Lidar-Radiometer Inversion Code (LIRIC) for the retrieval of vertical aerosol properties from combined lidar/radiometer data: development and distribution in EARLINET

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Abstract.

This paper presents a detailed description of LIRIC (LIdar-Radiometer Inversion Code) algorithm for simultaneous processing of coincident lidar and radiometric (sun photometric) observations for the retrieval of the aerosol concentration vertical profiles. As the lidar/radiometric input data we use measurements from European Aerosol Research Lidar Network (EARLINET) lidars and collocated sun-photometers of Aerosol Robotic Network (AERONET). ~~The LIRIC data processing provides sequential inversion of the combined lidar and radiometric data by the estimations of column-integrated aerosol parameters from radiometric measurements followed by the retrieval of height-dependent concentrations of fine and coarse aerosols from lidar signals using integrated column characteristics of aerosol layer as *a priori* constraints.~~ The LIRIC data processing provides sequential inversion of the combined lidar and radiometric data. The algorithm starts with the estimations of column-integrated aerosol parameters from radiometric measurements followed by the retrieval of height dependent concentrations of fine and coarse aerosols from lidar signals using integrated column characteristics of aerosol layer as a priori constraints. The use of polarized lidar observations allows us to discriminate between spherical and non-spherical particles of the coarse aerosol mode.

The LIRIC software package was implemented and tested at a number of EARLINET stations. Intercomparison of the LIRIC-based aerosol retrievals was performed for the observations by seven EARLNET lidars in Leipzig, Germany on 25 May, 2009. We found close agreement between the aerosol parameters derived from different lidars that supports high robustness of the LIRIC algorithm. The sensitivity of the retrieval results to the possible reduction of the available observation data is also discussed.

1 Introduction

1   The aerosol impact on the radiation balance of the atmosphere is an important climate ~~forming~~

2   forcing factor. In addition, aerosol particles are among the unhealthiest pollutants of the air. It is

3   dramatized by rapid propagation of pollutants in the atmosphere that expands local ecocatastrophes

4   to a global scale. Therefore, the monitoring of the aerosol evolution and transport in the atmosphere

5   is an obligatory prerequisite for predicting climatic and ecological changes.

6

7   Sun-radiometer and lidar networks contribute to aerosol remote sensing. The global Aerosol

8   Robotic Network (AERONET) of ground-based sun/sky-scanning radiometers (e.g. Holben et al.,

9   1998) provides reliable data on columnar aerosol properties from more than 200 world-wide

10   distributed sites. The results of AERONET observations are the aerosol optical thickness (AOT)

11   obtained from direct sun observations and additional microphysical and optical properties of aerosol

12   particles (single scattering albedo, volume distribution of aerosol particles, complex refractive

13   index, fraction of spherical particles, etc.) derived by the inversion of direct and scattered radiation

14   measurements (Dubovik and King, 2000; Dubovik et all, 2002, 2004). The regional radiometer

15   network SKYNET was established in the South Eastern Asian regions (Takamura et al., 2004) and

16   it employs its own equipment and processing procedure (Hashimoto et al., 2012).

17

18   The lidar measurements are ~~basically~~ used to provide information on the vertical variability of the

19   aerosol characteristics. Currently, lidar networks, such as the European Aerosol Research Lidar

20   Network (EARLINET) (Bösenberg et al., 2000, Pappalardo et al., 2014), the micro-pulse lidars

21   network (MPL-Net) (Welton et al., 2002), the Asian Dust Network (AD-Net) (Murayama et al.,

22   2001), the lidar network in Former SU countries CIS-LiNet (Chaikovsky et al., 2005), the northeast

23   American CREST Lidar Network (CLN) (Hoff et all, 2009), and the Latin-American Lidar Network

24   LALINET (Antuña et al., 2012), monitor aerosol vertical distributions in the atmosphere over the

25   vast regions of the Earth. The Global Atmosphere Watch (GAW) Aerosol Lidar Observation

26   Network (GALION), also known as the "network of networks" (e.g. Bösenberg and Hoff, 2007),

27   was established under the aegis of GAW to coordinate lidar activity all over the world. The

28   outcome of the lidar observations are presented in the lidar network databases as vertical profiles of

29   aerosol backscatter and extinction coefficients.

30

Aerosol columnar properties from AERONET and ~~pointed above~~ aforementioned vertical profiles of aerosol parameters from lidar networks are complementary pieces of information characterizing aerosol properties. Nowadays, lidars and sun/sky-scanning radiometers are among the basic tools in comprehensive experiments aimed at studying the transformation and transport of smoke (e.g. Lund Myhre et al., 2007; McKendry et al., 2011; Colarco et al., 2004), dust (e.g. Ansmann et al., 2009; McKendry et al., 2007; Müller et al., 2003; Papayannis et al., 2008) and volcanic ash (e.g. Ansmann et al., 2010, 2011, 2012; ; Papayannis et al., 2012; Gasteiger et al., 2010). A number of SKYNET sites (Takamura et al., 2004) and most of the EARLINET stations are equipped with lidar and radiometer instruments. Further enhancement of the aerosol characterization is expected from the synergy of co-located radiometer and lidar observations. Namely, the coordination of measurement procedures of the two systems and the derivation of aerosol parameters from combined measurements results in advanced characterization of the aerosol layer with a superior performance compared to the aerosol information that would have been obtained from independent processing of lidar and radiometer data.

The idea of combined lidar and radiometer sounding (LRS) for retrieving vertical distributions of aerosol characteristics was first proposed by Chaikovsky et al. (2002), and it gave rise to the development of the lidar-radiometer synergetic algorithms (e.g. Chaikovsky et al., 2004a, 2004b). Later, in 2012 under the ACTRIS Research Infrastructure project within the European Union Seventh Framework Programme, the algorithm and software package, named LIRIC (LIdar-Radiometer Inversion Code), was developed for processing data of EARLINET measurements. LIRIC is based on processing co-located lidar and radiometer measurements by using a two-step *sequential inversion*. First, the radiometer data was processed ~~in tune to~~ according to the standard AERONET inversion algorithm. Then, first-step results ~~were~~ are used as *a priori* constrains on aerosol properties for lidar data processing.

First application of LRS technique to the actual data processing was presented by Chaikovsky et al. (2004a). In that study, the technique was adapted to the EARLINET-AERONET stations in Minsk (Belarus) and Belsk (Poland) (e.g. Chaikovsky et al., 2004c, 2010a; Pietruczuk and Chaikovsky, 2007). Results of the LRS observations were of interest for studying of long range aerosol transport

in the East European region (Kabashnikov et al., 2010; Chaikovsky et al., 2010b; Papayannis et al., 2014).

Another algorithm for data processing in combined lidar-and-radiometer experiments exploits the decomposition of the AERONET column-integrated aerosol size distribution into log-normal modes and selection of some of these modes for characterization of aerosol layers using measured lidar data (Cuesta et al., 2008).

The LRS technique for retrieving the aerosol concentration profiles from single wavelength lidar measurements at the MPLNET (Micro-pulse Lidar Network) stations collocated with the Sun/sky radiometer sites of AERONET was developed by Ganguly et al. 2009a. Then this method was applied to processing of the combined AERONET and space CALIOP lidar data (Ganguly et al. 2009b).

Besides, the single-wavelength POLIPHON technique was alternatively developed (e.g. Tesche et al., 2009; Ansmann et al., 2012). This technique retrieves particle volume concentration profiles of aerosol separately for fine and coarse fractions. The algorithm relies on the measured profiles of the particle linear depolarization ratio and lidar ratio and it does not require the assumption of a specific particle shape. Columnar concentrations of aerosol modes retrieved by AERONET are used in POLIPHON as additional input data. The algorithm POLIPHON is designed for the data processing in lidar sounding of the aerosol layers with coarse non-spherical particles (dust, volcano ash).

In the recent years, the LRS technique has been implemented within the advanced research network ACTRIS in the frame of EU 7th Framework Programme project. To date, a number of joint EARLINET/AERONET stations have implemented regular atmospheric observations using LIRIC for processing combined sun-radiometer and lidar-measured data (e.g. Chaikovsky et al., 2012; Papayannis et al., 2014; Tsekeri et al., 2013). The aerosol model and mathematical basis of the LIRIC algorithm became the prerequisite for further development of algorithms for parallel processing simultaneous inversion of combined lidar-radiometer measurements, e.g. GARRLIC (Generalized Aerosol Retrieval from Radiometer and Lidar Combined data) (Lopatin et al., 2013) and the results of ground/satellite bond closure experiments (Dubovik et al., 2014). Note that

1 LIRIC technique should not be regarded only as a basis for new algorithms (e. g. POLIPHON or
2 GARRLIC). LIRIC might be superior to them for many aerosol scenarios: it allows one, for
3 example, to distinguish between fine and coarse spherical fractions (unlike POLIPHON) or
4 distinguish between spherical and non-spherical coarse particles (unlike GARRLIC). At the same
5 time, a comprehensive description of the LIRIC algorithm has not been yet documented in detail.
6
7 This paper describes the basic physical and mathematical aspects of LIRIC algorithm with all
8 necessary equations, thus filling up this gap. The appendices contain the details of the inversion
9 scheme and can be useful for advanced users to modify and improve this code.
10

11 **2 The algorithm concept and structure**
12 The aerosol retrievals from combined lidar and radiometer measurements belong to a class of "ill-
13 posed" inverse problems that, in particular, is characterized by non-unique and highly unstable
14 solutions arising even under small measurement or simulation errors. In practice, solution of the
15 "ill-posed" problems requires to introduce *a priori* information (e.g. Turchin et al., 1971; Tikhonov
16 et al., 1977; Twomey, 1977; Tarantola, 1987; Rodgers, 2000). LIRIC algorithm was designed on the
17 basis of multi-term LSM (Least Square Method) (Dubovik, 2004). This method was implemented in
18 AERONET data processing (Dubovik and King, 2000) and then it was refined in the retrieval
19 algorithms for the data processing of the combined optical measurements (e.g. Dubovik, 2004;
20 Dubovik et al., 2011, 2014; Lopatin et al., 2013).
21

22 The inversion algorithm LIRIC can be divided into three key procedures (e.g. Tarantola, 1987): (i)
23 *parameterization* of the object under study (i.e., development of the aerosol layer model); (ii)
24 *forward modeling*, i.e. derivation of the equations that relate observed signals with specified
25 parameters of the aerosol model; and (iii) *inverse modeling* or retrieval of the target parameters of
26 the aerosol model that minimize discrepancies between the measured and the calculated input
27 signals.
28

29 **2.1 Combined lidar/radiometer experiment and aerosol model**
30 The lidar/radiometer input data assumed to come from measurements of EARLINET lidars (e.g.
31 Matthias et al., 2004; Freudenthaler et al., 2010) and spectral-scanning sun-radiometers of

AERONET (Holben et al., 1998). The majority of EARLINET lidar stations provides day-time measurements of elastic backscatter signals at three wavelengths (355, 532 and 1064 nm) and two cross/parallel polarization components of the signal at a single wavelength. Additional information on aerosol parameters is expected to come from day-time Raman lidar measurements. (Current version of LIRIC algorithm is not designed for using Raman lidar data.)

Radiometric data includes results of direct-Sun and almucantar (scanning) measurements (Holben et al., 1998; Dubovik and King, 2000). Direct-Sun measurements are carried out in 15 min intervals. Almost clear-sky measurements are required to obtain almucantar data and about 2-6 successful measurements are made during the daytime under favorable meteorological conditions at EARLINET/AERONET stations. Under these circumstances time synchronization of lidar and radiometric observations usually means nearly simultaneous measurements within the same 1-hour interval.

These radiometric measurements enable the retrieval of the aerosol properties over the entire atmospheric column. Thus, except for volcanic events, the maximum lidar sounding height, h_{max}, can be limited to the tropopause level because the stratospheric aerosol layer does not significantly contribute to columnar aerosol optical parameters. In contrast, aerosols in the lower troposphere are key contributors to the observed columnar characteristics. Consequently, it is desirable to perform the lidar sounding from the lowest possible altitude. Likewise, the contribution of the bottom layer (which is not observed by lidar) to the columnar optical parameters must be small enough to be modeled by a homogeneous layer with the same aerosol parameters as at the lowest level of lidar sounding. In practice, the lower sounding limit for most of the lidar measurements in EARLINET is about 200 m or more that can be too high especially for low boundary layers in winter seasons. It should be decreased in winter to compensate reduction of the boundary layer height. Therefore, lidar measurements in the lower layer have to be carried out by a ~~co-operative~~ second, complementary receiving system with smaller objective and larger field of view or by sounding the atmosphere along a slant trajectory.

The choice of the optical aerosol model is a key step of the retrieval algorithm. The optical model should be constructed following the principle of parsimony or "Occam's razor": the number of

1 aerosol parameters has to be minimal but complete in order to provide unbiased retrieval from

2 available measurements.

3

4 In this work, we use the AERONET model approach to characterize the aerosol layer of the

5 atmosphere (Dubovik and King, 2000): aerosols are modeled by several modes with a certain

6 aerosol particle size distribution, wherein each mode is a mixture of homogeneous spherical

7 particles and randomly-oriented spheroids (Dubovik et al., 2002, 2006). The distribution of the

8 spheroid aspect ratio is fixed. The number of aerosol modes, K, depends on specification of the

9 lidar data. If we use only total (scalar) backscatter lidar measurements, the aerosol model includes

10 fine and coarse modes ($K = 2$). There is boundary size between fine and coarse fractions in the

11 algorithm, which is determined as the value in 0.194–0.576 µm range that corresponds to a

12 minimum of the column particle volume size distribution, $dV(r)/d\ln r$. If measurements of cross

13 and parallel co-polarized components are available, spherical and non-spherical particles of the

14 coarse mode are considered as two different fractions ($K = 3$).

15

16 Thus, two sets of parameters characterize the aerosol layer:

17   i. A number of columnar aerosol parameters retrieved from radiometer measurements

18   (Dubovik et al., 2000a, 2002, 2006). This set of parameters is formed by: (1) the total

19   content of each aerosol mode, (i.e. columnar volume concentrations),
20

21
$$\hat{C}_k^V = \int_{r_{\min,k}}^{r_{\max,k}} \frac{dV_k(r)}{d\ln r} d\ln r, \tag{1}$$

22

23   where $r_{\min.k}$ and $r_{\max,k}$ is the minimum and the maximum radius of the kth aerosol

24   mode ($k = 1,..,K$), respectively; (2) the particle volume size distribution $dV_k(r)/d\ln r$

25   for each aerosol mode, (3) complex refractive indices at the wavelength λ,

26          $m(\lambda) = n(\lambda) + i\kappa(\lambda)$, (4) the "sphericity", ζ_{sph} (the ratio of spherical particle's volume to

27   the total volume), (5) aerosol optical thickness (AOT) of the kth aerosol mode, $\hat{E}_k(\lambda_j)$,

28   (6) the single scattering albedo for the kth aerosol mode, $\varpi_k(\lambda)$, (7) the elements of the

backscattering matrix, $P_{x,x}^{k}(\lambda,180^{0})$, and (8) coefficients $\mathbf{a}_k$ and $\mathbf{b}_k$, which determine optical extinction and backscatter characteristics of aerosol particles for the k-aerosol mode (see Sect. 3.1). ==Parameters (1)-(4) are the independent "state" variables whilst parameters (5)-(8) are derived from the state variables.== Parameters $\mathsf{m}(\lambda)$ and ζ_{sph} are assumed the same for particles of all sizes. Definitions and detailed description of the columnar aerosol parameters are available at the AERONET information system; cloud screening and quality control algorithms were described by Holben et al. (2006).

ii. The height, h, distributions of particle volume concentrations (PVC) for each of aerosol mode, $\mathsf{c}_k(h)$, which define the vertical variability of the aerosol features.

A lack of lidar data to resolve height-variation of aerosol microstructure motivates the assumption of altitude-independent microphysical parameters of the aerosol modes.

