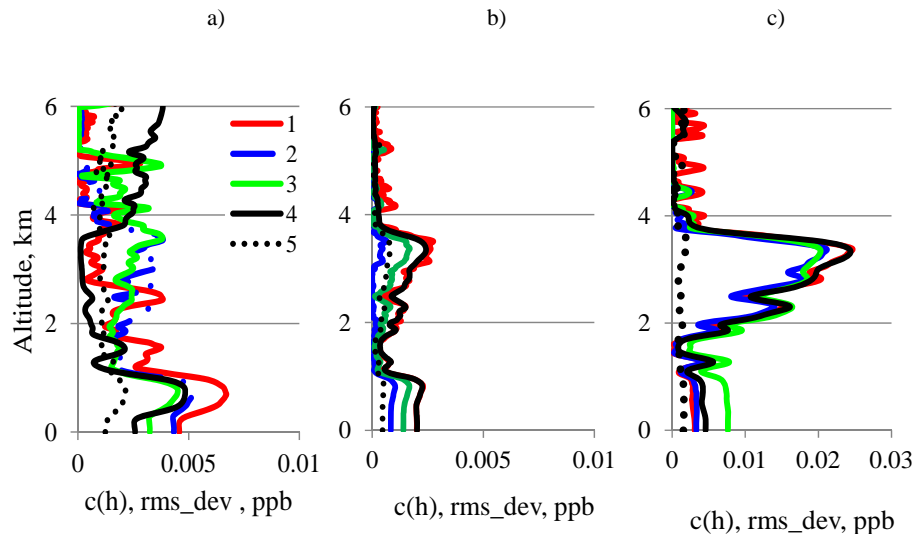


Figure 7. PVC profiles, $c_k(h)$, and their rms-deviations retrieved in response to disturbed data from of the München lidar, EARLI09 intercomparison campaign, 14:30–15:30 UTC, 25 May 2009, Leipzig, Germany: **(a)** – fine, **(b)** – coarse spherical, **(c)** – coarse non-spherical modes; 4 – for the original signal, 1–3 – for disturbed signals; 5 – rms-deviation.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

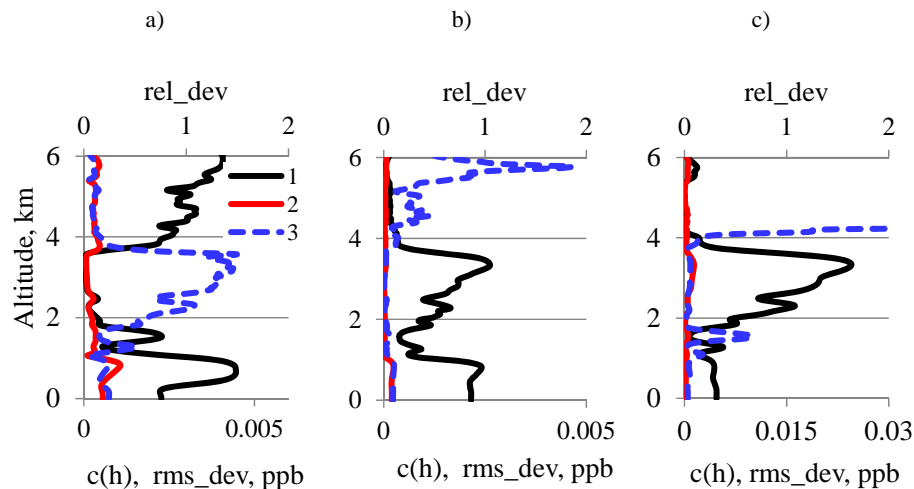


Figure 8. Variations of PVC profiles, $c_k(h)$, retrieved with 20% uncertainties in the aerosol lidar ratios; data of München lidar, EARLI09 intercomparison campaign, 14:30–15:30 UTC, 25 May 2009, Leipzig, Germany are used; **(a)** fine, **(b)** coarse spherical, **(c)** coarse non-spherical modes; 1–average value, 2 rms-deviation, 3 relative deviation.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

