



# Evidence of a significant rotational non-LTE effect in the CO<sub>2</sub> 4.3 μm PFS-MEX limb spectra

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**Abstract.** Since January 2004, the planetary Fourier spectrometer (PFS) on board the Mars Express satellite has been recording near infrared limb spectra of high quality up to  $\approx 150$  km, with potential information on density and thermal structure of the upper atmosphere. We present first results of our modeling of the PFS short wavelength channel (SWC) daytime limb spectra for the altitude region above 90 km. We applied a ro-vibrational non-LTE model based on the stellar astrophysics technique of accelerated lambda iteration (ALI) to solve the multi-species and multi-level CO<sub>2</sub> problem in the Martian atmosphere. We show that the long standing discrepancy between observed and calculated spectra in the cores and wings of 4.3 μm region is explained by the non-thermal rotational distribution of molecules in the upper vibrational states 10011 and 10012 of the CO<sub>2</sub> main isotope second hot (SH) bands above 90 km altitude. The redistribution of SH band intensities from band branch cores into their wings is caused (a) by intensive production of the CO<sub>2</sub> molecules in rotational states with  $j > 30$  due to the absorption of solar radiation in optically thin wings of 2.7 μm bands, and (b) by a short radiative life time of excited molecules, which is insufficient at altitudes above 90 km for collisions to maintain rotation of excited molecules thermalized. Implications for developing operational algorithms for massive processing of PFS and other instrument limb observations are discussed.

## 1 Introduction

For more than six Martian years the Planetary Fourier Spectrometer (PFS) on board the Mars Express satellite (Formisano et al., 2005) has been measuring the near-infrared radiances in both, nadir and limb geometries. For the study of Martian middle and upper (60-130 km) atmosphere limb measurements of the CO<sub>2</sub> 4.3 μm emission are of particular interest as the source of information on density and thermal structure of this layer. These measurements are achieved by the PFS' short wavelength channel (SWC) [1.25–5 μm] with the spectral resolution allowing to unambiguously identify many of the CO<sub>2</sub> emission bands. Details on the instrument description, its calibration and in-flight performance can be found in (Formisano et al., 2005; Giuranna et al., 2005; Formisano et al., 2006).

The strong daytime CO<sub>2</sub> 4.3 μm emission measured by PFS SWC are formed in the non-local thermodynamic equilibrium (non-LTE), which requires detailed accounting for the variety of processes influencing the ro-vibrational populations. The excitation of the CO<sub>2</sub>(ν<sub>3</sub>) vibrations is facilitated by absorption of solar radiation in the range of 1 – 2.7 μm by the CO<sub>2</sub>



molecules. This excitation is followed by the cascade radiative one-quantum  $4.3 \mu\text{m}$  transitions, which escaping the atmosphere form the daytime atmospheric emission around  $4.3 \mu\text{m}$ . In addition, the collisional quenching of molecular vibrations (V-T processes) and the inter- and intra-molecular exchange of vibrational energy (V-V processes) also strongly influence the populations.

5 A number of studies (López-Valverde et al., 2005; Formisano et al., 2006; López-Valverde et al., 2011) (hereafter “previous studies”) were undertaken in recent years aimed at interpretation of the PFS SWC limb spectra. These studies consider only vibrational non-LTE in  $\text{CO}_2$  assuming molecular rotations completely thermalized (rotational LTE), and reiterate main features of daytime  $4.3 \mu\text{m}$   $\text{CO}_2$  emission formation in the upper atmosphere, which is dominated by the second hot (SH) and third hot (TH) bands of the main  $\text{CO}_2$  isotopologue (hereafter 626). Despite the good general understanding of how the absorption of  
10 solar radiation and subsequent redistribution of vibrational energy generate the hot  $\text{CO}_2$  band emissions, significant features present in the measured PFS spectra remain unexplained. Specifically, the majority of the daytime limb radiance measurements above approximately 90 km exhibit substantially stronger emissions than predicted by the non-LTE models in the spectral ranges  $2290\text{-}2305 \text{ cm}^{-1}$  and  $2345\text{-}2355 \text{ cm}^{-1}$  relative to the maxima of radiation at  $2317 \text{ cm}^{-1}$  and  $2335 \text{ cm}^{-1}$ , respectively (López-Valverde et al., 2011). An approach to explain this feature used in the previous studies was to treat the collisional rate  
15 coefficients as free parameters. However, the desired result was not reached, and it was speculated that remaining discrepancies might be due to contributions of some unidentified weak  $\text{CO}_2$   $4.3 \mu\text{m}$  bands not yet present in the HITRAN/HITEMP and GEISA spectroscopic databases. This issue remains unresolved until now. It prevents the science community to apply suitable retrieval algorithms for obtaining parameters of the Martian middle atmosphere from the PFS SWC limb spectra around  $4.3 \mu\text{m}$ .