2.2 Algorithm's structure

Two options of the retrieval procedure for the processing LRS data have been developed:

1. *First one deals with sequential inversion* of lidar and radiometer data. It is carried out by preliminary calculation of the ~~i type~~ ==column== parameters defined in sec. 2.1 from radiometric measurements by using the AERONET inversion algorithm (Dubovik and King, 2000), followed by subsequent inversion of the ~~ii type~~ ==height distribution== parameters by using lidar data with columnar characteristics of aerosol layer passed as *a priori* data (Chaikovsky et al., 2012);

2. *Second option suggests* ~~parallel~~ ==simultaneous *inversion*== approach for retrieving optimal ~~i- and ii-type~~ parameters of the aerosol model by using a joint inversion procedure from combined lidar and radiometer data.

While the sequential algorithm could be considered as an unsophisticated inversion procedure to combine lidar and AERONET data, the parallel inversion method leads, in principle, to more effective estimation of aerosol parameters because it allows simultaneously retrieved columnar aerosol parameters to be specified in accordance with the additional lidar data. Currently, the ~~parallel~~ ==simultaneous inversion== algorithm for a two-component aerosol model is implemented in

9

GARRLIC (Lopatin et al., 2013). Similar aerosol mode concentration profiles and residual discrepancies between measured and calculated input signals are obtained from both retrieval procedures when processing experimental data (Lopatin et al., 2013).

Advantages of the "~~parallel~~ simultaneous inversion approach" are expected for more involved measurements, such as in the unified algorithm GRASP (Generalized Retrieval of Aerosol and Surface Properties), which aimed at characterizing atmospheric properties from remote ground and satellite observations (Dubovik et al., 2014).

LIRIC algorithm described below was created on the base of the sequential inversion approach. Figure 1 shows the structure of the algorithm.

The algorithm is divided into several rather independent modules to provide flexibility of the software package. Module 1 (*preprocessing of lidar data*) creates a set of smoothed and normalized lidar signals, $\mathbf{L}^*$, covariance matrix, $\mathbf{\Omega}_L$, and setting parameters (type of lidar measurement, sounding wavelength, geographical coordinates of lidar station and date of measurement, etc.) for modeling aerosol and molecular layers. Module 2 (*recalculation of radiometer data*) estimates ~~i)-type~~ columnar parameters of the aerosol model for lidar sounding wavelengths. Level 1.5 or Level 2.0 AERONET data are acceptable as input data in LIRIC (These data are inputs to Module 2). Initial profiles of the aerosol-mode concentrations, $c_k^0(h)$, as well as molecular (Rayleigh) extinction, $\sigma_r(\lambda, h)$, and molecular backscatter coefficients, $\beta_r(\lambda, h)$, are generated by Module 3 (*atmospheric model*). Module 4 (*forward model*) calculates arrays of lidar signals, $\mathsf{L}_j\left(c_k^{m-1}(h)\right)$, and columnar volume concentrations, $\hat{C}_k^{V,m-1}$, given aerosol concentration profiles, $c_k^{m-1}(h)$, in iterative inversion procedure, where "m" stands for the m-th retrieval iteration and "j" is the number of receiving channel.
Inversion parameters, constraints on the smoothness characteristics, and error signals for the sensitivity test are passed to the algorithm by Module 5 (*inversion settings&error modeling*). The sensitivity test (see Sect. 6) was designed to estimate the responds of the retrieval results to measurement errors and/or uncertainties of input data.

Module 6 (*numerical inversion*) is responsible for fitting aerosol-mode concentration profiles for the retrieved aerosol model, $c_k^{m-1}(h)$, given measured data and *a priori* information.

3 Forward modeling of LRS experiment

Range-corrected normalized lidar signals and columnar-aerosol parameters retrieved from radiometer measurements are the input data to the LRS processing procedure (see Fig. 1). Below we define a set of basic equations that are needed for the forward modeling of the measured quantities as well as to estimate the error-covariance matrix.

3.1 Basic lidar equations

The multichannel lidar carries out J "different" lidar measurements ($j \in 1,..J$) that yields a set of lidar signal records, P_j^*, $j \in 1,..J$. The term "different" means that different kinds of lidar measurements are performed, such of total intensity, cross- and parallel-polarized signal components at different wavelengths. Here we consider that each "different" lidar measurement is provided by a specific j-th channel. Parameter J stands for the number of lidar channels irrespective of the actual implementation of the lidar system.

Range-corrected normalized lidar signals are calculated at the preprocessing stage of the inversion procedure (Module 1 in Fig.1):

$$L_j^*(h) = \frac{S_j^*(\lambda_j, h)}{\hat{S}_j^*(\lambda_j, h_{ref})} \exp(-2\tau_r(\lambda_j, h, h_{ref})), \tag{2}$$

where $S_j^*(\lambda_j, h) = P_j^*(\lambda_j, h)h^2$; $\hat{S}_j^*(\lambda_j, h_{ref})$ is the value of $S_j^*(\lambda_j, h)$ at the reference point, h_{ref} is usually defined in the end of the sensing range; $\tau_r(\lambda_j, h, h_{ref})$ is the molecular optical thickness related to the range of (h, h_{ref}), λ_j is the wavelength, and h is the height. The set of lidar signals, $L_j^*(h)$, constitutes the input lidar vector, $\mathbf{L}^*$.

The lidar system provides measurements from the lowest to the highest altitude levels specified by h_{min} and h_{max}, respectively. Currently, it is assumed that the radiometer is co-located at a height

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1   of $h_0 < h_{min}$, so columnar aerosol optical properties of the layer $h_0 < h < h_{min}$ are to be taken into

2   consideration. If there is no information on the aerosol parameters in the surface layer, this layer is

3   assumed to be homogeneous. Under this assumption, scattering parameters for the altitude range

4    $h_0 < h < h_{min}$ of the lidar vector $\mathbf{L}^*$ are set equal to the values at h_{min}.

5

6   The relationship between the measured lidar signals $\mathbf{L}^*(\lambda)$, and the aerosol mode concentration,

7    $c_k(h)$, can be written as follows:

8    $$\mathbf{L}^* = \mathbf{L}\left(\lambda, c_k(h), \mathbf{a}_k, \mathbf{b}_k\right) + \mathbf{\Delta}_L,$$ (3)

9   where $\mathbf{\Delta}_L$ is the vector of measurement uncertainties. Here, an asterisk (*) denotes "measured" and

10   no-asterisk denotes "model estimated".

11

12   Since function $L(...)$ in Eq. (3) depends on the type of lidar measurement, it is expedient to

13   introduce special parameter, $p_j \in 1,2,...,U$ that indicates the type of measurement associated to the

14   j-channel of the lidar and U is a number of the types. In our case, $p_j \in 1,2,3$, indicates total

15   intensity, cross-polarized, and parallel-polarized measurements, correspondingly.

16

17   The lidar functions, $L_{j,p_j}(...)$, for the p_j-type measurements are defined by the following

18   equations:

19   -   Equation for the total backscatter signal:

20   $$L_{j,1}\left(\lambda_j, h\right) = \frac{\beta_{a,1}(\lambda_j, h) + \beta_r(\lambda_j, h)}{R_{j,1}(\lambda_j, h_{ref})\beta_r(\lambda_j, h_{ref})} \exp\left(-2\int_{h_{ref}}^{h} \sigma_a(\lambda_j, h)dh\right),$$ (4)

21   where

22   $$R_{j,1}(\lambda_j, h) = \frac{\beta_{a,1}(\lambda_j, h) + \beta_r(\lambda_j, h)}{\beta_r(\lambda_j, h)};$$ (5)

23   -   Equation for the parallel-polarized signal component:

1 $$L_{j,3}(\lambda_j, h) = \frac{\beta_{a,3}(\lambda_j, h) + \frac{1}{1+\chi}\beta_r(\lambda_j, h)}{\frac{1}{1+\chi}\beta_r(\lambda_j, h_{ref})R_{j,3}(\lambda_j, h_{ref})} \exp\left(-2\int_{h_{ref}}^{h} \sigma_a(\lambda_j, h)dh\right)$$ (6)

2 where

3

4 $$R_{j,3}(\lambda_j, h) = \frac{\beta_{a,3}(\lambda_j, h) + \frac{1}{1+\chi}\beta_r(\lambda_j, h)}{\frac{1}{1+\chi}\beta_r(\lambda_j, h)} \ ;$$ (7)

5

6   -   Equation for the cross-polarized signal component:

7

8 $$L_{j,2}(\lambda_j, h) = \frac{\left(\beta_{a,2}(\lambda_j, h) + \mu\beta_{a,3}(\lambda_j, h_i) + \frac{\chi+\mu}{\chi+1}\beta_r(\lambda_j, h)\right)}{\frac{\chi+\mu}{\chi+1}\beta_r(\lambda_j, h)R^{eff}(\lambda_j, h_{ref})} \exp\left(-2\int_{h_{ref}}^{h} \sigma_a(\lambda_j, h)dh\right)$$ (8)

9 where

10

11 $$R^{eff}(\lambda_j, h) = \left(\frac{\chi}{(\chi+\mu)}\frac{\beta_{a,2}(\lambda_j, h) + \frac{\chi}{\chi+1}\beta_{,r}(\lambda_j, h)}{\frac{\chi}{\chi+1}\beta_r(\lambda_j, h)} + \frac{\mu}{(\chi+\mu)}\frac{\beta_{a,3}(\lambda_j, h) + \frac{1}{1+\chi}\beta_r(\lambda_j, h)}{\frac{1}{1+\chi}\beta_r(\lambda_j, h)}\right).$$ (9)

12

13 In Eqs. (4)-(9), $\beta_{a,1}$, $\beta_{a,3}$, and $\beta_{a,2}$ denote the aerosol backscatter coefficient and its parallel- and

14 cross- polarized components, respectively; $\sigma_a(\lambda_j, h)$ is the aerosol extinction coefficient;

15 $\chi(\lambda_j) = \frac{\beta_{r,2}(\lambda_j)}{\beta_{r,3}(\lambda_j)}$ is the ratio of cross- and parallel-polarized components of the molecular

16 backscatter coefficient.

17

18 Different cross-talk factors contribute to the spurious signal in the cross polarized receiving

19 channel. These factors include the residual of cross polarized component of the laser beam, non-

20 ideal adjustment of the polarization planes between transmitter/receiver channels and depolarization

by optical elements. Equations (6) and (8) allow for these cross-talk effects in a similar manner to Chaikovsky (1990) and Biele et al. (2000). Thus, parameter μ in Eqs. (8)-(9) represents the leakage of the parallel component of the sounding beam into the cross polarized lidar receiving channel. Parameter μ is instrument characteristic that is assumed to be known quantity, i.e. it is not updated by retrieval procedure.

The aerosol extinction and backscatter coefficients in the Eqs. (3)-(9) are expressed as a function of the parameters of the aerosol modes:

$$\sigma_a(\lambda_j, h) = \sum_k c_k(h) a_k(\lambda_j), \tag{10}$$

$$\beta_{a,1}(\lambda_j, h) = \sum_k c_k(h) b_{k,1}(\lambda_j), \tag{11}$$

$$\beta_{a,2}(\lambda_j, h) = \sum_k c_k(h) b_{k,2}(\lambda_j), \tag{12}$$

and

$$\beta_{a,3}(\lambda_j, h) = \sum_k c_k(h) b_{k,3}(\lambda_j) \tag{13}$$

The coefficients $a_k(\lambda_j)$ and $b_{k,x}(\lambda_j)$, pointed out in Section 2.1, are determined by columnar optical parameters of aerosol modes:

$$a_k(\lambda_j) = \frac{\hat{E}_k(\lambda_j)}{\hat{C}_k^V}, \tag{14}$$

$$b_{k,1}(\lambda_j) = \frac{1}{4\pi} \varpi_k(\lambda_j) a_k(\lambda_j) P_{1,1}^k(\lambda_j, 180^0), \tag{15}$$

$$b_{k,3}(\lambda_j) = \frac{1}{4\pi} \varpi_k(\lambda_j) a_k(\lambda_j) \frac{P_{1,1}^k(\lambda_j, 180^0) + P_{2,2}^k(\lambda_j, 180^0)}{2}, \tag{16}$$

$$b_{k,2}(\lambda_j) = \frac{1}{4\pi} \varpi_k(\lambda_j) a_k(\lambda_j) \frac{P_{1,1}^k(\lambda_j, 180^0) - P_{2,2}^k(\lambda_j, 180^0)}{2}, \tag{17}$$

where $\hat{E}_k$ is aerosol optical thickness for the kth aerosol mode; $\varpi_k(\lambda)$ is the single scattering albedo for the kth aerosol mode, and $P^k_{x,x}(\lambda, 180^0)$ are the elements of the backscattering matrix.

3.2 Forward model of radiometer data

In accordance with the multi-term LSM approach (Dubovik, 2004), the columnar concentrations of aerosol modes, $\hat{C}^V_k$, obtained from radiometer measurements are formally considered in LIRIC as a result of additional independent measurements.

The equation for the vector, $\hat{\mathbf{C}}^{*V}$, which is defined as the "measured" columnar volume concentrations of the aerosol modes given vector of aerosol modes concentration, $c(h_i)$, $i \in 1,...I$, can be written in the following form

$$\hat{\mathbf{C}}^{*V} = \mathbf{H}\mathbf{c} + \mathbf{\Delta}_V \ , \tag{18}$$

where $\mathbf{H}$ is convolution matrix for summing the height-resolved concentration over the column; $\mathbf{\Delta}_V$ is the vector of $\hat{\mathbf{C}}^{*V}$ uncertainties.

The k-th component of the vector $\hat{\mathbf{C}}^{*V}$ is defined by the equation

$$C^{*V}_k \left(c_k(h_i) \right) = \sum_{i=1}^{I} c_k(h_i)\Delta h_i + \Delta_{V,k} \ . \tag{19}$$

The structure of the vectors $\hat{\mathbf{C}}^{*V}$, $\mathbf{c}$ and matrix $\mathbf{H}$ is considered in Appendix C.

4 Numerical inversion

Statistical regularization technique (e.g. Turchin et al., 1971; Tarantola, 1987; Rodgers, 2000) considers errors, $\mathbf{\Delta}_L$ and $\mathbf{\Delta}_V$, in Eqs. (3) and (18) as random variables. Under the additional assumption that errors have independent normal distributions, the multidimensional conditional probability density function (PDF) (or "likelihood function") is defined by (Chaikovsky et al., 2004a)

$$F\left(\mathbf{L}^*, \hat{\mathbf{C}}^{*V} | \mathbf{c}\right) \sim \exp\left[\left(-\frac{1}{2}\left(\left(\mathbf{L}^* - \mathbf{L(c)}\right)^T \mathbf{\Omega}_L^{-1}\left(\mathbf{L}^* - \mathbf{L(c)}\right) + \left(\hat{\mathbf{C}}^{*V} - \mathbf{H}\mathbf{c}\right)^T \mathbf{\Omega}_V^{-1}\left(\hat{\mathbf{C}}^{*V} - \mathbf{H}\mathbf{c}\right)\right)\right)\right] \tag{20}$$

1

2 Here, $F\left(\mathbf{L}^*, \hat{\mathbf{C}}^{*V} \middle| \mathbf{c}\right)$ is the PDF of measurement vectors $\mathbf{L}^*$ and $\hat{\mathbf{C}}^{*V}$, $\mathbf{L}(\mathbf{c})$ is the vector function in

3 Eq. (3), $\mathbf{H}$ is the matrix in Eq. (18), $\mathbf{c}$ is the target retrieval vector of aerosol modes concentration,

4 and $\mathbf{\Omega}_L$ and $\mathbf{\Omega}_V$ are the covariance matrices of error vectors $\mathbf{\Delta}_L$ and $\mathbf{\Delta}_V$, respectively.

5

6 An extensively used tool for the regularization of an "ill-posed" problem is the application of *a*

7 *priori* constraint on the smoothness of retrieved characteristics. LIRIC restricts the norms of the

8 second differences of functions $c_k(h_i)$. Following the statistical regularization approach (Turchin

9 et al., 1971) we included *a priori* probability function,

10
$$F_{apr}(\mathbf{c}) \sim \exp\left(-\frac{1}{2}\left(\mathbf{c}^T \mathbf{\Omega}_S \mathbf{c}\right)\right) \tag{21}$$

11 into the retrieval procedure as the additional constraint. Here, $\mathbf{\Omega}_S = \mathbf{S}_2^T \mathbf{Q}_2^{-1} \mathbf{S}_2$ is smoothing

12 matrix, $\mathbf{S}_2$ is the matrix of the second-order differences, and $\mathbf{Q}_2$ is diagonal weighting matrix

13 (Twomey, 1977; Dubovik et al., 2011).

