



## Abstract

Measurements of land-surface emission rates of greenhouse and other gases at large spatial scales (10 000 m<sup>2</sup>) are needed to assess the spatial distribution of emissions. This can be more readily done using spatial-integrating micro-meteorological methods than the widely-utilized small chamber measurements. Several micro-meteorological flux-gradient methods utilizing a non-intrusive path-averaging measurement method were evaluated for determining land-surface emission rates of trace gases under stable boundary layers. Successful application of a flux-gradient method requires confidence in the gradients of trace gas concentration and wind and in the applicability of boundary-layer turbulence theory. While there is relatively high confidence in flux measurements made under unstable atmospheres with mean winds greater than 1 m s<sup>-1</sup>, there is greater uncertainty in flux measurements made under free convective or stable conditions. The study involved quality-assured determinations of fluxes under low wind, stable or night-time atmospheric conditions when the continuous “steady-state” turbulence of the surface boundary layer breaks down and the layer has intermittent turbulence. Results indicate that the Monin-Obukhov similarity theory (MOST) flux-gradient methods that assume a log-linear profile of the wind speed and concentration gradient incorrectly determine vertical profiles and thus fluxes in the stable boundary layer.

## 1 Introduction

The determination of greenhouse gas (GHG) emission rates from open land is important to understand our climate and current and future changes in climate. Nitrous oxide (N<sub>2</sub>O) has one of the longest lifetimes of the GHG (114 yr). Its efficiency in absorbing infrared radiation emitted by the Earth’s surface is 310 times larger compared to the radiative efficiency of the same amount of CO<sub>2</sub> (Forster et al., 2007). Consequently, a complete understanding of the atmospheric concentration variations of this

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GHG is necessary. Even though man-made emissions of N<sub>2</sub>O, mainly from agriculture (e.g. use of fertilizers) and industry (e.g. combustion processes), are known, the quantified source and sink fluxes of N<sub>2</sub>O by different soils are not well known. Furthermore, there is high spatial variability of GHG fluxes from soils of the same type and under identical management (Smith et al., 1994; Laville et al., 1999; Denmead et al., 2000; Turner et al., 2008).

A better knowledge of the temporal and spatial variations of the trace gas fluxes from open land requires an up-scaling of source and sink measurements near the surface (Bange et al., 2006; Schmid et al., 2006; Denman et al., 2007; Forster et al., 2007; Chapuis-Lardy et al., 2007; Holst et al., 2007). Such variability then needs to be expressed in emission inventories and processes parameterized for use in chemistry-transport models at grid scales of about 1 000 000 m<sup>2</sup> (Ganzeveld et al., 2007). In addition, measurements are needed to verify satellite measurements of Earth surface fluxes at their pixel scales (e.g. MODIS about 1 km<sup>2</sup> or CO<sub>2</sub> and CH<sub>4</sub> from GOSAT about 100 km<sup>2</sup>). If the up-scaling of point, line and area data to grid or pixel scale is not correct the resultant data can include tremendous errors.

Up to now, a large amount of trace gas emission measurements have been done by time-dependent accumulation of gasses in small (less than 1 m<sup>2</sup>) but variously-sized chambers (see e.g. Butterbach-Bahl et al., 2004). A possible larger-scale measurement technique of this is the time-dependent accumulation of trace gases in a measuring tunnel with open-path optical methods like FTIR (Fourier Transform InfraRed spectrometry), DOAS (Differential Optical Absorption Spectroscopy) or TDLAS (Tunable Diode Laser Absorption Spectroscopy) along the longitudinal tunnel axis, which yields path-integrated concentration (PIC) values (see e.g. Galle et al., 1994). Measurements of fluxes at larger spatial scales (10 000 m<sup>2</sup>) require the application of micro-meteorological methods. In general, micro-meteorological emission or flux measurement methods have limited applicability during stable atmospheric conditions due to the low turbulence levels and non-linear interactions of different processes on several temporal and spatial scales (Fernando and Weil, 2010).

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Non-intrusive measurement methods like open-path optical techniques are well suited to study the up-take of trace gases by soils (Chapuis-Lardy et al., 2007) because they avoid any influences upon the soil emission processes by wind, temperature, humidity and radiation or heat fluxes. Ideally, path-integrated wind and turbulence measurements with identical control volumes to the concentration measurements would be used for flux determinations. Such path-integrated measures have been developed using laser-scintillometers (Nakaya et al., 2007) and acoustic travel-time tomography (Barth et al., 2007). Unfortunately, these measurement techniques possess wind speed detection limits of about  $0.2 \text{ m s}^{-1}$ , and thus are not sensitive enough to describe the turbulence during stable and very stable conditions (Andreas, 2000).

The objective of this paper is to present a method for areal-averaged emission rates of  $\text{N}_2\text{O}$  over flat grassland (Fuhrberg experiment) made during periods of low wind speed and stable boundary layers using open-path gas measurement techniques. The methods evaluated substitute integrated open-path concentration measurements for point measurements used in the usual flux-gradient methods (Wichink Kruit et al., 2007, 2010; Foken, 2008; Denmead, 2008; Luhar et al., 2009). The open path lengths should be bigger than 100 m which are spatially representative of 10 000 to 100 000  $\text{m}^2$  areas. Successful application of this method requires confidence in the gradients in trace gas concentration and wind and confidence in the applicability of boundary-layer turbulence theory and consequently the procedures to qualify measurements that can be used to determine the flux is critical.

## 2 Spatial and temporal averaging

The necessary step from measurements of concentrations and meteorological parameters to fluxes is determining the surface area of influence or source area of the sensor. In the case of quantities that are subject to turbulent diffusion, this depends on the flow and turbulence conditions. The “footprint function” describes the relationship between the spatial distribution of surface sources (or sinks) and the measured flux at height in

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the surface layer. Schmid (1994) described a numerical model which provided a user-friendly estimate of the surface area responsible for measured concentrations or fluxes at a point. It is the objective of the study to extend the surface area influencing the flux measurement by extending from point to path-integrated measurements.

5 However, existing footprint analysis are based on Monin-Obukhov similarity theory (MOST) they cannot be applied during stable and very stable boundary layer conditions common at night. During such stable conditions, the near-surface wind will cease even in cases with moderate winds in the free atmosphere because the stabilization of the surface layer decouples the surface layer and the free atmosphere above. Although  
10 no continuous turbulence production takes place and “steady-state” conditions cease and the layer can have intermittent and wave-like turbulent structures. While intermittent turbulence may be representable in path-integrated measurements and contribute to the flux under stable conditions, the non-transporting wave structure must be removed from the computation of fluxes. The quality-assurance procedures presented  
15 attempt to include the intermittent turbulence and exclude the wave structures from flux determinations.

### 3 Flux measurement methods

The focus of this study is on the non-intrusive determination of areal-averaged trace gas emission rates using open-path measurement techniques.

#### 3.1 Measurement campaigns

20 Two experiments were conducted: one at Fuhrberg, Germany and one at Indiana, United States of America. All measurement data were averaged over 30-min intervals.

### 3.1.1 N<sub>2</sub>O emission rates over flat grassland (Fuhrberg experiment)

The campaigns were conducted in the flat catchment area of the Fuhrberger Feld aquifer, 30 km northeast of the city of Hannover in Northern Germany, of 20–40 m thick sand. The soil type was classified as a Gleyic-Podzol with an Ap, Go-Bsh and a Go horizon (FAO-UNESCO, 1994). The mean precipitation rate in that area is 680 mm yr<sup>-1</sup> and the mean annual temperature is 8.9 °C. The land use was tillage until July 2005. At the time cultivated grass species *Festuca rubra* remained on the plot.

The concentration differences of N<sub>2</sub>O, CO<sub>2</sub>, CH<sub>4</sub> and H<sub>2</sub>O were measured by two bi-static open-path FTIR spectrometers K300 from Kayser-Threde (Haus et al., 1994): one at 2.70 m above ground level (a.g.l.),  $z^+ = 2.70$  m, and one at 0.5 m a.g.l.,  $z^- = 0.50$  m (Fig. 1). Both FTIR paths (installed with globars) were 98 m in length and parallel to each other. The lower measurement height was determined to be about 10 times the roughness length of the surface which is grassland. Three-dimensional (3-D) sonic anemometers (USA-1) from Metek GmbH measured the wind and turbulence at the same heights as the FTIR open paths. Corresponding to the upper height the fetch or footprint i.e. the required homogeneity in upwind direction was about 270 m. Considering the path lengths of about 100 m the investigated area was about 27 000 m<sup>2</sup>.

