

**1.57 μm differential
laser absorption
spectrometry**

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**An airborne amplitude-modulated 1.57 μm
differential laser absorption spectrometry:
simultaneous measurement of partial
column-averaged dry air mixing ratio of
 CO_2 and target range**

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Abstract

Simultaneous measurements of the partial column-averaged dry air mixing ratio of CO₂ (\bar{q}) and target range were demonstrated using airborne amplitude-modulated 1.57 μm differential laser absorption spectrometry (LAS). The LAS system is useful for discriminating between ground and cloud return signals and has a demonstrated ability to suppress the impact of integrated aerosol signals on differential absorption optical depth ($\Delta\tau$) measurements. A high correlation coefficient (R) of 0.99 between $\Delta\tau$ observed by LAS and $\Delta\tau$ calculated from in-situ measurements of CO₂ was obtained. The averaged difference in \bar{q} obtained from LAS (\bar{q}_{LAS}) and validation data (\bar{q}_{val}) was within 1.5 ppm for all spiral measurements. A significant profile was observed for both \bar{q}_{LAS} and \bar{q}_{val} , in which lower altitude CO₂ decreases compared to higher altitude CO₂ attributed to the photosynthesis over grassland in the summer. In the case of an urban area where CO₂ and aerosol are highly distributed in the lower atmosphere in the winter, the difference of \bar{q}_{LAS} to \bar{q}_{val} is -1.5 ppm, and evaluated \bar{q}_{LAS} is in agreement with \bar{q}_{val} within the measurement precision of 2.4 ppm (1σ).

1 Introduction

A global carbon cycle study using higher spatial resolution than an $8^\circ \times 10^\circ$ grid is currently required to improve the knowledge of the carbon cycle. Transport models and observational data sets improve evaluations of regional carbon fluxes. Evaluation of the spatial and temporal distribution of natural carbon fluxes over land and ocean continues to be difficult, hindering improvements in the quantifying and understanding of the mechanism of the fluxes (Kawa et al., 2010; Kaminske et al., 2010). Uncertainty in flux evaluations is a major contributor to uncertainty in climate predictions (Randall et al., 2007). However, confirmation of the consistency between the sum of the regional and global budgets of carbon fluxes is expected to provide a unique index of the level of confidence in the outcomes of climate mitigation policies (IPCC, 2007). A sustainable

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technique for CO₂ remote sensing from space is one of the greatest challenges and necessities for understanding the global carbon cycle, as well as for predicting and validating its evolution under future climate changes.

The Greenhouse gases Observing SATellite (GOSAT) is the first step in dealing with the above-mentioned issue (Yoshida et al., 2011; Palmer et al., 2011). The sensors on-board the GOSAT, GOSAT-2, the Orbiting Carbon Observatory (OCO), and OCO-2 are based on a passive remote sensing technique (Kuze et al., 2009; Crisp et al., 2012; Eldering et al., 2012). The GOSAT sensor was developed to derive the column-averaged mixing ratio of CO₂ (XCO₂) with a precision better than 1 % for an 8° × 10° grid without any biases or with uniform bias (Rayner and O'Brien, 2001; Houweling et al., 2004; Miller et al., 2007; Morino et al., 2011). However, there are unavoidable limitations imposed by the measurement approach: (1) the best performance for CO₂ total column measurements can only be obtained under clear-sky conditions; (2) seasonal dependence, such as in the case of the Northern Hemisphere in winter, reduces its global coverage, and (3) CO₂ measurements are highly sensitive to unknowns and variations in cloud and aerosol contamination.

In contrast, active optical remote sensors are valuable for near-future trace gas measurements from space. Ground-based measurement techniques using a high-energy pulse laser have been developed to measure vertical CO₂ mixing ratios (Amediek et al., 2008; Sakaizawa et al., 2009; Ishii et al., 2010; Gibert et al., 2011). Airborne systems to observe partial column-averaged CO₂ have also been reported in earlier studies (Browell et al., 2011; Abshire et al., 2010; Spiers et al., 2011) to demonstrate technology feasible for future space-borne missions. Although in a pulsed system aerosol or cirrus clouds have less impact on total column measurements, the pulsed-laser wavelength must be stabilized at a specific position less than 100 kHz to reduce error due to wavelength stability, which requires large resources. Errors due to variations in the surface reflectivity along the track also increase the impact, unless the transmitter has a double pulse system (Yu et al., 2003). Our fiber-based continuous laser approach to measure the differential absorption optical depth (DAOD = $\Delta\tau$) allows for compact

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storage of the components, including the electronics and optics. Moreover, the system achieves matching of the optical axes of multi-transmitted laser beams, which can contribute to reducing error due to footprint overlap. In this paper, we evaluate the performance of airborne 1.57 μm amplitude-modulated differential laser absorption spectrometry (LAS) for obtaining the partial column-averaged mixing ratio of CO_2 with simultaneous range detection. In addition, the impact of integrated aerosol signals on CO_2 measurements is described in an area where aerosol has accumulated (e.g. over a city).

2 Partial column-averaged CO_2

A LAS system on an aircraft platform was utilized for measuring the light scattered or reflected by a target (land or sea surface). Our system employed three narrow linewidth lasers, which are based on continuous-wave distributed feedback diode lasers and a fiber amplifier (Kameyama et al., 2011a). The system used two laser wavelengths during measurements, and the output of each laser was amplitude modulated by different sinusoidal waves. The details of amplitude modulation, frequencies, and phase shift are discussed in Sect. 3.

One wavelength (offline, λ_{off}), for which there was weak or no gas absorption, was selected as a reference. The other wavelength (online, λ_{on}) was selected for strong gas absorption. In this airborne test, the online wavelength could be selected from the center (λ_{onc}) or edge (λ_{one}) position of the absorption curve (as shown in Fig. 1). The online laser power was attenuated relative to the offline wavelength in the atmosphere. By taking the ratio of online to offline signals, we could measure the differential absorption optical depth $\Delta\tau$.

The performance of $\Delta\tau$ was evaluated using the ratio of the mean value to deviation and the magnitude of the bias error. Our system obtained round-trip $\Delta\tau$ and the range (z) from the height of the aircraft to the target (Sakaizawa et al., 2010; Kameyama et al., 2011a).