In this paper we address this long standing discrepancy and provide its clear physical explanation.

## 20 2 Modeling the $4.3 \mu\text{m}$ emissions

### 2.1 The ALI-ARMS code

In this paper the  $\text{CO}_2$  non-LTE populations and limb spectra are calculated using the ALI-ARMS (for Accelerated Lambda Iterations for Atmospheric Radiation and Molecular Spectra) non-LTE code package (Kutepov et al., 1998; Gusev and Kutepov, 2003; Feofilov and Kutepov, 2012). ALI-ARMS utilizes the accelerated lambda iteration (ALI) technique developed in stellar  
25 astrophysics (Werner, 1986, 1987; Rybicki and Hummer, 1991; Pauldrach et al., 1994, 2001; Pauldrach, 2003; Hubeny and Lanz, 2003) for calculating the non-LTE populations of a very large number of atomic and ionic levels in optically thick atmospheres. ALI has become a standard technique for spectrum formation calculations and for the computation of the non-LTE model stellar atmospheres. The ALI-ARMS code has been successfully applied to the diagnostics of a number of space infrared Earth’s and Martian observations, both spectrally resolved (Kaufmann et al., 2002; Kaufmann et al., 2003; Maguire  
30 et al., 2002; Gusev et al., 2006) and broadband signals (Kutepov et al., 2006; Feofilov et al., 2009, 2012; Rezac et al., 2015), as well as to study the infra-red radiative cooling/heating in planetary atmospheres (Hartogh et al., 2005; Kutepov et al., 2007; Feofilov and Kutepov, 2012).



## 2.2 Rotational LTE and non-LTE

In order to establish a nominal result for a pure vibrational non-LTE model (complete rotational LTE is assumed) as used in previous studies we run ALI-ARMS for our extended reference CO<sub>2</sub> model for a dust-free atmospheric model retrieved from the Mars Climate Database (hereafter MCD)<sup>1</sup>, which matches the observation conditions. This reference model accounts for 150 vibrational levels of 5 CO<sub>2</sub> isotopologues and about 40k lines in calculating the non-LTE populations. The radiative transfer and solar absorption terms are calculated for all the 376 bands present in the task. We apply the same collisional rate coefficients for the lower vibrational levels as those used in previous studies, however, a different scaling of these basic rates for higher vibrational levels is performed based on the first-order perturbation theory as suggested by Shved et al. (1998).

Once the CO<sub>2</sub> populations are known we simulate the limb spectra in the range 2200 - 2400 cm<sup>-1</sup> for different tangent altitudes both with and without accounting for the overlapping, and compare them to the PFS observations. In our model we include in total 60 bands of the 5 isotopologues that contribute to the CO<sub>2</sub> 4.3 μm limb emission.