A comparison of both open-path FTIR instruments was performed (22 and 23 November 2007 as well as 10 and 11 July 2008) following each flux-gradient measurement campaign (15–17 October 2007 and 17–19 June 2008). An 1-h inter-comparison of the mean winds measured by the two 3-D sonic anemometers was also performed, but this was completed in the field during the campaigns. The comparison-measured bias between FTIR spectrometers was used to correct the measured differences for each measurement campaign. No bias corrections were needed for the 3-D sonic anemometers.

N<sub>2</sub>O measurements by four chambers (area of 0.045 m<sup>2</sup> each) at the same soil area on a plot of 1.5 m by 1.5 m were performed at the following day of the open-path flux-gradient measurements around 22:00 for the comparison with these measurements.

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N<sub>2</sub>O soil emission studies were performed at the experimental field from 2006 until 2008 (Deurer et al., 2008; Weymann et al., 2009; von der Heide et al., 2009).

### 3.1.2 NH<sub>3</sub> emission rates over cattle lagoons (Indiana experiment)

Gaseous emissions of NH<sub>3</sub> from a cattle lagoon (dairy waste lagoon) were determined from wind measurements made by 3-D sonic anemometers (R. M. Young sonic 8100) and open-path gas concentration measurements using TDLAS (GasFinder 2.0, Boreal Laser, Inc). These measurements were made as part of the National Air Emissions Monitoring Study (NAEMS) at representative swine, dairy, egg layer, and broiler production facilities in the United States of America. Measurements were made for a total of 321 days between September 2008 and August 2009.

Concentration gradients of NH<sub>3</sub> were determined from scanning TDLAS measurements. The TDLAS units were set up at opposite corners of the 85 m by 116 m lagoon with 16 m towers placed at the other two corners of the lagoon (see Fig. 2). Open-path NH<sub>3</sub> measurements were made along each adjacent side of the lagoon at 0.9 m a.g.l. along paths extending about 1/3rd and 2/3rds and slightly more than the full distance down the side of the source and at two vertical angles to retro-reflectors between 85 m and 116 m away and at approximately 7.6 m a.g.l. and 15.2 m a.g.l. Thus, each side of the source has three near-surface paths and two elevated paths. The scanner sequentially aimed the TDLAS at one of the ten retro-reflectors on the two adjacent sides of the source such that a complete set of measurements occurred every 100 s.

Quality control (QC) checks for the TDLAS measurements included checks for path obstruction every measurement, internal calibration checks every 100 measurements, spectral feature check and single point calibration verification approximately every 21 days, and multi-point calibrations every six months. Based on twice the precision, the TDLAS units had a minimum detection limit of 2 ppm m or less. The TDLAS calibration was verified at 50 ppm m using a 0.5 m calibration cell in an approximately 4 m path length in the trailer. Spectral feature checks were made using a digital oscilloscope.

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3-D sonic anemometers were mounted at 2.5 m ( $z^-$ ), 4.4 m ( $z_F$ ), and 16.2 m ( $z^+$ ) height (see Fig. 2). The wind speed and temperature gradient measurements were assured by frequent inter-comparisons of the three sensors in the field and three “reference” (unused except for calibration checks) anemometers. The comparison-measured bias between the temperature measurements of the sonic anemometers was used to correct the measured temperature gradients. No bias corrections were needed for the sonic anemometer wind measurements.

### 3.2 Open-path flux-gradient method

Following MOST, the flux  $F$  ( $\mu\text{g m}^{-2} \text{s}^{-1}$ ) at the height  $z_F$  (m) can be determined for steady-state turbulence of the surface boundary layer and height-constant conditions using the vertical gradient of the concentration  $c$  ( $\mu\text{g m}^{-3}$ ) and horizontal wind speed  $u$  ( $\text{m s}^{-1}$ ) approximated with the measurements at the two heights  $z^-$  (m) (below) and  $z^+$  (m) (above) and the friction velocity  $u_*$  ( $\text{m s}^{-1}$ ).  $u_*$  is estimated according to Lege (1981) as

$$u_* = \kappa \cdot \frac{u(z^+) - u(z^-)}{\ln(z^+/z^-) - \psi_m^+ + \psi_m^-}, \quad (1)$$

with  $\kappa$ : Karman’s constant,  $\kappa = 0.4$  and where  $\psi_m$  is the momentum stability correction function for the wind profile. It was assumed that the stability correction for trace gas  $c$  ( $\psi_c$ ) was identical to that for momentum (see e.g. Holtslag and de Bruin, 1988). The empirically-derived values are from Businger et al. (1971), Dyer (1974) and Hicks (1976) (termed BDH).  $\psi_c$  was estimated according to the suggested equations of Holtslag and de Bruin (1988). During the campaigns the stability correction function  $\psi_c$  showed the expected relationship to  $\zeta = (z - d)/L$  where  $L$  is the Monin-Obukov length and  $d$  is the displacement height with the exception that the values crossed the  $\zeta = 0$  line below the expected value of  $\psi_c = 0$  (Indiana experiment shown in Fig. 3). This suggests that either turbulence was advected into the profiles, variance in the motion of the boundary layer was associated with non-turbulent structures, and/or the wind



profile was not log-normal. Since the flow often has the lagoon upwind and the lagoon berm is both a roughness and height transition, we might expect convergence with turbulence regularly advected across the measurement planes. Although non-turbulent variance in motion associated with waves at high  $\zeta$  also contribute to deviations from the theoretical relationship, periods with such motion were largely excluded from the analysis using procedures described below.

A “parameterized” flux-gradient method can be defined by multiplying the  $u_*$  (Eq. 1) by the concentration scaling parameter  $c_*$  as

$$F(z_F) = -u_* \cdot c_* \quad (2)$$

where

$$c_* = \kappa \cdot \frac{c(z^+) - c(z^-)}{\ln(z^+/z^-) - \psi_c^+ + \psi_c^-} \quad (3)$$

Alternatively, measured values can replace parameterizations for  $c_*$  and  $u_*$ : if the wind and concentration measurements are at the same heights,  $\psi_c$  can also be based on the measured values of  $u_*$  by solving Eq. (1) for  $\psi_c$ , and if the wind and concentration measurements are at different heights,  $\psi_c$  can be parameterized while  $u_*$  is determined directly from the sonic anemometer measurements. This approach is termed “measured  $u_*$ ” flux-gradient method from here on.

The flux can also be calculated via the turbulent diffusivity as exchange coefficient. This more generic flux estimation method (termed “similarity” method) uses fewer assumptions than the “parameterized” or “measured  $u_*$ ” flux-gradient methods in that there is no assumed MOST correction for stability and no assumed neutral MOST logarithmic wind profile. The turbulent diffusivity  $K_c$  was calculated from the measured friction velocity  $u_*$  (mean of the measured friction velocities in the heights  $z^+$  and  $z^-$ , see also Stull, 1997) and the measured height difference of the horizontal wind speeds

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$$K_c = \frac{1}{Sc} \frac{u_*(z_F)^2}{\frac{u(z^+) - u(z^-)}{z^+ - z^-}}, \quad (4)$$

where  $Sc$  is the Schmidt number. The  $Sc$  is 1.079 for  $N_2O$  and 0.63 for  $NH_3$  given the molecular diffusion coefficient ( $\nu_c$ ) of  $1.436 \times 10^{-5} \text{ m}^2 \text{ s}^{-1}$  for  $N_2O$  and  $1.978 \times 10^{-5} \text{ m}^2 \text{ s}^{-1}$  for  $NH_3$  (Massman, 1998) and a kinematic viscosity of  $1.346 \times 10^{-5} \text{ m}^2 \text{ s}^{-1}$  (List, 1963) for a temperature of 298 K and barometric pressure of 1 atm. Then the flux is calculated as

$$F(z_F) = -K_c \cdot \frac{dc}{dz} = -K_c \cdot \frac{c(z^+) - c(z^-)}{z^+ - z^-}. \quad (5)$$

If there is no turbulent transport nor any coherent wind structure influencing the transport of the gas,  $K_c = \nu_c$  and the flux is described by the Fick's law.

