



1 **Real time retrieval of volcanic cloud particles and SO₂ by**
2 **satellite using an improved simplified approach**

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9
10 **Abstract**

11 Volcanic Plume Removal (VPR) is a procedure developed to retrieve the ash optical depth, effective
12 radius and mass, and sulphur dioxide mass contained in a tropospheric volcanic cloud from the
13 thermal radiance at 8.7, 11, and 12 μm . It is based on an estimation of a virtual image representing
14 what the sensor would have seen in a multispectral thermal image if the volcanic cloud were not
15 present. Ash and sulphur dioxide were retrieved by the first version of the VPR using a very simple
16 atmospheric model that ignored the layer above the volcanic cloud. This new version takes into
17 account the layer of atmosphere above the cloud as well as thermal radiance scattering along the line
18 of sight of the sensor. In addition to improved results, the new version also offers easier and faster
19 preliminary preparation and includes other types of volcanic particles. As in the previous version, a
20 set of parameters regarding the volcanic area, particle types, and sensor are required to run the
21 procedure. However, in the new version, only the mean plume temperature is required as input data.
22 In this work a set of parameters have been computed for different types of plume particles (andesite,
23 obsidian, pumice, ice, water, and sulphuric acid droplets), for both the Mt. Etna (Italy) and
24 Eyjafjallajökull (Iceland) volcanoes, and for the MODIS Terra and Aqua instruments. Two different
25 synthetic images, one for Mt. Etna and one for Eyjafjallajökull, are used to compare the results from
26 the new and old procedures. Finally, a sensitivity analysis was conducted to investigate variations in
27 VPR ash and sulphur dioxide retrievals as a function of plume altitude and particle type.

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1 Introduction

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3 The large volumes of ash mixed with various gases that can be released into the atmosphere during
4 explosive volcanic eruptions sometimes form clouds that travel great distances from the source over
5 long periods, carried by the wind. These ash clouds can be generated at any time from the eruption
6 of any one of more than 1,200 active volcanoes scattered over the Earth's surface (Prata, 2009) and
7 pose a real threat to air safety (Casadevall, 1994).

8 An effective global monitoring system today depends on the use of satellite data to detect and monitor
9 the evolution of volcanic ash clouds. Timely information on the location, size, height, and ash content
10 of potentially hazardous eruption clouds derived from satellite data are generated and used by the
11 Volcanic Ash Advisory Centres (VAACs) to mitigate this type of threat and improve aviation safety
12 (Francis et al., 2012).

13 Satellite sensors operating in the thermal infrared range are particularly effective for this purpose,
14 when the interaction of volcanic ash with electromagnetic radiation makes it possible to detect and
15 monitor volcanic clouds even at night. The algorithms developed exploit in various ways the reverse
16 absorption of the brightness temperature observable in the channels centred at 10.8 and 12 microns.
17 This feature is used both for discriminating ash and meteorological clouds (Prata 1989a, 1989b), and
18 for quantifying the mass, optical thickness, and effective radius of the ash contained in volcanic
19 clouds (Wen and Rose, 1994).

20 Several algorithms were developed in the early efforts to detect volcanic clouds and retrieve
21 the ash and SO₂ contents, as discussed in a recent critical review (Clarisse and Prata, 2016). Among
22 the new algorithms the simplified approach of the VPR is distinguished by its ease of use and speed
23 of calculation, making it highly effective for monitoring. Another advantage of the VPR approach is
24 that it only requires the plume temperature as additional input, providing fresh estimates of ash and
25 SO₂ as soon as new satellite images of an ongoing eruption become available (Pugnaghi et al., 2013;
26 Guerrieri et al., 2015).

27 The VPR procedure was developed using thermal infrared (TIR) data collected by the Moderate
28 Resolution Imaging Spectroradiometer (MODIS) instrument on board the Terra and Aqua polar
29 platforms, and by the Spinning Enhanced Visible and Infra Red Imager radiometer (SEVIRI) on board
30 meteorological satellites positioned on MSG geostationary orbits.

31 This paper aims to present the VPR procedure in an improved and simplified form as developed for
32 the selected case studies of the Mt. Etna (Italy) and Eyjafjallajökull (Iceland) eruptions. Section 2 is
33 dedicated to a theoretical description of the novel improvements of the VPR procedure, while section
34 3 presents and discusses the results obtained for the validation case studies. Section 4 provides



1 conclusions. Further theoretical details are included in Appendix A, while the VPR coefficients are
2 tabulated in Supplement for different types of plume particles (andesite, obsidian, pumice, ice, water,
3 and sulphuric acid droplets), for both the Mt. Etna and Eyjafjallajökull volcanoes, and for the MODIS
4 Terra and Aqua instruments.

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7 **2. Theory**

8 The Volcanic Plume Removal (VPR) procedure (Pugnaghi et al. 2013; Guerrieri et al. 2015) is a
9 linearization of the radiative transfer equation developed to retrieve, from multispectral satellite
10 images at 8.7, 11, and 12 μm , the ash optical depth at 550 nm (δ^*), effective radius (R_e), mass (M_a),
11 and sulphur dioxide mass (M_s) of a tropospheric volcanic cloud. The parameters required to apply
12 the VPR are specific for a given volcano, type of plume particles, and sensor on board the satellite
13 and these are easily determined *a priori* using the MODTRAN radiative transfer model. Once they
14 have been computed, the only additional inputs required are the multispectral image and the mean
15 plume temperature.

16 Figure 1 shows the VPR procedure flowchart (dashed rectangle). The land-sea mask is usually
17 available with the radiance data while the operator has to define the plume mask and possibly the
18 meteorological cloud mask. For the multispectral sensors the plume mask can be derived from ash
19 detection techniques based on Brightness Temperature Difference (BTD, see Prata, 1989b) and
20 successive improvements (see Millington et al., 2012, Pavolonis et al., 2013), principle components
21 analysis (Hillger and Clark 2002a,b), or neural networks (Picchiani et al., 2014). The volcanic cloud
22 temperature input data can be obtained from VIS/TIR ground-based cameras (Scollo et al., 2014),
23 ground radar (Montopoli et al., 2014; Marzano et al., 2006; Corradini et al., 2015), lidar system
24 (Scollo et al., 2012) measurements, or from multispectral satellite data using different techniques like
25 *dark pixels* (Prata and Grant, 2001; Corradini et al., 2010), CO₂ slicing (Menzel et al., 1983; Platnick
26 et al., 2003), H₂O intercept method (Nieman et al., 1993), tracking of volcanic cloud centre of mass
27 (Guerrieri et al., 2015), or inversion schemes based on Optimal Estimation (Francis et al., 2012).

28 The first step of the VPR is definition of the virtual image with the removed volcanic cloud and
29 computation of plume transmittances for the three bands considered (8.7, 11, and 12 μm). In the
30 earlier VPR approach, the atmosphere above the plume was assumed to be negligible and the results
31 were adjusted with a cubic relationship, derived by fitting an adequate set of MODTRAN simulations
32 (Pugnaghi et al. 2013; Guerrieri et al. 2015). The transmittance values at 11 and 12 μm were used to
33 define maps of R_e , δ^* and M_a , while the sulphur dioxide abundance map was estimated from the



1 transmittance at 8.7 μm . Finally, the wind speed at the plume altitude was used to reconstruct the flux
2 at the vents, considering both the ash mass and SO_2 maps (Merucci et al., 2013; Guerrieri et al., 2015;
3 Merucci, 2015).

4 The novel VPR procedure described here applies a new atmospheric model for estimating
5 volcanic cloud transmittance (white box, inside the dashed square in Fig. 1). Here both the
6 transmittance τ'' and the up-welling radiance L''_{uo} of the layer of atmosphere above the plume are
7 considered (as shown in the scheme in Fig. 2). The term representing the surface thermal radiance
8 scattered by the volcanic particles along the line of sight of the sensor is now also considered (not
9 shown in the scheme of Fig. 2).