A random sample of the PFS measured spectrum is shown as a black curve in Fig.1. This and other spectra have been kindly provided for this study by the PFS PI M. Giuranna. The shown spectrum is an average of 27 single scans taken for the following conditions:  $L_s = 254$ , latitude = -67°, local time = 18 h, SZA = 69°, and tangent height of 115 km. As we already noted the effect addressed here is present in the majority of daytime spectra, therefore, no special criteria in latitude/longitude or local-time were applied to select these particular spectra for comparison with calculations here. The (turquoise) curve in Fig.1 presents the spectrum for our nominal run. We show here the synthetic radiance for tangent height of 120 km since for a given input atmospheric model it provides better agreement with measured spectrum at 115 km.<sup>2</sup> Comparison in Fig.1 of simulated and measured signals obviously confirms the general conclusion of previous studies, i.e. above the altitude of 90 km, no matter the location and local time, *taking into account all known weak bands of CO<sub>2</sub> does not improve the mismatch* with the observed radiance in the “shoulders” (2290-2305 cm<sup>-1</sup> and 2345-2355 cm<sup>-1</sup>) relative to the maximum at 2317 cm<sup>-1</sup> and 2335 cm<sup>-1</sup>.

The problem of “incorrect” radiance distribution between cores and wings of molecular bands is not a new one and has been theoretically investigated in series of papers by Kutepov et al. (1985); Ogibalov and Kutepov (1989); Kutepov et al. (1991) for CO<sub>2</sub> bands for atmospheric conditions of Mars and Venus, and for CO bands by Kutepov et al. (1997) for the Earth’s atmosphere. In these studies it has been shown that in Martian atmosphere the Boltzmann rotational distribution for CO<sub>2</sub> molecules in the  $\nu_3$ -vibrationally excited states brakes down for altitudes above 80 km. At these altitudes the lifetime  $\tau = 1/A$  of these vibrational levels, where  $A \sim 200 \text{ s}^{-1}$  is the Einstein coefficient for the spontaneous 4.3 μm transitions, is not long enough for collisions to keep the rotation of excited molecules thermalized. Therefore, the need for complete vibrational-rotational non-LTE consideration has been identified already at that time.

<sup>1</sup><http://www-mars.lmd.jussieu.fr/mars/access.html>

<sup>2</sup>The mismatch of emission absolute values between measured and simulated spectra for the tangent altitude of 115 km (see also orange curve in Fig.1, discussed below) may be attributed both to the mismatch of input model and to the simplifying assumptions applied in this and previous studies, namely, the missing field-of-view averaging.



In this study we relax the assumption of rotational LTE for two vibrational levels 10011 and 10012 of main CO<sub>2</sub> isotope 626. These levels are strongly pumped by the absorption of solar radiation at 2.7 μm and give origin to the two strong SH bands, which dominate the 4.3 μm limb emissions (Gilli et al., 2009; López-Valverde et al., 2011) of the Martian middle and upper atmosphere discussed here. We accounted for rotational sub-levels up to  $j = 102$  for each of these vibrational levels. The state-to-state rate coefficients describing rotational relaxation of vibrationally excited CO<sub>2</sub> were calculated as described by Kutev et al. (1985) who applied the model of Preston and Pack (1977) and Pack (1979) based on infinite order sudden (IOS) approximation results. The total rotational relaxation rate used in these calculations was estimated as  $k_{rot}(T) = 2\pi k(T/p)\Delta\nu_L(T_o)(T/T_o)^{0.7}$ , where  $\Delta\nu_L$  is the Lorentz line width. The validity of this approach is discussed in detail for linear and non-linear molecules by Hartmann et al. (2008). For our model we used mean Lorentz width  $\Delta\nu_L$  of CO<sub>2</sub> rotational-vibrational lines from HITRAN12 (Rothman et al., 2013). With accounting for temperature dependence of these widths  $k_{rot} \approx 3 \times 10^{-10} \text{ cm}^3\text{sec}^{-1}$  for  $T = 200 \text{ K}$ .

