Application of Oxygen A-band Equivalent Width to 1 **Disambiguate Downwelling Radiances for Cloud Optical** 2 Depth MeasurementApplication of Oxygen A-band 3 Equivalent Width for Cloud Optical Depth Measurement 4 5 Edward R. Niple¹, Herman E. Scott¹, John A. Conant¹, Stephen H. Jones¹, Frank 6 J. lannarilli¹, and Wellesley E. Pereira² 7 [1]{Aerodyne Research, Inc., Billerica, MA, USA} 8 9 [2]{Air Force Research Lab, Albuquerque, NM, USA} Correspondence to: F. J. Iannarilli (franki@aerodyne.com) 10 11 12 Abstract 13 This paper presents the Three Waveband Spectrally-agile Technique (TWST) for measuring Cloud Optical Depth (COD). TWST is a portable field-proven sensor and retrieval method 14 15 offering a unique combination of fast (1 Hz) cloud-resolving (0.5 ° field of view) real-timereported COD measurements. It entails ground-based measurement of visible/near-infrared 16 (VNIR) zenith spectral radiances much like the AERONET Cloud-Mode sensors. What is 17 18 novel in our approach is that we employ absorption in the oxygen A-band as a means of 19 resolving the COD ambiguity inherent in using up-looking spectral radiances. We describe the TWST sensor and algorithm, and assess its merits by comparison to AERONET Cloud-20

Mode measurements collected during the Department of Energy, Atmospheric Radiation

<u>Measurements (ARM), Two Column Aerosol Project (TCAP).</u> a new technique for measuring Cloud Optical Depth (COD). It is based on ground based visible band zenith spectral

radiances much like the AERONET Cloud-Mode sensors. What is novel in our approach is

that we employ absorption in the oxygen A band as a means of resolving the COD Ambiguity inherent in using up looking spectral radiances. We describe the algorithm and a sensor that

implements it, and compare its performance to AERONET Cloud Mode measurements collected during the Two Column Aerosol Project (TCAP). Spectral radiance agreement was

excellent (better than 1%) while COD agreement was good.

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1 Keywords

- 2 Cloud optical depth, radiative transfer, oxygen A-band, liquid water path, equivalent width
- 3

4 **1** Introduction

5 <u>1.1 Motivations and contribution</u>

Accurate global climate model (GCM) predictions are absolutely essential. Not only are they 6 7 needed to help mitigate damage from unwanted climate changes, but also to determine what, 8 if any, interventions can reverse those changes directly. Yet clouds, because of their 9 complexity and inherently random nature, remain a primary challenge to GCM accuracy. 10 Direct optical measurements of cloud optical depth (COD) serve a number of atmospheric and 11 climate change science and monitoring purposes. In fundamental cloud process and 12 observational studies, the column-integrated nature of COD imposes a strong calibration 13 constraint on measurements by vertical-profiling instruments such as LIDAR and cloud radar 14 (Kikuchi et al, (2006)). Localized COD measurements serve as concise validation checks on 15 cloud radiative transfer (RT) models. These purposes are often cost-effectively served by ground-based COD instruments. Global distributed, routine time series observations of COD 16 17 are similarly useful in GCM validation. Although such observation grids are routinely available from satellite retrievals from short-wavelength reflectances (Nakajima and King, 18 (1990)), there remains substantial added value from ground-based COD instrument networks, 19 20 such as AERONET (in Cloud Mode). This added value is ascribed to the continuing need for independent validation of satellite-retrieved COD (Liu et al, (2013)), since satellite sensor 21 22 calibrations degrade in-orbit and suffer their own measurement biases. 23 A further use that we have made for ground-based COD measurements is providing ground 24 truth for field experiments of optical contrast propagation through clouds. In these 25 applications we have placed emphasis on: (1) fast measurement rate, e.g. 1 Hz; (2) cloud-26 resolving narrow field of view applicable to scattered clouds as opposed to overcast-only 27 situations; and (3) real-time reporting. As this combination of features was to our knowledge unavailable from existing instruments (circa 2010), we developed the TWST sensor and its 28

29 COD retrieval algorithm.

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1	COD and cloud droplet effective radius r_e together are the minimal required parameters to	
2	determine a Liquid Water Path (LWP), and thus a connection between cloud macro	
3	observables and microphysical parameters. To retrieve COD and r_{e} from short-wavelength	
4	radiances, a longstanding approach (Nakajima and King, (1990)) has been to employ two	
5	wavelengths, one at a non-absorbing (e.g. visible/NIR) and the other at an absorbing	
6	wavelength for liquid water, e.g. longer than 1500 nm. Indeed, the AERONET Cloud Mode	
7	has recently adopted a longer-wavelength channel for this purpose (Chiu et al, (2012)), and	
8	other ground-based optical sensors employ a sufficiently broad spectral range (McBride et al,	
9	(2011), Liu et al, (2013), Fielding et al (2014)). For the TWST sensor's present state of	
10	development, we purposely chose not to operate beyond 1100 nm due to cost and complexity	
11	burdens we wished to avoid. As the primary data product of TWST is COD, in section 3.2.2	
12	we provide evidence of retrieved COD's relative insensitivity to r_e for low to moderate COD	
13	values.	
14	Our TWST technical approach was inspired by the development of AERONET Cloud Mode,	
15	circa 2010 (prior to its employment of the 1640 nm channel), which we summarize next.	
16	1.2 Prior art of COD retrieval from zenith spectral radiance	
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1 1.3 TWST synopsis and paper outline

The TWST sensor is a zenith staring narrow field of view (NFOV) VNIR spectral radiometer 2 3 built around an inexpensive commercial compact grating spectrometer (CGS) with a nominal 4 2.5 nm resolution. The technological sophistication and robustness of the TWST instrument 5 derives almost entirely from its commercial components; we neither depend upon nor make 6 any remarkable claims about sensor design or suitability. In section 2, we describe the TWST 7 sensor and its field-worthiness, present example data, and discuss its calibration, including 8 dark current correction. In section 4.2.1, within our measurements section, we establish 9 TWST radiometric veracity and stability by comparison to coincident AERONET spectral 10 radiance observations over a period of several weeks. 11 Although we customarily record the full spectral record spanning about 350-1000 nm, the TWST COD retrieval presently makes use of a sparse set of spectral bins composed into 3 12 13 spectral factors: the spectral radiances at 440 nm and 870 nm (SR440 and SR870) and the 14 equivalent width (EQW) of the oxygen A-band centered near 760 nm. TWST is not using A-15 Band spectrometry to retrieve a numerical COD. Like AERONET Cloud Mode, TWST 16 employs model-generated look-up tables of spectral radiance to COD. In particular, TWST 17 tables relate SR440 to COD. TWST differs from Cloud Mode in its resolution of the 18 aforementioned COD ambiguity. Cloud Mode (circa 2010) employs a 2-dimensional ordinate 19 space involving the sum and difference of SR440 and SR870, which is bijective to pairing of 20 COD and effective cloud fraction. TWST instead first determines the cloud optical thickness 21 regime, Thin or Thick, and thus whether to reference the thin or thick branch of the SR440 to 22 COD look-up table. The novelty of TWST is its determination of thickness regime from a 23 multivariate temporal filter employing a color index (the SR440/SR870 ratio) and the slope of the plot of SR440 versus A-Band EQW. TWST development status presently circumscribes it 24 25 to the regime of low-moderate altitude water clouds and small-moderate solar zenith angles (SZA). Defining TWST's precise operational boundaries is a future task. 26 27 In section 3.1 we discuss the TWST COD retrieval, including its physical basis, COD error 28 sensitivity to primary uncertainties, implementation details, and an explanatory example of 29 the retrieval technique operating on a time sequence of sensor data. In section 4, in addition

- 30 to the aforementioned spectral radiance comparison, we establish TWST COD retrieval
 31 accuracy by comparing to coincident AERONET Cloud Mode COD. We present and discuss
- 32 the quality of agreement in both an illustrative several-hour time series and the cumulative

1	correlation over a several week period of the Two Column Aerosol Project (TCAP) field
2	campaign.
3	2 The TWST sensor
4	2.1 Design and characteristics
5	The heart of the TWST sensor is a zenith-pointing calibrated spectroradiometer. We elected
6	to design the sensor around a commercial compact grating spectrometer (CGS), given the
7	significant advances in miniaturization, rugged monolithic construction, and linear array
8	detectors. Several advantages accrue from this design choice, the most important to our COD
9	measurement application being the acquisition of spectral radiances at high signal-to-noise-
10	ratio (SNR) and high temporal resolution, attributed to the multiplex advantage provided by
11	the CGS.
12	The key specifications for the TWST COD sensor are listed in Table 1. Here we are excluding
13	extreme ambient conditions outside the range of -10 C to +40 C that require special
14	temperature control. Our design has proven to be field worthy, easily transportable, and
15	stable over a wide range of environmental conditions. We have experienced no instances of
16	eondensation inside the sealed TWST enclosure while operating in humidity and temperature
17	eonditions well below the dew point. The spectral resolution defined by the spectrometer
18	configuration is currently ~2.5 nm (including convolution with slit function) and the sampling
19	interval is ~0.3nm. With integral order-sorting filter, the stray light level is cited to be $<0.1\%$.
20	The temporal resolution determined by the available sunlight, spectrometer throughput, and
21	linear FPA detector characteristics is a variable sampling interval from 1 msec to 60 sec (1 sec
22	typical). The relative precision of the TWST COD measurements determined from field data
23	is estimated as follows. Each rawA typical TWST spectrum recorded at 1 sec interval consists
24	of 400 co-added snapshots scans each of 2.5 msec integration time. The SNR for a single
25	snapshots scan-is limited to 400:1 due to photon-electron noise based on the electron well
26	depth of 160,000. When readout noise is included, this drops to 275:1. With 400 co-adds, the
27	<u>1 second maximum-signal SNR is therefore limited to about 5,500:1.</u>
28	We have now built and tested a few different configurations of the TWST sensor, but each
29	includes the basic elements represented in the schematic design in Figure 1; a companion
30	photograph looking inside a recent model is shown in Figure 2. The entrance window (A) is

31 slanted to shed rain drops. A simple mechanical shutter (S) for recording dark spectra,

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1	selected for its reliability and effective light blocking performance, is located just inside the
2	input window and driven by an inexpensive stepping motor. An incoming light baffle (B)
3	limits the total field-of-view (TFOV) to 0.5 ° FWHM. At the end of the baffle sits a 400 nm
4	long pass filter. A collecting lens (C) then focuses the filtered light onto the end of a 400 µm
5	dia. optical fiber (D) which feeds the light into the CGS (E). The entire system is contained in
6	an IP66 (NEMA 4X) rated sealed enclosure with desiccant to prevent water condensation
7	over deployment periods of several months. Our design has proven to be field-worthy, easily
8	transportable, and stable over a wide range of environmental conditions as supported in
9	section 4.2.1. We have experienced no instances of condensation inside the sealed TWST
10	enclosure while operating in humidity and temperature conditions well below the dew point.
11	2.2 Example spectral data
12	The TWST retrieval algorithm uses three spectral factors (Figure 3): the spectral radiances at
13	440 nm and 870 nm (SR440 and SR870) and the equivalent width (EQW) of the oxygen A-
14	band (section 3.3.1). Figure 3 shows example calibrated spectral measurements for nearly
15	identical SZA, but for clear sky and a range of COD values in the Thin optical thickness
16	regime. The overall radiance level as well as the depth of the oxygen A-band absorption are

17 observed to increase with COD.

