

1 **Altitude Registration of Limb-Scattered Radiation**

2  
3  
4  
5  
6  
7  
8  
9  
10  
11  
12  
13  
14  
15  
16  
17  
18  
19  
20  
21  
22  
23  
24  
25  
26  
27  
28  
29  
30  
31

Leslie Moy<sup>1</sup>, P.K. Bhartia<sup>2</sup>, Glen Jaross<sup>2</sup>, Robert Loughman<sup>3</sup>, Natalya Kramarova<sup>1</sup>, Zhong Chen<sup>1</sup>, Ghassan Taha<sup>4</sup>, Grace Chen<sup>1</sup>, and Philippe Xu<sup>5</sup>

[1] Science Systems and Applications, Inc. (SSAI), 10210 Greenbelt Road, Suite 600, Lanham, Maryland 20706 USA

[2] NASA Goddard Space Flight Center, Greenbelt, Maryland, USA

[3] Hampton University, Hampton, Virginia USA

[4] GESTAR, Columbia, MD USA

[5] Science Applications International Corporation (SAIC), Lanham, MD

Correspondence to: Leslie Moy ([leslie.moy@ssaihq.com](mailto:leslie.moy@ssaihq.com))

32 **Abstract**

33  
34  
35  
36  
37  
38  
39  
40  
41  
42  
43  
44  
45  
46  
47  
48  
49  
50  
51  
52  
53  
54  
55  
56  
57  
58  
59  
60  
61  
62

One of the largest constraints to the retrieval of accurate ozone profiles from UV backscatter limb sounding sensors is altitude registration. Two methods, the Rayleigh Scattering Attitude Sensing (RSAS) and Absolute Radiance Residual Method (ARRM), have been developed to determine the altitude registration to the accuracy necessary for long-term ozone monitoring. The methods compare model calculations of radiances to measured radiances, and are independent of onboard tracking devices. RSAS determines absolute altitude errors, but because the method is susceptible to aerosol interference, it is limited to latitudes and time periods with minimal aerosols. ARRM can be applied across all seasons and altitudes. However, it is only appropriate for relative altitude error estimates. The application of these methods to Ozone Mapping and Profiler Suite (OMPS) Limb Profiler (LP) measurements showed that, at launch, the OMPS LP instrument had a 1-2 km altitude registration error, resulting in a 50% error in the derived ozone density at some altitudes. Though some of the error has been attributed to thermal shifts in the focal plane of the instrument, most of it appears to be due to misalignment of the spacecraft star trackers or the OMPS LP focal plane with respect to the spacecraft axes. In addition, there are  $\pm 200$  m seasonally varying errors that could either be due to errors in the spacecraft pointing information or in the geopotential height (GPH) data that we use in our analysis.

**Keywords: altitude registration, OMPS Limb Profiler, RSAS, ARRM, ozone profile, backscattered ultraviolet**

## 63 **1 Introduction**

64 Instruments that measure the solar radiation scattered by the earth's atmosphere in the limb  
65 direction provide a low cost way of measuring stratospheric ozone and aerosol profiles from  
66 satellites. The technique provides daily full coverage of the sunlit earth from commonly used  
67 polar sun-synchronous satellites. To meet long-term ozone monitoring needs (3% precision  
68 between 15 and 50 km) requires the altitude registration of the radiances to be accurate to within  
69 ~100 m. For a sensor orbiting at 800 km, this translates into ~6 arcsec accuracy in the pointing  
70 direction of the instrument line-of-sight (LOS) with respect to Earth's horizon. This is often a  
71 difficult goal to achieve.

72 In this paper we critically examine the performance of two methods of altitude registration that  
73 compare measured and simulated radiances. We discuss the inherent strengths and limitations of  
74 each method and then assess their performance using data from the OMPS Limb Profiler (LP),  
75 launched onboard the Suomi NPP (SNPP) satellite on October 28, 2011.

76 One of these techniques, known as Rayleigh Scattering Attitude Sensing (RSAS), is relatively  
77 insensitive to instrument radiometric errors and drift. However, since the 350 nm/20 km limb  
78 radiances are greatly affected by aerosols, it works best where there is minimal aerosol loading.  
79 Under these conditions the accuracy of the method is limited by the accuracy of the geopotential  
80 height (GPH) data near 3 hPa (~40 km). Since aerosol contamination limits the range of latitudes  
81 and seasons where RSAS can be applied, we developed the Absolute Radiance Residual Method  
82 (ARRM). Although ARRM can be applied more broadly than RSAS, it is more suitable to  
83 analyzing relative rather than absolute errors in limb altitude registration.

84 We describe the theoretical basis of these two techniques in Section 2, and move on to results for  
85 the OMPS LP instrument in Section 3. Finally, we present several validations of our uncertainty  
86 estimates in Section 4.

87

## 88 **2 Theoretical Basis**

89 Most scene-based altitude registration methods applied to limb-scattering instruments take  
90 advantage of the fact that the atmospheric Rayleigh scattering measured by these instruments  
91 varies by 12-14%/km in the absence of particulate scattering by aerosols and clouds and  
92 absorption by trace gases. For wavelengths longer than 310 nm, the limb-scattered radiance has a  
93 significant contribution from diffuse upwelling radiance (DUR), which is affected by

94 tropospheric clouds, aerosols and surface albedos from inside a circular cone whose base extends  
95 hundreds of km to the horizon. At non-ozone absorbing wavelengths DUR can be as much as  
96 half of the measured radiance. Since DUR is challenging to model accurately, all successful  
97 altitude registration methods must be relatively insensitive to variations within the cone.

98 The RSAS method, described in Sect. 2.1, employs signal ratios in which the DUR effects  
99 largely cancel. The ARRM, described in Sect. 2.2, uses 295 nm radiances for which ozone  
100 absorption screens the DUR signal. The Knee method, described in Sect. 2.3, has been used  
101 extensively by others (Sioris et al., 2003; Kaiser et al., 2004; Rault et al., 2005; von Savigny et  
102 al., 2005, Taha et al., 2008), but our analysis indicates that it has no advantages over RSAS and  
103 ARRM.

104

## 105 **2.1 Rayleigh Scattering Attitude Sensor (RSAS)**

106 This technique is named after a sensor that was flown on the Space Shuttle STS-72 in January  
107 1996 (Janz et al., 1996) to test the concept originally proposed by Bhartia in 1992. The technique  
108 takes advantage of the fact that the gradient in the log of the limb-scattered (LS) radiance  $I$  with  
109 altitude  $z$ ,  $d\ln I/dz$ , changes by a factor of 3 between 40 km and 20 km for wavelengths near 350  
110 nm (Fig. 1). This is caused by the exponentially increasing attenuation of Rayleigh scattering  
111 with pressure. At 40 km this attenuation is small and when aerosol loading is minimal,  $d\ln I/dz$  is  
112 largely determined by  $d\ln P/dz$ , where  $P$  is the atmospheric pressure at altitude  $z$ . However, at 20  
113 km the extinction and scattering nearly cancel where the line of sight (LOS) intersects the Earth  
114 radius vector at a right angle, called the tangent point (TP). Therefore the radiances at 20 km are  
115 relatively insensitive to the exact altitude of the TP. Though several variations of the RSAS  
116 technique have been developed (McPeters et al., 2000; Rault et al., 2005; Taha et al., 2008), we  
117 find that the simplest formulation described below works as well as any other.

118 If  $r$  is the ratio of radiances for wavelength  $\lambda$  at altitudes  $z_1$  and  $z_2$ , and  $s_1$  and  $s_2$  are the vertical  
119 slopes  $d\ln I/dz$  at those altitudes, then the error in tangent height (TH) can be calculated as  
120 follows:

$$121 \Delta z = - \frac{\ln(r)_m - \ln(r)_c}{s(z_1) - s(z_2)} \quad (1)$$

122 where the subscript  $m$  refers to the measured radiance ratios, and  $c$  to the ratio calculated using a  
123 radiative transfer model. To get the most accurate estimate of TH error the denominator should  
124 be as large as possible and the uncertainties in estimating the numerator should be small. The

125 smallest uncertainties in the numerator occur at wavelengths near 350 nm, where trace gas  
126 absorption and aerosol scattering effects are small. Setting  $z_1$  to be near 40 km and  $z_2$  to be at or  
127 below 20 km maximizes the value of the denominator, typically near 0.10/km. So an accuracy of  
128 0.01 (equal to 1% in radiance ratios) is needed to estimate TH within 100 m.

129 As Fig. 2 shows, the largest source of noise in estimating the numerator comes from the  
130 mismatch in cloud sensitivity of radiances at the two altitudes. However, this noise is random  
131 and can be reduced by averaging data from multiple orbits.

132 Aerosols in the instrument's LOS are a more significant source of error. Though the effect of  
133 aerosols near 350 nm is small compared to longer wavelengths, it is difficult to model (Fig. 3)  
134 due to the determination of subtle differences between two large effects: the reduction of  
135 Rayleigh scattering by aerosol extinction and the enhancement of limb radiances by aerosol  
136 scattering. In addition, 350 nm LS radiances at 20 km are significantly affected by variation of  
137 aerosols along the LOS because of large Rayleigh attenuation; aerosols in the LOS close to the  
138 sensor contribute more to the radiance than those far away. Though this effect is similar to the  
139 cloud effect mentioned earlier, it is not random because aerosols tend to have systematic  
140 latitudinal variability. Given this complexity, the RSAS method works best in latitudes and  
141 months where the 350 nm aerosol extinction at 20 km is relatively small.

142 Another potential source of uncertainty in applying the RSAS technique comes from uncertainty  
143 in simulating  $r_c$  at 40 km; one needs to have accurate pressure profiles for altitudes at and above  
144 40 km. If the pressure profiles are obtained from geopotential height (GPH) profiles provided by  
145 meteorological data assimilation systems, a one-to-one relationship exists between the two errors:  
146 a 100 m error in GPH at 3 hPa translates into ~100 m error in determining TH altitude.

147

## 148 **2.2 Absolute Radiance Residual Method (ARRM)**

149 This method uses radiances measured by a limb instrument near 295 nm at ~65 km to determine  
150 altitude error. The main advantage of ARRM is that it reduces aerosol contamination effects  
151 because, with the exception of polar mesospheric clouds (PMCs), the atmosphere is typically  
152 free of particulate matter at 65 km. PMCs, which form in the polar summer and are typically  
153 located at 80 km, can significantly affect 65 km limb radiances if they are in the LOS of the  
154 instrument. Fortunately most of the PMC contamination is screened using a 353 nm channel  
155 radiance residual flag at 65 km. Though 295 nm radiances are very ozone sensitive, this

156 sensitivity drops to less than 0.2% for a 10% change in ozone above 65 km. At this altitude the  
157 ozone concentration typically changes by 25% per km, so model radiance errors will be within  
158 0.5% provided the OMPS TH is accurate to 1 km. A second iteration of ARRM will remove any  
159 vestige of ozone sensitivity.