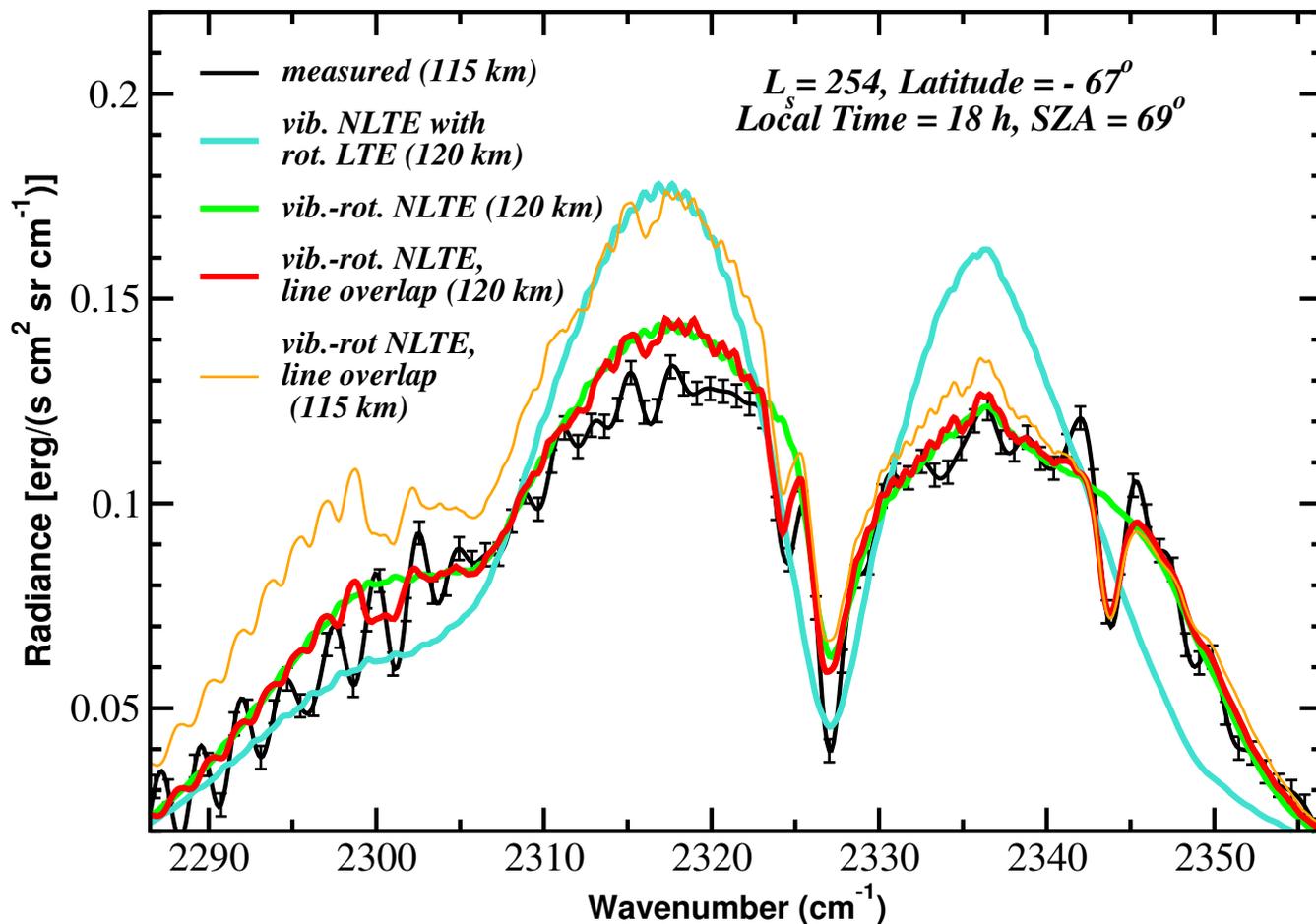
### 3 Results and Discussion

#### 3.1 Importance of rotational non-LTE

The green curve in Fig.1 shows the spectrum calculated with accounting for the rotational non-LTE at vibrational levels 10011 and 10012 of main CO<sub>2</sub> isotope. One may see that this spectrum shows the redistribution of the radiances from SH band branch cores into there wings and, as a result, reproduces the measured signal shape significantly better than calculations based on the assumption of rotational LTE (turquoise curve). The reason for this alternation of the band shape is demonstrated in Fig.2. In the left panel we show for the altitude of 100 km the production rate (normalized over the rotational quantum number  $j$ ) of excited 626 molecules in vibrational state 10011 due to the absorption of the 2.7 μm solar radiation for two different SZA. Similar production rate was also obtained for level 10012. We investigated these production rates as well as corresponding rotational distributions (right panel) over the altitude range of 80-140 km. In general the molecules in level 10011 are most efficiently generated in rotational states with  $j = 30 - 50$ . The low production for  $j < 30$  is caused by a strong absorption of solar radiance in the cores of the 2.7 μm band branches in the upper atmosphere, which, however, remain transparent in the branch wings. At the altitudes above 90 km the collisions are not able to completely restore the Boltzmann rotational distribution of these vibrational levels, and the latter takes form shown in the right panel of Fig.2. Compared to the Boltzmann distribution (green) both rotational non-LTE curves (red and blue) demonstrate enhanced tails for  $j > 30$ , which reflects intensive radiative pumping of these levels shown in the left panel.

