

Retrieving Vertical Ozone Profiles from Measurements of Global Spectral Irradiance

Germar Bernhard¹, Irina Petropavlovskikh^{2,3}, Bernhard Mayer⁴

¹Biospherical Instruments Inc., San Diego, CA 92110, USA

²Cooperative Institute for Research in Environmental Sciences (CIRES), University of Colorado, Boulder, CO 80309, USA

³NOAA Earth System Research Laboratory, Global Monitoring Division, Boulder, CO 80305, USA

⁴Ludwig-Maximilians-Universität, Munich, 80333, Germany

Correspondence to: Germar Bernhard (bernhard@biospherical.com)

Abstract. A new method is presented to determine vertical ozone profiles from measurements of spectral global (direct Sun plus upper hemisphere) irradiance in the [ultraviolet](#). The method is similar to the widely used Umkehr technique, which inverts measurements of zenith sky radiance. The procedure was applied to measurements of a high-resolution spectroradiometer installed near the centre of the Greenland ice sheet. Retrieved profiles were validated with balloon sonde observations and ozone profiles from the space-borne Microwave Limb Sounder (MLS). Depending on altitude, the bias between retrieval results presented in this paper and MLS observations ranges between -5% and $+3\%$. The magnitude of this bias is comparable, if not smaller, to values reported in the literature for the standard Dobson Umkehr method. Total ozone columns (TOCs) calculated from the retrieved profiles agree to within $0.7\pm 2.0\%$ ($\pm 1\sigma$) with TOCs measured by the Ozone Monitoring Instrument onboard the Aura satellite. The new method is called the “Global-Umkehr” method.

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1 Introduction

The “Umkehr” method for determining the vertical distribution of ozone in the atmosphere was first introduced in the 1930s (Götz et al., 1934) and is now routinely applied to measurements of Dobson (e.g., Dütsch, 1959, Mateer and DeLuisi, 1992; Petropavlovskikh et al, 2005) and Brewer (McElroy and Kerr, 1995; Petropavlovskikh et al., 2011) spectrophotometers. The method is typically based on analyzing ratios of zenith-sky radiances at two wavelengths in the ultraviolet (UV), one strongly and one weakly attenuated by ozone, that are measured at solar zenith angles (SZAs) between 60° and 90° . Here we explore a similar optimal statistical approach to obtain vertical ozone information from measurements of spectrally resolved global irradiance, i.e., the irradiance received by a horizontal “cosine” collector from direct Sun and sky (upper hemisphere, from zenith to horizon). Such measurements [were started by several groups in the early 1990s to monitor changes in UV radiation at the Earth’s surface. These activities were motivated by concerns that decreases in atmospheric ozone concentrations, which were caused by ozone depleting substances released by man into the atmosphere, could lead to increases in UV radiation with detrimental effects on human health, and terrestrial and aquatic ecosystems \(e.g., Bais et al.,](#)

1 | [2015](#). [Measurements of global spectral irradiance have been](#) routinely performed by several UV monitoring networks
2 | sponsored by the NSF (<http://uv.biospherical.com/>), NOAA (<http://www.esrl.noaa.gov/gmd/grad/antuv/>), the Network for
3 | the Detection of Atmospheric Composition Change (NDACC; <http://www.ndsc.ncep.noaa.gov/>), Environment Canada
4 | (<http://exp-studies.tor.ec.gc.ca/e/ozone/ozonecanada.htm>), [the European Union](#) (<http://uvdb.fmi.fi/uvdb/>), [and](#) others. The
5 | proposed method has the potential to make these long-term data sets available for assessing vertical ozone information in an
6 | approach similar to [standard zenith-sky](#) Umkehr retrievals. This is particularly interesting for locations where zenith-sky
7 | observations are not available.

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8 | Compared to other methods (e.g., Lidar observations (Megie et al., 1977), balloon sondes, and microwave spectrometers
9 | (Parrish et al., 1992; Waters et al., 2006)), the Umkehr technique provides a relatively inexpensive way of measuring the
10 | vertical distribution of ozone in the atmosphere. The method is most sensitive to the altitude range of 20 to 45 km and has a
11 | resolution of about 10 km within this range. For mid-latitude sites, Brewer Umkehr data have a precision of about 15 % in
12 | the 20- to 40-km region, with larger departures outside this altitude range (McElroy and Kerr, 1995). Umkehr data are
13 | routinely used for monitoring the drift of sensors measuring the vertical distribution from ozone from space (Newchurch et
14 | al. 1987; DeLuisi et al, 1994; Miller et al. 1997; Krzyściński et al, 2009; Petropavlovskikh et al., 2005; Petropavlovskikh et al,
15 | 2011).

16 | The use of measurements of global irradiance instead of zenith-sky radiance for Umkehr retrievals is of no advantage *per se*.
17 | First, global irradiance includes the direct solar beam, which is attenuated according to Beer's law and therefore does not
18 | contain information on the profile. Second, global irradiance includes photons received from directions close to the horizon
19 | and multiple-scattering effects are therefore not negligible. We will show that both challenges can be overcome, resulting in
20 | profiles of similar accuracy than those inverted from zenith-sky observations. The main advantage of the method presented
21 | here is that the vertical distribution of ozone can be derived for locations where no other ground-based data exist from which
22 | profiles could be calculated. The new method is called the "Global-Umkehr" method.

23 | The Global-Umkehr method was tested using data from the NSF UV Monitoring Network (Booth et al., 1994), which has
24 | been measuring UV and visible global spectral irradiance (290 – 600 nm) at six high-latitude sites since 1990. For this study
25 | we used data from Summit, Greenland (72° 35' N, 38° 27' W, 3,202 m a.s.l.) where ozone profiles have been routinely
26 | measured also by balloon sondes. The method can also be applied to measurements at lower latitude sites. [We estimate that
27 | about 25 spectroradiometers that are part of the various UV monitoring networks mentioned earlier provide data of sufficient
28 | quality for the Global-Umkehr method. Some of these instruments were established in the early 1990s at locations around the
29 | globe, including the Arctic, North America, Hawaii, Europe, New Zealand, Australia, and Antarctica.](#)

2 Method

2.1 Retrieval method

The retrieval method is based on the optimal estimation approach (Gauss-Newton method) developed by Rodgers (2000). In brief, the solution (i.e., the ozone concentration as a function of altitude or pressure) is determined iteratively with the matrix equation:

$$\mathbf{x}_{i+1} = \mathbf{x}_i + \mathbf{S}_{i+1} [\mathbf{K}_i^T \mathbf{S}_\varepsilon^{-1} (\mathbf{y} - \mathbf{F}(\mathbf{x}_i)) - \mathbf{S}_a^{-1} (\mathbf{x}_i - \mathbf{x}_a)] \quad (1)$$

where

$$\mathbf{S}_{i+1} = (\mathbf{S}_a^{-1} + \mathbf{K}_i^T \mathbf{S}_\varepsilon^{-1} \mathbf{K}_i)^{-1}. \quad (2)$$

Eqs. (1) and (2) contain the following parameters:

- \mathbf{x}_i is the “state vector” of iteration i . In our implementation, it is defined as the average ozone concentration in eleven layers with a layer thickness of 5 km.
- \mathbf{y} is the “measurement vector,” which is composed of ratios of global spectral irradiance $E(\lambda)$ measured at 310 nm (a wavelength strongly attenuated by ozone) and 337 nm (a wavelength weakly attenuated by ozone) for SZAs ranging between 70° and 90°.
- $\mathbf{F}(\mathbf{x}_i)$ is the solution of the forward model (Sect. 2.3), which simulates the measurements using the state vector as input.
- \mathbf{K}_i is the Jacobian matrix of the partial derivatives of the forward model results and the state vector.
- \mathbf{S}_ε is the covariance matrix quantifying the uncertainty of the measurements.
- \mathbf{x}_a is the *a priori* state vector. The iteration starts by setting $\mathbf{x}_0 = \mathbf{x}_a$.
- \mathbf{S}_a is the covariance matrix pertaining to the *a priori* state vector.
- \mathbf{S}_{i+1} is the solution error covariance matrix at iteration $i+1$, which can be exploited to calculate the uncertainty of the retrieval.

