



**Simultaneous
retrieval of effective
refractive index and
density**

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Simultaneous retrieval of effective refractive index and density from size distribution and light scattering data: weakly absorbing aerosol

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Abstract

We propose here a novel approach for retrieving in parallel the effective density and real refractive index of weakly absorbing aerosol from optical and size distribution measurements. Here we define “weakly absorbing” as aerosol single-scattering albedos that exceed 0.95 at 0.5 μm . The required optical measurements are the scattering coefficient and the hemispheric backscatter fraction, obtained in this work from an integrating nephelometer. The required size spectra come from a Scanning Mobility Particle Sizer and an Aerodynamic Particle Sizer. The performance of this approach is first evaluated using a sensitivity study with synthetically generated but measurement-related inputs. The sensitivity study reveals that the proposed approach is robust to random noise; additionally the uncertainties of the retrieval are almost linearly proportional to the measurement errors, and these uncertainties are smaller for the real refractive index than for the effective density. Next, actual measurements are used to evaluate our approach. These measurements include the optical, microphysical, and chemical properties of weakly absorbing aerosol which are representative of a variety of coastal summertime conditions observed during the Two-Column Aerosol Project (TCAP; <http://campaign.arm.gov/tcap/>). The evaluation includes calculating the root mean square error (RMSE) between the aerosol characteristics retrieved by our approach, and the same quantities calculated using the conventional volume mixing rule for chemical constituents. For dry conditions (defined in this work as relative humidity less than 55 %) and sub-micron particles, a very good (RMSE \sim 3 %) and reasonable (RMSE \sim 28 %) agreement is obtained for the retrieved real refractive index (1.49 ± 0.02) and effective density (1.68 ± 0.21), respectively. Our approach permits discrimination between the retrieved aerosol characteristics of sub-micron and sub-10-micron particles. The evaluation results also reveal that the retrieved density and refractive index tend to decrease with an increase of the relative humidity.

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problem using density values obtained from other methods, such as alignment methods.

Alignment methods are designed to estimate the effective density from size distributions measured simultaneously by different aerosol sizing instruments (e.g., Hand and Kreidenweis, 2002; Khlystov et al., 2004). The solution is achieved by merging two measured size distributions with different size ranges into a single combined spectrum with a wider size range. During this merging process, the measured spectra are compared in the instruments' size overlap region. The best-fit solution is determined by finding the minimum of a merit function, which is often proportional to the squared difference between the two measured size distributions in the overlap region. The design of these schemes is driven by selection of the spectrum type and may include data consisting of both mobility and aerodynamic sizes (Khlystov et al., 2004). For spherical singly charged particles, the mobility size is equal to the geometric size D_p , while the aerodynamic size D_a is linked to the geometric size D_p through the well-known relationship: $D_p = D_a / \sqrt{\rho}$ (e.g., Baron and Willeke, 2001, p.51). The distinguishing feature of the alignment methods is the use of the effective density to convert the aerodynamic size into the geometric size. Note that measurements of the mobility and aerodynamic size distributions in tandem with the optical size spectra enable simultaneous retrieval of the effective density and real refractive index of weakly absorbing aerosol (Hand and Kreidenweis, 2002). The term "weakly absorbing" will be used throughout our paper for describing situations when values of aerosol single-scattering albedo are large (>0.95). These values are typical for marine aerosol (e.g., Fujitani et al., 2007; Kassianov et al., 2013; Berg et al., 2014).

In this work we put forth a novel technique for inferring in parallel the effective density and real refractive index of weakly absorbing aerosols from simultaneously measured size distributions (with mobility and aerodynamic sizes), and two measured optical properties, namely the scattering coefficient (σ_s) and hemispheric backscatter fraction (β). The backscatter fraction is defined as the ratio of the hemispheric backscattering and total scattering coefficients. The choice of these two widely measured optical

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properties is motivated by their different sensitivities to the effective density, as explained in Sect. 2 below. The inclusion of these two optical constraints (σ_s and β) coupled with measured size distributions enables estimation of the effective density and real refractive index simultaneously. Moreover, it is well-known that the aerosol optical depth (τ_a) and asymmetry factor (g) of weakly absorbing aerosols are the most important aerosol optical parameters for numerous climate-relevant applications, including the Earth's radiation budget (e.g., McComiskey et al., 2008). It is widely recognized that τ_a and g are strongly related to the scattering coefficient and hemispheric backscatter fraction, respectively (e.g., Andrews et al., 2006). Therefore, the estimation of n and ρ can be extended to the calculation of these two climate-relevant aerosol optical properties. Finally, our technique allows for determining the difference between the retrieved aerosol characteristics of sub-micron ($D_p < 1 \mu\text{m}$) and sub-10-micron ($D_p < 10 \mu\text{m}$) particles. Given that there is a dearth of routinely deployed instruments that provide information on large ($1 \mu\text{m} < D_p < 10 \mu\text{m}$) particles, our technique gives an opportunity for obtaining this important information from widely used observations.

In Sect. 2 we provide details of our method. In Sect. 3 we evaluate the performance of our approach through a sensitivity study utilizing simulated but measurement-based inputs. In Sect. 4, we describe an integrated dataset for its evaluation. This dataset includes the optical, microphysical, and chemical properties of weakly absorbing aerosol collected during the Two-Column Aerosol Project (TCAP; Berg et al., 2014), which was conducted with support from the US Department of Energy's (DOE's) Atmospheric Radiation Measurement (ARM) Program (<http://www.arm.gov/>). Also, we compare aerosol characteristics retrieved by our method with those obtained from the conventional volume mixing rule applied to chemical composition measurements. We summarize the evaluation results in the last section.

2 Approach

As mentioned previously, our retrieval approach utilizes two complementary and simultaneously obtained datasets. The first dataset includes two observational quantities: the total scattering coefficient ($\sigma_{s, \text{obs}}$) and hemispheric backscatter fraction (β_{obs}) measured by a TSI 3563 integrating nephelometer. Throughout this work we use the subscript “obs” to indicate observed quantities. Although the nephelometer provides the light scattering measurements at three wavelengths (0.45, 0.55 and 0.70 μm), here we only consider aerosol optical properties at a wavelength of 0.55 μm . The spectral dependence of the retrieved real refractive index will be the subject of our future studies. The second dataset includes observed particle number distributions measured by a TSI Scanning Mobility Particle Sizer (SMPS) with a 0.01–0.48 μm electrical mobility size range, and a TSI 3321 Aerodynamic Particle Sizer (APS) with a 0.52–19.8 μm aerodynamic size range. TCAP data from these instruments are discussed in Sect. 4.

