



# 1 Improved OSIRIS NO<sub>2</sub> retrieval algorithm: Description and 2 validation

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15

## 16 **Abstract**

17 A new retrieval algorithm for OSIRIS (Optical Spectrograph and Infrared Imager System)  
18 nitrogen dioxide (NO<sub>2</sub>) profiles is described and validated. The algorithm relies on spectral  
19 fitting to obtain line-of-sight (LOS) column densities of NO<sub>2</sub> followed by inversion using an  
20 algebraic reconstruction technique and the SaskTran spherical radiative transfer model to obtain  
21 vertical profiles of local number density. The validation covers different latitudes (tropical to  
22 polar), years (2002-2012), all seasons (winter, spring, summer, and autumn), different  
23 concentrations of nitrogen dioxide (from deNOxified polar vortex to polar summer), a large  
24 range of solar zenith angles (68.6 to 90.5°) and altitudes between 10.5 and 39 km, thereby  
25 covering the full retrieval range of a typical OSIRIS NO<sub>2</sub> profile. The use of a larger spectral  
26 fitting window than used in previous retrievals reduces retrieval uncertainties and the scatter in  
27 the retrieved profiles due to noisy radiances. Improvements are also demonstrated through the



1 validation in terms of bias reduction at 15-17 km relative to the OSIRIS operational v3.0  
2 algorithm. By accounting for the diurnal variation along the LOS in the two-dimensional  
3 radiative transfer model, the scatter of the differences relative to the correlative balloon NO<sub>2</sub>  
4 profile data is reduced.

## 5 **1. Introduction**

6 Nitrogen oxides, such as NO and NO<sub>2</sub> are the reactive nitrogen-containing species in the middle  
7 atmosphere and are produced mainly from the breakdown of nitrous oxide in the stratosphere.  
8 Oxides of nitrogen dominate ozone loss in the middle stratosphere, whereas in the lower  
9 stratosphere, they react with oxides of chlorine and bromine, such as ClO and BrO, to reduce the  
10 halogen-catalyzed destruction of ozone.

11 The partitioning between NO and NO<sub>2</sub> depends on several factors such as the local ozone  
12 concentration and the photolysis frequency of NO<sub>2</sub>. Reactive nitrogen is chemically converted at  
13 night to N<sub>2</sub>O<sub>5</sub> and, upon hydrolysis, can be further sequestered into unreactive ‘reservoir’ species  
14 such as HNO<sub>3</sub>. NO<sub>2</sub> increases steadily during daylight hours due to the UV photolysis of N<sub>2</sub>O<sub>5</sub>  
15 (e.g. Wetzel et al., 2012).

16 The photochemistry of NO<sub>2</sub>, which is particularly rapid near the day-night terminator, leads to  
17 spatial gradients within the field of view, particularly for limb sounders such as OSIRIS (Optical  
18 Spectrograph and Infrared Imager System) (Llewellyn et al., 2004), or for solar occultation  
19 instruments operating in either the UV-visible or mid-infrared (e.g., Kerzenmacher et al., 2008).

20 The OSIRIS operational NO<sub>2</sub> retrieval algorithm was developed and validated by Haley and  
21 Brohede (2007) and the current version is 3.0. The work of Kerzenmacher et al. (2008) is the  
22 only other publication comparing version 3.0 OSIRIS NO<sub>2</sub> data to correlative profile  
23 measurements. Earlier versions (e.g., 2.x, Haley et al., 2004) were more thoroughly validated, for  
24 example by Brohede et al. (2007). The pseudo-spherical forward model used in the operational  
25 OSIRIS NO<sub>2</sub> retrieval algorithm is less accurate than the SaskTran spherical radiative transfer  
26 model (RTM, Bourassa et al., 2008, Zawada et al., 2015), currently used as a forward model in  
27 the operational retrieval algorithm for OSIRIS ozone and aerosol extinction data products.  
28 Recently, Bourassa et al. (2011) developed an alternative NO<sub>2</sub> algorithm which relied on 4  
29 wavelengths covering a single NO<sub>2</sub> absorption band. This was later modified to 13 wavelengths



1 covering three adjacent NO<sub>2</sub> bands (438–450 nm) and used to process the entire OSIRIS data  
2 record and is referred to here as the ‘fast’ OSIRIS NO<sub>2</sub> product. While the spectral information  
3 content is reduced relative to the operational OSIRIS algorithm and the algorithm described  
4 below (see Table 1), the ‘fast’ algorithm has two key common elements to the one described  
5 herein:

- 6 1) The forward model is the successive-orders-of-scattering version of SaskTran.
- 7 2) MART (Multiplicative Algebraic Reconstruction Technique) (Degenstein et al., 2003) is  
8 used for inversion.

9 In this work, we provide a detailed description of the new retrieval algorithm whose heritage is  
10 the ‘fast’ algorithm (Bourassa et al., 2011) as well as the algorithm developed in a series of  
11 papers (Sioris et al., 2003; Sioris et al., 2004; Sioris et al., 2007). The current algorithm was  
12 developed to demonstrate that improved accuracy is possible through the combination of a better  
13 forward model, better forward model inputs, and additional, longer wavelengths, with a  
14 particular focus on the lower stratosphere (and upper troposphere).

15 We then compare NO<sub>2</sub> profiles retrieved from OSIRIS observations to balloon-borne NO<sub>2</sub> profile  
16 measurements during the decade (2002–2012) when several balloon-borne limb measurements  
17 were performed. Balloon profiles are chosen for the validation for many reasons, of which the  
18 most important is the expected accuracy of this data. The accuracy is due to two main factors:  
19 very high signal-to-noise afforded by the luxury of long exposure times (which can be traded for  
20 higher vertical and/or spectral resolution) and superior altitude determination, which for balloon-  
21 borne limb geometry is due to the sensor being an order of magnitude closer to the tangent point  
22 than for satellite limb sounders, thereby reducing the impact of imperfect viewing angle  
23 knowledge. Profile measurements from balloons are preferable to those from satellite for the  
24 purpose of validation because of their high vertical resolution, generally matching or exceeding  
25 the ~2 km vertical resolution of OSIRIS NO<sub>2</sub> (e.g. Sioris et al., 2003) (see below). Balloon  
26 measurements exploit a greater diversity of methods as in-situ techniques are possible in addition  
27 to remote sensing. Furthermore, one balloon-borne remote sensing technique relies on  
28 occultation during balloon ascent/descent, which is not possible for satellite instruments, and  
29 provides very accurate altitude registration, offers potentially finer vertical resolution than limb



1 occultation for instruments observing the full solar disk (see Sect 2.3) and smaller errors due to  
2 the neglect of diurnal NO<sub>2</sub> gradients.

## 3 2. Method

### 4 2.1 Algorithm settings

5 The algorithm is a classic 2-step approach of spectral fitting using absorption cross-sections to  
6 determine slant column densities (SCDs) and vertical fitting to invert the SCDs. This approach is  
7 used by many groups including some of the balloon remote sensing teams providing data used in  
8 this study. The reference spectrum is the co-addition of spectra at tangent heights (THs) in the  
9 50-70 km range (Sioris et al., 2003). Spectral fitting refers to a multiple linear regression  
10 including the following basis functions: a fourth order closure polynomial which is justified  
11 based on an adjusted-R<sup>2</sup> test, and temperature-dependent NO<sub>2</sub> (Vandaele et al., 1998) and O<sub>3</sub>  
12 (Serdyuchenkov et al., 2014) absolute absorption cross sections interpolated to the temperature  
13 (*T*) of the tangent layer using the European Centre for Medium-Range Weather Forecasting  
14 (ECMWF) analysis and convolved with a Gaussian to OSIRIS spectral resolution. Water vapour  
15 absorption is neglected despite maximal absorptions of >0.1% in the fitting window (434.8-476.7  
16 nm) in the upper troposphere as it is not spectrally correlated with NO<sub>2</sub> absorption over this  
17 window and would greatly increase the forward modelling computational burden. The spectral  
18 fitting is exactly the same for observed and simulated normalized radiances. For example, the  
19 actual OSIRIS wavelengths and tangent heights are inputs into the simulation and any spectral  
20 pixels which are rejected due to radiation hits or detector saturation at the Level 1 processing  
21 stage (i.e. observed radiances) are also omitted from the fit of the SaskTran-simulated radiances.  
22 The longest wavelength in the fitting window is extended to 476.7 nm, raising the number of  
23 spectral pixels to 107, thereby increasing the spectrally-integrated signal-to-noise ratio and the  
24 penetration of the lower atmosphere relative to both existing OSIRIS NO<sub>2</sub> algorithms mentioned  
25 above. Both of these benefits of an extended fitting window had been demonstrated for  
26 SCIAMACHY limb scattering (Sioris et al., 2004) but OSIRIS has a glass filter that prevents the  
27 detection of higher orders of light reflected off the grating that was positioned such that the 477-  
28 530 nm region is not usable (Warshaw et al., 1996). Thus the 434.8-476.7 nm window is used  
29 here and spectral fitting residuals are shown in Fig. 1. Using a discontinuous fitting window that  
30 included wavelengths greater than 530 nm is not beneficial. The NO<sub>2</sub> SCD uncertainties are