### 3.3 Sonic anemometers and quality assurance checks

While 3-D sonic anemometer measurements were aligned to the natural wind coordinates for the Fuhrberg experiment conducted over a flat open field they were not aligned in the Indiana experiment because the measurements were made in close proximity to a lagoon with a 1–3 m drop in elevation. Turbulence statistics (variances and co-variances) were calculated for all heights except 0.5 m where it was expected that significant fractions of the turbulence would not be measureable by the sonic anemometer. The momentum flux and sensible heat flux were determined by the eddy-covariance method (Foken, 2008).  $L$  was determined with  $\theta_v'$  (perturbation in the virtual potential temperature), approximated by  $\theta_s'$  (perturbation in sonic anemometer temperature), and  $w'$  (perturbation of the vertical wind speed). Since there is often a strong self-correlation between the turbulence quantities used in the calculation of  $L$

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and the gradient Richardson number is vulnerable to significant errors when the vertical wind profile is complex, the bulk Richardson number  $Ri_B$  was used as an additional measure of overall stability (Mahrt, 2010).

The nature of the flow variability during each measurement period was evaluated since stable boundary layers may have periods of continuous turbulence, intermittent turbulence and waves. Tests used to evaluate the nature of the flow variability are summarized in Table 1 together with the acceptable ranges for the various parameters. The inter-comparisons during the Indiana experiment showed a mean difference between the field-used 3-D sonic anemometers and the reference anemometers of  $0.03 \text{ m s}^{-1}$  and up to 3 K. Temperature bias between sonic anemometers was corrected prior to determination of  $Ri_B$ . During the Fuhrberg experiment typically differences between sensors were less than  $0.1 \text{ m s}^{-1}$ , zero flow check resulted in a flow of less than  $\pm 0.05 \text{ m s}^{-1}$ .

Influences of non-turbulent motions and non-homogeneous, unsteady flow were partially excluded from the data set by restrictions on the acceptable values of  $Ri_B$  (Indiana experiment shown in Fig. 4). Distinct non-turbulent waves were evident in many of the most stable periods in the Indiana experiment: a  $Ri_B$  of less than 0.2 corresponded to boundary layers with distinct (although sometimes weak) turbulence, consistent with theory (Luhar et al., 2009). The removal of all measurement periods where  $Ri_B > 0.2$  (Fig. 4) reduced the range of  $\zeta$  to values less than 0.8 (Fig. 3) and the quality-assured dataset to 358 30-min measurement periods. Even with this exclusion, the deviation of  $\psi_c$  from the theoretical relationship with  $\zeta$  (Fig. 3) implies non-stationarity of the flow associated with waves, meandering flow and other coherent flow structures (Mahrt, 2007). Consequently, we also then expected differences between the diffusivity-based flux calculation and the MOST-theory based flux-gradient calculations since non-stationarity often results in more linear (not logarithmic) wind profiles (Mahrt, 2007).

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### 3.4 FTIR measurement campaigns and quality assurance checks

FTIR spectrometry (Griffith, 2002) is a factor of 10 more sensitive than commercial laser techniques for N<sub>2</sub>O measurements and was used in these campaigns. The spectral resolution of the FTIR measurements was 0.2 cm<sup>-1</sup> allowing the determination of interferences in the spectra and thus a high accuracy of the concentration measurements (Sussmann and Schäfer, 1997). The telescope of the spectrometers for the detection of the IR radiation of the radiation source (globar) had a diameter of 15 cm and a field of view of 3 mrad. The measurement time was about 5 min. The measurements were running continuously interrupted by short time periods for internal alignment procedures of the spectrometers (20 up to 60 min) every 3 up to 6 h. The minimum detection limit (MDL) of the K300 spectrometers at a 100 m open path was 5 ppb (10 µg m<sup>-3</sup>) for N<sub>2</sub>O. The statistical concentration measurement error was less than 3 % (Haus et al., 1994) and was calculated from the signal-to-noise-distance in the measured spectra for each measurement. This detection limit and measurement precision were appropriate to detect concentration differences of more than 1 %. Griffith et al. (2002) and Esler et al. (2000) who used FTIR for micro-meteorological techniques as closed-path measurement method found similar values. The temporal variations of the measured greenhouse gas concentrations at both heights followed the expected daily courses. The data quality filtering criteria are given in Table 1.

The MDL of each flux-gradient methodology was evaluated assuming that the measurement errors are independent (Taylor, 1982). An analysis of the propagation of measurement errors for the “parameterized” flux-gradient method indicated a MDL of approximately 5 %. An analysis of the propagation of measurement errors for the “similarity” flux-gradient method was approximately 7 % with a K<sub>c</sub> MDL of approximately 0.01 m<sup>2</sup> s<sup>-1</sup>. It is important to understand that these errors are only the propagation of measurement errors through the equations and not associated with the ability of a given method to describe the true flux given the variability in the environment.

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### 3.5 TDLAS measurement campaign and quality assurance checks

Ammonia TDLAS are robust and were used to measure  $\text{NH}_3$  in this campaign. The TDLAS  $\text{NH}_3$  MDL was determined to be approximately 30 ppb ( $13.9 \mu\text{g m}^{-3}$ ) for the paths measured in the campaign. The statistical concentration measurement error across the measurement domain was measured at less than 5% based on calibration verifications. Measurements were running nearly continuously for 321 days. Emissions of  $\text{NH}_3$  during the Indiana experiment were determined at 30-min intervals from wind profiles based on the three anemometers and concentration profiles based on the multiple TDLAS-measured PIC using the vertical radial plume mapping method (Hashmonay et al., 2001). This method utilizes the product of an interpolated mass concentration distribution and interpolated wind speed normal to the measurement plane and is therefore fundamentally an integrated horizontal flux (IHF) method. Only measurement periods with valid measurements for all five PICs of a measurement plane for the entire 30-min period were included in the analysis. Comparisons between the IHF and the various flux-gradient methods discussed above were limited to periods of measurement when only one of the four planes had a measured flux. Because of the mismatch in sampling intervals between the sonic anemometer and the TDLAS, some of the flux measured by the IHF method would be a result of the covariance between horizontal wind perturbations and gas concentration perturbations (Denmead, 2008). We expect an overestimation of the flux using the IHF method by 5 to 20%. Background concentrations of  $\text{NH}_3$  and the corresponding flux were assumed negligible.

Using the “parameterized” and “measured  $u_*$ ” flux-gradient methods the gradient at 4.4 m height was determined from the wind measurements made at heights of 2.5 m and 16.2 m, the stability  $\zeta$  determined from the 4.4 m sonic anemometer height, and the concentration gradient at 4.4 m height determined from the gradient approximated from horizontal gradients of  $\text{NH}_3$  based on the TDLAS measurements. The vertical  $\text{NH}_3$  gradients were derived from the slant-path measurements by defining three layers (0–2.6 m, 2.6–8 m, and 8–15.2 m) and assuming the gas is uniformly mixed within each

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layer. The concentrations of each layer were then determined by sequentially solving three equations with three unknowns where the fractional length of the overall measurement path through each layer is known, the concentration through the entire layer is assumed uniform and the cumulative concentration of gas along the entire path is measured (see also Table 1). An analysis of the propagation of measurement errors for the MDL of the “similarity” and the “parameterized” flux-gradient methods were the same as for the Fuhrberg campaigns. The  $K_c$  MDL was approximately  $0.002 \text{ m}^2 \text{ s}^{-1}$ .

**3.6 Chamber measurements and quality assurance checks**

Measuring soil  $\text{N}_2\text{O}$  emissions using chambers is widely accepted. Gas samples were collected in ambient air before closing of the chambers and every 20 min over 1 h after the closing. Further information describing the applied method is given in Weymann et al. (2009) and von der Heide et al. (2009). Concentrations of  $\text{N}_2\text{O}$  were analysed from the gas samples using a gas chromatograph (Fisons GC 8000, Milan, Italy) equipped with a split-injector, an electronic capture detector and a HP-PLOT Q column (30 m length \* 0.32 mm ID; Agilent Technologies, Santa Clara, USA). The gas samples were analysed within 3 weeks after the collection.

The  $\text{N}_2\text{O}$  fluxes were determined and qualified for use by the linearity of the increases/decreases of the concentration measurement. As the  $\text{N}_2\text{O}$  fluxes were very low, we chose a linearity criterion of the square of regression coefficient of 0.8. If this coefficient was smaller than 0.8, the flux rate of the chamber was set to zero. The NDFE method of Livingston et al. (2006) was also used to determine the flux.