10 With this atmospheric model, the plume radiance L_p measured by the sensor can be approximated by
11 the parabolic trend (see Appendix A for a detailed description):

12

$$13 \quad L_p = -\alpha \cdot \tau_p^2 + [L_o + \alpha - B_p \cdot \tau'' - L''_{uo}] \cdot \tau_p + [B_p \cdot \tau'' + L''_{uo}] \quad (1)$$

14

15 where α is a term mainly proportional to $\varepsilon \cdot B(T_s) \cdot \tau$; ε is the surface emissivity, $B(T_s)$ is the Planck
16 emission at the surface temperature T_s , and $\tau = \tau' \cdot \tau''$ is the transmittance of the whole atmosphere
17 (α also depends on the aerosol optical depth, but this effect is important mainly for very optically
18 thick pixels); L_o is the radiance at the sensor with the plume removed; B_p is the Planck emission at
19 the mean plume temperature T_p ; $\tau_p = \tau_a \cdot \tau_s$, is the plume transmittance where τ_a is the aerosol
20 transmittance, and τ_s is the part due to sulphur dioxide. From these definitions it follows that if SO_2
21 is absent then $\tau_s = 1$; and if the aerosol optical depth $\delta = 0$, then $\tau_a = 1$.

22

23

24 **2.1 Absence of sulphur dioxide**

25 If sulphur dioxide is absent or if only the thermal bands not affected by SO_2 are considered, in
26 Eq. (1) τ_p can be substituted with τ_a representing only the ash component.

27 Fig. 3a shows a series of MODTRAN simulated radiances at the sensor versus the plume
28 transmittance obtained specifically for the band at 11 μm of the MODIS-Aqua sensor, pumice (Volz,
29 1973) ash type, and a set of possible plume configurations (see Supplement for details).

30 The parameter values of the parabolic fit of the radiance L_p versus the plume transmittance τ_a shown
31 in Fig. 3a, $L_p \cong \sum_{i=0}^2 a_i (\tau_a)^i$, are reported in Table 1. By definition: $a_0 + a_1 + a_2 = L_o$. In this case
32 the sum is 8.57 ($\text{W m}^{-2} \text{sr}^{-1} \mu\text{m}^{-1}$), as can be seen in Fig. 3 for $\tau_a = 1$.



1 The parabolic trend shown in Fig. 3a changes according to the state of the atmosphere, the surface
2 characteristics and, of course, the position of the volcanic cloud, composition, and ash content.
3 Approximating the radiance L_p expressed as a function of the plume transmittance τ_a with two linear
4 trends, for high radiance values (i.e. the transparent pixels of the plume) and for low values (more
5 opaque plume pixels), if the surface characteristics do not vary excessively over time, it can be
6 observed that the linear trends always intersect close to the same transmittance value (named τ_t).
7 Figure 3b shows that $\tau_t \approx 0.3$. Clearly, the gains and offsets of these two linear trends also change
8 according to the state of the atmosphere, plume temperature, and so on. These two linear fits, are
9 characterised by four parameters. However, only the offset (named B_{up}) is required to fit the
10 transparent part because the radiance L_o is known from the plume removal part of the procedure.
11 Similarly, if the intersection point of the two linear trends is known, the offset of the opaque part
12 (named B_{dn}) is sufficient to determine the linear fit.

13 Summarising, by knowing the air temperature T_p at the mean plume altitude (thus B_p term) and
14 the radiance L_o with the plume removed, it is possible to estimate the aerosol plume transmittance τ_a
15 directly from the radiance measured by the satellite (without atmospheric correction or radiative
16 transfer models) using Eq. (2) (red line) and, if necessary, Eq. 3 (blue line) of Fig. 3b:

17

$$18 \quad L_p = [L_o - B_{up}] \cdot \tau_a + B_{up} \quad (2)$$

19

20 If the computed transmittance τ_a is lower than τ_t (intersection point), then the plume transmittance
21 is recomputed by:

22

$$23 \quad L_p = [(L_t - B_{dn})/\tau_t] \cdot \tau_a + B_{dn} \quad (3)$$

24

25 where L_t is the radiance at the sensor computed using Eq. (2) for a plume transmittance $\tau_a = \tau_t$.

26 Figure 4a shows that in the 11 μm band there is a linear relationships between the two aforementioned
27 offsets B_{up} , B_{dn} and the Planck emission of the plume B_p . A similar relationship also exists for the
28 other two bands (obviously, for the band centred at 8.7 μm , sulphur dioxide must be absent) and for
29 other volcanic particle types (see Supplement). Figure 4a also shows that the plume transmittance at
30 the intersection point τ_t is almost constant with only a small dependence on B_p .

31 Therefore:

32

$$33 \quad B_{up} = a_{up} \cdot B_p + b_{up} \quad (4)$$



1
2 $\tau_t = a_{tt} \cdot B_p + b_{tt}$ (5)

3
4 $B_{dn} = a_{dn} \cdot B_p + b_{dn}$ (6)

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7 **2.2 Presence of sulphur dioxide**

8 The presence of sulphur dioxide complicates transmittance retrieval at 8.7 μm because weak SO_2
9 absorption affects this band. If the aerosol component of the plume transmittance at 8.7 μm is known,
10 then the radiance at the sensor without the presence of sulphur dioxide L_a can be computed using the
11 Eqs. (2) and (3). A knowledge of radiance due only to aerosols makes it possible to define the
12 following simple equation:

13
14 $L_p = [L_a - B_s] \cdot \tau_s + B_s$ (7)

15
16 where L_p is the total plume radiance measured by the sensor, L_a is radiance due to the aerosol
17 components of the plume, and τ_s is the plume transmittance due to SO_2 absorption. B_s is a constant
18 which takes into account the plume temperature, plume position, and state of the atmosphere above
19 the plume and it is computed using a linear function of B_p :

20

21 $B_s = a_s \cdot B_p + b_s$ (8)

22

23 Figure 4b shows the trend of B_s versus B_p derived from a complete dataset of MODTRAN
24 simulations for Mt. Etna measured with a MODIS-Aqua instrument.

25 Therefore, to compute τ_s from Eq. (7), it is necessary to know L_a which is derived from Eqs. (2) and
26 (3) when τ_a for the band at 8.7 μm is known. The transmittance τ_a can easily be computed for
27 pumice-type ash particles because a very good correlation exists between $\tau_{a,8.7}$ and $\tau_{a,11}$ (see Fig. 5a,
28 and Pugnaghi et al. 2013). The fit is a cubic polynomial: $\tau_{a,8.7} = \sum_{i=0}^3 a_{i,8.7} (\tau_{a,11})^i$ and the
29 parameter values from the MODIS sensors on board the Terra and Aqua satellites are reported in
30 Table 2. These parameters are an improved version of those reported in Pugnaghi et al. (2013),
31 because in the first version of the VPR the thermal radiance scattered along the line of sight of the
32 sensor was ignored.



1 Unfortunately, for other particle types (see Supplement), the correlation between $\tau_{a,8.7}$ and $\tau_{a,11}$ is
2 not always good, as in the example of Fig. 5b showing the scatter plot for water droplets.
3 Nevertheless, it should be noted that this correlation becomes very good if only particles of the same
4 effective radius R_e are considered.

5 In these cases with non-pumice ash types, the aerosol transmittance $\tau_{a,8.7}$ at 8.7 μm can be obtained
6 from the formula:

7

$$8 \quad \tau_{a,8.7} = e^{-\mu \cdot \delta_{8.7}} = e^{-\mu \cdot m_{8.7} \cdot \delta^*} \quad (9)$$

9

10 where μ is the optical air mass factor, $\delta_{8.7}$ is the vertical optical depth, δ^* is the vertical optical depth
11 at 550 nm, and $m_{8.7}$ is the gain of the linear relationship which gives the optical depth $\delta_{8.7}$, when δ^*
12 is known; the gain $m_{8.7}$ is a function of the effective radius R_e and is known from the MODTRAN
13 simulations (Guerrieri et al., 2015).