The general concept behind the simultaneous retrieval of n and ρ is enumerated as follows:

(1) *Shifting the APS number distribution to form a combined size distribution.* By changing the effective density, the APS number distribution is shifted horizontally along the abscissa $\log D_p$ – a procedure similar to the alignment method (Fig. 1a). The SMPS spectrum remains unchanged. Such shifting of APS number distribution modifies the relative contribution of particles with moderate (0.5–0.7 μm) and large (>1 μm) sizes to the light scattering. These relative contributions as a function of D_p control the total scattering coefficient σ_s and hemispheric backscatter fraction β quite differently. For example, a reduction in the number of particles with moderate size, which scatter sunlight most efficiently, results mainly in a decrease of the total scattering coefficient σ_s . In contrast, a reduction in the number of particles with large size, which are generally responsible for the forward scattering of sunlight, results mostly in an increase of hemispheric backscatter fraction β . These distinct dependences of σ_s and β form the basis of our retrieval. We obtain combined SMPS-APS distributions for a set of assumed ρ

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values within a wide range ($1.0 \leq \rho \leq 2.6$), which represents reasonable values of ρ appearing in the literature (e.g., Barnard et al., 2010).

(2) *Calculate modeled optical properties from combined size distribution.* Next we apply a technique similar to the optical closure methods (e.g., Liu and Daum, 2000; Mack et al., 2010) by using the size distributions found in step (1) above, and Mie theory, to calculate modeled optical properties $\sigma_{s, \text{mod}}$ and β_{mod} (Fig. 1b and c); here the subscript “mod” refers to calculated, modeled quantities. These properties are found for an assumed imaginary refractive index (0.001) and the range of ρ values mentioned above and for a set of assumed n . For n , we specify a range in values ($1.3 \leq n \leq 1.7$) that is representative of atmospheric aerosol (e.g., Barnard et al., 2010). The calculated optical properties can be considered as two-dimensional (2-D) functions of the two independent variables considered here (n and ρ) (Fig. 1b and c). The near-orthogonal isolines seen in these 2-D diagrams of $\sigma_{s, \text{mod}}$ and β_{mod} suggest that a unique solution for n and ρ can be obtained. For example, isolines in Fig. 1b and c correspond to constant values of $\sigma_{s, \text{mod}}$ and β_{mod} . At a given time, observed values of these variables can be associated with a single isoline in each figure. By superimposing these two isolines over one another, and finding their intersection, n and ρ are determined. However, the issue of “ripple” in Fig. 1c must be addressed to find a unique solution.

(3) *Finding a unique solution using observed optical properties.* Once observed values of the optical properties are fed to the algorithm, the process of finding a unique solution for n and ρ needs to address distinctive features of these 2-D diagrams: “smooth” (Fig. 1b) vs. “rippled” (Fig. 1c) structures for $\sigma_{s, \text{mod}}$ and β_{mod} , respectively. These features are attributable to the different responses of $\sigma_{s, \text{mod}}$ and β_{mod} to the relative contributions of moderate and large particles. Unlike β_{mod} , $\sigma_{s, \text{mod}}$ is obtained via integration of the scattered sunlight over all angles, and this full angular integration washes out any ripple structure. The apparent difference between smooth and rippled structure is consistent with results from previous studies (e.g., Chamaillard et al., 2006; Chartier and Greenslade, 2012), which have examined differences between the backscattering and total scattering efficiencies as a function particle size.

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We emphasize that finding a unique solution for a domain with a rippled structure is often problematic because, in our case, the isolines can intersect in more than one place. To circumnavigate this possible problem, we select two unique populations of points from these 2-D diagrams. The first of these defines points where the observed and modeled values of the scattering coefficient agree within a given threshold (e.g., 3%). The second population represents similar points for the backscatter fraction. These two populations of points are approximated by the corresponding polynomial fits (curves) and the solution for n and ρ is found by locating the intersection of these fits. Illustrative examples of such fits and the determination of the unique solution are given in the next section.

3 Sensitivity study

To illustrate the performance of our method, we devise a sensitivity study using the SMPS-APS size distributions and assumed aerosol characteristics, n_{mod} and ρ_{mod} , which we treat as reference values. These assumed characteristics are used to calculate aerosol optical properties, $\sigma_{\text{s, mod}}^{\varepsilon}$ and $\beta_{\text{mod}}^{\varepsilon}$. Artificial errors – devised to mimic measurement uncertainties – are added to the optical properties, and these properties along with the added errors are treated as “observations”, which serve as input for our retrieval. Here and below we use the superscript “ ε ” to indicate that errors have been added. We then retrieve aerosol characteristics from these “observed” optical properties. Finally, we compare the “retrieved” characteristics ($n_{\text{mod}}^{\varepsilon}$ and $\rho_{\text{mod}}^{\varepsilon}$) with their reference values (n_{mod} and ρ_{mod}).

3.1 Simulation of input data

We assume that the diurnal changes of the original aerosol characteristics can be approximated by simple functions with sine terms:

$$\begin{aligned} \rho_{\text{mod}} &= \rho_0(1 \pm \sin(t)) \\ n_{\text{mod}} &= n_0(1 \pm \sin(t)) \end{aligned} \quad (1)$$

The amplitudes ρ_0 and n_0 are selected is such way that the span of modeled optical properties (Fig. 2a and b) matches roughly the wide range of optical properties observed during the TCAP study. Note that in our sensitivity study we generate “observations” for a 24 h period to cover a full diurnal cycle.

To estimate the sensitivity of our method to measurement uncertainties, we add “observational” errors (ε_σ and ε_β) to the corresponding “error-free” modeled optical properties:

$$\begin{aligned} \sigma_{\text{s, mod}}^\varepsilon &= \sigma_{\text{s, mod}}(1 \pm \varepsilon_\sigma) \\ \beta_{\text{s, mod}}^\varepsilon &= \beta_{\text{s, mod}}(1 \pm \varepsilon_\sigma) \end{aligned} \quad (2)$$

These errors are related to the uncertainties in the light scattering measurements and are defined for two extreme cases. The first case represents a situation where the errors have the same magnitudes and signs (positively correlated errors): $\sigma_{\text{s, mod}}^\varepsilon = \sigma_{\text{s, mod}}(1 + \varepsilon)$, $\beta_{\text{s, mod}}^\varepsilon = \beta_{\text{s, mod}}(1 + \varepsilon)$. The second case describes situation when these errors have the same magnitudes and but different signs (negatively correlated errors): $\sigma_{\text{s, mod}}^\varepsilon = \sigma_{\text{s, mod}}(1 + \varepsilon)$, $\beta_{\text{s, mod}}^\varepsilon = \beta_{\text{s, mod}}(1 - \varepsilon)$. The assumed error ε is distributed uniformly over the interval $(\alpha, -\alpha)$, where α is the relative magnitude (0 %, 5 %, or 10 %), defined as a percentage of the “error free” optical properties. Note that the uncertainties in the light scattering measurements are about $\pm 10\%$ (e.g., Anderson et al., 1996).

Figure 2 (top panel) shows examples of simulated optical properties ($\sigma_{\text{s, mod}}^\varepsilon$ and $\beta_{\text{s, mod}}^\varepsilon$) for two extreme cases with different magnitudes of the observational errors,

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while Fig. 2 (middle panel) and Fig. 3 illustrate the sensitivity of the solution to these errors. There are small differences between the aerosol characteristics retrieved for the “error-free” and “noisy” conditions with *positively* correlated errors (Figs. 2d and 3a and b). This encouraging result indicates clearly that the obtained solution (n^ε and ρ^ε) depends weakly on the positively correlated observational errors (ε_σ and ε_β). This sensitivity becomes relatively stronger for the “noisy” conditions with *negatively* correlated errors (Figs. 2c and 3a and c). However, the error-induced differences between the retrieved aerosol characteristics for the “error-free” and “noisy” conditions do not exceed 10% and 35% for the real refractive index and density, respectively (Fig. 3a and c). Note that illustration considered here represents an “unfavorable” case with large error-induced differences. On average, these differences are smaller (Fig. 2c and d).