160 The main difficulty in applying ARRM is the inaccuracy of GPH data near 0.1 hPa needed to  
161 calculate 295 nm radiances at 65 km. To reduce this error we developed a variation of a  
162 technique that has been used for many years to derive mesospheric temperature profiles from the  
163 vertical slope of Rayleigh-scattered radiances measured by ground-based UV lidars (McGee et  
164 al., 1991). Temperatures were computed using the relative density differences between  
165 successive altitudes where the scattering mechanism was purely Rayleigh. Since our technique  
166 relies on this region of Rayleigh dominance, their technique can be applied to GPH (which is  
167 related to temperature assuming hydrostatic balance). The 350 nm and 295 nm residuals are  
168 affected similarly by the errors in the GPH with altitude so we use the 350 nm residual to correct  
169 for the GPH errors at 295 nm. Similarly to the extent that stray light is wavelength independent,  
170 this correction will correct for stray light.

171 The residual at wavelength  $\lambda$  at altitude  $z$ , defined as  $d(\lambda,z) = \ln I_m(\lambda,z) - \ln I_c(\lambda,z)$ , is corrected  
172 using 350 nm residuals:

$$173 \quad d_{corr}(\lambda, z) = d(\lambda, z) - [d(350, z) - d(350, z_0)] \quad (2)$$

174 where  $z_0$  is a normalization altitude.

175 The 350 nm differential residuals on the right side of equation (2) provide an estimate of the  
176 relative error in calculating radiances using meteorological data between  $z$  and  $z_0$ . Since this  
177 error should be wavelength independent, we can use this term to correct the residuals at any  
178 wavelength, assuming that the meteorological data at  $z_0$  is accurate and that the 350 nm  
179 wavelength is well calibrated. The large response of OMPS LP at 350 nm results in signals that  
180 are the least affected by out-of-band stray light.

181 The TH error estimated using this method is given by:

$$182 \quad \Delta Z = \frac{d_{corr}(\lambda, z)}{s(\lambda, z) - [s(350, z) - s(350, z_0)]} \quad (3)$$

183 We are minimizing ozone profile sensitivity by applying this method to radiances at wavelengths  
184 shorter than 300 nm. At longer wavelengths DUR makes the LS radiances sensitive to total  
185 column ozone at all altitudes. At 295 nm, the use of  $z$  near 65 km provides low ozone sensitivity.

186 Though it is best to set  $z_0$  as low as possible to minimize GPH caused errors, aerosol  
187 contamination limits the value to around 40 km.

188 ARRM has two primary uncertainties. Since 1% error in radiance calibration produces ~70 m  
189 error in determining the TH, this method requires accurate radiances (or sun-normalized  
190 radiances) and may be affected by instrument degradation. Though the absolute accuracy of  
191 ARRM may not be as good as RSAS, this method can be applied at latitudes/seasons where  
192 RSAS cannot be applied reliably because of aerosol contamination. And like RSAS, this method  
193 is also sensitive to errors in GPH profile near 3 hPa, which are used for calculating 350 nm  
194 radiances at 40 km.

195

### 196 **2.3 “Knee Method”**

197 The name of this method is derived from the characteristic knee shape of the limb radiance  
198 profiles (Fig. 4). Above the knee the radiances decrease with altitude due to exponential decrease  
199 in Rayleigh scattering and ozone density. Below the knee ozone absorption becomes so large that  
200 it essentially blocks most of the Rayleigh-scattered radiation from reaching the satellite, making  
201 the radiances insensitive to atmospheric pressure. This characteristic shape allows estimations of  
202 altitude registration error in a manner very similar to that of RSAS. An advantage of this method  
203 is the ability to use shorter wavelengths which are less sensitive to aerosols. However, the  
204 method requires accurate ozone and pressure profiles near and above the knee region. Radiative  
205 transfer calculations using climatological ozone profiles indicate that a 10% error in assumed  
206 ozone density (at all altitudes) will produce about a 250 m error in altitude registration (Fig. 5).  
207 The method also has a sensitivity to GPH errors that are similar to RSAS and ARRM. In our  
208 view this method provides no compelling advantage over comparing the ozone profiles retrieved  
209 from a limb scattering sensor with other ozone sensors to determine altitude registration errors.  
210 Indeed, direct ozone comparisons are simpler and more reliable if the altitude registration error is  
211 the largest error source, and we use this technique to evaluate the results of RSAS and ARRM in  
212 Sect. 3.

213

## 214 **3 Results**

215 In this section we discuss altitude registration errors in OMPS LP radiances determined first by  
216 “Slit Edge” analysis of the instrument focal plane image and then by the application of the RSAS  
217 technique. The remaining errors are analyzed using ARRM.

218

### 219 **3.1 Slit Edge Results**

220 The OMPS LP sensor utilizes a two-dimensional charge coupled device (CCD) detector to  
221 capture spectrally dispersed (along the 740 pixel row dimension) and vertically distributed (along  
222 the 340 pixel column dimension) radiation (Fig. 6). Three long vertical entrance slits spaced  $4.25^\circ$   
223 apart produce three distinct images of the atmosphere that are collected simultaneously on the  
224 single CCD. The resulting limb radiance profile from the center slit is aligned very closely to the  
225 satellite ground track with tangent points trailing approximately 3000 km south of the sub-  
226 satellite point. The east and west slit images are separated in longitude by  $2.25^\circ$  (250 km at their  
227 tangent points) from that of the center slit.

228 An unexpected thermal sensitivity was discovered in the LP instrument soon after launch (Jaross  
229 et al., 2014). Expansion of the LP instrument’s entrance baffle as the sun illuminates it midway  
230 through the northern hemisphere causes mirrors in the telescope to rotate slightly, which in turn  
231 moves the limb radiance image on the detector. Since there are separate mirrors for each entrance  
232 slit, the three slit images move independently. These image motions cause misregistration of  
233 both the vertical pointing and center wavelength of each pixel. Vertical pointing changes are  
234 detected most clearly by observing the location (detector column) of the lower slit edge, which  
235 has a sharp signal gradient. Figure 7 contains plots of the average edge locations in the vertical  
236 (altitude) dimension along the orbit. These pointing shifts are very repeatable (ranging only  $\pm 15$   
237 m at a given point in the orbit over a year).

238 Since the same slit edge analysis can be applied to pre-launch test data, it is possible to obtain the  
239 pixel line of sight shift relative to its calibrated value in the spacecraft reference frame. There is  
240 no evidence of image distortion so this shift is the same for all detector pixels within a slit image.  
241 The edge analysis indicates the three slit edges shifted by the equivalent of 570/470/950 m  
242 (east/center/west slits, respectively) at the middle of an orbit relative to pre-launch measurements.  
243 A mean sensor temperature decrease exceeding  $25^\circ\text{C}$  from ground to on-orbit conditions is the  
244 suspected cause. We believe there are no additional uncorrected pointing shifts arising from



245 within the LP instrument. An error or change in the alignment of the instrument with respect to  
246 S/C axes is not detectable using this method.

247

### 248 **Section 3.2 RSAS results**

249 We use the radiative transfer code described by Loughman et al. (2015) to estimate 350 nm  
250 radiances. Since the 40/20 km radiance ratio is not sensitive to polarization effects, we use the  
251 faster scalar code rather than the full vector one to calculate DUR. The calculations are done  
252 assuming a pure Rayleigh atmosphere bounded by a Lambertian reflecting surface at 1013.25  
253 hPa. The reflectivity of this surface is calculated using limb measurements at 40 km. However,  
254 both measurements and calculations show that the ratio of 40/20 km radiances is not affected by  
255 reflectivity or surface pressure and there is no discernible cloud effect. Since NO<sub>2</sub> only has a  
256 very small (<0.5%) effect on 350 nm radiances, climatological NO<sub>2</sub> profiles are sufficient for the  
257 calculation. We use OMPS LP retrieved ozone profiles. The RSAS analysis is not sensitive to  
258 ozone assumptions because it uses the ozone insensitive 350 nm radiances.

259 We estimate pressure and temperature versus altitude at the LP measurement locations and time  
260 from the Modern-Era Retrospective Analysis for Research and Application (MERRA) data  
261 (GEOS-5 FP\_IT Np) from the Global Modeling and Assimilation Office (GMAO) at NASA  
262 Goddard Space Flight Center (GSFC). The data are provided as geopotential heights (GPH) at 42  
263 pressures from the surface to 0.1 hPa, on a 0.5° latitude x 0.625° longitude horizontal resolution  
264 grid, and at a 3 hour interval. The GPH is converted to geometric height using a standard formula  
265 that takes into account the variation of gravity with latitude and elevation.

266 As discussed in Sect. 2.1, RSAS results are affected by aerosols near 20 km. Aerosol profiles  
267 derived from the Optical Spectrograph and InfraRed Imaging System (OSIRIS) data (Llewellyn  
268 et al, 2004; Bourassa et al, 2007) indicate that tropical aerosols reached a minimum value (during  
269 the OMPS lifetime) just before the eruption of the Kelud volcano in Indonesia on February 14,  
270 2014 (Fig. 8). We have therefore chosen to use equatorial RSAS data before the eruption to  
271 estimate the altitude registration errors (listed in Table 1). These TH errors range between ~1 and  
272 ~1.5 km for the three slits, and have been applied to the Version 2 OMPS LP Ozone data set.  
273 Radiative transfer calculations using OSIRIS-derived aerosol profiles indicate that the aerosol  
274 caused errors in the results shown are less than 100 m.

275

276 **Section 3.3**    **ARRM results**

277 We utilized the same radiative transfer code and profile inputs used for RSAS to calculate  
278 radiances at 295 nm. Ozone concentrations from the OMPS LP retrievals were used. As  
279 mentioned previously, though 295 nm radiances are very ozone sensitive, this sensitivity drops to  
280 less than 0.2% for a 10% change in ozone above 65 km. As discussed in Sect. 2.2, the absolute  
281 accuracy of ARRM may not be as good as RSAS, but this method can be applied at latitudes and  
282 seasons where RSAS cannot be applied reliably because of aerosol contamination.

283 Time dependent plots (Fig. 9) show negative pointing trends of approximately 100 m over the  
284 four years of data, and even larger seasonal variations depending on the latitude band and slit  
285 over the four years of data. Much of this trend is the result of a 6 arcsec (a TH change of ~100m)  
286 spacecraft pitch adjustment that occurred on 25 April 2013. Figure 9 clearly shows this abrupt  
287 change.

288 The largest disagreement between the 3 slits (~400 m) occurs in the high northern hemisphere. In  
289 the southern hemisphere the disagreement is about 100 m. In addition we see  $\pm 200$  m seasonally  
290 varying errors that could be due to either true pointing changes or errors in the GPH data that we  
291 used in our analysis. Such variations in 295 nm radiances cannot be explained by known  
292 seasonality in ozone concentration at 65 km.