14

15 The Bayes' strategy (Turchin et al., 1971; Tarantola, 1987; Rodgers, 2000) for solving an "ill-

16 posed" problem combined with multi-term LSM technique (Dubovik, 2004; Dubovik et al., 2011)

17 defines the solution $\hat{\mathbf{c}}$ in accordance with the maximum a posteriori rule

18

19
$$\hat{\mathbf{c}} = \arg \min_{\mathbf{c}} \{\Psi(\mathbf{c})\},$$

20

21 where the objective or cost function, $\Psi(\mathbf{c})$, has the following multi-term representation (Dubovik,

22 2004; Dubovik et al., 2011)

23
$$\Psi(\mathbf{c}) = \left(\mathbf{L}^* - \mathbf{L}(\mathbf{c})\right)^T \mathbf{\Omega}_L^{-1}\left(\mathbf{L}^* - \mathbf{L}(\mathbf{c})\right) + \left(\hat{\mathbf{C}}^{*V} - \mathbf{Hc}\right)^T \mathbf{\Omega}_V^{-1}\left(\hat{\mathbf{C}}^{*V} - \mathbf{Hc}\right) + \left(\mathbf{c}^T \mathbf{S}_2^T \mathbf{Q}_2^{-1} \mathbf{S}_2 \mathbf{c}\right) \tag{22}$$

24

1 We assume that the errors $\mathbf{\Delta}_L$ in Eq. (3) and $\mathbf{\Delta}_V$ in Eq. (18) are uncorrelated. In this case the non-

2 zero diagonal elements of the covariance matrices $\mathbf{\Omega}_L$ and $\mathbf{\Omega}_V$ are the variances of the elements of

3 the vectors $\mathbf{\Delta}_L$ and $\mathbf{\Delta}_V$, respectively.

4

5 Since the minimization procedure does not prescribe a residual value for $\Psi(\mathbf{c})$, it is convenient to

6 reformulate weight matrices as follows (Dubovik, 2004):

7
$$\breve{\mathbf{\Omega}}_L = \frac{1}{\varepsilon_L^2}\mathbf{\Omega}_L \; ; \; \breve{\mathbf{\Omega}}_V = \frac{1}{\varepsilon_V^2}\mathbf{\Omega}_V \; ; \; \breve{\mathbf{\Omega}}_S = \frac{1}{\varepsilon_S^2}\mathbf{\Omega}_S \; , \tag{23}$$

8 where ε_L^2, ε_W^2, and ε_S^2 are the first elements of the corresponding covariance matrices.

9

10 After substitution of the covariance matrices expressed through the weight matrices into Eq. (22)

11 and multiplication it by ε_L^2, the $\Psi(\mathbf{c})$ takes the form of the sum of three components:

12
$$\Psi(\mathbf{L}^*,\hat{\mathbf{C}}_V,\mathbf{c}) = \Psi_L(\mathbf{L}^*,\mathbf{c}) + \gamma_V\Psi_V(\hat{\mathbf{C}}^{*V},\mathbf{c}) + \gamma_S\Psi_S(\mathbf{c}) \tag{24}$$

13 where

14
$$\breve{\Psi}_L(\mathbf{L}^*,\mathbf{c}) = \left(\mathbf{L}^* - \mathbf{L}(\mathbf{c})\right)^T \breve{\mathbf{\Omega}}_L^{-1}\left(\mathbf{L}^* - \mathbf{L}(\mathbf{c})\right), \tag{25}$$

15 is related to "lidar-measured" data, Eq. (3),

16
$$\Psi_V(\hat{\mathbf{C}}^{*V},\mathbf{c}) = \left(\hat{\mathbf{C}}^{*V} - \mathbf{Hc}\right)^T \breve{\mathbf{\Omega}}_V^{-1}\left(\hat{\mathbf{C}}^{*V} - \mathbf{Hc}\right), \tag{26}$$

17 is related to radiometer-measured data, Eq. (18),

18
$$\Psi_S(\mathbf{c}) = \left(\mathbf{c}^T\mathbf{S}^T\breve{\mathbf{Q}}_2^{-1}S\mathbf{c}\right), \tag{27}$$

19 is related to *a priori* information, Eq. (21),

20
$$\gamma_V = \frac{\varepsilon_L^2}{\varepsilon_V^2} \; ; \; \gamma_S = \frac{\varepsilon_L^2}{\varepsilon_2^2} \tag{28}$$

21

22 The coefficients γ_V and γ_S are so-called *Lagrange* multipliers that determine the weight of

23 different contributors from each source of information (i.e., "measurements" and "*a-priori*"

24 contribution) to the retrieval solution relative to the contribution of the first data source ((since

25 $\gamma_L = 1$). Equations (22) and (24) are equivalent; however Eq. (24) is more convenient for the

26 analysis of the relative contribution from different data source.

1

2 If $\gamma_V, \gamma_S \to 0$, we return to a non-regularized solution for vector **c** that is based solely on

3 measured lidar data with the minimum discrepancy between measured and calculated input signals.

4 This solution, however, could be non-physical, multivalued, and unstable. The possible solution

5 space should be restricted by increasing the *Lagrange* multipliers despite the fact that it results in

6 increasing of discrepancy between measured and model signal. The algorithms to determine the

7 *Lagrange* multipliers by finding a reasonable compromise between the solution quality and the

8 closeness of the measured and model signals are described in Hansen, 2001; Vogel, 2002; and

9 Doicu et al., 2010. The set of Lagrange multipliers is provided to LIRIC's users along with software

10 package. However, we do not consider this set as the ultimate one allowing its modification to meet

11 user's specifications.

12

13 The final step of the retrieval procedure is calculation of the concentration profiles $c_k(h_i)$ for each

14 aerosol mode. Initial approximations $c_k^0(h_i)$ are set and stepwise improved to provide the minimum

15 of the objective function (Eq. 25). Increments are calculated by means of the Levenberg-Marquardt

16 method (Levenberg, 1944; Marquardt, 1963).

17

18 The analytical expressions of the terms of Eq. (25), the covariance matrices, as well as the details of

19 the inversion procedure, are described in Appendices A, B and C.

20

21 **5 Program package for processing combined lidar and radiometer data**

22 Figure 2 shows the structure of the software package that implements the LIRIC algorithm. A set of

23 specific programs are joined in three sub-packages.

24

25 The sub-package *LiOpt* implements module (2) of the LIRIC algorithm (Fig.1), which provides

26 preprocessing of the AERONET retrieval products. Program *AERLID* recalculates the columnar

27 optical characteristics for the lidar sounding wavelengths, including the elements of the scattering

28 matrices for the spherical and non-spherical particles as well as for fine and coarse aerosol modes.

29 Then, this code writes data down to the Radiometer Database.

30 The preprocessing of lidar data is carried out by the *SignalSuite* sub-package. It contains several

programs. Among them:

- *ULIS* is an operational program that provides measurement procedures and record of raw lidar data to Microsoft ACCESS database (*DB Lidar raw*);

– *nc2mdb* is a program to convert EARLINET standard raw-lidar nc-files into mdb-files to process by LIRIC;

– program *Synthesizer* averages the series of lidar signals, converts the profiles to the optimal altitude scale and, then, "glues" signals (i.e. synthesizes single signal) for the upper and lower troposphere, which were measured with different receiving systems, as well as provides the "dead-time" correction, i.e. the correction for the finite time resolution of the photo-counting system;

– program *Tropoexport* calculates a normalized smoothed lidar signal and its variance, and generates molecular and aerosol atmospheric models; this program aims at implementing modules 2 and 3 of the algorithm.

Finally, the main sub-package *ProfileRetriever* implements the LIRIC inversion procedure. Program *ConcentRetriever* retrieves profiles $c_V^{k,m}(h)$ of the aerosol mode concentrations and writes data down to Access database, *DB-processed*. Module *Inversion setting&Errors modeling* generates a set of noise-corrupted input data files by adding "white noise" and amplitude distortions to the initial lidar signals and perturbing aerosol model parameters retrieved from radiometer measurements in order to provide the error sensitivity analysis. Real measurement conditions, technical features of the lidar system and the accuracy of columnar aerosol parameters retrieved from the radiometer measurements (Dubovik et al., 2000a, b) are taken into account in setting parameters of the module. The user can upgrade default instrumental noise parameters to meet real measurement conditions and technical features of the lidar system; the accuracy of columnar aerosol parameters retrieved from the radiometer measurements (Dubovik et al., 2000a, b) is also taken into account in setting parameters of the module. Program *OutputViewer* allows viewing the output data and their conversion from mdb-files into other formats.

6 Verification of operability and sensitivity tests

The LRS technique uses the aerosol model that was initially developed in AERONET to describe column-averaged aerosol properties and generalized it to the case of the height-resolved aerosol concentrations. This model assumes that aerosol consists of fine and coarse modes and that both are

mixtures of spherical particles and randomly-oriented homogeneous spheroids. The advanced T-matrix code (Mishchenko al., 2000; 2002) provides computation of scattering matrices of the aerosol particles. Thus, any optical characteristic of the aerosol layer can be calculated using data of the LRS experiment.

The applicability analysis of the AERONET spheroid model to aerosol particles is beyond the scope of this paper. We only note that this model was validated by the comparison of calculated optical parameters and laboratory measurements of light scattering matrices for mineral dust particles (Volten et al., 2001). Incorporation of the spheroid model into AERONET operational retrieval code has significantly improved AERONET products when evaluating parameters of coarse non-spherical particles (Cattrall et al., 2005; Dubovik et al., 2006). This model has also been incorporated when processing data from ground-based polarimetric measurements (e.g. Li et al., 2009), lidar sounding data (e.g. Veselovskii et al., 2010; David et al., 2013; Müller et al., 2013), and satellite-base observations (e.g. Levy et al., 2007a,b; Dubovik et al., 2011; Schuster et al., 2012).

6.1 Verification of LIRIC program package: EARLI09 intercomparison experiment

EARLI09 intercomparison experiment was held in May, 2009 at Leibniz Institute for Tropospheric Research in Leipzig, Germany (Wandinger et al., 2015). This campaign provided excellent opportunity to validate the LRS technique for network measurements. The results of the LIRIC data processing for simultaneous measurements by seven lidars of different scientific teams on May 25, 2009 in Leipzig were compared.

Total optical depth distribution (Fig. 3a) and back-trajectories analysis (Fig. 3b) indicates that LRS measurements were carried out during the Saharan dust event in the Leipzig region and the dust was transported in the layer above 2 km.

Figures 4 and 5 show PVC profiles, $c_k(h)$, retrieved from lidar data of the different EARLINET teams combined with the same AERONET information, as well as their root mean square deviations and relative deviations for the two types of input data set, namely, with and without depolarization measurements.

It is evident from Fig. 4a-c and Fig. 5a-b that $c_k(h)$ profiles have similar structure over the troposphere except for the lower layer. The relative deviations increase mainly when values of the aerosol concentration become negligible. The discrepancies are also possible in the near-surface atmospheric layer due to overlap effect (e. g., for the Hamburg lidar system, Fig. 4a).

We explain the discrepancy between $c_k(h)$ profiles in the near-surface atmosphere by the uncertainty in geometrical overlap factors and the differences in lower-boundary heights of the considered lidar systems. Also some differences in the retrieved concentration profiles $c_k(h)$ are due to measurement errors and uncertainties in aerosol modeling as well as specificities on the inverse operator (see Appendix C).

The potential errors in the PVC profiles for the specific combined lidar/radiometer experiment were estimated by using the *Errors modeling* module of the LIRIC package (Fig. 2). Figs. 6 and 7 illustrate the sensitivity of the retrieved aerosol concentration profiles to the errors of the lidar measurements. The original lidar signals were taken as they measured by München lidar (curves 4 in Fig. 6) and have been perturbed by adding "white noise" with different root-mean-square deviations (rms-deviations), α_j, and have been distorted by multiplying them by the coefficient,

$$k_j(h_i) = 1 + \frac{\Delta_j}{100} \frac{h_{ref} - h_i}{h_{ref}} \tag{29}$$

where percentage parameter Δ_j determines the amount of non-linearity.

In response, the program module generated twelve disturbed lidar signal sets that allowed us to estimate the impact of measurement errors. As an illustration, Fig. 6and 7 simulate higher errors that typical ones in most EARLINET lidars. Four realizations of the disturbed signals are shown in Fig. 6. Coefficient $k_j(h_i)$ increases/decreases from referent to start point that results in divergence of the lidar signals in Fig. 6.

1   PVC profiles, $c_k(h)$, corresponding to the lidar signals in Fig. 6 and their rms-deviations calculated

2   for full ensembles of input data are shown in Fig. 7. Changes in the PVC profiles of the dominant

3   coarse non-spherical mode are shown by the Fig. 7 to be minor (Fig. 7c). Although profiles $c_k(h)$

4   of fine and coarse spherical particles (Fig. 7a and 7b) are not very stable; they qualitatively retain

5   similarity with the initial distributions.

6

7   Figure 8 illustrates the effect of uncertainties in columnar aerosol parameters retrieved from

8   radiometer data. Variations of the columnar aerosol characteristics lead to changes in coefficients a

9   and b of lidar-related Eqs. (14)-(17) (Sect. 3.1). Statistical characteristics of aerosol concentration

10   profiles retrieved with relative deviation of the parameter $\vartheta_{k,p}^j$ (effective lidar ratio of the aerosol

11   fraction, see Appendix B) in the range $\pm 20\%$ (the full range) are presented in Fig. 8. Relative

12   deviation of aerosol concentration profile becomes significant only for small values of the

13   concentration.

14

15   **6.2 Dependence of retrieved aerosol concentration profiles on the content of the**

16   **input data set**

17   Three types of data set related to different sources of information compose the LIRIC input data-

18   file: three or four measured lidar signals, column-aerosol parameters from radiometer

19   measurements, and *a priori* smoothness constraints. Two- or three-mode aerosol models are used

20   according to the type of the measured lidar signals. Formally, we deal with redundant input

21   information and, hence, for the accepted aerosol model the number of input data set can be

22   decreased. Consequently, the significance of the different information components in retrieval

23   procedure is of interest as well as variations of the retrieved profiles, $c_k(h)$, in the absence of some

24   input data

25

26   As pointed in Sect. 4, the objective functions of LIRIC regularization algorithm (Eq. 22), consists of

27   a set of terms, which implement contribution of different types of input data into the retrieval

28   process. Setting the variance of the specific kind of measurement to a large value implies neglecting

29   of the correspondent term in the objective function (Eq. 22) and elimination of this part of the input

30   data in estimation of the final aerosol parameters. Program package implements this option and

1   makes allowing one to analyze the contribution of different measured data in the processing

2   procedure of specific experiment.

3

4   Below we shortly examine sensitivity of the retrieved profiles, $c_k(h)$, to the input data selection for

5   the case of combined lidar/radiometer sounding of the atmospheric aerosol during the last period of

6   Eyjafjallajökull volcano ash transport to the European area in Lille, France, on the 19th of May,

7   2010. Air mass back-trajectories (Fig. 9) forecasted the possibility of appearance of volcanic ash in

8   the layer between 1300 and 2500 m. The structure of the retrieved profiles, $c_k(h)$, shown in Fig.

9   10a agrees well with the forecast. Deviations (by "deviations" hereinafter we mean "standard

10   deviation") ($\delta(c_k(h_i))$ associated to the profiles $c_k(h)$ have been calculated by an "error

11   modelling" procedure similar to the one described in section 6.1.

12

13   A mixture of spherical and non-spherical particles constitutes the aerosol layer at the height of

14   about 2000 m. The profile of particle depolarization ratio at 532 nm and its deviation have been

15   calculated from the retrieved aerosol mode concentrations, $c_k(h)$. The profiles are shown in Fig.