We chose 310 nm as the lower wavelength because measurements at this wavelength are at least a factor of 50 larger than the spectroradiometer’s detection limit of 0.001 mW m⁻² nm⁻¹ for all SZAs and ozone columns of interest. The upper wavelength of 337 nm was chosen because the temperature sensitivity of the ozone absorption cross section has a local minimum at about this wavelength (Bass and Paur, 1985). We also tested other wavelength pairs or combinations of several pairs of wavelengths (e.g., combinations of $E(305)/E(337)$; $E(310)/E(337)$; $E(325)/E(337)$) when developing the method. We found that the use of multiple pairs improved the information content only minimally but increased the computational time considerably.

The SZA range chosen for Umkehr observation is a trade-off between the additional information content resulting from a larger range and the risk that environmental conditions (e.g., clouds, ozone profile) may change substantially over the longer observation time that a larger SZA range requires. During development, we tried several SZA ranges and found that a range

1 of 70° to 90° is a good compromise. This observation is consistent with the conclusion of Petropavlovskikh et al. (2005) that
 2 information in the upper layers is not degraded by changing the SZA range from 60°–90° to 70°–90° in the standard [zenith-](#)
 3 [sky](#) Umkehr method. We also omitted observations with SZAs larger than 90° because of potential systematic errors in the
 4 forward model results (Sect. 2.3) when the Sun is below the horizon. At the latitude of Summit, a SZA range of 70° to 90° is
 5 available in spring between 27 March and 8 May and in fall between 4 August and 15 September.

6 The Jacobian matrix \mathbf{K}_i has the elements $[\mathbf{K}_i]_{mn} = [\partial \mathbf{F}(\mathbf{x}_i)]_m / [\partial \mathbf{x}_i]_n$ and is calculated for every iteration step.

7 The measurement error covariance matrix \mathbf{S}_ε is a diagonal matrix and is constructed by assuming that elements of the
 8 measurement vector have an uncertainty of $\sigma_\varepsilon = 3\%$ and are independent of wavelength and SZA:

$$9 \quad [\mathbf{S}_\varepsilon]_{mn} = \begin{cases} \sigma_\varepsilon^2 [\mathbf{y}]_m [\mathbf{y}]_n & \text{for } m = n \\ 0 & \text{for } m \neq n \end{cases} \quad (3)$$

10 The value of 3% was chosen based on the uncertainty budget of the spectroradiometer installed at Summit (Sect. 2.2). The
 11 choice of 3% was further supported by analyzing the residuals of the retrieval results $(\mathbf{y} - \mathbf{F}(\hat{\mathbf{x}}))$ where $\hat{\mathbf{x}}$ indicates the
 12 solution state vector after the final iteration.

13 *A priori* state vectors \mathbf{x}_a were constructed by combining balloon sonde profiles for altitudes below 10 km and profiles
 14 measured by the Microwave Limb Sounder (MLS) on NASA's Aura satellite for altitudes above 10 km (see Sect. 2.5 for
 15 additional information on these profiles). Separate *a priori* profiles were used for processing data from spring (27 March – 8
 16 May) and fall (4 August – 15 September). Profiles for both seasons were constructed by calculating the median of a large
 17 number of sonde and MLS profiles measured during the two periods using data from the years 2004 to 2014.

18 The covariance matrix pertaining to the *a priori* state vector, \mathbf{S}_a , was constructed as suggested by Bhartia et al. (2013):

$$19 \quad [\mathbf{S}_a]_{mn} = \sigma_a^2 [\mathbf{x}_a]_m [\mathbf{x}_a]_n \exp(-|m - n|/d). \quad (4)$$

20 The parameter σ_a specifies the anticipated variability of the retrieved profiles about the *a priori* profile and can be
 21 interpreted as the relative standard deviation of the profiles' distribution. The correlation length d was set to two, which is
 22 equivalent to 10 km for our definition of the state vector.

23 When σ_a is set to a small value (e.g., 0.1), the solution of the inversion becomes very sensitive to the *a priori* profile. In
 24 contrast, when σ_a is set to a large value, the solution is mostly determined by the measurements. Choosing the optimum
 25 value for σ_a is a trade-off between two competing effects: a large value of σ_a ensures correct inversion result even if the
 26 true profile deviates greatly from the *a priori* profile. On the other hand, a small value of σ_a reduces the risk that the
 27 retrieval result is grossly incorrect if measurements are affected by unanticipated errors.

28 We calculated profiles for $\sigma_a = 0.1$ and 0.4, and compare the results in Sect. 3. The value of $\sigma_a = 0.1$ was chosen by
 29 analyzing the variability of MLS profiles relative to the spring and fall *a priori* profiles introduced above. For Umkehr layers
 30 3 through 7 (the layers for which the Umkehr method is most sensitive) the relative standard deviations calculated from the

1 MLS profiles vary between 0.05 and 0.15; averaged over layers 3 though 7, the relative standard deviation is 0.12 for the
 2 spring and 0.09 for the fall period. The value of $\sigma_a = 0.4$ was chosen as the other extreme. With this value, the *a priori*
 3 profile has little influence on the inversion result and the effect of errors in the measurement vector \mathbf{y} becomes more
 4 prominent. Of note, the retrieval results depend technically on the ratio $\gamma \equiv (\sigma_\varepsilon / \sigma_a)^2$ as opposed to σ_a (Bhartia et al.,
 5 2013). Because the measurement uncertainty σ_ε is well defined, we discuss the results using σ_a instead of γ .

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6 The iteration is repeated until two conditions are met: first, the norms of \mathbf{x}_{i+1} and \mathbf{x}_i must differ by less than 0.5 %, and
 7 second, the values of consecutive results of the cost function $\Psi(\mathbf{x})$ must agree to within 5.0 %, where

$$8 \quad \Psi(\mathbf{x}) = (\mathbf{y} - f(\mathbf{x}))^T \mathbf{S}_\varepsilon^{-1} (\mathbf{y} - f(\mathbf{x})) + (\mathbf{x}_a - \mathbf{x})^T \mathbf{S}_a^{-1} (\mathbf{x}_a - \mathbf{x}). \quad (5)$$

9 These convergence criteria were adopted from Tzortziou et al. (2008). We confirmed that these criteria are also appropriate
 10 for our application by analyzing changes of the two convergence metrics as a function of iteration i . The two criteria are
 11 always met in two to four iterations.

12 The uncertainty e_m of each element of the solution's state vector was calculated according to Goering et al. (2005) from the
 13 diagonal elements of the solution error covariance matrix and the solution state vector:

$$14 \quad e_m = \frac{\sqrt{[\hat{\mathbf{S}}]_{mm}}}{[\hat{\mathbf{x}}]_m}, \quad (6)$$

15 where the caret (^) above the symbols \mathbf{x} and \mathbf{S} indicates the values of \mathbf{x}_i and \mathbf{S}_i of the final iteration.

16 The performance of an inversion based on the optimal estimation approach is often assessed with the averaging kernel matrix
 17 $\mathbf{A} \equiv \hat{\mathbf{S}} \mathbf{K}_i^T \mathbf{S}_\varepsilon^{-1} \mathbf{K}_i$, which quantifies the sensitivity of the retrieved state $\hat{\mathbf{x}}$ to perturbations in the true state \mathbf{x} . For an ideal
 18 observing system, \mathbf{A} is the identity matrix. In reality, the rows of the averaging kernel matrix are peaked with a finite width,
 19 which can be regarded as a measure of the vertical resolution of the retrieved profile. Similarity to the identity matrix
 20 indicates that the retrieval solution has been determined using the observations rather than the *a priori* information, and as
 21 such, the retrieval has provided new information about the actual state.