18 2.3 Calibration and dark current correction

19 There are two forms of calibration that must be managed for any technique that uses spectroradiometers: wavelength and radiometric. The wavelength calibration of the compact grating 20 21 spectrometers used in our TWST sensors has proven stable over periods of months. 22 Furthermore, the TWST approach does not rely on resolving spectral line structure. The 440 23 nm and 870 nm spectral radiance levels, due to their shallow spectral slopes (Figure 3), and 24 A-Band EQW value, due to its accumulation over many spectral bins, are relatively 25 insensitive to foreseeable thermal shifting of the spectral sampling grid. 26 TWST spectral radiance calibration is performed at the beginning and end of every field

deployment and more frequently as needed. Our calibration source is a Labsphere Uniform
Radiance Standard integrating sphere. It is well-known that the standard incandescent source
lamps age and need to be replaced periodically. Like other long time users, we find these
lamps to be the largest source of uncertainty and absolute error in our radiometric calibration

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1 procedure; that uncertainty is ~5%. During each calibration we set the integration period, 2 number of snapshot coadds, and aperture radiance to span the range of field conditions 3 anticipated for sunlit clouds, and then we derive a linear photoresponsivity coefficient in the 4 usual manner. These radiometric calibration records for each TWST unit are kept and 5 compared over periods of years to monitor the stability of each unit for its lifetime. Having 6 records for some units over 2-4 years, we find changes in the calibration of 1-3%, well within 7 the uncertainty of our calibration lamps which as noted above is on the order of 5%. 8 The spectrometer's silicon CCD detector outputs are susceptible to offset drift, typically 9 driven by changes in ambient temperature, but the detector array contains light-shielded dark-10 reference detectors intended to automatically track and subtract such drift. In addition, TWST 11 employs a mechanical shutter for frequent collection of dark spectra, typically a 1-second 12 dark spectrum every 60 seconds. The dark correction (offset subtraction) applied to each 13 recorded spectrum is spline-extrapolated from earlier collected dark spectra. In section 4.2.1, 14 the effectiveness of these calibration methods is evaluated by comparison to coincident 15 AERONET spectral radiance observations over a period of several weeks.

16 <u>3 TWST cloud optical depth retrieval method</u>

17 3.1 Section outline

The TWST retrieval algorithm employs model-generated look-up tables to convert zenith 18 19 spectral radiance at 440 nm (SR440) to a numerical value of COD. However, the TWST 20 algorithm first determines the cloud optical thickness regime, Thin or Thick, and thus whether 21 to reference the thin or thick branch of the SR440 to COD look-up table. We first discuss the 22 somewhat conventional spectral radiance to numerical COD lookup, including table 23 generation, COD error sensitivity to principal uncertainties via radiative transfer simulations, 24 and technique of interpolation between table entries. Then we discuss the determination of 25 the optical thickness regime. This entails discussion of why the Oxygen A-band and its 26 equivalent width (EQW) metric are informative of the thickness regime. We introduce the 27 "nose" plot of SR440 vs EQW, and its generic slope characteristics are revealed as a key to 28 resolving the thin/thick ambiguity; the algorithm does not use *model-generated* nose plots. 29 We explain the need and basis for the SR440/SR870 ratio as color index. The color index and 30 nose plot slope metrics are combined in a multivariate temporal filter that continually updates

the estimate of optical thickness regime. To illustrate its operation, and how it copes with 3D
 cloud effects, we discuss an example nose plot time sequence and filtered results.

3 3.2 Numerical spectral radiance to COD lookup

4 1.1.13.2.1 Radiative transfer construct

5 **2** Principle of Operation

The general measurement consists of a zenith staring spectral radiometer recording the 6 7 absolute power coming from a narrow field of view (FOV) in a collection of spectral bands. 8 Assuming the sensor's field of view (FOV) does not include the sun, the zenith spectral 9 radiance this consists of solar radiation scattered by the molecules, aerosols, and cloud 10 droplets in the FOV, which may include radiation that has been scattered multiple times from the atmosphere and the terrain. The visible/NIR spectral band (Figure 3), at a moderate 11 12 spectral resolution of 2 nm, shows a broad baseline with multiple narrow absorption features. Many of these are due to water vapor, as well as Fraunhofer linesconstituents of the sun. The 13 14 spectral radiance at 440 nm is in a region relatively free from atmospheric gaseous absorption, 15 and thus suitable as a COD proxy. We chose 440 nm for TWST radiance-to-COD look-upthis effort because that is a wavelength used by AERONET Cloud-Mode sensors, which serve as a 16 17 source for comparative validation. The model used for generating 440 nm radiance-to-COD look-up tables is theis MODTRAN5 18

10	The model used for generating 440 min radiance-to-COD look-up tables is the MODTRANS		
19	atmospheric radiative transfer code (Berk et al., 2006). MODTRAN incorporates the DISORT		
20	code (Stamnes et al, 1988) for plane-parallel stratified media, i.e., idealized 1-dimensional		
21	radiative transfer (1DRT). Calculations are done for a typical water stratus cloud above a		
22	stated ground albedo, for a stated nominal aerosol profile, over a rangegrid of COD and solar		
23	zenith angles (SZA). Figure 4 is a graphical depiction of sample tables. For any solar zenith		
24	angleSZA, there is a "bright point" radiance where the idealized 1DRT cloud radiance		
25	isreaches a maximum, typically occurring for a COD between 2 and 8, as seen in Figure 4.		
26	Real clouds manifest 3-dimensional radiative transfer (3DRT) effects, including radiances		
27	exceeding the idealized 1DRT bright point radiances (Marshak et al, (2000)). When faced		
28	with such exceedances, the TWST retrieval algorithm reports the COD corresponding to the		
29	bright point radiance but flags an out-of-bounds ("3D Cloud") condition.		

1 3.2.2 COD error sensitivity to radiative transfer parameter uncertainties

2 The TWST algorithm currently operates without any information on the droplet size 3 distribution or the cloud base height, and with a prior estimate of the ground albedo and aerosol loading profile. We do not consider deviation of the actual from nominal aerosol 4 5 profile, as such perturbation from the baseline aerosol optical depth (AOD) is typically a 6 small contributor to reported COD. We performed some initial sensitivity studies on these 7 remaining parameters. The albedo sensitivity findings below will prove of value in helping to 8 explain the minor disagreement bias between coincident TWST and AERONET COD 9 observations (section 4.2.2). 10 Water cloud drop size distributions typically vary from an effective radius of 1 - 20 µm (see 11 e.g. Chiu et al. (2006)). Because our implementation of the TWST algorithm uses a radiance 12 database generated with the MODTRAN model, we studied 440 nm radiances from four 13 different cloud types included parameterized within MODTRAN, which assume Mie 14 scattering, log-normal droplet size distribution, and liquid water refractive index. These types 15 have effective radii of 12.0 (cumulus), 7.2 (altostratus), 8.3 (stratus), and 6.7 (stratocumulus) 16 <u>µm.</u> Water cloud drop-size distributions typically vary from an effective radius of $1 - 20 \mu m$ 17 (see e.g. Chiu et al, (2006)). We modeled clouds with fixed base height of 0.5 km and fixed physical thickness of 0.5 km. For each cloud type we varied the Liquid Water Path (LWP) 18 19 enough to achieve 550 nm CODs between 0 and 100; LWP was used because it is an input to 20 MODTRAN. COD values were estimated from LWP using the Wood and Hartmann, (2006) modification to the Stephens, (1978) formula as described in Chiu et al, (2012). Figure 5 21 22 shows the computed 440 nm vertical radiances plotted versus LWP and COD. The different size distributions lead to different radiances for the same LWP as expected. However, in the 23 24 COD vs. radiance plot the curves overlay closely, at least for COD<=20. We verified this 25 intuitive result using a simple standalone two-stream computation as an alternative to the 26 DISORT algorithm included within MODTRAN (Stamnes et al, 1988). These results are 27 corroborated by the more extensive sensitivity results of McBride et al (2011). This relative 28 insensitivity of COD with effective radius relation-gives us some confidence in reporting 29 COD values in face of the variety of water cloudsfrom radiances at the non-absorbing 440 nm 30 wavelength. 31 Inter-reflections between the ground and a thick cloud can be significant unless the Earth

32 <u>albedo is low. The TWST algorithm uses a look up table of radiance computations for 440</u>

1 mm. At the 440 nm look-up table that-wavelength most Earth cover types have albedo of 0.2 or less as shown in Figure 6 with samples of the ASTER Database (Baldridge et al. 2009); 2 3 notable exceptions are for white sand, fresh snow, and ocean ice. To characterize the 4 sensitivity of retrieved COD due to albedo uncertaintyvariability, we used MODTRAN to 5 compute 440 nm zenith radiances for a low-altitude stratus cloud for albedos of 0.0, 0.1, 0.2 6 and 0.5, over a range of solar zeniths and cloud optical thicknesses. Figure 7 presents a first-7 order indication of sensitivity, and plots 440 nm radiances versus COD for albedo bounds of 8 0.0 and 0.5. The thin/thick ambiguityduality, the strong variation with solar zenith, and the 9 weaker variation with ground albedo are evident in these plotted results. 10 The plot in Figure 8(a) further explores this sensitivity, and shows the absolute signed change 11 in retrieved COD value for an unexpected increase in the ground albedo from 0.1 (for which 12 the radiance-to-COD lookup tables are computed) to 0.2. Each curve, for either thick or thin 13 cloudplotted versus SZA, pertains to some fixed is for a different percentage of the aforementioned 1DRT bright point radiance "Lbrt" (which varies with SZA, cf. Figure 4), for 14 15 either thick or thin cloud. Note that we used linear interpolations of our data table in 16 computing the CODs for this study, but near the bright point (red curves in Figure 7) the COD 17 vs radiance is quite non linear, and thus the red curves are less precise, and jagged. (The 18 TWST retrieval algorithm doesn't suffer this limitation, as it employs non-linear spline 19 interpolation).- These curves show that a higher than expected albedo implies retrieval of a 20 higherlower (lowerhigher) COD in the Thick (Thin) region. This is because an elevated 21 albedo increases the cloud radiance, but decreases its percent brightness relative to the higher 22 bright point radiance (Figure 6). The largest change we found was $\triangle COD$ of 5, but that 23 occurred for very thick clouds, such that the relative change was only 10%. Near the bright point (red curves in Figure 8(a)) the COD vs radiance is quite non-linear (Figure 4), and thus 24 25 the red curves, approximated by linear interpolations for this study, are less-precise, and 26 jagged. 27 To provide analytic support to these albedo sensitivity findings, we performed calculations 28 employing asymptotic radiative transfer (ART) theory relations as elucidated by King, (1987) 29 and Melnikova et al, (2000). Both MODTRAN and our ART calculations compute the 440 30 nm radiance and optical thickness for cloud with phase function asymmetry parameter g=0.86