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$$\Delta\tau = \ln \left(\frac{P_r(\lambda_{\text{off}}) P_m(\lambda_{\text{on}})}{P_r(\lambda_{\text{on}}) P_m(\lambda_{\text{off}})} \right) \quad (1)$$

$$z = \frac{\Delta\phi T c}{4\pi} \quad (2)$$

where $P_r(\lambda_{\text{on}})$ and $P_r(\lambda_{\text{off}})$ are the online and offline laser powers received from the surface of the ground, $P_m(\lambda_{\text{on}})$ and $P_m(\lambda_{\text{off}})$ are the monitored transmitted online and offline laser powers, $\Delta\phi$ is the phase difference between the monitored and the received sinusoidal signals, T is the period of a modulated sinusoidal signal, and c is the speed of light.

The phase difference between transmitted and received sinusoidal signals corresponds to the range at which a target is acquired. Each laser power is amplitude-modulated 10 kHz for λ_{on} and 11 kHz for λ_{off} using LiNbO₃ devices. The phase identification is performed by the fast Fourier transform (FFT). This range of the detection technique is ambiguous at the inverse of the modulation frequency. However, the height of an elevated layer, such as cirrus or water clouds, can be compared with ground returns. In addition, images taken under the moving platform help in filtering signals with cloud returns, especially over complex terrain.

The partial column-averaged dry air mixing ratio of CO₂ (\bar{q}) can be described by the following equation (Ehret et al., 2008):

$$\bar{q} = \Delta\tau_{\text{ob}}/2i\text{wf} \quad (3)$$

$$i\text{wf} = \int_{z_{\text{ac}}}^{z_{\text{trg}}} \text{wf}(r) dr \quad (4)$$

$$\text{wf}(r) = \Delta\sigma_{\text{CO}_2}(r) n_{\text{air}}(r) (1 - V_{\text{H}_2\text{O}}(r)). \quad (5)$$

Here, $i\text{wf}$ is the integrated weighting function; z_{ac} , the altitude of the aircraft; z_{trg} , the target height; $\text{wf}(r)$, the weighting function at a specific altitude r ; $\Delta\sigma$, the differential

absorption cross section of CO₂ between the online and offline wavelengths; n_{air} , the air molecular number density, and $V_{\text{H}_2\text{O}}$, the water vapor mixing ratio. n_{air} and $V_{\text{H}_2\text{O}}$ are calculated from meteorological observation or mesoscale re-analysis data.

The measurement uncertainty is a quadratic summation of the precision and bias factors ($(\delta\bar{q}/\bar{q})^2 = \text{precision}^2 + \text{bias}^2$). In this paper, we evaluate the precision and bias separately. The precision is evaluated using following equations:

$$\text{precision}^2 = \text{SNR}(\Delta\tau)^{-2} + (\partial\text{iwf}/\partial z)^2 + (\partial\text{iwf}/\partial\lambda)^2 \quad (6)$$

$$\text{SNR}(\Delta\tau) = \delta\Delta\tau/\Delta\tau \quad (7)$$

where $\delta\Delta\tau$ represents the fluctuation in measured $\Delta\tau$, and δiwf depends on the evaluation errors of the range accuracy ($\partial\text{iwf}/\partial z$) and the wavelength stability ($\partial\text{iwf}/\partial\lambda$). If mesoscale data re-analysis is employed in Eq. (4), the error due to temporal and spatial differences compared with radiosonde measurements ($(\partial\text{iwf}/\partial T)^2 + (\partial\text{iwf}/\partial P)^2 + (\partial\text{iwf}/\partial U)^2 = (0.16\%)^2$) is added to Eq. (6), which T means temperature, P means pressure, U means relative humidity.

Considering the sensitivity of near-surface CO₂ and $\text{SNR}(\Delta\tau)$, they depend on the stabilized position of the online wavelength (Fig. 1, top panel). As shown in Table 1 and Fig. 1 (top panel), $\Delta\tau$ taken by the λ_{onc} is greater than that taken by the λ_{one} . Operation at the λ_{onc} can mitigate the random error ($\delta\Delta\tau$) to decrease the required $\text{SNR}(\Delta\tau)$. However, the weighting function of the λ_{one} indicates a moderate peak less than an altitude of 2 km (Fig. 1, bottom panel), and yields a higher sensitivity at lower altitude. Therefore, in the case of Table 1, the difference of $\Delta\tau$ of highly distributed CO₂ at lower altitude to $\Delta\tau$ of constant profile of CO₂ is +1.5% for the λ_{one} , +0.4% for the λ_{onc} . This implies that the λ_{one} can more easily indicate a contrast between urban and vegetated areas. In addition, use of both edge and wing wavelength (5 GHz offset to the center wavelength) provides better surface constraint and $\geq 50\%$ improvement in carbon flux evaluations over vegetated land areas at ~ 500 km resolution for spaceborne measurements (Baker et al., 2011).

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The bias error of $\Delta\tau$ is evaluated from the impact of the elevated particulate layer on the measured bias (τ_{bias}) and the spectroscopic parameter. The bias factor due to spectroscopic parameters is calculated using the Voigt profile function and the uncertainty parameters from earlier studies (Devi et al., 2007; Rothman et al., 2009; Predoi-Cross et al., 2009). Measured τ_{bias} is related to the path-integrated intensity of the aerosol layer or cirrus clouds. A narrow field of view and employment of range detection can allow ground and cloud return signals to be distinguished. Assuming the backscatter coefficients of suburban aerosol data (Sakaizawa et al., 2009), we found that the bias from the integrated backscatter depends on the surface albedo, for example, 0.13 % for a surface albedo of 0.1 sr^{-1} and 0.05 % for an albedo of 0.3 sr^{-1} . This evaluation indicates that higher surface albedo (such as for deserts) can suppress the impact of path-integrated aerosol intensity. Mitigation due to amplitude modulation is described in more detail in Kameyama et al. (2011b).

3 Airborne instruments

We first manufactured an LAS system for ground-based measurements (Kameyama et al., 2011a; Sakaizawa et al., 2009). Figure 2 shows the block diagram of the $1.57 \mu\text{m}$ prototype LAS system and other instrument settings in the aircraft cabin. The specifications for the instruments are summarized in Table 2.