1 improved relative to those obtained by Sioris et al. (2003) who used a 434.8–449.0 nm fitting  
2 window and the ‘tilt’ pseudo-absorber. The improved NO<sub>2</sub> number density precision is illustrated  
3 below. The ‘tilt’ basis function is now excluded from the spectral fitting. We also tested an  
4 alternative approach of fitting an NO<sub>2</sub> spectral weighting function to the normalized radiances, as  
5 is used to fit the ozone absorption signal in nadir reflectance spectra (Coldewey-Egbers et al.,  
6 2005) and the spectral fits did not improve significantly.

7 The retrieval upper altitude limit is defined by the lowest tangent height in a limb scan for which  
8 the NO<sub>2</sub> SCD error is <100% for all tangent heights below. This can be as high as 49 km, but is  
9 typically ~40 km and typically a few kilometres higher than the upper altitude of the ‘fast’ NO<sub>2</sub>  
10 product. The lower limit of the retrieval is often determined by cloud tops as Odin generally  
11 scans the limb into the upper troposphere. Cloud-contaminated observations are excluded as they  
12 can lead to biases and difficulty in retrieval convergence. Cloud top detection is particularly  
13 important for MART since spectra from successive tangent heights immediately below a given  
14 altitude are used in a weighted fashion to retrieve the NO<sub>2</sub> number density at that altitude. We  
15 assume, for OSIRIS geometry (scattering angles of 90±30°), that an observed scene near a cloud  
16 top sampled by the ~1 km tall instantaneous-field-of-view has larger limb radiance at ~810 nm  
17 than if it were cloud-free. A scene is deemed to be cloudy if the ~810 nm radiance scale height is  
18 < 2.4139. This threshold is lowered from an overly stringent value of 3.84 (Sioris et al., 2007).  
19 The new threshold is chosen to allow NO<sub>2</sub> retrievals (and validation) to extend below 2–3 month  
20 old volcanic aerosol layers due to Sarychev Peak and Nabro, two of the eruptions during the  
21 OSIRIS mission which led to the largest stratospheric aerosol optical depth.

22 Algebraic reconstruction techniques (Fesen and Hays, 1982) have been used to recover the  
23 vertical and along-track distribution of atmospheric constituents for over four decades (Thomas  
24 and Donahue, 1972). Chahine’s (1968) relaxation method, used by Sioris et al. (2007) for the  
25 retrieval of OSIRIS NO<sub>2</sub> vertical profiles, is a variant of MART in which only the tangent layer  
26 is used to retrieve the local number density and is also tested (see Sect. 4). The MART retrieval  
27 in this work uses a 0.6:0.3:0.1 weighting following the OSIRIS aerosol extinction retrieval (see  
28 Eq. 8 of Bourassa et al., 2007). Exactly 15 retrieval iterations are used, which appears to be  
29 adequate for most cases and, with MART, does not lead to overfitting. Five orders of scattering  
30 are used (Sioris et al., 2004). The retrieval uses a 1 km altitude grid, although the NO<sub>2</sub> vertical



1 resolution remains ~2 km, a consequence of the ~2 km vertical sampling provided by Odin.  
2 Above the retrieval range, NO<sub>2</sub> profiles are scaled every iteration using a Chahine-like update  
3 based on the highest tangent height within the retrieval range. Below the retrieval range, the NO<sub>2</sub>  
4 profile is assumed to have a constant number density down to the ground, equal to the number  
5 density at the lowest retrieved altitude. The air number density profile is from ECMWF.  
6 The retrieval uses OSIRIS-retrieved ozone and aerosol extinction profiles (version 5.07, v5.07  
7 hereafter) and the 675 nm scene albedo (Bourassa et al., 2007) to provide more realistic forward  
8 model inputs into SaskTran.  
9 The NO<sub>2</sub> retrieval uncertainty is obtained by perturbation as described by Sioris et al. (2010),  
10 with the NO<sub>2</sub> SCD standard errors and an altitude-independent number density perturbation  
11 serving as inputs. The 1D profile retrieval including the error calculation takes ~5 minutes on a  
12 desktop computer with eight 3.4 GHz processors.

## 13 **2.2. Modelling diurnal gradients within SaskTran**

14 Zawada et al. (2015) describe the high spatial resolution capability which is required for  
15 modelling (horizontal) diurnal gradients of NO<sub>2</sub> within SaskTran. SaskTran is capable of  
16 modelling radiation fields with the atmosphere varying along the line-of-sight (LOS), and also  
17 along the incoming solar beam (referred to as ‘2D mode’ and ‘3D mode’, respectively,  
18 hereafter). The 1D forward model and associated retrieval (Table 1) completely neglects diurnal  
19 gradients in NO<sub>2</sub>. In 2D mode, the atmosphere consists of sectors along the LOS but the diurnal  
20 variation of NO<sub>2</sub> along the incoming solar beam is not considered. In 3D mode, the atmosphere  
21 essentially consists of stacked triangular prisms of increasing horizontal extent with increasing  
22 distance from the tangent point as illustrated by Zawada et al. (2015) and diurnal gradients are  
23 simulated for any light path from its point of entry into the atmosphere to its exit.  
24 To provide realistic diurnal gradients in the atmosphere of the forward model, SaskTran has been  
25 linked to the PRATMO stratospheric gas-phase photochemical box model (McLinden et al.,  
26 2000). PRATMO is configured to converge to 0.5% between the start and end of each one day  
27 run. The atmosphere is aerosol-free for photolysis frequency calculations. The latitudinal and  
28 vertical resolutions are 5° and ~2 km, respectively. The default number of time steps per day is  
29 35. More details are available in McLinden et al. (2006) and references therein.



### 2.3 Validation approach and datasets

Even though the OSIRIS 2D and 3D profiles are retrieved accounting for the diurnal gradients expected for the OSIRIS viewing geometry, photochemical modelling is also required to scale all of the OSIRIS NO<sub>2</sub> profiles to the local time of the balloon measurement. The PRATMO box model is also used for this purpose.

Balloon correlative data are used from the following instruments:

Differential Optical Absorption Spectroscopy (DOAS, Butz et al., 2006; Kreycky et al., 2013, and references therein), miniDOAS (Weidner et al., 2005), Limb Profile Monitor of the Atmosphere (LPMA, Butz et al., 2006), Spectromètre Infra-Rouge d’Absorption par Lasers Embarqués (SPIRALE, Kerzenmacher et al., 2008; Moreau et al., 2005), Système d’Analyse par Observation Zénithale (SAOZ, Pommereau and Piquard, 1994; Wetzel et al., 2007), SAOZ-BrO (Pundt et al., 2002), mini-SAOZ (Vicomte and Pommereau, 2011), MkIV (Toon et al., 2002), and Spectroscopie d’Absorption Lunaire pour l’Observation des Minoritaires Ozone et NO<sub>x</sub> – Nacelle 2 (SALOMON-N2) operating in solar occultation mode (Jégou et al., 2013). Table 2 provides information on these correlative datasets. The five balloon-borne instruments not previously used in OSIRIS NO<sub>2</sub> validation are LPMA, miniDOAS, SPIRALE, SALOMON-N<sub>2</sub> and MkIV. DOAS and SAOZ had been used previously (Brohede et al., 2007; Haley et al., 2004; Sioris et al., 2003). SPIRALE, LPMA, and SALOMON-N2 data were obtained from the CNES/CNRS-INSU Ether web site (<http://www.pole-ether.fr>).