16   10b, curves D(2) and rms_dev(2). The results of the direct calculation of depolarization ratio and

17   their deviations from lidar measurements are presented by curves D(1) and rms_dev(1). It should be

18   noted that the lidar measurements included additional calibration measurement that was not used by

19   the retrieval procedure. Profiles D(1) and D(2) show rather close agreement in magnitude and

20   vertical structure that could confirm the efficiency of the aerosol modeling used in this study.

21

22   The curves in Fig. 11 show the deviations in the retrieved concentration profiles, $c_k(h)$, after

23   elimination one of the lidar signals or columnar volume concentrations of aerosol modes, $\mathbf{C}^V$, from

24   the input data set. As can be seen from Fig.11, the concentration profile of the fine-particle mode

25   undergoes minor changes upon elimination of a single lidar signal or columnar volume

26   concentrations. This ~~implies~~ could imply that our experiment well-defined with respect to the fine-

27   mode concentration. On the other hand, concentrations of coarse modes are sensitive to input

28   information. Thus, lidar data at 1064 nm wavelength plays a crucial role in the retrieval of the

29   coarse spherical mode. In the same manner, lidar depolarization measurement is the key factor in

the retrieval of the coarse spheroid particle mode. Evaluations of columnar volume concentrations from radiometer measurement are necessary for all cases.

Fig. 12a shows concentration profiles, $c_k(h)$, which were retrieved for two- and three-mode aerosol models and characterized the aerosol layer in the same LRS experiment. The fine-mode concentration profiles for two aerosol models are practically coincident. Profiles $c_k(h)$ of coarse modes for two-mode aerosol model, coarse(2), and the sum of two coarse components for three-mode aerosol model, coarse(3), are similar in shape but quantitatively are a bit different. The column concentrations of the course (2) and (3) modes are equal.

The curves in Figs. 12b and c show the deviations of the concentration profiles, $c_k(h)$, for the two-mode aerosol model after reduction of the input data set. Deviations of $c_k(h)$ profiles are rather similar to those for the three-mode aerosol model in Fig. 11. Deviations of fine-mode concentration profile are small, even if any single sub-set of input data is eliminated. Coarse-mode concentration profiles preserve original forms when one of the lidar signals at the 355 or 532 nm wavelength is excluded from the processing procedure.

Generally, for measurement conditions that characterize the experiment under discussion, two-wavelength lidar sounding (at 355 and 1064 or at 532 и 1064 nm) combined with radiometer measurement provides retrieving concentration profiles of fine and coarse aerosol modes for two-mode aerosol model.

7 Discussion and conclusions

The active process of dissemination of the LIRIC in EARLINET started in 2012. Nowadays, 11 EARLINET teams participate in implementation of LRS technique (see Fig. 13). New scientific teams beyond EARLINET join the LIRIC user group. The detailed description of LIRIC algorithm and software in this paper should contribute to the effective implementation of the LRS technique by advanced users.

Retrieval of the aerosol parameters from the LRS measurements is "ill-posed" inverse problem and its solution should be tested on stability to the measurement errors and variations of the regularization parameters, which are set by the module *Inversion setting&Errors modeling* of the software package (Fig. 2). Results of the EARLI09 intercomparison experiment presented in Sect. 6.1 demonstrate rather small scatter in $c_k(h)$ profiles that were retrieved from the data of different lidar systems with significantly corrupted input lidar signals and big uncertainties of the aerosol lidar ratio. This scatter is characterized by standard deviations of $5 - 20\%$ of the maximum aerosol layer concentration. Increase in $c_k(h)$ deviation in the bottom layer results from uncertainties of the overlap function of the lidar systems.

The uncertainties in the retrieved aerosol parameters for different aerosol types, aerosol loads, overlap characteristics of the lidar systems and regularization parameters that are defined by the LIRIC operator were evaluated by Granados-Muñoz et al. (2014). The analysis covered combined lidar and radiometer measurements that were carried out during dust, smoke, and anthropogenic pollution events. This analysis mostly supports our conclusions on the stability of LIRIC solutions that retrieve basic aerosol features even under significant measurement errors. In particular, variations of the regularization parameters within one order interval from the original set lead to minor deviations of the retrieved $c_k(h)$ profiles. Usually, it is unnecessary to change recommended utility regularisation parameters while homogeneous input data sets are processed. The requirements to pre-processing of lidar signals along with the set of recommended regularisation parameters are provided in the LIRIC user guide. However, the utility parameters for error modelling menu should be defined by the LIRIC user with regard to the specific lidar system.

The requirement of having possibly minimal "full overlap" height of lidar sensing is important technical problem for LRS measurements, because the near-surface aerosol layer contributes strongly to the radiometric data. In the absence of lidar data, the surface aerosol layer is assumed to be homogeneous in the LIRIC aerosol modelling. Obviously, aerosol parameters can vary within the near-surface layer resulting in significant uncertainties in the LIRIC product, especially when the lidar "dead zone" becomes comparable to the boundary-layer thickness. The effective solution

of this problem is the set-up of a double lidar receiving block with special near-range channels for the detection of near ground aerosol.

The analysis of the aerosol parameters that are retrieved from the incomplete sets of lidar data in Sect. 6.2 supports the possibility to use LIRIC for processing data of two-wavelength lidar systems. Aerosol sounding by two-wavelength lidars, usually at 532 and 1064 nm wavelengths, is a widespread practice in atmospheric investigations. Simulation results in Sect. 6.2 show the possibility to retrieve $c_k(h)$ for two-mode aerosol model. The uncertainties of such evaluated $c_k(h)$ are expected to surpass ones of three-wavelength lidar sounding.

LIRIC implementation for the special lidar data set (532-cross, 532-parallel and 1064 nm) for retrieving parameters of the three-mode aerosol model is of interest for the satellite lidar CALIOP that provides similar lidar data (Winker et al, 2006).

Since the beginning of LIRIC dissemination in EARLINET community, experimental works on the validation of the LIRIC product for different aerosol types have being carried out. Comparisons of aerosol backscatter coefficients and depolarization ratios directly derived from lidar data against similar characteristics calculated from the aerosol optical and microphysical parameters retrieved by LIRIC (e.g. Tsekeri et al., 2012, 2013; Wagner et all, 2013; Kokkalis et al., 2013; Granados-Muñoz et al., 2014) as well as LIRIC against modeled or airborne *in situ* measured profiles of aerosol mode concentrations (e.g. Kokkalis et al., 2012, 2013; Nemuc et al., 2013) have shown reasonable agreement.

The LIRIC concentration profiles of aerosol fractions during dust and volcano ash events have been compared with those for spherical and non-spherical particles derived from polarization measurements using the POLIPHON technique (e.g. Wagner et al., 2013; Nemuc et al., 2013, Papayannis et al., 2014). In spite of the noticeable difference between the aerosol models and independent processing algorithms, the retrieved aerosol concentration profiles have proved to be similar. This is quite natural because both approaches use the depolarization of backscatter signal to distinguish between spherical and non-spherical particles.

26

1   The number of aerosol studies using LIRIC algorithm increases. They focus on the investigation of

2   the dynamics of aerosol microstructure during transport of air masses polluted by dust (e.g.

3   Chaikovsky et al., 2010b; Tsekeri et al., 2013; Binietoglou et al., 2015; Granados-Muñoz et al.,

4   2015a) , fire smoke (e.g. Chaikovsky et al., 2010b; Granados-Muñoz et al., 2015a), volcano ash

5   (Kokkalis et al., 2013) and anthropogenic pollution (Granados-Muñoz et al., 2015b). LIRIC has

6   become a tool for validation of the modeling of aerosol transport in atmosphere (Binietoglou et al.,

7   2015; Granados-Muñoz et al., 2015b). EARLINET teams form the data-base of the results of

8   combined lidar and radiometer sounding.

9

10   The list of lidar teams that take advantage of LIRIC is still expanding. The LIRIC software package

11   is open and free-distributed within the EARLINET community and beyond it. The EARLINET

12   teams provide continuous improvement of the software and cooperate on the implementation of the

13   LRS measurements at new sites.

14

15   **Appendix A: General equation for received lidar signal**

16

17   Using general formula for received lidar signal instead of Eqs. (4), (6), and (8) allows us to derive

18   compact and explicit expression for the covariance matrices, $\mathbf{\Omega}_L$, and regularizing term, $\breve{\Psi}_L(\mathbf{L}^*,\mathbf{c})$

19   (Sect. 4).

20

21   We will use the utility function

22
$$\delta_{p_j,u}^{j} = \begin{cases} 1...if.....p_j = u \\ 0...f...p_j \neq u \end{cases} , \tag{A1}$$

23   along with the following definitions of combinations of aerosol and molecular optical parameters in

24   Eqs. (4) – (9):

25
$$\beta_a^{ef}(\lambda_j, p_j, h) = \left(\beta_{a,p_j}(\lambda_j, h) + \delta_{p_j,2}^{j} \mu \beta_{a,3}(\lambda_j, h) \right) =$$
$$= \left(\sum_k c_k(h) b_{k,p_j}(\lambda_j) + \delta_{p_j,2}^{j} \mu \sum_k c_k(h) b_{k,p_j}(\lambda_j) \right) \tag{A2}$$

1

$$\beta_r^{ef}(\lambda_j, p_j, h) = \left(\delta_{p_j,2}^j(p_j)\left(\frac{\mu-1}{\chi+1}\right) + \frac{1}{1+\delta_{p_j,3}^j\chi}\right)\beta_r(\lambda_j, h) \tag{A3}$$

2

$$\beta^{ef}(\lambda_j, p_j, h) = \beta_a^{ef}(\lambda_j, p_j, h) + \beta_r^{ef}(\lambda_j, p_j, h) \tag{A4}$$

3

4

$$\hat{R}_j^{ef}(\lambda_j, p_j, h) = \frac{\beta_a^{ef}(\lambda_j, p_j, h) + \beta_r^{ef}(\lambda_j, p_j, h)}{\beta_r^{ef}(\lambda_j, p_j, h)} \tag{A5}$$

5

$$\tau_a(\lambda_j, h, h_{ref}) = \int_h^{h_{ref}} \sigma_a(\lambda_j, h)dh \tag{A6}$$

6

7   This permits Eqs. (4), (5) and (8) to be written in general form:

8

9

$$L_j(p_j, \lambda_j, h) = \frac{\beta^{ef}(\lambda_j, p_j, h)\exp(2\tau_a(\lambda_j, h, h_{ref}))}{\beta_r^{ef}(\lambda_j, p_j, h_{ref})\hat{R}_j^{ef}(\lambda_j, p_j, h_{ref})}. \tag{A7}$$

10

11

12   Therefore, the related to the lidar objective function, $\breve{\Psi}_L(\mathbf{L}^*, \mathbf{c})$, (Eq. 25), is given by the equation:

13

$$\Psi_L(\mathbf{L}^*, \mathbf{c}) = \sum_j \sum_i \frac{\Delta h_i}{\breve{\Omega}_{L_j}(i,i)} \left(L_{j,i}^* - \frac{\left(\sum_k c_k(h_i)b_{k,p_j}(\lambda_j) + \delta_{p,2}^j\mu\sum_k c_k(h_i)b_{k,p_j}(\lambda_j)\right)}{\beta_r^{ef}(\lambda_j, p_j, h_{ref})\hat{R}_j^{ef}(\lambda_j, p_j, h_{ref})} \times \exp\left(2\sum_k \sum_i c_k(h_i)a_k(\lambda_j)\Delta h_i\right)\right)^2, \tag{A8}$$

14   $$i \in 1,...I.$$

15

16   Equation (26), $\breve{\Psi}_V(\hat{\mathbf{C}}^{*V}, \mathbf{c})$, which brings radiometer data into the processing procedure can be

17   expressed as follows:

18

$$\breve{\Psi}_V(\hat{\mathbf{C}}^{*V}, \mathbf{c}) = \sum_k \frac{1}{\breve{\Omega}_V(k,k)} \left(\hat{\mathbf{C}}^{*V} - \sum_i c_k(h_i)|\Delta h_i|\right)^2. \tag{A9}$$

28

1   Calculation of the "smoothness" part of the objective function is described in details in Dubovik,

2   2004; and Dubovik et al., 2011.

3

4   **Appendix B: Evaluation of covariance matrix Ω_L**

5   The covariance matrixes, $\mathbf{\Omega}_L$, $\mathbf{\Omega}_V$, and $\mathbf{\Omega}_2$, defined in Sect. 4 characterize uncertainties of the

6   complex input vector, $\left(\mathbf{L}^*, \hat{\mathbf{C}}^{*V}, \hat{\mathbf{0}}\right)$, where $\mathbf{0}^*$ is "zero" vector that is defined to formalize *a priori*

7   smoothness restrictions on concentration profiles (e.g. Dubovik, 2004). These matrices determine

8   the "weights" of different parts of input information through the minimization procedure of the

9   objective function (Eq. 22).

10

11   In our case the measure of the smoothness for concentration profiles, $c_k(h_i)$, should be chosen as *a*

12   *priori* evaluated parameters. Aerosol columnar volume concentrations, $\hat{\mathbf{C}}^{*V}$, and variances,

13    $\Omega_V(k,k)$, are the parts of input radiometer data. Thus, only evaluation of covariance matrix, $\mathbf{\Omega}_L$, is

14   to be done.

15

16   The assumption of independent normal distribution for variations of "lidar" vector, $\mathbf{L}^*$, at different

17   heights implies the diagonal covariance matrix. The non-zero diagonal elements, $\Omega_{L_j}(h_i, h_i)$, of the

18   covariance matrix are the variances of differences between the components, $L^*_{j,i}$, of the lidar vector

19   and the appropriate modeled function, $L_j(c_k, p_j, \lambda_j, h_i)$, in Eq. (A7).

20

21   Given Eqs. (2), (3) and (A1)-(A-8), the elements of vector, Δ_{L_j}, are defined by:

22
$$\Delta_{L_j}(h_i) = L^*_{j,i} - L_j\left(p_j, \lambda_j, h_i\right) =$$
$$= \frac{S^*\left(\lambda_j, h_i\right)}{\hat{S}^*\left(\lambda_j, h_{ref}\right)} \exp(-2\tau_r(\lambda_j, h_i, h_{ref})) - \frac{\beta^{ef}\left(\lambda_j, p_j, h_i\right)\exp\left(2\tau_a(\lambda_j, h_i, h_{ref})\right)}{\beta_r^{ef}(\lambda_j, p_j, h_{ref})\hat{R}_j^{ef}(\lambda_j, p_j, h_{ref})} \quad \text{(B1)}$$

23

1   Using the finite differences technique (e.g. Russell et al., 1979) one can expand $\Delta_{L_j}(h_i)$ in *Taylor*

2   series, and then neglect all the terms of the second or higher order. As a result, variation $\delta(\Delta_{L_j}(h_i))$

3   can be expressed as a function of variations related with the input parameters, $\delta\!\left(S^*\!\left(\lambda_j, h_i\right)\right)$,

4    $\delta(\beta^{ef}(\lambda_j, p_j, h_i))$, $\delta(\tau_a(\lambda_j, h_i, h_{ref}))$, and $\delta(\tau_r(\lambda_j, h_i, h_{ref}))$:

$$
\begin{aligned}
\delta(\Delta_{L_j}(h_i)) ={}& -2\,\mathrm{L}^*_{j,i}\,\delta(\tau_r(\lambda_j, h_i, h_{ref})) + \mathrm{L}^*_{j,i}\,\frac{\delta\!\left(S^*\!\left(\lambda_j, h_i\right)\right)}{S^*\!\left(\lambda_j, h_i\right)} + \\[2mm]
&+\frac{\beta^{ef}\!\left(\lambda_j, p_j, h_i\right)\exp\!\left(2\tau_a(\lambda_j, h_i, h_{ref})\right)}{\beta_r^{ef}(\lambda_j, p_j, h_{ref})\hat{\mathrm{R}}_j^{ef}(\lambda_j, p_j, h_{ref})}\,\frac{\delta\!\left(\beta^{ef}\!\left(\lambda_j, p_j, h_i\right)\right)}{\beta^{ef}\!\left(\lambda_j, p_j, h_i\right)} + \\[2mm]
&-2\frac{\beta^{ef}\!\left(\lambda_j, p_j, h_i\right)\exp\!\left(2\tau_a(\lambda_j, h_i, h_{ref})\right)}{\beta_r^{ef}(\lambda_j, p_j, h_{ref})\hat{\mathrm{R}}_j^{ef}(\lambda_j, p_j, h_{ref})}\,\delta(\tau_a(\lambda_j, h_i, h_{ref})) \approx \\[2mm]
\approx{}& \mathrm{L}^*_{j,i}\left(\frac{\delta\!\left(S^{*j}(h_i)\right)}{S^{*j}(h_i)} - \frac{\delta(\beta^{ef}(\lambda_j, p_j, h_i))}{\beta^{ef}(\lambda_j, p_j, h_i)} - 2\delta\!\left(\tau_r(\lambda_j, h_i, h_{ref})\right) - 2\delta\!\left(\tau_a(\lambda_j, h_i, h_{ref})\right)\right).
\end{aligned}
\tag{B2}
$$

6

7   Under the assumption of independent variations of different parameters, the variance $\Omega_L(h_i, h_n)$ is

8   expressed as follows

$$
\begin{aligned}
\Omega_L(h_i, h_i) = {}& \left\langle \delta(\Delta_{L_j}(h_i))\delta(\Delta_{L_j}(h_i)) \right\rangle = \\[2mm]
={}& \mathrm{L}^{*2}_{j,n}\left(\frac{\delta^2\!\left(P^*_{j,i}\right)}{\left(P^*_{j,i}\right)^2} + \frac{\delta^2\!\left(\beta^{ef}(\lambda_j, p_j, h_i)\right)}{\left(\beta^{ef}(\lambda_j, p_j, h_i)\right)^2} + 4\delta^2\!\left(\tau_r(\lambda_j, h_i, h_{ref})\right) + 4\delta^2\!\left(\tau_a(\lambda_j, h_i, h_{ref})\right)\right),
\end{aligned}
\tag{B3}
$$

10   where $\langle ... \rangle$ denotes ensemble averaging over measurement realizations, and $P^*_{j,i} = P^*_j(h_i)$,

11   The terms in the large round parentheses in Eq. (B3) determine contributions of measurement errors

12   and uncertainties of *a priori* defined optical characteristics. We aim at approximate estimation of

13   $\Omega_{L_j}(h_i, h_i)$ at the preprocessing stage without involving of retrieved parameters. This feedback-free

14   approach greatly simplifies the structure of the inversion algorithm.