The online and offline sources were polarization-maintained, fiber-coupled diode lasers. The other system is based on optical fiber circuits. Laser-1 (as the center of online wavelength) was stabilized within a root mean square (RMS) value of 12 MHz at the peak of the R(12) line in the $30012 \leftarrow 00001$ absorption band using a gas cell filled with pure CO_2 instead of the reference cavity used in the Pound-Drever-Hall method (Drever et al., 1983). The gas cell was sealed with a gas pressure of 0.1 atm. Laser-2 (as the edge of online wavelength) was stabilized at a position of 2.55 GHz offset from the center position. Laser-1 and Laser-2 were combined using a fiber combiner and detected using a photodiode (InGaAs-PD). The PD generates a heterodyne signal in

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which the beat signal was controlled at a constant 2.55 GHz. Laser-3 was stabilized at offline wavelength within 48 MHz RMS by controlling its temperature and injection currents. The fiber-coupled outputs were amplitude modulated with LiNbO_3 devices. Each modulation signal had a different sinusoidal frequency. The modulated outputs were combined and amplified using a fiber amplifier. Almost all of the amplified power (99 %) was expanded and transmitted through an anti-reflection coated window. The diameter of the transmitted $1/e^2$ beam was 60 mm, and the full angle beam divergence was 120 μrad . The total transmitted power at the fiber end was 1.2 W. The residual 1 % was monitored as a reference for received signals. Scattered signals from the ground surface were collected using a receiving telescope with a field of view of 0.2 mrad and a 110 mm active aperture. The receiving and transmitting optics were fixed on a rigid base plate. The signals were focused on a multi-mode fiber with a 200 μm core diameter and detected using a 0.5 mm diameter InGaAs-PIN PD. The received signals were digitized using a high-speed digitizer (60 MSs^{-1} , 14 bit). The wavelength identification and power evaluation were performed by means of fast Fourier transform (FFT) on a laptop computer.

Airborne in-situ CO_2 measurements were carried out using a module consisting of a commercial CO_2 analyzer (LI-COR Inc., Type: Li-840) modified for airborne operation (Machida et al., 2008). In addition, other trace gases, such as carbon monoxide, methane, etc., were also determined by air analysis using flask sampling devices. Both systems collected air from outside the aircraft using stainless steel sampling tubes facing the direction of flight. The flask sampling was only performed during spiral flight measurements. The time resolution of the in-situ data was 2 s. The precision of in-situ measurements was 0.12 ppm (1σ) in 2 s data. The end-to-end performance was additionally affected by a change in the instrument sensitivity. Consequently, highly accurate calibrated gases were used to compensate the instrumental drift. Hence, the total uncertainty of the in-situ CO_2 measurement was estimated within ± 0.5 ppm.

4 Experiment

Nine flights were conducted for evaluating the value of \bar{q} during August 2009 and February 2010. The aircraft used was a Beechcraft King Air (Type: 200T), Diamond Air Service Inc. and each flight lasted approximately 4 h, including the spiral and level flights. These flights were performed under various conditions, such as over the land and sea, in clear skies, and on partial cloudy days. The flight paths taken are depicted in Fig. 3.

The measurements in August were taken over Hokkaido prefecture in northern Japan, while those in February were taken over the Tsukuba and Koganei sites. The Tsukuba site is approximately 50 km northeast, and the Koganei site is approximately 10 km west of Tokyo. The LAS system provides the $\Delta\tau$ and the range from the aircraft to the target. The amplitude modulation frequencies were 10 kHz for the online wavelength and 11 kHz for the offline wavelength in this measurement. The transmitting online wavelength was set to the edge position of absorption in August 2009 (as shown in Fig. 1), and the center position was used as the online wavelength in February 2010. The accumulation time was 2 s during both measurements. An additional 3 s were required for signal processing on the laptop computer. A visible CCD camera (ARTRAY Inc., Model ARTCAM 150pIII) also monitored the landscape under the aircraft every 5 s. These temporal images were capable of detecting cloud cover over both land and sea.

To validate the LAS measurements, atmospheric CO₂ was taken from 1500 ft (0.5 km) to 23 500 ft (7 km) using flask sampling and in-situ CO₂ devices. Simultaneous radiosonde measurements were carried out by the Japan Weather Association under a contract with the National Institute for Environmental Studies during the spiral flight measurements over the Moshiri site (Hokkaido) in August and the Tsukuba site in February. Other radiosonde measurements at the Koganei site, corresponding to the path of the aircraft, were also performed by the National Institute of Information Communications Technology. Spiral flight measurements were taken over the Tsukuba

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(0.5 km to 2 km) and Koganei (0.5 km to 3 km) sites owing to air traffic control regulations. We employed flask sampling data for validation in August, as the in-situ data were unusable owing to a gas leak in the instrument.

5 Results

Figure 4 graphically illustrates the return signal intensity for the August and February measurements. Various return signals were obtained over grassland, urban areas, and the surface of the sea. The return signals were consistent with z^{-2} . Both offline signals were attenuated by weak CO₂ absorption. It is clear that CO₂ absorption strength varies with online wavelength.

Figures 5 and 6 show the temporal $\Delta\tau$, its fluctuations, and the range from the aircraft to the targets. The graphs at the top of Figs. 5 and 6 depict $\Delta\tau$, while the graphs immediately below it (in both figures) depict $\delta\Delta\tau$. The third graphs from the top depict the optical path length from LAS, the geometrical height from airborne GPS minus the digital elevation model (DEM) from ASTER (Yamaguchi et al., 1998). The graphs at the bottom illustrate the difference between the range obtained from LAS and the geometrical height. The results in August 2009 include two sets of spiral flight measurements. The first set of spiral flight measurements was taken over the Moshiri site (basin in a mountainous area, rough field). The second set was taken over the Taiki site (grassland, plain field). Meteorological data from radiosonde were obtained over the Moshiri site only. The results in February 2010 also included two sets of spiral measurements over the Tsukuba site. Simultaneous radiosonde measurements were also taken.