22 Elements of \mathbf{A} can have large positive and negative values for layers where the ozone concentration is close to zero. To
 23 prevent this predicament, Bhartia et al. (2013) suggested to illustrate the performance of the algorithm with “relative
 24 averaging kernels” (RAK or \mathbf{A}_R), which quantify the relative change of the retrieved state $\hat{\mathbf{x}}$ to the perturbations in the true
 25 state \mathbf{x} . \mathbf{A}_R is defined by

$$26 \quad [\mathbf{A}_R]_{mn} = [\mathbf{A}]_{mn} \frac{[\hat{\mathbf{x}}]_n}{[\hat{\mathbf{x}}]_m}. \quad (7)$$

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27 The optimal estimation technique provides several diagnostics [in addition to the averaging kernels](#) about the quality of the
 28 retrieved profile. The diagnostic used here is d_s , which expresses the “number of degrees of freedom for signal” and

1 indicates the number of useful independent observations in the measurement vector \mathbf{y} . d_s was calculated as suggested by
2 Rodgers (2000) and Goering et al. (2005) from the singular values λ_m of the “error-weighted weighting function matrix”

3 $\tilde{\mathbf{K}} \equiv \mathbf{S}_\varepsilon^{-1/2} \mathbf{K} \mathbf{S}_a^{-1/2}$ via:

$$4 \quad d_s = \sum_m \frac{\lambda_m^2}{m + \lambda_m^2}. \quad (8)$$

5 The diagnostic d_s depends on \mathbf{S}_a and in turn on σ_a . We will show in Sect. 3 that d_s is considerably smaller for profiles
6 calculated with $\sigma_a = 0.1$ than 0.4.

7 2.2 Measurements

8 The method was tested using measurements of global spectral irradiance performed at Summit with a SUV-150B
9 spectroradiometer designed by Biospherical Instruments Inc. The instrument has a spectral resolution of 0.63 nm, is part of
10 the U.S. National Science Foundation’s Arctic Observing Network, and contributes data to NDACC. The expanded
11 uncertainty (coverage factor $k = 2$, equivalent to uncertainties at the 2σ -level or a confidence interval of 95 %) of global
12 spectral irradiance measurements for wavelengths between 310 and 337 nm is between 6.0 and 6.7 %. More information on
13 the instrument is provided by Bernhard et al. (2008) and a detailed uncertainty budget is available at
14 http://uv.biospherical.com/Version2/Uncertainty_SUV150B.pdf.

15 Data used in this paper are from the “Version 2” data edition (Bernhard et al., 2004) and are corrected for the cosine error of
16 the instrument’s entrance optics. The wavelength mapping was determined with a Fraunhofer-line correlation method and
17 the wavelength uncertainty ($k = 2$) of processed data is 0.02 nm. Measured spectra and spectra calculated with the forward
18 model (Sect. 2.3) were convolved with a triangular function of 2 nm bandwidth to further reduce uncertainties resulting from
19 potential wavelength shifts between measured and modelled spectra.

20 The SUV-150B is a scanning instrument, which measures each wavelength at a different time. The time required to scan
21 between 310 and 340 nm is about 140 seconds. Changing cloud condition will therefore affect the ratio of measurements at
22 these wavelengths, and in turn the accuracy of the retrieval result. The effect of clouds on the ratio of $E(310)/E(337)$ can be
23 reduced using measurements of a filtered photodiode, which is illuminated via a beam splitter located between the entrance
24 optics and monochromator of the SUV-150B system. The sensitivity of the diode is centred at 330 nm and measurements are
25 performed continuously during the recording of spectra. Because attenuation of thin clouds is fairly uniform in the 310 to
26 337 nm range (Seckmeyer et al., 1996), measurements of the photodiode can be used to correct for variable cloud
27 attenuation. Specifically, spectral measurements at $\lambda = 310$ nm or $\lambda = 337$ nm are multiplied with a correction factor
28 $C(\lambda, t)$, defined as:

$$29 \quad C(\lambda, t) = \frac{D_C(\theta(t))}{D(t)}, \quad (9)$$

1 where t is the time of the spectral measurement, $\theta(t)$ is the SZA at time t , $D(t)$ is the measurement of the photodiode at
2 time t . $D_C(\theta(t))$ is the hypothetical clear-sky photodiode measurement at time t . The function was parameterized as a
3 function of SZA using measurement of the photodiode obtained during clear skies. Clear-sky periods were determined based
4 on temporal variability using the method described by Bernhard et al. (2008). The correction takes into account that the SZA
5 changes between measurements at 310 and 337 nm. Of note, this technique cannot be applied in the presence of optical thick
6 clouds which enhance ozone absorption of tropospheric ozone due to path length enhancement (Mayer et al., 1998). This
7 restriction does not apply to Summit where clouds are always optically thin (Bernhard et al., 2008). Measurement vectors
8 were inverted both with and without the cloud correction and results are compared in Sect. 3.2.
9 Spectral irradiances at 310 and 337 nm were calculated for all spectra measured during a given period of Umkehr
10 observations and interpolated to a common SZA grid (70, 75, 80, 85, 87, 88, 89, and 90°) using an approximating
11 (smoothing) spline. Compared to an interpolating spline, an approximating spline has the advantage to reduce noise in the
12 measurement vector further. Tests indicated that retrieval results do not change significantly by adding measurements at
13 additional SZAs.
14 The measurement vector is only constructed from spectra measured in the afternoon (between 15:00 and 20:00 UTC)
15 because solar measurements have gaps in the morning when the system performs diagnostics scans with internal lamps
16 (wavelength and irradiance standards).

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 $D_C(\theta(t))$ was parameterized as a function of SZA by filtering measurements for clear skies

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17 2.3 Forward model

18 Forward modelling was performed with Version 1.01 of the pseudospherical discrete ordinate (SDISORT) radiative transfer
19 solver of the UVSPEC/libRadtran model (Mayer and Kylling, 2005). The number of streams was set to 12. The model's
20 results are spectra of global irradiance. Model input parameters include the extraterrestrial spectrum as defined by Bernhard
21 et al. (2004) and available at http://uv.biospherical.com/Version2/Paper/2004JD004937-ETS_GUEYMARD.txt; surface
22 albedo; atmospheric pressure; and the ozone absorption cross section (Bass and Paur, 1985). While more accurate ozone
23 absorption cross sections are now available (Gorshchev et al., 2014; Orphal et al., 2016), we used Bass and Paur (1985) data
24 to facilitate validation with OMI total ozone column measurements, which are also based on Bass and Paur (1985). The
25 surface albedo at Summit was set to 0.97 in good agreement with recent measurements (Carmagnola et al., 2013). Aerosol
26 optical depth was set to stratospheric background conditions. Atmospheric pressure and profiles of gases other than ozone
27 (O_2 , H_2O , CO_2 , and NO_2) were taken from the [Air Force Geophysics Laboratory \(AFGL\)](#) atmospheric constituent profile for
28 subarctic summer (Anderson et al., 1986), which defines the atmosphere at 51 levels. The vertical distribution of ozone in
29 this standard profile was replaced with the profile defined by the state vector \mathbf{x}_i and updated in every iteration.
30 The SDISORT solver has been successfully validated using data of the NSF UV Monitoring Network (e.g., Bernhard et al.,
31 2004, 2008) and for a large range of conditions at other sites (e.g., Mayer and Kylling, 2005, and references therein).
32 However to the best of our knowledge, a rigorous validation for the large SZAs required for Umkehr retrievals has not been

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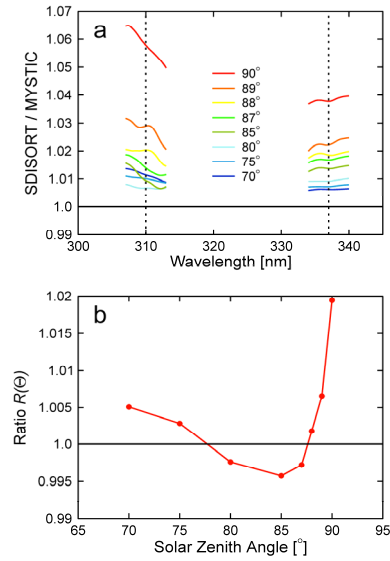
1 conducted. The pseudospherical approximation used by SDISORT correctly describes the attenuation of the direct beam in
2 spherical geometry but the diffuse radiance is computed in plane-parallel geometry (Mayer et al., 2015). This approximation
3 can lead to significant errors at large SZAs (Petropavlovskikh et al., 2000; Emde and Mayer, 2007). To quantify these errors
4 for our application, we have compared spectra of global irradiance calculated with SDISORT with the spherical solver of the
5 MYSTIC (Monte Carlo code for the phYSically correct Tracing of photons In Cloudy atmospheres) model, which fully
6 solves the spherical geometry without any approximations (Mayer, 2009). Both models were run with the same set of input
7 parameters (AFGL subarctic summer with *a priori* ozone profile for spring at Summit) for wavelengths between 307 and
8 313 nm and between 334 and 340 nm in 0.5 nm steps. The MYSTIC model was run with 84 million photons per wavelength
9 and per SZA, resulting in photon noise of less than 0.5 % at SZA=90° (worst case). Resulting spectra of both models were
10 convolved with a triangular function of 2 nm bandwidth to further reduce noise and to be consistent with the method used in
11 the Umkehr code.

