



# Improvements of the OMI O<sub>2</sub>-O<sub>2</sub> Operational Cloud Algorithm and Comparisons with Ground-Based Radar-

# **3 Lidar Observations**

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8 Abstract. The OMI (Ozone Monitoring Instrument) OMCLDO2 cloud product supports trace gas 9 retrievals of for example ozone and nitrogen dioxide. The OMCLDO2 algorithm derives the effective 10 cloud fraction and effective cloud pressure using a DOAS fit of the O<sub>2</sub>-O<sub>2</sub> absorption feature around 477 11 nm. A new version of the OMI OMCLDO2 cloud product is presented that contains several 12 improvements, of which the introduction of a temperature correction on the O2-O2 slant columns and the 13 updated look-up-tables have the largest impact. Whereas the differences in the cloud fraction are limited 14 to approximately 0.1, the differences of the cloud pressure can be up to 200 hPa, especially at cloud 15 fractions below 0.3. As expected, the temperature correction depends on latitude and season. The updated 16 look-up tables have a systematic effect on the cloud pressure at low cloud fractions. The improvements 17 at low cloud fractions are very important for the retrieval of trace gases in the lower troposphere, for 18 example for nitrogen dioxide and formaldehyde. The cloud pressure retrievals of the improved algorithm 19 are compared with ground-based radar-lidar observations for three sites in the mid-latitudes. For low 20 clouds that have a limited vertical extent the comparison is favorable. For higher clouds, which are 21 vertically extensive and often contain several layers, the satellite retrievals give a lower cloud-height. 22 For high clouds mixed results are obtained.

#### 23 Introduction

24 The Ozone Monitoring Instrument (OMI) is a imaging spectrometer developed by The Netherlands and 25 Finland that has been launched in 2004 on board of the NASA EOS Aura satellite (Levelt et al., 2006). 26 OMI has a continuous spectral coverage from 270-500 nm, with a resolution of approximately 0.5 nm. 27 The primary data products from OMI are concentrations of trace gases, including ozone, nitrogen dioxide 28 and formaldehyde. The trace gas retrieval algorithms rely on a priori information of cloud properties. For 29 tropospheric trace gas retrievals, clouds are among the leading error sources in the retrieval (e.g. Boersma 30 et al., 2011). 31 The OMI O<sub>2</sub>-O<sub>2</sub> cloud product (OMCLDO2) contains information on the cloud fraction and cloud 32 pressure for each ground pixel. The OMCLDO2 product has been designed to support the trace gas 33 retrieval algorithms and is therefore driven by what these algorithms need for cloud information. The 34 trace gas retrieval algorithms use the independent pixel approximation (Zuidema and Evans, 1998) 35 representing clouds as Lambertian reflectors with a fixed albedo of 0.8 (Stammes et al., 2008). To be 36 consistent with the trace gas retrievals, the OMCLDO2 product uses the same cloud model. The initial

37 OMCLDO2 algorithm has been described by Acarreta et al. (2004). Because the amount of information





38	in the OMI spectral range is limited, the algorithm derives an effective cloud fraction and an effective		
39	cloud pressure. The cloud fraction and cloud pressure are derived from the continuum radiance		
40	and the depth of the $\mathrm{O_2\text{-}O_2}$ absorption feature around 477 nm. The algorithm does not distinguish		
41	between clouds and aerosols. Cloud-free conditions with significant thick aerosols layers will be		
42	represented by small cloud fractions. Similarly, thin clouds, for instance cirrus, will also be represented		
43	by a small cloud fraction. The main a-priori information that is used is the surface reflectance and the		
44	surface altitude, which are obtained from static look-up tables. Validation studies (Sneep et al., 2008)		
45	have shown that the effective cloud fraction compares well with effective cloud fractions derived from		
46	the cloud optical thickness observed by MODIS (Moderate Resolution Imaging Spectroradiometer) and		
47	that the	derived cloud pressure determines a level somewhere near the middle of the clouds. This is	
48	differen	t from the cloud pressures derived from the thermal infrared, which are very sensitive near the	
49	actual c	loud tops. The OMCLDO2 retrieval is very similar to the FRESCO algorithm (Wang et al., 2008)	
50	with the	e difference that it is based on $O_2$ - $O_2$ rather than $O_2$ absorption lines. The reason for using $O_2$ - $O_2$	
51	is that t	he OMI spectral range doesn't cover the oxygen absorption bands. An important difference of	
52	using th	ne oxygen dimer is that its absorption scales with the oxygen density squared, which makes it	
53	increasi	ngly more sensitive to the lower altitudes in the atmosphere. Besides the OMCLDO2 algorithm,	
54	there is also an OMI product based on the information from Raman scattering (Joiner et al., 2012; Joiner		
55	and Vas	ssilkov, 2006). It has been demonstrated that this product is also sensitive to the middle of cloud	
56	layers, v	which has been referred to as the optical centroid pressure.	
57	This pa	per describes version 2.0 of the OMCLDO2 product. Whereas updates and reprocessing was	
58	perform	ed regularly in the past, the version 2.0 contains the following improvements and extensions:	
59	1.	A temperature correction is implemented which is needed because of the density-squared nature	
60		of the $O_2$ - $O_2$ absorption;	
61	2.	Besides the independent pixel approximation, a second cloud model is implemented, which	
62		represents the scene as a Lambertian surface at a certain pressure level. The retrieved parameters	
63		are the scene albedo and scene pressure;	
64	3.	The look-up-tables that are used to derive the cloud fraction and pressure have a higher number	
65		of nodes, especially for the surface albedo and the surface altitude;	
66	4.	A method has been implemented to remove outliers from the spectral fitting;	
67	5.	The resolution of the a priori surface altitude is brought in line with the average OMI spatial	
68		resolution;	
69	6.	The gas absorption cross-sections are made consistent with the OMI $\mathrm{NO}_2$ retrieval algorithm	
70		(Geffen et al., 2014).	
71	This pa	per is organized as follows: in section 2 we describe the OMCLDO2 algorithm, focusing on the	
72	improvements that have been introduced in this version. In section 3 we discuss differences compared to		
73	the previous version. In section 4 we present comparisons of the cloud pressure to ground based radar		
74	observations.		