#### 3.2 Importance of line overlapping along the limb LOS

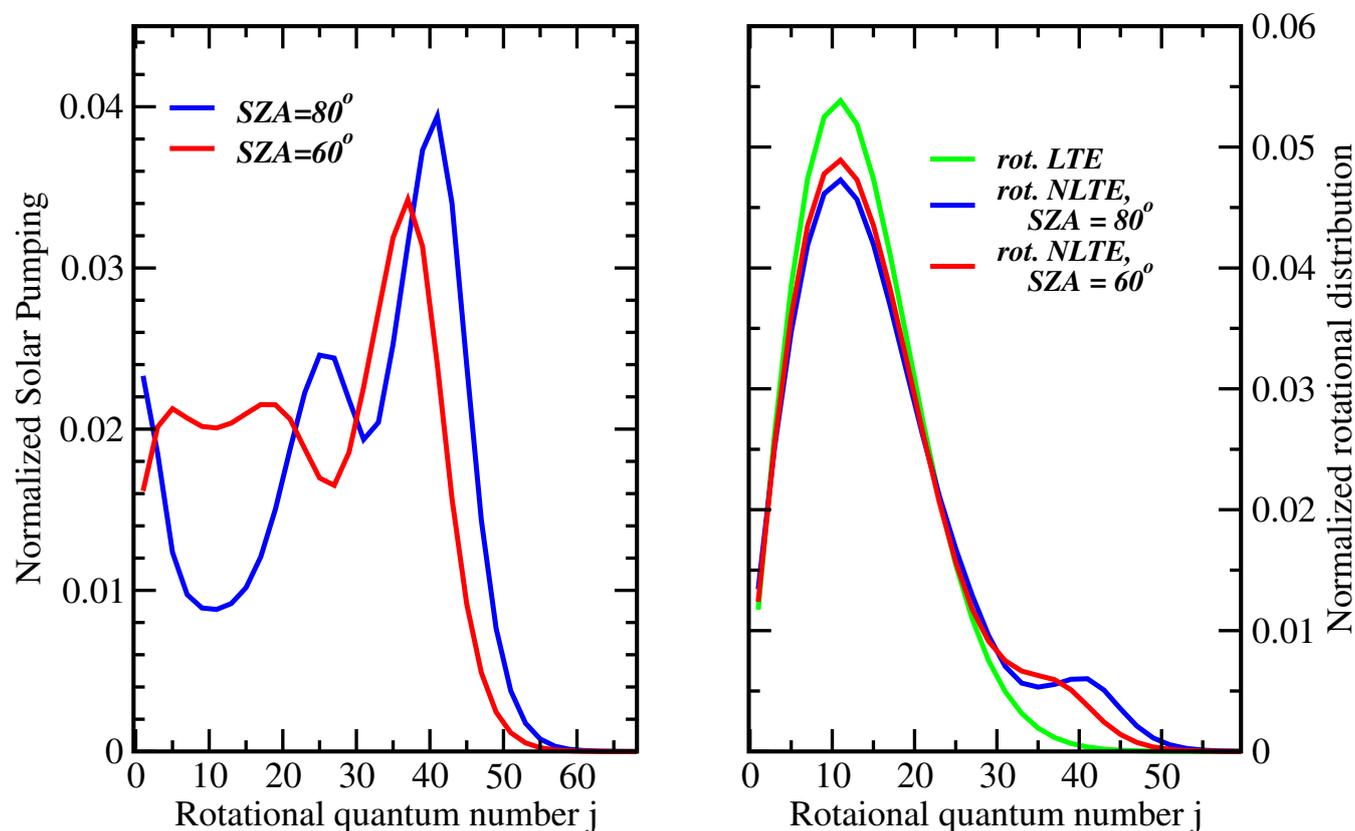
The rotational non-LTE spectrum in Fig.1 (green curve) significantly better reproduces general shape of measured signal (black curve) compared to that calculated with the rotational LTE assumption (turquoise curve). However, there are a number of small and medium size spectral features visible in the measured spectrum not present in the simulated one (green curve). So far