12 Fig. 1a shows the ratio of SDISORT and MYSTIC spectra calculated for the eight SZAs used in our [implementation of the](#)
13 [Global-Umkehr method](#). SDISORT overestimates spectral irradiances relative to MYSTIC at all wavelengths and SZAs. For
14 SZA $\leq 88^\circ$, the bias is less than 2 % but increases to up to 6.5 % for SZA = 90°. For the Umkehr retrieval, only the ratio

15 $q(\theta) \equiv \frac{E(310, \theta)}{E(337, \theta)}$ is important where θ indicates again the SZA. The ratio $R(\theta) \equiv \frac{q_{SDISORT}(\theta)}{q_{MYSTIC}(\theta)}$ resulting from calculating

16 $q(\theta)$ with SDISORT and MYSTIC is shown in Fig. 1b. $R(\theta)$ ranges between 0.998 at 80° to 1.019 at 90°. Calculations
17 with the MYSTIC model can be considered [as](#) the most accurate results attainable because the Monte Carlo code does not
18 use approximations. The model has been validated by comparison to other spherical radiative transfer models and by
19 simulating the radiance distribution of the sky during a total solar eclipse. For such calculations, a spherical solver without
20 approximations is required because light entering the atmosphere more than 1000 km away may impact the radiance in the
21 centre of the umbral shadow (Emde and Mayer, 2007).

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2 Fig. 1. Comparison of results calculated with the SDISORT and MYSTIC
 3 spectra calculated for eight SZAs (see legend). (b) Ratio $R(\theta)$. See text for definition.

4

5 Relative to MYSTIC, SDISORT overestimates $q(\theta)$ for SZA larger than 88°. In our Umkehr code, we scale the results of
 6 the forward model with $1/R(\theta)$ to account for the bias of the SDISORT model. Note that the MYSTIC model is too slow to
 7 be used for Umkehr retrievals: the calculation of the eight spectra used for Fig. 1a required a run time of over three days.

8 The forward model requires that the vertical structure of the atmosphere is defined as a function of altitude. The association
 9 between altitude and pressure is defined in the AFGL profile and this relationship may differ from the actual pressure profile
 10 at the time of Umkehr observation. Because our measurements do not allow to reconstruct the pressure profile, we report all
 11 ozone profiles as a function of pressure, and compare the retrieved profile with sonde and MLS profiles, which are also
 12 provided as a function of pressure. The standard [zenith-sky](#) Umkehr technique (Petropavlovskikh et al., 2005) uses a similar
 13 approach. Table 1 provides altitude and pressure ranges for each Umkehr Layer. Note that Layer 0 starts at the elevation of
 14 Summit (3,202 m).

1 Table 1. Assignment of Umkehr Layers.

Umkehr Layer	Altitude range forward model [km]	Pressure range [hPa]
10	50.0 – 55.0	0.987 – 0.537
9	45.0 – 50.0	1.82 – 0.987
8	40.0 – 45.0	3.40 – 1.82
7	35.0 – 40.0	6.61 – 3.40
6	30.0 – 35.0	13.4 – 6.61
5	25.0 – 30.0	27.8 – 13.4
4	20.0 – 25.0	59.0 – 27.8
3	15.0 – 20.0	126.0 – 59.0
2	10.0 – 15.0	267.7 – 126.0
1	5.0 – 10.0	541.0 – 267.7
0	3.202 – 5.0	664 – 541

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3 **2.4 Validation method**

4 The retrieved Umkehr profiles were validated using ozone profiles measured at Summit with balloon sondes by
 5 NOAA/GMD (Oltmans et al., 2010) and ozone profiles provided by MLS on Aura. Sondes are typically launched between
 6 12:00 and 20:00 UTC. MLS measures thermal emissions from rotational lines of ozone through the limb of the atmosphere.
 7 Ozone measurements have a vertical range of 12–73 km with a vertical resolution of 2–3 km below 65 km. The horizontal
 8 resolution is about 200 km and the accuracy is about 5–10 % between 16 and 60 km (Froidevaux et al., 2008). [The average](#)
 9 [horizontal distance between the locations of Summit and MLS data is 160 km.](#) Sonde and MLS profiles were downloaded
 10 from <ftp://ftp.cmdl.noaa.gov/ozwv/Ozonesonde/Summit,%20Greenland/> and
 11 http://avdc.gsfc.nasa.gov/pub/data/satellite/Aura/MLS/V04/L2GPOVP_Prof/O3/Summit/, respectively. Sonde profiles are
 12 only available for 2 to 4 days per month whereas MLS profiles are available on a daily basis. [MLS measurements at Summit](#)
 13 [take place either between 5:28 and 6:26 UTC or between 14:11 and 15:10 UTC. There is only one data file per day in the](#)
 14 [NASA archive.](#)

15 The total ozone column (TOC) was calculated from the retrieved Umkehr profiles and compared with measurements of the
 16 Ozone Monitoring Instrument (OMI) on board NASA’s Aura spacecraft. OMI overpass data were downloaded from
 17 <http://avdc.gsfc.nasa.gov/index.php?site=1593048672&id=28>. OMI data use the Bass and Paur (1985) ozone absorption
 18 cross section (pers. comm., David Haffner, NASA), like the forward model.

19 Good validation results can only be expected if the actual ozone profile does not change over the period of Umkehr
 20 observations. We therefore only considered periods where the TOC measured by OMI did not change by more than 20 DU

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1 | between 15:00 UTC on the day of the comparison and the first observation on the following day. [This criterion ensures that](#)
2 | [changes in the ozone profile remain below about 4% for all Umkehr layers.](#) Retrieved Umkehr profiles were compared with
3 | the sonde profile measured on the same day (if available) and with the MLS profiles measured on this day (labelled “MLS 1”
4 | in the following) as well as the next day (labelled “MLS 2”).

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5 | **3 Results**

6 | We first show retrieval results for three sample days with greatly different conditions and compare these results with profiles
7 | measured by balloon sondes and MLS (Sect. 3.1). We then discuss in Sect. 3.2 statistics for all profiles that were retrieved
8 | under sufficiently stable conditions (variation in total ozone of less than ± 20 DU).