Coincidence criteria are within 1000 km (Brohede et al., 2007) and on the same calendar day (using UT time). Only daytime correlative measurements are considered. The closest spatial coincidence is used if located within a 1000 km range. Comparisons between OSIRIS NO<sub>2</sub> and balloon correlative data are performed only down to the lower altitude limit of each profile retrieved with the v3.0 algorithm in an effort to keep the number of balloon-coincident altitudes the same between that algorithm and the algorithm debuting here. The upper limit for validation is determined in all cases by the balloon float altitude.

The best opportunities for validation of a deNO<sub>x</sub>ified profile come on March 4<sup>th</sup> and 16<sup>th</sup>, 2003 in Kiruna, Sweden when the NO<sub>2</sub> number density measured by LPMA and SAOZ, respectively, did not exceed 10<sup>9</sup> molec/cm<sup>3</sup> at any altitude. These peak number densities are the lowest in the



validation dataset. Note that OSIRIS does not measure in the northern polar region until late February due to a lack of sunlight at ~6 AM/PM. NO<sub>2</sub> profiles at southern high-latitudes have not been measured by balloon-borne instruments during the OSIRIS mission.

Averaging kernels are not taken into account since the vertical resolution is similar between the NO<sub>2</sub> profiles from most of the selected validation instruments and from the three OSIRIS algorithms, namely the operational v3.0, ‘fast’, and the algorithm described herein. Results from the current algorithm will be treated separately for each of the forward model modes (1D, 2D, and 3D) described in section 2.2.

In section 3, the 1D profiles used in the validation are those processed by a network of computers (using a solar zenith angle cutoff of SZA<90°). The entire OSIRIS data record has now been processed and is currently available at <ftp://odin-osiris.usask.ca/Level2/NO2/> following registration. One time-saving approximation is used to process the entire record: multiple scattering (MS) is only calculated at 21 wavelengths corresponding to the peaks and troughs of the NO<sub>2</sub> absorption cross-section across the spectral fitting window rather than the entire set of 107 wavelengths. The single-scattering radiances at the remaining 86 wavelengths are scaled by the ratio of radiances between multiple scattering and single scattering simulations at these 21 wavelengths, linearly interpolated to the remaining wavelengths. The radiance error due to the MS approximation is typically <0.05% at all wavelengths and all tangent heights used in the NO<sub>2</sub> retrieval.

20

### 21 **3 Results**

One difference in the method described above (Sect. 2) compared to previous OSIRIS NO<sub>2</sub> algorithms (Haley and Brohede, 2007; Bourassa et al., 2011) is the use of spectral information from longer wavelengths. Two sources of random error in retrieved NO<sub>2</sub> are shot noise, a consequence of the finite number of electrons generated by the detector by impinging limb-scattered photons, as well as radiation hits by energetic particles (e.g. protons) that are not filtered in the Level 1 data. The extended fitting window tends to reduce the susceptibility of the retrieved NO<sub>2</sub> to these noise sources. Using the 1D inversion approach described above, the 33 OSIRIS limb scans used in validation are processed with the extended fitting window and the



one used in the operational algorithm (Haley and Brohede, 2007). The standard deviation in retrieved NO<sub>2</sub> number density for each set (i.e. fitting window) is calculated over the common altitude range for each of the 33 scans. Natural variability of NO<sub>2</sub> is identical since the same limb scans and altitude ranges are used, so any difference in the standard deviation is due solely to the different spectral fitting windows. Figure 2 shows the slight reduction in NO<sub>2</sub> variability at all altitudes with the extended fitting window. Note that the NO<sub>2</sub> profiles retrieved from the two fitting windows do not have a significant bias ( $\pm 1$  standard error) between themselves, which points to the self-consistency of NO<sub>2</sub> spectral cross-section between the two windows.

Other benefits of the extended fitting window are that the NO<sub>2</sub> retrieval uncertainty is improved and the upper altitude limit of the retrieval moves slightly higher, as shown for a sample case in Fig. 2. For this tropical case, only at the NO<sub>2</sub> number density minimum at 15.5 km is the current retrieval unable to measure NO<sub>2</sub> with uncertainty of <100%, whereas with the fitting window of the operational algorithm, the measurements are below the lower detection limit between 11.5 and 17.5 km and at 7.5 km (Fig. 2). The significant upper tropospheric NO<sub>2</sub> enhancement at ~10 km is also observed by the mini-SAOZ (not shown).

Figure 3 shows that the 1D retrieval overestimates NO<sub>2</sub> by 29% relative to the 3D one at 17.5 km. This overestimate is due to the larger NO<sub>2</sub> concentrations that are present on the far side of the limb (yet neglected by a 1D model) for a SZA (86.6°) that is small enough to allow significant far side contribution to the radiance. The increasing NO<sub>2</sub> gradient toward the far side of the limb occurs because of the decreasing photolysis frequency of NO<sub>2</sub> with increasing SZA. The 2D retrieval underestimates slightly since the slightly lower NO<sub>2</sub> concentrations along the incoming solar beam are not included in that model version. The 3D retrieval clearly reduces the bias versus SAOZ relative to the 1D case at altitudes below 18 km where the relative diurnal gradients are expected to be largest. This improved precision is demonstrated below for a large ensemble of profiles by contrasting the correlation of the 1D and 2D retrievals and the standard error of the biases relative to coincident balloon data.

In order to compare the various OSIRIS NO<sub>2</sub> products to the balloon correlative data, all profiles are linearly interpolated onto a 1 km grid (12.0 to 39.0 km) commonly used in the balloon data. The ‘fast’ product tends to be limited in its vertical range at both extremes of the profile, not only relative to the other OSIRIS NO<sub>2</sub> products but also relative to the coincident balloon profiles and



1 also the ‘fast’ retrieval is not available for some coincidences (Fig. 4). The sample size is  $\geq 20$   
2 between 13 and 31 km for all OSIRIS NO<sub>2</sub> products and the validation discussion focusses  
3 mostly on this altitude range.

4 The v3.0 operational OSIRIS NO<sub>2</sub> has proven to be of high quality in the 15 to 42 km range from  
5 previous satellite intercomparisons (e.g. Haley and Brohede, 2007) and the same can also be  
6 inferred for the upper troposphere from the ability to detect small lightning-generated NO<sub>2</sub>  
7 enhancements with the expected latitudinal and longitudinal distribution (e.g. Sioris et al., 2007).  
8 Thus, in Fig. 5, we examine the standard error of the coincident profiles for OSIRIS and balloon  
9 data to compare the variability of each dataset, similar to Kerzenmacher et al. (2008). The upper  
10 altitude limit of the data obtained from balloons launched by CNES (Centre National d’Etudes  
11 Spatiales) is  $31 \pm 2$  km (Table 2). This excludes MkIV and SAOZ-BrO. From 31 down to 28 km,  
12 the balloon-borne NO<sub>2</sub> profiles show more scatter than any of the OSIRIS data products. This is  
13 likely due to the need to assume the NO<sub>2</sub> vertical distribution above float altitude for the remote  
14 sensors which is particularly problematic in the tropics where the NO<sub>2</sub> profile typically peaks at a  
15 higher altitude than for the extratropics. A second factor is the smaller absorption signal for the  
16 remotely-sensing balloon instruments due to the short path lengths when observing altitudes near  
17 float. Below 16 km, the balloon data clearly exhibits less scatter than OSIRIS, thereby  
18 quantitatively supporting the choice of balloon validation data for reasons discussed in Sect. 1.  
19 Kerzenmacher et al. (2008) show that this is not true for NO<sub>2</sub> at ~15 km measured by the  
20 (Atmospheric Chemistry Experiment) ACE satellite instruments versus OSIRIS v3.0 NO<sub>2</sub>.  
21 Between the three OSIRIS products, ‘fast’ exhibits the largest scatter at all altitudes, whereas the  
22 2D algorithm described here offers noticeably less scatter between 15 and 20 km as compared to  
23 the v3.0 product. The 3D algorithm takes 1.5 hours to retrieve a profile on the computer  
24 specified above and thus the entire set of balloon-coincident OSIRIS limb scans was not  
25 processed. The 2D algorithm takes half of the processing time of the 3D algorithm.