14 To sum up, the novel VPR procedure first computes the 11 and 12 μm band transmittances (as
15 indicated in the flowchart of Fig. 1), and from these the aerosol optical depth at 550 nm (δ^*) and the
16 effective radius (R_e) of each pixel of the plume (Pugnaghi et al. 2013); then the aerosol transmittance
17 at 8.7 μm ($\tau_{a,8.7}$) is obtained using Eq. (9).

18 Finally, the transmittance $\tau_{s,8.7}$ (derived from Eq. 7) is used to estimate the SO_2 columnar abundance
19 C_s , given the proper absorption coefficient β (Pugnaghi et al. 2013) and the optical air mass μ factor.

20

$$21 \quad \tau_{s,8.7} = e^{-\mu \cdot \beta \cdot C_s} \quad (10)$$

22

23 The subsequent steps of the VPR procedure have not been changed and can be found in Pugnaghi
24 et al., (2013). Nevertheless, to conclude the theoretical discussion, it is important to note the
25 superposition effect of ash and sulphur dioxide on the radiance measured by the sensor. This means
26 that the proposed VPR procedure can also work well in cases of a *double-plume* at different
27 temperatures, for example if an ash plume is located directly above or below a sulphur dioxide plume.

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30 **3. Validation test cases**

31 To test the improved version of the VPR procedure two trial synthetic images were defined as
32 described in Corradini et al. (2014), for both the MODIS-Aqua effective wavelengths, depicting a
33 uniform ocean surface under a cloudless sky, and a plume of known spherical particles.



1 The first image is characterized by an atmospheric situation, ocean temperature, and the particle
2 type typical of the Sicilian Mt. Etna volcano, while the second is adapted to match the Icelandic
3 Eyjafjallajökull volcano. Figure 6 shows the two RGB colour composite synthetic images with the
4 radiances of the channels centred at 8.7, 11, and 12 μm respectively. The left plate shows the Mt.
5 Etna scenario with a plume of the same shape as the volcanic cloud detected by MODIS-Aqua during
6 the 26 October 2013 eruption at 12:20 GMT. The right plate shows the Eyjafjallajökull scenario,
7 depicting a portion of the Eyjafjallajökull plume detected by MODIS-Aqua on 11 May 2010 at 14:05
8 GMT.

9
10

11 **3.1 The Mt Etna-Pumice scenario**

12 The synthetic atmosphere in the Mt. Etna image in Fig. 6a is derived from the radiosonde pressure,
13 temperature, and humidity (PTH) profiles measured by the Trapani (western tip of Sicily, Italy) WMO
14 station 16429 on 26 October 2013 at 12 GMT, while the plume mask and the vertical zenith angles
15 used to prepare the synthetic image are derived from the actual MODIS-Aqua data collected on
16 October, 26, 2013 at 12:20 GMT. The plume in the synthetic image is defined as 1 km thick, located
17 between 7 and 8 km, containing pumice ash (Volz, 1973) and SO_2 . It has a Gaussian shape moving
18 from the centre to the edge and ranging from 10 to 1 g m^{-2} columnar SO_2 abundance, and from 1.5 to
19 0.1 ash optical depth δ^* (AOD at 550 nm). Therefore, a minimal quantity of sulphur dioxide and ash
20 is always present in the plume, and the effective radii R_e of the particles have a uniform distribution,
21 on a logarithmic scale, in the range 0.8-7 μm .

22 Table 3 shows that by excluding the SO_2 total mass, all the retrieval values of the new version of the
23 VPR are closer to the true values than the old version. Both versions estimate a lower mass of ash in
24 the volcanic cloud, but this probably also implies a greater burden of SO_2 detected with the old VPR.
25 The retrieval of the total ash mass computed with the new VPR is better not only because it is closer
26 to the true value, but also because both the estimated effective radius and optical depth used in mass
27 estimation are closer to the true values.

28 Figure 7 shows the scatter plots of R_e , δ^* , M_a , and M_s versus the true values (synthetic image). All
29 the scatter plots show a widening dispersion with increasing values. Fig. 8 reports (on the left) the
30 trends of R_e and δ^* mean values retrieved with VPR using different input plume altitudes, and on the
31 right the trends of ash and SO_2 total mass. As described in Pugnaghi et al. (2013), Guerrieri et al.
32 (2015), and Merucci (2015), the effective radius and aerosol optical depth at 550 nm are derived from
33 the transmittances retrieved at 11 and 12 μm , and then the ash mass is computed in each pixel with



1 the Wen and Rose (1994) simplified formula. The trend of R_e versus volcanic cloud altitude is almost
2 flat, while the optical depth δ^* shows a clear drop with height. The best retrieval (closest to true
3 values) is at 7 km rather than the height used of 7.5 km. This is also true for SO₂ total mass.
4 Figure 9 shows the VPR retrievals for the Mt. Etna scenario giving as input the right plume
5 temperature and all the types of particles reported in Tables S1, S2, S3, S5, S6, S7 (see Supplement).
6 The upper plates show the mean effective radius (R_e , left) and the mean optical depth (δ^* , right). The
7 lower plates show the retrievals of ash (left) and SO₂ (right) total mass. Among the different types of
8 ash, andesite gives the worst effective radius and optical depth results, with respect to the true values.
9 Nevertheless, because the two retrieved variables R_e and δ^* compensate each other, all the ash types
10 considered give good estimations of the ash total mass. Conversely, for ice and water the results
11 retrieved for the total mass are much higher and divergent from true values. Finally, by varying the
12 ash type, the total ash mass exhibits a much lower variability when compared to that of SO₂.

13
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15 **3.2 The Eyjafjallajökull-Andesite scenario**

16 The second synthetic image shown in Fig. 6b was created considering the state of the atmosphere
17 derived from the PTH vertical profiles measured at Keflavik (WMO station 04018) on 11 May 2010
18 at 12:00 GMT. On that day, MODIS-Aqua captured the Eyjafjallajökull eruption during its transit at
19 14:05 GMT. As in the previous Mt. Etna eruption, the plume mask and the vertical zenith angles used
20 in the synthetic image derive from the actual MODIS-Aqua data. The plume was again 1 km thick,
21 but located lower than the previous case at between 4 and 5 km, containing spherical particles of
22 andesite (Pollack et al. 1973) and SO₂. The same ranges and distributions of SO₂ columnar content,
23 ash optical depth δ^* at 550 nm, and effective radius R_e were used, as for Mt. Etna.

24 Once again Table 4 demonstrates that the new version of the VPR generates better estimations
25 compared to the old one, with all the parameters exhibiting difference percentages lower than 10%.
26 The older version detects the presence of ash in a smaller number of pixels, but its greater effective
27 radius and optical depth partly compensate for the smaller number of detected ash plume pixels in the
28 final mass estimation.

29 The numbers of detected pixels for the sulphur dioxide component are in close agreement with the
30 true value in both versions, and they retrieve a similar total mass of SO₂.

31 Figure 10 shows the scatter plots of R_e , δ^* , M_a , and M_s versus the true values. The correlation is quite
32 good up to about $R_e = 5 \mu\text{m}$, $\delta^* = 1$, $M_a = 10 \text{ gm}^{-2}$, and $M_s = 7 \text{ gm}^{-2}$; wider dispersions can be
33 observed for higher values.



1 Fig. 11 shows the trends of the VPR R_e and δ^* mean values (left), and total M_a and M_s (right)
2 retrieved as functions of the input plume altitude. R_e and δ^* retrievals exhibit opposite trends which
3 compensates each other, making the final retrieval of total ash mass less sensitive to plume altitude.
4 In fact, from 3 to 5 km the ash mass ranges between 11 and 15 kt with a true value of about 13 kt.
5 Only the optical depth at the right plume altitude (4.5 km) is very close to the true value. The best
6 effective radius is at a mean plume altitude greater than 5 km and the best total mass is at about 4 km.
7 However, the result obtained appears to be the best compromise. A greater plume height (e.g. 5 km)
8 would mean a better R_e but a worse δ^* and also a worse total mass. Conversely, a lower plume height
9 (e.g. 4 km) yields to very good total mass, but worse R_e and δ^* values.