3.2 Comparison of retrieved and original aerosol characteristics

To assess this sensitivity quantitatively, we calculate differences between the retrieved and original characteristics for each 1 h time step within the 24 h diurnal cycle:

$$\begin{aligned}\Delta_\rho &= \rho_{\text{mod}} - \rho_{\text{mod}}^\varepsilon \\ \Delta_n &= n_{\text{mod}} - n_{\text{mod}}^\varepsilon\end{aligned}\quad (3)$$

We calculate standard deviations of the above difference for different values of α for the same 24 h period, and in Fig. 2 (bottom panel) we display the statistics $\text{SD}(\Delta_\rho)$, $\text{SD}(\Delta_n)$ and $\text{SD}(\varepsilon)$ as a function of α magnitude.

Differences of the retrieved aerosol characteristics, Δ_ρ and Δ_n , and the standard deviations of these differences increase almost linearly with the magnitude of the errors. This increase is more pronounced for the negatively correlated errors than for the positively correlated errors (Fig. 2: left column vs. right column). In comparison with Δ_n , Δ_ρ is more sensitive to errors (Fig. 2). Such sensitivity is governed mostly by different impacts of these aerosol characteristics on the aerosol optical properties. Overall, the results of the sensitivity study demonstrate that: (1) our retrieval has the benefit of being

robust to random noise, (2) the uncertainties of the aerosol retrieval are almost linearly proportional to the measurement errors, and (3) these uncertainties are smaller for the retrieved real refractive index compared to the retrieved effective density.

4 Case study

Thus far, we have discussed the application of our retrieval to realistic, but artificial inputs. Here, we consider its application to actual TCAP measurements. The TCAP program was designed to study the evolution of optical, microphysical and chemical properties of both marine aerosol and aerosol transported from continental North America over the Atlantic Ocean and their impact on the radiation energy budget. During this project the ground-based ARM Mobile Facility (AMF) was deployed at the Highlands Center in the Cape Cod National Seashore, which is located on the easternmost portion of Massachusetts, along the east coast of the US.

The TCAP data represent coastal summertime observations and therefore values of ambient relative humidity (RH) for the sample period can be quite high (70–85%). The temporal changes of RH inside the nephelometer are uncontrolled due to instrumental setup, and the average RH value within the nephelometer is about 60%. Below, the term “RH” refers to the RH inside the nephelometer. To account for the RH changes (at least partially) and their impact on optical properties and particle size, the light scattering measurements and ASP data are adjusted to reference “dry” conditions (RH = 40%) using standard corrections (Appendix A). In particular, we apply these corrections to both the measured sub-micron ($D_p < 1 \mu\text{m}$) and sub-10-micron ($D_p < 10 \mu\text{m}$) optical properties.

The AMF site (41.87° N; 70.28° W) was equipped with numerous instruments for sampling aerosol, cloud and radiative properties; a comprehensive overview of the full instrumentation is given by Berg et al. (2014). Here we consider only the measurements that are relevant in applying and evaluating our method (Figs. 4, 5 and 6a). For example, TCAP aerosol chemical composition measurements, obtained from an

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Aerodyne Aerosol Chemical Speciation Monitor (ACSM) and a Single Particle Soot Photometer (SP2), are shown Fig. 4. The large contribution of the organic aerosol to the total mass suggests that aerosol considered here is weakly absorbing at mid-visible wavelengths and larger (Barnard et al., 2008). High values of the single-scattering albedo (>0.95) obtained from the ground-based retrievals (Kassianov et al., 2013), in situ aircraft (e.g., Berg et al., 2014) and near-surface (Titos et al., 2014) measurements confirm this conjecture. Similar high values of the single-scattering albedo for typical marine conditions have been reported previously (Fujitani et al., 2007 and references therein).

The retrieval is then evaluated using independently obtained values of density and refractive index obtained from measured chemical constituents (Fig. 4) and volume averaging, as described in Barnard et al. (2010). When using this method there are several factors that must be considered, which we briefly mention here. First, this method assumes that the size distributions for each chemical constituent are identical. Second, the inferred density and refractive index are necessarily found from measured chemical species, and to the extent that we cannot measure all species – such as crustal materials – some error will occur. This error is expected to be small because a coarse mode is not prominently seen at the TCAP site. Finally, because the chemical measurements do not include water uptake, and because these measurements are for sub-micron aerosol, we can only use the chemical method to find aerosol properties for these dry conditions. The values found in this manner are assigned the symbols $n_{\text{mod, chem}}$ and $\rho_{\text{mod, chem}}$. The values inferred from the measured optical properties are assigned the symbols $n_{\text{mod, opt}}$ and $\rho_{\text{mod, opt}}$; note the added subscripts “chem” and “opt”. To estimate uncertainties of the retrieved aerosol characteristics, we introduce the root mean squared error (RMSE), which is defined as the root mean squared difference (%) between the aerosol characteristics retrieved from the measured optical properties and their counterparts calculated from the chemical composition data. Before considering these uncertainties, we review features of the TCAP measurements relevant for retrieving the aerosol characteristics (n and ρ) and their temporal variability.

4.1 Temporal changes of observed and retrieved aerosol properties

There are strong temporal changes of the observed RH-corrected sub-10-micron scattering coefficient and hemispheric backscatter fraction (Fig. 6a). However, they do not follow, in general, the observed RH variations (Fig. 6b). This confirms that the RH-corrected optical properties depend weakly on the RH variability and that the standard RH corrections (Appendix A) work reasonably well when applied to the challenging observational conditions considered here. Similar to the sub-10-micron aerosol optical properties (Fig. 6a), the ratio (f_{obs}) of sub-micron and sub-10-micron scattering coefficients (Fig. 6b), also depends weakly on the RH changes. This ratio, called the fine-mode fraction of light scattering, defines the relative contribution of the sub-micron particles to the measured sub-10-micron scattering coefficient (which includes the sub-micron scattering) and is commonly used as a proxy of the aerosol source (anthropogenic vs. natural). The frequently observed large values of f_{obs} (Fig. 6b) suggest that the light scattering is often dominated by sub-micron particles. However, there are several events where values of f_{obs} are quite small (less than 0.5) (Fig. 6b). For these events, the relative contribution of large ($1 \mu\text{m} < D_p < 10 \mu\text{m}$) particles to the light scattering can be substantial.