31 and single scattering albedo $\omega_0 > 0.9999997$ (i.e. conservative scattering). ART-computed

32 <u>sensitivities are plotted in Figure 8(b), and compare well in both trend and magnitude against</u>

1	the Thick regime curves of Figure 8(a), to which ART theory pertains (here COD>9). These
2	ART sensitivities are processed from the more directly obtained ART calculations plotted in
3	Figure 9. The connection between Figure 9 and Figure 8(b) is depicted by the dotted path
4	shown in Figure 9. The reader is directed to graphically determine the error in retrieved COD
5	value by first starting with a given true COD value, tracing rightward from that y-intercept
6	parallel to the x-axis and reaching a curve pair for a given SZA. The right curve of the pair
7	registers the actual radiance measured for the unexpected 0.2 albedo, but the left curve is
8	radiance-to-COD lookup table computed for the expected 0.1 albedo. So at intersection with
9	right curve, the reader traces down parallel to the y-axis to intercept the left curve, then traces
10	leftward back to the y-axis and reads out the lower COD value. The solid curves of Figure
11	8(b) are not exactly comparable to those of Figure 8(a), in that the former are for constant-
12	COD whereas the latter are for constant relative radiance (with respect to the 1DRT bright
13	point radiance). To corroborate that the COD error in fact decreases with SZA, the dashed
14	curves of Figure 8(b) better (but not exactly) correspond to constant relative radiance. These
15	dashed curves have COD decrease with decreasing SZA, which referring to Figure 4 yields a
16	more stationary relative brightness (downward-sloping line marker) than for COD constant
17	with SZA (horizontal line marker).
18	At this point it is important to note that the spectral radiance chosen at some other wavelength
19	than 440 nm could be used for the radiance-to-COD look-up. In principle the choice of
20	wavelength depends upon freedom from atmospheric gaseous absorption and on the ground
21	albedo of the measurement site. One wants a wavelength with the lowest absolute albedo
22	uncertainty to minimize errors in the COD due to errors in the assumed albedo for the
23	MODTRAN5 computations. This flexibility in the choice of wavelength is basis of the term
24	"Spectrally-agile" within the TWST acronym.
25	For low-altitude water clouds, uncertainty in the cloud base height has a negligible impact on
26	COD retrieval. We ran MODTRAN for cloud base heights of 500 m and 2 km, iterating over
27	10 COD and 11 solar zeniths for each. The 440 nm radiances were nearly identical, as shown
28	in the scatter plot of Figure 10.
29	3.2.3 Lookup table interpolation

30 <u>Various table lookup algorithms were investigated. A to reach a reasonable tradeoff between</u>
 31 <u>accuracy and speed. The MODTRAN5 tables are preprocessed as follows. Referring to</u>

Comment [fji3]: This paragraph is largely original text, which should have been shown as Moved.

1	Figure 4, for each SZA _i entry, a cubic spline curve SR440(COD;SZA _i) is fit to its SR440-vs-		
2	COD table for both optical thickness regimes. A bright point radiance vs SZA spline curve		
3	Lbrt(SZA), depicted by the dotted black curve in Figure 4, is fit through the bright point		
4	radiances across the SZA _i entries. During operation, the algorithm identifies the tabulated		
5	SZA _j closest to the current solar zenith angle SZA _{obs} . Then a working copy of its spline curv		
6	SR440(COD;SZA _j) is linearly scaled in radiance so that its bright point matches Lbrt(SZA _{obs} ,		
7	This scaled curve is then used to look up the COD value for the measured SR440 is		
8	combination of 1 D spline interpolations and searching for the tabulated solar zenith angle		
9	(SZA) closest to the current solar zenith angle based on the current time, day of the year and		
10	latitude/longitude of the sensor. The first step involves preprocessing the MODTRAN5 results		
11	to find the COD (CODmax) of the brightest SR440 (SR440max) for each SZA. These are		
12	then used to compute spline coefficients for SR440max as a function of SZA. Spline		
13	coefficients are then computed for both optical states for each COD as a function of SR440.		
14	All these spline coefficients are computed once the MODTRAN5 data is generated and then		
15	stored for use during instrument operation.		
16	<u>During operation, the measured spectral radiance $L_{obs}(\theta_{obs})$ at 440 nm is scaled to a</u>		
17	<u>corresponding L_{sc} at the tabulated solar zenith angle θ_{table} nearest the current solar zenith</u>		
	angle θ_{obs} according to		
18	angle θ_{obs} according to		
18 19			
19	$L_{se} = \frac{L_{abs}(\theta_{abs}) L_{tabte}(\theta_{tabte})}{L_{max}(\theta_{abs})} $ (1)		
19 20	$L_{se} = \frac{L_{obs}(\theta_{obs}) L_{table}(\theta_{table})}{L_{max}(\theta_{obs})} $ (1) where $L_{table}(\theta_{table})$ is the tabulated spectral radiance and $L_{max}(\theta_{obs})$ is the interpolated		
19 20 21	$L_{sc} = \frac{L_{obs}(\theta_{obs}) L_{tabte}(\theta_{tabte})}{L_{max}(\theta_{obs})} $ (1) where $L_{tabte}(\theta_{tabte})$ is the tabulated spectral radiance and $L_{max}(\theta_{obs})$ is the interpolated maximum spectral radiance at the current solar zenith angle. The corresponding COD is then		
19 20	$L_{se} = \frac{L_{obs}(\theta_{obs}) L_{table}(\theta_{table})}{L_{max}(\theta_{obs})} $ (1) where $L_{table}(\theta_{table})$ is the tabulated spectral radiance and $L_{max}(\theta_{obs})$ is the interpolated		
19 20 21	$L_{sc} = \frac{L_{obs}(\theta_{obs}) L_{tabte}(\theta_{tabte})}{L_{max}(\theta_{obs})} $ (1) where $L_{tabte}(\theta_{tabte})$ is the tabulated spectral radiance and $L_{max}(\theta_{obs})$ is the interpolated maximum spectral radiance at the current solar zenith angle. The corresponding COD is then		
19 20 21 22	$L_{se} = \frac{L_{obs}(\theta_{obs}) L_{table}(\theta_{table})}{L_{max}(\theta_{obs})} $ (1) where $L_{table}(\theta_{table})$ is the tabulated spectral radiance and $L_{max}(\theta_{obs})$ is the interpolated maximum spectral radiance at the current solar zenith angle. The corresponding COD is then computed by spline interpolation of L_{sei}		
 19 20 21 22 23 	$L_{ge} = \frac{L_{abs}(\theta_{obs}) L_{cabte}(\theta_{tabte})}{L_{max}(\theta_{obs})} $ (1) where $L_{tabte}(\theta_{tabte})$ is the tabulated spectral radiance and $L_{max}(\theta_{obs})$ is the interpolated maximum spectral radiance at the current solar zenith angle. The corresponding COD is then computed by spline interpolation of L_{ges} 3.3 Optical thickness regime determination		
 19 20 21 22 23 24 	$L_{ge} = \frac{L_{obs}(\theta_{obs}) L_{table}(\theta_{table})}{L_{max}(\theta_{obs})} $ (1) where $L_{table}(\theta_{table})$ is the tabulated spectral radiance and $L_{max}(\theta_{obs})$ is the interpolated maximum spectral radiance at the current solar zenith angle. The corresponding COD is then computed by spline interpolation of L_{ge1} 3.3 Optical thickness regime determination 3.3.1 Cue from Oxygen A-band equivalent width (EQW)		
 19 20 21 22 23 24 25 	$L_{ge} = \frac{L_{abs}(\theta_{obs}) L_{cabte}(\theta_{tabte})}{L_{max}(\theta_{obs})} $ (1) where $L_{tabte}(\theta_{tabte})$ is the tabulated spectral radiance and $L_{max}(\theta_{obs})$ is the interpolated maximum spectral radiance at the current solar zenith angle. The corresponding COD is then computed by spline interpolation of L_{ses} . 3.3 Optical thickness regime determination 3.3.1 Cue from Oxygen A-band equivalent width (EQW). The well-known Θ Oxygen A-band atcentered near 760 nm, already referred to as "the well-		
 19 20 21 22 23 24 25 26 	$L_{see} = \frac{L_{obs}(\theta_{obs}) L_{cabte}(\theta_{tabte})}{L_{max}(\theta_{obs})} $ (1) where $L_{tabte}(\theta_{tabte})$ is the tabulated spectral radiance and $L_{max}(\theta_{obs})$ is the interpolated maximum spectral radiance at the current solar zenith angle. The corresponding COD is then computed by spline interpolation of L_{ges} 3.3 Optical thickness regime determination 3.3.1 Cue from Oxygen A-band equivalent width (EQW) The well-known Θ oxygen A-band atcentered near 760 nm, already referred to as "the well- known atmospheric band system" by (Mulliken, (1928), Wark and Mercer, (1965)) when he		
 19 20 21 22 23 24 25 26 27 	$L_{ge} = \frac{L_{obs}(\theta_{obs}) L_{tabte}(\theta_{tabte})}{L_{max}(\theta_{obs})} $ (1) where $L_{tabte}(\theta_{tabte})$ is the tabulated spectral radiance and $L_{max}(\theta_{obs})$ is the interpolated maximum spectral radiance at the current solar zenith angle. The corresponding COD is then computed by spline interpolation of L_{ge} : 3.3 Optical thickness regime determination 3.3.1 Cue from Oxygen A-band equivalent width (EQW) The well-known Θ xygen A-band atcentered near 760 nm, already referred to as "the well- known atmospheric band system" by (Mulliken, (1928), Wark and Mercer, (1965)) when he first published the interpretation of the bands, has been used for many years to study the		
 19 20 21 22 23 24 25 26 27 28 	$L_{ge} = \frac{L_{obs}(\theta_{obs}) L_{cubte}(\theta_{cubte})}{L_{max}(\theta_{obs})} $ (1) where $L_{rabte}(\theta_{rabte})$ is the tabulated spectral radiance and $L_{max}(\theta_{obs})$ is the interpolated maximum spectral radiance at the current solar zenith angle. The corresponding COD is then computed by spline interpolation of L_{ges} . 3.3 Optical thickness regime determination 3.3.1 Cue from Oxygen A-band equivalent width (EQW) The well-known ρ Oxygen A-band atcentered near 760 nm, already referred to as "the well- known atmospheric band system" by (Mulliken, (1928), Wark and Mercer, (1965)) when he first published the interpretation of the bands, has been used for many years to study the atmosphere from satellite and ground-based sensors. Wark and Mercer, (1965) first proposed		