The values of SNR($\Delta\tau$) at an altitude of 2 km were 147 in August and 270 in February. The corresponding errors due to the value of $\delta\Delta\tau$ were 0.68 % ($2\Delta\tau = 0.18$) in August and 0.37 % ($2\Delta\tau = 0.54$) in February. The error at an altitude of 7 km was 0.85 % ($2\Delta\tau = 0.51$, SNR($\Delta\tau$) = 118) in August 2009. It was found that the SNR($\Delta\tau$) in February (2 km altitude) was 2.5 times greater than that in August (7 km altitude), despite the fact that $\Delta\tau$ was nearly unity. This resulted from attenuated return signal intensities from more

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distant targets. As illustrated in Fig. 4, the return signals at 7 km altitude in August were smaller than those obtained at a 2 km altitude in February. Furthermore, when $\Delta\tau$ is small, it is associated with significant online fluctuation at the edge position. The error due to fluctuation of the operating wavelength ($\partial iwf/\partial\lambda$) was evaluated as being less than 0.58 % at the edge of the online wavelength and 0.05 % at its center. These errors can be reduced by optimizing transmitted laser power, receiving aperture, and detector dark current noise.

To validate LAS altitude, we extracted the geometric height from the on-board GPS and the ASTER-GDEM. Cloud screening was performed in August. The resolution of the DEM was approximately 30 m per pixel, 7–14 m ($= 1\sigma$) vertical precision over a flat field, and 20–30 m over complex terrain, such as mountain slopes (Hirano et al., 2003). The altitude of the aircraft obtained from LAS was consistent with that from GPS-DEM: the difference between the LAS and geometric heights was less than ± 15 m ($1\sigma = 4.9$ m) over a flat field and ± 15 –30 m during rotating movements. The precision of the range measurement and the accuracy of the ground-based measurements were confirmed as 2 m and 5 m, respectively (Sakaizawa et al., 2010). The measured phase difference when calculating the target range was based on averaged return signals coming from groups of trees, buildings, and ground surface over a range from 150 to 200 m in these airborne measurements. The probability of detecting elevations from ground surfaces varies because of the presence of trees over the integration range, and the effective optical path also changes during rotating movement, which may be sources of potential bias in measured mean aircraft altitude. The error due to range measurement ($\partial iwf/\partial z$) was 0.12 %. The bias error due to the spectroscopic data for the CO₂ R(12) line was estimated to be 0.13 % (Devi et al., 2007; Rothman et al., 2009; Predoi-Cross et al., 2009).

$\Delta\tau$ was compared with validation data ($\Delta\tau_{\text{val}}$) calculated from CO₂ concentrations from 1500 ft (0.5 km) to 23 500 ft (7 km). The values of CO₂ concentrations are collected through the airborne in-situ or flask sampling devices (Fig. 7). $\Delta\tau_{\text{val}}$ can be evaluated using the following equation:

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$$\Delta\tau_{\text{val}} = \int_{z_{\text{ac}}}^{z_{\text{trg}}} n_{\text{CO}_2}(r)wf(r)dr \quad (8)$$

where $n_{\text{CO}_2}(r)$ is the dry air mixing ratio from the flask sampling or in-situ data. Figure 8 indicates the linear relation between $\Delta\tau$ and $\Delta\tau_{\text{val}}$. The CO_2 profiles are calculated with a third order polynomial fitting. The CO_2 concentration from the ground to an altitude of 0.5 km was assumed to be constant, due to a lack of surface CO_2 measurements. Note that $\Delta\tau$ and $\Delta\tau_{\text{val}}$ are in agreement, as their correlation coefficient (R) is greater than 0.99.

The graph on the left in Fig. 9 indicates \bar{q}_{LAS} at various elevations, while the graph on the right indicates the difference between the validation data (in-situ and flask sampling) and the measured data. The results indicate a maximum difference of 4 ppm and an averaged difference of 1.5 ppm. \bar{q}_{LAS} for the August measurements shows lower CO_2 levels below 2 km than above 2 km, as seen in the in-situ data in Fig. 7a. \bar{q}_{LAS} for the February measurements shows dense CO_2 below the boundary layer (as shown in Fig. 7b–e) and a tendency to decrease monotonically with height. Note that the August measurements were impacted by photosynthesis in the biosphere, while the February measurements were impacted by a high CO_2 mixing ratio.

The difference between \bar{q}_{LAS} and \bar{q}_{val} can be attributed to bias sources due to aerosol return signals from nearby airplanes, the impact of signal averaging over structured terrain from the ground to the aircraft, and spectroscopic parameters. We evaluated the impact of distributed aerosol on spaceborne CO_2 measurement and found that the bias was less than 0.27 %, as described in Kameyama et al. (2011b). The effect of the other bias factors was evaluated as 0.13 % due to spectroscopic parameters and 0.12 % due to structured terrain (corresponding to range measurement error). The total bias error (τ_{bias}) is at least 0.52 %, which is reasonable compared with the difference between \bar{q}_{LAS} and \bar{q}_{val} .

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In the case of the Tsukuba site, aerosol distributions were measured using a 532 nm ground-based LIDAR from the Meteorological Research Institute. We could not use data analysis of aerosol distribution at other sites, but the atmosphere above Moshiri site is generally clear compared with urban area such as Tsukuba site. The aerosol optical depth (AOD) from an altitude of 0.4 km to 2 km was also evaluated from the LIDAR data. Values of AOD were found to be 0.11 on 14 February, 0.07 on 20 February and 0.12 on 23 February during airborne measurements, for which the corresponding \bar{q}_{LAS} and \bar{q}_{val} at an altitude of 2 km are summarized in Table 3. The difference of \bar{q}_{LAS} to \bar{q}_{val} is -1.5 ppm and evaluated \bar{q}_{LAS} is in agreement with \bar{q}_{val} within the measurement precision of 2.4 ppm (1σ); nevertheless, not only CO_2 concentrations but also aerosols are highly distributed in the lower atmosphere.

The global distribution of AOD values, obtained from space-borne measurements by extraction from the 5 km mesh of the CALIPSO level 2 aerosol layer through the year 2008, ranges from 0.02 to 2. The AOD range without any thick clouds indicated that AOD values of less than 0.12 account for 72 % of the total observed data, while AOD values of less than 0.2 account for more than 84 %. The \bar{q}_{LAS} measurements listed in Table 3 were observed under the above probability for column-integrated AOD, for which the corresponding AOD at 532 nm is partial column-integrated. In addition, the error is considerably smaller than in the case of the airborne measurement where the modulation frequency is higher than 30 kHz (Kameyama et al., 2011b).