26 To determine if OSIRIS NO<sub>2</sub> is biased relative to the balloon data, we studied medians and  
27 means of individual profile differences over all coincidences (as a function of altitude). These  
28 results are shown in Figs. 6-7. It is clear that all of the OSIRIS NO<sub>2</sub> algorithms have a  
29 statistically significant bias near the NO<sub>2</sub> peak (typically ~30 km). This overestimate near the  
30 peak is similar to the overestimate by OSIRIS v3.0 relative to ACE Fourier Transform



1 Spectrometer and MAESTRO (Measurements of Aerosol Extinction in the Stratosphere and  
2 Troposphere Retrieved by Occultation) which coincided with the NO<sub>2</sub> peak altitude. Local biases  
3 were +17% and 14%, respectively (Kerzenmacher et al., 2008). Haley and Brohede (2007) found  
4 no such overestimate at the peak versus SAGE (Stratospheric Aerosol and Gas Experiment) III  
5 and POAM (Polar Ozone and Aerosol Measurement) III, but a similar positive bias near ~28 km  
6 versus HALOE (Halogen Occultation Experiment). The ‘fast’ product has the largest  
7 overestimate at the peak of ~16% with a sharp gradient in its bias, swinging to a statistically  
8 significant ~18% underestimate at 18 km. The retrieved profiles using the 2D retrieval described  
9 above have a similar bias profile shape with ‘fast’ (Figs. 6-7) but of smaller amplitude, with a  
10 ~+10% typical bias at the number density peak for the 2D retrieval and a ~10% negative bias at  
11 18 km which is statistically insignificant in terms of the median (Fig. 6). The bias of the v3.0  
12 product is similar to the alternative OSIRIS NO<sub>2</sub> products above 20 km, but between 15 and 17  
13 km, there is a significant positive bias using both central tendency statistics (Figs. 6-7). In  
14 contrast to ‘fast’ and v3.0 products, there are no altitudes in the lower stratosphere (below 24  
15 km) with statistically significant average and median biases for the 1D and 2D products. This  
16 conclusion is not sensitive to the use of the multiple scattering approximation used for the 1D  
17 product (see Sect. 2) (not shown).

18 While Figs. 6-7 address the systematic errors of the various OSIRIS products, it is also important  
19 to consider the precision of these data since a product whose biases in individual profiles average  
20 to zero could still fail to adequately capture the variability. This is potentially a greater concern  
21 for the v3.0 product since it relies slightly on a priori NO<sub>2</sub> and vertically smooths the retrieved  
22 NO<sub>2</sub> when measurement precision is lacking. However, Haley and Brohede (2007) and Haley et  
23 al. (2004) show that these are very minor concerns as the measurement response and vertical  
24 resolution do not deteriorate significantly down to an altitude of 12 km. Figure 8 shows that the  
25 v3.0 product captures the variability of the balloon measurements down to ~22 km as well as the  
26 2D and 1D products and better than ‘fast’ NO<sub>2</sub>, but then has larger scatter below 20 km than the  
27 2D and ‘fast’. The ‘fast’ retrieval benefits from more accurate forward modelling than the v3.0  
28 algorithm. The reduced scatter in the lower stratosphere is unlikely to be a difference between  
29 inversion approaches (MART versus the optimal estimation approach used in the v3.0 algorithm)  
30 as discussed above. The 2D retrieval also may have an edge over the v3.0 algorithm by virtue of  
31 using SaskTran as its forward model. There is also the specific benefit of its forward model



1 accounting for diurnal gradients in NO<sub>2</sub>. This can be seen clearly between 13 and 16 km where  
 2 the 2D retrieval is more precise than the 1D retrieval owing to this built-in photochemical  
 3 modelling capability. This confirms the better precision of the 2D retrieval suggested by Fig. 3,  
 4 and is only expected for large SZAs (McLinden et al., 2006). Note that errors due to the neglect  
 5 of diurnally varying chemical gradients can alternate in sign depending on the viewing geometry  
 6 (McLinden et al., 2006; Brohede et al., 2007).

7 Correlation can be used as an alternative statistic to verify whether an OSIRIS product captures  
 8 the variability observed by coincident balloon data. It is different from the standard error statistic  
 9 used in Fig. 8 since multiplicative or additive biases do not affect the correlation but do affect the  
 10 standard error of the individual biases. The correlation is calculated over the 46 coincident  
 11 balloon profiles at each available altitude (Fig. 4). Figure 9 shows a general decrease in  
 12 correlation between OSIRIS and coincident balloon data with decreasing altitude. There is higher  
 13 correlation with balloon NO<sub>2</sub> data for the 2D retrieval over the 1D retrieval at the lowest  
 14 altitudes, consistent with Fig. 3. The correlation of the v3.0 product is comparable with the 1D  
 15 and 2D product down to the lowest validated altitudes, whereas the ‘fast’ product has generally  
 16 lower correlations above 22 km than the OSIRIS NO<sub>2</sub> products relying on spectral fitting.

17 Finally, since the coverage of balloon is limited to latitudes north of -22°S, Fig. 10 is included to  
 18 illustrate the climatological (2001-2015) vertical and meridional distribution of stratospheric  
 19 NO<sub>2</sub> volume mixing ratio in the southern hemisphere during austral spring. For each NO<sub>2</sub>  
 20 vertical profile, the conversion to mixing ratio uses the local air density profile from the  
 21 ECMWF analysis.

## 22 **4 Discussion**

23 Next, we review and discuss the sensitivity of the retrieval to forward model parameters. There is  
 24 a slight sensitivity of the retrieved NO<sub>2</sub> to changes in aerosol extinction. Previous sensitivity  
 25 studies (Sioris et al., 2003; Haley et al., 2004) are consistent with the current findings. The scene  
 26 albedo in the visible can vary from <0.03 for calm ocean to almost unity for fresh snow, although  
 27 the NO<sub>2</sub> retrieval is, by design, insensitive to surface albedo by virtue of the high-TH  
 28 normalization of the radiance spectra (Sioris et al., 2003). Use of the retrieved scene albedo  
 29 (instead of the default value of 0.3) and of the v5.07 aerosol extinction profile further reduce  
 30 these sensitivities. Finally, use of OSIRIS-observed (v5.07) ozone instead of the default ozone



1 climatology in SaskTran (McPeters et al., 2007) is expected to have a minor impact on retrieved  
2 NO<sub>2</sub> via errors in modelling the atmospheric extinction based on Sioris et al. (2007).

3 With a method that involves simulating spectra at high tangent heights for the purpose of  
4 normalizing radiances simulated for tangent heights within the retrieval range, the forward model  
5 must accurately compute the radiance over a large range of tangent heights. The NO<sub>2</sub> retrieval  
6 error due to the pseudo-spherical approximation is expected to have its largest impact at low  
7 altitudes, because of the vertical gradient in pseudo-spherical RTM errors (Griffioen and  
8 Oikarinen, 2000). At the top of the retrieval range (40 km), errors cancel with the high altitude  
9 reference.

10 While not included in Figs. 5-9, Chahine's relaxation method tends to produce sharper extrema  
11 in the retrieved NO<sub>2</sub> profile than MART that tend to be slightly displaced vertically from those in  
12 the balloon data. This may stem from the coincidence criterion of 1000 km in distance, but may  
13 also relate to the ~2 km vertical sampling of OSIRIS and the 1 km altitude grid used for the  
14 retrieval. If the vertical grid of the retrieval is finer than the vertical sampling, a vertically narrow  
15 layer at 19 km, for example, would be retrieved as peaking at 18 km when the available radiance  
16 spectra are measured at THs of 18 and 20 km. The comparison of MART and Chahine inversion  
17 approaches is worth revisiting with the 1 km vertical sampling offered by the Ozone Mapping  
18 Profiler Suite (Jaross et al., 2014).