1 lengths for skylight transmitted from clear and cloudy skies to the ground, following the 2 suggestion of Pfeilsticker et al, (1996) and Harrison and Min, (1997). Our technique is merely 3 one more application of the A band; in this case, to help determine COD but one particular 4 feature at 760 nm is the well known O₂ A band. The A-band is virtuallyThis feature is free 5 from absorption by other atmospheric constituents (Pfeilsticker et al, (1998)) except for 6 aerosol and cloud continua extinction plus a very small amount of line absorption by water 7 vapor (Figure 11). Thus its continua-normalized (section 3.3.2) spectral-average quantity, termed the Equivalent Width (EQW), and provides a direct measurement of the average 8 9 amount of oxygen-density-weighted photon path length from the sun to the sensor. Since 10 oxygen is uniformly mixed in the atmosphere, this is related to the photon's' physical path 11 lengths. Therefore the EQW supplies useful information about whether a zenith radiance 12 measurement is in the optically-thin or optically-thick regime. A virtue of EQW is that it may 13 be stably computed from low-resolution spectral data such as from the TWST sensor, as 14 detailed in section 3.3.2. 15 Of course, other factors cause EQW to change besides COD. Changes in the SZA produce 16 decreases in EQW with time during the morning and increases in the afternoon. Changes in 17 the density-weighted average cloud thickness and cloud altitude also affect the EQW 18 independent of the COD. 19 3.3.2 Calculation of EQW 20 Our algorithm uses three spectral factors: the spectral radiances at 440 nm and 870 nm 21 (SR440 and SR870) and the equivalent width of the oxygen A band (EQW). The equivalent 22 width is computed from the spectral radiances between 750 to 785 nm by fitting a straight line 23 to the continuum baseline from 750 - 760 nm and 770 - 785 nm (Figure 11), then, dividing 24 each measured spectral radiance by the corresponding linear fit baseline to produce a 25 transmittance value, and then summing these values acrossintegrating the area under the absorption band. This calculation normalizes away the continuum transmittance. This requires 26 an accurate determination of the zero radiance level. The veracity of these calculations 27 28 depends on accurate spectrometer dark current calibration and subtraction (discussed in 29 section 2.3). Otherwise, a dark bias of the spectral radiance would falsely alter the computed 30 transmittances and EQW value.

1 <u>3.3.3 "Nose" plot of SR440 vs EQW</u>

2 COD is a two-valued function of up-looking spectral radiance while oxygen equivalent width 3 is a monotonic function of COD for COD~>1., so the EQW supplies useful information about 4 whether a measurement is in the optically thin or optically thick regimes. By plotting SR440 5 versus EQW as COD increases from no cloud to thick clouds one traces out a "nose-like" 6 shape (Figure 12). For the very lowest COD values, EQW decreases with increasing COD and 7 the slope is negative. Beyond about COD=1, <u>The lower portion of the "nose" where the slope</u> 8 is positive corresponds to the optically thin regimestate; the upper portion where the slope is negative corresponds to the optically thick regimestate. Within a span of several seconds, 9 10 realpassing clouds most often do not trace out a complete nose but only a small segment of it as the cloud evolves and drifts in the wind over the sensor. Notwithstanding the lowest COD 11 12 values discussed further below, Wwhether the cloud changes involve increases or decreases of 13 the COD, the slope of the corresponding segment indicates the cloud's optical thickness 14 regimestate. This is the TWST basis for the new algorithm as it allows us to resolvinge the 15 COD ambiguity. 16 The nose plot in Figure 12 includes a smooth curve, based on MODTRAN5 computations but 17 elastically stretched to fit the depicted data points over a 4 minute measurement where the 18 COD varied strongly between the indicated blue sky, thin and thick regimes, as well as points deviating well away from the ideal 1DRT smooth curve. These deviating points are classified 19 as either "3D Cloud" based on their SR440 exceedance of the 1DRT bright point radiance 20 value (section 3.2.1) or as "Mixed" points attributed to heterogeneous cloud structure within 21 the field of view, itself a 3DRT effect. The classification of the remaining data points into 22 23 optically thin, thick or blue sky regimes was corroborated against coincident all sky camera 24 video. Although the MODTRAN5 computed nose plot curve supports these regime 25 classifications, it is important to note that the TWST algorithm does not employ model-26 generated nose plot curves to guide its thickness regime determination. Indeed, the particular 27 shape and slope of a computed nose plot curve varies, as it should, with the unknown physical 28 cloud thickness. Instead, the algorithm exploits the aforementioned positive-slope:Thin, 29 negative-slope: Thick generic properties of the nose plot.

1 3.3.4 Thickness regime filter

2 The cloud optical thickness regime determination operates in two distinct radiance domains. 3 When the COD is very low, e.g. COD<=1, the amount of radiation in the NIR is very small and the signal to noise ratio (SNR) of the equivalent widthEQW is low (e.g. "clear sky" in 4 5 Figure 3) and thus the nose plot slope SNR is too low to resolve the thickness regimestate. 6 The ratio of SR440 to SR870, termed the color indexwhich we call the Blueness factor, has a 7 much higher SNR and is more reliable in this regime. This, of course, is a simple consequence 8 of the wavelength dependence of Rayleigh scattering. For low-moderate altitude water clouds, 9 small-moderate SZA, and typical 440 nm ground albedos less than 0.2 (Figure 6), our data 10 analyses We have found a hard threshold of 54 isto be a sure indicator of optically very thin clouds (e.g. "clear sky" in Figure 3) and a soft threshold of 2<index<4 to be a strong indicator 11 12 (e.g. "COD<1" in Figure 3). When the color index Blueness factor is less than 2, the cloud's 13 optical thickness state-is not well correlated with the color index Blueness factor, and wethe 14 algorithm must rely on the nose plot slope. 15 Figure 13 re-depicts the nose plot data points of Figure 12, this time connecting a subset of 16 points with line segments to indicate adjacent samples in a 2-minute time series. If measured 17 nose plots followed an idealized 1DRT curvewere as pristine as indicated in Figure 12, the determination of thickness regime state-would be nearly trivial. A linear regression over a 18 19 short time segment would suffice. Clearly more complex logic is required, yet it is visually evident that local coherence could be exploited. However, measured nose plots deviate from 20 the ideal due largely to time varying structure in the field of view but also to temporal 21 22 variation of the "bright point" (the "nose tip"). Figure 3 shows a typical nose plot for over 13 23 hours of data with a 10 minute segment highlighted. In the top plot, it is not too difficult to 24 see that the 10 minute segment lies in a positive slope region and thus is indicative of an optically thin cloud. However, this assessment takes into account a great deal of context that 25 26 is not evident if the 10 minute segment is examined by itself as can be seen in the bottom part 27 of the figure. Clearly, a simple linear regression would yield highly erroneous results. It 28 should be noted here that the signal to noise ratio of the measurements is very high, often 29 >1000. Consequently, increased integration time will not improve the situation. 30 The qualitative reduction in ambiguity afforded by examining a sufficient the much larger time 31 record suggests the use of a filter with memory. For example, for passing or evolving clouds 32 spatially well-resolved within a narrow field of view, the thickness regime should we know

1 that clouds do not switch rapidly between the thick and thin states except possibly near the 2 bright point (thick/thin regime boundary) where thea switch is inconsequential to the retrieved 3 CODanyhow. We implemented a time varying hysteresis filter to effectively avoid this 4 unwanted switching. The hysteresis action is achieved by keeping track of the maximum and 5 minimum values of equivalent width over a predetermined time interval, typically about three 6 minutes. In order to drive the filter toward a different state, the hysteresis limits must be 7 exceeded. The output of the hysteresis filter is discrete ternary: -1, 0 or 1, corresponding to 8 thick, indeterminate or thin. Finally, this ternary variable is input to a linear, single pole 9 autoregressive (AR(1)) filter to afford additional smoothing. The output of this filter is 10 thresholded and used as the thickness regime state estimate.

3.3.5 Example operation of thickness regime filter 11

12 Figure 13 includes time series plots of SR440 and algorithm retrieved COD corresponding to 13 the connected data points in the nose plot. The dark- and medium- blue points indicate thick-14 and thin-cloud, while light cyan points indicate clear sky. Yellow points are for SR440 values 15 greater than the 1DRT bright point value and are indicative of 3D cloud effects. The green 16 aforementioned Mixed points (section 3.3.3) were determined manually, and are those where 17 the values deviated strongly from the overall nose curve locus; red line segments serve to 18 indicate where along the time series those deviations start and end. Using an instrument with 19 small field-of-view and fast time-response we expect to see good temporal coherence in the 20 data, and in fact the radiance time plot shows that the cloud optical thickness regime does not 21 change randomly. One can see that the green-labeled points are always transitions between thin and thick cloud which didn't follow the idealized nose curve through the bright point (at 22 a radiance of about 23). The EQW values of those points are reasonable, but the radiances are 23 lower than expected. Our supposition at this time is that some of these are due to further 3D 24 25 cloud structure effects, but leading to darker radiances rather than the bright radiances of the 26 yellow points, while others points may be due to spatial averaging over the field-of-view.

- 27
- 28

4 Measurements and comparisons to AERONET

29 The spectral radiances at 440 nm and 870 nm are in regions relatively free from atmospheric absorption. These three factors are then used to determine the COD: all three to determine 30 31 cloud state (thick, thin, or blue sky), and SR440 for numerical COD by comparison to model

predictions. The model used is MODTRAN5 (Berk et al., 2006) for a typical water stratus
 cloud over a range of COD and solar zenith angles.