6 Conclusions

We demonstrated an airborne measurement system for simultaneous detection of the column-averaged dry air mixing ratio of CO_2 (\bar{q}) and target range using a 1.57 μm laser absorption sensor. The observed differential absorption optical depth and validation data were in good agreement and showed a high correlation coefficient (R) greater than 0.99. The difference between the value of \bar{q}_{LAS} and the validation data \bar{q}_{val} had a maximum value of 4 ppm and an average value of 1.5 ppm. In the dense aerosol

environment over a city, the values of \bar{q}_{LAS} and \bar{q}_{val} were in agreement within the measurement precision of 2.4 ppm, with the corresponding aerosol optical depth in the range 0.07–0.12. In addition, the observed \bar{q} profiles indicated a significant similarity to the validation data. Even though LAS employed a small effective aperture and had a low transmitting laser power, a precision better than 1 % for simultaneous measurements of CO₂ and altitude could be demonstrated. These results indicate that, with an improved system (telescope diameter greater than 1.2 m, total transmitting power of 10 W), a precision better than 0.3 % can be obtained from a space platform.

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Table 1. Estimated $\Delta\tau$ and sensitivity of lower altitude CO_2 for the center and edge wavelength. Atmospheric parameters are based on the AFGL mid-latitude winter. Two CO_2 vertical profiles are assumed, which one of the CO_2 mixing ratio is a constant 385 ppm along the height, the other profile assumes the mixing ratio of 410 ppm on the ground, 398 ppm at 1 km height, and 385 ppm above 2 km height, respectively. Those parameters are 1 km resolution.

	385 ppm constant		Urban area	
	Center	Edge	Center	Edge
$\Delta\tau$ at 7 km	0.970	0.261	0.974	0.265
q at 7 km	385.0	385.0	387.1	398.4

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Table 2. Specifications for instrumental and spectroscopic data.

Transmitter	
Online/ cm^{-1}	6357.31113
Offline/ cm^{-1}	6356.49917
Transmitter power/W	1.2 (Fiber end)
Frequency stability/MHz	On:12, Off:48
Modulation frequency/kHz	On:10, Off:11
Beam diameter/mm	60
Beam divergence/mrad	0.12
Laser spectrum/MHz	0.8
Receiver	
Receiver diameter/mm	110
FOV/mrad	0.2
Detector	InGaAs-PIN
Bandwidth/kHz	200
Detector size/ μm	200
Distinguish on/off	FFT
Integration time/s	2
A/D speed/ MSs^{-1}	1
A/D resolution/bit	14
Spectroscopic data*	
Center wavenumber/ cm^{-1}	6357.31157
Line intensity/ $\text{cm}^{-1}/\text{molec}/\text{cm}^{-2}$	1.6613×10^{-23}
Air-broadening coeff./ $\text{cm}^{-1}/\text{atm}$	0.07781
Air-pressure shift coeff./ $\text{cm}^{-1}/\text{atm}$	-4.30×10^{-3}
Self-pressure shift coeff./ $\text{cm}^{-1}/\text{atm}$	-4.82×10^{-3}
Temperature index	0.695
Lower state energy/ cm^{-1}	60.8709

* (Rothman et al., 2009; Devi et al., 2007; Predoi-Cross et al., 2009)

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Table 3. Partial column-averaged CO_2 from the ground to airplane height and aerosol optical depth from the 0.5 km to 2 km (February 2010). z is the airplane height. 1σ means measurement precision.

	14 Feb	20 Feb	23 Feb
z [m]	1966	1925	1973
Aerosol OD	0.11	0.07	0.12
$q_{\text{LAS}} (1\sigma)$	398.7 (2.4)	398.8 (2.3)	400.6 (2.4)
q_{val}	397.41	397.36	398.85

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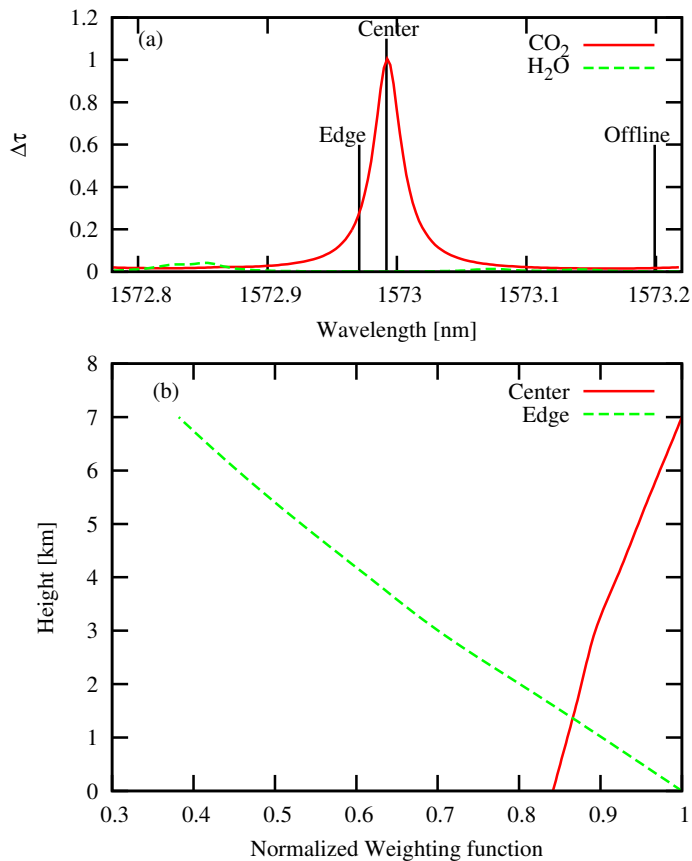



Fig. 1. (a) Operating laser wavelengths and CO_2 optical depth versus wavelength and (b) vertical weighting function dependent on online wavelength. Both plots were calculated using the R(12) line parameters from HITRAN 2008 and some updated data for the two-way path from the ground to an altitude of 7 km. Atmospheric parameters are based on the AFGL mid-latitude winter. CO_2 mixing ratio is assumed to be 385 ppm for all heights.