### 75 Algorithm

76	The OMCLDO2 retrieval consists of two main steps: first a DOAS (Differential Absorption
77	Spectroscopty) fit is performed on the spectral region between 460 and 490 nm to derive the O2.O2 slant
78	column amount $N_{s,O2-O2}$ and the continuum reflectance $R_c$ . In the second step these parameters are

converted into cloud fraction  $c_{f}$ , cloud pressure  $p_{cld}$ , scene albedo  $A_{scn}$  and scene pressure  $p_{scn}$ .

#### 80 DOAS fit

The DOAS fit is performed on the Earth's reflectance. OMI measures the Earth's radiance and once per day the solar irradiance. The wavelength grids of the Earth radiance and solar irradiance differ, because of the Doppler shift and because of non-homogeneous filling of the slit for partly cloudy scenes (Voors et al., 2006). For each ground pixel, the radiance (*I*) and irradiance (*F*) are brought on the same spectral grid (see Van Geffen et al., 2015) and the reflectance is calculated as  $R(\lambda) = \frac{\pi I(\lambda)}{\cos \theta_0 F(\lambda)}$ , where  $\lambda$  is the wavelength and  $\theta_0$  is the solar zenith angle. Next, the following equation is used for the DOAS fit:

87 
$$R(\lambda) = P(\lambda)e^{-(N_{s,0202}\sigma_{0202}(\lambda) + N_{s,03}\sigma_{03}(\lambda))} \cdot (1 + c_R \frac{I_R(\lambda)}{F(\lambda)})$$
(1)

where  $P(\lambda)$  a polynomial of the first order,  $N_{s,O2O2}$  the slant column of  $O_2$ - $O_2$ ,  $\sigma_{O2O2}(\lambda)$  the  $O_2$ . $O_2$  cross section convolved with the OMI slit function,  $N_{s,O3}$  the slant column of  $O_3$ ,  $\sigma_{O3}(\lambda)$  the  $O_3$  cross section convolved with the OMI slit function,  $I_R(\lambda)$  a synthetic radiance Raman spectrum convolved with the OMI slit function and  $c_R$  a scale parameter for the amount of Raman scattering. For the reference cross sections for  $O_2$ - $O_2$  we use (Thalman and Volkamer, 2013) at 293 K and for  $O_3$  we use (Bogumil et al., 2000) at 220 K.

94 We solve Eq. 1 using a modified Levenberg-Marquardt method, using the errors for the radiance and 95 irradiance as weights. The fit parameters are the slant columns  $N_{s,O2O2}$  and  $N_{s,O3}$ , and  $c_{R}$ , and the coefficients for the polynomial  $P(\lambda)$ . In addition, also the diagnostics of the fit is obtained, including the 96 97 residuals and error estimates for all fit parameters. The residuals are analyzed for possible outliers. Such 98 outliers may be caused by high-energy particles hitting the detector or by varying dark current. Although 99 al the information in the OMI Level 1B product is used to remove bad spectral pixels, some may remain. 100 For outlier detection several methods have been used (e.g. Richter et al., 2011), which are mostly based 101 on Gaussian statistics, i.e. by using the mean and standard deviation of the residual. Because particle hits 102 will cause only increases in detected radiance and because the mean and standard deviation themselves 103 are strongly affected by outliers, we selected the so-called box-plot method for outlier detection (http://www.itl.nist.gov/div898/handbook/prc/section1/prc16.htm). This method determines lower and 104 upper values based on the 25<sup>th</sup> and 75<sup>th</sup> percentile of a distribution. If the lower quartile is Q1 and the 105 106 upper quartile is Q3, then the difference (Q3 - Q1) is called the interquartile range or IQ. We define 107 outliers as those values smaller than Q1 - 1.5 IQ or larger than Q3 + 1.5 IQ. After removal of the outliers, 108 we redo the fitting of the spectrum to provide the final fit parameters. We have noted that the outlier 109 removal is not stable; continuing iterating and each time applying the outlier removal procedure will 110 often result in more and more removed spectral pixels. We therefore iterate only one time, thus removing 111 the largest outliers.





### 113 Conversion to Cloud and Scene Parameters

### 114 Radiative transfer modelling

115 For the conversion of the DOAS fit parameters into respectively cloud fraction and pressure, and scene 116 albedo and scene pressure, we use radiative transfer modeling. The difference between the independent 117 pixel approximation (IPA) (Zuidema and Evans, 1998) that determines the cloud fraction and pressure, 118 and the Lambertian equivalent reflectance (LER) model that determines the scene albedo and pressure, 119 is illustrated in Fig 1. It is noted that the clouds and the ground surface in the IPA model are treated as 120 opaque Lambertian reflectors. Therefore, the name LER maybe somewhat confusing, but is used for 121 consistency with the existing literature. For each ground pixel, both the IPA and LER method is applied. 122 The IPA requires a-priori information on the surface reflectance and surface pressure. The clouds are 123 represented as Lambertian reflectors with an albedo of 0.8. Different studies have found that this is an 124 optimal choice for the purpose of cloud corrections in trace retrieval schemes (see (Stammes et al., 2008) and references therein). Using such a high albedo for the clouds will represent thin clouds covering the 125 126 entire ground pixel as small cloud fraction. Thus, the cloud-free part will implicitly model the 127 transmission of light through the cloud, which is otherwise absent in the Lambertian cloud model. 128 For very small cloud fractions the cloud pressure derived using the IPA will become undetermined. In 129 case of surface albedo's close to 0.8, e.g. over snow and ice, the IPA retrieval for both cloud fraction and 130 pressure will become undetermined. In such cases, the LER method may be a good fallback. 131 For both the IPA and LER model, we use the same set of forward model simulations of the reflectance 132 between 460 and 490 nm, see Table 1. These simulations are performed for a mid-latitude summer 133 standard atmosphere. The correction for different temperature profiles is discussed later on in this section. On the simulated reflectance the same DOAS fit is performed as for the measured OMI spectra (Eq. 1). 134 135 For all the nodes listed in Table 1, we obtain the slant column  $O_2O_2$  as well as the continuum reflectance at 475 nm. The continuum is computed by evaluating the polynomial  $P(\lambda)$  for this wavelength. 136

