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Improving accuracy and precision of ice core $\delta D(CH_4)$ analyses using methane pre- and hydrogen post-pyrolysis trapping and subsequent chromatographic separation

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Abstract

Firn and polar ice cores offer the only direct paleoatmospheric archive. Analyses of past greenhouse gas concentrations and their isotopic compositions in air bubbles in the ice can help to constrain changes in global biogeochemical cycles in the past.

- ⁵ For the analysis of the hydrogen isotopic composition of methane $(\delta D(CH_4))$ 0.5 to 1.5 kg of ice was previously necessary to achieve the required precision. Here we present a method to improve precision and reduce the sample amount for $\delta D(CH_4)$ measurements on (ice core) air. Pre-concentrated methane is focused before a high temperature oven (pre pyrolysis trapping), and molecular hydrogen formed by pyroly-
- ¹⁰ sis is trapped afterwards (post pyrolysis trapping), both on a carbon-PLOT capillary at -196 °C. A small amount of methane and krypton are trapped together with H₂ and must be separated using a short second chromatographic column to ensure accurate results. Pre and post pyrolysis trapping largely removes the isotopic fractionation induced during chromatographic separation and results in a narrow peak in the mass
- ¹⁵ spectrometer. Air standards can be measured with a precision better than 1 ‰. For polar ice samples from glacial periods we estimate a precision of 2.2‰ for 350 g of ice (or roughly 30 mL (at standard temperature and pressure (STP)) of air) with 350 ppb of methane. This corresponds to recent tropospheric air samples (about 1900 ppb CH₄) of about 6 mL (STP) or about 500 pmol of pure CH₄.

20 1 Introduction

25

Methane (CH_4) is a potent greenhouse gas showing increased atmospheric concentrations since the industrial revolution (IPCC, 2007). A recent assessment of the present day methane budget is presented in Kirschke et al. (2013). However, the atmospheric load of CH_4 has varied on various time scales. A wealth of information has been gained from concentration measurements regarding annual (Dlugokencky et al., 1995), decadal (Mitchell et al., 2011), millenial up to glacial/interglacial (Loulergue et al., 2008)



 CH_4 variability. Stable isotope data of methane on recent air samples (e.g. Quay et al., 1999) and on the past atmosphere using ice cores (e.g. Ferretti et al., 2005; Fischer et al., 2008; Sowers, 2010; Sapart et al., 2012; Möller et al., 2013) provide further insight into processes and sources controlling the global methane cycle. For instance,

- ⁵ the temporal evolution of the hydrogen isotopic composition of methane ($\delta D(CH_4)$) over the termination of the last ice age (14 000–18 000 yr before present) (Sowers, 2006) as well as rapid warming events between 32 000–42 000 yr before present (Bock et al., 2010b) made it possible to reject the "clathrate gun hypothesis" proposed by Kennett et al. (2003) as the trigger for the steep atmospheric methane increases. Bock et al.
- ¹⁰ (2010b) could also show that the precipitation signal ($\delta D(H_2O)$), which changes from cold stadial to warm interstadial conditions, is traced into the paleo hydrogen isotopic signature of methane.

However, we are still far from a complete picture of the biogeochemistry of methane in the past. Ice core isotope studies on $\delta D(CH_4)$ have the potential to improve our ¹⁵ understanding of the global CH₄ cycle but are still scarce due to analytical difficulties (e.g. Bock et al., 2010a; Sapart et al., 2011) and the large sample amount needed. To date the few published ice core $\delta D(CH_4)$ studies used 0.5 kg (Bock et al., 2010b) and more than 1 kg (Sowers, 2006, 2010; Mischler et al., 2009) of ice from multi parameter deep ice cores with a typical precision of around 3 to 4 ‰. This study presents new ²⁰ developments based on (Bock et al., 2010a) to improve precision and accuracy and reduce the sample size for (ice core) $\delta D(CH_4)$ measurements significantly.

2 Experimental

We present an improved continuous-flow gas chromatography (GC) pyrolysis (P) isotope ratio monitoring mass spectrometry (irmMS) system (GC/P/irmMS) designed to analyze $\delta D(CH_4)$ from (ice core) air samples (Fig. 1) with high precision. In the following we give a short summary of our previous instrumentation (Bock et al., 2010a) and new developments concerning the physical system and data processing.



The most important new features presented here are pre and post pyrolysis trapping (pre&postPT) of CH₄ and molecular hydrogen (H₂), respectively, and post pyrolysis gas chromatographic separation using a Porous Layer Open Tubular (PLOT) column, which improve accuracy and precision and reduces the required sample amount considerably.

A systematic dependency of the δD(CH₄) values on the amount of CH₄ (signal dependency) is observed but can be precisely corrected for. The successful implementation of pre&postPT demands for better purification of helium as a prerequisite. At the same time we created a new calibration software tool which enables a one step correction of system drifts over time and signal dependency (linearity) in an iterative way.