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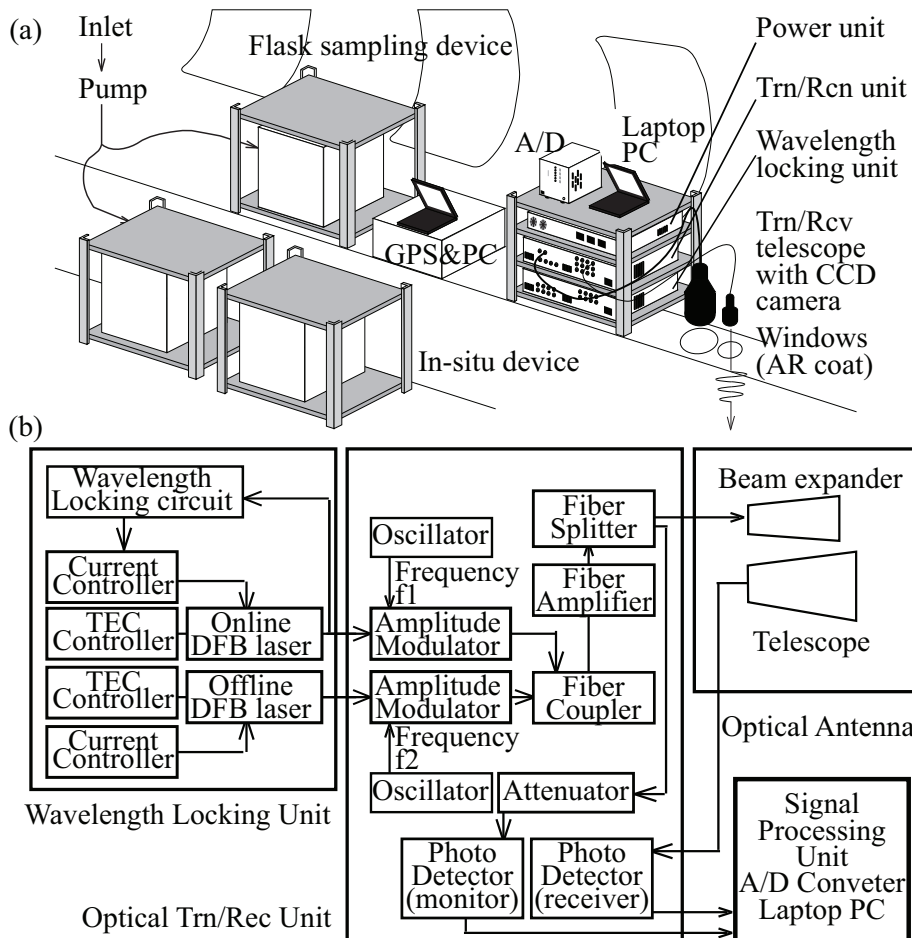


Fig. 2. (a) Airborne instrument setup and (b) block diagram of the 1.57 μm prototype LAS system. TEC: thermo electric cooler, DFB laser: distributed feedback laser.

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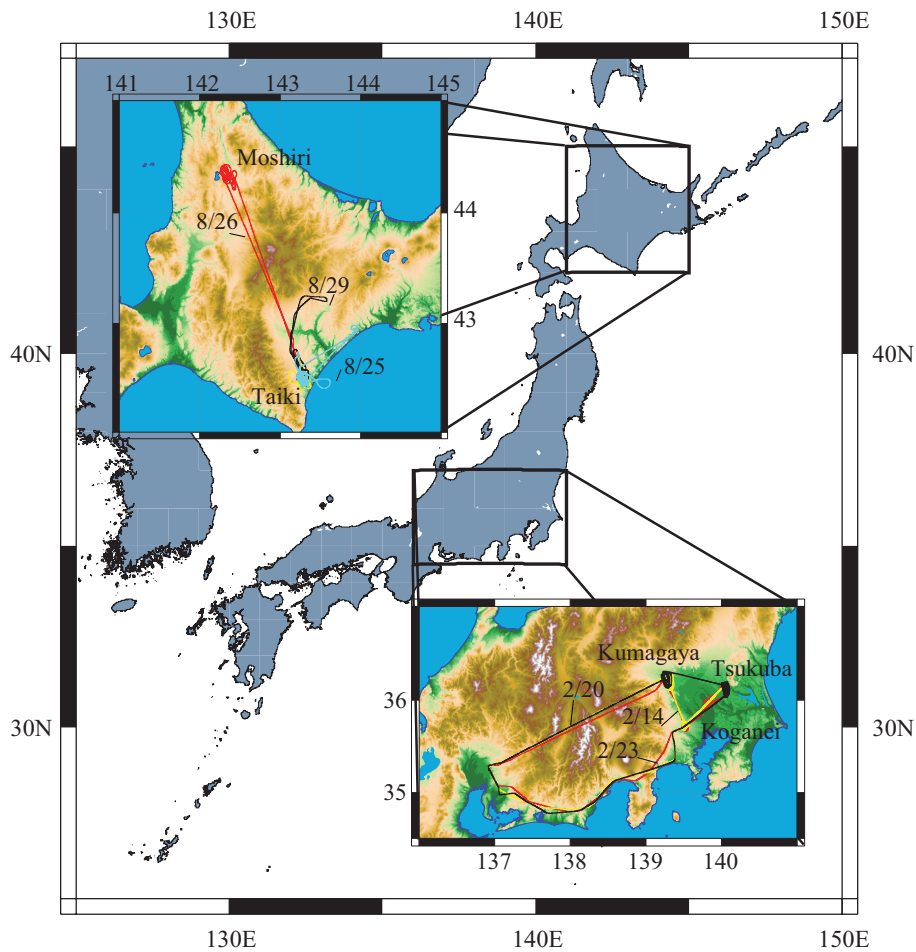


Fig. 3. Flight paths of the aircraft measurements in August 2009 and February 2010.