**4 Results****4.1 Atmospheric dynamics aspects**

The atmospheric conditions during the Fuhrberg experiment were similar to the stable boundary layer measurement periods of the Indiana experiment. The range of  $\zeta$

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differed some at the stable extreme to 0.81 in the Indiana experiment and to 1.16 in the Fuhrberg experiment. In both experiments, the concentration gradient range in  $\text{NH}_3$  (0.6 to  $114.6 \mu\text{g m}^{-3} \text{m}^{-1}$ ) and  $\text{N}_2\text{O}$  (4.4 to  $22.0 \mu\text{g m}^{-3} \text{m}^{-1}$ ) increased with decreasing wind speed (Fig. 5). There was no distinct relationship between the vertical  $\text{NH}_3$  and  $\text{N}_2\text{O}$  concentration gradients and the vertical gradient in horizontal wind speed, however there was a tendency for greater concentration gradients corresponding to low horizontal wind speed gradients less than  $0.2 \text{ m s}^{-1} \text{m}^{-1}$ .

The turbulent conditions during the quality-assured 30-min periods were also similar in both experiments (Fig. 6a) with the calculated  $K_c$  for both experiments decreased with decreasing wind gradient and  $u_*$  and increasing  $\zeta$  for all stable measurement periods (normalized by  $Sc$  in Fig. 6b,d). The wide range in possible  $K_c$  for a given wind speed gradient is a result of low gradients developing either under high winds under near-neutral stability or low winds. Note that the  $Sc K_c$  does not approach the molecular diffusivity  $K_m$  of the order  $1 \times 10^{-5} \text{ m}^2 \text{s}^{-1}$  as  $\zeta$  increases to 1.0 (Fig. 6c). This was probably largely a result of the MDL for the measured  $K_c$  although it may also indicate the predominance of turbulent mixing over molecular mixing for the weak, often intermittent, turbulence of the stable boundary layer. Note that the quality assurance minimizes the influence of waves and other coherence structures on the calculated  $Sc K_c$ .

### 4.2 $\text{NH}_3$ emission fluxes from a lagoon

It was assumed that the flux determined from the IHF method was closest to the actual flux. The  $\text{NH}_3$  flux calculated by IHF varied from near zero to  $330 \text{ mg m}^{-2} \text{h}^{-1}$  (mean  $90 \text{ mg m}^{-2} \text{h}^{-1}$ ) with 49 30-min measurements with  $\zeta$  between 0 and 0.01, 139 with  $\zeta$  between 0.01 and 0.1, 86 with  $\zeta$  between 0.1 and 1.0. A negative flux indicated flux to the lagoon from the surroundings. The IHF estimates of  $\text{NH}_3$  flux in the Indiana experiment are homogeneously spread around a relationship from logarithmic  $\zeta$ .

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There is an IHF positive bias due to in part to the inability of the method to quantify a “plume” when the concentration field is near background values and in part to the inclusion of turbulent diffusive fluxes (Denmead, 2008).  $\text{NH}_3$  flux calculated using the three flux-gradient methods were generally lower than the IHF method (Fig. 7). This tendency was probably in part due to errors in the calculated vertical  $\text{NH}_3$  concentration gradients described previously. Since this error was present for all three flux-gradient methods, a comparison of the three methods can still give useful information concerning the ability of each method to approximate the IHF. The correlation ( $r$ ) between the parameterized flux-gradient and the IHF method was 0.53, with root mean square error (RMSE) of  $43 \text{ mg m}^{-2} \text{ h}^{-1}$  (solid line, Fig. 7). The regression for the parameterized flux comparison is driven by three outlying values that are not indicated in Fig. 7: the dashed line indicates the linear regressions without outliers less than zero and greater than  $180 \text{ mg m}^{-2} \text{ h}^{-1}$ . Removing these outliers dramatically improves the correlation between the IHF and the parameterized methods to 0.61 with negligible change in the RMSE ( $40 \text{ mg m}^{-2} \text{ h}^{-1}$ ).

Calculating the flux based on the measured wind profile resulted in a better relationship to the IHF method than the parameterized method:  $r = 0.68$ ,  $\text{RMSE} = 38 \text{ mg m}^{-2} \text{ h}^{-1}$  (Fig. 7). The use of the measured  $u_*$  in place of the parameterized  $u_*$  prevented the outlying flux estimates and consequently is an improvement on the parameterization of  $u_*$  (Fig. 7) with the greatest flux related to less stable conditions (Fig. 8).

The difference between the flux calculations using the parameterization of  $u_*$  and the measurement-based  $u_*$  suggested the wind profile was not log-linear as implied in the analysis of  $\psi_c$  (Fig. 3), reducing the utility of the parameterization to describe the gradient. This is consistent with the finding that the wind profile is commonly more linear than log-linear in the stable boundary layer (Mahrt, 2007). These non-log-linear wind profiles within the measured layer strongly influenced the flux.

The  $\text{NH}_3$  flux calculated using the similarity method assumed linear profiles rather than a log-linear profile of wind and concentration. This method resulted in fluxes closer



to those determined by the IHF method than the other two methods that assumed log-linear profiles (Fig. 7). The correlation between the diffusivity method and the IHF method was 0.80 with a RMSE of  $31 \text{ mg m}^{-2} \text{ h}^{-1}$  (Fig. 7).

All flux-gradient methods resulted in estimated fluxes with the range of values decreasing with increasing  $\zeta$  (Fig. 8). This was consistent with mixing theory where mechanical turbulent mixing is suppressed to an increasing extent as the stability increases.

### 4.3 N<sub>2</sub>O emission fluxes from soil

The calculated  $K_c$  values for this experiment were within the range of those calculated from the Indiana experiment and showed similar relationships to the vertical gradient in horizontal wind speed,  $u_*$ , and  $\zeta$  (Fig. 6). The calculated diffusivity for the Fuhrberg experiment showed a decrease with the increasing stability limited by the measurement errors to values of  $0.01 \text{ m}^2 \text{ s}^{-1}$  (Fig. 6).

During the first campaign (16–17 October 2007) wind speeds varied from  $2.3 \text{ m s}^{-1}$  to  $0.6 \text{ m s}^{-1}$  and the surface boundary layer was characterised by stable and very stable conditions with  $\zeta$  up to +8. This extreme stability resulted in minimal valid emissions determinations. Between 16:30 and 19:30 h of 16 October, all wind speeds (above and below) were relatively small with corresponding small wind speed gradients. During this time the difference between the concentration measurements was not statistically significant (Fig. 9). Between 19:30 and 23:30 h the measured vertical concentration differences were statistically significant resulting in valid flux determinations. After this time, the difference in the two concentrations was again not statistically significant. There are some gaps in the N<sub>2</sub>O measurement time series due to the necessity for internal optical adjustment of the spectrometers (see Figs. 9 and 10).

During the second campaign, the surface boundary layer had very stable conditions and very low horizontal wind speeds with  $\zeta$  up to +11. The N<sub>2</sub>O concentrations in ambient air (17 June 2008) at both heights including the statistical errors (as

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± values) are shown in Fig. 10. Most measured concentration differences were not significantly different.

Chamber measurements of the N<sub>2</sub>O flux were made on 17 October 2007 after the first campaign when the wind speed at 2.7 m a.g.l. was 0.5 m s<sup>-1</sup>, the air temperature was 10 °C and there was no precipitation. These conditions were similar to those during the flux-gradient measurements. Of the four chambers used, one had an unusually high concentration value after 40 min (value excluded), one failed the linearity criterion and had one unusually high concentration (flux set to 0), one indicated an up-take of N<sub>2</sub>O (linear regression -9.2 μg m<sup>-2</sup> h<sup>-1</sup> and NDFE -13.7 μg m<sup>-2</sup> h<sup>-1</sup>), and one indicated flux of 155.5 μg m<sup>-2</sup> h<sup>-1</sup> by both data analysis methods. The mean of the four chamber measurements was 37 μg m<sup>-2</sup> h<sup>-1</sup> by the linear regression and 36 μg m<sup>-2</sup> h<sup>-1</sup> by the NDFE method.

The quality assurance condition of a significant difference in the concentrations for the flux-gradient method was fulfilled in only 22 30-min mean measurements of the two campaigns. Exclusions of measurement periods with high wind shear or strong waves during the averaging period resulted in valid flux-gradient measurements for only 8 of 56 30-min measurement periods.