10 2.1 Instrumentation

The system is fed by helium (He) (Alphagaz I, 99.9990 % purity; Carbagas, Switzerland), which is purified by a high-capacity gas purifier and an inline gas purifier (both Supelco, Bellefonte, PA, USA). In addition to the description given in Bock et al. (2010a), we further purify the He used for the complete system in a 3m long 1/4" stainless steel (SST) capillary (i.d. 5.3 mm) filled with charcoal (grain size 0.3–0.5 mm, 0.41 g cm⁻³, Fluka 05112; Sigma-Aldrich Chemie, Steinheim, Germany). This trap is immersed in liquid nitrogen (LN) during the week and can be vented at room temperature during weekends using valve V0 in Fig. 1, a pneumatic six-port, two-position valve (1/16" fittings, 0.4 mm port diameter, Valcon M rotor; Valco, VICI AG, Schenkon,

- ²⁰ Switzerland). The additional He purification cold trap preceding the system lowers the blank CH_4 contribution considerably compared to Bock et al. (2010a). Extracting remnant gas in the sample cylinder (containing melt water) after an ice sample results in a H_2 peak 0.5% the area of a 200 g B34 ice sample (mimicking the CH_4 amount of the last glacial maximum (LGM) about 20 ka before present (BP), where present is defined
- as 1950 (Loulergue et al., 2008; Clark et al., 2009)). Trapping only He on T2 bypassing the sample cylinder results in a peak area of only 0.2%. Thus, blank CH_4 contribution appears to be reduced by a factor of 2 compared to the old set-up.



The following steps are similar to Bock et al. (2010a): in a nutshell, a glass vessel containing an ice core sample is evacuated, and the enclosed air is released upon melting. In a high flow (He, 500 mLmin⁻¹) water vapour is removed using a cooled Nafion membrane and a cold trap (T1) while the air sample is transferred to a trap filled ⁵ with charcoal (T2) immersed in LN.

Contrary to Bock et al. (2010a), T1 is made up of an empty 1/8" tube of 3 coils that enter or leave a dewar maintained at -90 °C. Temperature controlled cooling of the dewar is achieved using LN droplets projected into the dewar (Schmitt, 2006; Bock et al., 2010a). T1 now removes only residual water vapour, while CO₂ is already adsorbed on an Ascarite trap, made of a 10 cm 1/4'' stainless steel tube. In this new set-up N₂O is passed through the system to be measured in the mass spectrometer.

Air reference injections are realized by switching V1, either mimicking an ice sample by introducing the air into the glass extraction vessel or by bypassing the sample vessel. depending on the position of V2. Following a switch of V3, the air sample is transferred

from the charcoal trap to a trap filled with Hayesep D (T3, at -100° C), where methane 15 is quantitatively trapped, while the bulk air (N_2, O_2, Ar) is vented. Residual air components and CH_4 are focussed on T4 (-196 °C) and injected onto a GC column. Valve V5 is switched to route the sample through a new cold trap (T5, -90°C), replacing Nafion-2 of the old set-up, and towards the pyrolysis furnace only for the time window

in which CH_4 is leaving the GC column. 20

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In the following we describe the main new developments. Eluting CH_4 from the GC is focused on T6 for 18 s (pre pyrolysis trapping, prePT) before it is released by passive warming to room temperature. Subsequently, CH_4 is pyrolysed as described in Bock et al. (2010a), but the produced H_2 is not allowed to enter the mass spectrometer directly. Instead, H₂ is trapped on T7 for 40 s (post pyrolysis trapping, postPT). Both 25 traps – T6 and T7 – are U-shaped, 20 cm long GC columns (GS-CarbonPLOT, ID 0.32 mm, film 1.5 μm, Agilent Technologies, part number 113-3112) retaining CH₄ and H₂ at LN temperature. After post pyrolysis trapping is complete, T7 is lifted out of LN and warmed to room temperature allowing H_2 to enter the mass spectrometer via an



open split. After T7 the line has been extended by two meters of the same PLOT column (2nd GC in Fig. 1) at room temperature to separate H_2 from other gases, which produce signals visible in both the m/z 2 and 3 traces (Figs. 2 and 3). A short piece of larger inner diameter capillary can be added to broaden the peaks. Valve V6 is also used to bypass the pyrolysis furnace and traps T6 and T7 in order to vent water eluting from a warm T5.

Note that in an early version, our system was only extended by post pyrolysis trapping, while pre pyrolysis trapping was implemented later. As CH₄ and CDH₃ are held on T6, this pre pyrolysis trapping step resets the time shift introduced by the GC (Meier-Augenstein, 1999; Bock et al., 2010a) allowing pyrolysis of a non fractionated methane peak. Furthermore prePT allows for a shorter post pyrolysis trapping time as H₂ from CH₄ pyrolysis elutes during a shorter time interval. This is advantageous because H₂ cannot be held on T7 for an extended time under the described conditions. Strong chromatographic separation between H₂ and HD in a system using only postPT leads to strong intra-peak fractionation visible in a large time shift. This strong chromatographic effect makes the peak very sensitive to biases during ionisation and peak integration leading to amount effects or interferences.