9 | **3.1 Comparison with balloon sonde and MLS profiles – sample profiles**

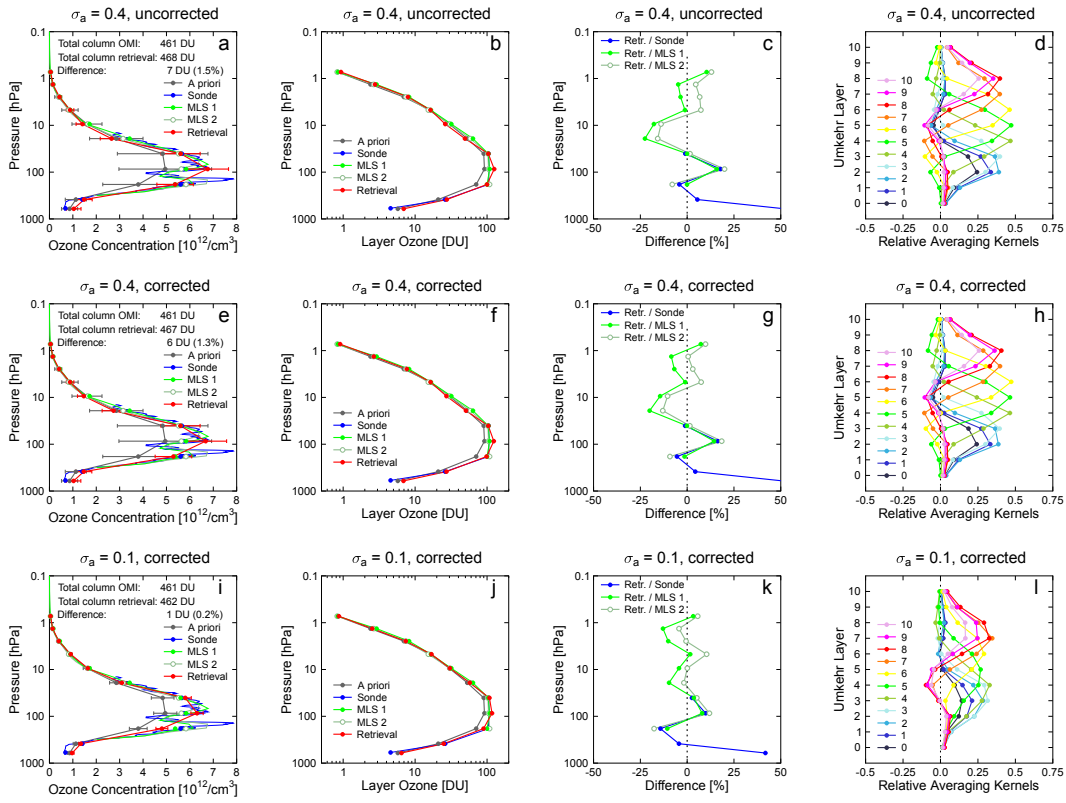
10 | **3.1.1 Validation for 19 April 2014**

11 | Fig. 2 compares the retrieved ozone profile for 19 April 2014 with the *a priori*, balloon sonde and MLS profiles. OMI
12 | measured a TOC of 461 DU on this day, which was the third highest TOC of the dataset and the highest TOC of days when
13 | balloon sonde data were available. [Therefore, the profile represents one of the highest departures from the spring *a priori*](#)
14 | profile.

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3 Fig. 2. Validation of ozone profile retrieved for 19 April 2014. Top row: results for $\sigma_a = 0.4$ and uncorrected forward
 4 model. Centre row: results of $\sigma_a = 0.4$ and corrected forward model. Bottom row: $\sigma_a = 0.1$, and corrected forward model.
 5 1st column: ozone concentration as a function of pressure for *a priori* profile (grey), balloon sonde profile (blue), MLS
 6 profile for day of retrieval (MLS 1, dark green), MLS profile of the following day (MLS 2, light green), and retrieved profile
 7 (red). Solid or open circles indicate for each dataset ozone concentrations averaged over each of the eleven Umkehr layers
 8 defined in Table 1. Grey error bars indicate the diagonal elements of S_a . Red error bars indicate the uncertainty of the
 9 retrieval e_m . TOCs measured by OMI and calculated from the retrieved profile are indicated in the legend. 2nd column: layer
 10 ozone as a function of pressure for *a priori* profile, balloon sonde profile, MLS profile for day of retrieval, MLS profile of
 11 the following day, and retrieved profile. 3rd column: difference between the retrieval and sonde, MLS 1 and MLS 2
 12 datasets averaged over each Umkehr layer. 4th column: relative averaging kernels.

1 Results are shown for three sets of retrieval parameters: (1) $\sigma_a = 0.4$, forward model not corrected (top row of Fig. 2); (2)
2 $\sigma_a = 0.4$, forward model corrected by scaling with $1/R(\theta)$ (centre row of Fig. 2); and (3) $\sigma_a = 0.1$, forward model
3 corrected (bottom row of Fig. 2). For each set of parameters, we show profiles of ozone concentrations (1st column of Fig.
4 2), layer ozone (2nd column of Fig. 2), the difference between the retrieved profile and the profiles measured by sondes and
5 MLS (3rd column of Fig. 2), and the relative averaging kernels (RAKs) of the retrieval (4th column of Fig. 2).

6 Layer ozone (Fig. 2b, f, and j) was calculated by integrating average ozone concentrations of each Umkehr layers over
7 height. Note that ozone concentrations (Fig. 2a, e, and i) are plotted on a linear scale to highlight differences in the
8 troposphere and lower stratosphere, while layer ozone (Fig. 2b, f, and j) is plotted on a logarithmic scale to better distinguish
9 differences in the upper stratosphere.

10 Fig. 2c, g, and k show differences of the average ozone concentrations for the 11 Umkehr layers. Two MLS datasets are
11 considered. The dataset labelled “MLS 1” is from the same day as the retrieval while the dataset labelled “MLS 2” is from
12 the following day.

13 When plotting ozone concentrations on a linear scale (Fig. 2a, e, i), results for the three sets of parameters look similar. As
14 expected, the resolution of the retrieval is not sufficient to capture the large fluctuation in the ozone concentrations between
15 about 100 and 300 hPa indicated by sonde and MLS measurements. Furthermore, the retrieved profiles overestimate the
16 ozone concentration at the peak of the profile at about 100 hPa and underestimates the profile in the 7 to 28 hPa range
17 (Layers 5 and 6). The difference of -22.5% between the retrieval and MLS 1 seen in Fig. 2c for Layer 5 is one of the largest
18 negative biases of all profiles processed. This large bias may partially be caused by errors in the measurement vector due to
19 clouds (The photodiode used for cloud correction was not available on this day). The large deviation for Layer 0 of 52%
20 is not surprising considering that this layer is only 1.8 km thick and the sensitivity of the Umkehr method to ozone
21 concentrations close to the surface is poor.

22 The bias of the retrieval becomes smaller when the forward model is corrected for the systematic error resulting from the
23 pseudospherical approximation (compare Fig. 2c and Fig. 2g), indicating that the correction is appropriate.

24 The smallest difference between the retrieval on one hand and sonde and MLS measurements on the other is observed for
25 $\sigma_a = 0.1$ (Fig. 2k). This suggests that a relatively small value for σ_a is advantageous even though the sample profile
26 deviates considerably from the *a priori* profile. For Layers 5 to 9, the magnitude of the bias is comparable in magnitude to
27 the difference between the two MLS profiles, suggesting that a portion of the bias could be due to changes in the ozone
28 profile occurring during the period of Umkehr observations.

29 When σ_a is set to 0.4, the RAKs of Layers 3 to 7 peak at the correct layer and drop to zero within two layers, suggesting
30 that ozone concentrations in this altitude range can be well resolved (Fig. 2d, h). In contrast, RAKs for layers 0, 1, and 2 are
31 similar and peak at about the same altitude. Hence, ozone concentrations in these layers cannot be separated well. The
32 altitude resolution of the standard [zenith-sky](#) Umkehr method is also poor in these layers, and results for layers 0 and 1 are
33 typically combined when reporting data. RAKs for layers 8 – 10 peak at the same altitude, indicating that ozone

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1 concentrations above the 3 hPa level (about 45 km) cannot be resolved and the retrieval is predominantly driven by the *a*
2 *priori* profile. This is not surprising considering the small ozone concentrations in these layers. Also the traditional [zenith-](#)
3 [sky](#) Umkehr method has little sensitivity at these altitudes.
4 When σ_a is set to 0.1, the RAKs become rather broad (Fig. 2l). The solution is therefore more determined by the *a priori*
5 profile than the observations. The reduced importance of the measurements is also reflected in the value of d_s : d_s is 3.02
6 for $\sigma_a = 0.4$ and 2.15 for $\sigma_a = 0.1$.
7 TOCs calculated from the retrieved profiles agree well with the OMI measurements and depend only little on the choice of
8 retrieval parameters: absolute and relative biases are 7 DU (1.5 %) for parameter set (1), 6 DU (1.3 %) for set (2), and 1 DU
9 (0.2 %) for set (3).