293 The ARRM method is designed to accommodate stray light errors that are independent of  
294 wavelength. No additional TH errors occur when stray light at 65km is the same at 295 and  
295 350nm. The ground characterization of Limb sensor stray light indicates a small wavelength  
296 dependence (Jaross, 2014), and this is removed in ground processing. Our subsequent  
297 comparisons with RTM predictions indicate that residual stray light errors at 65 km have a daily  
298 mean bias that translates to less than 100 meters in TH.

299 The ARRM analysis shows a distinct latitude dependence (Fig. 10) with some seasonal  
300 differences. While it is tempting to attribute this entirely to TH error, we conservatively do not  
301 apply the ARRM results to our data since the uncertainties are of the same magnitude. As with  
302 RSAS, this method is sensitive to errors in GPH profile near 3 hPa used for calculating 350 nm  
303 radiances at 40 km (see Sect. 4.1). Further analysis is needed to determine the precise cause of  
304 the remaining errors.

305

306 **Section 4 Validation**

307 In this section we consider uncertainties in the parameters used to derive TH errors to indirectly  
308 validate the results shown in Sect. 3, and to estimate the remaining uncertainties in the LP TH  
309 due to errors in these parameters. Section 4.1 focuses on the validation of 3 hPa GPH  
310 information from MERRA that was assumed as the truth in our calculations. In Section 4.2 we  
311 consider the reflectivity measured by the OMPS nadir sensor to validate LP-measured radiances  
312 at 350 nm, which vary by  $\sim 14\%/km$  around 40 km. Finally, in Sect. 4.3 we compare the LP-  
313 derived ozone mixing ratio at 3 hPa with the Microwave Limb Sounder (MLS). For the  
314 validation studies in this section, the OMPS LP TH has been corrected with the errors listed in  
315 Table 1.

316

317 **Section 4.1 GPH comparison**

318 Errors in the GPH profile assumptions directly translate into TH errors. Although the 3 hPa  
319 GPH varies over 4 km along an orbit, a comparison of daily averaged values from MLS and  
320 MERRA show differences that are usually less than 200 m (Fig. 11). These differences do not  
321 directly explain the latitude dependence of TH errors shown in Fig. 10, but do provide an  
322 estimate of the magnitude of errors caused by the use of MERRA GPH in our radiative transfer  
323 calculations. There is better agreement at the poles, but this may be due to the reliance on  
324 climatology where there are scant measurements. As a result, it is not clear how these GPH  
325 errors influence the ARRM results. However, in Section 4.3 we discuss some suggestive but  
326 inconclusive results untangling GPH errors from TH errors.

327

328 **Section 4.2 Radiances comparison**

329 We previously described (in Section 2) the difficulty modeling DUR caused by scene  
330 heterogeneity and aerosols. Both the RSAS method and ARRM depend upon an accurate model  
331 for DUR at 40 km relative to other altitudes, and any model errors translate directly (Equations 1  
332 and 3) into false estimates in the TH errors.

333 We estimate the DUR modeling error by comparing LP measurements and modeled 353 nm  
334 radiances at 3 hPa. The OMPS Nadir instrument makes nearly simultaneous measurements from  
335 a smaller field-of-view (50 x 50 km at nadir). The surface reflectivity derived from its 340 nm

336 radiances are therefore relatively insensitive to DUR effects (compared to measurements derived  
337 from the LP measurements). With better reflectivity assumptions the model/measurement  
338 comparisons offer a lower bound of the effect of DUR modeling errors.

339 The radiance comparison, shown in Fig. 12, suggests model or calibration errors of 2-3% on  
340 average, plus structures caused by the limb and nadir scene mismatch. If this error were  
341 attributed solely to the limb modeled DUR effect, the resulting TH error would be less than +/-  
342 200 m. There is no evidence of either a seasonal or a latitude dependence in the four days of  
343 comparisons, meaning that DUR effects cannot explain the robust seasonal and latitudinal  
344 variations seen in ARRM results (Fig. 9 and Fig. 10). These model/measurement comparisons  
345 provide an estimate of errors related to incomplete modeling of DUR and inhomogeneous  
346 surface albedo included in our RSAS and ARRM results. We therefore conclude those variations  
347 arise from errors in the GPH scale or from true TH variations.

348

### 349 **Section 4.3 Ozone comparison**

350 At 3 hPa limb ozone retrievals are very sensitive to TH errors, with 20 to 25% per km change in  
351 ozone concentration (see Fig. 5). Similar to the Knee Method, we can use this sensitivity to  
352 gauge the residual TH errors. We compare LP ozone retrievals against Aura MLS v4 ozone  
353 retrievals at 3 hPa (near 40 km) (Fig. 13). While the latitudinal patterns of differences  
354 significantly vary with season, we find agreement within  $\pm 10\%$  over all seasons and latitude  
355 bands. If completely interpreted as due to TH error, a 10% difference would translate to less than  
356 500 m error. These comparisons confirm a residual uncertainty in our scene-based altitude  
357 registration techniques of  $\pm 200\text{m}$ .

358 The ARRM method has displayed the ability to track any drifts or sudden changes of 50 m (Sect.  
359 3.3), and time series of TH error derived from the ARRM method track very closely to the time  
360 series of the LP/MLS 3 hPa ozone differences (Fig. 14). The highest correlation (0.76) was found  
361 at  $45^\circ$  south latitude, with considerable smaller values in the northern hemisphere (0.30 at  $60^\circ$   
362 north). Whether this suggests the ARRM results can be attributed solely to TH errors has not  
363 been determined yet.

364 Both ARRM and LP/MLS ozone comparison depend upon accurate TH and MERRA  
365 information, and in the same way. So, while these results suggest some confidence in the  
366 ARRM technique, we cannot assign the correlation shown in Fig. 14 to only a TH error or a

367 MERRA error. It is important to note that MLS ozone profiles are reported as volume mixing  
368 ratio on a vertical pressure grid, while the LP algorithm retrieves ozone as number density on an  
369 altitude grid. Thus, in order to compare LP and MLS ozone retrievals we had to convert ozone  
370 number densities to mixing ratios using MERRA temperature and GPH profiles. This conversion  
371 inevitably introduces errors in MERRA GPH and temperature into the ozone comparisons.  
372 Therefore ozone differences between LP and MLS ozone retrievals not only depend on the LP  
373 TH error, but on errors in MERRA GPH as well as on errors in the retrieval algorithms and  
374 instrumental sampling (geophysical noise). Furthermore, analysis of LP and MLS ozone  
375 retrievals indicates a large daily ozone variability within a 5-degree latitude bin at 3 hPa that  
376 ranges from 2% in the tropics to 20% at high latitudes with the seasonal maximum during austral  
377 winters (results are not shown here), which can give readers a sense of geophysical ozone  
378 variability. In consideration of all of the above factors, we remain cautious in making definite  
379 conclusions and applying time-dependent corrections for the LP TH at this time; further analysis  
380 and comparisons with independent ozone observations (like SAGE III) are needed to confirm the  
381 results.

382

## 383 **Section 5      Conclusions**

384 Accurate altitude registration is key to the success of the limb scattering measurement  
385 technique. We have described two scene-based techniques that together provide highly precise  
386 and accurate estimates of the tangent height. These altitude registration techniques are  
387 inexpensive and more comprehensive than external sources of altitude information, such as star  
388 trackers mounted on the spacecraft. Star trackers are typically used on spacecraft when accurate  
389 pointing knowledge is desired, but this accuracy does not necessarily transfer to the limb  
390 measurements. In Section 3.1 we've described pointing shifts that occur within the sensor optics.  
391 Though we were able to calibrate thermal sensitivities within the OMPS instrument, we have yet  
392 to identify the source of 1-1.5 km pointing errors derived from RSAS (see Table 1). These may  
393 arise from mounting offsets of the instrument and star trackers, or from spacecraft flexure  
394 between the two.

395 The RSAS and ARRM techniques are complementary because the former is  
396 accurate to  $\pm 200$  m, but only under limited conditions. The accuracy of  
397 ARRM cannot be easily established, but it has a precision within

398 ±200 m. We believe this results in small, less than 100 m, trend  
399 uncertainties for sufficiently long time series, as demonstrated by the  
400 OMPS ARRM record.

401 The single largest source of uncertainty in both techniques is knowledge of the vertical profiles  
402 of pressure, which must be provided from external sources. Given uncertainties in GPH data, as  
403 seen in the MLS comparison, as well as uncertainties in our ozone retrieval algorithm (not  
404 related to TH error), it is currently not possible to tell if the latitudinal and seasonal variations  
405 seen in ARRM results are caused by TH error. Further work will be needed to understand their  
406 cause. We have shown, however, that ARRM is capable of deriving multi-year trend  
407 uncertainties that are on the order of 100m or smaller. Furthermore, the two TH registration  
408 methods discussed in the paper allow us to track any drifts or sudden changes in our altitude  
409 registration to better than 50 m, which is the minimum level necessary to derive accurate ozone  
410 trends from a limb technique.

411

412

413

414 Acknowledgements: The authors gratefully acknowledge the assistance of NASA's Limb  
415 Processing Team in providing the data used in this paper. We would also like to thank Dave  
416 Flittner, Ernest Nyaku and Didier Rault, helped with the development and updates of the RT  
417 model. Finally, we'd like to acknowledge the role Didier played in laying the groundwork for the  
418 OMPS limb retrieval algorithm

419

420

## 421 **References**

422 Bourassa, A. E., Degenstein, D. A., Gattinger, R. L., and Llewellyn, E. J.: Stratospheric aerosol  
423 retrieval with optical spectrograph and infrared imaging system limb scatter measurements, *J.*  
424 *Geo-phys. Res.*, 112, D10217, doi:10.1029/2006JD008079, 2007.

425

426 Jaross, G., P. K. Bhartia, G. Chen, M. Kowitt, M. Haken, Z. Chen, P. Xu, J. Warner, and T.  
427 Kelly: OMPS Limb Profiler instrument performance assessment, *J. Geophys. Res. Atmos.*, 119,  
428 doi:10.1002/2013JD020482, 2014.

429

430 Janz, S. J., E. Hilsenrath, D. E. Flittner, and D. F. Heath, Rayleigh scattering attitude sensor,  
431 Proc. SPIE 2831, Atmospheric Ultraviolet and Space Remote Sensing: Methods and  
432 Instrumentation, 146, doi:10.1117/12.257207, 1996.

433

434 Kaiser, J. W., C. Von Savigny, K.-U. Eichmann, S. Noel, H. Bovensmann, and J. P. Burrows,  
435 Satellite-pointing retrieval from atmospheric limb-scattering of solar UV-B radiation, Can. J.  
436 Phys., 82, 1041-1052, 2004.