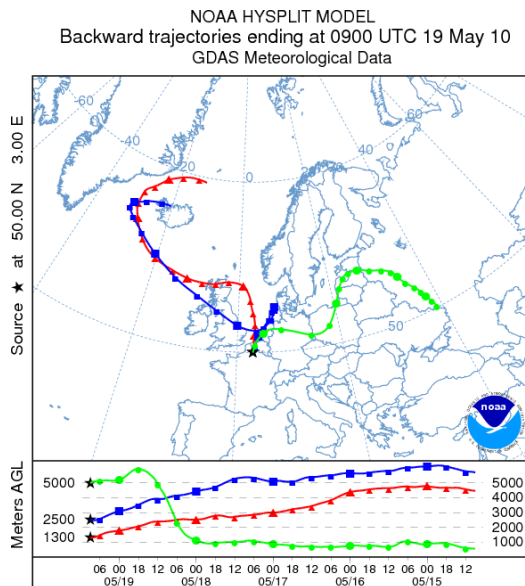


Figure 9. Air-mass back-trajectories for Lille at 08:00 UTC, 19 May 2010, (NOAA HYSPLIT model).

Title Page	
Abstract	Introduction
Conclusions	References
Tables	Figures
◀	▶
◀	▶
Back	Close
Full Screen / Esc	
Printer-friendly Version	
Interactive Discussion	



LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

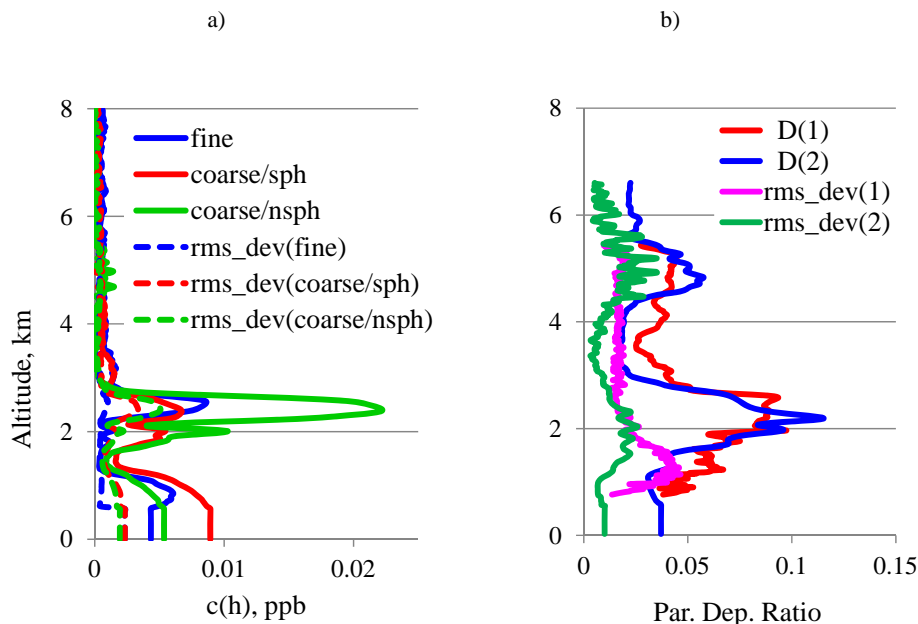


Figure 10. (a), PVC profiles, $c_k(h)$, of the fine, course-spherical (course/sph) and course-nonspherical (coarse/nsph) aerosol modes, and their rms-deviations (rms_dev(fine), rms_dev(coarse/sph), and rms_dev(coarsw/nsph)); (b), particle depolarization ratio, D(1) and (2), and their rms-deviations, rms_dev(1) and rms_dev(2). Profiles were retrieved from the data measured in Lille, 19 May, 2010, 09:17–09:58 UTC. Profiles D(1) and rms_dev(1) are the results of the direct calculation of depolarization ratio and their rms-deviations from lidar measurements, as well as D(2) and rms_dev(2) were calculated from retrieved aerosol mode concentrations, $c_k(h)$.

LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

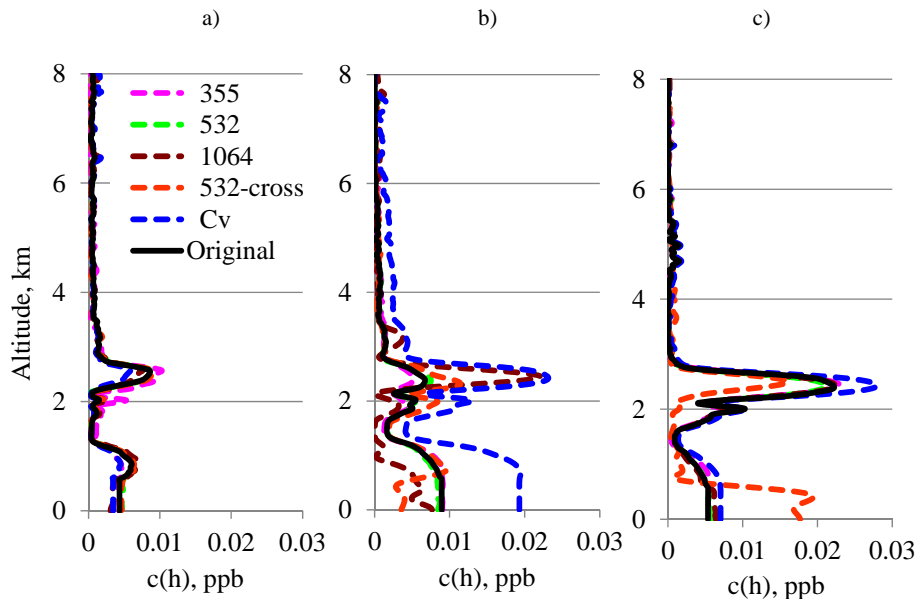


Figure 11. Variation of aerosol concentration profiles, $c_k(h)$, for fine **(a)**, coarse spherical **(b)** and coarse non-spherical **(c)** aerosol modes in response to elimination of different parts of input information. Tag “Original” denotes complete set of input data; tag “355” (or 532, 1064, 532-cross) denotes that lidar signal at 355 nm (or 532, 1064, 532-cross) wavelength is excluded; tag “ C^V ” denotes that columnar volume concentrations of aerosol modes are excluded. Lille, 08:00 UTC, 19 May 2010.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