10 3.1.2 Comparison for 11 April 2007

11 Fig. 3 shows results for 11 April 2007. On this day, ozone concentrations measured by sonde and MLS were consistently
12 below the *a priori* profile between 5 and 100 hPa, but between 100 and 300 hPa, the actual profile exceeded the *a priori*. Fig.
13 3a – d show results calculated with $\sigma_a = 0.4$ while calculations used for Fig. 3e – h used $\sigma_a = 0.1$. The forward model was
14 corrected by scaling with $1/R(\theta)$ in both cases. Note that the MLS 1 and MLS 2 datasets are almost identical, indicating that
15 the actual ozone profile was constant over the observation period. RAKs are very similar to those for 19 April 2014
16 ([Compare Fig. 2h with Fig. 3d and Fig. 2l with Fig. 3h](#)).
17 For both settings of σ_a , the retrieved profile is narrower than the *a priori* profile and matches the MLS profile almost
18 ideally for Layers 3 – 9. This is an example that the retrieval result is not simply the *a priori* profile scaled with a constant
19 factor. Instead, the information contained in the measurement vector is sufficient to modify the shape of the profile to match
20 the actual, narrower shape. However, like in the case of the first example, the resolution of the Umkehr method is not
21 sufficient to reproduce the fluctuation of the actual ozone profile between 70 and 300 hPa. The most obvious difference
22 between the results calculated with $\sigma_a = 0.4$ and 0.1 is the difference at 183 hPa (Layer 2). Because the Umkehr method has
23 little sensitivity at this pressure level, the retrieved ozone concentration is mostly determined by the *a priori* profile for $\sigma_a =$
24 0.1 (Fig. 3g). In contrast, when setting $\sigma_a = 0.4$, measurements “pull” the retrieval to the higher concentrations of the actual
25 profile, resulting in a smaller bias relative to sonde and MLS data (Fig. 3c). The TOCs of both retrievals agree to within 7
26 DU (or 2.1 %) with OMI.

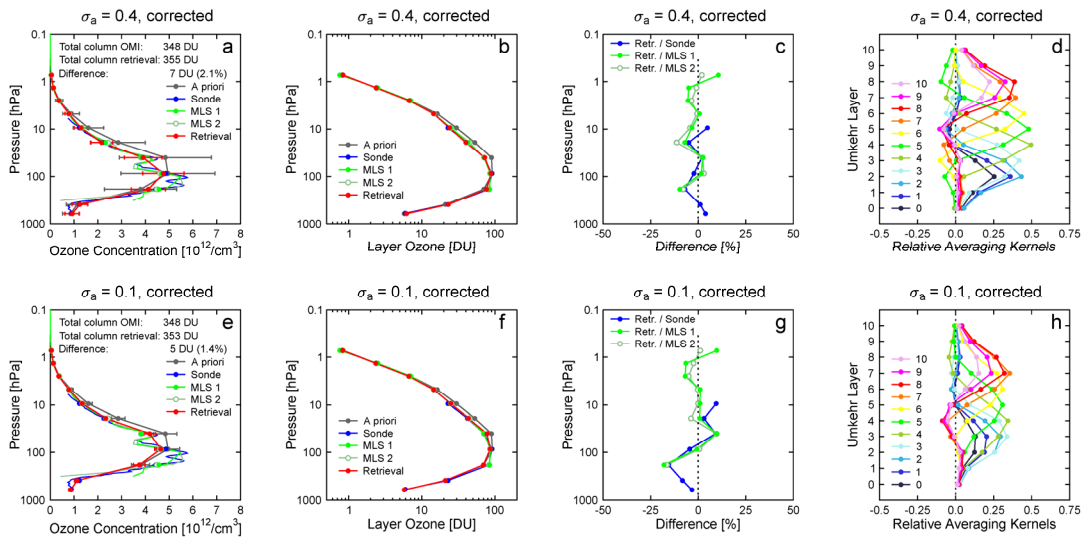
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Fig. 3. Validation of ozone profile retrieved for 11 April 2007. Top row: results of $\sigma_a = 0.4$ and corrected forward model. Bottom row: $\sigma_a = 0.1$, and corrected forward model. 1st column: ozone concentration as a function of pressure. 2nd column: layer ozone as a function of pressure. 3rd column: difference between the retrieval and sonde, MLS 1 and MLS 2 datasets averaged over each Umkehr layer. 4th column: [relative averaging kernels](#). Labelling of the different datasets is identical to that of Fig. 2.

8 3.1.3 Comparison for 14 August 2009

9 The third example (Fig. 4) shows results from 14 August 2009 when the ozone profile was almost identical with the fall *a priori* profile. Note that this *a priori* profile is considerably below that for spring (e.g., Fig. 3d). Calculations were performed
10 with $\sigma_a = 0.1$ and the corrected forward model. Results agree with sonde and MLS profiles to within $\pm 13\%$ for Layers 1 –
11 10 and the TOC of the retrieval is [virtually](#) identical to the OMI measurement. [The effect of changing \$\sigma_a\$ from 0.4 to 0.1 are similar for spring and fall profiles and results for \$\sigma_a = 0.4\$ were therefore omitted in Fig. 4.](#)

12 In summary, Umkehr profiles replicate the general pattern in the sonde and MLS data but cannot resolve the fine structure in
13 the ozone distribution, in particular below 100 hPa. The relatively poor resolution in the troposphere and lower stratosphere
14 is similar for the standard [zenith-sky](#) Umkehr method.

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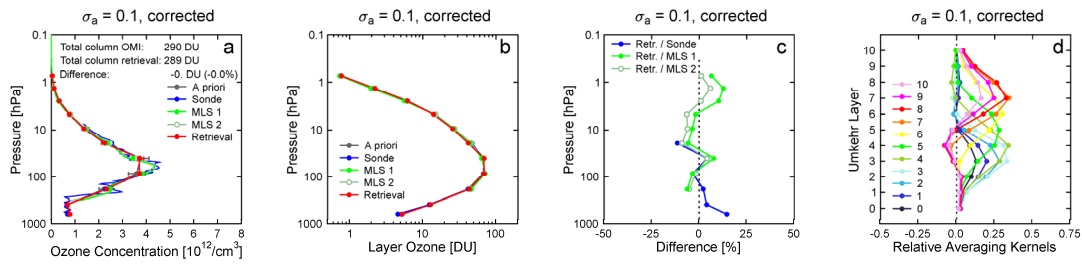


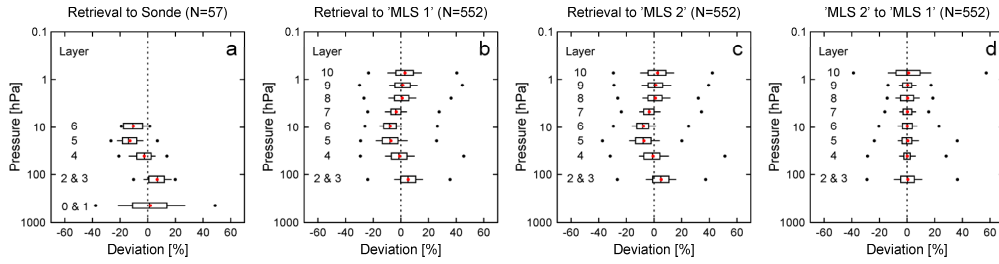
Fig. 4. Validation of ozone profile retrieved for 14 August 2009. The retrieved profile was calculated with $\sigma_a = 0.1$ using the corrected forward model. (a) ozone concentration. (b) layer ozone. (c) difference between the retrieval and sonde, MLS 1 and MLS 2 datasets averaged over each Umkehr layer. (d) [relative averaging kernels](#). Labelling of the different datasets is identical to that of Fig. 2.