Pyrolysis of CH₄ is achieved in a custom-made high-temperature furnace (Bock et al., 2010a). A brand new thermocouple suggested that optimal pyrolysis temperature is 1400 °C in our case. However due to ageing of the thermocouple, the read out of the temperature is reduced considerably over months. To determine the optimal pyrolysis temperature we introduce pure CH₄ peaks via V7 (e.g. 3 times 10 µL loop

with roughly 500 ppb CH₄ in He) at different temperatures. We observe two plateaus at different temperatures, one for $\delta D(CH_4)$ values and one for peak areas. The plateau

of $\delta D(CH_4)$ at higher temperature is favoured for high precision isotope measurements because the small but inevitable temperature fluctuations in the reactor then lead to smaller scatter in isotope numbers. Too high temperatures lead to shortened lifetimes of pyrolysis reactors which becomes noticeable through higher backgrounds of nitrogen and argon caused by ambient air. Typically, a reactor (stone-ware GmbH,



Switzerland, DEGUSSIT[®] Al_2O_3 , length = 420 mm, ID = 0.5 mm, OD = 1.5 mm) facilitates reproducible results for about four months. When a new reactor has to be installed, it is heated up using a ramp of 5 h and pre-conditioned over a day injecting 10 µL loops every 3 min (without using any trap).

- As a second major improvement to the system following a development by Schmitt et al. (2013b) we can now measure N₂O concentration, δ^{15} N and δ^{18} O of N₂O and xenon on the same sample. Therefore, after the H₂ acquisition for methane is completed, the pyrolysis reactor is bypassed using valve V6 and a peak jump is performed in order to tune the mass spectrometer to the N₂O configuration measuring *m*/*z* 44, 45,
- ¹⁰ 46 using the triple collector. After a second peak jump we measure xenon (as 132 Xe²⁺ and 136 Xe²⁺) using beams m/z 66 and 68. Xenon is considered a proxy for total air content and is used to calculate CH₄ and N₂O concentrations. For detailed descriptions of N₂O and Xe analytic we refer the reader to a companion publication by Schmitt et al. (2013b) reporting on a new system to simultaneously measure δ^{13} CH₄, isotopes of N₂O, Xe and other trace gas concentrations.

2.2 Data processing

We use custom made Python (http://www.python.org/) scripts to process the raw beam data, to organize peak data of reference, standards and samples in specific libraries and to perform the calibration to the international VSMOW (Vienna Standard Mean Ocean Water) scale. The peak integration method is similar to that described in Bock et al. (2010a). Integration limits are found based on the major beam time series and also applied to the minor beam. Instead of detecting the integration limits according to a slope threshold, we determine the peak maximum and set the peak start and end points to fixed numbers of data points before and after the peak maximum, i.e. we use
²⁵ a fixed peak width. In contrast to our previous procedure, pre&postPT removes the

²⁵ a fixed peak width. In contrast to our previous procedure, pre&postP1 removes the isotopic fractionation induced by the chromatographic separation resulting in nearly unfractionated H₂ peaks in the current set-up. Hence, we no longer perform a time



shift correction of the m/z 3 beam, nor do we shift the left integration limit according to a peak size dependent value to correct for signal dependency as we did for our old system (Bock et al., 2010a). Generally, the background is determined as the median of the data points 6 s before the peak start (see Fig. 2).

- In order to calibrate samples, it is essential to compare samples to standard measurements that are sufficiently stable over time and match the sample size. If this cannot be achieved, one has to correct for any drift and signal dependency (e.g. Schmitt et al., 2003; Potter and Siemann, 2004; Bock et al., 2010a; Brass and Röckmann, 2010). In our case this is essential, because we observe a clear signal dependency
- ¹⁰ of the $\delta D(CH_4)$ values (Fig. 4). The signal dependency, however, is stable and reproducible over long time intervals and can therefore be precisely corrected for without compromising the overall precision of the measurement (see Sect. 3.2). When a new pyrolysis reactor is installed, the signal dependency changes, and a new interval of our data analysis has to be started to account for this change. We developed a new
- ¹⁵ tool to correct for any system (time) drift and signal dependency at the same time, which is presented in detail in the appendix of this article. It takes standard measurements of known isotopic signature and iteratively fits parameters for (temporal) drift and signal dependency at the same time in order to minimize the standard deviation of $\delta D(CH_4)$ of our reference air. This assumes constant signal dependency within a cer-
- tain time period (typically some weeks). The same assumption holds for laboratories determining signal dependency on a periodic schedule, but we see two advantages of our approach: (1) no extra day is needed to examine signal dependency and (2) if signal dependency changes slightly during the chosen time interval, this change is already accounted for by our standard measurements covering this interval. The fit pa-
- rameters and daily mean values of our reference "Air Controlé" are used to calibrate the samples.



3 System performance

3.1 Accuracy

Our reference used to calibrate all samples is "Air Controlé", a recent clean air tank $(CH_4 \text{ concentration} = [CH_4] = 1971 \pm 7 \text{ ppb})$ for medical purposes (bottle 541659, filled

- ⁵ February 2007 in Basel, Switzerland, Carbagas). "Air Controlé" was cross-referenced to -93.6 ± 2.2% with respect to Vienna Standard Mean Ocean Water (wrt VSMOW) using bottled air from Alert station "Alert 2002/11" (Bock et al., 2010a; Poss, 2003). According to thorough analyses performed at the University of Heidelberg, Germany, the Centre for Ice and Climate at the Niels Bohr Institute of the University of Copenhagen,
- Denmark and the Max Planck Institute in Jena, Germany, we are confident that our values exhibit a deviation from the VSMOW scale smaller than 3.5 ‰ (I. Levin, P. Sperlich, W. Brand, personal communication, 2013, and Sperlich et al., 2012).