15

16   *Uncertainties of the optical parameters*

1   The term $\delta^2\left(\beta^{ef}\right)\big/\left(\beta^{ef}\right)^2$ in Eq. (B3) is the relative variance of the total backscatter coefficient. It

2   can be transformed into the sum of relative variances of aerosol and molecular backscatter

3   coefficients:

4

$$\frac{\delta^2\left(\beta^{ef}(\lambda_j,p_j,h_i)\right)}{\left(\beta^{ef}(\lambda_j,p_j,h_i)\right)^2} = \frac{\delta^2\left(\beta_r^{ef}(\lambda_j,p_j,h_i)\right)}{\left(\beta_r^{ef}(\lambda_j,p_j,h_i)\right)^2}\frac{1}{\left(\hat{R}_j^{ef}(\lambda_j,p_j,h_i)\right)^2} + \frac{\delta^2\left(\beta_a^{ef}(\lambda_j,p_j,h_i)\right)}{\left(\beta_a^{ef}(\lambda_j,p_j,h_i)\right)^2}\frac{\left(\hat{R}_j^{ef}(\lambda_j,p_j,h_i)-1\right)^2}{\left(\hat{R}_j^{ef}(\lambda_j,p_j,h_i)\right)^2}$$

5   (B4)

6

7   The International Standard Atmosphere ISO 2533 and seasonal latitudinal changed model CIRA

8   (Committee on Space Research (COSPAR), 2012; Fleming et al., 1988), as well as measurements

9   by radiosondes are applied in LIRIC for the calculation of molecular optical parameters. The

10   relative variance of calculated molecular backscatter coefficient

11

$$\alpha_1^2 = \frac{\delta^2\left(\beta_r^{ef}(\lambda_j,p_j,h_i)\right)}{\left(\beta_r^{ef}(\lambda_j,p_j,h_i)\right)^2} \tag{B5}$$

12   is assumed to be a constant and its value can be reduced to $\alpha_1 = 0.01$ (e.g. Russell et al., 1979) if

13   data of coordinated radiosonde measurements is available.

14

15   The aerosol backscatter coefficients, $\beta_a^{ef}(\lambda_j,p_j,h_i)$, are estimated by using Eqs. (10)-(17).

16   Uncertainties of $\beta_a^{ef}(\lambda_j,p_j,h_i)$ basically follow from estimation errors of the coefficient

17    $b(v,j,p_j)$ in Eqs. (15) – (17) that can be written by equation:

18

$$b(j,p_j,k) = \frac{1}{\vartheta_{k,p}^j}\frac{E_k^*(\lambda_j)}{\hat{C}_k^V} \tag{B6}$$

19   where

20

$$\frac{1}{\vartheta_{k,p}^j} = \frac{1}{4\pi}\varpi_k(\lambda_j)A_{k,p}^j \ , \tag{B7}$$

$$
1 \qquad A_{k,p}^{j} = \begin{cases} P_{1,1}^{v}(\lambda_{j}, \gamma = 180^{0}) & \text{if} \quad p_{j} = 1 \\ \dfrac{P_{1,1}^{k}(\lambda_{j}, \gamma = 180^{0}) - P_{2,2}^{k}(\lambda_{j}, \gamma = 180^{0})}{2} & p_{j} = 2 \\ \dfrac{P_{1,1}^{k}(\lambda_{j}, \gamma = 180^{0}) + P_{2,2}^{k}(\lambda_{j}, \gamma = 180^{0})}{2} & p_{j} = 3 \end{cases} \qquad \text{(B8)}
$$

2

3   Parameter $\vartheta_{k,p}^{j} = \sigma_{a}^{k}(\lambda_{j}, h_{i}) / \beta_{a}^{ef,k}(\lambda_{j}, h_{i})$ in Eq. (B8) is the extinction-to-backscatter ratio or "lidar

4   ratio" of the k-aerosol mode.

5

6   Parameters $\vartheta_{k,p}^{j}$ are retrieved from the data of radiometric direct Sun and almucantar measurements

7   that are usually performed with the maximum scattering angle less than 150^{0}. The range of the

8   scattering angles decreases with decreasing the sun zenith angle. Retrieval of optical parameters in

9   the backscatter direction, in a certain sense, is an extrapolation procedure out of the measured range

10 with possible increasing of estimation uncertainties. One assumes that the errors of the estimation of

11 $\vartheta_{k,p}^{j}$ are the main reason of the incorrect calculation of backscatter coefficients $\beta_{a}^{ef}(\lambda_{j}, p_{j}, h_{i})$ and

12 introduces parameter α_{2} for characterization of the standard deviation of coefficients $1 \big/ \vartheta_{k,p}^{j}$ in

13 LRS measurements.

14

15 Thus, Eq. (B4) is transformed to

$$
16 \qquad \frac{\delta^{2}\left(\beta^{ef}(\lambda_{j}, p_{j}, h_{i})\right)}{\left(\beta^{ef}(\lambda_{j}, p_{j}, h_{i})\right)^{2}} = \frac{\alpha_{1}^{2}}{\left(R_{j}^{ef}(\lambda_{j}, p_{j}, h_{i})\right)^{2}} + \alpha_{2}^{2}\frac{\left(R_{j}^{ef}(\lambda_{j}, p_{j}, h_{i}) - 1\right)^{2}}{\left(R_{j}^{ef}(\lambda_{j}, p_{j}, h_{i})\right)^{2}} \; . \qquad \text{(B9)}
$$

17

18 The backscatter ratio $R_{j}^{ef}(\lambda_{j}, p_{j}, h_{i})$ in Eq. (B9) under assumption $\mu = 0$ is approximately

19 calculated at the pre-processing stage using the Klett algorithm (Klett, 1981).

20

21 Basically, the variance of aerosol optical thickness, $\delta^{2}\left(\tau_{a}(\lambda_{j}, h_{i}, h_{ref})\right)$, arises from altitude

22 variations of aerosol modes that are not assumed by the aerosol model. Relative error of

23 $\tau_{a}(\lambda_{j}, h_{i}, h_{ref})$ is zero at the reference point and is equal to α_{3}^{2} at the start point h_{1}, where α_{3} is

close to the error of AOT calculation from radiometer measurements. Thus, the following approximation is used in the LIRIC algorithm:

$$\delta^2\left(\tau_a(\lambda_j, h_i, h_{ref})\right) = \alpha_3^2 \tau_a^2(\lambda_j, h_i, h_{ref}).$$ (B10)

Term $\delta^2\left(\tau_r(\lambda_j, h_i, h_{ref})\right)$ in Eq. (B3) denotes the variance of molecular optical thickness of the atmospheric layer (h_n, h_{ref}). Only long scale or systematic deviations of molecular density contribute to the variance $\delta^2\left(\tau_r(\lambda_j, h_i, h_N)\right)$. Similar to (B10),

$$\delta^2\left(\tau_r(\lambda_j, h_i, h_{ref})\right) = \alpha_4^2 \tau_r^2(\lambda_j, h_i, h_{ref}).$$ (B11)

Measurement errors

Optical signals, detected by the lidar data acquisition system consists of backscatter $P_{j,i}^*$ and background B_j^* components. A suitable algorithm for estimating the measurement errors is described by Slesar et al. (2013, 2015). Regardless of the type of the photo-receiving sensor, three factors determine the measurement errors:

- non-linearity of the recording channel, which consist of nonlinearity of the photodetector and electronic units;
- "nonsynchronous" noise (non-correlated with the sounding pulse);
- "synchronized" noise (correlated with the sounding pulse).

Non-linearity of a receiving channel basically originates from saturation of an output signal at high incident light because of photo-sensor or electronic unit limitations. Likewise, deviations of an amplifier gain cause linear distortions of the detecting signal within the working range of photo-receiving module.

Basic difference between two types of noise is that "nonsynchronous" noise can be reduced by accumulation of input signals or by decreasing frequency bandwidth of the receiving channel, while this method is ineffective for "synchronized" noise. The main type of the "nonsynchronous" noise is the Schottky noise. "Synchronous" noise is basically caused by the interference of the electrical impulses from the laser power supply, synchronous with the sounding optical pulse. It is predominantly a low-frequency noise, and acceptable limitation of the frequency band of the photo-receiving channel does not lead to its decline.

1

2 We assume that the accumulation of the receiving lidar signal with A sounding pulses and averaging

3 of the lidar signal over 2M +1 bins are carried out at the measurement and pre-processing stages.

4

5 Summing up the contributions of the noise components, one can write the following expression for

6 the variances of the receiving analog and photon-counting signals (Slesar et al., 2013, 2015):

7 - for analog channel:

8
$$\frac{\delta^2(P_{j,i}^*)}{(P_{j,i}^*)^2} = \omega_j^2 \frac{\left(P_{j,n}^* + B_j^*\right)^2}{(P_{j,n}^*)^2} + \frac{\left(G_j^*\right)^2 + q_j^2\left(P_{j,n}^* + B_j^*\right)}{A(2M+1)(P_{j,n}^*)^2} + \frac{\left(U_j^*\right)^2}{(P_{j,n}^*)^2} , \qquad (B12)$$

9 where ω is the coefficient of nonlinearity, G_j^* is the amplitude of electrical noise, q_j^2 is the

10 coefficient characterizing the power of the Schottky noise, U_j^* is the amplitude of "synchronized"

11 noise;

12 - for counting channel:

13
$$\frac{\delta^2(N_{j,i}^*)}{(N_{j,i}^*)^2} = \omega_j^2 \frac{\left(N_{j,i}^* + N_{j,B}\right)^2}{(N_{j,i}^*)^2} + \frac{\left(N_{j,G}^*\right)^2 + N_{j,i}^* + N_{j,B}}{A(2M+1)(N_{j,i}^*)^2} + \frac{\left(N_{j,U}^*\right)^2}{(N_{j,i}^*)^2} , \qquad (B13)$$

14 where $N_{j,i}^*$ is the detected lidar signal, $N_{j,B}$ is the background signal, $N_{j,G}$ is the external

15 "nonsynchronous" noise, and $N_{j,U}^*$ is "synchronized" noise.

16

17 Parameters ω_j, G_j^*, $N_{j,G}^*, q_j, U_j^*$, $N_{j,G}^*$, $N_{j,U}^*$ for specific photo-receiving module can be

18 evaluated on a dedicated test bench by means of special calibration procedures (Slesar et al., 2013,

19 2015).

20

21 **Appendix C: Details of inversion procedure**

22 One can understand intuitively that optical parameters of aerosol modes, which constitute the

23 aerosol model (see Sect.1.1), should be different to allow retrieving aerosol mode concentrations by

24 means of algorithm described in Sect. 4. More corrected definition of this requirement is deficiency

25 of linear relation that there should not be a linear relationship between the sets of coefficients

26 $\{\mathbf{a}_k, \mathbf{b}_k\}$ which define optical characteristics of k th aerosol mode. This conclusion results from the

1 linear approximation of the Eq. (2). It means that we seek the solution, $\mathbf{c}(h_i)$, only from data of

2 multi-wavelength lidar sounding. The linear least squares solution of Eq. (3) can be written as

3 $$\mathbf{c} = \left(\mathbf{K}_L^T \boldsymbol{\Omega}_L^{-1} \mathbf{K}_L\right)^{-1} \mathbf{K}_L^T \boldsymbol{\Omega}_L^{-1} \mathbf{L}^*, \tag{C1}$$

4 where $\mathbf{K}_L$ is the *Jacobi* matrix of the first partial derivatives $\left\{\mathbf{K}_L\right\}_{x,y} = \delta L_x / \delta c_y \big|_{\mathbf{c}}$. The following

5 definitions are used in Eq. (C1) for *measured vector,* $\hat{\mathbf{L}}^*$, and *state vector*, $\mathbf{c}$, with dimensions

6 $JI \times 1$ and $KI \times 1$, correspondently:

7
$$L_{j,i}^* = \begin{cases} L_1^*(h_1) \\ L_2^*(h_1) \\ ... \\ L_j^*(h_i) \\ ... \\ L_J^*(h_I) \end{cases}, \qquad c_{k,i} = \begin{cases} c_1(h_1) \\ c_2(h_1) \\ ... \\ c_k(h_i) \\ ... \\ c_K(h_I) \end{cases}, \tag{C2}.$$

8 The formula (C1) is valid if $\det\left(\mathbf{U}_L = \mathbf{K}_L^T \boldsymbol{\Omega}_L^{-1} \mathbf{K}_L\right) \neq 0$.

9

10 We use additional requirements that optical thickness of the aerosol layer is small, and the variances

11 of the measured errors, $(\varepsilon_{L,1}^2, \varepsilon_{L,i}^2, \varepsilon_{L,I}^2)$, do not depend on h_i. So the matrix $\mathbf{U}_L = \mathbf{K}_L^T \boldsymbol{\Omega}_L^{-1} \mathbf{K}_L$ with

12 dimensions $KI \times KI$ takes the block-diagonal form

13
$$\mathbf{U}_L = \begin{bmatrix} \frac{1}{\varepsilon_1^2} \mathbf{U}_{k,k^*} & ... & & \\ & ... & & \\ ... & \frac{1}{\varepsilon_i^2} \mathbf{U}_{k,k^*} & ... & \\ & ... & & \\ & & ... & \frac{1}{\varepsilon_I^2} \mathbf{U}_{k,k^*} \end{bmatrix}, \tag{C3}$$

14 where matrix $\mathbf{U}_{k,k^*}$, ($k \in 1,..,K$), does not depend on the superscript i. For 3-mode aerosol model

15 ($K = 3$) and 4-channel lidar measurements ($J = 4$) matrix $\mathbf{U}_{k,k^*}$ can be written

1
$$\mathbf{U}_{k,k^*} = \begin{bmatrix} \sum_j b^2[j,p_j,1] & \sum_j b[j,p_j,1]b[j,p_j,2] & \sum_j b[j,p_j,1]b[j,p_j,3] \\ \sum_j b[j,p_j,1]b[j,p_j,2] & \sum_j b^2[j,p_j,2] & \sum_j b[j,p_j,2]b[j,p_j,3] \\ \sum_j b[j,p_j,1]b[j,p_j,3] & \sum_j b[j,p_j,2]b[j,p_j,3] & \sum_j b^2[j,p_j,3] \end{bmatrix} . \quad \text{(C4)}$$

2

3 Thus, results of the retrieval depend on the specifics of matrix $\mathbf{U}_{k,k^*}$. The well-conditioned matrix

4 $\mathbf{U}_{k,k^*}$ provides suitable solution of Eq. (3). On the analogy with Veselovsky et al. (2005), the

5 eigenvalue decomposition technique has been used to evaluate the "condition number" of matrix

6 $\mathbf{U}_{k,k^*}$

7
$$Cond\{\mathbf{U}_{k,k^*}\} = |\psi_{\max}|/|\psi_{\min}| \qquad \text{(C5)}$$

8 where $\psi_{\max}$ and $\psi_{\min}$ are the maximum and minimum eigenvalues of matrix $\mathbf{U}_{k,k}$, respectively.