3.2 Comparison with balloon sonde and MLS profiles – statistics

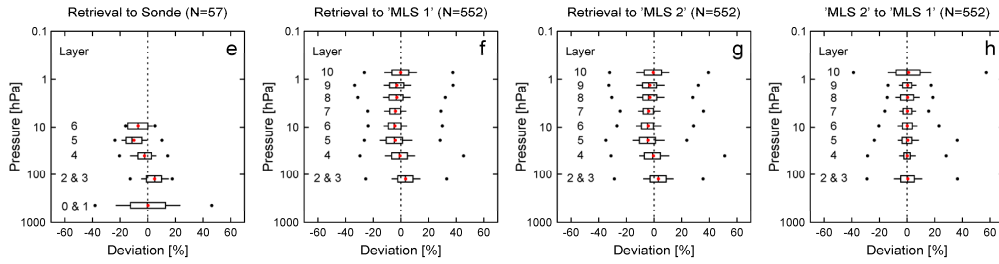
While the results for the three profiles discussed above are promising, they do not allow [a comprehensive assessment of the Global-Umkehr technique](#). To fully validate the method, we compared a large number of sonde and MLS profiles with our retrievals using measurements from the years 2004 to 2014, and calculated statistics. We only considered periods when the TOC was constant to within ± 20 DU as indicated by OMI. This criterion restricted the number of comparisons with sonde profiles to 57 and with MLS profiles to 552. Data were processed with and without the model correction discussed in Sect. 2.3 and with and without the cloud correction discussed in Sect. 2.2. The latter correction requires measurements of the photodiode internal to the SUV-150B instrument. Unfortunately, these measurements were not available during all days, reducing the number of retrieval/sonde and retrieval/MLS comparisons to 38 and 396, respectively. Results from Layers 0 and 1 and Layers 2 and 3 were combined because of the poor vertical resolution of the Umkehr methods in the troposphere and lower stratosphere discussed earlier. Differences between retrieval and sonde, MLS 1, and MLS 2 data are illustrated with box-whisker plots (Fig. 5), which show the minimum and maximum difference (black dots), median (black line), average (red dot), interquartile (i.e., 25th – 75th percentile) range (box), and the 10th – 90th percentile range (whiskers) for each layer or combination of layers. We also plotted statistics for the difference of the MLS 1 and MLS 2 datasets to indicate the variability of the actual ozone profile over the course of one day. Fig. 5 includes results from spring and fall combined. Table 2 provides statistics calculated separately for spring and fall.

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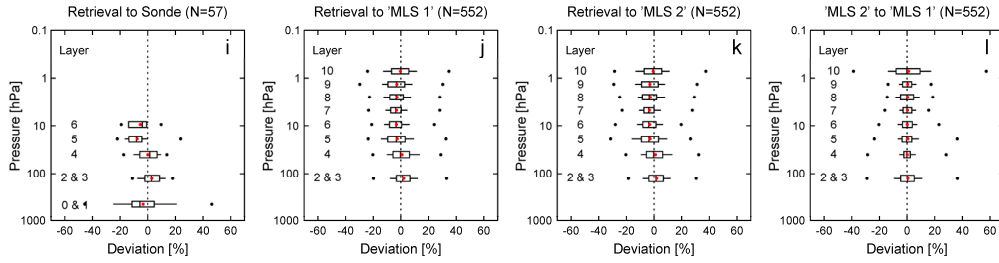
No model correction, no cloud correction, $\sigma_a = 0.4$



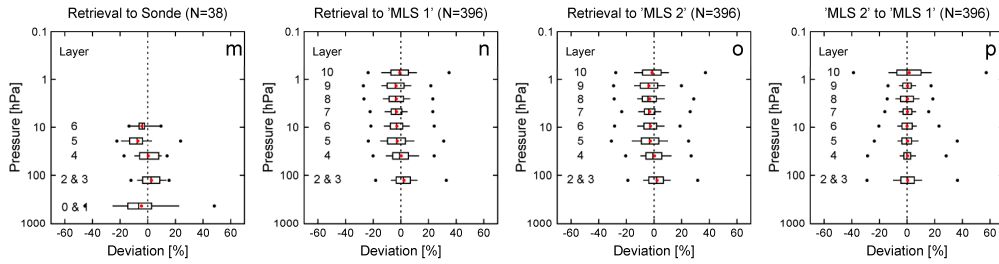
Model correction, no cloud correction, $\sigma_a = 0.4$



Model correction, no cloud correction, $\sigma_a = 0.1$



Model correction, cloud correction, $\sigma_a = 0.1$



1 Fig. 5. Box-whisker plots showing the difference between Umkehr retrieval results and sonde measurements (1st column),
2 MLS observations for day of retrieval ('MLS 1' dataset, 2nd column), and MLS observations for the following day ('MLS 2'
3 dataset, 3rd column). The 4th column illustrates the difference of the 'MLS 2' and 'MLS 1' datasets. Each plot shows the
4 minimum and maximum difference (black dots), median (black line), average (red dot) interquartile range (box) and the 10th
5 – 90th percentile range (whiskers) for each layer. Results for Layers 0 and 1 and Layers 2 and 3 were combined. The N-
6 number in the headers of each plot indicates the number of profiles used for computing the statistics. Results in each row
7 were calculated with a different set of parameters: 1st row (panels a–d): forward model not corrected, no cloud correction,
8 $\sigma_a = 0.4$. 2nd row (panels e–h): forward model corrected by scaling with $1/R(\theta)$, no cloud correction, $\sigma_a = 0.4$. 3rd row
9 (panels i–l): forward model corrected, no cloud correction, $\sigma_a = 0.1$. 4th row (panels m–p): forward model corrected, cloud
10 correction using data of photodiode, $\sigma_a = 0.1$.

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13 Table 2. Bias and interquartile range (in parenthesis) of retrieval/MLS 1 comparison; average and standard deviation of the
14 difference between total ozone calculated from retrieved profiles and measured by OMI (TOC); and average number of
15 degrees of freedom for signal ($\langle d_s \rangle$) for spring and fall periods. The second column provides the number of profiles (N)
16 contributing to the statistics.

Season	N	Bias and interquartile range of retrieval/MLS 1 comparison for Layer								TOC	$\langle d_s \rangle$
		2 & 3	4	5	6	7	8	9	10		
<i>No model correction, no cloud correction, $\sigma_a = 0.4$</i>											
Spring	197	4% (14%)	-1% (10%)	-8% (9%)	-10% (8%)	-4% (9%)	-1% (11%)	0% (10%)	4% (15%)	0.2% (1.9%)	3.1
Fall	355	6% (10%)	-1% (12%)	-8% (12%)	-7% (9%)	-3% (7%)	1% (9%)	1% (11%)	3% (11%)	0.7% (1.8%)	3.0
<i>Model correction, no cloud correction, $\sigma_a = 0.4$</i>											
Spring	197	2% (13%)	-1% (10%)	-6% (10%)	-6% (8%)	-4% (10%)	-5% (11%)	-5% (10%)	0% (14%)	0.0% (1.9%)	3.1
Fall	355	4% (10%)	-1% (12%)	-4% (13%)	-4% (9%)	-4% (7%)	-3% (9%)	-3% (11%)	0% (11%)	0.5% (1.8%)	3.0
<i>Model correction, cloud correction, $\sigma_a = 0.4$</i>											
Spring	142	3% (13%)	-2% (12%)	-6% (10%)	-6% (8%)	-5% (10%)	-6% (12%)	-6% (10%)	-1% (15%)	0.0% (2.0%)	3.1
Fall	254	3% (10%)	-1% (11%)	-4% (12%)	-3% (9%)	-4% (8%)	-3% (10%)	-3% (10%)	-1% (11%)	0.5% (1.9%)	3.0
<i>Model correction, no cloud correction, $\sigma_a = 0.1$</i>											
Spring	197	1% (12%)	-1% (12%)	-4% (12%)	-5% (10%)	-4% (10%)	-4% (12%)	-4% (13%)	1% (14%)	-0.2% (1.8%)	2.2
Fall	355	2% (9%)	0% (12%)	-3% (14%)	-3% (8%)	-4% (7%)	-3% (9%)	-3% (12%)	0% (12%)	0.3% (1.7%)	2.1
<i>Model correction, cloud correction, $\sigma_a = 0.1$</i>											
Spring	142	1% (12%)	-1% (13%)	-4% (12%)	-5% (10%)	-4% (11%)	-4% (12%)	-5% (13%)	-1% (16%)	-0.2% (1.9%)	2.2
Fall	254	2% (8%)	-1% (11%)	-2% (14%)	-2% (8%)	-3% (7%)	-4% (9%)	-4% (12%)	-1% (11%)	0.3% (1.7%)	2.1

5 The 1st row (panels a–d) of Fig. 5 shows results calculated without the model and cloud corrections; σ_a was set to 0.4. The average and median biases between retrieval and MLS data vary between –8 % and +5 % (Fig. 5b, c). The largest negative bias is observed for Layers 5 and 6 while the largest positive bias of 5 % is observed closest to the surface (Layer 2&3). Biases relative to the sonde measurements (Fig. 5a) are by and large consistent with biases relative to MLS data, although the comparatively small number of sonde observations makes statistics less robust. Fig. 5d confirms that there is no systematic difference between the MLS measurements on the day of Umkehr observations (MLS 1) and the following day (MLS 2).