The aerosol characteristics $n_{\text{mod, opt}}$ and $\rho_{\text{mod, opt}}$ retrieved from the measured optical properties $\sigma_{\text{s, obs}}$ and β_{obs} show strong temporal variability (Fig. 6c and d). Occasional prolonged gaps in time series of the retrieved aerosol characteristics (Fig. 6c and d) are periods when no solution can be found. These periods typically represent optically thin conditions when values of the observed scattering coefficient ($\sigma_{\text{s, obs}}$) are relatively small (e.g., $\sigma_{\text{s, obs}} < 10 \text{ M m}^{-1}$) (Fig. 6a), and are characterized by larger values of the measurement uncertainties when compared to other periods. Also, occasional prolonged gaps occur for events where the relative contribution of large particles to the light scattering is substantial (Fig. 6b). Commonly, these large particles may have an irregular (non-spherical) geometry, which is not taken into account by the current version of our retrieval. The observed temporal variations of the retrieved aerosol properties

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the corresponding RMSE is about 28 %. This agreement is even more pronounced for the real refractive index: the mean sub-10-micron ($\bar{n}_{\text{mod, opt}} \sim 1.53$) and sub-micron ($\bar{n}_{\text{mod, opt}} \sim 1.49$) values of the real refractive index overestimate slightly their chemical-based counterpart ($\bar{n}_{\text{mod, chem}} \sim 1.48$). The corresponding values of RMSE are about 6 % and 3 %, respectively (Table 2b). The better agreement for the real refractive index compared to the effective density (Table 2b; Fig. 7) is consistent with the sensitivity study results. A related marine study (Quinn et al., 2004) has demonstrated that both the density and real refractive index of sampled aerosols are size dependent, and their dry values (RH = 55 ± 5 %) are typically higher for large particles (1.1–10 μm range) than for sub-micron particles. The differences observed in the TCAP case study between the retrieved dry sub- and sub-10-micron aerosol characteristics (Table 2b) are in line with these findings.

In contrast with the size factor, the humidity-related factor can modify both the magnitude and sign of the relative differences between mean values of the optical- and chemical-based aerosol characteristics (Table 2b and c), and both the shape and peak position of the corresponding probability distribution functions of the retrieval errors (Fig. 7). The observed RH-sensitivity of the retrieved aerosol characteristics suggests that the RH corrections of the input parameters (Appendix A) only partially remove the corresponding RH-dependence. For humid conditions, most of relative differences tend to have positive values; for dry conditions, the opposite is true (Fig. 7). The observed underestimation of retrieved characteristics under humid conditions is consistent with results from previous studies (e.g., Erlick et al., 2011) and can be explained as follows. The aerosol characteristics (n and ρ) have smaller values for the pure water than for other chemical components, and the fraction of aerosol water increases with an increase of the RH due to water uptake. This fraction depends on the hygroscopicity of aerosol and therefore on its chemical composition (e.g., Zieger et al., 2013 and references therein). Comparison of results (Table 2 and Fig. 6; measured chemical data vs. retrieved optical estimations) suggest that our retrieval can estimate the average values

of the sub-micron aerosol characteristics quite accurately, yet the retrieval uncertainties can be large for the corresponding instantaneous values.

Although the large ($1 \mu\text{m} < D_p < 10 \mu\text{m}$) particles can be of significant importance in modeling aerosol radiative effects (e.g., Eck et al., 2010; Kassianov et al., 2012), in situ information of their physical and optical properties is often lacking. Therefore, there is a strong need to obtain the aerosol characteristics for large particles, and our approach can meet this need by retrieving sub-micron and sub-10-micron aerosol characteristics separately.

4.3 Retrieved size distributions

The main objective of our approach is to retrieve in parallel the two aerosol characteristics n and ρ , which provide closure for two optical properties $\sigma_{s, \text{obs}}$ and β_{obs} . However, this optical closure could be obtained even for size distributions with poor alignment in the overlap region between the SPMS and APS size distribution measurements. To make sure that the obtained size distributions are aligned reasonably well, we consider them for both dry and humid conditions (Fig. 8). For dry conditions, the large particles barely contribute to the combined size distributions, and hence these distributions have weakly defined coarse modes (Fig. 8a and c). For the size distributions that occur under wet conditions, contributions of the small ($D_p < 0.5 \mu\text{m}$), moderate ($0.5 \mu\text{m} < D_p \leq 1.0 \mu\text{m}$) and sub-10-micron particles to the volume distribution are comparable, and hence these volume distributions may have multi-modal structure with well-defined fine, intermediate and coarse modes with peaks near $0.30 \mu\text{m}$, $0.55 \mu\text{m}$ and $2.5 \mu\text{m}$, respectively (Fig. 8d). Since the size of particles increases substantially during water uptake, the observed multi-modal structure of the wet distributions can be associated with measured particle composition/size and the RH-related particle growth. Most importantly, however, visual inspection indicates that there are no major alignment problems in the overlap region and therefore our retrieval shows the capability to align smoothly the measured SMPS and APS spectra. In this regard, the

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effective density (1.68 ± 0.21). Moreover, the retrieved effective density and real refractive index exhibit a tendency to decrease gradually with increase of the relative humidity, and the observed tendency is consistent with the present knowledge of the hygroscopic nature of aerosols.

The majority of existing methods, such as frequently used techniques based on chemical composition measurements, offer essential information on the aerosol characteristics of sub-micron particles. However, such information for the large particles ($1 \mu\text{m} < D_p < 10 \mu\text{m}$) from many current field instruments is mostly lacking. Our retrieval overcomes this limitation using light scattering measurements of both sub-micron and sub-10-micron aerosol optical parameters, which in turns permits discrimination between retrieved aerosol characteristics of small and large particles. This size-resolved information, as a new aspect of future studies, will be vital for better understanding the sources, formation, and fate of aerosol particles, and for evaluating and improving regional and large-scale climate models.

Appendix A: Relative humidity corrections

A1 Coarse mode size distribution (PNNL-APS) adjustment to relative humidity of 40 %

The relative humidity of the sample flow changed due to temperature changes as the sample propagated from outside (ambient conditions) into the Aerodynamic Particle Sizer (APS), which was installed inside a climate controlled trailer. The APS was placed as close as practical to the main inlet stack of the trailer; no control on the sample line relative humidity was exercised. The relative humidity of the APS sample line is not normally measured; so, strictly speaking, size distribution measurements were performed under unknown relative humidity.

We have estimated the relative humidity inside the APS sample line using the temperature inside the APS (reported in parameter “Box Temperature”) and the absolute

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humidity of ambient air calculated from the ambient temperature and RH (measured by the MAOS AOS meteorological station). The APS measured size distributions were adjusted for hygroscopicity following a generally accepted approach (e.g., Swietlicki et al., 2000): the hygroscopic diameter growth factors $G(\text{RH}) = \frac{D_{\text{RH}}}{D_{\text{dry}}}$ are parameterized as

$$G(\text{RH}) = \left(1 - \frac{\text{RH}}{100}\right)^{-\gamma}$$

A value of 0.233 for the parameter γ was adopted from Swietlicki et al. (2000), recommended for clean marine air mass with RH within $0 < \text{RH} < 95\%$ and wind speed below 8 m s^{-1} (for which local sea-salt aerosol production is small); relative humidity of 40% was assumed for “dry” size. In particular, the World Meteorological Organization recommends RH levels of 40% or below (WMO/GAW, 2003). Note that the selected value ($\gamma = 0.233$) was obtained for a particle diameter of 166 nm; the data on hygroscopic growth of super micron aerosol is virtually nonexistent.