437

438 Llewellyn, E. J., N. D. Lloyd, D. A. Degenstein, R. L. Gattinger, S. V. Petelina, A. E. Bourassa,  
439 J. T. Wiensz, E. V. Ivanov, I. C. McDade, B. H. Solheim, J. C. McConnell, C. S. Haley, C. von  
440 Savigny, C. E. Sioris, C. A. McLinden, E. Griffioen, J. Kaminski, W. F. Evans, E. Puckrin, K.  
441 Strong, V. Wehrle, R. H. Hum, D. J.W. Kendall, J. Matsushita, D. P. Murtagh, S. Brohede, J.  
442 Stegman, G. Witt, G. Barnes, W. F. Payne, L. Piche, K. Smith, G. Warshaw, D. L. Deslauniers,  
443 P. Marc-hand, E. H. Richardson, R. A. King, I. Wevers, W. McCreath, E. Kyrölä, L. Oikarinen,  
444 G. W. Leppelmeier, H. Auvinen, G. Megie, A. Hauchecorne, F. Lefevre, J. de La Noe, P.  
445 Ricaud, U. Frisk, F. Sjöberg, F. von Scheele, and L. Nordh: The OSIRIS instrument on the  
446 Odin spacecraft, Can. J. Phys., 82, 411–422, 2004.

447

448 Loughman, R., D. Flittner, E. Nyaku, and P. K. Bhartia, Gauss–Seidel limb scattering (GSLs)  
449 radiative transfer model development in support of the Ozone Mapping and Profiler Suite  
450 (OMPS) limb profiler mission, Atmos. Chem. Phys., 15, 3007-3020, 2015.

451

452 McGee, T. J., D. Whiteman, R. Ferrare, J. J. Butler, and J. Burris, STROZ LITE: Stratospheric  
453 ozone lidar trailer experiment, Opt. Eng., 30, 31-39 (1991).

454

455 McPeters, R. D., S. J. Janz, E. Hilsenrath, T. L. Brown, D. E. Flittner, and D. F. Heath, The  
456 retrieval of O<sub>3</sub> profiles from limb scatter measurements: Results from the Shuttle Ozone Limb  
457 Sounding Experiment, Geophys. Res. Lett., 27, 2597-2600, 2000.

458

459 Rault, D.F., Ozone profile retrieval from Stratospheric Aerosol and Gas Experiment (SAGE III)  
460 limb scatter measurements, J. Geophys. Res., 110, D09309, doi:10.1029/2004JD004970, 2005.

461

462 Sioris, C. E., et al., Stratospheric profiles of nitrogen dioxide observed by Optical Spectrograph  
463 and Infrared Imager System on the Odin satellite, J. Geophys. Res., 108(D7), 4215,  
464 doi:10.1029/2002JD002672, 2003.

465

466 Taha, G., G. Jaross, D. Fussen, F. Vanhellefont, E. Kyrölä, and R. D. McPeters, Ozone profile  
467 retrieval from GOMOS limb scattering measurements, J. Geophys. Res., 113, D23307,  
468 doi:10.1029/2007JD009409, 2008.

469

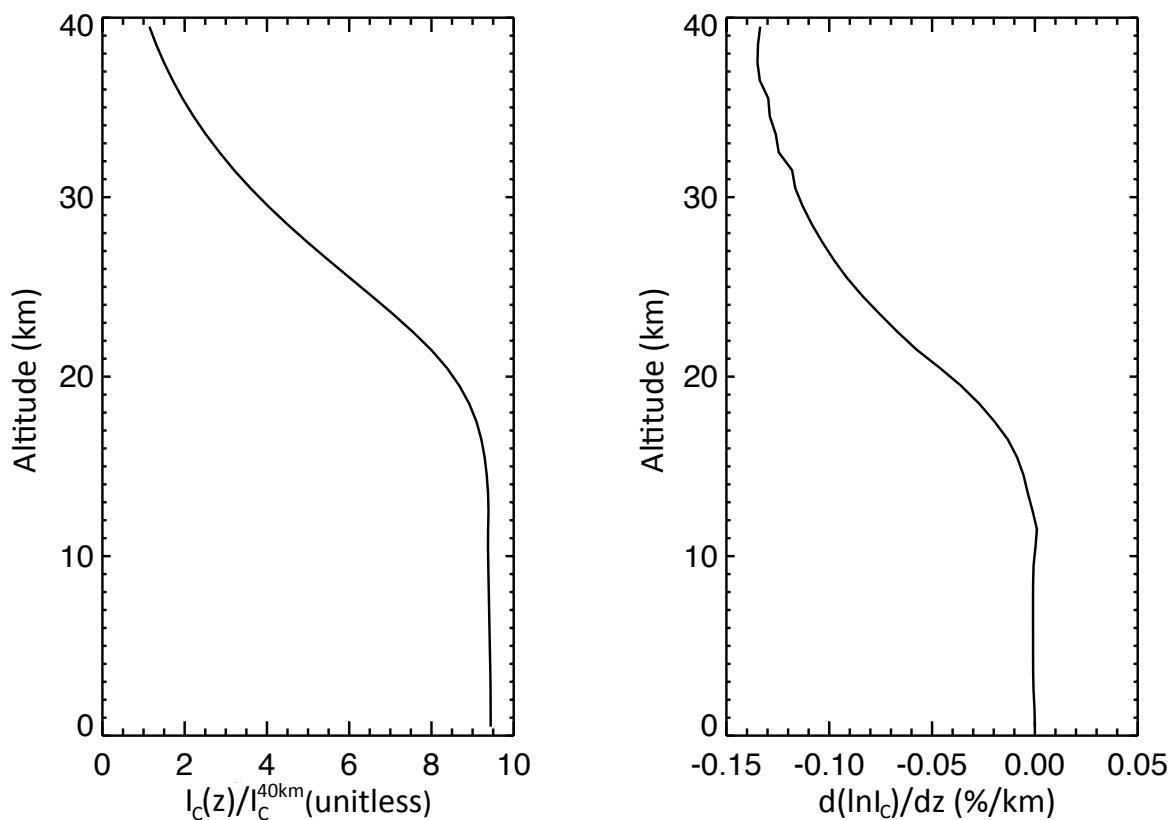
470 von Savigny, C., J. W. Kaiser, H. Bovensmann, J. P. Burrows, I. S. McDermid, and T. Leblanc  
471 (2005), Spatial and temporal characterization of SCIAMACHY limb pointing errors during the  
472 first three years of the mission, Atmos. Chem. Phys., 5, 2593-2602, doi:10.5194/acp-5-2593-  
473 2005.

475  
476  
477  
478  
479  
480  
481  
482  
483  
484  
485  
486  
487

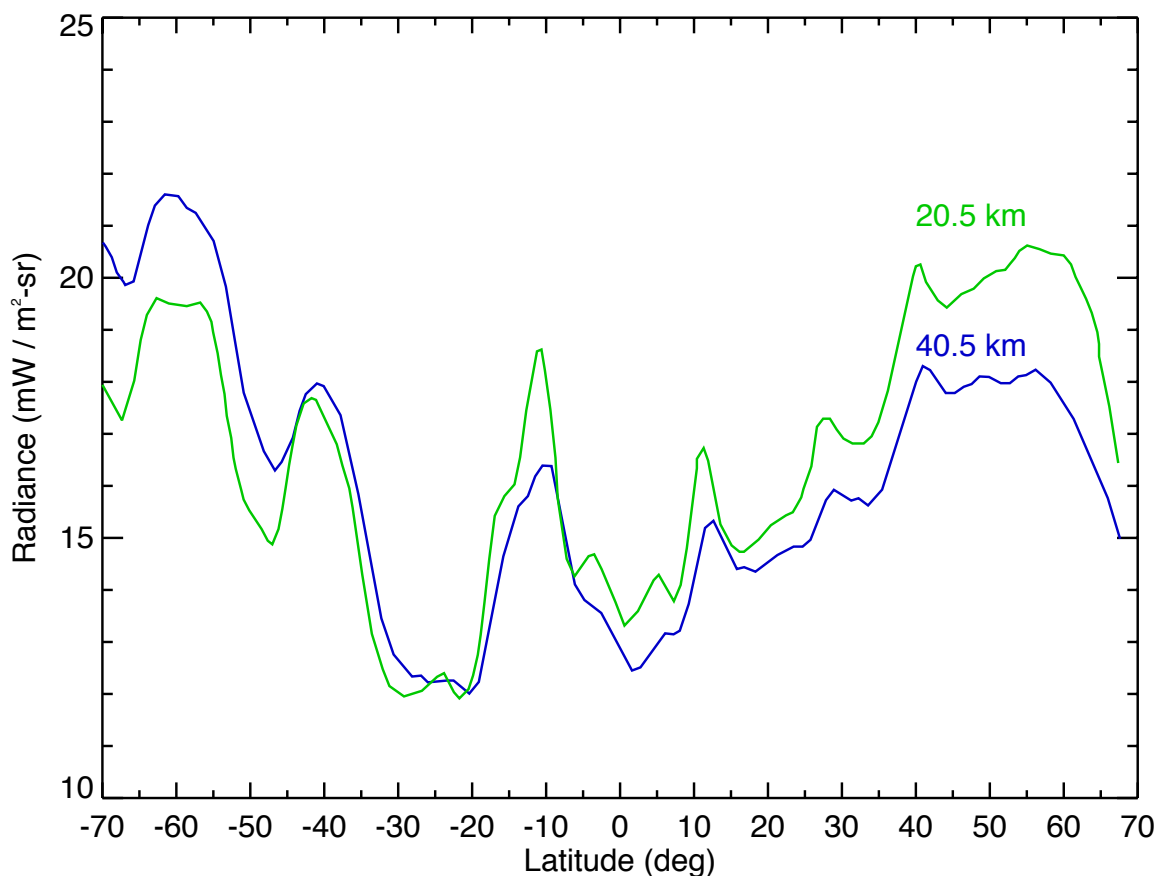
Table 1: RSAS results at the equator before the Kelud eruption 2014 February. The time period had a minimum value (during OMPS life time) and was chosen using OSIRIS measurements (Fig.8).