The general dependence of the calculated emission fluxes on ζ is given in Fig. 11. As with the Indiana experiment (Fig. 8), the similarity method resulted in the highest fluxes. Also in both the Indiana and Fuhrberg experiments the flux based on a measured  $u_*$  was similar but generally greater than that based on the MOST parameterization. All three approaches to calculate the flux using flux-gradient methods showed similar tendencies. As with the upper emissions bound of the Indiana experiment (Fig. 8), the emission rates during near-neutral conditions were much higher than during stable conditions (Fig. 11). This was likely caused by changes in weather conditions such as higher wind speeds with correlated greater mixing causing enhanced emission processes of the soil or by higher soil and lagoon temperatures, which will change equilibrium gas concentrations. However it is important to remember that near-neutral conditions do not necessarily result in higher emissions, as evidenced by the results

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for the  $\text{NH}_3$  emissions in Indiana (Fig. 8). The negative exponential tendency affects only the maximum possible flux as a smaller concentration gradient would reduce the flux under near neutral conditions. The range of possible flux values across the range of  $\zeta$  was evident in the  $\text{NH}_3$  fluxes in Indiana (see Fig. 8).

The  $\text{N}_2\text{O}$  flux was measured by the chamber method on 19 June 2008 after the second campaign. The meteorological conditions were similar as on 17 and 18 June: no precipitation, a wind speed at 2.7 m a.g.l. of  $0.7 \text{ m s}^{-1}$ , and an air temperature of  $13^\circ\text{C}$ . Of the four chambers measured, one set of concentration measurements failed the linearity criteria and was assumed to be zero but calculated as  $18.4 \mu\text{g m}^{-2} \text{ h}^{-1}$  by the NDFE method. The other chambers indicated: 1) a flux of  $267.7 \mu\text{g m}^{-2} \text{ h}^{-1}$  by both methods, 2) a flux of  $97.8 \mu\text{g m}^{-2} \text{ h}^{-1}$  by linear regression and  $113.6 \mu\text{g m}^{-2} \text{ h}^{-1}$  by NDFE method, and 3) a flux of  $146.2 \mu\text{g m}^{-2} \text{ h}^{-1}$  by linear regression and  $1029.6 \mu\text{g m}^{-2} \text{ h}^{-1}$  by NDFE. The mean measured flux was  $128 \mu\text{g m}^{-2} \text{ h}^{-1}$  using the linear regression method and  $357 \mu\text{g m}^{-2} \text{ h}^{-1}$  for the NDFE method.

The difference between the mean of the  $\text{N}_2\text{O}$  emission rate results using the open-path flux-gradient methods during stable conditions averaged  $236 \mu\text{g m}^{-2} \text{ h}^{-1}$  using the parameterized approach,  $498 \mu\text{g m}^{-2} \text{ h}^{-1}$  using the measured  $u_*$  approach and  $903 \mu\text{g m}^{-2} \text{ h}^{-1}$  using the similarity approach. Although similar, the average emission from the chamber measurements using the NDFE simulation was  $313 \mu\text{g m}^{-2} \text{ h}^{-1}$  (Fig. 11) were performed at a time difference of one day in October 2007 and two days in June 2008.

The calculated  $\text{N}_2\text{O}$  emission rates from the chambers and the flux-gradient methods agree with the order of magnitude of chamber and box flux measurements from other grasslands. Dittert et al. (2005) found  $\text{N}_2\text{O}$  fluxes of 2 to  $19 \mu\text{g m}^{-2} \text{ h}^{-1}$  over similar soils on a Northern Germany grassland. Goossens et al. (2001) determined  $\text{N}_2\text{O}$  fluxes from 10 to  $2200 \mu\text{g m}^{-2} \text{ h}^{-1}$  from a Belgian grassland.

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## 5 Data interpretation and discussion

The prerequisites for the application of a flux-gradient measurement method are: (1) negligible advection, (2) negligible chemical changes during the transport of air from the footprint to the measurement system, (3) vertical turbulent exchange in a not too stable, neutral or slightly unstable lower atmosphere ( $\zeta \leq +10$ ). These conditions were fulfilled in the measurement periods reported here due to the horizontal homogeneity within the fetch and evident presence of vertically-consistent turbulence regardless of wind direction (see Foken, 2008, p. 26; Schmid, 1994). The horizontal homogeneity condition was well-fulfilled at the Fuhrberg location but only generally fulfilled at the Indiana location. At the Indiana experiment location, the proximity of the lagoon source to the measurement planes and the resulting changes in local surface roughness and stability associated with the lagoon appeared to influence the universal wind profile (and the stability influence on the profile) and may have decreased the correlations between the IHF and flux-gradient methods. Fundamentally the transport from the surface to the measurement height requires a turbulent linkage between the surface and the height of measurements. This condition was assured under the stable conditions studied by restricting analysis to periods with low vertical wind shear, velocity gradient (Indiana experiment only), stability, and consistency of turbulent structure even if those periods are intermittently occurring within the turbulence time history.

Since the fluxes calculated using the similarity method correlated better with the IHF method than the MOST flux-gradient methods based on the measured  $u_*$  or parameterized theory in the Indiana experiment, it appears that the MOST flux-gradient methods incorrectly assumed log-linear profile of the wind speed and concentration gradient in the stable boundary layer. Consequently it was assumed that the similarity method provided the best estimate of  $N_2O$  flux in the Fuhrberg experiment.

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## 6 Conclusions

The use of open-path flux-gradient methods in non-intrusive up-scaling of trace gas fluxes near the surface was demonstrated here for  $\text{N}_2\text{O}$  and  $\text{NH}_3$  in stable boundary layers and path lengths of 100 m. The evaluation of three flux-gradient methods in the Indiana experiment indicate that a similarity approach that does not assume log-linear relationship of wind speed and concentration with height is the best method to approximate the flux under stable conditions ( $0 < \zeta < 1$ ). The determination of flux under these stable conditions however requires careful evaluation of the continuity of turbulence between the surface and the height of measurement. The near-simultaneous chamber measurements made at Fuhrberg showed similar mean emission values as the open-path flux-gradient measurement under stable conditions. The measured fluxes are comparable to other  $\text{N}_2\text{O}$  flux measurements reported in the literature.

### 6.1 Advantages and disadvantages of the open-path flux-gradient method

In general the application of flux-gradient methods for trace gas flux determinations is difficult to use when either there are: (1) low wind speeds with poorly defined wind speed gradients and typically high vertical concentration gradients or (2) high wind speeds with small concentration gradients. Under stable conditions the weak and intermittent turbulence further restricts the use of micro-meteorological measurement of fluxes to conditions in which this turbulence is consistent across the layer and that the variability in vertical motion is not a result of waves or other structures that do not transfer mass vertically in the boundary layer. This study provided an evaluation of three flux-gradient methods against an integrated horizontal flux (IHF) method under conditions of positive  $\zeta$  and vertically-consistent turbulence. Of the flux-gradient methods considered, the approach using the similarity flux-gradient method resulted in the smallest difference to the IHF calculation implying that the wind gradient and/or concentration gradient is not log-linear under stable conditions and MOST cannot be applied.

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## 6.2 Limitations to large-scale long-term emissions determinations

Careful quality assurance of sonic anemometers provided sufficient accuracy, precision, and detection limits for flux determination and qualification in the stable boundary layer using a flux-gradient method. A primary limitation on the yield of valid flux measurements from these methods was instrument sensitivity; the precision of the sonic anemometers was sufficient to define the vertical gradient in wind speeds and turbulence parameters but the available FTIR spectrometer for the N<sub>2</sub>O measurements did not routinely provide the necessary precision and detection limits to determine a statistically significant concentration gradient needed for flux determination. The use of a wider range in height may improve the flux estimate provided the lower measurement can be made closer to the surface. The FTIR also required internal optical adjustments which were dependent on the temperature environment and these adjustments significantly interrupted the measurement series. This was a first attempt at using flux-gradient methods in stable boundary layers with local perturbation partially averaged out by using open-path gas concentration measurements. A more sophisticated, longer term demonstration of the performance of trace gas flux determination by open-path gas measurements and flux-gradient approaches is needed to confirm applicability, ideally with spatially averaged turbulence measures matching the optical path of the gas measurement.

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**Table 1.** Data, which have been filtered at all heights for the quality-assured determinations of fluxes under low wind, stable or night-time atmospheric conditions and acceptable ranges for the various parameters.