A2 Aerosol scattering coefficients (NOAA AOS “dry” Nephelometer) adjustment to relative humidity of 40%

Due to technical difficulties, the relative humidity inside the NOAA AOS “dry” Nephelometer was not kept at 40% or below (as recommended by the WMO). We have adjusted the scattering coefficients to $\text{RH}_{\text{dry}} = 40\%$ using a generally accepted approach (e.g., Gasso et al., 2000; Crahan et al., 2004): the dependence of aerosol scattering on RH was parameterized through aerosol hygroscopic exponent α :

$$\sigma_{\text{sp}}(\text{RH}) = \sigma_{\text{sp}}(\text{RH}_{\text{dry}}) \left(1 - \frac{\text{RH}}{100}\right)^{-\alpha} \quad (\text{A1})$$

The relative humidity inside the nephelometer was estimated using the temperature inside the nephelometer (reported in parameter “Box Temperature”) and the absolute

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humidity of ambient air calculated from the ambient temperature and RH (measured by the MAOS AOS meteorological station). Gasso et al. (2000) reported α values of 0.57 for polluted marine and 0.73 for clean marine aerosol; Crahan et al. (2004) reported α values between 0.23 and 0.48 for marine air masses, with an average for coastal air of 0.45. Here we use value of 0.5 for parameter α .

Appendix B: Examples of “dry” and “wet” sub-periods

This section shows time series of the measured and retrieved aerosol parameters for the dry (Fig. B1) and wet (Fig. B2) conditions. The selected dry and wet sub-periods have similar mean values of the scattering coefficient ($\bar{\sigma}_{s, \text{dry}} \sim 60 \text{ M m}^{-1}$ vs. $\bar{\sigma}_{s, \text{wet}} \sim 50 \text{ M m}^{-1}$), hemispherical backscatter fraction ($\bar{\beta}_{s, \text{dry}} \sim 0.11$ vs. $\bar{\beta}_{s, \text{wet}} \sim 0.08$) and fine-mode fraction of light scattering ($\bar{f}_{\text{dry}} \sim 0.9$ vs. $\bar{f}_{\text{wet}} \sim 0.8$). In contrast, these sub-periods have very different mean values of the RH within the nephelometer ($\text{RH}_{\text{dry}} \sim 45\%$ vs. $\text{RH}_{\text{wet}} \sim 70\%$). These plots (Figs. B1 and B2) illustrate two important points: (1) the retrieved optical-based aerosol characteristics ($n_{\text{mod, opt}}$ and $\rho_{\text{mod, opt}}$) overestimate and underestimate their chemical-based counterparts ($n_{\text{mod, chem}}$ and $\rho_{\text{mod, chem}}$) for the dry (Fig. B1c and d) and wet (Fig. B2c and d) conditions, respectively, and (2) the obtained overestimation is more pronounced for the sub-10-micron optical-based characteristics compared to the sub-micron characteristics.

Acknowledgements. This work has been supported by the Office of Biological and Environmental Research (OBER) of the US Department of Energy (DOE) as part of the Atmospheric Radiation Measurement (ARM) and Atmospheric System Research (ASR) Programs. The Pacific Northwest National Laboratory (PNNL) is operated by Battelle for the DOE under contract DE-A06-76RLO 1830.

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Table 1. Examples of the typical light scattering data, measured particle size distributions, and assumed values of particle density used for the optical closure studies under marine conditions. The listed examples illustrate a variety of instruments for measuring the light scattering properties and size spectra, as well as a wide range of assumed values of particle density (from 1 to 2.2 g cm^{-3}).

Light scattering observations	Size distribution measurements	Assumed density	Environmental conditions	Reference
Integrating nephelometer 3563 ^a Raman lidar Sun photometer Cimel	SMPS 3936 ^b APS 3321 ^c PCASP-100X ^d	2.2 g cm^{-3}	Mixture of marine aerosol with small fraction of absorbing particles	Guerrero-Rascado et al. (2011)
Integrating nephelometer 3563 ^a	ELPI ^e APS 3320	1.0 g cm^{-3}	Marine aerosol	Chamaillard et al. (2006)

^a TSI *Integrating Nephelometer* 3563 measures the total and hemispheric backscatter light-scattering coefficients at three wavelengths (0.45, 0.55 and $0.70 \mu\text{m}$) (<http://www.tsi.com/Integrating-Nephelometer-3563/>).

^b TSI *Scanning Mobility Particle Sizer* (SMPS) 3936 measures aerosols with $0.003\text{--}1.0 \mu\text{m}$ electrical mobility size range (<http://www.tsi.com/scanning-mobility-particle-sizer-spectrometer-3936/>).

^c TSI *Aerodynamic Particle Sizer* (APS) 3321 measures aerosols with $0.5\text{--}20.0 \mu\text{m}$ aerodynamic size range (<http://www.tsi.com/aerodynamic-particle-sizer-spectrometer-3321/>).

^d *Passive Cavity Aerosol Spectrometer* (PCASP)-100X measures aerosols with $0.1\text{--}3.0 \mu\text{m}$ optical equivalent radius range (<http://www.dropletmeasurement.com/products/airborne/PCASP-100X>).

^e *Electrical low-pressure impactor* (ELPI) measures aerosols with $0.007\text{--}10.0 \mu\text{m}$ aerodynamic size range (<http://dekati.com/cms/elpi>).

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Table 2a. Mean and standard deviation (SD) of the retrieved effective density (ρ_{mod}) and real refractive index (n_{mod}) obtained for all retrievals under “wet” and “dry” conditions (237 points). The RH is 58.4 ± 8.1 (mean \pm SD). The chemically derived aerosol properties represent dry conditions only. The corresponding Root Mean Squared Error (RMSE) is defined as the root mean squared difference (%) between the aerosol characteristics retrieved from the measured optical properties and their counterparts calculated from the chemical composition data. These basic statistics (mean, SD and RMSE) represent both the sub- μm and super- μm aerosol characteristics. As discussed in Barnard et al. (2010), uncertainties in the chemically derived density and refractive index are at least 5 %.

Parameter	Model	Mean	SD	RMSE
ρ_{mod}	Opt (sub-10 μm)	1.605	0.271	31.3
	Opt (sub- μm)	1.602	0.252	28.9
	Chem (sub- μm)	1.543	0.055	–
n_{mod}	Opt (sub-10 μm)	1.502	0.034	4.8
	Opt (sub- μm)	1.467	0.026	3.8
	Chem (sub- μm)	1.481	0.017	–

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Table 2b. The same as Table 2a, but for retrievals under dry conditions (104 points; $RH = 53.7 \pm 4.2$).

Parameter	Model	Mean	SD	RMSE
ρ_{mod}	Opt (sub-10 μm)	1.769	0.222	35.0
	Opt (sub- μm)	1.677	0.208	27.8
	Chem (sub- μm)	1.522	0.049	–
n_{mod}	Opt (sub-10 μm)	1.528	0.026	6.2
	Opt (sub- μm)	1.487	0.022	3.3
	Chem (sub- μm)	1.477	0.015	–

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Table 2c. The same as Table 2a, but for retrievals under wet conditions (133 points; $RH = 65.7 \pm 3.9$).