19 In order to rigorously validate this new retrieval algorithm in the upper troposphere, tropospheric  
20 chemistry must be added to the photochemical model used to scale the OSIRIS observations to  
21 balloon local time. The reaction of NO<sub>2</sub> with the hydroxyl radical to form HNO<sub>3</sub> drives the  
22 diurnal variation of tropospheric NO<sub>2</sub> more than N<sub>2</sub>O<sub>5</sub> photolysis (Boersma et al., 2009).  
23 Accurate knowledge of the seasonal variation of the OH concentration is required for modelling  
24 the diurnal variation of upper tropospheric NO<sub>2</sub>.

## 25 **5 Conclusions**

26 Profiles of NO<sub>2</sub> retrieved from OSIRIS have been improved in terms of reduced bias and scatter  
27 versus the current v3.0 operational algorithm in the lower stratosphere (e.g. 15 km) as  
28 determined using highly accurate balloon measurements as truth. The bias is within ~±10%  
29 between 14 and 37 km.



1 The benefits of spectral fitting and extending the fitting window to longer wavelengths are  
2 evident at the highest altitudes where the photoelectron shot noise tends to be large relative to the  
3 NO<sub>2</sub> absorption signal, but also at the lowest altitudes where the retrieved profile shape is largely  
4 driven by small differences in NO<sub>2</sub> SCDs obtained from spectra at adjacent tangent heights.  
5 Algorithms that exploit the richness of the NO<sub>2</sub> absorption spectrum are shown to better capture  
6 the mid-stratospheric variability of this key constituent. The use of a fully spherical forward  
7 model is an important advantage, particularly at lower altitudes since, while the pseudo-spherical  
8 RTM errors are largest at the top of the retrieval range, the largest NO<sub>2</sub> retrieval errors due to the  
9 use of a pseudo-spherical RTM occur at low altitudes because a high tangent height reference is  
10 used in the retrieval.

11 A model capable of modelling diurnal gradients in NO<sub>2</sub> also helps to improve the precision of the  
12 retrieved number density profile, particularly at the lowest altitudes where the horizontal  
13 gradients in NO<sub>2</sub> are sharpest and where the radiance is predominantly from the near side of the  
14 limb.

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21 for validation. SALOMON-N<sub>2</sub>, SPIRALE, SAOZ, mini-SAOZ, and LPMA data were obtained  
22 from CNES/CNRS-INSU Ether web site.

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1 Table 1. Main features of OSIRIS NO<sub>2</sub> algorithms compared in this work.

	<i>v3.0 operational</i>	<i>'fast'</i>	<i>1D, 2D, 3D</i>
Wavelength range (nm)	435-451	438-450	435-477
Number of spectral pixels	41	13	107
RT model	LIMBTRAN	SaskTran (1D)	SaskTran (1D, 2D, 3D)
Inversion scheme	Optimal estimation	MART (0.5: 0.3: 0.2)	MART (0.6: 0.3: 0.1)
Cloud detection	Yes	No	Yes
NO <sub>2</sub> absorption cross- sections	220 K (Vandaele et al., 1998)	(Burrows et al., 1998)	220, 298 K (Vandaele et al., 1998)
Interpolation of fitted cross-section to effective <i>T</i>	No	yes, using local ECMWF temperature at each TH	yes, using local ECMWF temperature at each TH

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2 Table 2. Summary of validation data based on OSIRIS coincidences. The minimum relative  
 3 uncertainty is reported on the native vertical grid. Vertical resolution is quoted or calculated for  
 4 an altitude of 16 km. The vertical resolution is provided by Butz (2006) for LPMA and DOAS,  
 5 Weidner et al. (2005) for mini-DOAS, available at <http://mark4sun.jpl.nasa.gov/m4data.html> for  
 6 MkIV, and calculated for limb occultation for the SAOZ-type instruments, including  
 7 SALOMON-N2 (Jégou et al., 2013) assuming vertical resolution equal to the vertical extent of  
 8 the solar disk at the tangent point. Uncertainties are for a 1 km vertical grid, except for SPIRALE  
 9 (0.005 km) and mini-DOAS (2 km).

Instrument	Number of coincidences	Relative uncertainty (%)	Vertical Resolution (km)	Average lower altitude (km)	Average upper altitude (km)
MkIV	6	2	~2	9	38
SPIRALE	4	<1	0.005	16	33
SALOMON-N2	1	15	2	15	33
DOAS asc	3	5	1.5	9	29
DOAS limb	4	<1	1.5	13	32
mini-DOAS	2	4	2	8	29
LPMA asc	1	10	6	10	33
LPMA limb	4	10	5	15	31
Mini-SAOZ	1	4	2	15	29
SAOZ-BrO asc	1	2	1	3	36
SAOZ-BrO SS	1	4	2	22	37
SAOZ asc/desc	9	<1	1	6	29
SAOZ limb	9	<1	2	11	29

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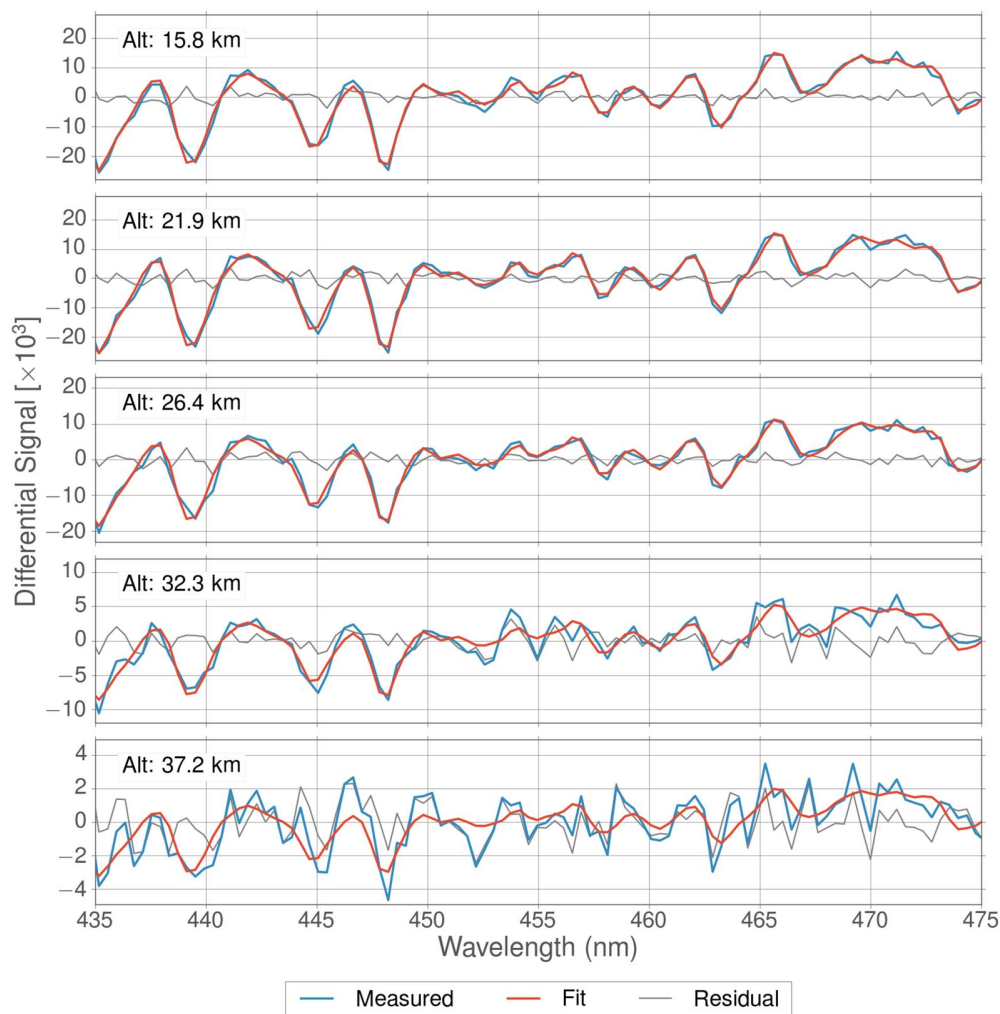
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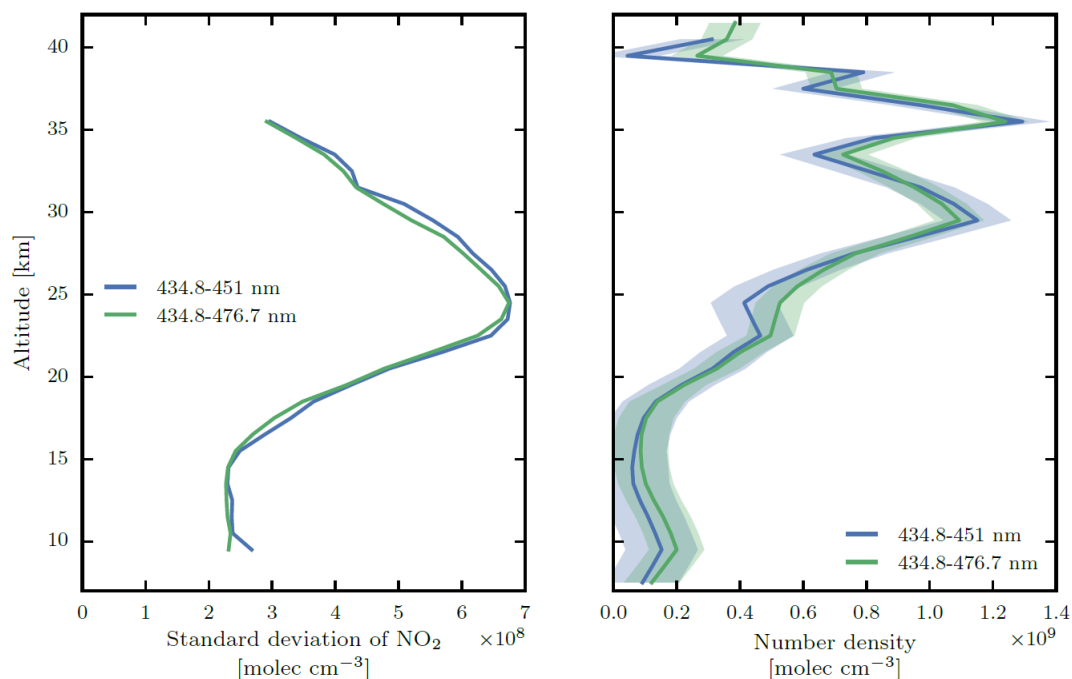