<b>TH error, km</b>	<b>EAST</b>	<b>CENTER</b>	<b>WEST</b>
RSAS results	1.12	1.37	1.52

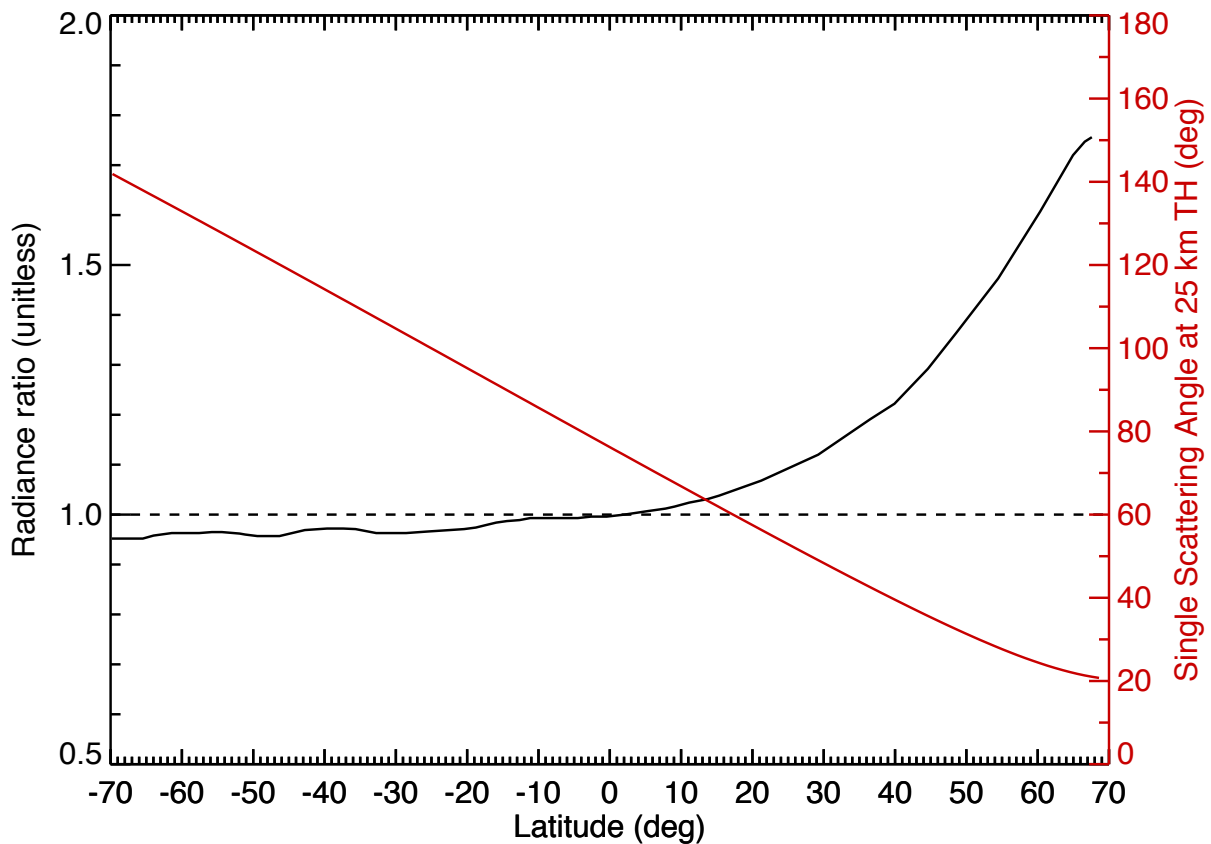




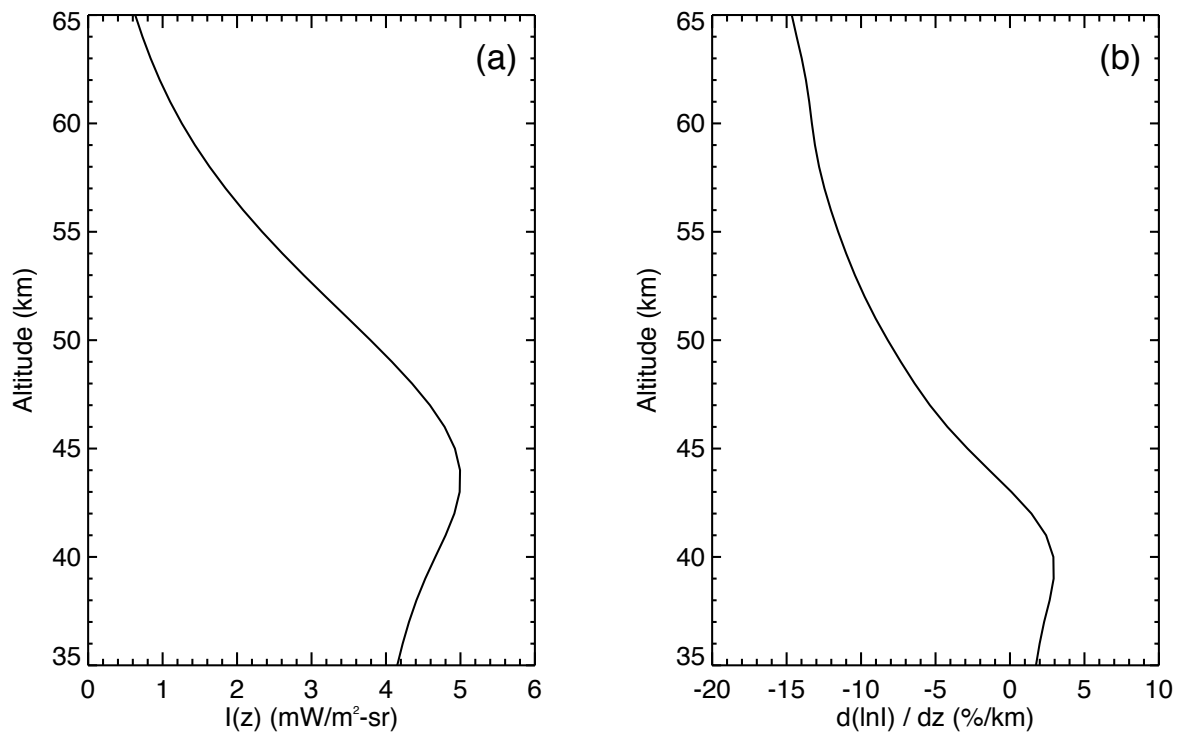
488  
 489 **Figure 1:** Figure 1a shows calculated 350 nm radiances as a function of altitude, normalized at  
 490 40.5km. The calculation models the OMPS LP field of view without aerosols. The shape of the  
 491 curve originates from the competition between molecular scattering, which increases roughly  
 492 linearly with pressure, and attenuation which becomes important when the Rayleigh optical  
 493 thicknesses near the tangent point start to become large. Attenuation causes the slope of 350  
 494 nm radiances to change sharply between 40 and 20 km (Fig. 1b), a ~9% /km difference between  
 495 20 and 40 km. The slope is used to estimate altitude registration errors by comparing measured  
 496 ratios with model simulated ratios.  
 497



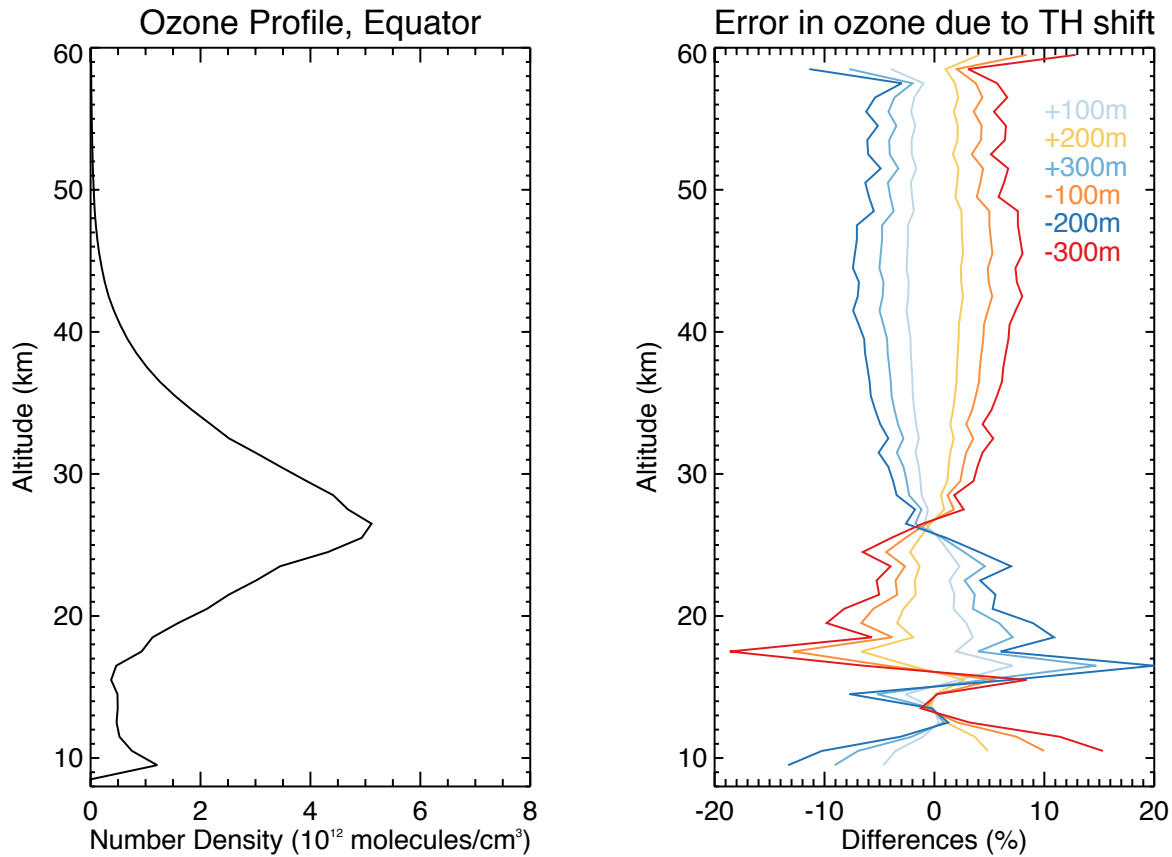
498  
 499 **Figure 2:** The 353 nm sun-normalized radiances from one orbit of OMPS LP (central slit) taken  
 500 on Feb. 2, 2012. The blue line shows 40.5 km values and the green line shows 20.5 km values  
 501 (divided by 8 to put both curves on a similar scale) versus latitude. Since the global aerosol  
 502 loading on this day was small, the short scale features in both curves are largely caused by  
 503 variations in cloud and surface albedo. The 20.5 km curve has sharper features and appears to  
 504 be shifted towards toward the South Pole. This is because large Rayleigh attenuation at 20.5 km  
 505 causes the radiances to have much higher sensitivity to the atmosphere on the satellite side of  
 506 the tangent point (TP), while 40.5 km radiances have similar sensitivities to both sides. This  
 507 effect creates large noise in applying the RSAS technique to orbital data. However, since the  
 508 noise varies randomly from orbit to orbit, it can be reduced by averaging data from multiple  
 509 orbits. Figure 6 of Loughman et al. (2015) is an example of how the contributions become  
 510 asymmetric about the tangent point at lower THs.  
 511