Parameter	Model	Mean	SD	RMSE
ρ_{mod}	Opt (sub-10 μm)	1.477	0.235	28.1
	Opt (sub- μm)	1.544	0.269	29.7
	Chem (sub- μm)	1.560	0.055	–
n_{mod}	Opt (sub-10 μm)	1.481	0.026	3.3
	Opt (sub- μm)	1.451	0.016	4.2
	Chem (sub- μm)	1.485	0.017	–

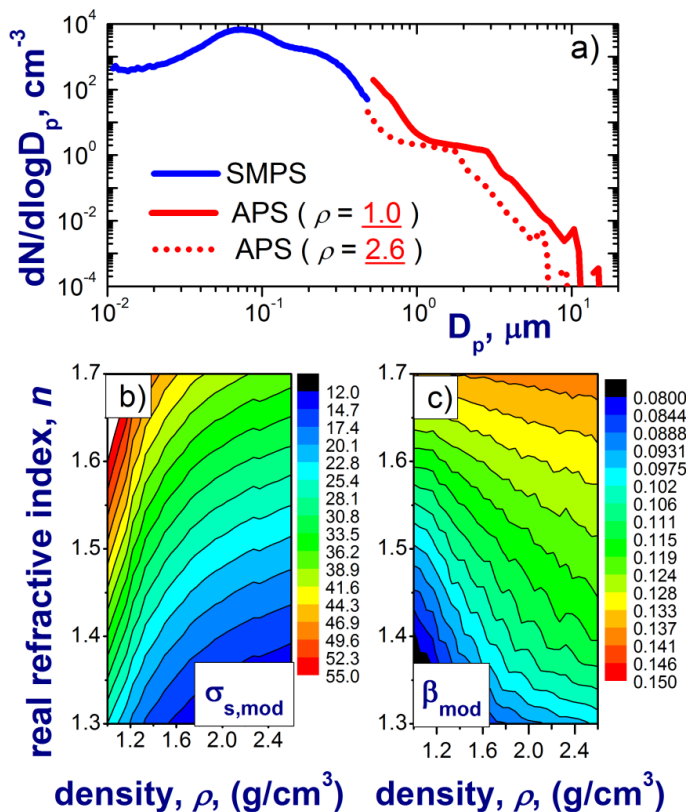


Figure 1. (a) Schematic illustration of merging the SMPS (size range 0.01 to 0.48 μm) and APS (size range 0.52 to 19.8 μm) distributions. (b and c) Example of 2-D diagrams of modeled scattering coefficient $\sigma_{s,\text{mod}}$ (b) and hemispheric backscatter fraction β_{mod} (c) as a function of the real refractive index (n) and effective density (ρ). The combined SMPS-APS distributions for a given ρ are used to obtain modelled optical properties as a function of n .

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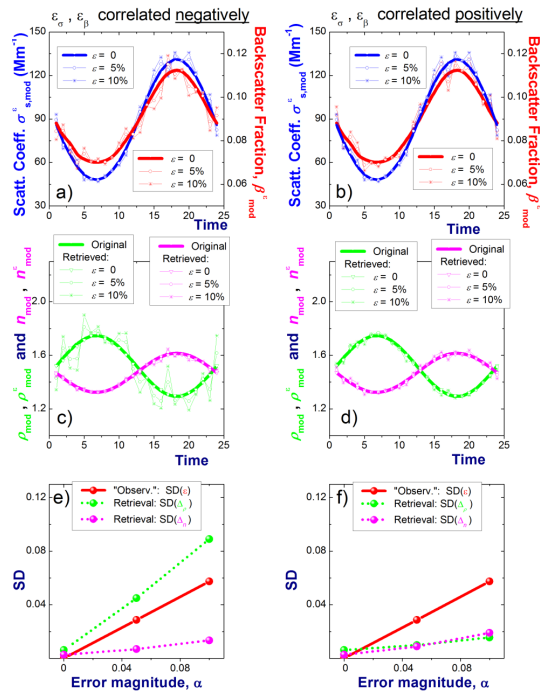


Figure 2. Example of (a and b) modeled aerosol optical properties (scattering coefficient σ_s^E and backscatter fraction $\beta_{s,mod}^E$) with different magnitudes (0%, 5% and 10%) of the added errors (ε_σ and ε_β); (c and d) the corresponding “retrieved” aerosol characteristics (effective density ρ_{mod}^E and real refractive index n_{mod}^E) and their original counterparts (ρ_{mod} and n_{mod}); and (e and f) the standard deviation (SD) of differences between the original and retrieved aerosol characteristics (Δ_ρ and Δ_n) and the SD of the “observational” errors for two extreme cases: the observational errors (ε_σ and ε_β) are correlated negatively (left column) and positively (right column). See text for details.

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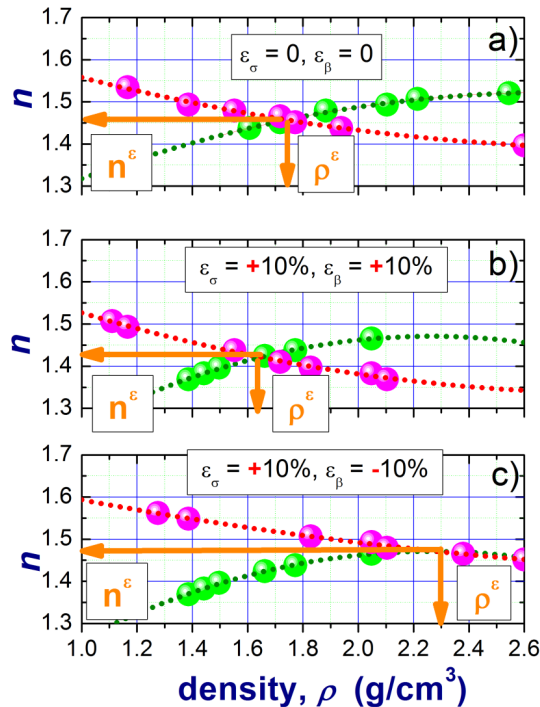


Figure 3. Schematic illustration of retrieving aerosol characteristics for three cases: **(a)** “error-free” conditions when the observational errors (ε_σ and ε_β) are zero, **(b)** and **(c)** “noisy” conditions when the observational errors are positively **(b)** or negatively **(c)** correlated. The family of green points is an isoline of the assumed scattering coefficient at a given time, and in a similar vein, the magenta points represent an isoline of the assumed backscatter fraction. The corresponding polynomial fits (dotted lines with blue and red colors) are used for obtaining the solution ($n_{\text{mod}}^\varepsilon$ and $\rho_{\text{mod}}^\varepsilon$). This illustration represents the “error-free” and “noisy” conditions from Fig. 2c and d for a given instant (Time = 2).