3 3 Prior Art

The use of zenith spectral radiances to measure COD began with the work of Marshak et al. 4 (2004) using a technique first suggested by Marshak et al, (2000) and Barker and Marshak, 5 (2001). In order to resolve the COD ambiguity they use spectral radiances at red (670 nm) 6 7 and NIR (870 nm) wavelengths which sit on opposite sides of the chlorophyll red edge feature of the albedo of vegetated terrain. The technique was validated at the ARM site in Oklahoma 8 9 by comparison to more conventional techniques (Microwave Radiometer and Multifilter Rotating Shadowband Radiometer). In the work of Chiu et al, (2006) the preliminary 10 11 validation was extended. In Chiu et al, (2010) the technique was improved by switching to 12 blue (440 nm) and NIR (870 nm) wavelengths.

13 The oxygen A-band at 760 nm, already referred to as "the well-known atmospheric band 14 system"-by Mulliken, (1928) when he first published the interpretation of the bands, has been used for many years to study the atmosphere from satellite and ground based sensors. Wark 15 and Mercer, (1965) first proposed using the A-band to study cloud top heights from space and 16 17 Pfeilsticker et al. (1998) first used the A-band to study the probability density function of geometrical path lengths for skylight transmitted from clear and cloudy skies to the ground, 18 19 following the suggestion of Pfeilsticker et al. (1996) and Harrison and Min, (1997). Our 20 technique is merely one more application of the A-band; in this case, to help determine COD.

21 4 "Nose" Plot

COD is a two-valued function of up-looking spectral radiance while oxygen equivalent width 22 23 is a monotonic function of COD, so the EQW supplies useful information about whether a measurement is in the optically-thin or optically-thick regimes. By plotting SR440 versus 24 25 EQW as COD increases from no cloud to thick clouds one traces out a "nose like"-shape (Figure 2). The lower portion of the "nose"-where the slope is positive corresponds to the 26 optically thin state; the upper portion corresponds to the optically thick state. Within a span of 27 28 several seconds, real clouds do not trace out a complete nose-but only a small segment of it-as the cloud evolves and drifts in the wind over the sensor. Whether the cloud changes involve 29 30 increases or decreases of the COD, the slope of the corresponding segment indicates the 31 eloud's optical state. This is the basis for the new algorithm as it allows us to resolve the COD 32 ambiguity.

1	5 The New Three Waveband Spectrally-agile Technique (TWST) Algorithm
2	Our new algorithm involves two steps. The first is to determine the optical state of the cloud.
3	We do this in two different regimes. When the COD is very low, the amount of radiation in
4	the NIR is very small and the signal to noise ratio of the equivalent width and thus the nose
5	plot slope is too low to resolve the state. The ratio of SR440 to SR870, which we call the
6	Blueness factor, has a much higher SNR and is more reliable. This, of course, is a simple
7	consequence of the wavelength dependence of Rayleigh scattering. We have found a
8	threshold of 5 is a sure indicator of optically thin clouds. When the Blueness factor is less
9	than 2, the cloud's optical state is not well correlated with the Blueness factor, and we must
10	rely on the nose plot slope.
11	If measured nose plots were as pristine as indicated in Figure 2, the determination of state
12	would be nearly trivial. A linear regression over a short time segment would suffice.
13	However, measured nose plots deviate from the ideal due largely to time varying structure in
14	the field of view but also to temporal variation of the "bright point" (the "nose tip"). Figure 3
15	shows a typical nose plot for over 13 hours of data with a 10 minute segment highlighted. In
16	the top plot, it is not too difficult to see that the 10 minute segment lies in a positive slope
17	region and thus is indicative of an optically thin cloud. However, this assessment takes into
18	account a great deal of context that is not evident if the 10 minute segment is examined by
19	itself as can be seen in the bottom part of the figure. Clearly, a simple linear regression would
20	yield highly erroneous results. It should be noted here that the signal to noise ratio of the
21	measurements is very high, often >1000. Consequently, increased integration time will-not
22	improve the situation.
23	The qualitative reduction in ambiguity afforded by examining the much larger time record
24	suggests the use of a filter with memory. For example, we know that clouds do not switch
25	rapidly between the thick and thin states except possibly near the bright point where the
26	switch is inconsequential anyhow. We implemented a time varying hysteresis filter to
27	effectively avoid this unwanted switching. The hysteresis action is achieved by keeping track
28	of the maximum and minimum values of equivalent width over a predetermined time interval,
29	typically about three minutes. In order to drive the filter toward a different state, the
30	hysteresis limits must be exceeded. The output of the hysteresis filter is discrete ternary: -1, 0

31 or 1, corresponding to thick, indeterminate or thin. Finally, this ternary variable is input to a

1	linear, single pole autoregressive (AR(1)) filter to afford additional smoothing. The output of	
2	this filter is thresholded and used as the state estimate.	
3	Once the cloud's optical state is determined, the COD is obtained from a table of	
4	MODTRAN5 computations of SR440 at a grid of COD and solar zenith angles. Various table	
5	lookup algorithms were investigated. A reasonable tradeoff between accuracy and speed is a	
6	combination of 1-D spline interpolations and searching for the tabulated solar zenith angle	
7	(SZA) closest to the current solar zenith angle based on the current time, day of the year and	
8	latitude/longitude of the sensor. The first step involves preprocessing the MODTRAN5 results	
9	to find the COD (CODmax) of the brightest SR440 (SR440max) for each SZA. These are	
10	then used to compute spline coefficients for SR440max as a function of SZA. Spline	
11	coefficients are then computed for both optical states for each COD as a function of SR440.	
12	All these spline coefficients are computed once the MODTRAN5 data is generated and then	
13	stored for use during instrument operation.	
14	During operation, the measured spectral radiance $L_{obs}(\theta_{obs})$ at 440 nm is scaled to a	
15	corresponding $L_{\rm res}$ at the tabulated solar zonith angle $\theta_{\rm result}$ nearest the current solar zonith	
15	corresponding and in the moduled solar zenitri unste of and the current solar zenitri	
15 16	angle θ_{abb} according to	
16	angle O _{gpp} according to	
16 17	$\frac{\text{angle } \theta_{\text{obs}} \text{ according to}}{L_{\text{max}}(\theta_{\text{obs}})} $ (1)	
16 17 18	$angle \theta_{abbc} according to$ $L_{abbc} = \frac{L_{abbc}(\theta_{cabbc}) L_{cabbc}(\theta_{cabbc})}{L_{max}(\theta_{cabc})}, \qquad (1)$ where $L_{cabbc}(\theta_{cabbc})$ is the tabulated spectral radiance and $L_{max}(\theta_{cabc})$ is the interpolated	
16 17 18 19	angle θ_{abg} according to $L_{abg} = \frac{L_{abg}(\theta_{abg}) L_{abg}(\theta_{abg})}{L_{max}(\theta_{abg})}, \qquad (1)$ where $-L_{capto}(\theta_{capto})$ is the tabulated spectral radiance and $-L_{max}(\theta_{abg})$ is the interpolated maximum spectral radiance at the current solar zonith angle. The corresponding COD is then	
16 17 18 19 20	angle θ_{abb} according to $I_{tree} = \frac{L_{cos}(\theta_{coss}) L_{caste}(\theta_{caste})}{L_{max}(\theta_{coss})},$ (1) where $-L_{caste}(\theta_{caste})$ is the tabulated spectral radiance and $-L_{max}(\theta_{coss})$ is the interpolated maximum spectral radiance at the current solar zenith angle. The corresponding COD is then computed by spline interpolation of L_{cost} .	
 16 17 18 19 20 21 	angle θ_{abb} according to $L_{abb} = \frac{L_{abb}(\theta_{abb}) L_{cubic}(\theta_{abb})}{L_{abb}}$ (1) where $-L_{cubic}(\theta_{cubic})$ is the tabulated spectral radiance and $-L_{max}(\theta_{abb})$ is the interpolated maximum spectral radiance at the current solar zenith angle. The corresponding COD is then computed by spline interpolation of L_{abb} . At this point it is important to note that the spectral radiance chosen at some other wavelength	
 16 17 18 19 20 21 22 	angle θ_{abb} according to $I_{abb} = \frac{L_{abb}(\theta_{abb}) L_{abb}(\theta_{abb})}{L_{abb}(\theta_{abb})}$ (1) where $-L_{capte}(\theta_{cabb})$ - is the tabulated spectral radiance and $-L_{max}(\theta_{abb})$ - is the interpolated maximum spectral radiance at the current solar zenith angle. The corresponding COD is then computed by spline interpolation of L_{abb} . At this point it is important to note that the spectral radiance chosen at some other wavelength than 440 nm could be used for the radiance to COD look up. We chose 440 nm for this effort	
 16 17 18 19 20 21 22 23 	angle θ_{gags} according to $\frac{1}{L_{gags}} = \frac{L_{gags}(\theta_{gags}) + L_{gags}(\theta_{gags})}{L_{gags}}$ (1) where $-L_{gagss}(\theta_{gags})$ is the tabulated spectral radiance and $-L_{gags}(\theta_{gags})$ is the interpolated maximum spectral radiance at the current solar zonith angle. The corresponding COD is then computed by spline interpolation of L_{ggs} . At this point it is important to note that the spectral radiance chosen at some other wavelength than 440 nm could be used for the radiance to COD look up. We chose 440 nm for this effort because that is a wavelength used by AERONET Cloud Mode sensors, which serve as a	
 16 17 18 19 20 21 22 23 24 	angle θ_{abb} according to $L_{abb} = \frac{L_{abb}(\theta_{abb}) L_{abb}(\theta_{abb})}{L_{abb}}$ (1) where $-L_{cable}(\theta_{cabb})$ is the tabulated spectral radiance and $-L_{max}(\theta_{abb})$ is the interpolated maximum spectral radiance at the current solar zenith angle. The corresponding COD is then computed by spline interpolation of L_{gg} . At this point it is important to note that the spectral radiance chosen at some other wavelength than 440 nm could be used for the radiance to COD look up. We chose 440 nm for this effort because that is a wavelength used by AERONET Cloud Mode sensors, which serve as a source for comparative validation. In principle the choice of wavelength depends on the	
 16 17 18 19 20 21 22 23 24 25 	angle θ_{abb} according to $\frac{1}{L_{get}} = \frac{L_{abb}(\theta_{abb}) \cdot L_{abal}(\theta_{abb})}{L_{abal}(\theta_{abb})}$ (1) where $L_{L_{abb}}(\theta_{abb})$ is the tabulated spectral radiance and $L_{max}(\theta_{abb})$ is the interpolated maximum spectral radiance at the current solar zenith angle. The corresponding COD is then computed by spline interpolation of L_{ge} . At this point it is important to note that the spectral radiance chosen at some other wavelength than 440 nm could be used for the radiance to COD look up. We chose 440 nm for this effort because that is a wavelength used by AERONET Cloud Mode sensors, which serve as a source for comparative validation. In principle the choice of wavelength depends on the ground albedo of the measurement site. One wants a wavelength with the lowest absolute	