10 Finally, as for Mt. Etna, the VPR results were considered using as input the actual plume temperature
11 and all the types of particles reported in Supplement, including the Eyjafjallajökull ash type (Peters,
12 2013, referred in the figures as Eyja ash). In Fig. 12 the upper plates show the mean effective radius
13 R_e (left) and optical depth δ^* (right) versus the different types of particles; the lower plates show the
14 retrievals of ash and SO₂ total mass.

15 Only andesite and Eyjafjallajökull ash types give good results in both mean effective radius and
16 optical depth, while obsidian is reasonably good only for the first parameter. The ash total mass
17 exhibits quite constant values between 10 to 15 kt for all the four ash types, close to the true value of
18 13 kt. Vice versa, the total mass values retrieved for ice and water droplets were much higher and
19 very different from the true values. As noted in the previous Etna scenario, their reciprocal difference
20 is mainly due to different radius thresholds used in the procedure for water droplets and ice. The
21 performance of sulphur dioxide total mass retrieval is seen to be strongly affected by the type of
22 particle used in the procedure. Only andesite gives a good result in this case. Finally, sulphuric acid
23 retrieval performance is different from both ash and the water-ice pair.

24

25

26 **4. Conclusions**

27

28 The new version of the VPR presented here is an approximated procedure. It uses only the mean
29 altitude cloud temperature as input to directly interpret MODIS-TIR multispectral images and retrieve
30 particle effective radius, optical depth, mass of the particle utilized, and mass of sulphur dioxide
31 contained in each pixel of the volcanic cloud. The VPR approach requires no atmospheric correction
32 because this is implicit in the procedure itself. The retrieval of effective radius, optical depth, and
33 sulphur dioxide abundance is derived from estimation of plume transmittances in the bands centred
34 at 8.7, 11, and 12 μm . This article presented a novel and effective improvement in the transmittance



1 estimation scheme. In the new VPR, plume transmittance is obtained from the radiance measured by
2 the sensor using two simple linear relationships; one for the most transparent part of the plume and
3 one for the most opaque. These two linear trends account for two minor terms not considered in the
4 previous version: the layer of atmosphere above the plume and the thermal radiance scattered along
5 the line of sight of the sensor. Approximation for very thick/opaque volcanic clouds (transmittances
6 lower than 0.05) is less effective. The improvement only involves the computation of volcanic cloud
7 transmittance, and no other parts of the previous procedure have been modified. Nevertheless, the
8 improvement has a dual positive effect: 1) it is simpler to use and provides more accurate results than
9 before; 2) the preliminary work to compute the parameters required by the procedure (the parameters
10 reported in Supplement) is easier than before and requires less processing time. The new VPR
11 procedure was assessed against the older version by applying it to synthetic images generated using
12 two real examples from the Mt. Etna (Italy) and Eyjafjallajökull (Iceland) volcanoes. The percentage
13 difference between the average input data of the synthetic images and the mean results of the
14 improved VPR ranges between 2-13 % for both Mt. Etna and Eyjafjallajökull, while the old VPR
15 produced ranges between 4-68 % (see Tables 3 and 4), confirming the improved performance of the
16 new version.

17 The correlation coefficient between the transmittance of the volcanic cloud simulated by MODTRAN
18 and the corresponding transmittance retrieved by VPR is reported in the last column of tables S1 to
19 S7, in nearly all cases this being close to one. However, the mean percentage errors of the retrieved
20 effective radius, optical depth, ash mass, and sulphur dioxide mass expected in a real example may
21 be greater than those reported in Tables 3 and 4. This is because the two synthetic images considered
22 here exhibit a uniform and perfectly clear sky, a uniform ocean surface, and a volcanic cloud
23 comprised of known spherical ash particles.

24

25 **Appendix A**

26 The radiance at the sensor when the volcanic cloud is absent is (see Fig. 2):

27

$$28 \quad L_o = [\varepsilon B(T_s) + (1 - \varepsilon) L_d] \cdot \tau + L_{uo} \quad (A1)$$

29



1 where: ε is the surface emissivity; $B(T_s)$ is the Planck function at the surface temperature T_s ; L_d is
2 the atmospheric down-welling radiance; τ is the total atmospheric transmittance; L_{uo} is the total
3 atmospheric up-welling radiance.

4

$$5 \quad \tau = \tau' \cdot \tau'' \quad (A2)$$

$$6 \quad L_{uo} = L'_{uo} \cdot \tau'' + L''_{uo} \quad (A3)$$

7

8 The radiance at the sensor when the volcanic cloud is present is:

9

$$10 \quad L_p = [\varepsilon B(T_s) + (1 - \varepsilon) L_d] \cdot \tau \cdot \tau_p + L_u + S \quad (A4)$$

11

12 where: τ_p is the volcanic cloud transmittance; L_u is the current atmospheric up-welling radiance; S is
13 the term accounting for the scattering of thermal radiance along the line of sight of the sensor.

14 Since surface emissivity is close to 1 (particularly above the ocean), the change of L_d in the presence
15 of a volcanic cloud was ignored.

16 The atmospheric up-welling radiance in the presence of a volcanic cloud is:

17

$$18 \quad L_u = L'_{uo} \cdot \tau_p \cdot \tau'' + L_{up} \cdot \tau'' + L''_{uo} \quad (A5)$$

19

20 and, assuming $L_{up} = B_p \cdot (1 - \tau_p)$ where B_p is the Planck function at temperature T_p (air
21 temperature at the mean plume altitude):

22

$$23 \quad L_u = \{L_{uo} - [B_p \cdot \tau'' + L''_{uo}]\} \cdot \tau_p + [B_p \cdot \tau'' + L''_{uo}] \quad (A6)$$

24



1 Indicating with P a degree of probability of the thermal radiation being scattered along the line of
2 sight of the sensor, the scattering term was modelled as:

3

$$4 \quad S = \left\{ \int_{\tau_p}^1 [\varepsilon B(T_s) \cdot \tau' \cdot \tau_p] \cdot P \, d\tau_p' \right\} \cdot \tau'' = \alpha \cdot \tau_p (1 - \tau_p) \quad (A7)$$

5

6 where: $\alpha = \varepsilon B(T_s) \cdot \tau \cdot P$.

7 Here P is assumed to be constant even if it is a function of the ash/particle characteristics and
8 therefore of τ_p itself. Clearly $P = 0$ if $\tau_p = 1$.

9 Inserting Eqs. (A6) and (A7) in (A4):

10

$$11 \quad L_p = [\varepsilon B(T_s) + (1 - \varepsilon) L_d] \cdot \tau \cdot \tau_p + \{L_{uo} - [B_p \tau'' + L''_{uo}]\} \cdot \tau_p + [B_p \tau'' + L''_{uo}] + \alpha \cdot \tau_p (1 - \tau_p) \quad (A8)$$

12

13

14 Finally, recalling Eq. (A1), Eq. (1) is obtained.

15

16

17

18

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25



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8

9



1 Table 1: Parameter values of the parabolic fit shown in Fig. 3a.

$(-\alpha)$	$(L_o + \alpha - B_p \cdot \tau'' - L_{uo})$	$(B_p \cdot \tau'' + L_{uo})$
a_2	a_1	a_0
-1.58	6.92	3.23

2

3



1 Table 2: Polynomial coefficients to compute $\tau_{a,8.7}$ from $\tau_{a,11}$ for the pumice (Volz, 1973) ash type.

Satellite	$a_{3,8.7}$	$a_{2,8.7}$	$a_{1,8.7}$	$a_{0,8.7}$
<i>Terra</i>	0.1645	-0.4249	1.2559	0.0050
<i>Aqua</i>	0.1412	-0.3483	1.2028	0.0046

2

3



- 1 Table 3: Main characteristics of synthetic image, indicated as “True”, together with the results of the
- 2 VPR procedure, both new and old versions. The air temperature at 7.5 km was used as input for the
- 3 VPR. The percentage differences are shown in brackets.