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3 Figure 1. Measured and fitted differential optical depth (DOD) as a function of tangent height  
 4 ('Alt') in the spectral window used for NO<sub>2</sub> retrieval. The residual is calculated as measured  
 5 DOD – fitted DOD.



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3 Figure 2. (left) Comparison of the standard deviation of the NO<sub>2</sub> number density profiles  
 4 retrieved with the algorithm described above using the default window (435-477 nm) versus the  
 5 fitting window used in the OSIRIS v3.0 operational algorithm (Haley and Brohede, 2007), (right)  
 6 Comparison of a single NO<sub>2</sub> profile retrieved from scan 20 of orbit 60346 with the algorithm  
 7 described above using the default window (435-477 nm) versus the fitting window used in the  
 8 OSIRIS v3.0 operational algorithm.

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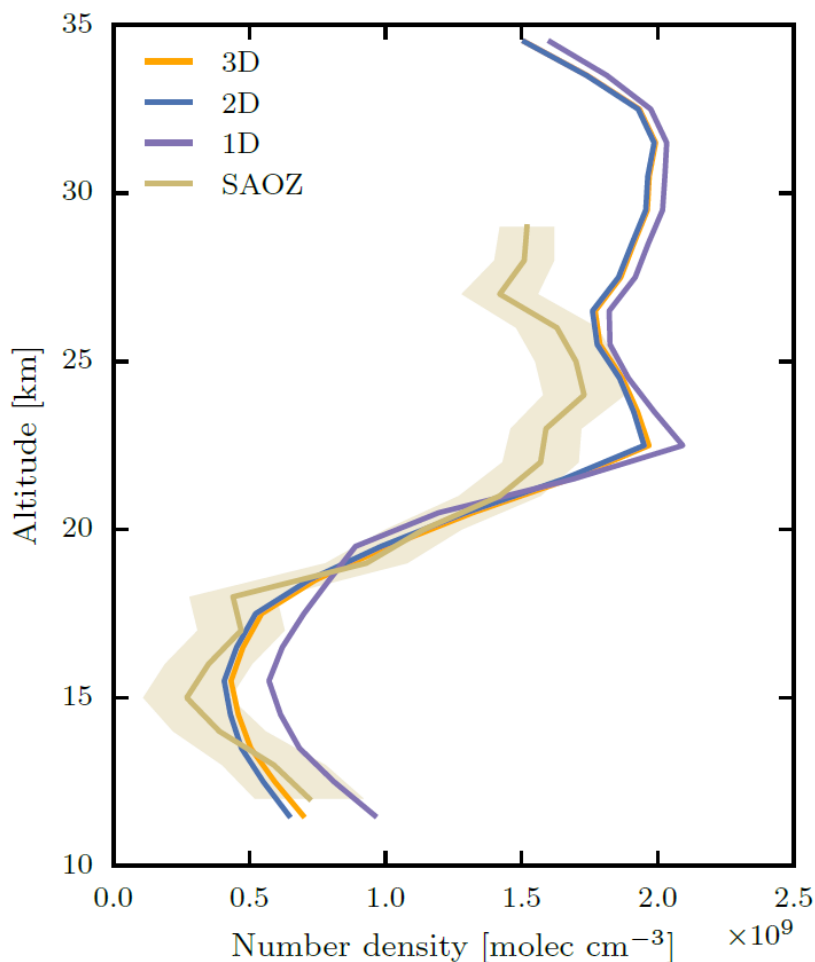
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 2 Figure 3. NO<sub>2</sub> profile retrieved from scan 44 of orbit 16011 using the 1D, 2D, and 3D retrievals,  
 3 all converted to sunset (tangent point SZA=90°), the local time of coincident SAOZ  
 4 measurements. OSIRIS is on the day side with SZA=86.6° (pm) and an azimuth difference angle  
 5 of 99° such that the far side of the limb is closer to the terminator.