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1 6 Sensitivity Studies

2	The TWST algorithm currently operates without any information on the droplet size			
3	distribution or the cloud base height, and with a prior estimate of the ground albedo. We			
4	performed some initial sensitivity studies on these parameters.			
5	Water cloud drop size distributions typically vary from an effective radius of $1-20 \ \mu m$ (see			
6	e.g. Chiu et al, (2006)). Because our implementation of the TWST algorithm uses a radiance			
7	database generated with the MODTRAN model, we studied 440 nm radiances from four			
8	different cloud types included with MODTRAN. These types have effective radii of 12.0			
9	(cumulus), 7.2 (altostratus), 8.3 (stratus), and 6.7 (stratocumulus) µm. We modeled clouds			
10	with fixed base height of 0.5 km and fixed physical thickness of 0.5 km. For each cloud type			
11	we varied the Liquid Water Path (LWP) enough to achieve 550 nm CODs between 0 and 100;			
12	LWP was used because it is an input to MODTRAN. COD values were estimated from LWP			
13	using the Wood and Hartmann, (2006) modification to the Stephens, (1978) formula as			
14	described in Chiu et al, (2012). Figure 4 shows the computed 440 nm vertical radiances			
15	plotted versus LWP and COD. The different size distributions lead to different radiances for			
16	the same LWP as expected. However, in the COD vs. radiance plot the curves overlay			
17	elosely. We verified this intuitive result using a simple two stream computation as an			
18	alternative to the DISORT algorithm included within MODTRAN (Stamnes et al, 1988). This			
19	relation gives us some confidence in reporting COD values from radiances at the non-			
20	absorbing 440 nm wavelength.			
21	Inter-reflections between the ground and a thick cloud can be significant unless the Earth			
22	albedo is low. The TWST algorithm uses a look-up table of radiance computations for 440			
23	nm. At that wavelength most Earth cover types have albedo of 0.2 or less as shown in Figure			
24	5 with samples of the ASTER Database (Baldridge et al, 2009); notable exceptions are for			
25	white sand, fresh snow, and ocean ice. To characterize the sensitivity of retrieved COD due to			
26	albedo variability, we used MODTRAN to compute 440 nm zenith radiances for albedos of			
27	0.0, 0.1, 0.2 and 0.5, over a range of solar zeniths and cloud thicknesses. Figure 6 presents a			
28	first-order indication of sensitivity, and plots 440 nm radiances versus COD for albedo			
29	bounds of 0.0 and 0.5. The thin/thick duality, the strong variation with solar zenith, and the			
30	weaker variation with ground albedo are evident in these plotted results.			
31	The plot in Figure 7 shows the absolute change in retrieved COD for an unexpected increase			
32	in the ground albedo from 0.1 to 0.2. Each curve, plotted versus SZA, is for a different			

1	percentage of the bright point radiance "Lbrt" (which varies with SZA), for either thick or
2	thin cloud. For any solar zenith angle, there is a "bright point" where the cloud radiance is a
3	maximum, typically occurring for a COD between 2 and 8, as seen in Figure 6. Note that we
4	used linear interpolations of our data table in computing the CODs for this study, but near the
5	bright point (red curves in Figure 7) the COD vs radiance is quite non-linear, and thus the red
6	eurves are less-precise, and jagged. (The TWST retrieval algorithm doesn't suffer this
7	limitation, as it employs non-linear spline interpolation). These curves show that a higher
8	than expected albedo implies retrieval of a higher (lower) COD in the Thick (Thin) region.
9	This is because an elevated albedo increases the cloud radiance, but decreases its percent
10	brightness relative to the higher bright point radiance (Figure 6). The largest change we found
11	was ACOD of 5, but that occurred for very thick clouds, such that the relative change was
12	only 10%.
13	For low altitude water clouds, uncertainty in the cloud base height has a negligible impact on
14	COD retrieval. We ran MODTRAN for cloud base heights of 500 m and 2 km, iterating over
15	10 COD and 11 solar zeniths for each. The 440 nm radiances were nearly identical, as shown
16	in the scatter plot of Figure 8.
17	7 The TWST Sensor
18	The heart of the TWST sensor is a zenith pointing calibrated spectroradiometer. We elected
19	to design the sensor around a commercial compact grating spectrometer (CGS), given the
20	
	significant advances in miniaturization, rugged monolithic construction, and linear array
21	significant advances in miniaturization, rugged monolithic construction, and linear array detectors. Several advantages accrue from this design choice, the most important to our COD
21 22	
	detectors. Several advantages accrue from this design choice, the most important to our COD
22	detectors. Several advantages accrue from this design choice, the most important to our COD measurement application being the acquisition of spectral radiances at high signal-to-noise-
22 23	detectors. Several advantages accrue from this design choice, the most important to our COD measurement application being the acquisition of spectral radiances at high signal-to-noise- ratio (SNR) and high temporal resolution, attributed to the multiplex advantage provided by
22 23 24	detectors. Several advantages accrue from this design choice, the most important to our COD measurement application being the acquisition of spectral radiances at high signal-to-noise- ratio (SNR) and high temporal resolution, attributed to the multiplex advantage provided by the CGS.
22 23 24 25	detectors. Several advantages accrue from this design choice, the most important to our COD measurement application being the acquisition of spectral radiances at high signal-to-noise- ratio (SNR) and high temporal resolution, attributed to the multiplex advantage provided by the CGS.
22 23 24 25 26	detectors. Several advantages accrue from this design choice, the most important to our COD measurement application being the acquisition of spectral radiances at high signal-to-noise- ratio (SNR) and high temporal resolution, attributed to the multiplex advantage provided by the CGS. We have now built and tested a few different configurations of the TWST sensor, but each includes the basic elements represented in the schematic design in Figure 9; a companion

32 long pass filter. A collecting lens (C) then focuses the filtered light onto the end of a 400 µm

input window and driven by an inexpensive stepping motor. An incoming light baffle (B)

limits the total field-of-view (TFOV) to 0.5° FWHM. At the end of the baffle sits a 400 nm

1 dia. optical fiber (D) which feeds the light into the CGS (E). The entire system is contained in 2 an IP66 (NEMA 4X) rated sealed enclosure with desiceant to prevent water condensation 3 over deployment periods of several months. 4 The key specifications for the TWST COD sensor are listed in Table 1. Here we are excluding extreme ambient conditions outside the range of -10 C to +40 C that require special 5 6 temperature control. Our design has proven to be field-worthy, easily transportable, and 7 stable over a wide range of environmental conditions. We have experienced no instances of condensation inside the sealed TWST enclosure while operating in humidity and temperature 8 9 conditions well below the dew point. The spectral resolution defined by the spectrometer 10 configuration is currently - 2.5 nm. The temporal resolution determined by the available sunlight, spectrometer throughput, and linear FPA detector characteristics is a variable 11 12 sampling interval from 1 msee to 60 see (1 see typical). The relative precision of the TWST COD measurements determined from field data is 13 estimated as follows. Each raw TWST spectrum consists of 400 co-added scans each of 2.5 14 msee integration time. The SNR for a single scan is limited to 400:1 due to photon noise 15 16 based on the electron well depth of 160,000. When readout noise is included, this drops to 275:1. With 400 co adds, the 1 second SNR is therefore limited to 5,500:1. This is the 17

relative precision of TWST COD results. Even higher precision can be achieved by
 combining the 125 independent spectral channels in the 425 475 nm blue spectral band.
 This yields a potential SNR of 60,000:1 for each 1 second spectral radiance value. In terms of
 COD SNR, for the optically thin state, the COD ranges from 0 to about 6 in an approximately
 linear relationship. This implies a theoretical noise limited sensitivity of 6/60,000 or OD of
 0.0001.

24 4.1 The Two Column Aerosol Project (TCAP)

TCAP was a one year measurement campaign directed by the Atmospheric Radiation Measurements (ARM) division of the U.S. Department of Energy. It was designed to quantify aerosol properties, radiation, and cloud characteristics, producing a database to assist climate modeling studies. The ground-based campaign involved the ARM Mobile Facility (AMF) suite of sensors deployed at the ARM Highlands in Cape Cod Massachusetts. The aerial campaign involved two aircraft loaded with remote and in situ sensors. Measurements were performed from July 2012 until June 2013.

With the kind permission and assistance of the TCAP project, the TWST sensor was set up on Cape Cod near the AERONET Cloud-Mode sensor and Total Sky Imager (TSI) which are part of AMF on 17 May 2013. Data were collected continuously for a period six days. Some minor adjustments were then made to the sensor configuration, and then data were collected for the next thirty days until 27 June when the AMF was taken down in preparation for its next deployment. During this period about 50,000 spectra were collected by TWST every day at one second intervals during the day.

8 7.14.2 AERONET Cloud-Mode and TWST Data data Comparison comparison

During the time TWST was deployed on Cape Cod, AERONET collected 266 COD values
that overlapped TWST measurements. In addition, 8,609 overlapping spectral radiance values
at 440 nm were collected. Since all TWST's COD values were based on SR440
measurements, it is important to compare the SR440 values before comparing the COD
values.

14 7.1.1<u>4.2.1</u> Spectral Radiance radiance Comparison comparison

This required careful time synchronization between the AERONET and TWST data times. A linear drift of 0.27 seconds/day was determined by least squares fits to the individual days with a 4 second difference between the high gain (x8) A and low gain (x1) K measurements from AERONET. The result (Figure 14) shows that both sensors were reporting SR440 values in very good agreement. The rms difference was 0.63 (μ W cm⁻² sr⁻¹ nm⁻¹). A simple linear fit without constant yielded a slope of 1.003 (0.0004). TWST values at high spectral radiance showed some evidence of nonlinear response.