Parameter	Phenomenon evaluated	Acceptable range/criterion	
		Fuhrberg	Indiana
30-min mean wind	Directional shear between the heights	< 10°	< 20°
	Wind direction variation	Within 45° of the perpendicular to a measurement plane	
	Wind speed limit	> 0.1 m s <sup>-1</sup>	
Perturbation of wind components $u'$ , $v'$ , $w'$	Spikes	> 5 standard deviations above the mean over 5 min, > 80 % of measurements had to be valid	
Perturbation of wind component $w'$	Intermittency and waves	Visual inspection of 30-min means	
Vertical turbulence (mean $w'^2$ ) $\sigma_w^2$	Analysis of spectral distribution	Similarity at both heights	
Calculated $\zeta$	Analysis of spectral distribution	Continuity, consistency across heights	
$Ri_B$	Dominance of turbulent production over destruction	–	< 0.2
$\Delta$ concentration	Consistency	Difference > 0 to statistical errors (combined errors with 0.90 levels, $t$ -statistics: factor 1.44)	Consistency across heights

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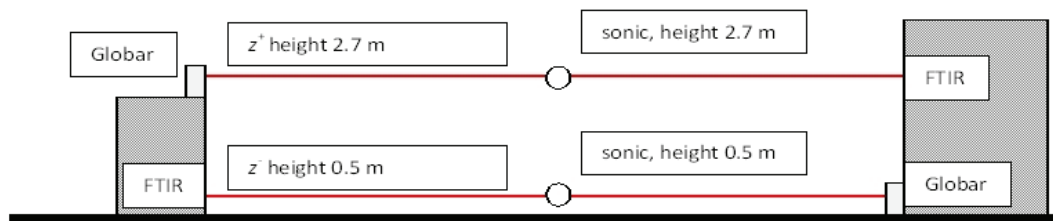
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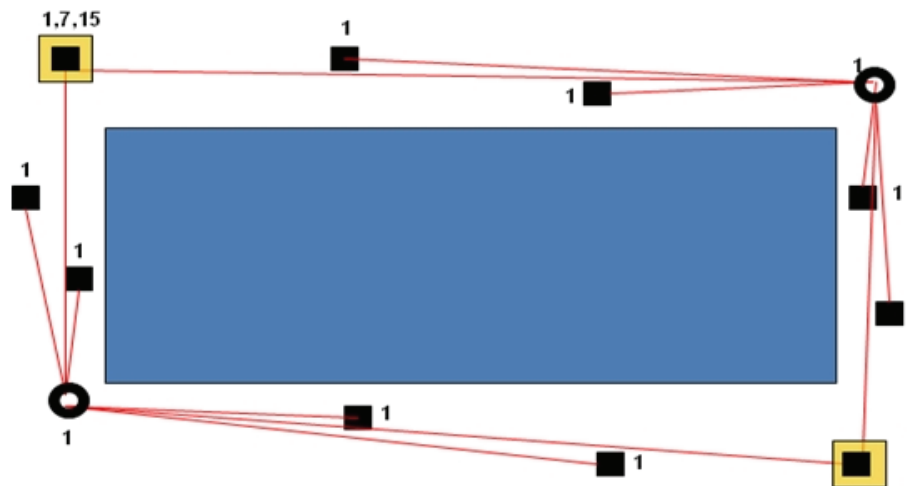


**Fig. 1.** Configuration of the Fuhrberg experiments with FTIR spectrometers and globars showing the measurement paths and sonic anemometers. Chamber measurements were made in about 10 m distance to the one end of the FTIR paths. The measurement volume was oriented in north-south direction.

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**Fig. 2.** Configuration of Indiana experiment. The scanning TDLAS (open circle), retro-reflectors (filled squares), towers (yellow squares) and open paths (solid lines) are shown with the height of the equipment indicated in meters.

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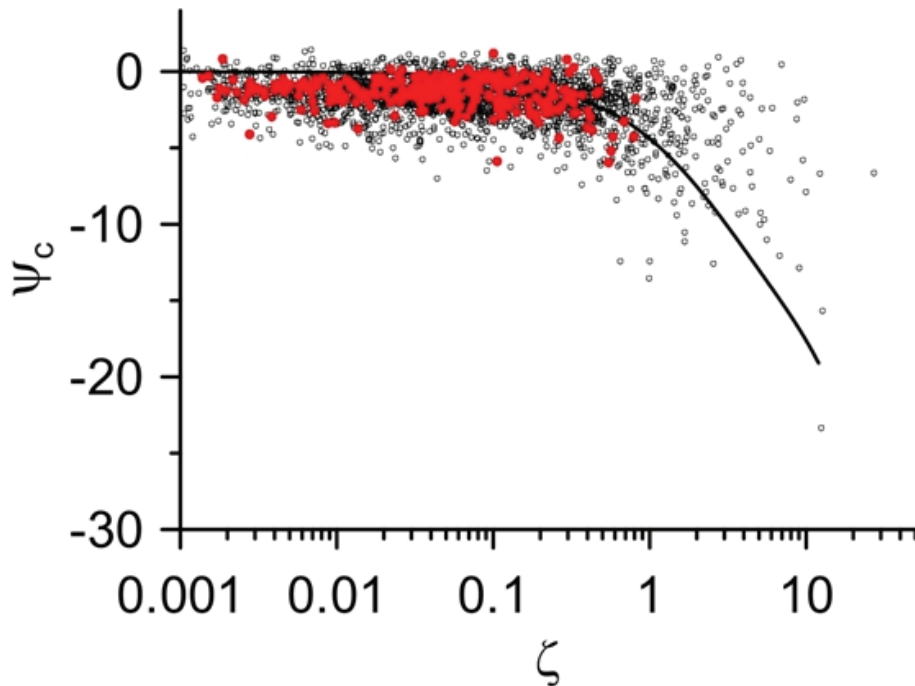
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**Fig. 3.** Relationship for the trace gas stability correction function  $\psi_c$  between the dependence of integrated profile at 4.4 m ( $z_F$ ) based on measurements at 2.5 m ( $z^-$ ) and 16.2 m ( $z^+$ ) on stability  $\zeta$ . The measured values (circles) and the nominal relationships according to Holtslag and de Bruin (1988) (solid line) are indicated. The remaining data after the quality assurance procedure stated in the text are red.

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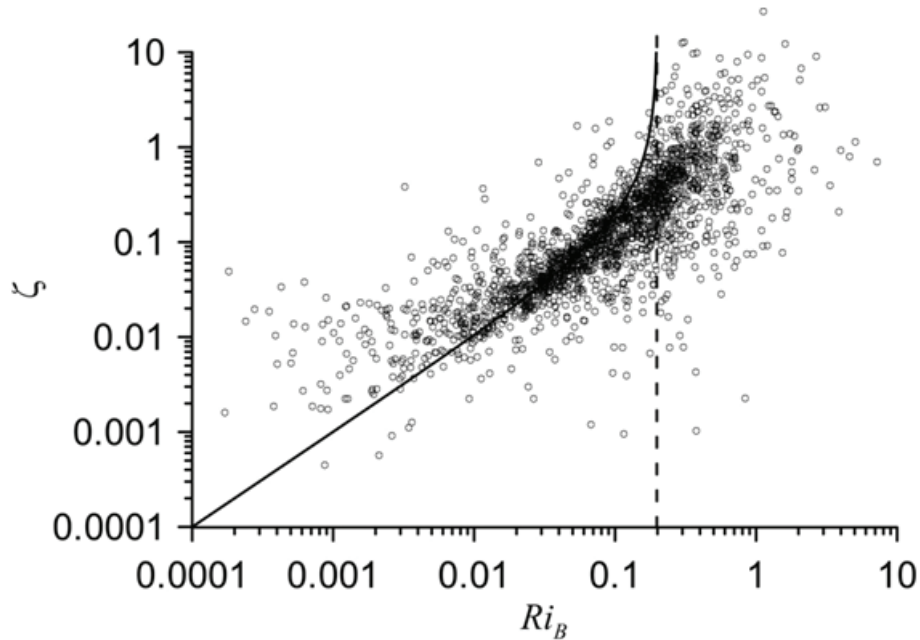
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**Fig. 4.** Relationship between the two measures of stability  $\zeta$  from bulk Richardson number  $Ri_B$  during the Indiana experiment. The critical  $Ri_B$  (dashed line) and BDH relationship below critical  $Ri_B$  (solid line) are indicated.