Mt. Etna – Pumice; 23 October 2013	True	VPR	VPR
plume altitude 7-8 km		new	old
Mean R_e (μm)	2.85	2.92	4.80
(% difference)		(2.5)	(68.4)
Mean δ^*	0.25	0.22	0.19
(% difference)		(-12)	(-24)
$R_e < 2 \mu\text{m}$ Fine particles (%)	42	42	21
$2 \mu\text{m} < R_e < 5 \mu\text{m}$ Mean particles (%)	42	46	42
$R_e > 5 \mu\text{m}$ Coarse particles (%)	16	12	37
Ash mass (t)	8336	7812	7166
(% difference)		(-6.3)	(-14.0)
Pixels detected with ash	7533	7533	7317
SO ₂ mass (t)	19636	17146	18880
(% difference)		(-12.7)	(-3.9)
Pixels detected with SO ₂	7533	7533	7533

4

5



- 1 Table 4: Main characteristics of synthetic image, indicated as “True”, together with the results of the
- 2 VPR procedure, both new and old versions. The air temperature at 4.5 km was used as input for the
- 3 VPR. The percentage differences are shown in brackets.

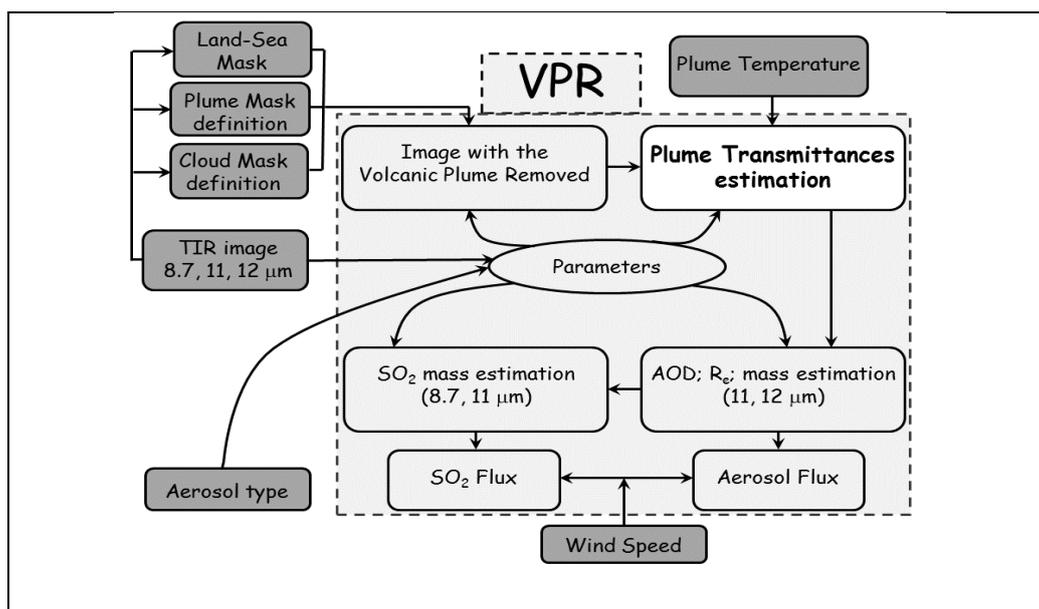
Eyjafjallajökull - Andesite	True	VPR	VPR
11 May 2010 14:05; plume altitude 4-5 km		new	old
Mean R_e (μm)	2.83	2.62	3.3
(% difference)		(-7.4)	(16.6)
Mean δ^*	0.28	0.29	0.39
(% difference)		(+3.6)	(39.3)
$R_e < 2 \mu\text{m}$ Fine particles (%)	43	46	38
$2 \mu\text{m} < R_e < 5 \mu\text{m}$ Mean particles (%)	42	44	46
$R_e > 5 \mu\text{m}$ Coarse particles (%)	15	10	16
Ash mass (t)	13227	12006	9674
(% difference)		(-9.2)	(-26.9)
Pixels detected with ash	10624	10624	6532
SO ₂ mass (t)	30724	28714	28235
(% difference)		(-6.5)	(-8.1)
Pixels detected with SO ₂	10624	10624	9827

4

5

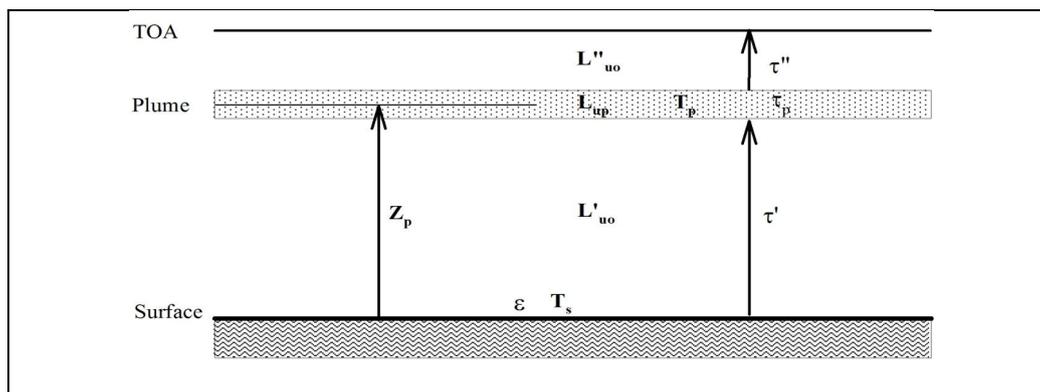


1



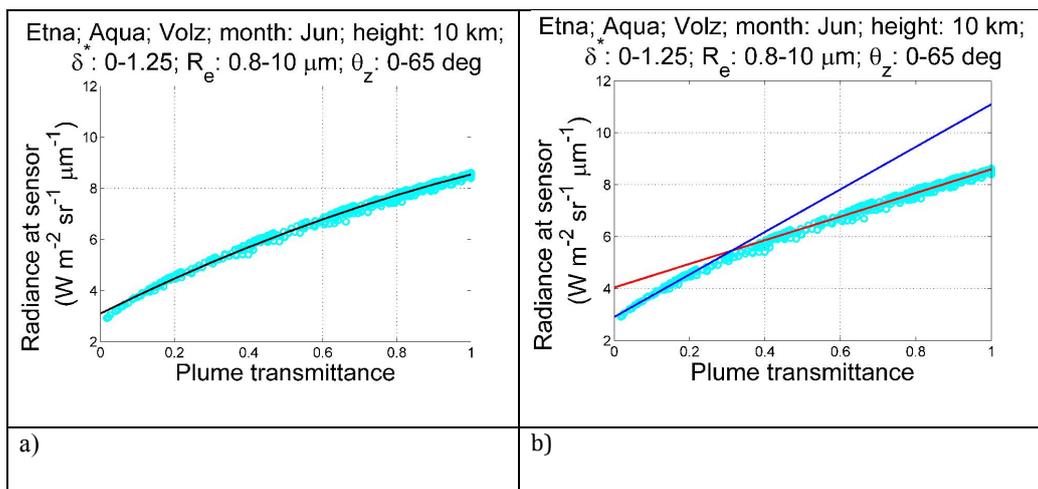
2 Figure 1. Flowchart illustrating the main steps of the VPR procedure.