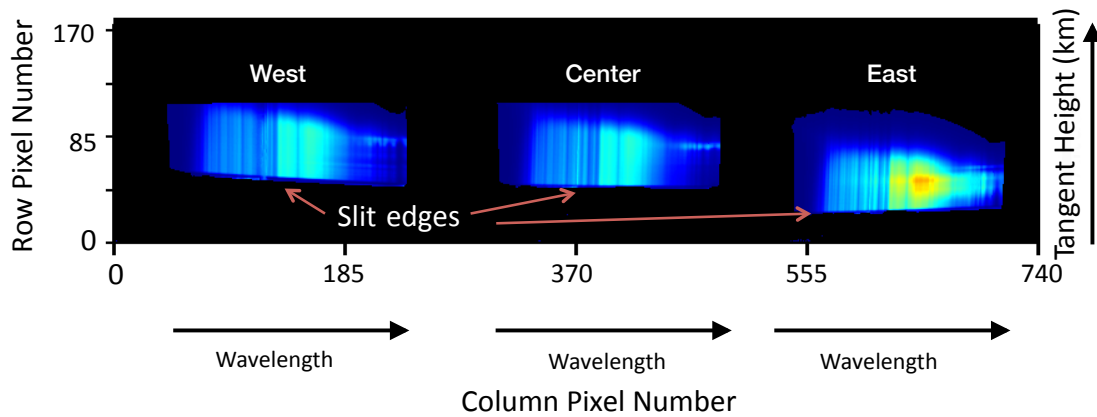
512  
 513 **Figure 3:** The ratio of 353 nm limb-scattered radiances at 20.5 km with and without aerosols  
 514 (left axis) and single scattering angle (right axis) as a function of latitude. A nominal latitude-  
 515 independent aerosol extinction profile was used in the calculation for the OMPS LP viewing  
 516 geometry on Feb 2, 2012. The strong latitude dependence is caused by an order of magnitude  
 517 change in aerosol scattering phase function with latitude combined with the attenuation of  
 518 Rayleigh-scattered radiation by aerosols along the line-of-sight (LOS). In the southern  
 519 hemisphere, where LP measures aerosols in the backscatter direction, the latter effect  
 520 dominates and the radiation decreases. The net effect is very sensitive to altitude, variation of  
 521 aerosol extinction profile along the LOS, and aerosol particle size distribution, and is therefore  
 522 difficult to calculate accurately.  
 523



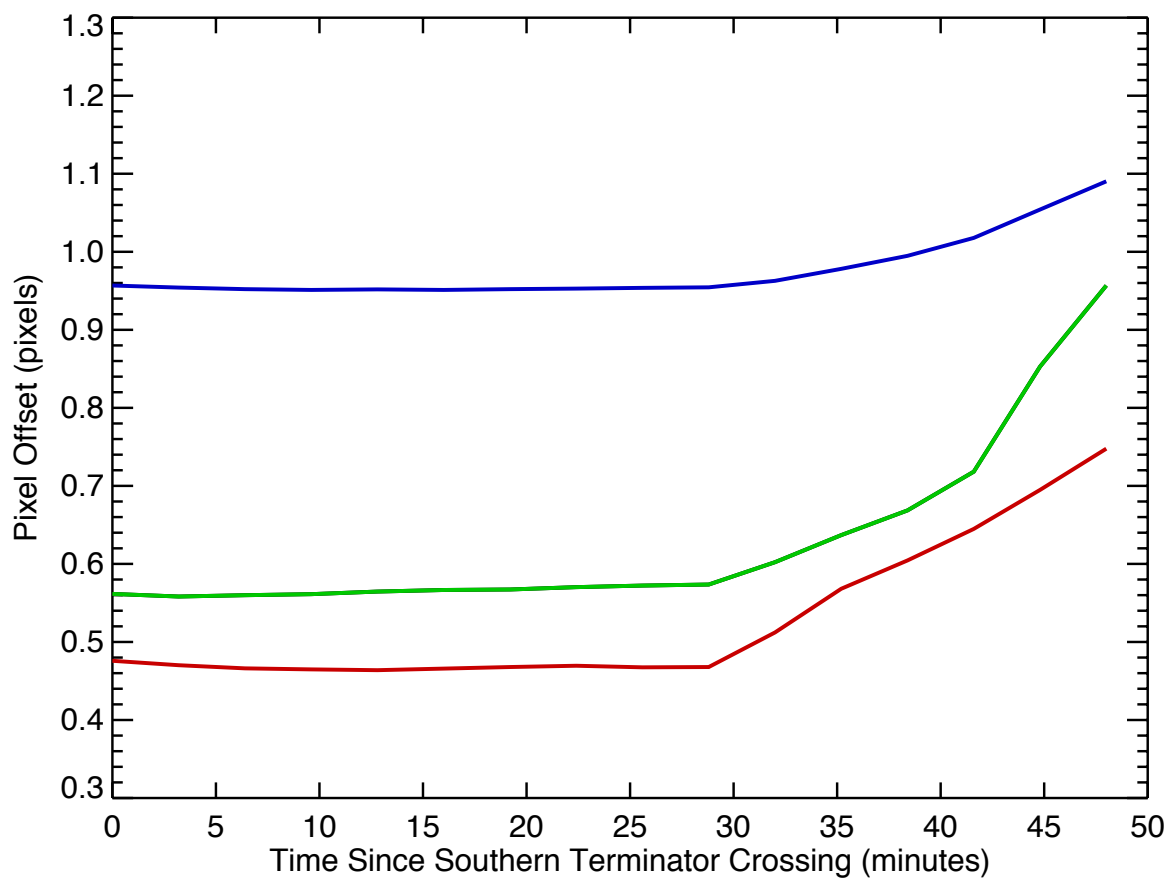
524  
 525 **Figure 4:** Figure 4a shows calculated 305 nm radiances as a function of altitude assuming an  
 526 atmosphere with no aerosols. The slope (Fig. 4b) is caused by competition between Rayleigh  
 527 scattering and ozone absorption near the altitude of maximum radiance,  $\sim 44$  km. Above 55 km  
 528 the sensitivity is nearly constant in height,  $\sim 13\%/km$  at 65 km. Above the knee the radiances  
 529 decrease with altitude due to the exponential decrease in Rayleigh scattering and ozone  
 530 density. Below the knee the ozone absorption becomes so large that it essentially blocks most  
 531 of the Rayleigh-scattered radiation from reaching the satellite, making the radiances insensitive  
 532 to atmospheric pressure. This characteristic knee shape allows one to estimate altitude  
 533 registration error in a manner very similar to that of RSAS, but also makes it is very susceptible  
 534 to ozone profile assumptions, as illustrated in Fig. 5.  
 535



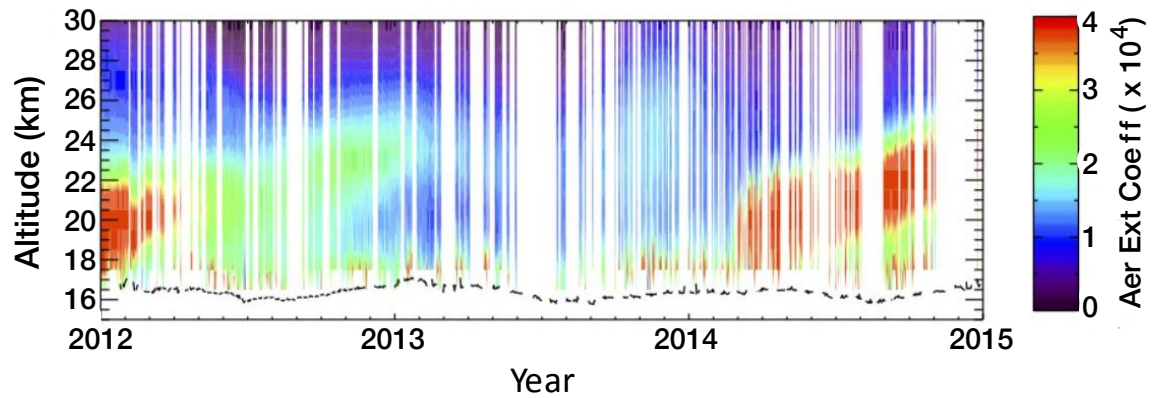
536  
 537 **Figure 5:** Typical ozone profile in the tropics (such as the one shown in the left panel) peak  
 538 between 25 and 30 km. By shifting the ozone profile we can estimate an order and pattern of  
 539 error in ozone profiles due to TH shift (right panel). Errors in ozone retrievals are within 8% at  
 540 40 km from TH errors of 300 m. Errors are least sensitive at the ozone peak, and are more  
 541 variable below.  
 542



543  
 544 **Figure 6:** OMPS LP CCD high gain earth viewing radiance images for the three slits  
 545 (east/center/west). The east and west slit images are separated in longitude by 2.25° from that  
 546 of the center slit. The wavelength range for each image is 270 to 1050 nm and the minimal TH  
 547 range is 0 to 80 km. The CCD has 740 pixels in the wavelength dimension. There are 340 pixels  
 548 in the spatial dimension; the high gain images occupy the lower half of the CCD (pixels 0 to 170).  
 549 The spatial extent of each slit's image on the detector is limited by the vertical length of that slit.  
 550 The lower slit edges (nearest the Earth surface) provide a high contrast signal cutoff that can be  
 551 monitored for movement.  
 552