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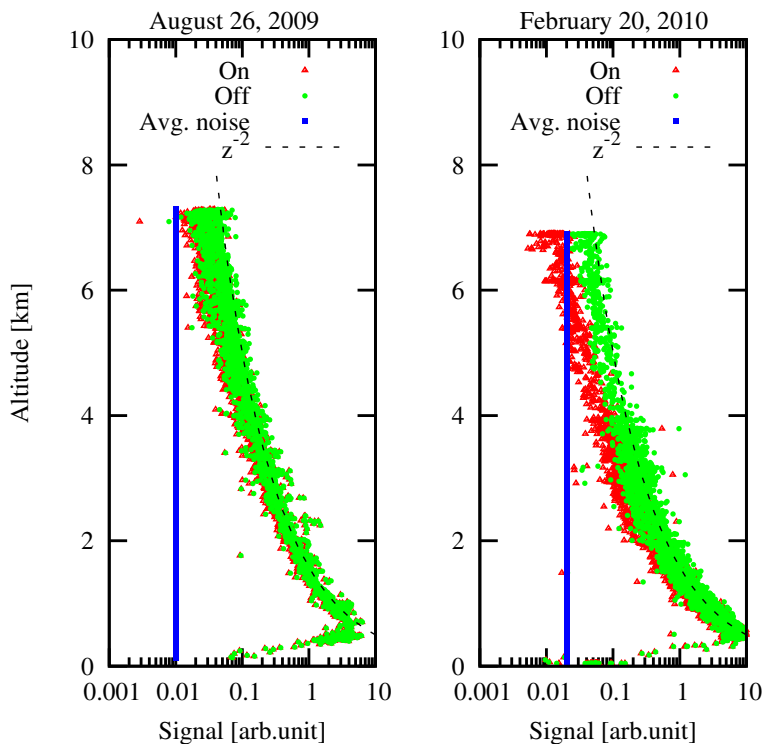


Fig. 4. Received return signals at various elevations in August 2009 (left panel) and February 2010 (right panel). Triangle (Δ): online signal, circle (\circ): offline signal, square (\square): averaged noise level, dashed curve: z^{-2} . The edge and center online wavelengths (Fig. 1) were used in August 2009 and February 2010, respectively.

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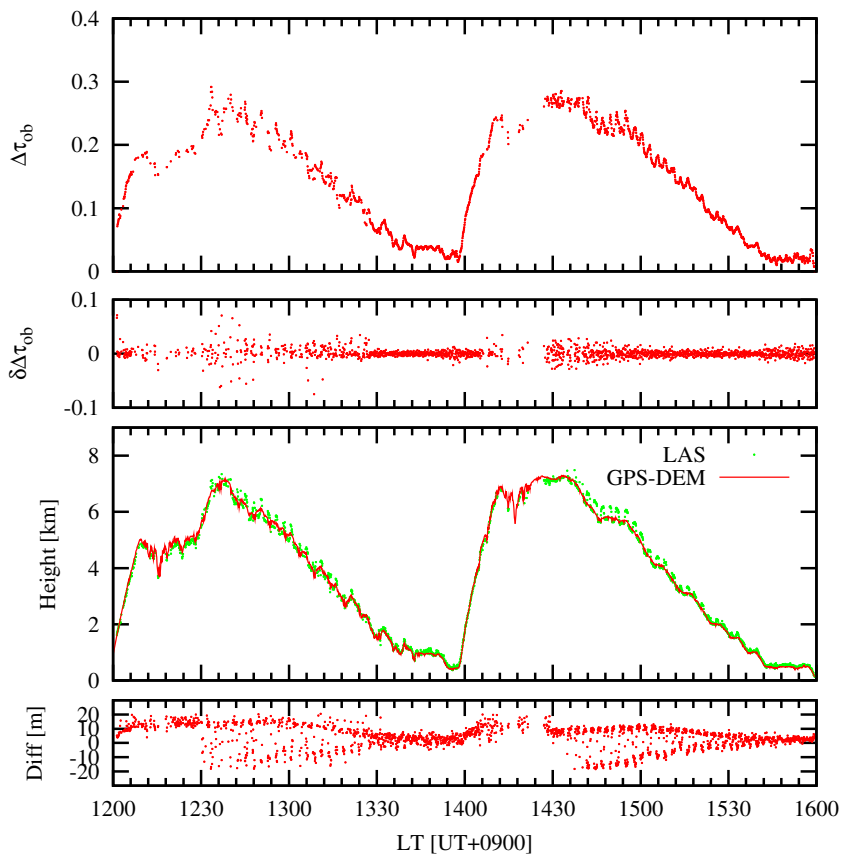


Fig. 5. Results of airborne flight measurements taken on 26 August 2009. The upper two panels show the differential absorption optical depth and its fluctuations versus time, and the lower two panels represent the heights and their differences obtained from LAS and airborne GPS and DEM. The LAS measurement was carried out with cloud screening.

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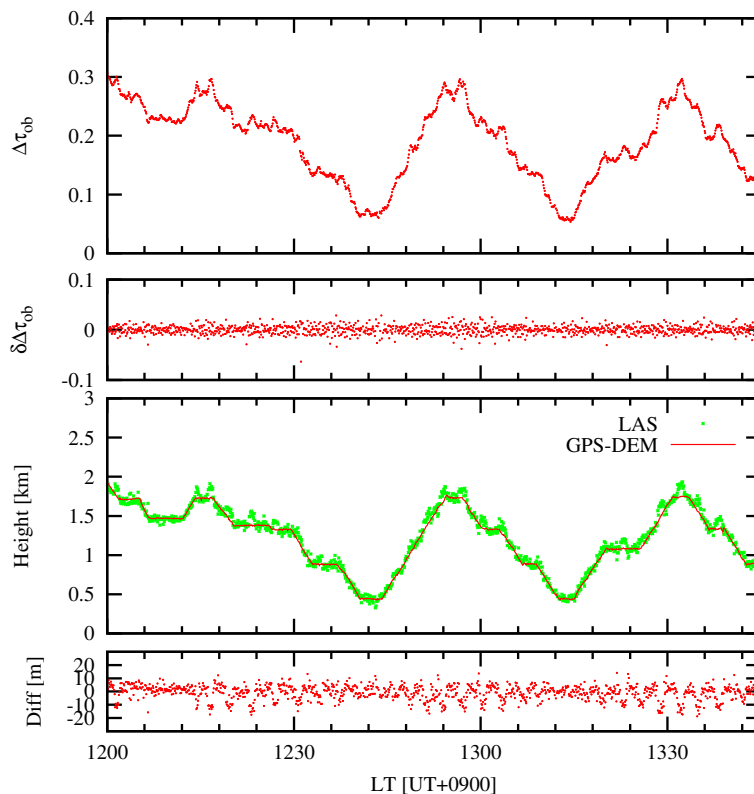


Fig. 6. Results of airborne flight measurements taken on 23 February 2010. These data were taken over the Tsukuba site (urban area).