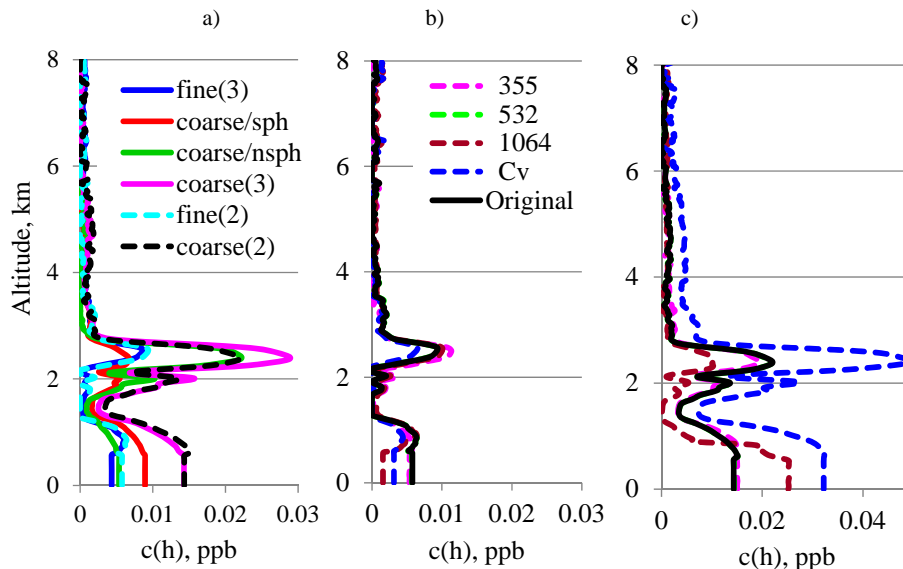


Figure 12. Comparison of PVC profiles, $c_k(h)$, for the two- and three-mode aerosol models (a), and variations of concentration profiles, $c_k(h)$, for fine (b) and coarse (c) aerosol modes of the two-mode aerosol model in response to elimination of different parts of input information. In Fig. 12a tags “fine(3)” and “coarse(3)” denote fine and coarse modes of two-mode aerosol model. Tags “fine(4)”, “coarse/sph”, “coarse/nsph” and “coarse(4)” denote fine, coarse spherical, coarse non-spherical and total coarse mode of three-mode aerosol model, correspondently. In Fig. 12b and c tag “Original” means complete set of input data; tag “355” (or 532, 1064) denotes that the lidar signal at 355 nm (or 532, 1064) wavelength is excluded; tag “ C^V ” denotes that columnar volume concentrations of aerosol modes are excluded. Lille, 08:00 UTC, 19 May 2010.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Figure 13. Map of the EARLINET stations (red dots). Green dots indicate the stations where LIRIC program package has been implemented.

AMTD

8, 12759–12822, 2015

LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

Title Page	
Abstract	Introduction
Conclusions	References
Tables	Figures
◀	▶
◀	▶
Back	Close
Full Screen / Esc	
Printer-friendly Version	
Interactive Discussion	



LIRIC for the retrieval of vertical aerosol properties from combined lidar/radiometer data

A. Chaikovsky et al.

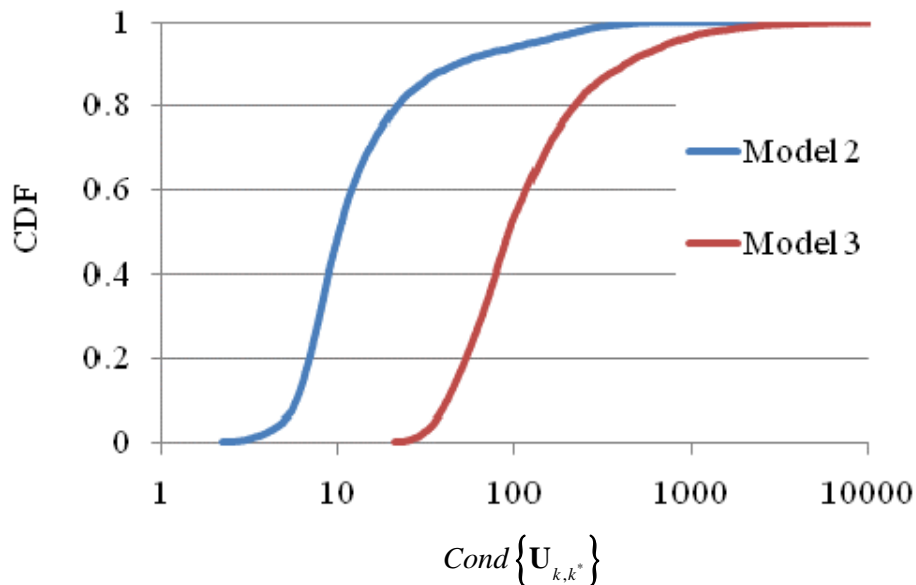


Figure C1. Cumulative Distribution Functions (CDF) of parameter $Cond \{U_{k,k^*}\}$ calculated from radiometer data of the AERONET station in Minsk for two- and three-fraction aerosol models, Model 2 and Model 3, respectively.

[Title Page](#)[Abstract](#)[Introduction](#)[Conclusions](#)[References](#)[Tables](#)[Figures](#)[◀](#)[▶](#)[◀](#)[▶](#)[Back](#)[Close](#)[Full Screen / Esc](#)[Printer-friendly Version](#)[Interactive Discussion](#)