We furthermore introduce a new standard gas here: "Saphir 4" (bottle 4405, Carbagas, artificial clean air mixture with 761 ppb CH_4 and no krypton) and ice core samples

of a core dry drilled next to the EPICA (European Project for Ice Coring in Antarctica) drill site in Dronning Maud Land, Antarctica (EDML, 75°0.15′ S, 00°4.104′ E, 2892 ma.s.l.), called "B34". The air occluded in this ice has an approximate age of 1500 a BP.

Figure 6 shows $\delta D(CH_4)$ measured on B34 ice samples on a depth scale. The same data are also presented in Table 2, which additionally shows the measurement date and the weight of samples. Overall, we are confident that the described system was stable in terms of accuracy over the past few years. However, for polar ice samples (B34) it turns out that significant variations in $\delta D(CH_4)$ can occur within small depth ranges. Note for example the interval between 190 and 191 m where duplicates (samples from exactly the same depths) indicate good reproducibility, while going down

²⁵ pies from exactly the same depths) indicate good reproducibility, while going down core a depletion of several permil becomes obvious, which is significantly larger than the measurement error.



Mean values for WAIS (Antarctica) and B30 (Greenland) from similar (pre industrial) time periods (around 410 and 670 a BP, respectively) are -73.0% and -91.5%; hence, we estimate an inter polar difference of 18.5% with a combined error of 1.9% (the square root of the sum of the squared standard deviations of samples and ref-

⁵ erence measurements). This is a larger difference compared to the assessment by Sowers (2010), who reconstructed 12 ± 6 ‰ for 550 to 960 a BP from the WAIS and GISP2 (Greenland Ice Sheet Project 2) ice cores, however still within the combined measurement uncertainties. Lower $\delta D(CH_4)$ values for the North are expected because high latitude methane emissions from boreal wetlands, thermokarst lakes and thawing permafrost are strongly depleted in deuterium and almost exclusively located

in the Northern Hemisphere (Walter et al., 2008).

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Saphir injections through the melt water of a previously extracted ice core sample are slightly depleted in deuterium (by 1.7%) compared to Saphir injections bypassing the sample container, but the mean values are within the combined error. It is not clear whether this offset prevails for ice samples or if the effect only occurs after ice sample

extractions. We therefore chose not to correct for this (potential) offset.

Note that our results for WAIS (West Antarctic Ice Sheet, core WDC05A, tube 184, depth range: 172.74–173.03 m, age approximately 410 a BP) are 15% more enriched in deuterium compared to data presented in Mischler et al. (2009). This offset is similar

- to the one observed for Boulder air (Bock et al., 2010a) compared to measurements performed at the Stable Isotope Lab of the Institute of Arctic and Alpine Research (IN-STAAR, University of Colorado, Boulder, CO, USA) and shows that the laboratories in the US and Europe are tied to different primary standard air bottles. Note that no internationally accepted isotope reference material for CH₄ from air samples is yet available.
- ²⁵ At the time of writing, the mentioned lab offsets are being addressed in a round robin organized by T. Sowers and E. Brook using WAIS ice and bottled air samples with varying methane concentrations.

Post pyrolysis trapping and subsequent gas chromatographic separation enables the measurement of a pure H_2 peak in the mass spectrometer. Recently Schmitt et al.



(2013a) demonstrated that there is krypton (Kr) interference during carbon isotopic analyses of CH₄. Conventional stable isotope analysis of CH₄ using GC-IRMS separation without post-conversion trapping leads to insufficient separation of CH₄ and Kr. In case of δ^{13} CH₄, the influence of Kr on δ^{13} C analysis leads to a significant alteration of the results. Meier-Augenstein et al. (2009) reported interference of N₂ for H₂ analyses. Based on that, we assess in the following the origin and influence of peaks showing up on *m*/*z* 2 and 3 after the CH₄-derived H₂ peak, subsequently referred to as post peaks (Figs. 2 and 3).

5

Focussing our mass spectrometer to m/z 43 and detecting Kr as 86 Kr²⁺ demonstrated that Kr enters the mass spectrometer (similar to Schmitt et al., 2013a) but is well separated from CH₄-derived H₂ due to the second chromatographic separation (Fig. 1). Methane derived hydrogen peaks are integrated over 2.5 s and have a full width at half maximum (FWHM) of 1 s, while Kr arrives in the mass spectrometer 4 s later. The FWHM for Kr is similar as for CH₄-derived H₂.

¹⁵ Focussing the mass spectrometer to m/z 16 and measuring clean CH₄ injections with the pyrolysis oven at normal operation temperature (1400 °C) and at room temperature shows that also a small amount (ca. 1 %) of injected CH₄ passes the hightemperature oven unpyrolized. This CH₄ eluting from T7 has a retention time 3.3 s longer than H₂, thus unpyrolized CH₄ and Kr overlap considerably in the ion source