#### 137 Look-up-table inversion

138 Although we now have the information needed to derive the cloud fraction/pressure and the scene 139 albedo/pressure, we invert the tables to improve the computational speed. Instead of having the cloud 140 fraction and cloud pressure as nodes of the tables, we want to have the slant column  $O_2O_2$  and continuum 141 reflectance as nodes. This conversion process involves interpolation and extrapolation, for which we use 142 radial functions (http://docs.scipy.org/doc/scipylinear basis 0.15.1/reference/generated/scipy.interpolate.Rbf.html). 143 144 Because the simulated spectra cover a very wide range of conditions, it is unlikely that the extrapolations 145 in this inversion procedure have a large effect on the final result. The inversion is illustrated in Fig. 2. 146 The final result of the inversion procedure are look-up tables (LUTs) for the cloud fraction, cloud 147 pressure, scene albedo and scene pressure on the nodes listed in Table 2. In the retrieval algorithm linear 148 interpolation is applied on all dimensions, except for the solar zenith angle, for which spline interpolation 149 is applied. This is implemented because of the non-linear behavior at large solar zenith angles.





#### 150 Temperature correction

151 As will be described in this section, the slant column amount of O<sub>2</sub>-O<sub>2</sub> depends on the temperature profile, 152 even if the cross section is not temperature dependent. This is due to the nature of the dimers, of which 153 the absorption scales with the pressure-squared instead of being linear with pressure. Because this effect 154 turns out to be significant, we have developed a temperature correction. This correction allows the use 155 of the LUTs described above, which have been derived for a single pressure-temperature profile. By 156 applying temperature correction, the O2-O2 slant columns are scaled to the values for the reference 157 temperature profile that has been used to construct the LUTs. 158 To understand the temperature effect of the O<sub>2</sub>-O<sub>2</sub> slant columns, we write the reflectance as: 159  $R(\lambda) = R_0(\lambda) \exp\left(-\int_{z_0}^{TOA} m(z,\lambda) n_{02}^2 \sigma_{02-02}(\lambda) dz\right),$ 160 (2)161 162 where  $R_0(\lambda)$  is the reflectance if absorption by O<sub>2</sub>-O<sub>2</sub> is ignored;  $z_0$  is the altitude of a Lambertian cloud 163 or the Earth surface; TOA is he top of the atmosphere;  $m(z, \lambda)$  is the altitude resolved air mass factor 164 which is weakly wavelength dependent;  $n_{O_2}(z)$  is the number density of oxygen and  $\sigma_{O2-O2}(\lambda)$  is the absorption cross section of O2-O2. 165 166 In hydrostatic equilibrium, the integral over the altitude can be replaced by an integral over the pressure, 167 using  $dp/dz = -\rho(z)g$ , where is the density of air. By expressing the density of air as  $\rho(z) = M p / (Rg T(z))$ , where M is the mean molecular mass of dry air and  $R_g$  is the gas constant, Eq. 2 becomes: 168 169  $R(\lambda) = R_0(\lambda) \exp\left(\int_{p_0}^{p_{TOA}} \frac{R_g}{Mg} T(p) m(p,\lambda) n_{02}^2(p) \sigma_{02-02}(\lambda) \frac{dp}{p}\right).$ (3) 170 171 172 Finally, we can express the number density of air in as  $n_{0_2} = 0.21 p/(k_b T(p))$ , where  $k_b$  is Boltzmann's 173 constant and we assume a mixing ratio of oxygen of 21%. Substituting this in Eq. 3 gives: 174  $R(\lambda) = R_0(\lambda) \exp\left((0.21)^2 \frac{R_g}{M \, g \, k_B^2} \sigma_{02-02}(\lambda) \int_{p_0}^{p_{TOA}} m(p,\lambda) \, \sigma_{02-02}(\lambda) \, \frac{p}{T(p)} dp \right),$ 175 (4)176 177 which shows that the reflectance and hence the slant column of  $O_2$ - $O_2$  changes when the temperature 178 profile changes. It is noted that this is due to the density-squared nature of the absorption of  $O_2$ - $O_2$ . For 179 "normal" absorbers (no collision complex) the slant column is independent of the temperature profile, 180 apart from temperature dependence of the absorption cross section. 181 182 In order to investigate the magnitude of the bias that is introduced if the temperature dependence is

ignored simulations of the retrieval were performed. In the retrieval the mid-latitude summer profile is used while for the simulations either a mid-latitude winter profile or a sub-arctic winter profile is used. The bias was calculated for different true pressure levels of the cloud and for different cloud fractions.

186 Fig 3 shows that the maximum bias in the retrieved cloud pressure ranges from less than 50 hPa at large





- 187 cloud fractions to 200 hPa at very small cloud fractions. Such biases will have a significant impact on
- 188 trace gas retrievals, which are commonly limited to scenes with small cloud fractions.
- 189
- 190 The OMCLDO2 retrieval is based on a LUT approach and generating LUTs for different temperature
- 191 profiles in not feasible. Therefore we introduce a correction factor  $\gamma$  that translates the measured slant
- 192 column into the slant column for the reference pressure-temperature profile. Using Eq 4., we can compute
- 193 γ as:

194 
$$\gamma = \frac{N_s^{ref}}{N_s^{meas}} = \frac{\int\limits_{p_c}^{p_{TOA}} m(p,\lambda) \frac{p}{T_{ref}(p)} dp}{\int\limits_{p_c}^{p_{TOA}} m(p,\lambda) \frac{p}{T(p)} dp}$$
(5)

195 where T(p) is the actual temperature profile taken and  $T_{ref}(p)$  is the temperature profile used in the