22 Several conclusions follow from the very good agreement among TWST and AERONET 23 spectral radiances. The first is the expectation of a COD comparison not influenced by TWST spectral radiance errors. The second is the radiometric stability of TWST during its TCAP 24 25 deployment. As a corollary, the COD comparison should not be unduly influenced Similarly 26 by different fields of view (1.2 ° for AERONET versus 0.5 ° for TWST) and zenith pointing 27 (robotic control for AERONET versus fixed tripod with bubble level for TWST), such-given the close agreement over many different cloud conditions. The sensors were laterally 28 29 displaced by about 3 m, and for a 1 km cloud base altitude their field of view footprints are 20 m and 8 m. implies that the different fields of view (1.2 ° for AERONET versus 0.5 ° for 30 TWST) and zenith pointing (robotic control for AERONET versus fixed tripod with bubble 31

level for TWST) were not significantly different. Of course, the agreement only proves
 consistency, not accuracy for either sensor. The second is the radiometric stability of TWST
 during its TCAP deployment. This is corroborated by the stability of the 4 pre- and 1 post-test
 radiometric calibrations, with the photoresponsivity coefficient at 440 nm for 9 July being
 98.1% of that for 17 May.

6 7.1.24.2.2 COD Comparison comparison

7 The comparison of the average COD (Figure 12) shows evidence of the two different types of 8 errors in the TWST and AERONET Cloud Mode algorithms: errors in cloud state and errors 9 in COD. A comparison of COD values between TWST and AERONET Cloud-Mode must 10 recognize the time-sampling differences between them. For theseis comparisons only the 90 11 second average COD was available for AERONET Cloud-Mode, which is a form of trimmed 12 mean based on up to ten instantaneous COD measurements during each measurement period 13 (see Chiu et al, (2010) [sec.2.3]).

14 A time series comparison of COD is shown in Figure 15. The agreement indicates that the 15 TWST thickness regime filter is able to track the rapidly changing COD. The ensemble 16 comparison of the average-COD values (Figure 16) shows evidence of the two different types 17 of errors in the TWST and AERONET Cloud-Mode algorithms: errors in cloud thickness 18 regime state-and errors in numerical COD. To attempt a comparison, each plot data point 19 represents the average of the Since TWST produces 90 instantaneous COD measurements 20 produced by TWST during the same 90 second period for AERONET., the trimmed mean process will have different statistical properties. 21

22 Of the 244 overlapping COD values, 235 (96%) showed the same cloud thickness 23 regimestate. Some analysis was done in an effort to determine whether TWST or AERONET Cloud-Mode was probably correct. For the nine cases where AERONET and TWST 24 25 disagreed on the thickness regime cloud state, detailed nose plots were generated to see if we 26 could visually extract more than the simple slope information used in the automated 27 algorithm. Four of the cases produced close to the ideal nose shape, indicating that the TWST 28 thickness regimecloud state was probably correct. For the other five cases, the nose plot was 29 too distorted to determine the thickness regimecloud state, indicating that the TWST thickness 30 regimecloud state was probably incorrect and should have been assigned the "unknown" 31 statelabel.

1 A linear fit of TWST to AERONET Cloud-Mode COD for the 235 cases of thickness 2 regimecloud state agreement, for fixed zero-intercept, found an average bias slope of 0.843 3 (TWST reporting higher COD values)for TWST with an rms difference of OD 3.2. This was 4 repeated while dropping the two high COD value outlier points (Figure 16), but the slope only 5 changed by 1%. No evidence of a constant offset between TWST and AERONET Cloud-Mode was found. Although However, the sparsity of such evidence is due to the relatively 6 7 few optically thin COD cases were-available from AERONET, due to the secondary mission status of its Cloud-Mode. No evidence of a constant offset between TWST and AERONET 8 9 Cloud Mode was found. (When skies are largely clear, AERONET executes its primary 10 mission of aerosol optical depth and microphysical property retrieval measurements). Therefore, another linear fit, this time with free-intercept, found a slope of 0.905 and constant 11 12 offset of -2.1. 13 The two primary candidates for causing the observed disagreements are differences in the

14 TWST and AERONET Cloud-Mode lookup tables and effects from the trimmed mean 15 process. There may also be some residual effects due to FOV and pointing differences, 16 although these are not expected to be large due to the very good spectral radiance agreement 17 (section 4.2.1). A partial explanation centering on the lookup tables is the difference in 18 assumed ground albedos between the sensors. The TWST SR440-to-COD lookup table 19 generated from MODTRAN used a weighted average of water, deciduous vegetation, dead 20 pine, and sand albedos, resulting in an earth albedo at 440 nm of 0.078545. On the other 21 hand, AERONET updates its ground albedo episodically every few days from MODIS data 22 products or a (rolling) 16-day average MODIS historical database (Chiu et al, (2012)). For this dataset, the AERONET-employed albedos were lower than that assumed for TWST, 23 varving between 0.02-0.04, 0.03 average. Most of the sample points are in the optically Thick 24 25 regime, and according to our albedo sensitivity discussion (section 3.2.2), a lower than expected albedo implies TWST retrieval of higher COD values in the Thick regime, 26 27 consistent with the linear fits. Figure 8 depicts an approximately constant relative COD retrieval error of about 10% per 0.1 albedo increment. The 0.05 average difference in 28 29 assumed albedo therefore explains about half (0.05) of the difference between a slope of unity 30 and the fitted slope (0.905).

<sup>It must always be kept in mind that the COD values determined by TWST and AERONET
Cloud-Mode are only equivalent 1D-cloud<u>RT</u> quantities. Violations of 1D assumptions are</sup>

present in nearly all our measurements to some extent. This includes (1) cases where the observed spectral radiance is greater than that possible for any 1D cloud, (2) cases where deviations from the ideal <u>Nose-nose Plot-plot</u> are too high for any 1D cloud and (3) the many cases where the rapid variation in spectral radiance is too high for 1D clouds.

5 85 Conclusions

6 Overall the good agreement between TWST and AERONET Cloud-Mode Cloud Optical 7 Depths values, across many weeks of coincident field deployment, validates the performance 8 of the Three Waveband Spectrally-agile Technique as well as the field-worthiness of the 9 TWST sensor. Although the spectrally-agile nature of TWST was not investigated in this 10 study, its advantage over fixed spectral bands for cases with high albedo at 440 nm will-may 11 be the focus of a future study. Although our error sensitivity studies in section 3.2.2 and the 12 agreement with AERONET over many weeks of the TCAP campaign lend confidence in 13 applying TWST for nominal conditions, future efforts will ascertain and, where possible, extend the operational limits (e.g., SZA, ground albedo) of the TWST retrieval algorithm. 14

One of the most notable results of our experience with TWST is the high signal-to-noise ratio available in the high temporal (1 second), spatial (0.5 ° field of view), and spectral (2.5 nm) resolution data TWST generates. At peak signal, <u>at a COD value of approximately 5</u>, the SNR <u>due to instrumental and photon noise</u> is estimated to be 5,000:1 <u>for 1 Hz reports</u>. This peak <u>occurs at a COD value of approximately OD 5</u>, so the implied rms COD error is OD 0.001.

The version of the TWST algorithm used in this report was limited to use of the slope of the Nose Plot. This is mostly the result of limitations in our current ability to model the variation of EQW with COD, cloud base altitude, cloud average physical depth, cloud phase function, and solar zenith angle. If these effects can be normalized out, either from fitting to TWST measurements or to auxiliary measurements (lidar, radar, etc.), we expect that more accurate cloud states can be determined and, therefore, more accurate COD. In addition the values of the variables determined by fitting TWST data provide additional atmospheric information.

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1 Table 1 TWST Cloud Optical Depth Sensor Specifications

TWST Cloud OD Sensor Specifications	
For Ambient Temperature Range -10°C to +40°C	
Weight	20 lbs
Power and Communication for Optical Head	5 Vdc, <250 mA via a single USB 2.0 cable connection to computer for power and data
Size	11 x 8 x 8 in plus 12 in external sun baffle; or 13 x 10 x 6 in with internal sun baffle
Operating Range	Blue Sky to Cloud OD 100
Cloud OD Sensitivity	0.001 for Optically Thin Clouds
Weatherproof Environmental Container	IP66, NEMA 4X sealed enclosure with desiccant
Data Logging Rate	1 Hz (typical), variable sampling interval from 0.1 to 60 seconds
Field of View	0.5 °
Spectral Range, Resolution	350 – 1000 nm, ~2.5 nm
Spectral Bands currently used in Cloud OD retrieval	440, 760, and 870 nm

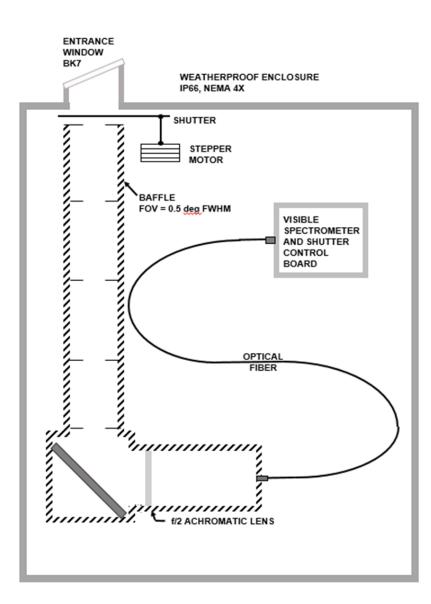


Figure 1 Simplified schematic of the TWST Cloud Optical Depth Sensor.

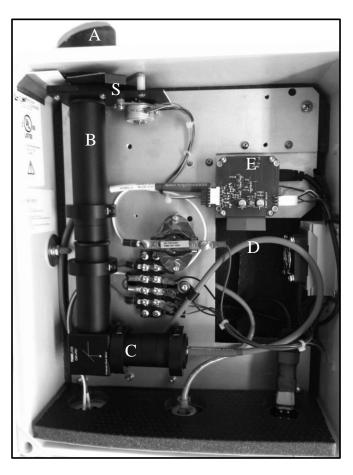
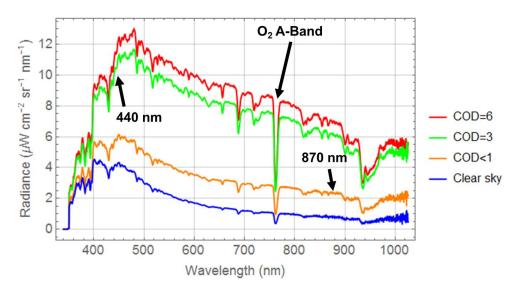


Figure 2 A view inside the TWST Cloud Optical Depth Sensor. See text for labeled component descriptions.