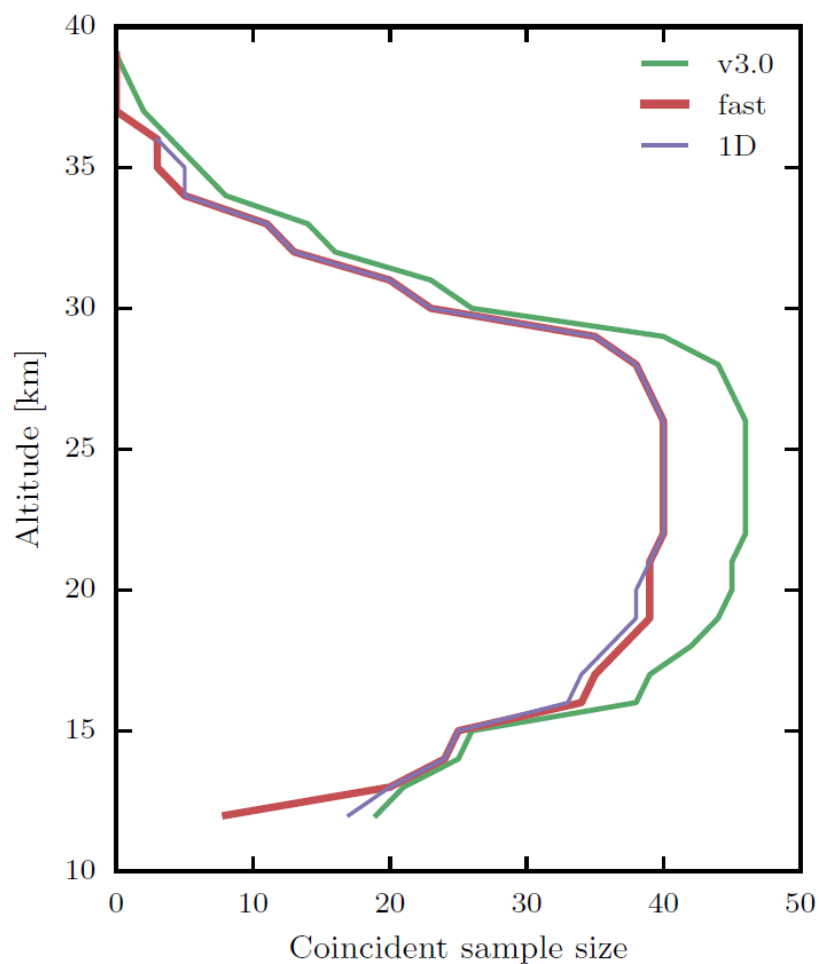
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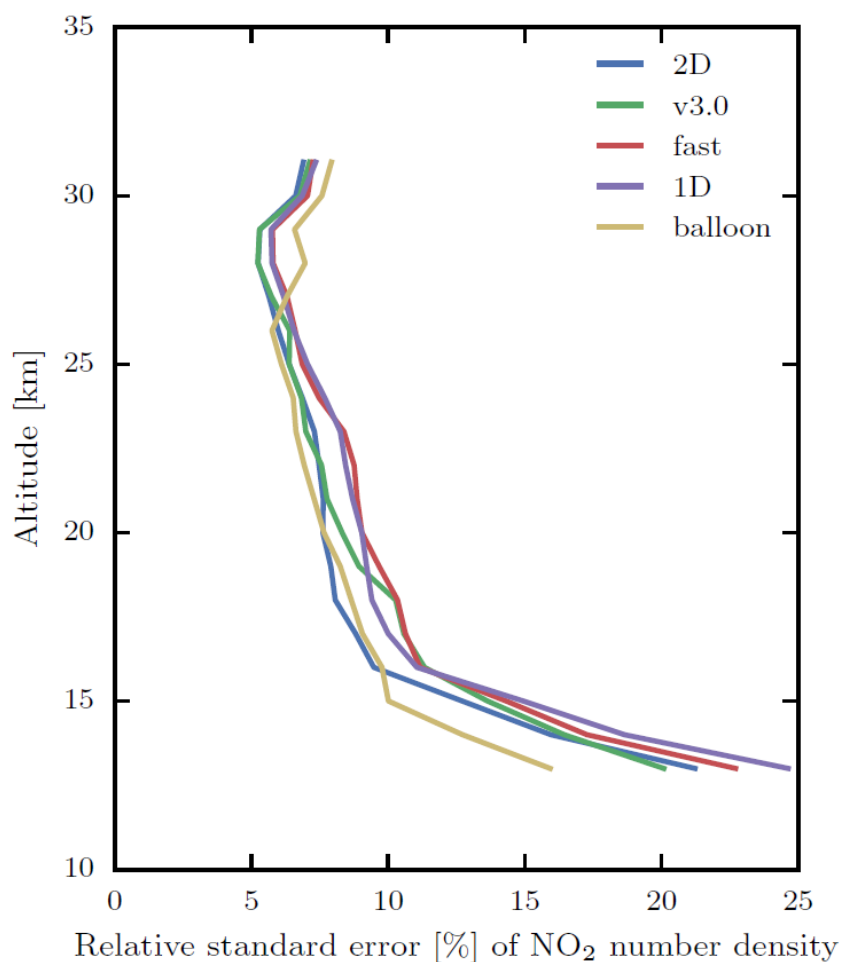
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 2 Figure 4. Sample size of coincident balloon data versus altitude and as a function of retrieval  
 3 algorithm.

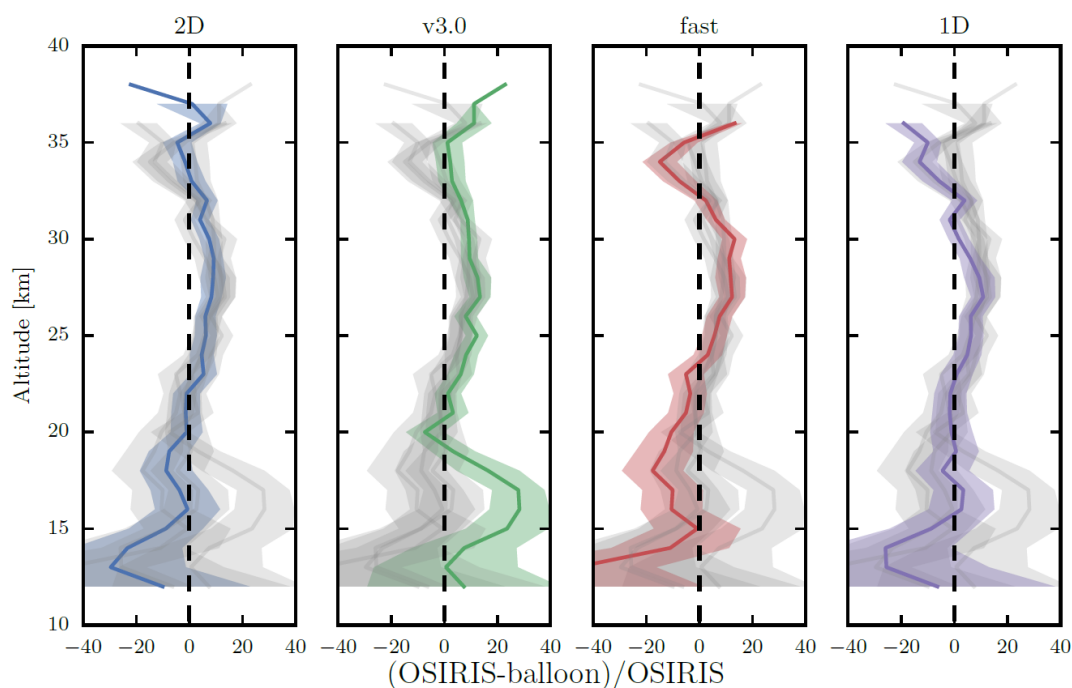
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 2 Figure 5. Relative standard error of coincident NO<sub>2</sub> profiles. Note that OSIRIS profiles have  
 3 been scaled to various local times of the balloon measurements which adds random error to the  
 4 OSIRIS profiles due to random variations between the photochemical model atmosphere and true  
 5 atmosphere.