196 creation of the look-up tables. In case of partial cloud cover and weak absorption we obtain

$$197 \qquad \gamma = \frac{N_s^{ref}}{N_s^{meas}} = \frac{(1-c_f)R_{clr}\int_{p_s}^{p_{TOA}} m_{clr}(p,\lambda) \frac{p}{T_{ref}(p)} dp + c_f R_{cld}\int_{p_c}^{p_{TOA}} m_{cld}(p,\lambda) \frac{p}{T_{ref}(p)} dp}{(1-c_f)R_{clr}\int_{p_s}^{p_{TOA}} m_{clr}(p,\lambda) \frac{p}{T(p)} dp + c_f R_{cld}\int_{p_c}^{p_{TOA}} m_{cld}(p,\lambda) \frac{p}{T(p)} dp}$$
(6)

where *R* is the reflectance at a representative wavelength in the fit window,  $p_s$  is the surface pressure and  $p_s$  the cloud pressure, and the subscripts *clr* and *cld* refer to the clear part and the cloudy part of the pixel, respectively.

To implement the temperature correction factor, new look-up-tables for the  $O_2$ - $O_2$  air mass factors  $m(p,\lambda)$ and the corresponding reflectance for a wavelength in the middle of the fit window have been generated. In the retrieval algorithm, the temperature correction is applied in an iterative manner because the cloud

204 fraction and pressure should be known to compute  $\gamma$ . As a default, we use three iterations to compute  $\gamma$ .

#### 205 A-priori information

206 The OMCLDO2 version 2 uses the following *a-priori* information.

207 For the absorption cross-sections for O2-O2, ozone and optionally NO2, as well as for the radiance Raman scattering, we use the spectra described in Van Geffen et al., (2015). For the surface reflectance, the 208 OMI derived monthly mean database described in Kleipool et al., (2008) extended to 5 years of OMI 209 210 data is used. For the temperature profiles needed for the temperature correction, we use a monthly mean 211 climatology at four times per day (00, 06, 12 and 18 UTC), derived from the NCEP reanalysis data for 212 the period 2005-2014. Actual temperatures maybe somewhat better than using a climatology. However 213 for practical reasons related to the operational data processing facility, we have decided to use a 214 temperature climatology. For detecting snow and sea-ice coverage, the Near-real-time Ice and Snow 215 Extent (NISE) product (Nolin et al., 1998) is used. 216





#### 217 Impact of algorithm updates

218 In this section we first compare the OMCLDO2 version 2 with the version 1.2.3 for one day of data. 219 Next, the impacts of each of the improvements are discussed separately. The impact of the improvements 220 are summarized in Table 2. 221 Figure 4 shows the OMCLDO2 retrieval results for 14 May 2005. This day has been selected arbitrarily 222 from the OMI data record. Note that we also have analysed other days, which show consistent results. 223 Figure 4a and b show the effective cloud fraction and the effective cloud pressure. Figures 4c and d show 224 the difference between version 2 and version 1.2.3. For areas with low effective cloud fractions, the 225 effective cloud fraction is approximately 0.01 higher in the version 2. Over the high latitudes in the 226 northern hemisphere considerably large positive and negative differences occur. These occur over snow 227 and ice, where the retrieval algorithm has problems to distinguish the clouds from the highly reflective 228 surface. Under such conditions, the accuracy of the retrieved effective cloud fraction will be very low. 229 Due to the assumed cloud albedo of 0.8, the cloud fraction will become undetermined when the surface 230 albedo is also close to this value. 231 The differences in effective cloud pressure are shown in Fig 4d. Version 2 shows higher cloud pressure 232 in the tropics and sub-tropics, and lower cloud pressures and mid and high latitudes. As discussed below, 233 this zonally dependent effect is caused by the temperature correction introduced in version 2. Especially

234 in the tropics, the differences in the cloud pressures are largest in regions with low cloud fractions. 235 Overall the uncertainty in the cloud pressure retrievals is a strong function of the effective cloud fraction. 236 This is illustrated in Figure 5, which shows the precision of the effective cloud pressure retrievals as a 237 function of the effective cloud pressure. The precision is calculated by the propagation of the DOAS fit 238 errors of the O<sub>2</sub>-O<sub>2</sub> slant columns and of the continuum reflectance. For cloud fractions below 0.1 the 239 average precision is larger than 20 hPa with a very large spread, whereas for cloud fractions above 0.9 240 the precision is less than 10 hPa with a much smaller spread. It is noted that other errors sources, for 241 example in the *a priori* surface albedo will also have a much stronger impact at low effective cloud 242 fractions.

#### 243 Temperature correction

244 The correction for the temperature dependence is described above. Based on a temperature climatology, 245 a correction factor is computed and applied to the O2-O2 slant columns. Figure 4g shows the temperature 246 correction factor for the OMI observations on 14 May 2005. Because the temperature correction factor 247 is computed relative to the midlatitude summer atmosphere, it is larger than 1 in the tropics and smaller 248 at the higher latitudes. On top of this general behavior there is spatial structure related to cloud structures, 249 especially when the clouds are at high altitudes and have significant optical thickness. The effect of clouds on the temperature correction factor is described in Eq. 6. For high and thick clouds the 250 251 temperature correction is in most cases closer to 1, indicating that the largest differences between the 252 climatological temperature and the mid-latitude summer atmosphere occurs at the lowest altitudes. 253 To test the impact of the temperature correction factor on the effective cloud fraction and pressure, we 254 produced datasets with and without the temperature correction applied for two days of OMI data in

different seasons (14 May 2005 and 15 November 2005). While the impact on the cloud fraction is





256 negligible, the impact on the cloud pressure can be significant. Fig. 6 shows the difference between the 257 retrievals without and with the correction applied, as a function of the effective cloud fraction. The impact 258 of the correction on the cloud pressure increases towards smaller cloud fractions. Depending on whether 259 the correction factor is smaller or larger than 1, the impact on the cloud pressure can be both positive or 260 negative. For cloud fractions below 0.2, the impact of the temperature correction can be as large as -100 261 to 150 hPa, whereas for cloud fractions larger than 0.2 the impact is in the range -20 to 40 hPa. For the 262 higher latitudes ( $\gamma$ >1) the clouds are at lower pressures (higher altitude) when the temperature correction 263 is applied, whereas in the tropics and sub-tropics the effects is reversed. 264 Fig. 6 can be compared to Fig. 2, which is based on retrieval simulations. Although Fig. 6 shows the

difference with and without the temperature corrections, and Fig. 2 shows the difference with the simulated truth, the behavior and magnitude of the bias is very similar. It is noted that for Fig. 2 only temperature profiles have been used which are colder in the troposphere than the reference mid-latitude summer atmosphere. Therefore, Fig 2 shows only positive biases, whereas in the tropics and sub-tropics Fig. 6 also shows negative values.