3 Figure 3. Example spectra measured by TWST sensor, delineating the three spectral factors

- 4 currently used in TWST retrieval algorithm. Spectra measured at SZA=65 $^{\circ}$ within 1 minute.
- 5

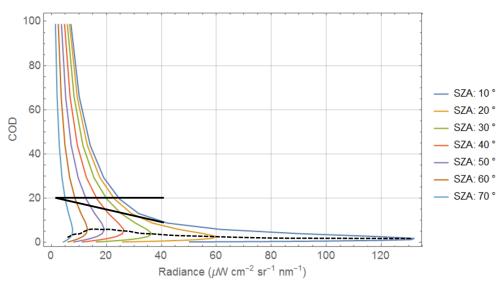
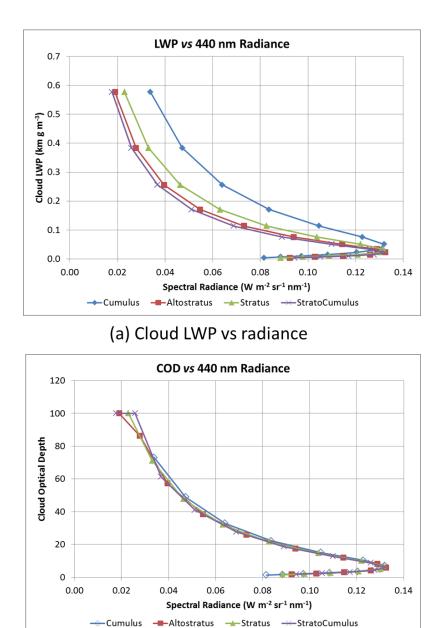


Figure 4 440 nm radiance to COD look-up tables for various SZA. The radiance peak for each curve is its 1DRT "bright point" radiance. The black solid and dotted line markers are referred to in various sections of the text.



(b) Cloud OD vs radiance

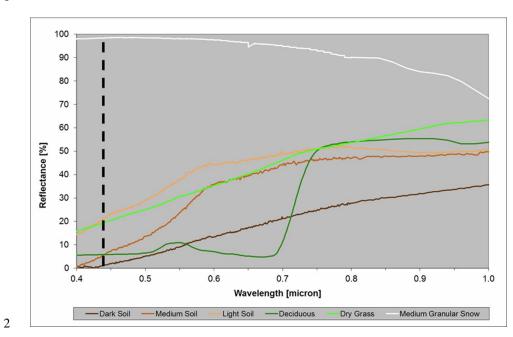


2 Figure 5. Relationship between spectral radiance at 440 nm wavelength and (a) liquid water

3 path (LWP) and (b) cloud optical depth, for four different cloud types: effective radii of 12.0,

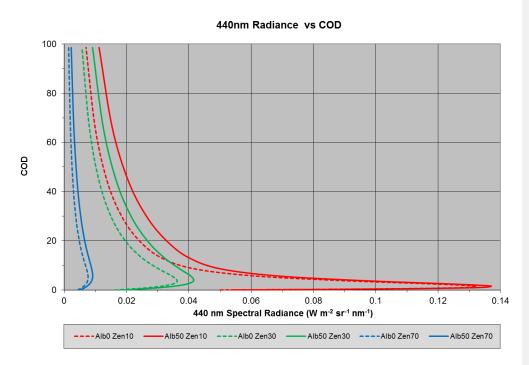
4 $\,$ (cumulus), 7.2 (altostratus), 8.3 (stratus), and 6.7 (stratocumulus) $\mu m.$

5





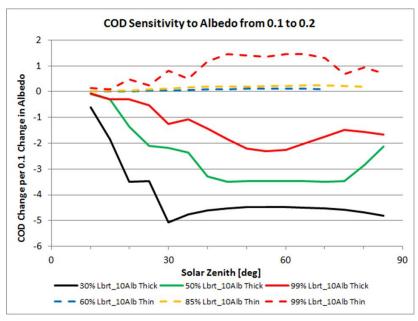
4 vertical dashed line indicates the 440 nm wavelength.





2 Figure 7. Cloud OD vs 440 nm spectral radiance for solar zenith angles of 10 $^\circ$, 30 $^\circ$ and 70 $^\circ$

- 3 and for ground albedos of 0 and 50%.
- 4



(a) Retrieved COD sensitivity: MODTRAN

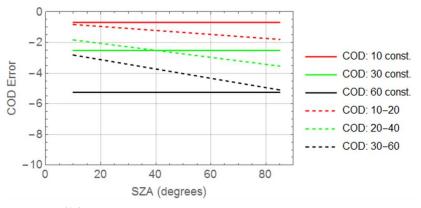
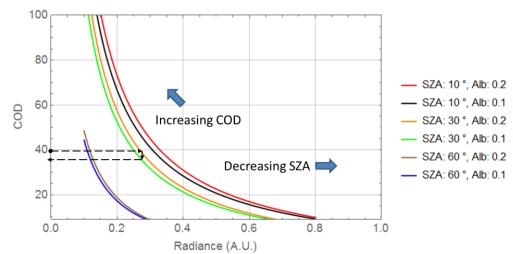




Figure 8. Retrieved COD sensitivity to change in albedo from 0.1 to 0.2 (a) for different
relative cloud radiance levels as computed using MODTRAN, and (b) as computed using
asymptotic RT relations. See text for details.



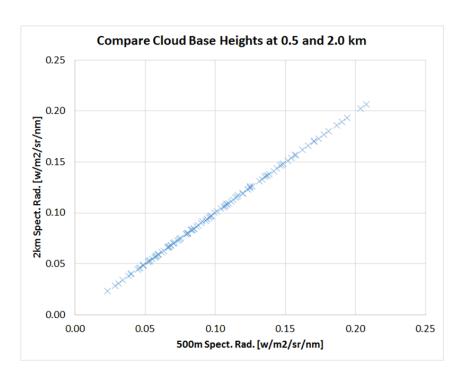


1 2 3 Figure 9 Retrieved COD sensitivity to change in ground albedo from 0.1 (left curves) to 0.2 (right curves) for 3 different SZA curve pairs, computed using asymptotic RT relations. See

⁴ text for details.









3 Figure 10. Scatter plot comparing 440 nm radiance computed for cloud base heights of 500 m

4 and 2 km, each varying over 10 COD and 11 solar zenith angles.

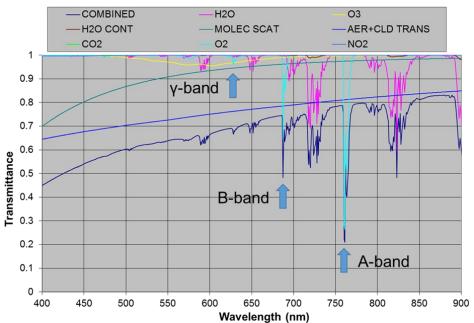
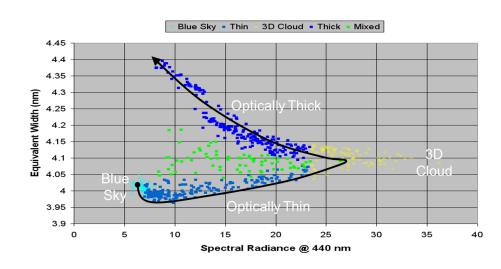


Figure 11 MODTRAN5 calculation of atmospheric transmittance for a ground-based zenith path to space. The Oxygen A-Band is virtually free of absorption by any other species except for aerosol and cloud continua extinction. At 0.1 cm⁻¹ spectral resolution, H2O (water vapor) has a minimum transmittance of 0.9972 across the A-band.





3 Figure 12. The "nose" plot of SR440 vs EQW, indicating a trajectory with increasing COD.

4 The data sample categories (colors) are described in the text.

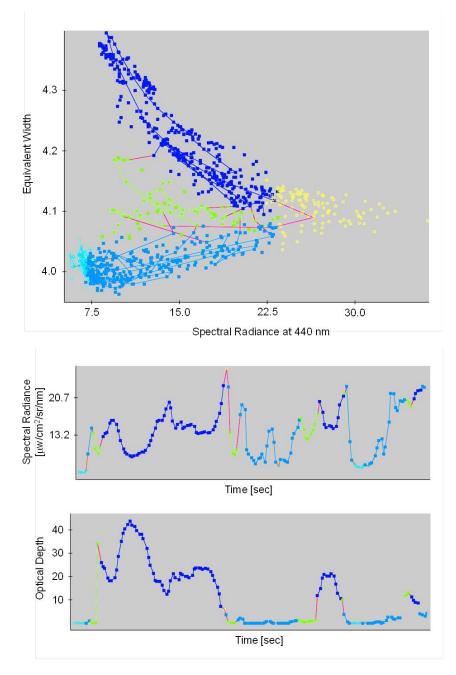
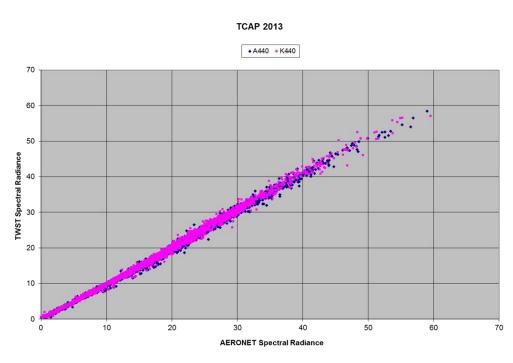


Figure 13. (TOP) Nose plot of data taken over several minutes. The colors are described in
the text. (MIDDLE, BOTTOM) Simultaneous time plots of spectral radiance and retrieved
optical depth corresponding to the TOP nose plot.







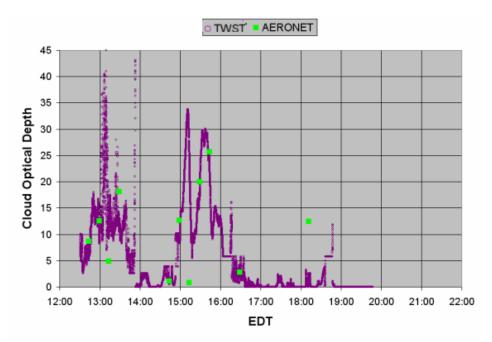
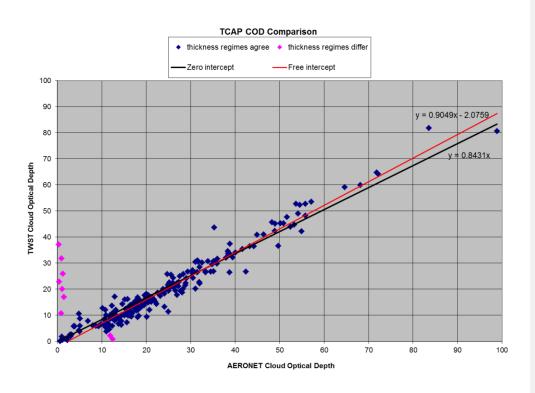


Figure 15 COD time series comparison between AERONET and TWST for 14 June 2013 at the ARM TCAP field campaign.



2 Figure 16 Comparison of TWST and AERONET Cloud-Mode Cloud O