**Figure 1.** Comparison of measured (black) and simulated PFS SWC spectra. The observed spectrum is an average of 27 single scans taken for the following conditions:  $L_s = 254$ , latitude =  $-67^\circ$ , local time = 18 h, SZA =  $69^\circ$ , and tangent height of 115 km. The one sigma random uncertainties for this measured spectrum are shown. The simulated spectra are shown for tangent height of 120 km including only vibrational non-LTE in  $\text{CO}_2$  (turquoise), with vibrational-rotational non-LTE (green), and with vibrational-rotational non-LTE plus line overlapping along the line-of-sight (red). The spectra shown in orange color are calculated with same assumptions as for red, but for tangent height 115 km as a demonstration. (see text for detailed description)

the synthetic spectra discussed here were calculated ignoring line overlapping along the LOS.<sup>3</sup> In this study we employed our radiance model in both, non-overlapping, and overlapping mode. In the latter case the calculations treat spectral line overlapping of all bands (within a band and between lines of different bands) of all isotopes in the line-by-line (LBL) fashion (technical discussion is detailed in Rezac et al. (2015)). The effect of accounting for line overlapping is clearly demonstrated

<sup>3</sup> We note here also that although, previous studies report using forward model capable of treating line overlapping in limb calculations, its effects on PFS spectra were not discussed.



**Figure 2.** Left - normalized distribution of radiative pumping at level 10011 of  $^{12}\text{C}^{16}\text{O}_2$  due to the absorption of the  $2.7 \mu\text{m}$  solar radiation, altitude 100 km. Right - normalized rotational distribution at level 10011 of  $^{12}\text{C}^{16}\text{O}_2$ , altitude 100 km

in spectrum plotted in red color. Two well developed absorption features appeared, one around  $2344 \text{ cm}^{-1}$  and another one around  $2324 \text{ cm}^{-1}$ , which overlay well the corresponding features of measured spectrum. We found that first of these features is caused by the absorption of the 10012 - 10002 SH band line R22e (here and below the line notations are taken from HITRAN) emission of main isotope with the nearly coincidental line R16e, which belongs to the fundamental band of isotope 628. Same is also true for the feature at  $2324 \text{ cm}^{-1}$ : here the emission line P4e of the 626 SH band 10012 - 10002 is blanketed by the line P15e of the first hot band 01111 - 01101 of same isotope.

In addition, the “rotational non-LTE + overlapping on limb” spectrum (red curve) has an absorption feature around  $2316 \text{ cm}^{-1}$ , which is, however, substantially less developed versus corresponding feature in measured spectrum. Compared to two other signatures discussed above the mechanism of this signature formation is significantly more complex. Here mainly the emission in the line R9e from the 10011 - 10001 SH band of isotope 628, which is pumped by the  $2.7 \mu\text{m}$  solar radiation, is attenuated by two close lines Q31f and Q74f of 628 and 626 FH bands, respectively, although both emission and absorption in a few other very closely located lines, which belong to various isotopic bands, may also contribute to the formation of



this signature. We note, however, that none of vibrational levels of bands, whose lines are involved in the formation of this absorption feature, were treated in this study with accounting for the rotational non-LTE. In general, rotational non-LTE at lower vibrational level can, however, influence the band absorption, whereas that at the upper level will impact the band emission. To get a better agreement with measurements both, for this spectral feature and for the region  $2285\text{-}2305\text{ cm}^{-1}$ , where the synthetic spectrum does not reproduce a number of fine “wavy” features, we plan (a) to apply more detailed rotational non-LTE model, and (b) to use in synthetic radiance calculations some elements of technical signal processing, among them “zeropadding” applied to the PFS interferograms, which contributes (Giuranna et al., 2005) to the formation of these “wavy” features of measured spectra.