#### 270 Look-up-tables

To test the impact of the LUTs that are used to derive the effective cloud fraction and effective cloud pressure, we produced a datasets using the version 2 algorithm with the new and the old LUTs. The cloud fraction with the new LUTs is about 0.01 larger than with the old version, except over snow and ice regions where the cloud fraction is in most cases significantly smaller. Because over snow and ice covered regions the cloud fraction is highly uncertain as the algorithm is not able to distinguish clouds from highly reflective surfaces, this impact is not unexpected.

277 The effect of the new LUTs on the effective cloud pressure is shown in Fig. 7c. This figure shows the 278 difference in the cloud pressure (old minus new) as a function of the effective cloud pressure. The 279 differences become significant at cloud fractions smaller than 0.25, where the difference shows an 280 oscillating behavior. At a  $c_f$  of approximately 0.125 a minimum is reached and at smaller cloud fractions 281 the mean difference reverses sign and increases towards lower  $c_{f}$ . To investigate the nature of this 282 behavior, Fig. 7a and 7b show the distribution of the retrieved cloud pressures as a function of cloud 283 fraction for the old and new LUT datasets. From these figures it is clear that the origin of the oscillating 284 behavior of the difference is in the retrievals with the old LUTs. Fig 7a shows that with the old LUTs the 285 cloud pressure increasing strongly towards lower cloud fractions, for which we have no physical 286 explanation. The results with the new LUTs (Fig. 7b) do not show this. We attribute the large 287 improvements with the new LUTs to the larger number of radiative transfer calculations on which it is 288 based, as well as the improved interpolation scheme that was used to produce it. 289 Figs. 7a and 7b also show that the effective cloud pressure for the largest  $c_f$  bin is significantly larger. A

290 further inspection showed that this is caused retrievals over snow and ice covered regions, for which the

291 cloud pressure retrievals are highly uncertain. For such cases the scene albedo and pressure provided by

the version 2 algorithm can be used.





#### 293 Outlier removal

294 The outlier removal procedure that was introduced in the version 2 of the algorithm removes spectral 295 pixels from the DOAS fit after evaluation of the fitting residuals. Outliers can have different behavior: 296 they can be transient, e.g. occurring only for spectral pixels for a few pixels, or they can occur systematically for certain spectral pixels. When outliers are detected they are removed from the data, 297 298 which will decrease the number of wavelengths used in the DOAS fit. Fig. 4h shows the number of 299 wavelengths used in the fit for 14 May 2005. The most prominent feature are the reduced values over the South America caused by the South Atlantic Anomaly (SAA). In this region the number of high 300 301 energetic particles hitting the OMI detectors is significantly increased (Dobber et al., 2006), resulting in 302 spikes in the data. It is noted that also the Level 0-1B processor flags transient pixels, so Fig. 4h is the 303 result of the Level 1B flags in combination with the outlier removal procedure. In addition to the SAA, 304 figure 4h also shows stripes in the along-track direction, as well as features related to geophysical 305 conditions (for example higher values of Australia and the India). 306 The impact of the outlier removal procedure was tested by running the algorithm with and without the

procedure switched on for 14 May 2005. The differences in the retrieved effective cloud fraction are negligible, whereas the impact on the effective cloud pressure depends on the cloud fraction. The mean difference is not significant, but the standard deviation of the difference varies for 16 hPa for  $c_f < 0.2$  to 3 hPa for  $c_f < 0.8$ .

We also inspected the root-mean-square error (RMSE) of the DOAS fit as a fit quality indicator.
Although the difference in RMSE with and without the outlier removal did not differ significantly from zero, the distribution is skewed towards larger RMSE values when the outlier removal is switched off.
This indicates that the outlier removal procedure improves the fit for cases with a high RMSE.

#### 315 Digital Elevation Model

316 The version 2 of the algorithm uses a DEM with a resolution of approximately 20 km, which is closer to 317 the spatial resolution of OMI compared to the 3 km resolution DEM used in previous versions. The 20 318 km resolution DEM is constructed from the Global Multi-resolution Terrain Elevation Data 2010 319 (Danielson and Gesch, 2011). 320 The impact of the new DEM will be largest in mountainous terrain. Fig. 8 illustrates the effect on the 321 retrieved effective cloud pressures over Europe for 14 May 2005. This is the same day as shown in Fig. 322 4. Fig. 8a shows that significant impacts of the new DEM are restricted to the main mountain ranges. 323 The difference between using the old and new DEM can be both positive and negative. The impact 324 increases towards the lower cloud fractions, when more signal comes from the surface and an accurate 325 knowledge of the surface altitude becomes more important. Fig 8b shows that for most pixels the impact 326 is smaller than ±50 hPa.

#### 327 Cross sections

In the new version of the algorithm, absorption cross-sections and the Raman radiance spectrum have been updated. The impact of this change was tested by running the algorithm with the old and the new cross sections. The impact on the cloud fraction was negligible. Using the new cross sections increased





the effective cloud fractions by 23±23 hPa. The difference in the root-mean-square error of the DOAS

332 was not significant. The new cross-sections didn't significantly reduce the residuals of the DOAS fit.