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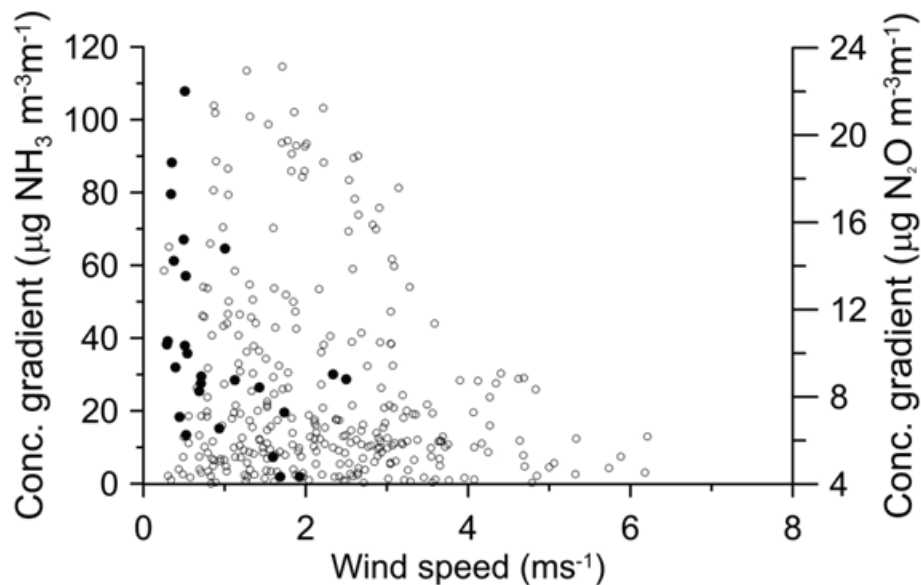
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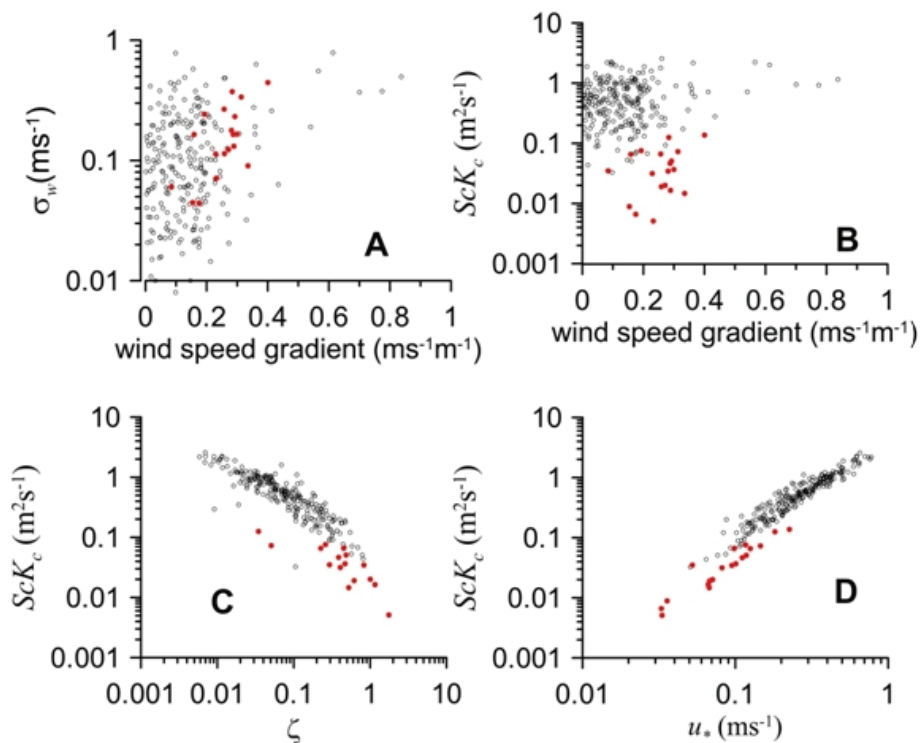
**Fig. 5.** Variation in concentration gradient relative to wind speed. Values in both the Indiana (NH<sub>3</sub>, open circles) and Fuhrberg (N<sub>2</sub>O, closed circles) experiments are indicated.

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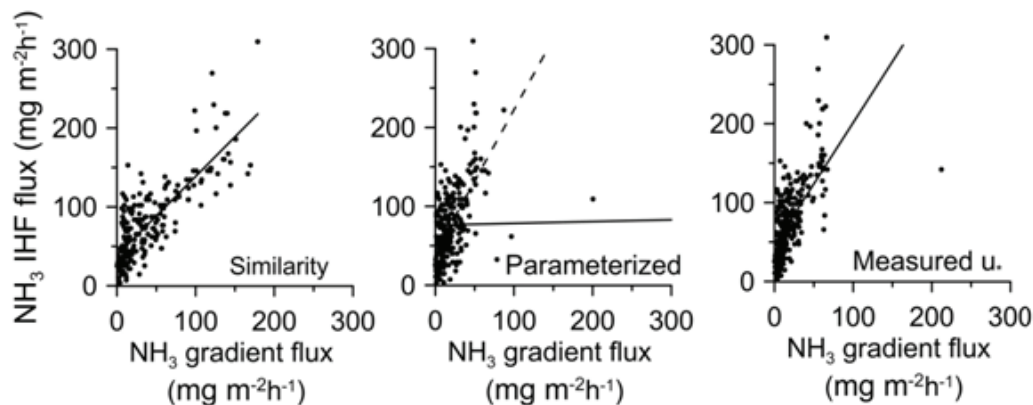


**Fig. 6.** Variation in  $\sigma_w$  and  $ScK_c$  at the flux calculation height with wind speed gradient below that height (**A** and **B**), and variation in  $ScK_c$  with stability  $\zeta$  (**C**) and friction velocity  $u_*$  (**D**) at that height. Values in both the Indiana (open circles) and Fuhrberg (closed red circles) experiments are indicated.

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**Fig. 7.** The comparison between the IHF method and the flux-gradient method using similarity  $K_c$  according to Eqs. (4) and (5), parameterised theoretical  $u_*$  and  $\psi_c$  according to Eqs. (1) to (3) and measured  $u_*$ . Linear regressions between the flux measures are indicated by solid lines. The dashed line at parameterized flux comparison indicates the linear regressions when outliers less than zero and greater than  $200 \text{ mg m}^{-2} \text{ h}^{-1}$  are neglected.

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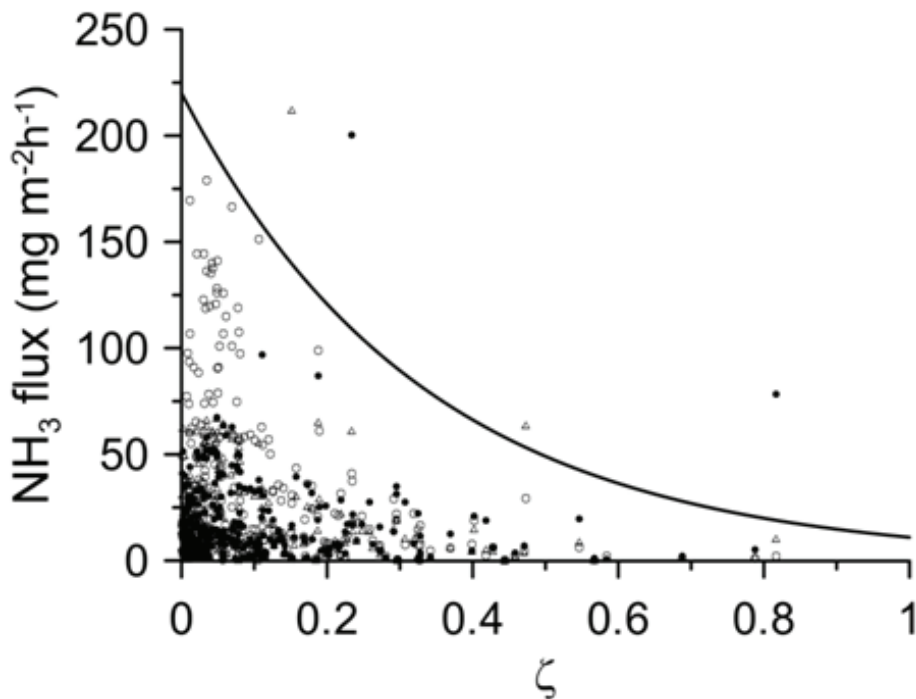
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**Fig. 8.** The relationship between stability  $\zeta$  and the computed fluxes using flux-gradient method using a similarity approach (open circles), the flux-gradient method assuming a MOST parameterised  $\psi_c$  (closed circles), and the flux-gradient method using measured  $u_*$  (triangles) are indicated. Three extreme outlier values of the parameterized method are not plotted. The line shows the upper emissions bound.