#### 4 Conclusions

We present first our modeling results of the Mars Express PFS SWC daytime limb  $4.3\text{ }\mu\text{m}$  spectra for the altitude region above 90 km. We show that the long standing discrepancy between observed and calculated spectra in the cores and wings of the  $4.3\text{ }\mu\text{m}$  region is explained by the non-thermal rotational distribution of the  $^{12}\text{C}^{16}\text{O}_2$  molecules in upper vibrational levels 10011 and 10012 of the strong SH bands. The enhancement of the SH band wing emissions is caused (a) by intensive production of the  $\text{CO}_2$  molecules in rotational states with  $j > 30$  due to the absorption of solar radiation in optically thin wings of  $2.7\text{ }\mu\text{m}$  bands, and (b) by a short radiative life time of molecules in these states, which is insufficient at altitudes above 90 km for collisions to maintain rotation of excited molecules thermalized. As a result redistribution of SH band intensities takes place going from band branch cores into their wings. This result confirms significant impact of rotational non-LTE on the  $\text{CO}_2$   $4.3\text{ }\mu\text{m}$  emissions of Martian and Venusian atmospheres, which was predicted by Kutepov et al. (1985); Ogibalov and Kutepov (1989). *The PFS spectra provide the first strong evidence of this effect.*

Although, the rotational non-LTE on the levels 10011, 10012 does excellent job explaining the radiance redistribution observed in the measured spectra, more detailed simulations are still needed to investigate if there are smaller order effects caused by the rotational non-LTE at other  $\text{CO}_2$  vibrational levels. However, accounting for rotational non-LTE significantly slows down the non-LTE calculations, by an order of magnitude for the problem discussed in this paper.

Additional improvement in matching the  $4.3\text{ }\mu\text{m}$  band shape/structure between simulated and PFS measured spectra was reached by accounting for spectral line overlapping (within each band and among lines of different bands) along the limb LOS. This allowed to reproduce some fine absorption features in measured spectra that were missing in previous calculations, caused, however, factor of 10 increase of the limb emission computing time. Together with the time increase caused by accounting for the rotational non-LTE this indicates the need of significant additional study aimed at optimizing the calculations to make the computational time acceptable for massive PFS data processing.

The presented results are also important for diagnostics of other similar observations. For instance, the Venus Express VIRTIS spectra around  $4.3\text{ }\mu\text{m}$  for tangent heights above 100-110 km, which are discussed in Gilli et al. (2009); López-Valverde et al. (2011), demonstrate remarkable resemblance of Mars Express PFS spectra analyzed in this study. Due to a well known similarity of the  $\text{CO}_2$  daytime emission formation mechanisms for both planets (Ogibalov and Kutepov, 1989) the



Venus spectra obviously require also detailed rotational non-LTE/line overlapping analysis. These results may also be relevant for the analysis of limb mode measurements of the Trace Gas Orbiter/NOMAD instrument, which will start science operation in late 2017.

At last, we note that, as in previous studies, the following simplifications were applied in our simulations, (a) infinite-narrow FOV, and (b) detailed spectra convolution was not performed, only simple ( $1 \text{ cm}^{-1}$ )-triangle window (to match the instrument spectral resolution) averaging of monochromatic spectra calculated on a fine grid was applied. Neither of these, however, influence main results of this study. Nevertheless, missing FOV averaging (in combination with in-sufficient matching of applied pressure/temperature model) may explain better agreement (see Fig.1) between absolute emission values of spectrum measured for the tangent altitude of 115 km and that simulated for 120 km (compare black, red and thin orange line in this figure). Additionally, simplified averaging of monochromatic spectra in combination with missing FOV averaging may cause certain fine structure inconsistencies of measured and simulated spectra.

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