#### 333 Scene albedo and scene pressure

334 As described in the algorithm section, for each ground pixel the scene albedo and scene pressure is 335 derived. The most important application of these parameters is over bright surfaces such as snow and 336 ice, where the surface albedo becomes close to the assumed cloud albedo of 0.8 and no meaningful cloud 337 fraction and pressure can be derived. Fig. 9 shows a comparison of the retrieved scene pressure with the 338 surface pressure derived from the DEM, assuming a sea level pressure of 1013 hPa. The figure shows a 339 very good agreement between the retrieved scene pressure and the DEM over Greenland. This figure 340 presents the comparison for the OMI cross track pixel 20, but other cross pixels show similar results. It 341 demonstrates the capabilities of the scene pressure for bright surfaces. Also, it is an indirect validation 342 of the retrieved O<sub>2</sub>-O<sub>2</sub> slant columns. A correction of the O<sub>2</sub>-O<sub>2</sub> slant columns, as is sometimes used in 343 ground based DOAS measurements (for a discussion see (Spinei et al., 2015)), is clearly not necessary 344 for the OMI retrievals. 345

345 Over dark surface, such as oceans, the scene pressure is less well understood. For some areas over the 346 ocean the retrieved scene pressure is significantly larger than the sea level pressure. Therefore, we 347 recommend using the scene albedo and scene pressure only for ground pixels which are covered with

348 snow and/or ice.

#### 349 Comparison with ground-based radar

350 The changes made in the version 2 of the OMCLDO2 algorithm have a stronger impact on the cloud 351 pressure retrieval then on the cloud fraction retrieval. Therefore, we focus in this section on comparisons 352 of the cloud pressure retrievals with correlative data. Because of the use of the IPA cloud model (Fig. 1), 353 it is not straightforward to compare the retrieved cloud pressure to profile information on cloud 354 parameters. As discussed below, we compare the OMI retrievals with ground based radar data, for which 355 the sensitivity to cloud droplet size is very different; the OMI retrievals are sensitive to the optical extinction with scales with droplet size to the power 2, whereas the radar reflectivity scales with droplet 356 357 size to the power 6. Thus, using these radar data it is not possible to compare the same quantity, which 358 is required in a validation study. Rather than a validation study, we focus on explaining the differences between the OMI retrievals and the radar data, given their different sensitivities. This comparison uses a 359 360 similar approach as was used for comparing SCIAMACHY cloud products with radar data (Wang and 361 Stammes, 2014).

- 362 We present comparisons for three sites: Cabuw, The Netherlands, Lindenberg, Germany and the ARM
- 363 Southern Great Plains, U.S.A., for the period January to June 2006. These datasets were selected because

364 of the continuous data availability for these sites in the Cloudnet datatabase.





## 365 Cloudnet data

366	The Cloudnet dataset is the Level 2 classification product (Illingworth et al., 2007), which is available		
367	approximately every 30 seconds. This product classifies each vertical layer as one of 11 classes, which		
368	distinguish ice and water clouds, precipitation, aerosols, insects, clear sky and combinations thereof. We		
369	attribute a value of 1 to layers that are classified as cloudy (classes 1-7) and 0 to layers identified as non-		
370	cloudy. For profiles containing at least one cloudy layer, we compute the cloud mid-height as the average		
371	of the altitude of the cloudy layers. Next we average all the profiles in the time window of +/- $30$ minutes		
372	of an OMI overpass. We also compute the average and standard deviation of the cloud mid-height over		
373	this time window and determine for the average cloud profile if it is single-layer or multi-layer.		
374	It is noted that this procedure for computing the cloud mid-height doesn't take the optical thickness of		
375	the layers into account; a optically thick cloud and optically thin cloud are weighted the same in the cloud		
376	mid-height. Weighting with the optical thickness - or even better, with the sensitivity of the $\mathrm{O_2-O_2}$ cloud		
377	algorithm- would make a comparison much more direct. Unfortunately, information on the full optical		
378	thickness profile is not available from the Cloudnet data. Alternatively, we could use the radar reflectivity		
379	as weighting parameter. However the radar reflectivity is very sensitive to cloud particle size, which is		
380	also not a good representation for the cloud extinction in the visible. We therefore decided to use the		
381	simple weighting described above. This weighting gives the same weight to optically thin cloud layer as		
382	to optically thick layers, whereas the $O_2$ - $O_2$ is cloud pressure retrieval is much more sensitive to the thick		
383	layers.		
384	Further filtering of the Cloudnet data was done using the following criteria:		
385	• The standard deviation of the cloud-mid height should not exceed 1.5 km, to avoid cases with		
386	large temporal variability during the OMI overpass;		
387	• At least one layer in the profile should be cloudy during at least 50% of the time averaging		
388	window.		
389	OMI collocated data		
390	For the OMI cloud data we average all the ground pixels of which the center is within 30 km distance of		
391	the ground station. For these pixels we determine the mean and standard deviation for the cloud fraction		
392	and pressure. We convert the cloud pressure to altitude using a scaling height of 8 km. We filter the OMI		
393	data using the following criteria:		

- The effective cloud fraction should exceed 0.2, because the cloud pressure for low cloud fraction
   has a large uncertainty;
- The standard deviation of effective cloud pressures should not exceed 1.5 km, to exclude cases
   with large horizontal variability.

#### 398 Results

Figure 10 shows a comparison between the Cloudnet data and the OMI effective cloud pressure for the collocations over Cabauw for the period January to June 2006. The cases presented in this figure are ordered by increasing mid-height of ground-based data. The following regimes can be distinguished in this data set:





403	1.	Case 1-50: These are low level clouds with limited vertical extent. The OMI effective cloud
404		height and the ground stations mid-height are in good agreement.