TCAP chemistry data 2012

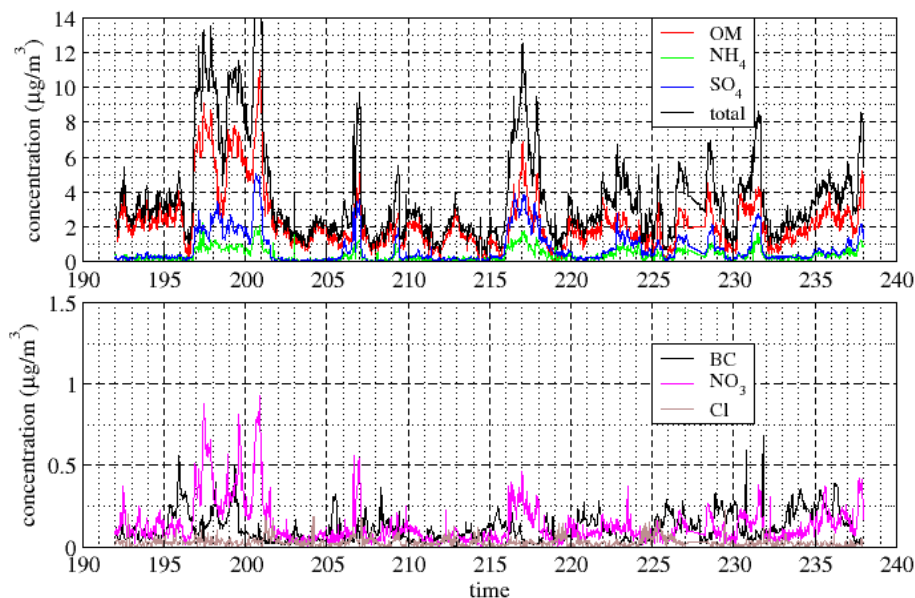


Figure 4. Time series of aerosol chemical components measured during the TCAP field experiment during summertime 2012. Here and on Figs. 5 and 6 time represents Julian day in 2012 (1 January = 1). Days 190 and 214 correspond to 8 July 2012 and 1 August 2012, respectively. Organic matter (OM), ammonium (NH_4), sulfate (SO_4), nitrate (NO_3), chloride (Cl) and total aerosol concentrations are from ASCM measurements. Black carbon (BC) concentrations are from SP2 measurements. The BC, OM, and ionic species are used to find the effective density and refractive index using the “chemical” method. Note the large mass fraction of organic matter (OM), indicating that the aerosol is weakly absorbing at mid-visible wavelengths.

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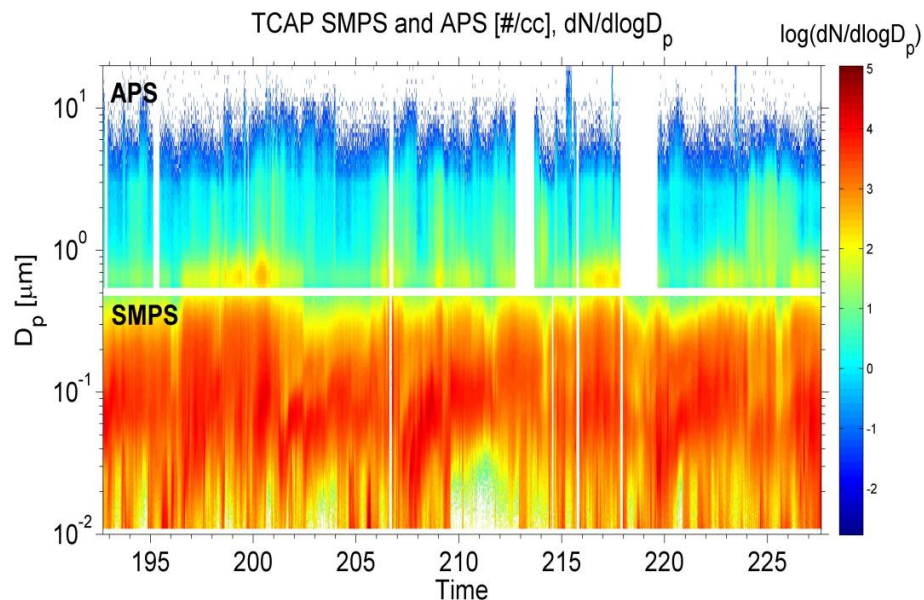


Figure 5. Time series of particle number distributions as a function of geometric diameter (y axis) measured by in situ ground-based instruments: SMPS (bottom panel) and APS (top panel). The measured SMPS and APS size distributions together with optical measurements are used for retrieving the aerosol characteristics (Fig. 6c and d).

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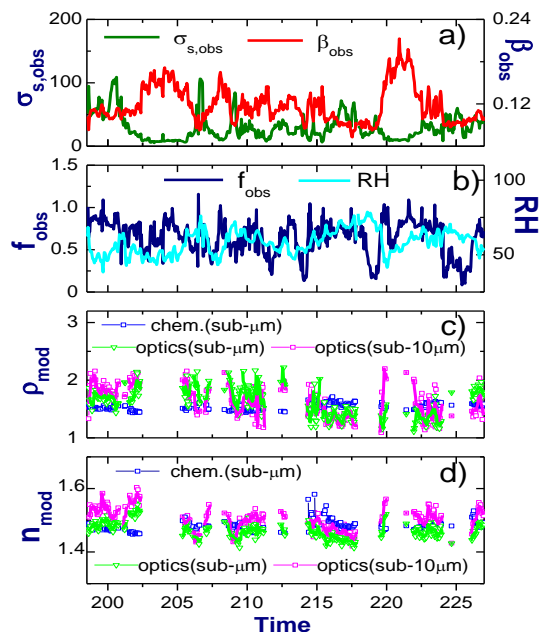


Figure 6. Time series of **(a)** the sub-10-micron ($D_p < 10 \mu\text{m}$) scattering coefficient ($\sigma_{s, \text{obs}}$) and hemispheric backscattering fraction (β_{obs}) measured by integrating nephelometer, **(b)** relative humidity (RH) and the fine-mode fraction of light scattering (f_{obs}), **(c)** the effective density calculated from the sub-micron ($D_p < 1 \mu\text{m}$) dry chemical composition ($\rho_{\text{mod, chem}}$; blue squares) and the retrieved ($\rho_{\text{mod, opt}}$) from measured sub-micron (green triangles) and sub-10-micron (magenta squares) optical properties, **(d)** the real refractive index calculated from the sub-micron dry chemical composition ($n_{\text{mod, chem}}$; blue squares) and retrieved ($n_{\text{mod, opt}}$) from measured sub-micron (green triangles) and sub-10-micron (magenta squares) optical properties **(c and d)**. Solid lines (blue, green and magenta) represent the smoothed versions (boxcar averaging; 5 points) of the corresponding calculated and retrieved aerosol characteristics (blue, green and magenta squares). Appendix B includes “zoom-in” versions of Fig. 6.