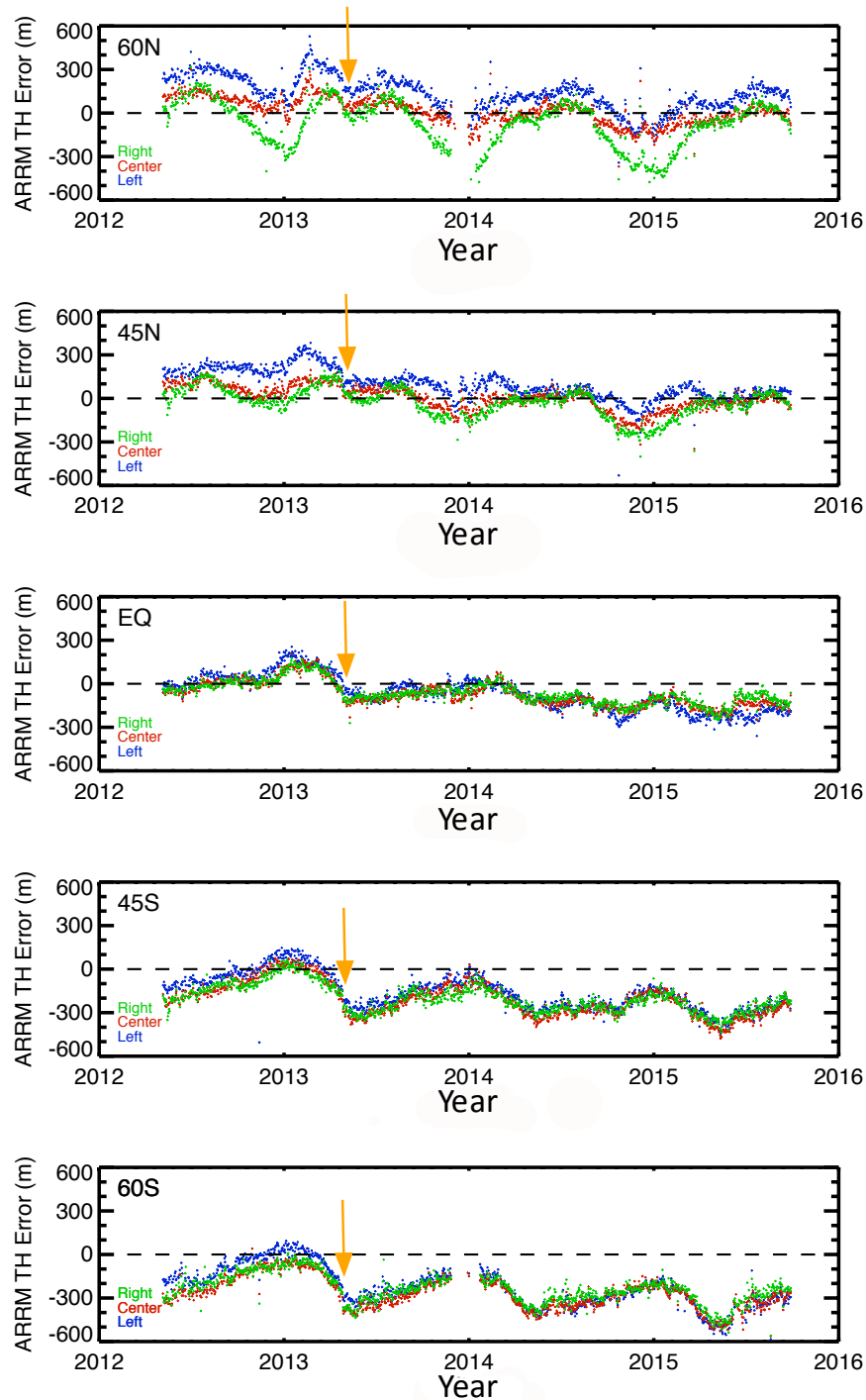


553  
 554 **Figure 7:** Slit edge results for the three slits (Green=East Slit, Red=Center Slit, Blue=West Slit)  
 555 plotted against time since Southern Terminator crossing. A 1 pixel shift corresponds to a 965 m  
 556 TH shift. The offsets are stable from the southern terminator to the mid latitude northern  
 557 hemisphere where the exposure to the sun increases thermal effects.  
 558

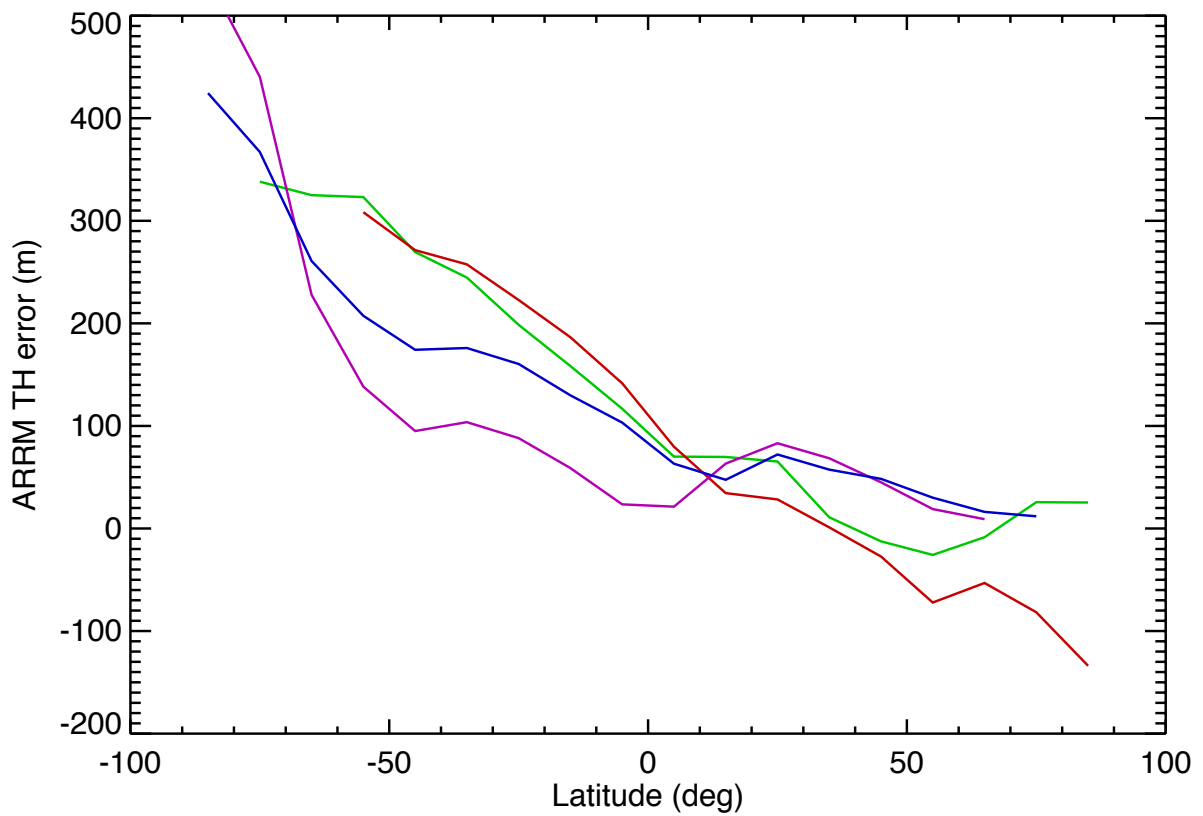


559  
560 **Figure 8:** Time series of OSIRIS aerosol extinction profiles above the tropopause (dashed line).  
561 The large concentration in 2012 is due to the June 2011 Nabro eruption in Eritrea. The aerosols  
562 at 20 km reached a minimum value (during OMPS life time) just before the eruption of the  
563 Kelud Volcano on 14 February 2014.  
564