9 Parameter $\sqrt{Cond\{\mathbf{U}_{k,k^*}\}}$ is a coefficient of increasing relative error of $c_{k,i}$ as compared to the

10 relative error of $L_{n,j}^*$ estimation (Trefethen and Bau, 1997).

11

12 The data of radiometric measurements in Minsk during 2002 – 2010 were used to calculate the

13 parameters $Cond\{\mathbf{U}_{k,k^*}\}$ for the aerosol models with 2 and 3 aerosol fractions (3 and 4 measuring

14 channels, correspondently). The Cumulative Distribution Functions (CDF) of parameter

15 $Cond\{\mathbf{U}_{k,k^*}\}$ is shown in Fig. C1.

16

17 Matrix $\mathbf{U}_{k,k}$ is sufficiently well conditioned for the two-fraction aerosol model, and solution (C1)

18 is applicable for the calculation of aerosol mode concentrations. In the case of the three-fractional

19 aerosol model, parameters $Cond\{\mathbf{U}_{k,k^*}\}$ increase approximately by 10, and the matrix $\mathbf{U}_{k,k^*}$

20 becomes ill- conditioned. In such case we have to involve the Eq. (18) in retrieving procedure, i.e.

21 to use information on parameter $\hat{\mathbf{C}}^{*V}$ from radiometric measurements. With our definitions, matrix

22 $\mathbf{H}$ in Eq. (18) is written as

$$\mathbf{H}_{K \times IK} = \begin{bmatrix} \Delta h_1 & ... & 0 & \Delta h_2 & ... & 0 & ... & \Delta h_I & ... & 0 \\ 0 & \Delta h_1 & 0 & 0 & ... & 0 & ... & 0 & ... & 0 \\ ... & ... & ... & ... & ... & ... & ... & ... & ... & ... \\ 0 & ... & \Delta h_1 & 0 & ... & \Delta h_2 & ... & 0 & ... & \Delta h_I \end{bmatrix} \qquad \text{(C6)}$$

Finally, *a priori* smoothness restrictions are used as the additional factor for regularizing "ill-posed" problem solution.

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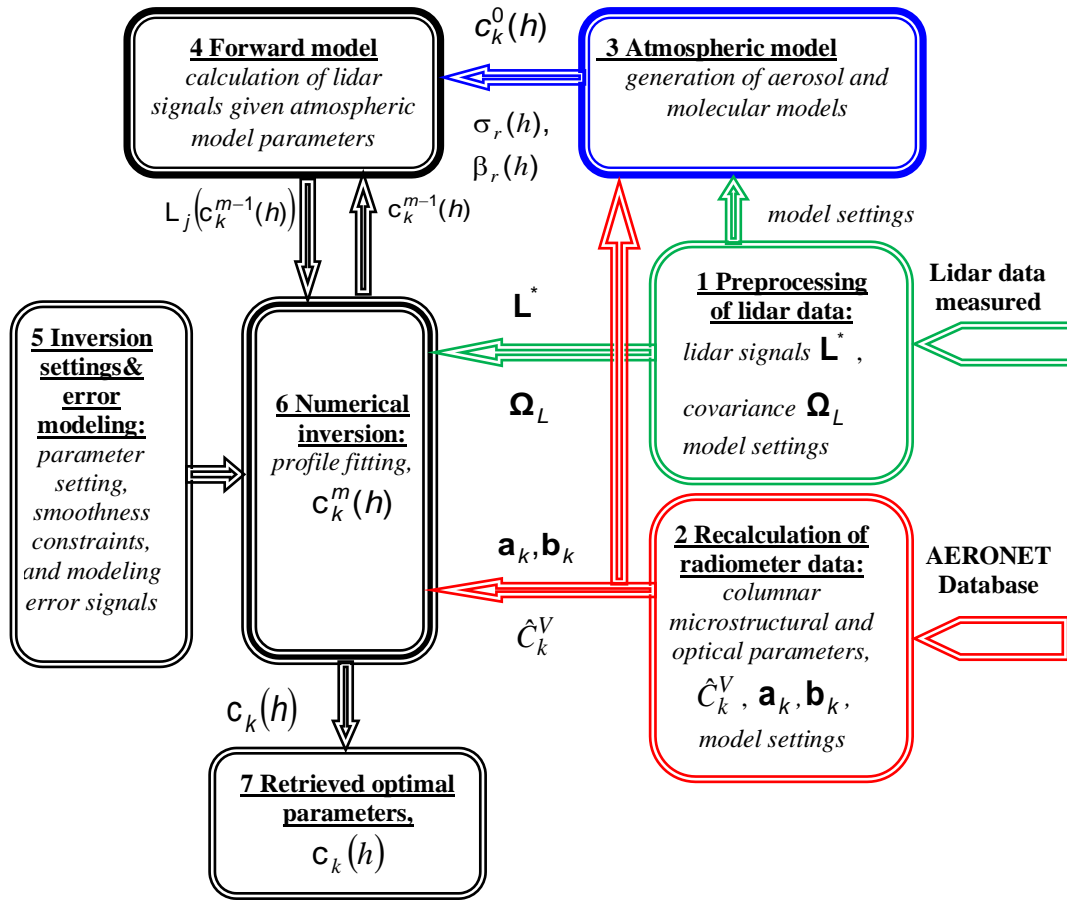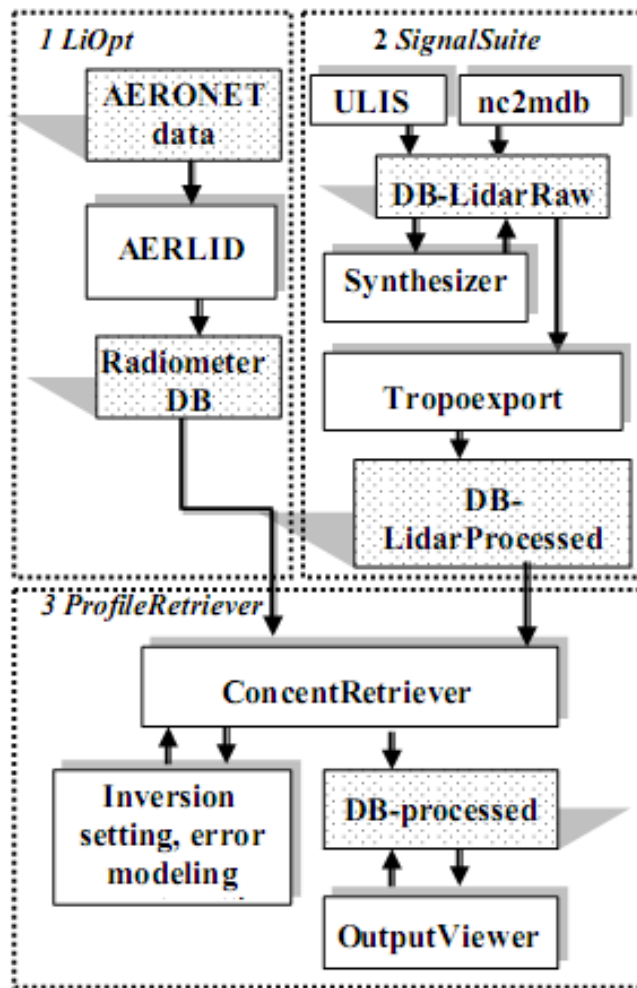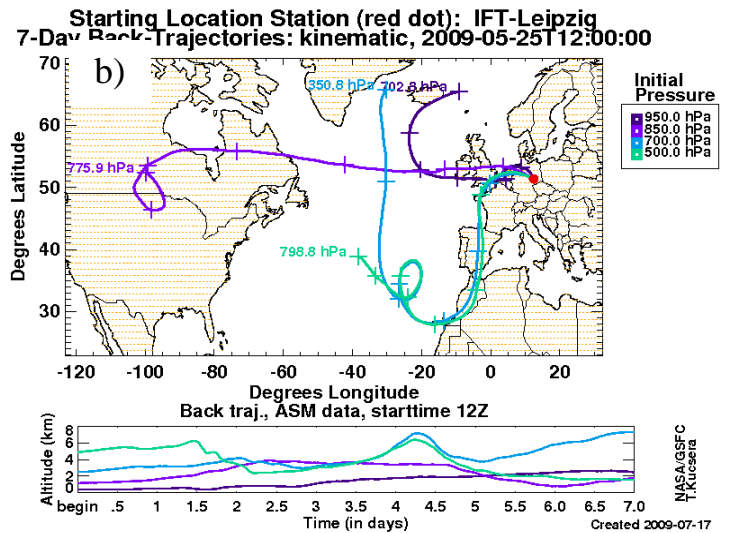Figure 1. Flowchart of LIRIC algorithm. Details are in section 2.2.

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Figure 2. Flowchart of the program package.
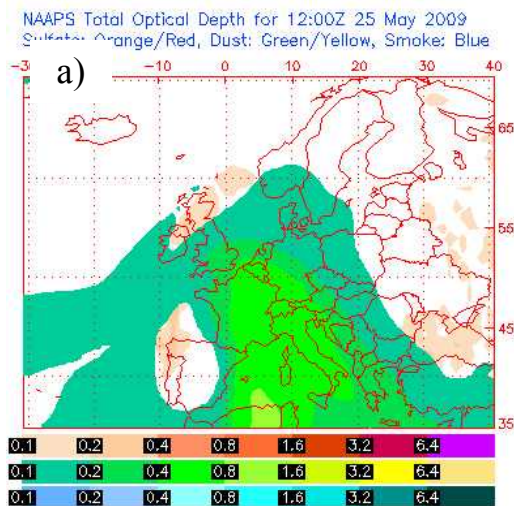
Figure 3. EARLI09 intercomparison experiment: (a) NAAPS Total Optical Depth forecast, 25 May, 2009 at 12:00 UTC; (b) 7-day backtrajectories ending over Leipzig, Germany at 12:00 UTC on 25 May, 2009.
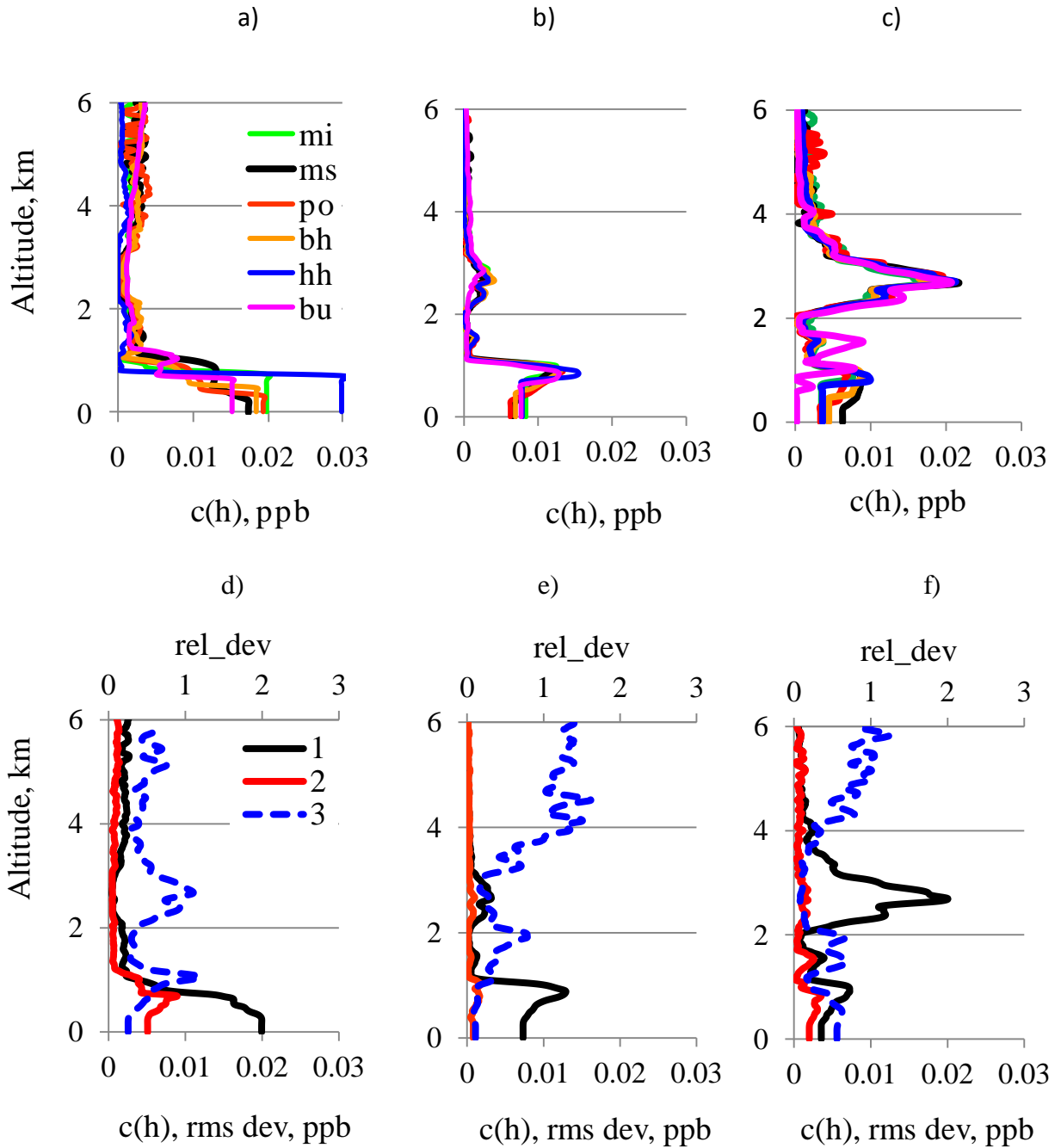
Figure 4. Particle Volume Concentrations (PVC) profiles, $c_k(h)$, and estimated deviations retrieved from data of EARLI09 intercomparison campaign, 10:20-11:40 UTC, 25 May, 2009, Leipzig, Germany, measured in Leipzig by six EARLINET lidars: mi – Minsk, ms – München, po – Potenza, bh – Bilthoven, hh – Hamburg, bu – Bucharest; (a,d) – fine, (b, e) – coarse spherical; (c, f) – coarse non-spherical; 1 – average PVC profile, 2 – rms-deviation (rms_dev), 3 – relative deviation (rel_dev). Measured data from four lidar channels (355, 532-parallel, 532 –cross, 1064 nm) and three-mode aerosol model were used.