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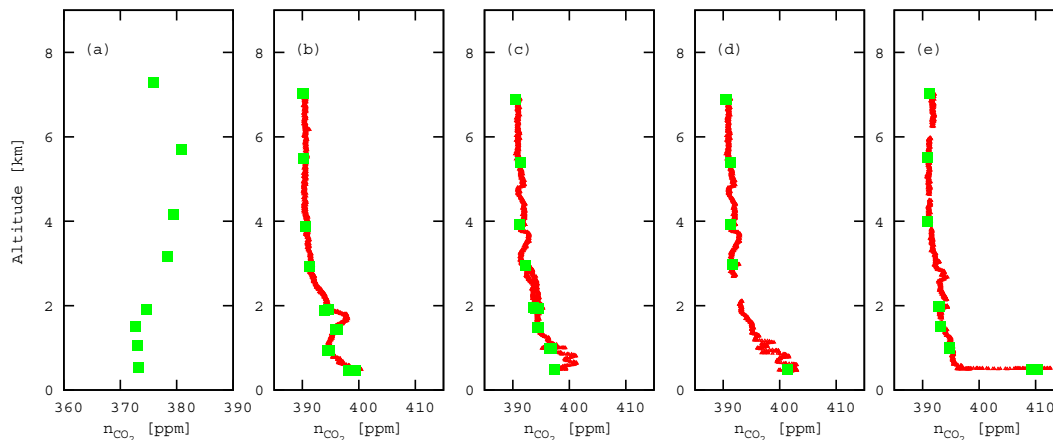


Fig. 7. Atmospheric CO_2 profiles from airborne flask (green squares) and in-situ devices (red points). **(a)** 26 August 2009, taken over Moshiri (0–7 km); **(b)** 14 February 2010; and **(c)** 20 February 2010, over Kumagaya (2–7 km) and Tsukuba (0.5–2 km); **(d)** 20 February 2010 over Kumagaya (2–7 km) and Koganei (0.5–3 km); **(e)** 23 February 2010, over Kumagaya (2–7 km) and Tsukuba (0.5–2 km).

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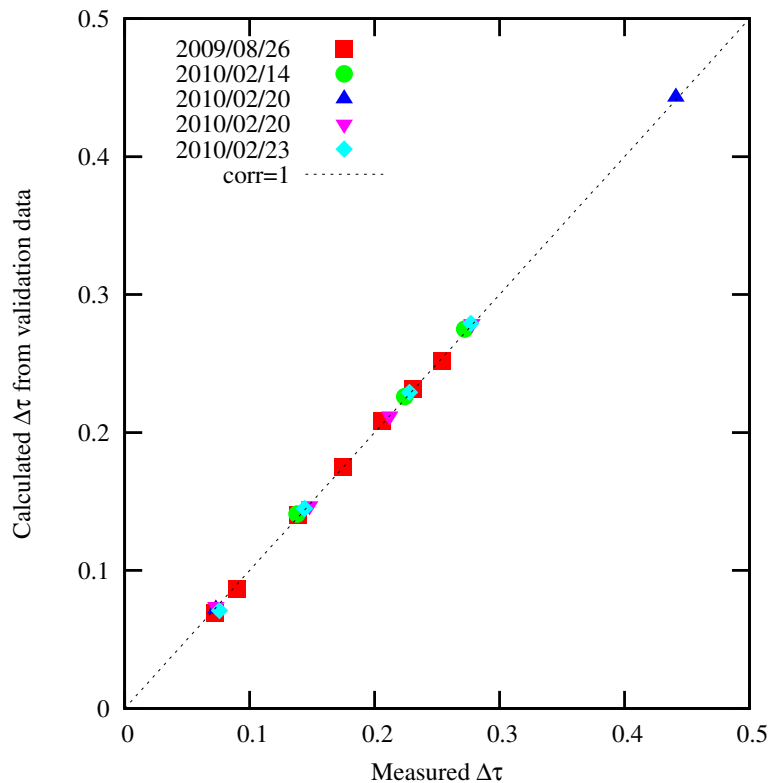


Fig. 8. Correlation between measured and calculated differential absorption optical depths ($\Delta\tau$).

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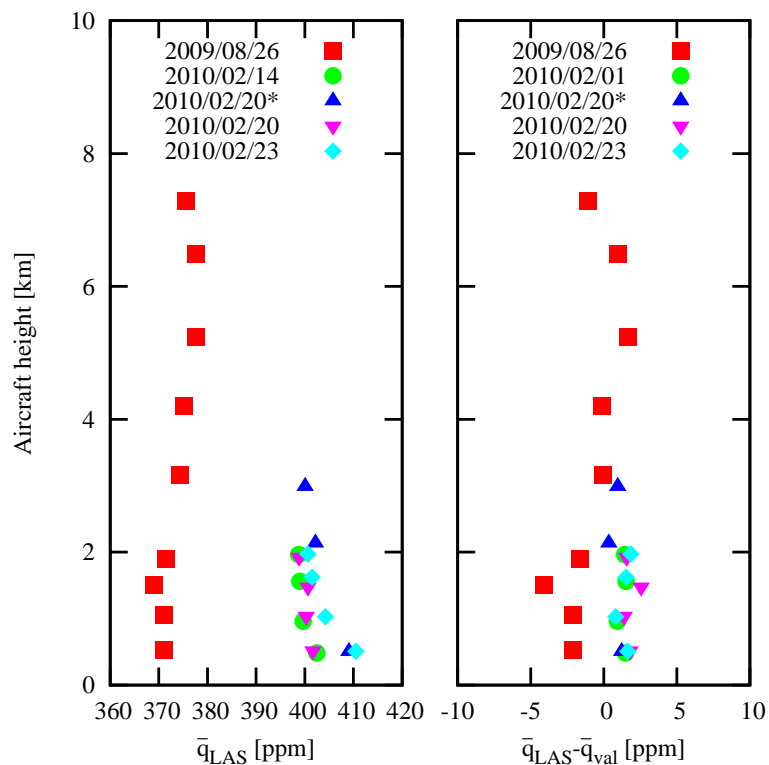


Fig. 9. Evaluated partial column-averaged column CO_2 (\bar{q}_{LAS}) obtained from LAS (left panel) and the difference compared with calculated values (\bar{q}_{val}) from flask and in-situ measurements (right panel).

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