3



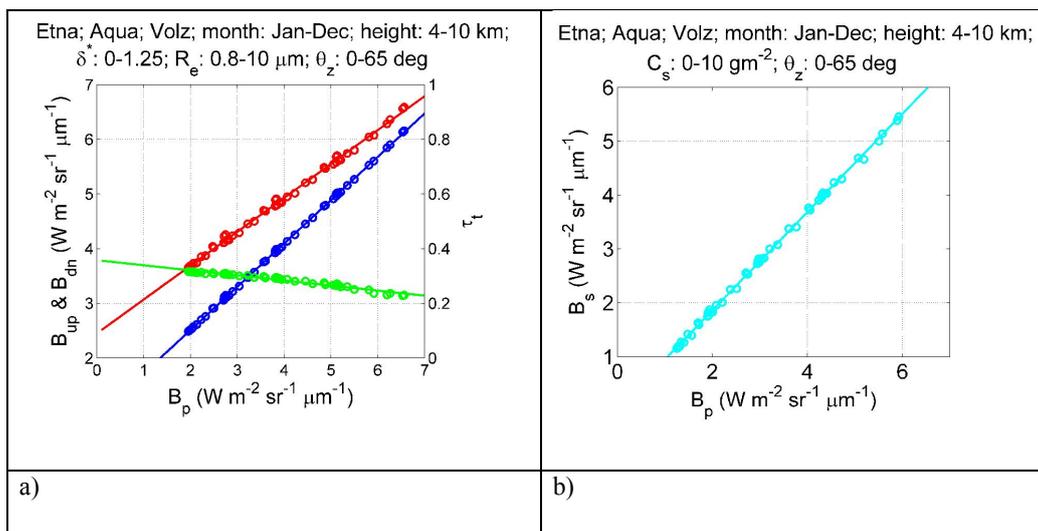
1 Figure 2. Scheme of the atmospheric model used in the improved VPR.

2

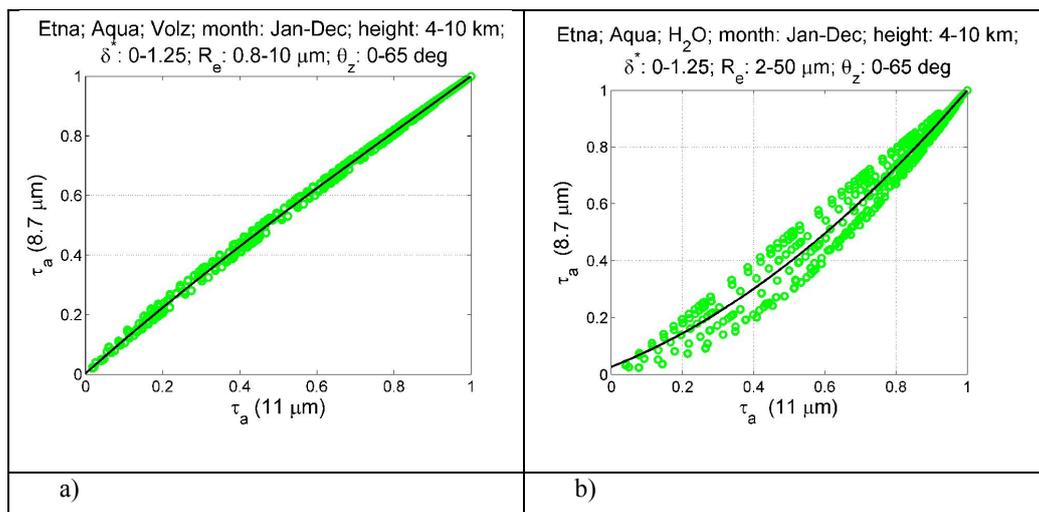


1 Figure 3. a) Radiances at the sensor (11 μm) vs. plume transmittances and their parabolic fit (black
2 line); b) Same radiances with the two linear fits of Eq. (2) for the more transparent part of the plume
3 (upper fit, red line), and Eq. (3) for the most opaque part of the plume (lower fit, blue line).

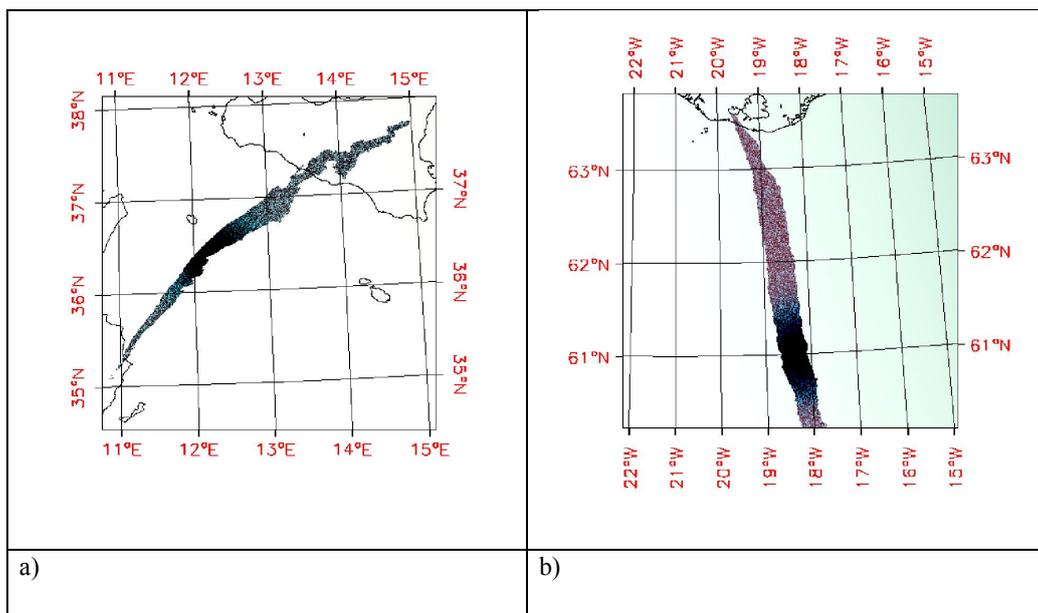
4



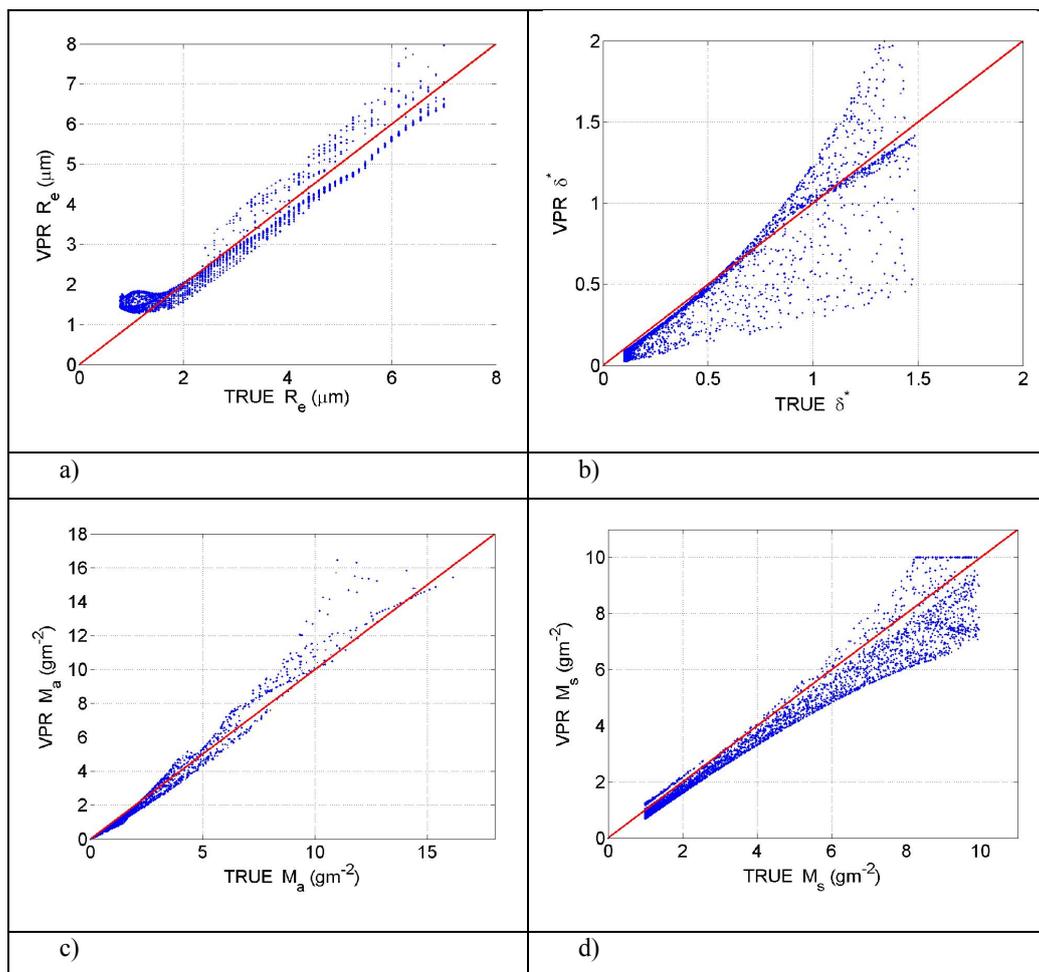
1 Figure 4. Linear trends of B_{up} (red), B_{dn} (blue), τ_t (green), and B_s (cyan) versus B_p for 48 different
 2 plumes (12 months and 4 heights) each obtained from a set of MODTRAN simulations.
 3