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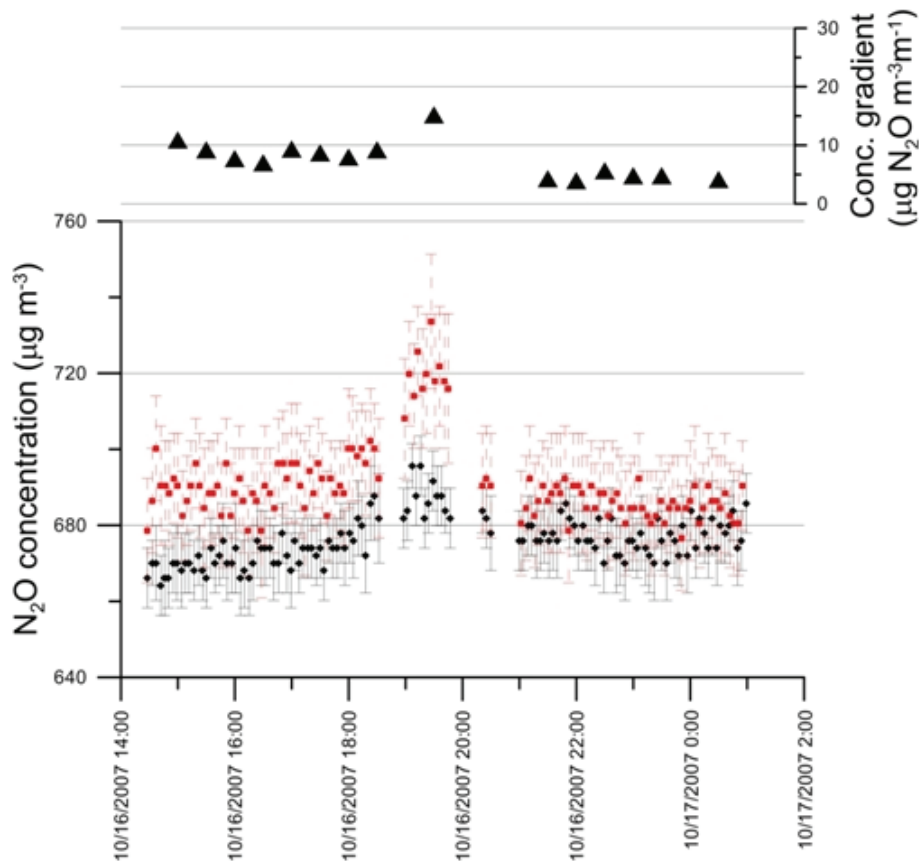
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**Fig. 9.** Temporal variation of the  $\text{N}_2\text{O}$  concentrations in ambient air (16–17 October 2007) at both heights (0.5 m: side-oriented squares, 2.7 m: edge-oriented squares) including the statistical errors (random absolute error as columns around the measurement value in both directions) and concentration gradients.

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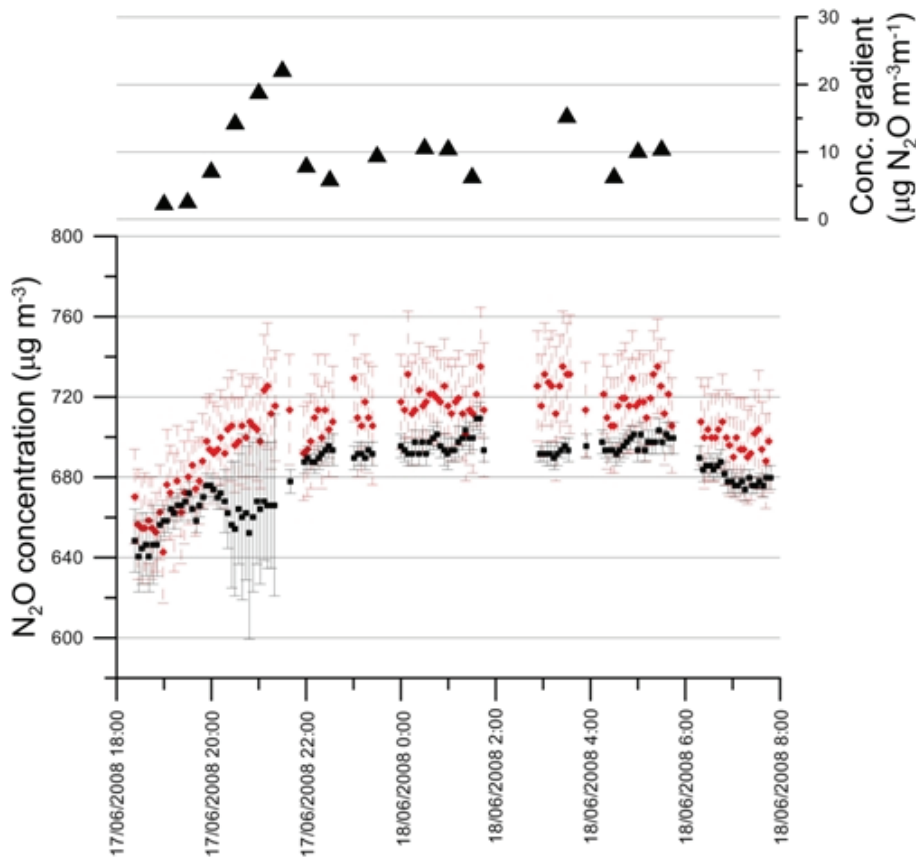
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**Fig. 10.** Temporal variation of the  $\text{N}_2\text{O}$  concentrations in ambient air (17–18 June 2008) at both heights (0.5 m: edge-oriented squares, 2.7 m: side-oriented squares) including the statistical errors (random absolute error as columns around the measurement value in both directions) and concentration gradients.

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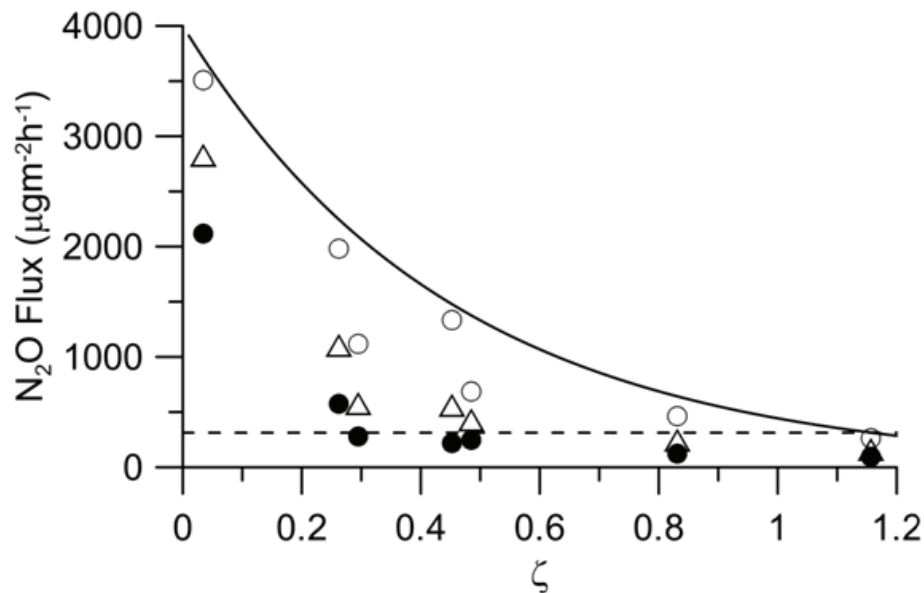
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**Fig. 11.** Dependence of the calculated fluxes on the stability of the boundary layer ( $\zeta > 0$ ). The flux measurement based on the similarity approach (open circles), MOST parameterized flux-gradient theory (closed circles), and measured  $u_*$  flux-gradient theory (open triangles) is indicated. The dashed line indicates the chamber-measured flux (see text). The line shows the upper emissions bound.

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