²⁰ (Fig. 3). We highlight that the described post peaks enter the mass spectrometer after the CH₄-derived H₂ peak in m/z 2 and 3. Therefore, we can exclude any interference during the IRMS measurement on $\delta D(CH_4)$ due to the described post peaks. Although the pyrolysis of CH₄ is not 100 %, no bias is expected for the reported calibrated $\delta D(CH_4)$ values as we strictly follow the "identical treatment principle" of samples and ²⁵ references (Werner and Brand, 2001) and, thus, correct for any potential fractionation

during the pyrolysis step. For our H₂ source settings, the unpyrolyzed CH₄ peak is seen on m/z 2 and 3, but most clearly on the minor trace (blue in Fig. 3) for clean CH₄ injections and Saphir 4, which is an artificial clean gas mixture (see above, no Kr). We speculate that fragments



of CH₄ reach the Faraday cups for m/z 2 and 3, producing the positive peaks observed on both beams 3.3 s after the CH₄-derived H₂ peak. One second later, samples containing krypton show a positive peak on m/z 2 while m/z 3 drops below the background level. Apparently, the presence of Kr draws down the m/z 3 trace just after the start of the first post peak (i.e. CH₄). We can only imagine of krypton-induced sputter ions that reach the Faraday cup for m/z 2, thus producing a positive peak, while for the case of the negative peak on the minor trace we speculate that Kr interferes with the (He-) background such that the electrostatic filter before the m/z 3 cup hinders slow background ions from reaching the collector cup. It is important to note that the observed signal reduction on m/z 3 is larger than our artificially added H₂ background (of < 0.05 nA on m/z 2, which is lower as in Bock et al. (2010a), where the level was

- increased by ca. 1 nA); thus, the signal drop at m/z 3 is clearly not due to the H₂background but more likely due to the He background. While we stress that this effect does not have an influence on our $\delta D(CH_4)$ analyses due to the good post pyrolysis
- ¹⁵ gas chromatographic separation, we speculate that the observed effect may also occur in a system where H₂ and Kr are simultaneously present in the mass spectrometer, potentially leading to biases in $\delta D(CH_4)$ as a function of the CH₄/Kr ratio as is the case for $\delta^{13}CH_4$ (Schmitt et al., 2013a). Note that no post peaks are observed for the blank measurement presented in Fig. 3d, indicating no measurable contamination that can
- ²⁰ be traced back to leaks. Therefore, any valid explanation of the described phenomena has to be based on sample derived matter.

Given that the described peaks show up only after postPT and a second separation step downstream were implemented, one might fear offsets in terms of accuracy of $\delta D(CH_4)$ for the old system. However, we found our system to be astoundingly stable over the past several years without pre&postPT (until 2010), using only postPT (2011 and 2012) and after the implementation of pre and post pyrolysis trapping (2013). Table 1 summarizes our results and shows no significant differences for Boulder, NAT332, Dome6 or B30 Greenland ice samples as compared to our previous set-up (Bock et al., 2010a). The only exception is Saphir 3, an artificial air sample containing no krypton

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(Table 1), which is now measured with a -5.8‰ offset and with a combined error of 3.5‰ (determined as above). Although this offset is still within 2σ of the error, we speculate that this is related to the Kr effect described above for measurements without post pyrolysis GC separation. The effect for Saphir is maximal as we standardize a non Kr and medium CH₄ containing sample (Saphir) with a standard containing recent Kr and CH₄ concentrations ("Air Controlé") (see also Schmitt et al., 2013a). As our ice core samples also contain Kr, the effect is much smaller, as reflected in the very good correspondence of analyses with and without pre&post pyrolysis trapping and subsequent GC separation (Table 1).

10 3.2 Precision

In this section we describe the improvements concerning precision and sample size due to pre and post pyrolysis (pre&postPT) trapping of methane and hydrogen, respectively. In our old system (without pre&postPT) a typical sample (up to 500 g of polar ice with CH_4 concentrations between 350 and 700 ppb) showed peak heights of the major beam between 0.6 and 1.3 nA (for ice core samples presented in Bock et al., 2010b). While peak areas are still in the same range for identical amounts of CH_4 , major peak heights are increased roughly fourfold due to postPT. For B34 ice core samples between 200 and 450 g (with a CH_4 concentration of roughly 630 ppb), we now obtain peak heights between 1.5 and 4.3 nA. To mimic glacial low CH_4 content only 200–220 g samples of B34 ice (corresponding to 350–400 g with 350 ppb CH_4 as

²⁰ only 200–220 g samples of B34 ice (corresponding to 350-400 g with 350 ppb CH₄ as found for the last glacial maximum) were used and are listed in Table 2.

Table 1 summarizes our isotope results for air standards and ice samples. It is clear that precision of the new set-up is improved as indicated by smaller standard deviations of air standards (1.8% or better) and pooled standard deviations of ice core samples (2.2% or better) determined according to

 $_{25}$ (2.3 ‰ or better), determined according to



$$\sigma_{p} = \sqrt{\frac{\sum_{i=1}^{k} (n_{i} - 1)\sigma_{i}^{2}}{\sum_{i=1}^{k} (n_{i} - 1)}}.$$