- Case 51-129: According to the Cloudnet the majority of these cases consist of vertically
   extended, and often multi-layered cases. For these cases the OMI effective cloud height is
   generally lower than the ground station mid-height.
- 408
  3. Case 130-135: Theses cases have high clouds with limited vertical extent. The OMI effective
  409
  409 cloud height compares well, except for the outlier for cases 131. However, the number of
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- 411 It is noted that the boundaries of these three regimes are not hard.
- 412 Figure 10 shows that for single layer clouds with a limited vertical extent, the  $O_2$ - $O_2$  effective cloud 413 height and the Cloudnet derived mid-height are in agreement. This shows that the OMI derived product 414 is capable of retrieving cloud height ranging from low clouds to high clouds. For vertically extended 415 clouds, the OMI derived cloud heights are generally lower than the radar-lidar derived heights. A 416 plausible explanation for this difference is that in these cases there are thin high clouds overlaying thicker 417 low-level water clouds. Whereas the radar-lidar mid-heights have equal sensitivity, the O2-O2 cloud 418 height will be more sensitive to the optically thick layers. 419 When we include not only Cabauw, but also Lindenberg and the ARM-SGP site, we get a similar picture. 420 Figure 11 shows a comparison for all these sites for the period January-June 2006, where the single and 421 multi-layer cloud cases are distinguished. Good correlation is observed for the cloud range of 0-2.5 km,
- where the single-cloud layers dominate. In the region between 2.5 and approximately 8 km the multilayer clouds dominate and the  $O_2$ - $O_2$  cloud-height is lower than the radar-lidar cloud mid-height. Above 8 km we find both good comparison but also very large differences, although the number of points is
- very limited. As we are interested in the average comparisons, we did not investigate individual caseswhere big differences occurred.

#### 427 Conclusions

428 We present a new version of the OMI OMCLDO2 Level 2 cloud product. This product is an important 429 input for several of the operational OMI Level 1-2 algorithms. The new version contains six major 430 improvements 431 1. The correction for the temperature sensitivity of the DOAS fit; 432 2. Improved look-up-tables for computing the effective cloud fraction and effective cloud 433 pressure; 434 3. Retrieval of the scene pressure and scene albedo for every ground pixel, using the Lambertian 435 Equivalent Reflector model; 436 4. Outlier removal procedure in the DOAS fit. 437 5. Updated of the reference cross sections; 438 6. Introduction of a DEM with a similar spatial resolution as the OMI ground pixels. 439 We show that the impact of these changes on the retrieved effective cloud fraction is for most ground 440 pixels less than 0.01. The impact on the effective cloud pressure is larger: especially for cloud fractions 441 less than approximately 0.3 the differences compared to the previous operational version can be as large





- 442 as 200 hPa. These differences are mainly caused by the temperature correction and the introduction of
- the new look-up tables. Due to the temperature the differences have a latitudinal and seasonal dependent
- 444 behavior, where the updated algorithm gives higher cloud pressures at higher latitudes and lower
- pressures in the tropics and sub-tropics. Also it was found that the new look-up-tables gives better resultsat low cloud fractions.
- 447 Cloud pressure retrievals have been compared to ground based radar-lidar observations in Cabauw,
- 448 Lindenberg and the ARM-SGP site. It was found that for low clouds, up to approximately 2.5 km, the
- satellite retrievals and ground-based results compare favorably. For clouds in the range between 2.5 and
- 450 approximately 8 km the ground-based observations indicate many multi-layer and vertically extensive
- 451 clouds. For these clouds the satellite retrieved cloud heights are generally lower, probably because the
- algorithm is more sensitive to the optically thick low-level clouds. For high clouds (>8 km) mixed results
  are found. The differences with the lidar-radar can be explained by the different sensitivity of the lidar-
- 454 radar observations versus the satellite observations.
- 455 We conclude that the new version of the OMCLDO2 product is a significant improvement of the previous
- versions, especially for the cloud pressure at cloud fractions smaller than approximately 0.3. This is very
- 457 important for cloud corrections in retrievals of gases like nitrogen dioxide, sulphur dioxide and
- 458 formaldehyde, which are very sensitive to the cloud pressure.
- 459 After the reprocessing of the entire OMI data record, the stability of the product should be investigated,
- 460 and the scene pressure and scene albedo should be validated.
- 461

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- 464 We acknowledge the EU Cloudnet and the ACTRIS projects for providing the cloud classification
- dataset, using Radar-Lidar cloud classification products for the Cabauw, Lindenberg and the ATM-SGPsites.
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532 Tables

533

- 534 Table 1: Nodes for the radiative transfer calculations. Note that cloud fractions smaller than 0 and larger
- 535 than 1 are included to enlarge the parameters space.

Parameter	Nodes
solar zenith angle [°]	0.0, 9.3, 21.2, 32.9, 44.2, 54.9, 64.8, 73.5, 80.8, 86.1
viewing zenith angle [°]	0.0, 9.3, 21.2, 32.9, 44.2, 54.9, 64.8, 73.5
relative azimuth angle [°]	0, 30, 60, 90, 120, 150, 180
surface albedo	0.0, 0.01, 0.025, 0.05, 0.075, 0.1, 0.15, 0.2, 0.25, 0.325,
	0.4, 0.5, 0.6, 0.7, 0.8, 0.9, 1.0
surface/cloud pressure [hPa]	1013, 963, 913, 863, 813, 763, 713, 663, 613, 563, 513,
	463, 413, 363, 313, 263, 213, 163, 113, 63
cloud fraction	-0.1, -0.05, 0., 0.01, 0.02, 0.04, 0.06, 0.08, 0.1, 0.125,
	0.15, 0.175, 0.2, 0.25, 0.3, 0.35, 0.4, 0.45, 0.5, 0.55, 0.6,
	0.65, 0.7, 0.75, 0.8, 0.85, 0.95, 1.0, 1.1, 1.2

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541Table 2: Nodes for the continuum reflectance and the slant Column O2-O2, for the cloud fraction/pressure542and scene albedo/scene pressure look-up-tables. The solar zenith angle, viewing zenith angle, relative azimuth

angle, surface albedo and surface/cloud pressure nodes are the same as given in Table 1.