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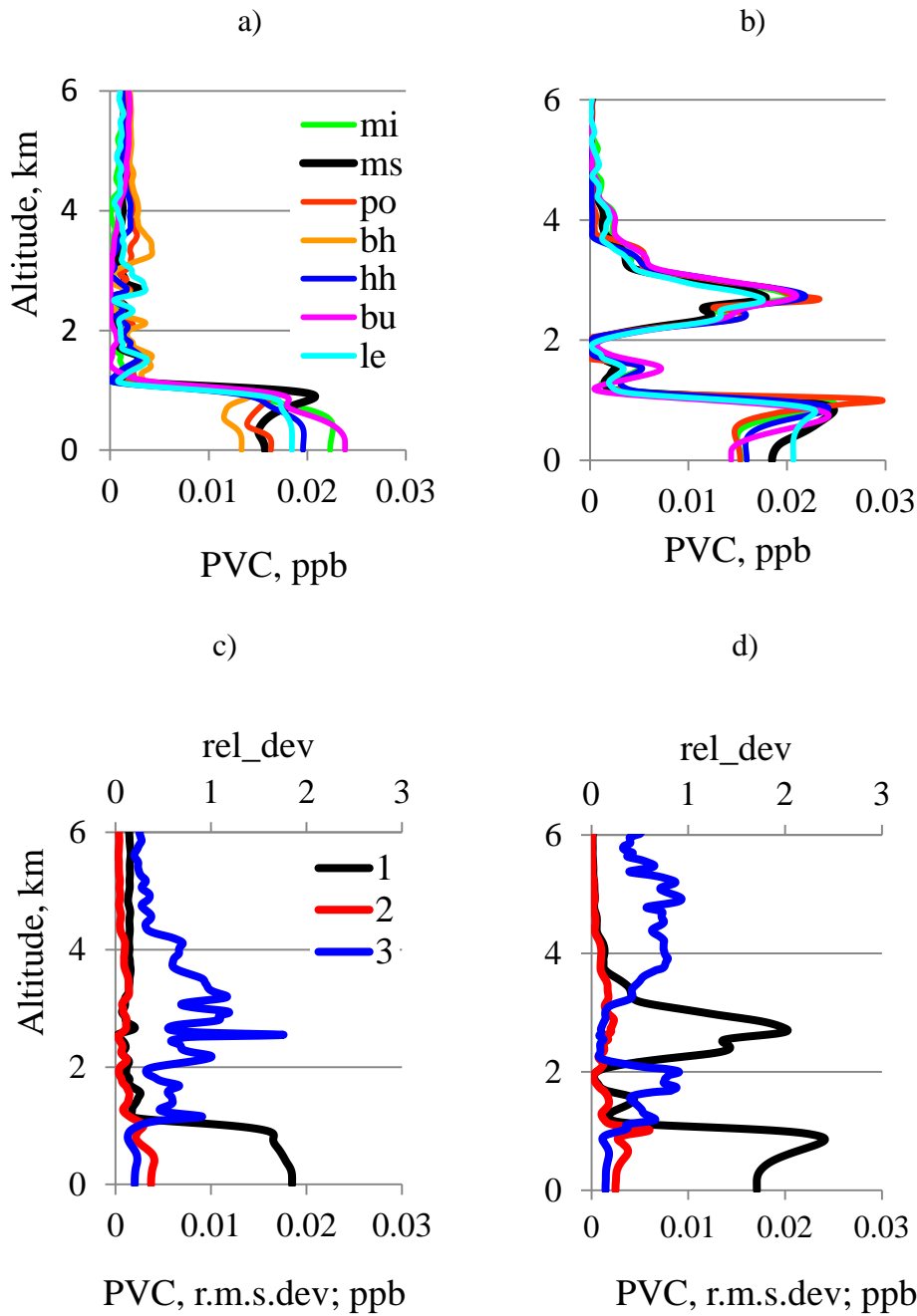
5   Figure 5. Same as Fig. 4 except for data from three lidar channels (355, 532 – intensity/parallel

6   polarized component, and 1064 nm) and two mode aerosol model were used. Label "le" stands for

7   lidar "PollyXT" of TROPOS, Leipzig. **(a, c)** – fine, **(b, d)** – coarse spherical aerosol mode.

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a)

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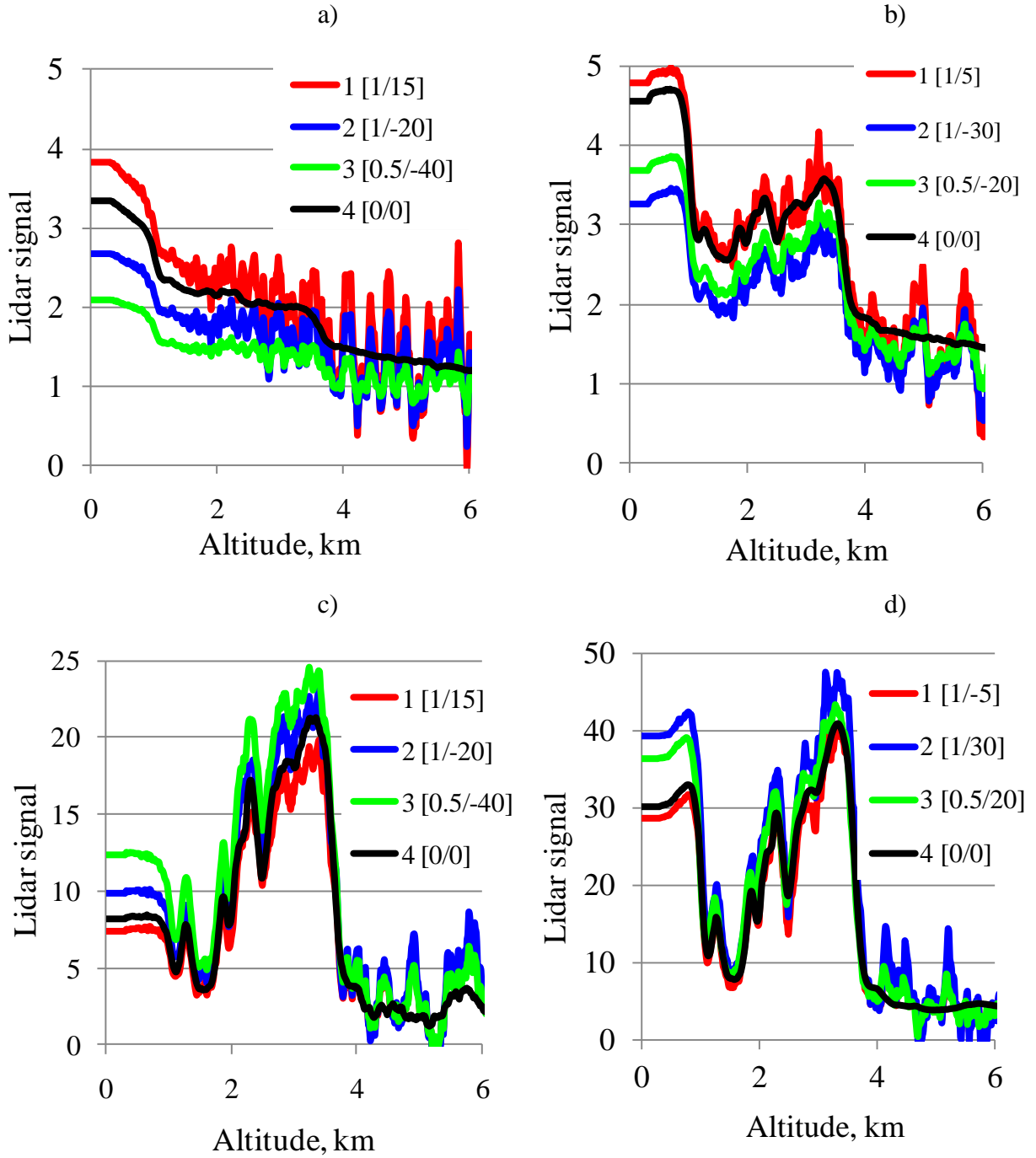
5 Figure 6. Range-corrected normalized lidar signals, $\mathbf{L}^*$, corrupted with noise and amplitude

6 distortions. Original data are provided by the München lidar team in the frame of EARLI09

7 intercomparison campaign, 14:30-15:30 UTC, 25 May, 2009, Leipzig, Germany: (a) – 355 nm, (b)

8 1064 nm, (c) – 532 nm, parallel polarized, (d) – 532 nm, cross polarized; 4 – original signal, 1 ÷ 3 –

9 corrupted signals. In square brackets distortion parameters α_j/Δ_j are given.

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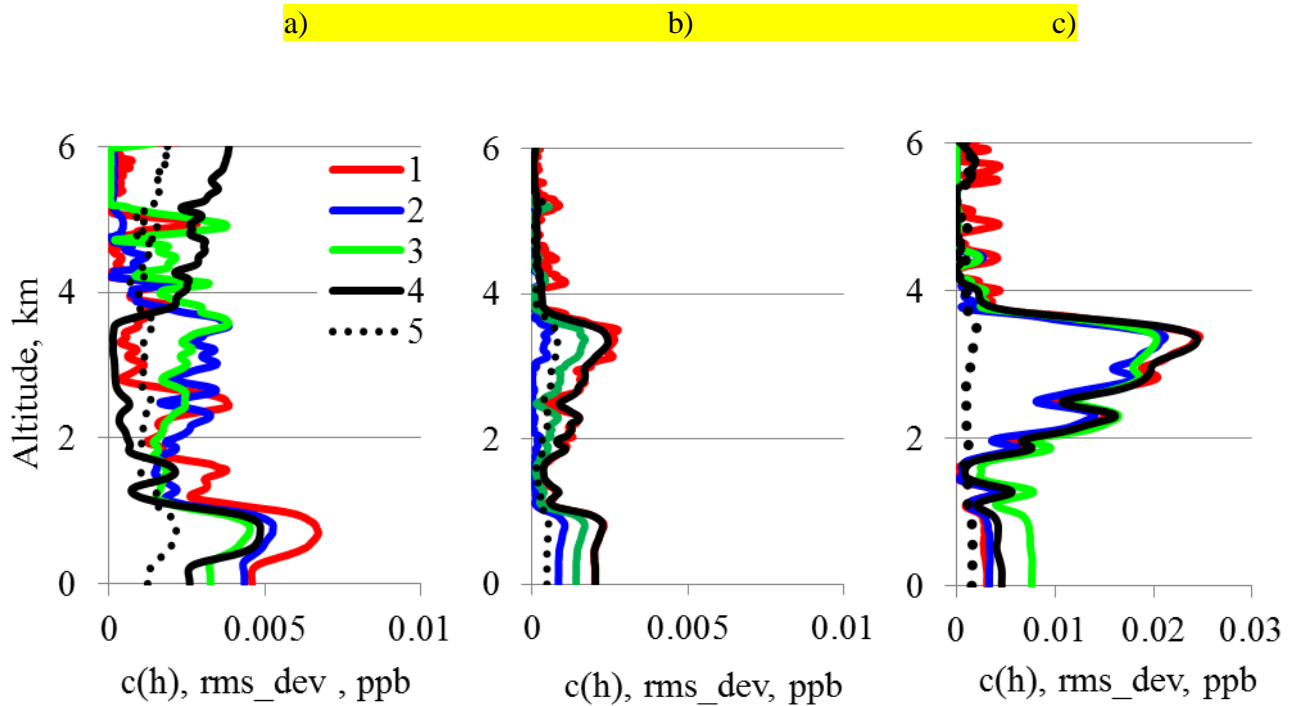
Figure 7. PVC profiles, $c_k(h)$, and their rms-deviations retrieved in response to disturbed data from of the München lidar, EARLI09 intercomparison campaign, 14:30-15:30 UTC, 25 May, 2009, Leipzig, Germany: (a) – fine, (b) – coarse spherical, (c) – coarse non-spherical modes; 4 – for the original signal, 1 - 3 – for disturbed signals; 5 – rms-deviation.
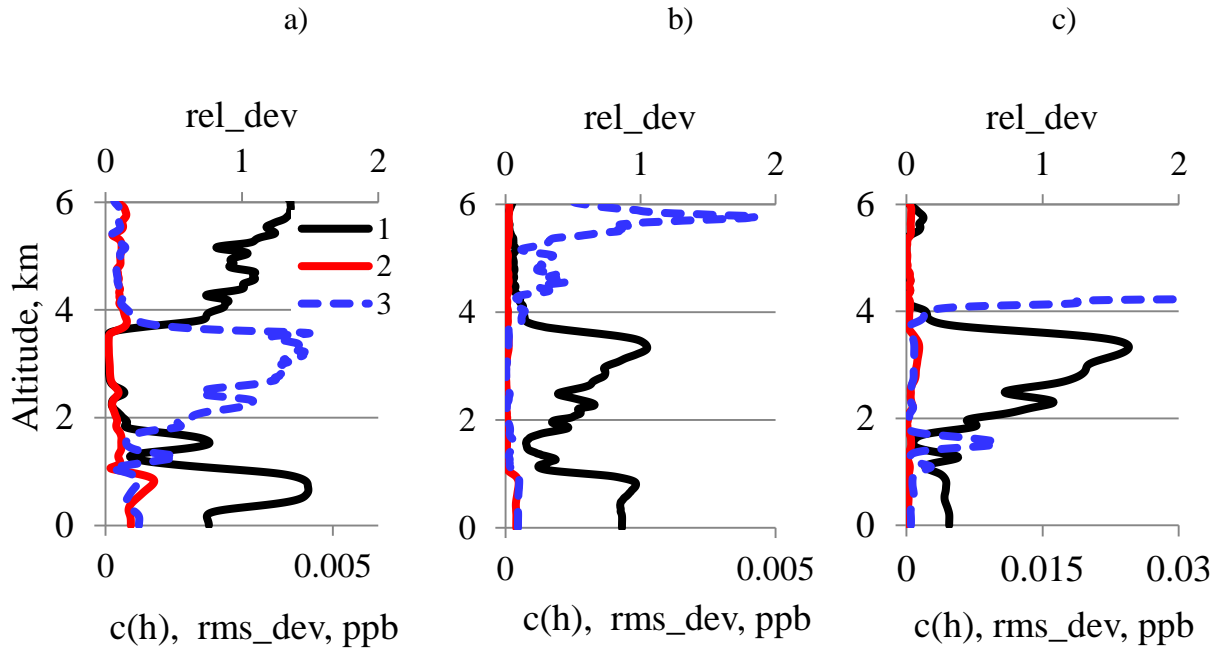
Figure 8. Variations of PVC profiles, $c_k(h)$, retrieved with 20% uncertainties in the aerosol lidar ratios; data of München lidar, EARLI09 intercomparison campaign, 14:30-15:30 UTC, 25 May, 2009, Leipzig, Germany are used; (a) fine, (b) coarse spherical, (c) coarse non-spherical modes; 1 - average value, 2 rms-deviation, 3 relative deviation.

NOAA HYSPLIT MODEL
Backward trajectories ending at 0900 UTC 19 May 10
GDAS Meteorological Data

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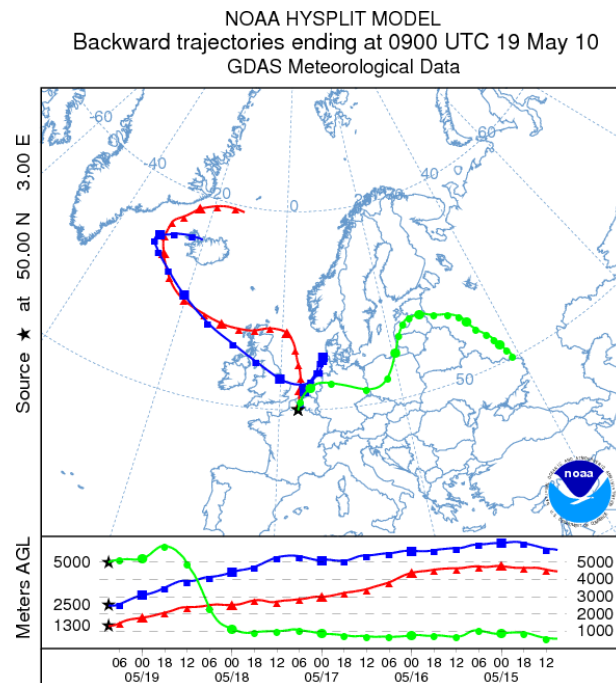
Figure 9. Air-mass back-trajectories for Lille at 08:00 UTC, 19 May 2010, (NOAA HYSPLIT
model).