Note that the precision is comparable for small and large B34 samples (Table 2: e.g. samples < 220 g ice: $1\sigma = 2.1\%$). Based on the pooled standard deviation of B34 samples from the same depths, we estimate that our system's precision for ice samples is ⁵ around 2.2 ‰. Note that with this method $\delta D(CH_4)$ of present day tropospheric air can be measured with a precision better than 1 ‰ on 18–40 mL (STP) samples (Table 2). We estimate that most of the gain in precision of the improved system is due to pre and postPT and only a small fraction can be attributed to our data processing routine. We assessed this by re-evaluating the standard measurements of our data set presented in Bock et al. (2010b) with the new python routine explained in the appendix. The standard deviation of all "Air Controlé" measurements using the new tool is 2.5 ‰ compared to 2.8% using the old procedure. Note, however, that even the smallest peaks of the old batch were larger by a factor of 1.6 compared to the tiny peaks that can now be measured with comparable precision using pre and post pyrolysis trapping. Furthermore, we acquired several runs of clean CH₄/He injections of varying methane amounts with and without pre&postPT as a second measure of the gain in precision. For the old system we obtained a standard deviation of 2.0% for peak areas between

- 1.8 and 13.3 nA. The smallest peaks between 1.8 and 3.0 nA could be measured with a precision of 2.6 ‰. After introducing pre&postPT we are able to achieve a precision of 1.6‰ for even smaller peaks between 1.3 and 1.7 nAs. As seen for the re-evaluated "Air Controló" measurements for larger peaks the gain is amaller; indicated by a step
- "Air Controlé" measurements, for larger peaks the gain is smaller: indicated by a standard deviation of 1.4% for peak areas between 1.3 and 6.9 nAs. We conclude that the new data processing tool presented here represents an efficient and robust way to handle time drifts and signal dependency in one step but the main benefit consider-
- ²⁵ ing precision is attributed to the implementation of pre and post pyrolysis trapping of methane and hydrogen, respectively.



4 Conclusions

We presented pre and post pyrolysis trapping of methane and hydrogen, respectively, combined with post trapping GC separation on a PLOT column to improve accuracy and precision and reduce sample amount in $\delta D(CH_4)$ analysis of atmospheric and ice

⁵ core samples. We showed that the precision for 350 g of ice (or roughly 30 mL of air) with 350 ppb of methane is approximately 2.2 ‰. This corresponds to recent tropospheric air samples (roughly 1900 ppb CH₄) of about 6 mL (STP) or about 500 pmol of pure CH₄. Vice versa 30 mL (STP) samples with recent tropospheric CH₄ concentration can be determined with a precision of better than 1 ‰. Compared to our old set-up (Bock et al., 2010a) this translates into improvement factors for sample size (350 g)/(500 g) and precision (2.3 ‰)/(3.4 ‰) of 0.7.

We note, however, that the high standard in accuracy and precision for such small samples is achieved at the cost of measurement time; the new set-up allows the analysis of only one to two ice samples or four atmospheric samples a day.

¹⁵ We showed that potentially the accuracy of systems without pre&postPT and subsequent chromatographic separation can be biased depending on pyrolysis efficiency and varying methane/krypton ratios in samples and the reference. However, for atmospheric samples (preindustrial ice and air samples) the updated method did not measurably change in terms of $\delta D(CH_4)$ values compared to our initial set-up described in ²⁰ Bock et al. (2010a).

Appendix A

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Correction for system drifts and signal dependency (linearity)

In order to calibrate samples measured on any isotope system, it is essential to compare samples to standard measurements that are sufficiently stable in time and match the sample size, or correct for any drift and signal dependency. As amount effects alter



isotopic results simultaneously with (time) drift effects, both errors should be corrected at the same time and not consecutively. A decoupling of the corrections is only possible when standards of constant peak size are measured to monitor the time trends only. Effects of signal dependency can be assessed by performing standard runs of different peak sizes, which is time consuming. Hence, we present an approach which allows simultaneous corrections of system drifts and signal dependency effects. For optimum conditions we choose size matching and bracketing standards for individual samples and pool standards measured over several days (assuming constant signal

dependency over this time period) to cover the samples' size range. To correct for both signal dependency and drift effects, we use the following approach.

Any measured isotope value δX^{meas} is composed of the true value δX^{true} , any signal dependency, which is a function of peak area *A*, and a drift correction, which is a function of time *t*

$$_{\epsilon} \delta X^{\text{true}} = \delta X^{\text{meas}} - f^{\text{lin}}(A) - f^{\text{drift}}(t).$$

In the following, signal dependency is characterized by a polynomial of order N

$$f^{\rm lin}(A) = \sum_{n=1}^N x_n A^n.$$

System drift is decomposed into two additive terms

$$f^{\text{drift}}(t) = f_1^{\text{drift}}(t) + f_2^{\text{drift}}(t).$$

The first term is a drift during a day, which is fitted to a polynomial of order M

$$f_1^{\text{drift}}(t) = \sum_{m=1}^M y_m t^m \Theta(t - [\overline{t_i} + \Delta t]) \Theta([\overline{t_i} - \Delta t] - t)$$
(A4)

(A1)

(A2)

(A3)

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dependent on time *t*. Therein t_i represents the time at day *i* the current sample was measured at. Thus, $\overline{t_i}$ describes the mean measurement time of all samples measured during one day. Since temporal system drifts occur typically on time scales of weeks to months, the size of the drift within a day is usually small. Accordingly, our software allows for the calculation of $\delta D(CH_4)$ values with or without a diurnal drift correction (usually this is our preferred setting). The Theta-function Θ is zero, if its argument is < 0 and one if its argument is > 0. This efficiently allows for the determination of the drift for each single measurement day in the program code. To discriminate between two adjacent laboratory days, Δt is defined as 0.4 days. The number of standard data points for each day should be larger but at least of the same size as *M*.