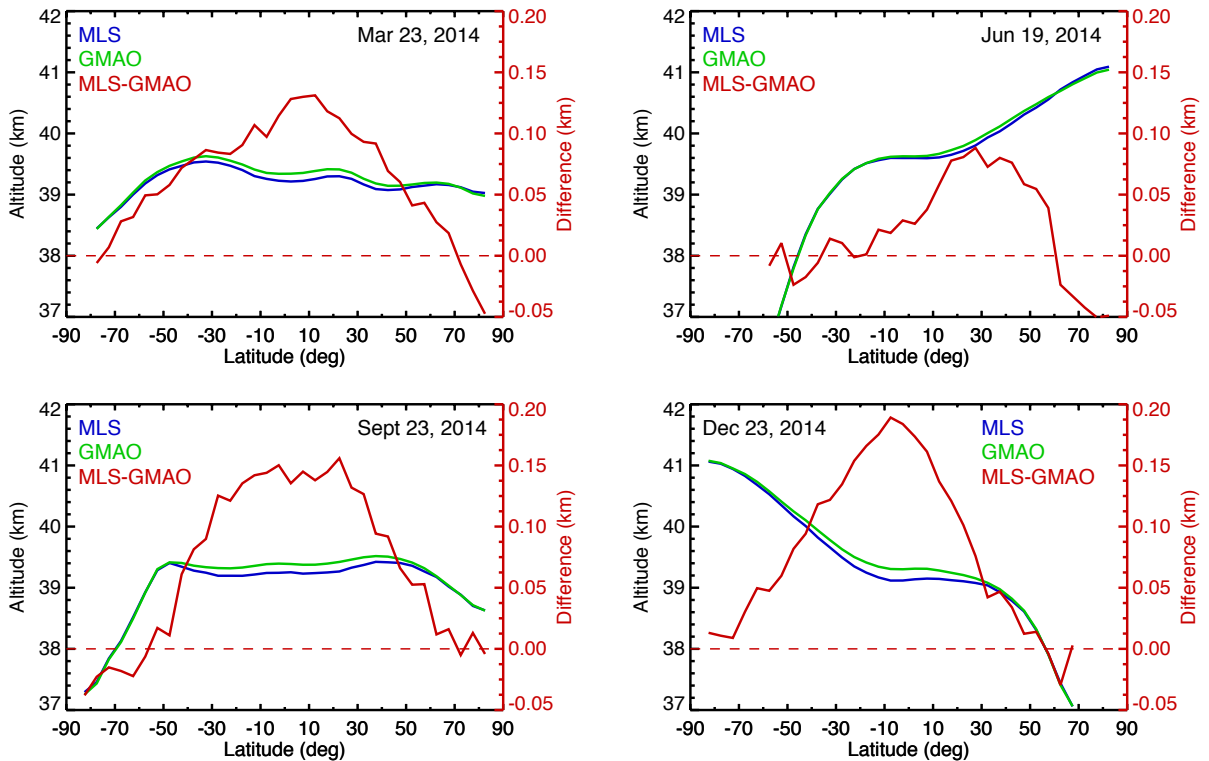




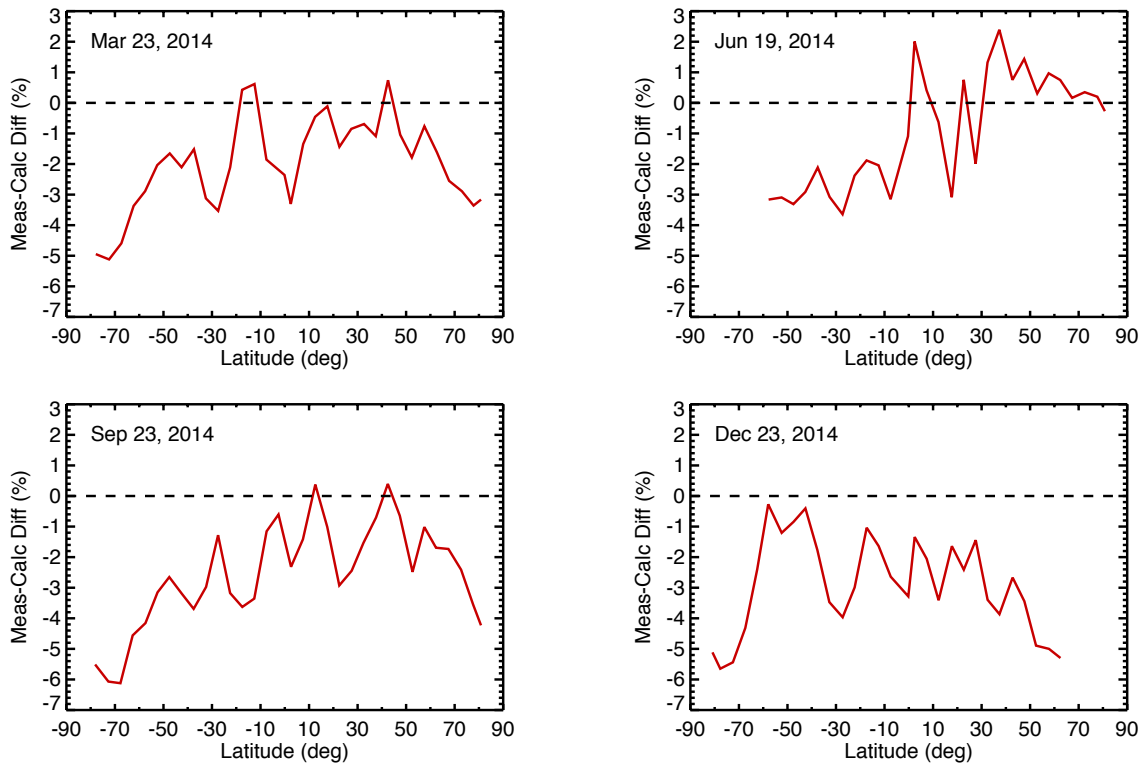
565  
 566 **Figure 9:** Time dependent plots of TH errors from ARR M analysis at five latitude bands for the  
 567 three slits. Values are normalized at the Equator just prior to the Kelud eruption on February  
 568 14, 2014 based on the RSAS results summarized in Table 1. Arrows indicate a 12 arcsec pointing  
 569 adjustment to one of the two spacecraft star trackers on April 25, 2013. The resulting 100 m TH  
 570 shift can be seen most clearly by comparing 2012 and 2013 results. Slit discrepancies and  
 571 seasonal dependencies of +/-200 m can also be seen.



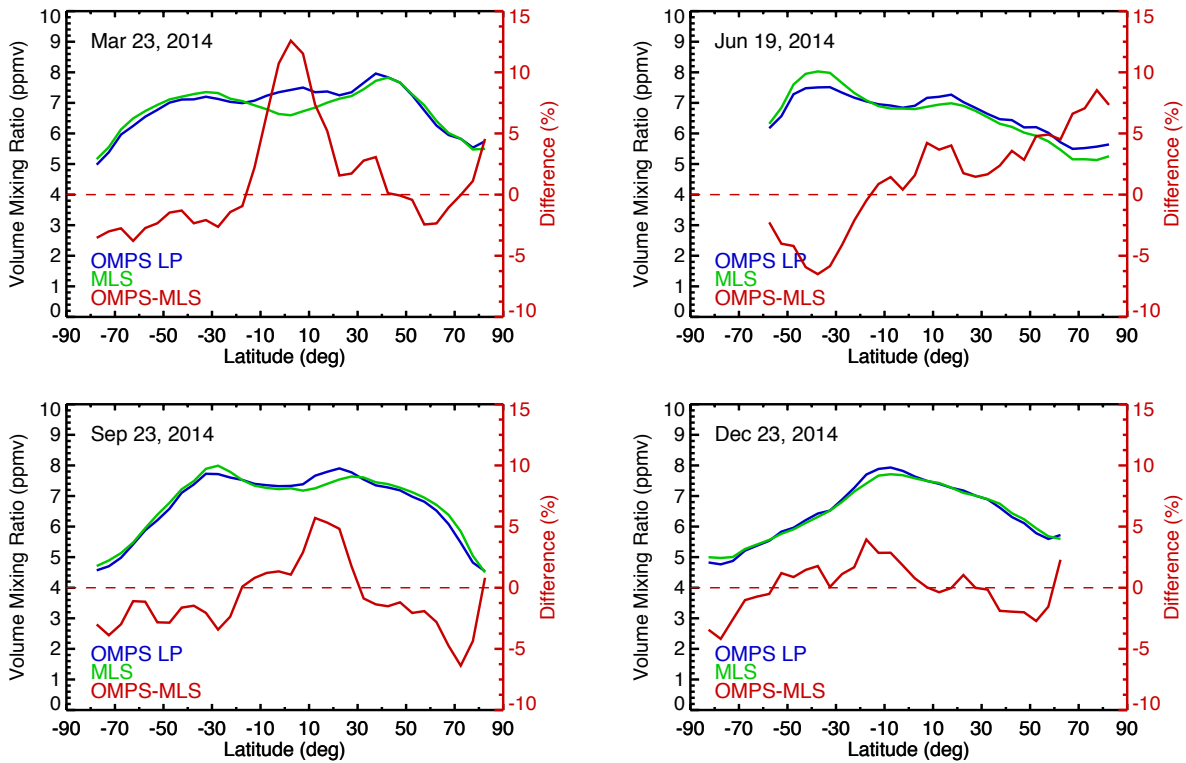
572  
 573 **Figure 10:** Average (over the ~4 year study period) ARRMM results by latitude and seasons  
 574 (MAM-green, JJA-red, SON-purple, DJF-blue) for Center Slit. Excluding the extreme polar  
 575 regions, there appears to be an average 300 m TH change over an orbit.  
 576



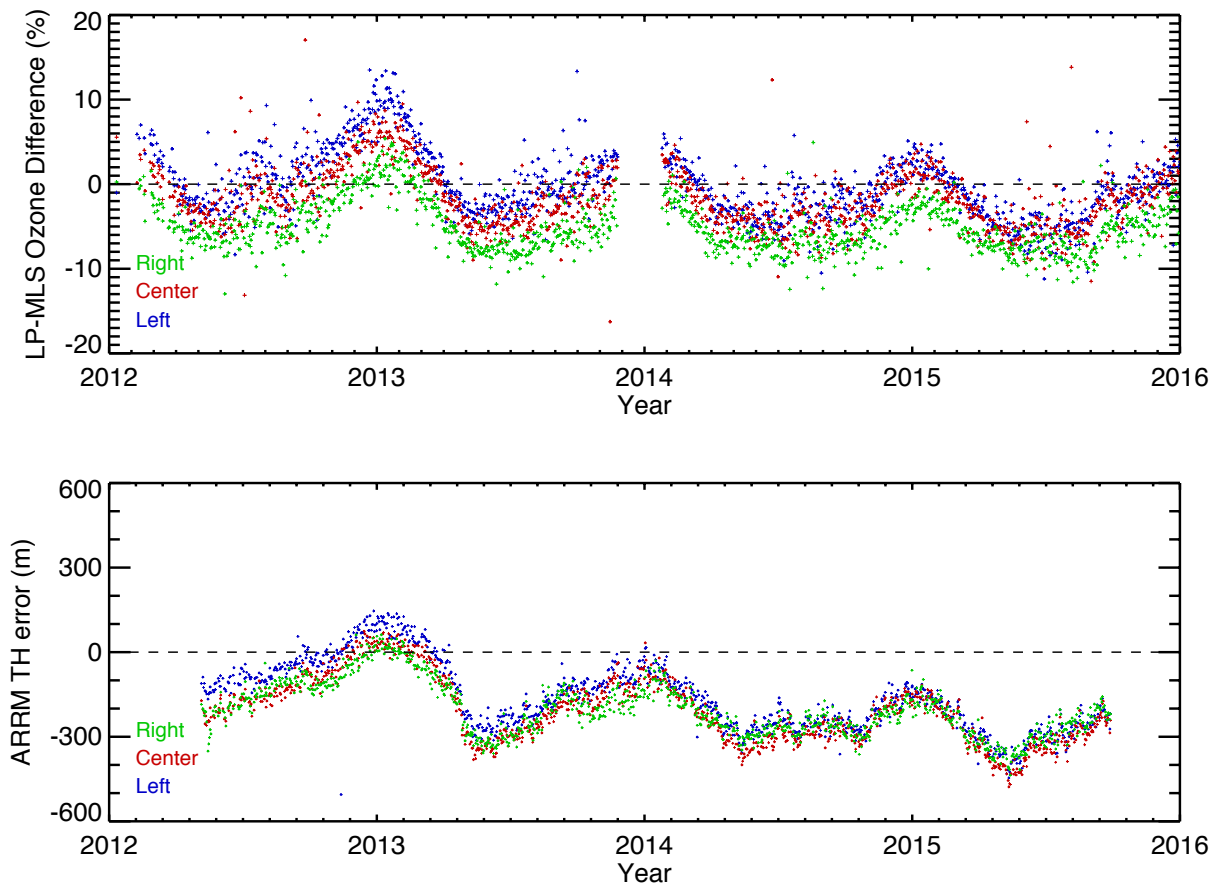
577  
 578 **Figure 11:** Daily five degree zonal means of GPH from MLS (blue), GMAO (green), and the  
 579 difference MLS-GMAO (red) at 3 hPa GPH for four cardinal days. Note that despite a 2 to 4 km  
 580 change over an orbit, the differences are generally within 200 m. These differences provide an  
 581 estimate of the errors caused by the use of MERRA GPH in our radiative transfer calculations.  
 582 Better agreement seen at the poles may simply be due to the fact that there are not many  
 583 measurements at these latitudes and both may be influenced by the same climatology. If this  
 584 error were attributed solely to the limb model and only at one altitude, the resulting TH error  
 585 would be less than +/-200 m. There is no evidence of either seasonal or latitude dependence in  
 586 the comparison, meaning that DUR effects cannot explain the variations seen in Fig. 9.  
 587



588  
 589 **Figure 12:** We have estimated the DUR modeling error by comparing 353 nm measured and  
 590 modeled radiances at 3 hPa. The radiances are modeled using an independent, nearly  
 591 simultaneous measure of surface reflectivity derived from the OMPS Nadir instrument at 340  
 592 nm. The 50x50 km nadir-view measurements are relatively insensitive to DUR effects. The  
 593 radiance differences (given for the same four cardinal days as in Fig. 11) suggest model or  
 594 calibration errors of 2-3% on average, plus structure caused by the contributing limb and nadir  
 595 scene mismatch.  
 596



597  
 598 **Figure 13:** Daily five degree zonal means of ozone from MLS (blue), GMAO (green), and the  
 599 MLS-GMAO differences (red) at 3 hPa GPH for four cardinal days. The differences are generally  
 600 within 6% which if completely attributed to TH error would be ~200 m.  
 601



602  
 603 **Figure 14:** The time series of daily zonal mean ozone differences (%) between OMPS LP and  
 604 Aura MLS for the three slits (top). OMPS LP profiles are corrected with the TH error shown in  
 605 Table 1. Note the similarity in time-dependent patterns for ozone differences and TH error  
 606 derived from AARM method (bottom, reprinted from Fig.9). The fact that these two completely  
 607 independent methods show very similar patterns give us additional confidence in the AARM  
 608 method.  
 609  
 610