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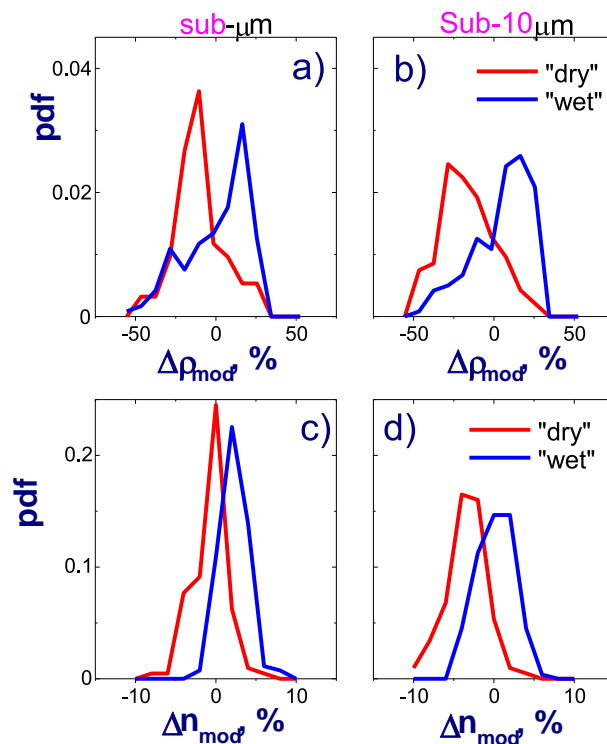


Figure 7. Probability distribution function (pdf) of the retrieval errors for the effective density ($\Delta\rho_{\text{mod}}$) and the real refractive index (Δn_{mod}) obtained for the sub-micron (left column) and sub-10-micron (right column) aerosol characteristics. Red and blue colors define the retrieval errors obtained for the “dry” and “wet” conditions, respectively. Note that different scales are used for $\Delta\rho_{\text{mod}}$ (top row) and Δn_{mod} (bottom row). Factors associated with particle size and humidity (see text for details) are the main sources of the observed differences between the pdfs (red vs. blue colors).

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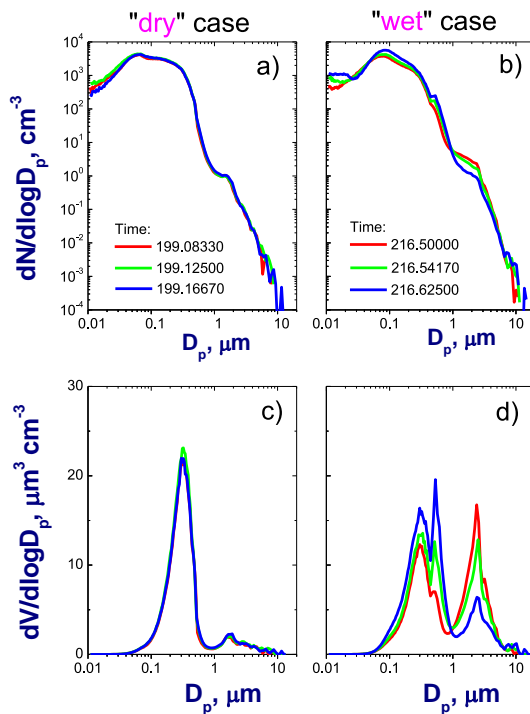


Figure 8. Aerosol number (top row) and volume (bottom row) distributions obtained from the retrieved effective density (ρ_{mod}) and the measured SMPS and APS spectra. Red, green and blue colors are used for different time steps within specified sub-periods (Appendix B). These sub-periods represent cases with similar optical properties, but very different values of the instrument RH. For the first sub-period (“dry” case; left column), the RH is relatively small (RH \sim 45 %). For the second sub-period (“wet” case; right column), the RH is relatively high (RH \sim 70 %).

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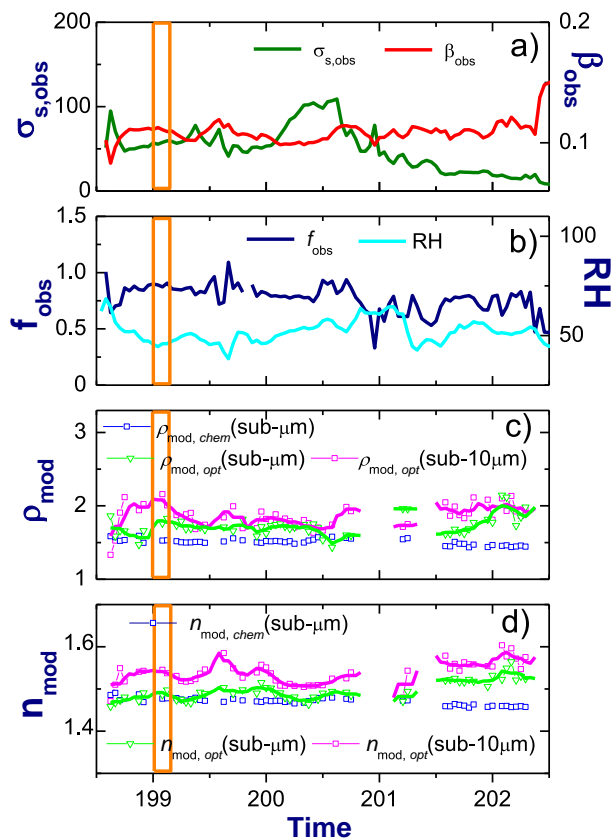


Figure B1. The same as Fig. 6, but for selected dry period (52 points, $\text{RH} = 51.9 \pm 5.8$). The vertical orange bar represents a sub-period where the dry optical properties (Fig. B1) are comparable with their wet counterparts (Fig. B2), but values of the RH are very different.

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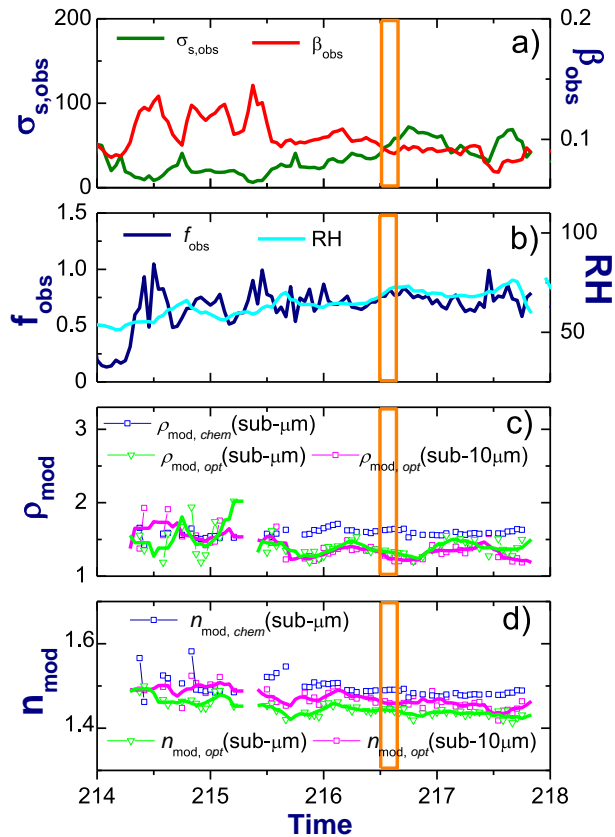


Figure B2. The same as Fig. 6, but for selected wet period (56 points, $RH = 64.2 \pm 6.3$). The vertical orange bar represents a sub-period where the wet optical properties (Fig. B2) are comparable with their dry counterparts (Fig. B1), but values of the RH are very different.

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