The second term represents the drift of the reference values between days. The mean isotopic reference signatures of all days are assumed to change in a stepwise linear fashion

$$f_2^{\text{drift}}(t) = \sum_{i=1}^{L} (m_i t + n_i) \Theta(\overline{t_i} - t) \Theta(t - \overline{t_{i-1}}),$$
(A5)

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where m_i quantifies the slope and n_i the intersection with the ordinate at a measuring day *i*, and *L* is the number of all measurement days.

Slope and intersection for each day *i* are calculated with respect to the previous day i - 1.

$$m_{i} = \frac{\overline{t_{i}} - \overline{t_{i-1}}}{\left[\overline{\delta X_{i}^{\text{meas}} - f^{\text{lin}}(A_{i})}\right] - \left[\overline{\delta X_{i-1}^{\text{meas}} - f^{\text{lin}}(A_{i-1})}\right]}$$

$$n_{i} = \left[\overline{\delta X_{i}^{\text{meas}} - f^{\text{lin}}(A_{i})}\right] - m_{i}\overline{t_{i}}.$$
(A6)
(A7)

Herein, influences of signal dependency have to be corrected for before calculating the mean standard isotopic signal of each day.



We can express all quantities given in Eq. (A1) as functions of on peak area A and isotopic signature δX^{meas} at every measured point in time *t*. The true value of the standard δX^{true} is known. Thus, Eq. (A1) can be used to fit all measured data points. The fit parameters determine both the signal dependency and the drifts during and between

⁵ the days by minimizing the standard deviation of all drift and signal dependency corrected standard values. The fitted parameters are then used to ultimately calibrate the samples by minimizing the standard deviation of all standard values.

Our routine is written in Python (www.python.org). The actual optimization uses the function scipy.optimize.fmin(). Figures 4 and 5 are produced by our routine and show uncalibrated and calibrated data, respectively. Each figure shows signal dependency in the left panel and time drift in the right panel.

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Table 1. Precision of the new $\delta D(CH_4)$ system. Mean values are given in column 3; columns 4 and 5 show standard deviations (1 σ) of samples and "Air Controlé" reference air measurements, respectively. Columns 6–8 show values obtained with the previous set-up presented in Bock et al. (2010a). "Air Controlé" measurements are used to calibrate the samples to the international VSMOW scale. "N" represents the number of measurements used. Ice sample results are not corrected for any firn diffusion process. Gas ages of the ice samples are estimated as follows: B30: 670 aBP, B34: 1500 aBP, WAIS: 410 aBP. The WAIS samples are from core WDC05A, tube 184, depth range: 172.74–173.03 m. NGRIP gas samples date from between 870 and 9000 aBP.

	this study				Bock et al. (2010a)		
sample description (sample size, origin CH₄ concentration)	N	δD(CH ₄) (‰)	1 <i>σ</i> sample (‰)	1 <i>o</i> reference (‰)	N	δD(CH ₄) (‰)	1 <i>σ</i> sample (‰)
Air reference and samples Air Controlé (all injections, 4–40 mL) Air Controlé (only larger loops (18–40 mL) Saphir 4 ([CH ₄] = 761 ppb) Saphir 4 (loop after sample) Saphir 3 ([CH ₄] = 1004 ppb) Boulder (CAO8289 [CH ₄] = 1500 ppb) NAT-332 ([CH ₄] = 2141 ppb) Dome 6 (firn air [CH ₄] = 1718 ppb)	544 69 36 34 2 14 3 2	-93.6 -93.6 -171.6 -173.2 -173.4 -81.0 -108.0 -71.0	1. 0. 1.2 1.4 0.4 1.1 1.8 0.8	3 8 0.9 0.9 0.7 0.8 0.2	343 86 0.9 18 8 6 2	-93.6 -93.5 -167.6 -80.8 -106.3 -71.0	2.8 2.3 2.4 1.3 1.2 0.1
B30 (Greenland, preindustrial, depth range 2 m) WAIS (Antarctica, preindustrial, parallel replicates) B34 (Antarctica, late Holocene, depth range 9 m) Ice core replicates B34 ice (parallel replicates, late Holocene) NGRIP (bag replicates of gas cut, Holocene)	2 4 47 35 27	-91.5 -73.0 -74.6 depth intervals 16 13	0.8 0.5 2.8 pooled 1 <i>o</i> 2.2 2.3	1.1 1.2 1.5	14	-94.7	-3.7



Table 2. Results of ice core samples from B34. Given depth is the middle of each sample. Depending on replicate shape and weight typical samples are between 5 and 15 cm long. The standard deviation of Air Controlé measurements used to calibrate the sample is given in the column named 1 σ . Samples with a weight < 220 g correspond to a methane amount comparable to samples from the glacial with lowest CH₄ concentration of around 350 ppb.