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a)

b)
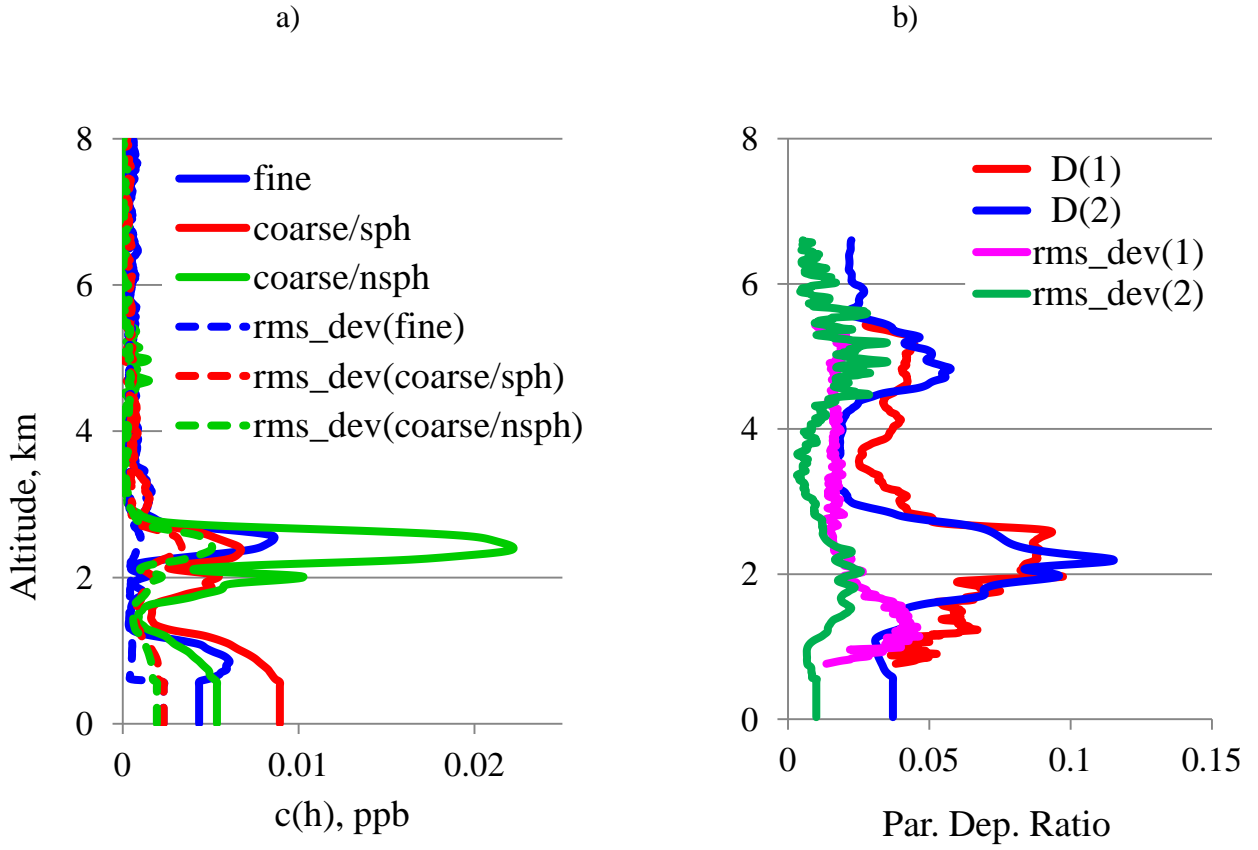
Figure 10. (a), PVC profiles, $c_k(h)$, of the fine, course-spherical (coarse/sph) and coarse-nonspherical (coarse/nsph) aerosol modes, and their rms-deviations (rms_dev(fine), rms_dev(coarse/sph), and rms_dev(coarse/nsph)); (b), particle depolarization ratio, D(1) and (2), and their rms-deviations, rms_dev(1) and rms_dev(2). Profiles were retrieved from the data measured in Lille, 19 May, 2010, 09:17–09:58 UTC. Profiles D(1) and rms_dev(1) are the results of the direct calculation of depolarization ratio and their rms-deviations from lidar measurements, as well as D(2) and rms_dev(2) were calculated from retrieved aerosol mode concentrations, $c_k(h)$.
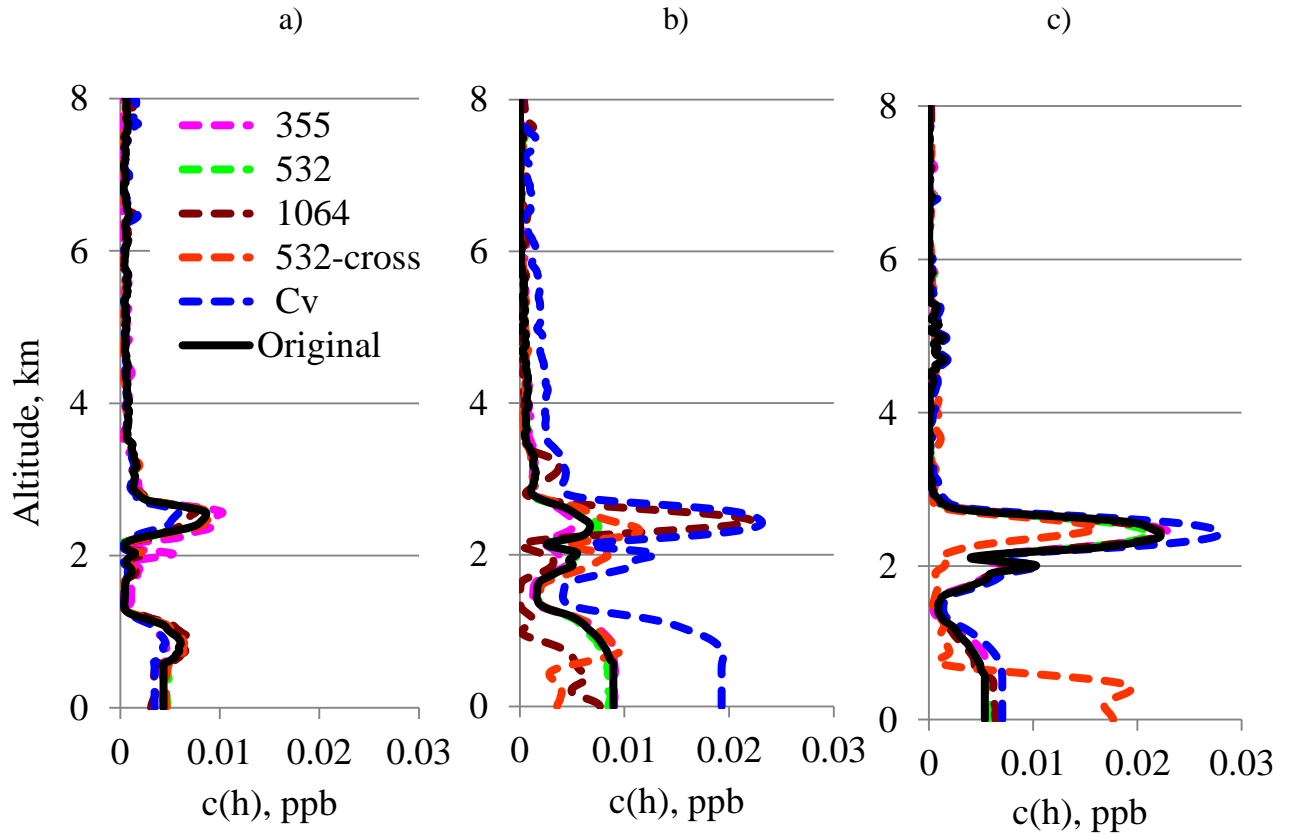
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5 Figure 11. Variation of aerosol concentration profiles, $c_k(h)$, for fine (a), coarse spherical (b) and

6 coarse non-spherical (c) aerosol modes in response to elimination of different parts of input

7 information. Tag "Original" denotes complete set of input data; tag "355" (or 532, 1064, 532-cross)

8 denotes that lidar signal at 355 nm (or 532, 1064, 532-cross) wavelength is excluded; tag

9 "C^V"denotes that columnar volume concentrations of aerosol modes are excluded. Lille, 08:00
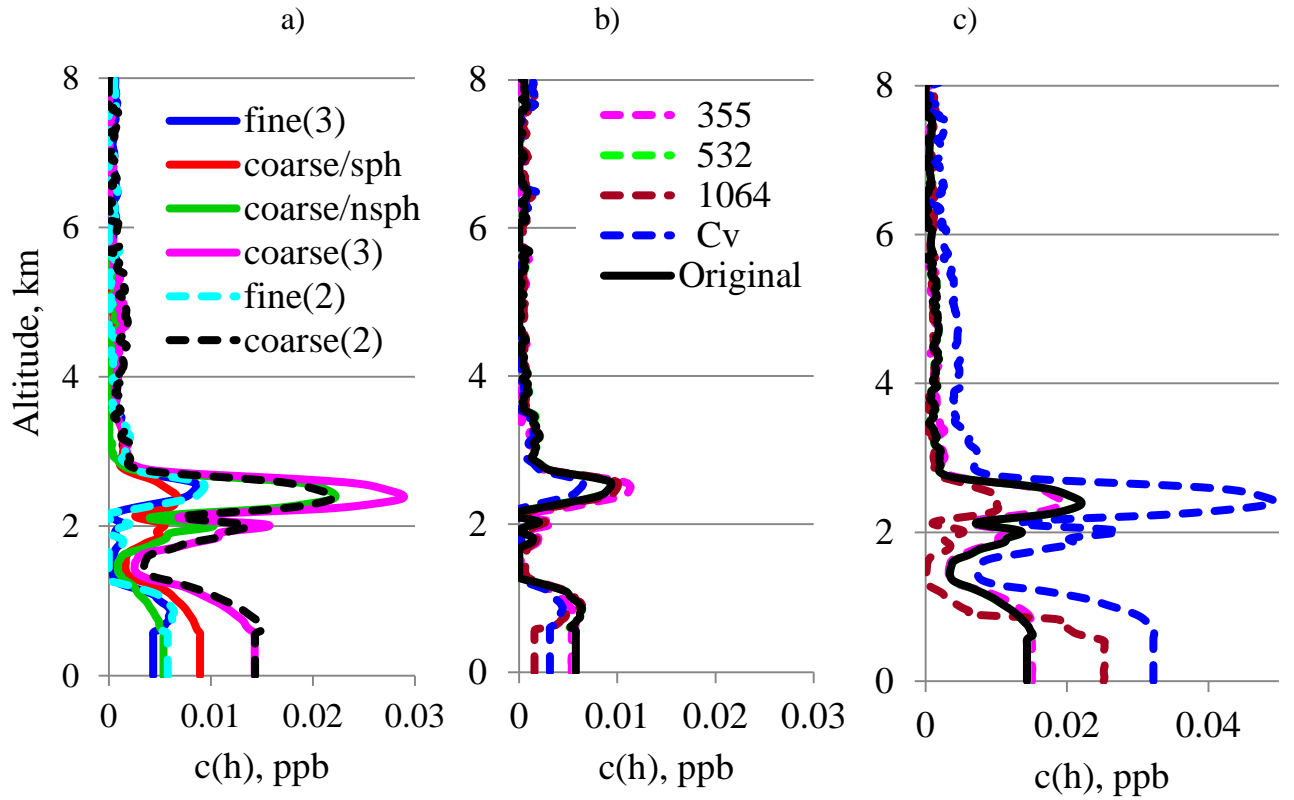
10 UTC, 19 May, 2010.

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Figure 12. Comparison of PVC profiles, $c_k(h)$, for the two- and three-mode aerosol models (a), and variations of concentration profiles, $c_k(h)$, for fine (b) and coarse (c) aerosol modes of the two-mode aerosol model in response to elimination of different parts of input information. In Fig. 12a tags "fine(2)" and "coarse(2)" denote fine and coarse modes of two-mode aerosol model. Tags "fine(3)", "coarse/sph", "coarse/nsph" and "coarse(3)" denote fine, coarse spherical, course non-spherical and total course mode of three-mode aerosol model, correspondently. In Fig 12b and 12c tag "Original" means complete set of input data; tag "355" (or 532, 1064) denotes that the lidar signal at 355 nm (or 532, 1064) wavelength is excluded; tag "C^V" denotes that columnar volume concentrations of aerosol modes are excluded. Lille, 08:00 UTC, 19 May 2010.
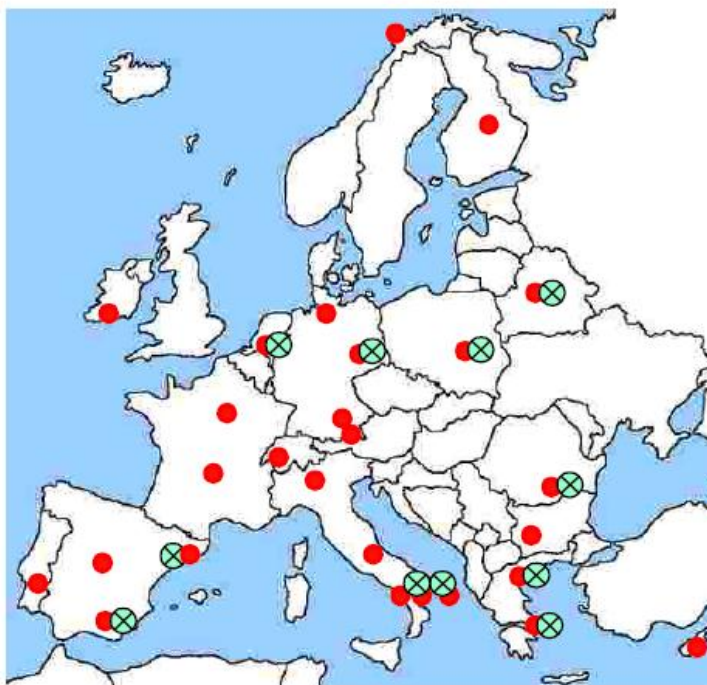
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2   Figure 13. Map of the EARLINET stations (red dots). Green dots indicate the stations where LIRIC

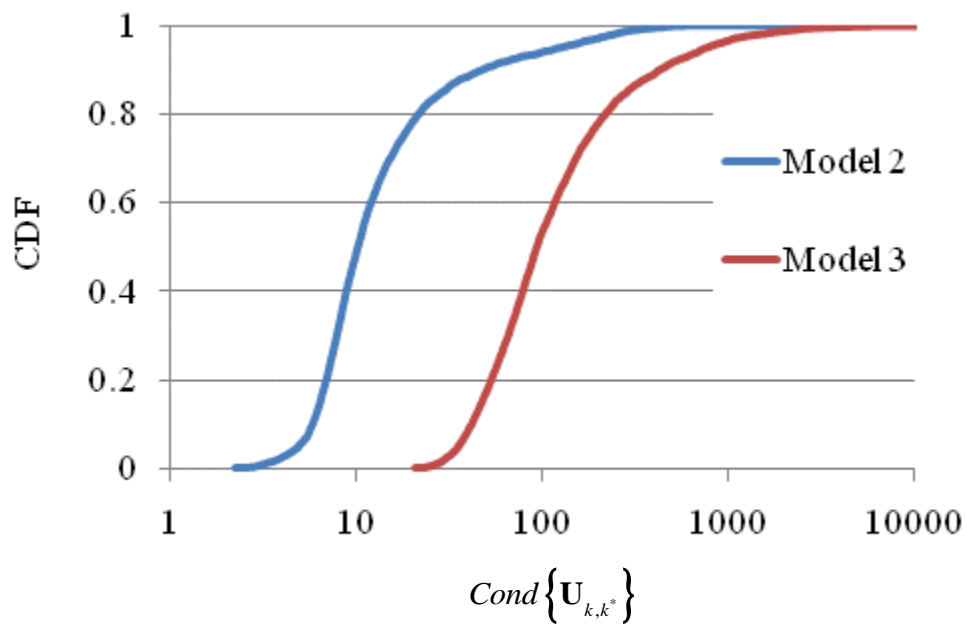3   program package has been implemented.

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Figure C1. Cumulative Distribution Functions (CDF) of parameter $Cond\left\{\mathbf{U}_{k,k^*}\right\}$ (condition number) calculated from radiometer data of the AERONET station in Minsk for two- and three-fraction aerosol models, Model 2 and Model 3, respectively.