Parameter	Nodes
Continuum reflectance	0.00, 0.05, 0.10, 0.15, 0.20, 0.25, 0.30, 0.35, 0.40, 0.45,
R <sub>c</sub> at 477 nm	0.50, 0.55, 0.60, 0.65, 0.70, 0.75, 0.80, 0.85, 0.90, 0.95,
	1.00, 1.05, 1.10, 1.15, 1.20, 1.25, 1.50, 1.75, 2.00
Slant Column O <sub>2</sub> -O <sub>2</sub>	0.00, 0.05, 0.10, 0.15, 0.20, 0.25, 0.30, 0.35, 0.40, 0.45,
[10 <sup>44</sup> molec <sup>2</sup> cm <sup>-5</sup> ]	0.50, 0.55, 0.60, 0.65, 0.70, 0.75, 0.80, 0.85, 0.90, 0.95,
	1.00, 1.10, 1.20

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- 547
- 548 Table 3. Impact of the improvements of the effective cloud fraction and effective cloud pressure
- 549 retrievals.
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Improvement	Impact on <i>p<sub>cld</sub></i>	Impact on <i>c</i> <sub>f</sub>
Temperature correction	Decreasing at higher latitudes Increasing in the tropics and sub-tropics $\Delta p_{cid}$ : -100 to 150 hPa for $c_f < 0.2$ $\Delta p_{cid}$ : -20 to 40 hPa for $c_f > 0.2$	negligible
New look-up- tables	Impact is non-significant for $c_f > 0.3$ $\Delta p_{cid}$ : -60 to 220 hPa for $c_f < 0.3$	$\Delta c_f$ : -0.01 except for high surface reflectivity for which $\Delta c_f > 0.05$
Outlier removal	No systematic impact	negligible
DEM	Impact restricted to mountainous terrain. $\Delta p_{cld}$ for most pixels smaller than +/- 25 hPa	negligible
Cross-sections	Δp <sub>cld</sub> : 23 +/- 23 hPa	negligible





# 553 Figures 554 Independent Pixel Approximation Lambertian Equivalent Reflector $p_{cld} \Rightarrow \boxed{A_{cld}} \Rightarrow \boxed{A_{sfc}} \Rightarrow p_{scn} \Rightarrow \boxed{A_{scn}} \xrightarrow{A_{scn}} \Rightarrow \boxed{A_{scn}} \Rightarrow \boxed{A_{s$

Figure 1: The Independent Pixel Approximation versus the Lambertian Equivalent Reflector model. In the IPA a ground pixel is modeled as the weighted sum of a cloudy part, (a Lambertian surface with an albedo of  $A_{cdd}$  at a pressure level  $p_{sfc}$ ) and a clear part (a Lambertian surface with an albedo of  $A_{sfc}$  at a pressure level  $p_{sfc}$ ). The effective cloud fraction  $c_f$  is used for the weighting of the cloudy and clear contributions. The IPA method uses a priori information on  $A_{sfc}$ ,  $A_{cld}$  and  $p_{sfc}$ . In the LER model the ground pixel is modeled as a Lambertian surface with a albedo  $A_{scn}$  at a pressure level  $p_{scn}$ . The LER method doesn't rely on a priori information. Note that the hatched areas below the opaque Lambertian indicate that these regions do not contribute in the radiative transfer calculations.





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Figure 2: Example of a slice of the effective cloud fraction LUT (top panel) and effective cloud pressure LUT (bottom panel), showing the LUT value as a function of the continuum reflectance  $\rho_c$  and the slant column O<sub>2</sub>-O<sub>2</sub>  $N_{s,0202}$ . The background colors show the values in the LUT derived from interpolation and extrapolation of the DOAS fit results, which are shown as the color-filled symbols. The other LUT nodes are fixed to the following values: solar zenith angle 44.2°; viewing zenith angle 21.2°; relative azimuth angle 0.0°; surface albedo 0.05; surface altitude 0 m.







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575 Figure 3: Bias in the retrieved pressure (pretr - ptrue) in hPa when in the retrieval a mid-latitude summer

576 temperature profile is used whereas in the simulation a mid-latitude winter profile (mlw) or a sub-arctic 577 winter profile (saw) is used. The results are plotted as a function of the cloud fraction and for different

578 579 pressure levels of the cloud used in the simulation. The surface albedo is fixed at 0.05, the cloud albedo is 0.80,

the solar zenith angle is 60 degrees and the viewing direction is nadir.







Figure 4: Results from the OMCLDO2 version 2 algorithm for 14 May 2005. a) effective cloud fraction, b)
 effective cloud pressure, c) difference of the effective cloud fraction (version 1.2.3 minus version 2), d)
 difference of the effective cloud pressure (version 1.2.3 minus version 2), e) scene albedo, f) scene pressure, g)
 SCD temperature correction factor γ, and h) number of wavelengths used in the DOAS fit.









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- Figure 5: Box-whisker plot of the precision of the effective cloud pressure as a function of the effective cloud fraction for 14 May 2005. 587
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- 592 Figure 6: Difference in the effective cloud pressure due to the temperature correction (without correction
- minus with correction plotted as function of the effective cloud fraction. The colors of the symbols indicate the SCD correction factor.

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Figure 7: Box-whisker plots of the effective cloud pressure as a function of the effective cloud fraction. The
 top plot is for the old LUTs, the middle for the new LUTs and the bottom plot for the difference of old minus
 new.







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- Figure 8: Difference in the effective cloud pressure (old DEM minus new DEM) for effective cloud fractions exceeding 0.1 over Europe for 14 May 2005. Left panel: map of the differences over Europe, right panel: histogram of the differences over Europe on a logarithmic scale. 601 602
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Orbit 4415, Xtrack 20 Orbit 4415, Xtrack 20 DEM DEM Retrieved



Fig 9: Top panel: map of the position of the ground pixels centers. Bottom panel: comparison of the retrieved scene pressure and the surface pressure derived from the DEM, plotted as a function of the longitude.

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610 Figure 10: The effective cloud altitude retrieved from OMI (red), compared to radar cloud information for 611 Cabauw (blue), The Netherlands. The grey background is the vertically resolved cloud occurrence derived 612 from the radar data for the period +/- 30 minutes of the OMI overpass. The cases are ordered according to 613 the ground station cloud mid-height.







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617 Figure 11: The retrieved effective cloud altitude from OMI, plotted as a function of the radar derived cloud

altitude. Closed symbols are for single-layer clouds, open symbols for multi-layer clouds.