middle depth	measurement date	δD(CH ₄)	1σ	weight
(m)	(Day Month Year)	(%)	(‰)	(q)
181 435	20 Feb 2013	-74.8	0.9	219.6
181.553	22 Feb 2013	-75.6	0.9	200.4
181,935	21 Feb 2013	-76.7	0.9	199.4
181 935	12 Mar 2013	-75.0	0.9	200.4
183.065	2 Nov 2011	-79.0	1.9	202.3
183.065	3 Nov 2011	-73.6	1.9	209.7
183 128	8 Nov 2011	-74.1	2.6	206.5
183 190	15 Feb 2012	-76.2	27	316.7
183,190	3 Nov 2011	-76.6	1.9	202.2
183.315	7 Mar 2012	-77.1	1.6	261.3
183.315	16 Feb 2012	-79.2	2.7	210.4
183,445	1 Nov 2011	-75.5	1.9	377.0
183,445	22 Feb 2012	-78.8	3.3	264.4
183.570	8 Dec 2011	-77.2	2.1	312.6
183.695	31 Oct 2011	-78.4	1.9	394.3
183.825	30 Sep 2011	-78.8	1.2	349.1
183.825	28 Sep 2011	-74.6	1.2	376.5
183.830	11 Sep 2013	-76.8	1.2	254.6
183.830	18 Jul 2013	-75.5	1.4	293.3
183.830	20 Aug 2013	-73.6	1.8	298.8
183.830	13 Sep 2013	-73.3	1.2	267.6
183.945	10 Feb 2012	-73.7	2.7	318.5
184.060	5 Aug 2013	-75.1	1.7	232.2
184.380	26 Jul 2013	-71.6	1.4	224.4
184.550	12 Jul 2013	-70.9	0.9	250.6
184.550	19 Aug 2013	-72.2	1.8	259.0
184.730	13 Jun 2013	-76.3	1.2	211.6
184.730	7 Jun 2013	-75.2	1.2	257.0
184.910	18 Jun 2013	-70.9	1.1	251.5
184.910	6 Jun 2013	-74.2	1.2	239.8
185.190	18 Jul 2011	-74.5	1.2	223.8
185.190	18 Jul 2011	-77.5	1.2	212.9
185.338	15 Jul 2011	-73.9	1.2	446.0
185.338	15 Jul 2011	-73.3	1.2	427.6
185.500	26 Apr 2011	-75.3	0.4	427.0
186.985	19 Mar 2013	-72.3	0.9	408.6
187.133	19 Feb 2013	-71.6	0.9	310.0
187.420	24 Oct 2012	-70.9	1.9	347.4
187.420	23 Oct 2012	-74.2	1.9	340.7
190.475	31 May 2013	-68.7	1.2	261.7
190.475	4 Jun 2013	-68.0	1.2	251.4
190.565	30 May 2013	-70.1	1.2	262.3
190.565	22 May 2013	-71.3	1.2	250.7
190.655	17 May 2013	-76.9	1.1	264.4
190.775	5 Jun 2013	-80.2	1.2	205.1
190.775	15 May 2013	-73.9	1.0	229.4
190.775	8 May 2013	-72.9	0.5	245.1



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Fig. 1. Flow scheme of the new $\delta D(CH_4)$ system including pre and post pyrolysis trapping of methane and hydrogen, respectively, and a short second chromatographic separation column (2nd GC). The orange areas highlight the major differences from Bock et al. (2010a).





Fig. 2. Chromatogram of a CH_4 derived H_2 peak of an ice core sample (B34). Peak integration limits are shown as red crosses. The data range used for the background determination is shown as magenta bullets and the subtracted background values as red line. Also indicated is the data range after the sample peak where we observe two post peaks, namely not pyrolized methane and krypton (compare text and Fig. 3).





Fig. 3. Chromatograms showing CH_4 -derived H_2 peaks and post peaks (methane and krypton). Peak data are normalized to the H_2 maximum, which is given in (nA) in the top left corner of each subplot. Shown are **(a)** a clean CH_4 injection (via V7), **(b)** Saphir-4 (artificial clean air mixture (no Kr and 761 ppb CH_4)), **(c)** Boulder (natural air with reduced CH_4 (1500 ppb) and ambient Kr), **(d)** pure helium extraction of an evacuated dry sample cylinder, **(e)** B34 ice (ambient Kr, about 630 ppb CH_4 , dry drilled next to EDML, the EPICA (European Project for Ice Coring in Antarctica) drill site in Dronning Maud Land, Antarctica: gas age is 1500 aBP and **(f)** EDML ice (ambient Kr, about 370 ppb CH_4): gas age is 18.2 ka BP.





Fig. 4. Uncalibrated $\delta D(CH_4)$ data of our reference Air Controlé (black bullets) and B34 ice core samples (green squares). The left panel shows $\delta D(CH_4)$ vs. major area, i.e. the observed signal dependency. The magenta line indicates the polynomial correction function for signal dependency. The right panel shows the same uncalibrated data plotted against time. The red dots and line show the fitted standard numbers, which are later used to calibrate the samples, as described in the Appendix. The closer the red and black symbols are to each other, the better the fit.





Fig. 5. Calibrated $\delta D(CH_4)$ data of Fig. 4 after correction for temporal drift and signal dependency. The left panel shows $\delta D(CH_4)$ vs. major area and the right panel shows the same data plotted against time. Black bullets show standard measurements (Air Controlé), green squares show B34 ice core samples.





Fig. 6. $\delta D(CH_4)$ of B34 ice core samples measured during the years 2011 to 2013 on a depth scale. Error bars represent the pooled standard deviation of B34 replicates (2.2 ‰) for $\delta D(CH_4)$ and the total depth range of each individual sample.