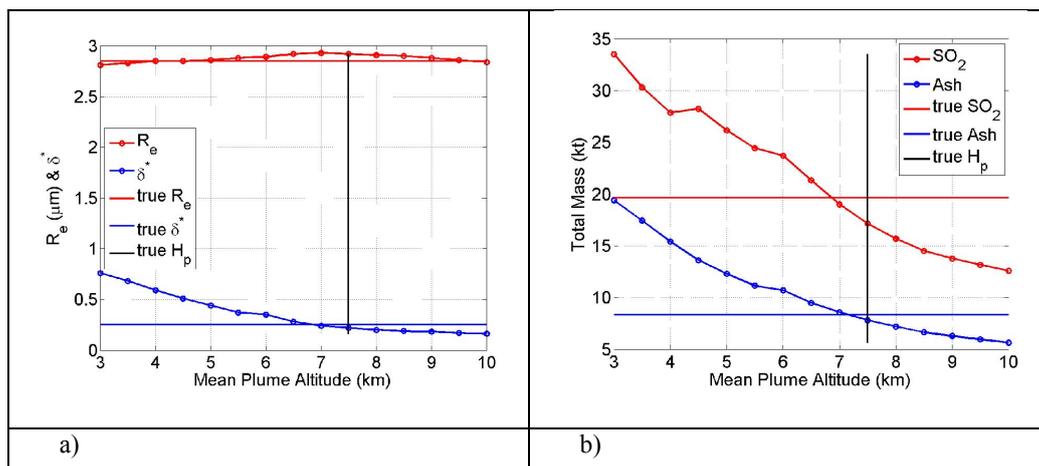
1 Figure 5. Scatter plots between the plume transmittance (obtained from a wide set of MODTRAN
2 simulations) for MODIS-Aqua bands at 11 and 8.7 μm for the pumice (Volz, 1973) ash type (a), and
3 for water droplets (b).
4



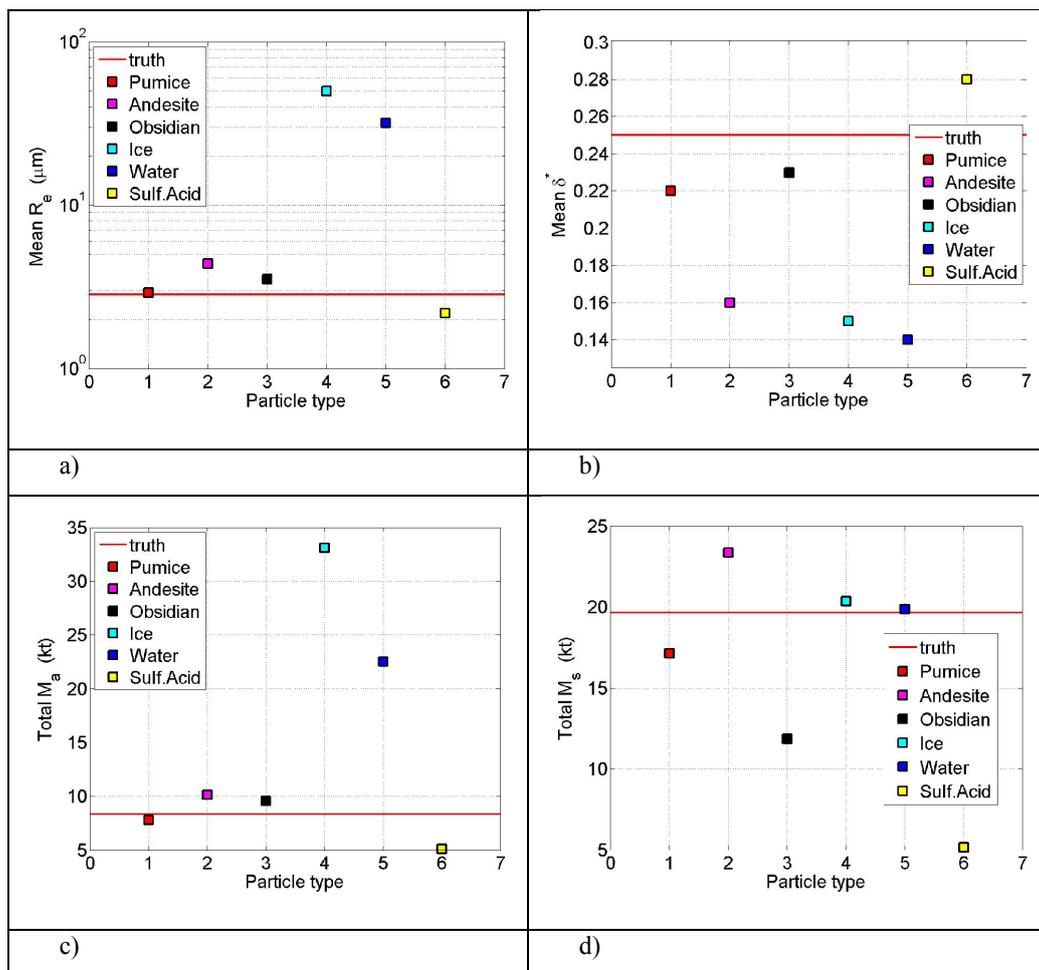
- 1 Figure 6. Synthetic images (radiance at the sensor); RGB: bands at 8.7, 11, and 12 μm respectively.
- 2 a) Mt. Etna 26 October 2013 at 12:20 GMT; b) Eyjafjallajökull 11 May 2010 at 14:05 GMT
- 3