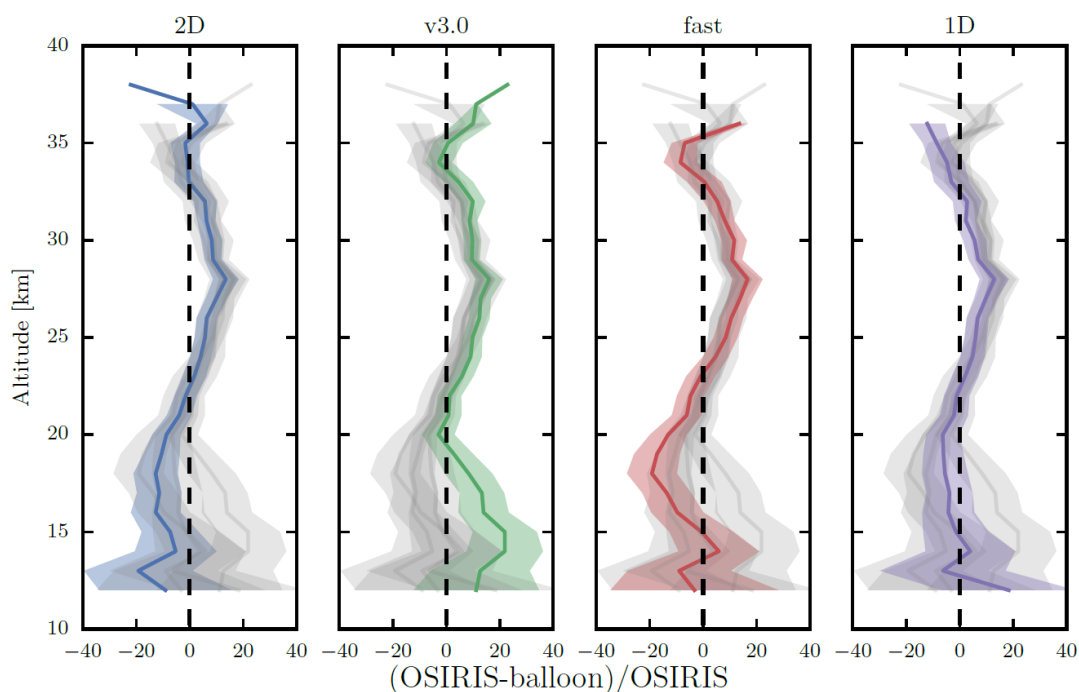
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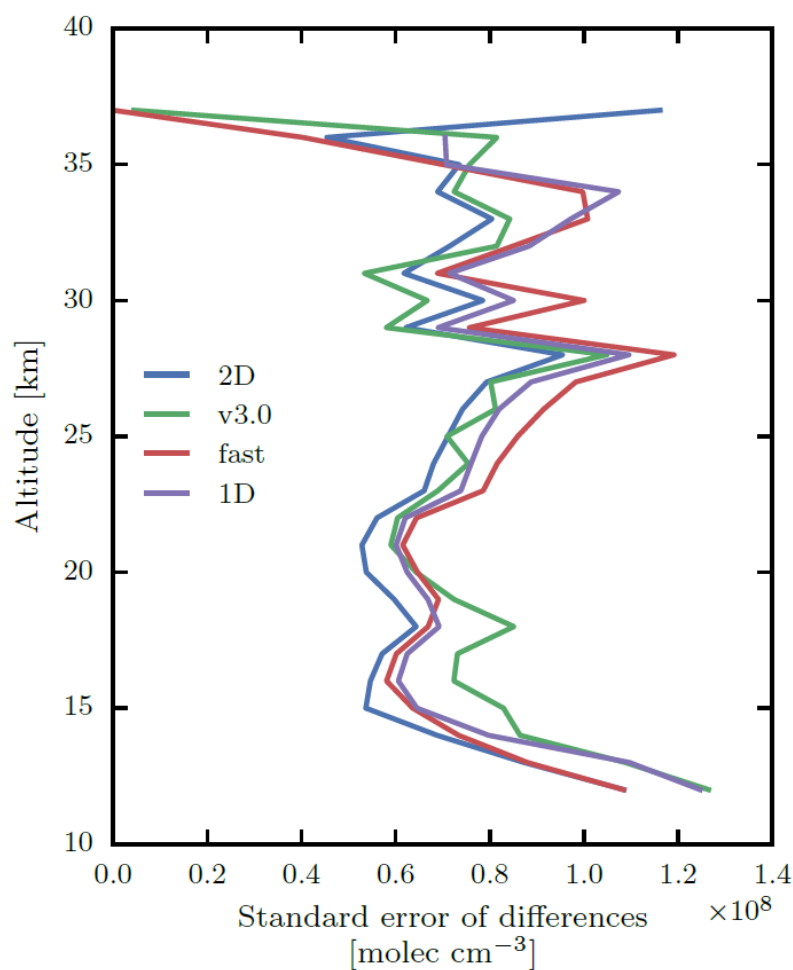


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 2 Figure 6. Median of individual biases versus balloon data. The error bar shows  $\pm 1$  standard error  
 3 of the median  $\text{NO}_2$  bias profile. The median and standard error of the individual biases are  
 4 converted to relative quantities by dividing by the corresponding median OSIRIS  $\text{NO}_2$  profile  
 5 (scaled to balloon local time). The relative median bias profiles for the other three algorithms are  
 6 shown in gray for comparison.

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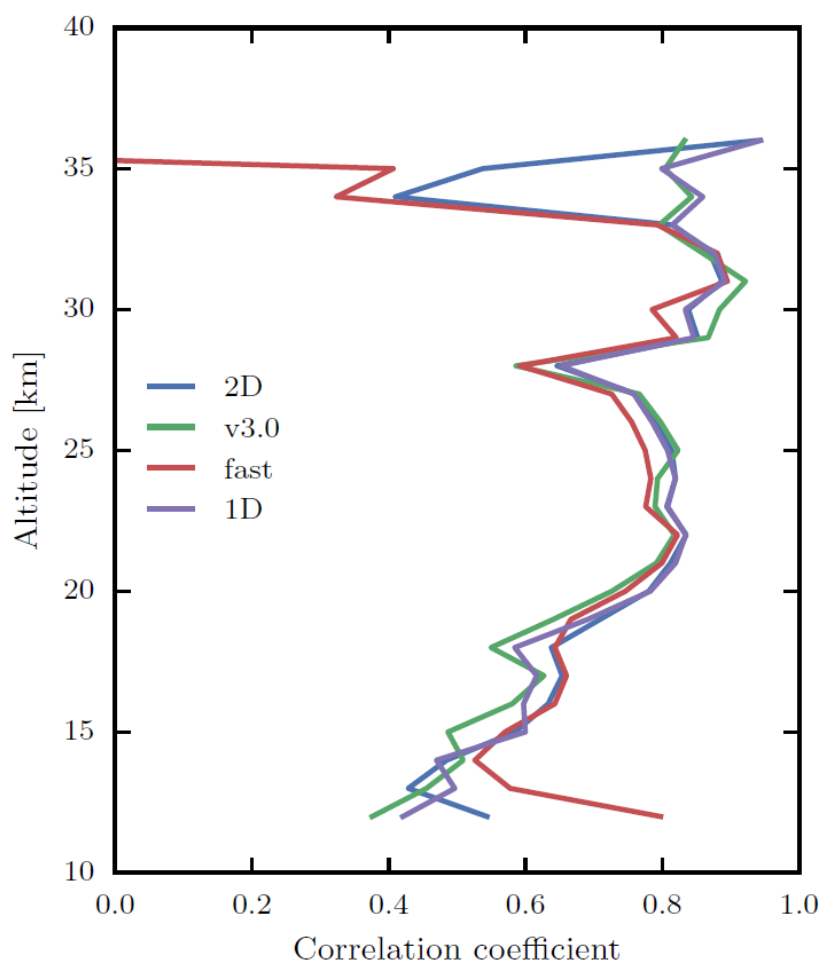


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 2 Figure 7. same as Fig. 6, but the bias is an average over individual differences between OSIRIS  
 3 and balloon correlative data. The average and standard error of the individual biases are  
 4 converted to relative quantities by dividing by the corresponding average OSIRIS NO<sub>2</sub> profile  
 5 (scaled to balloon local time).



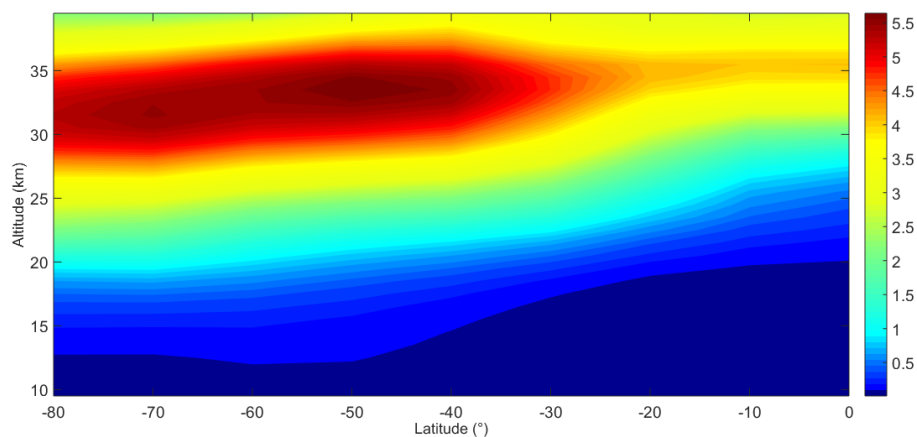
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 2 Figure 8. Standard error of the individual biases relative to balloon correlative data for different  
 3 OSIRIS NO<sub>2</sub> products. The curves here correspond to the half-widths of the error bars in Figs. 6-  
 4 7.

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2 Figure 9. Correlation between various OSIRIS NO<sub>2</sub> products and the balloon validation dataset.



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 2 Figure 10. Climatological map of zonal mean NO<sub>2</sub> volume mixing ratio (ppb) with 10°  
 3 latitudinal binning and 1 km vertical binning centered at altitudes between 9.5 and 39.5 km for  
 4 November (2001-2014).  
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