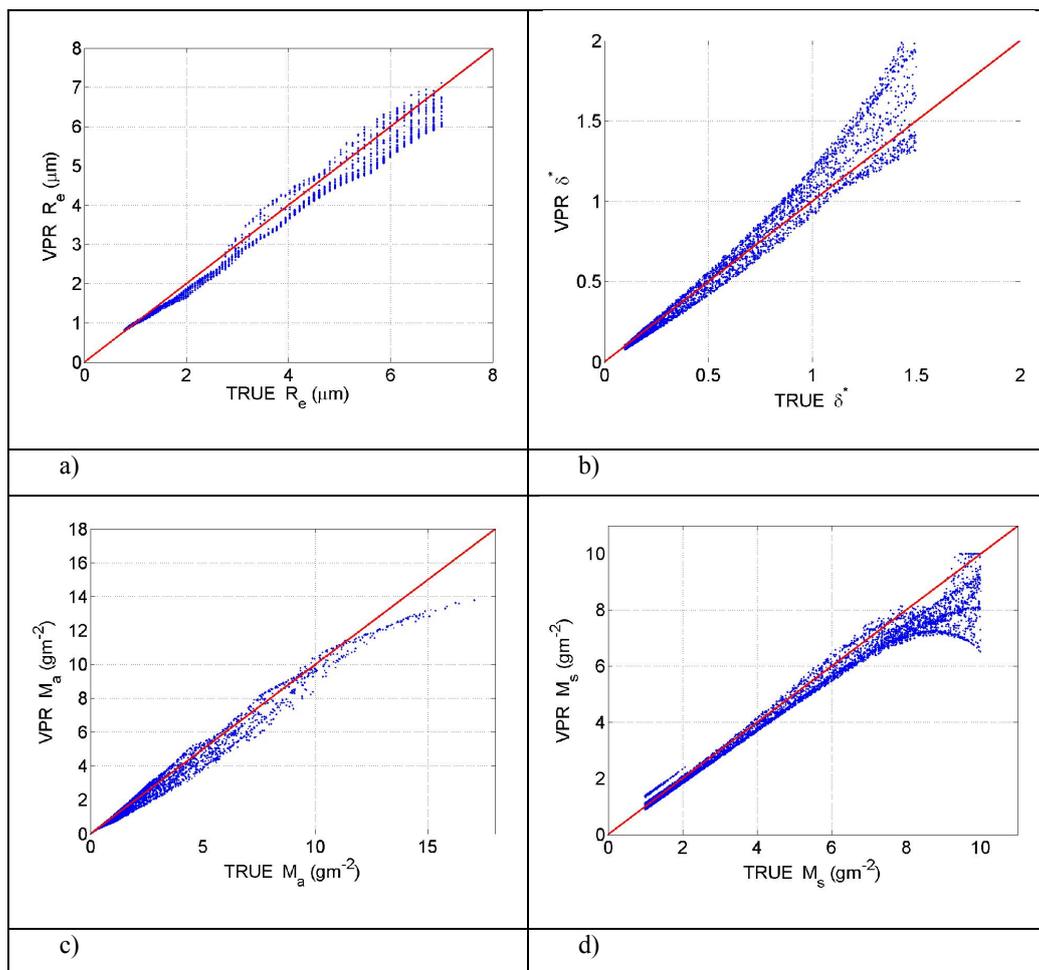
1 Figure 7. Etna-Pumice example: scatter plots between VPR results and true values: effective radius
2 (a), ash optical depth at 550 nm (b), ash mass (c), and SO_2 mass (d). Red line is the bisector.
3



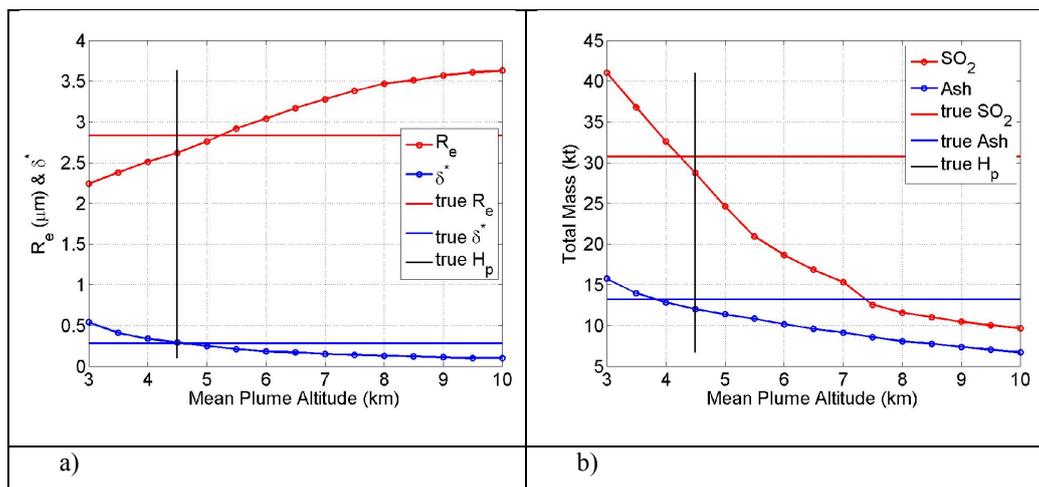
1 Figure 8. Etna-Pumice example: trends of R_e and δ^* mean values (a), and ash and SO_2 total mass (b)
 2 retrieved by VPR with different input plume altitudes.
 3



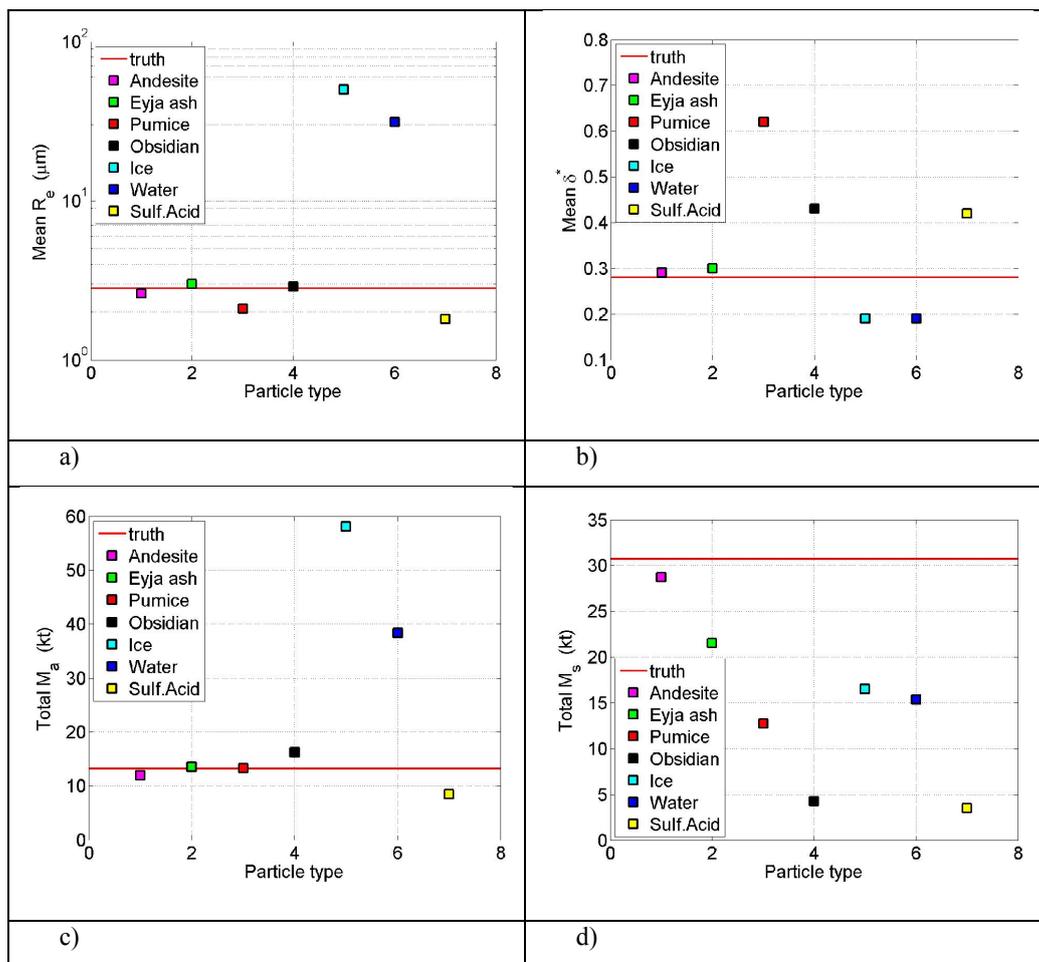
1 Figure 9. Etna-Pumice example. VPR results with different types of particles: mean effective radius
 2 (a), mean optical depth at 550 nm (b), total mass (c), and total SO₂ mass (d). The red lines are the true
 3 values.
 4



1 Figure 10. Eyjafjallajökull-Andesite example: scatter plots between VPR results and true values:
2 effective radius (a), ash optical depth at 550 nm (b), ash mass (c), and SO₂ mass (d). Red line is the
3 bisector.
4



1 Figure 11. Eyjafjallajökull-Andesite example: trends of R_e and δ^* mean values (a), and ash and SO_2
2 total mass (b) retrieved by VPR with different input plume altitudes.
3



1 Figure 12. Eyjafjallajökull-Andesite example. VPR results with different types of particles: mean
 2 effective radius (a), mean optical depth at 550 nm (b), total mass (c), and total SO₂ mass (d). The red
 3 lines are the true values.

4
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