10 For the retrieval/MLS comparisons, the interquartile ranges vary between 7 % and 12 % and depends only modestly on the layer. With the exception of the results for the highest layer, the interquartile ranges for the MLS 2 to MLS 1 comparison vary between 5 % and 10 %. Differences between the 10th and 90th percentiles vary between 14 % and 24 % for the retrieval/MLS comparisons (whiskers in Fig. 5b, c) and between 12 % and 17 % for the MLS 2 / MLS 1 comparison, excluding the highest layer (Fig. 5d). The similarity of the ranges for the retrieval/MLS and MLS 2 / MLS 1 comparisons suggests that a large portion of the observed retrieval/MLS differences can be attributed to changes in the actual ozone profile over the time periods relevant for these comparisons. Lastly, the large interquartile range for the retrieval/sonde comparison observed in Layer 0&1 (Fig. 5a) is again a manifestation of the fact that the Umkehr method has little sensitivity for the layers closest to the surface.

20 To assess the effect of the forward model correction on our Umkehr retrievals, we repeated the calculations with this correction. Results are presented in the 2nd row (panels e–h) of Fig. 5. As before, no cloud correction was applied and σ_a was set to 0.4. By comparing the original results (Fig. 5b, c) with the corrected results (Fig. 5f, g) it can be observed that the bias between retrieval and MLS data has diminished and now varies between –5 % (Layers 5 and 6) and +3 % (Layer 2&3), suggesting that the model correction is justified. The interquartile ranges with and without the correction are virtually indistinguishable. Note that the correction has no effect on the MLS 2 / MLS 1 comparison and Fig. 5d and h are therefore identical.

25 To explore the effect of σ_a on the results, we repeated the calculations using $\sigma_a = 0.1$ instead of $\sigma_a = 0.4$. Results are shown in the 3rd row (panels i–l) of Fig. 5. For $\sigma_a = 0.1$, the bias between retrieval and MLS data has decreased further and now varies between –4 % and +1 % (Fig. 5j, k). Differences between retrieval and sondes (Fig. 5i) have also decreased compared to calculations with $\sigma_a = 0.4$, except for Layer 0&1. The observation that biases are larger for a larger value of σ_a could be caused by systematic errors in the measurement vector or an incomplete correction of the forward model

results. Changing σ_a from 0.4 to 0.1 had almost no effect on the interquartile range, however, minimum and maximum differences (black dots) contracted somewhat.

Finally, the calculations were repeated with the cloud correction turned on (4th row of Fig. 5, panels m–p). For the retrieval/MLS comparison, biases and interquartile ranges with and without the cloud correction agree to within 1 %. Results for the retrieval to sonde comparison (Fig. 5m) are affected by the small sample size of N=38. (Note that results shown for Layer 6 are only based on eight samples because most balloons burst before they reach this layer).

The difference between uncorrected and cloud-corrected statistics is very small, suggesting that clouds affect the accuracy of the retrievals only marginally. However, this conclusion may not apply to locations with thicker clouds and should be tested if the method is used at other sites.

Table 2 allows the assessment of retrievals for spring and fall periods separately. Because statistics are more robust for the retrieval/MLS than retrieval/sonde comparisons, Table 2 only presents results for the former. Biases and interquartile ranges are provided with and without the model and cloud corrections, and with σ_a set to either 0.4 or 0.1. Biases for spring and fall agree to within 3 % for all layers. When no corrections are applied and $\sigma_a = 0.4$, biases range between –10 % (Layer 6 for spring) to +6 % (Layer 2&3 for fall). The model correction decreases this range to between –6 % to 4 %. By reducing σ_a from 0.4 to 0.1, the range decreases further to between –5 % to 2 %. The cloud correction has a negligible (≤ 1 %) effect on the biases. Interquartile ranges vary from 7 % to 16 % and depend only little (≤ 3 %) on σ_a and on whether or not corrections are applied.

Table 2 also includes a column comparing TOCs derived from the retrieved profiles with measurements by OMI. Depending on σ_a and the correction method, the average difference between the retrieved and OMI TOCs varies between –0.2 % and 0.7 %, and the standard deviation varies between 1.7 % and 2.0 %.

Lastly, the average value of d_s is about 3.0 for $\sigma_a = 0.4$ and 2.1 for $\sigma_a = 0.1$. A value of $d_s = 3.0$ may seem low, but it is consistent with values of d_s resulting from the standard zenith-sky Umkehr technique. For example, Stone et al. (2015) reported a value of $d_s = 3.1$ for Dobson zenith-sky Umkehr retrievals using the Dobson C wavelength pair (311.4 and 332.4 nm) and the standard Dobson SZAs ranging from 60° to 90°.

4 Discussion

When the forward model is corrected, the bias of our retrievals relative to MLS data is smaller than ± 6 % for all layers. This level of agreement compares favourably with published results of the standard zenith-sky Umkehr method. For example, McElroy and Kerr (1995) compared Umkehr profiles derived from a Brewer spectrophotometer with concurrent measurements of a lidar, a microwave radiometer and ozone sondes, which were performed during a one-month campaign at the Table Mountain Observatory in California. The mean bias between the Brewer zenith-sky Umkehr results and the mean

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of the other instruments varied to within $\pm 10\%$ for altitudes between 20 and 35 km. Between 37 and 47 km, the Brewer data were low by 15 to 20 % (Fig. 9 of McElroy and Kerr, 1995).

Nair et al. (2011) compared stratospheric ozone vertical distribution measured by a large number of ground-based and satellite sensors at the Haute-Provence Observatory, France. They found that zenith-sky Umkehr data from an automated Dobson spectrophotometer systematically underestimate the stratospheric ozone concentration with a near-zero bias at about 30 km, but increasing to 7 % at 21 km and 34 km, and to 14 % at 40 km (Fig. 8 of Nair et al., 2011). Despite of these large biases, Nair et al. (2011) concluded that Umkehr data are useful for studies of the long-term ozone evolution and for detecting drifts in satellite observations.

Miyagawa et al. (2014) compared Dobson zenith-sky Umkehr measurements with homogenized NOAA SBUV (Solar Backscatter Ultraviolet Instrument)(/2) 8.6 overpass data measured between 1977 and 2011. The mean bias between Dobson and SBUV partial ozone column varied between -12 % for Layer 7 to +3 % for Layer 2 (Fig. 1a of Miyagawa et al., 2014).

The biases reported in the three studies quoted above are comparable or larger than the differences between our Global-Umkehr retrievals and MLS and sonde measurements, suggesting that Umkehr results derived from global spectral irradiance can provide data of similar accuracy than the established zenith-sky method. A portion of the retrieval/MLS difference could also be caused by systematic errors in the MLS dataset, considering that the MLS accuracy specified by Froidevaux et al. (2008) is in the 5 to 10 % range.

Results presented in Fig. 5 illustrate that interquartile and 10th – 90th percentile ranges for the retrieval/MLS comparison on one hand and the MLS 2 / MLS 1 comparison on the other are similar for most layers. This suggests that a large portion of the observed retrieval/MLS differences can be attributed to actual changes in the ozone profile over the time periods relevant

for these comparisons. However, a portion of the change in the MLS profile from one day to the next may be caused by the relatively poor horizontal resolution of MLS profiles of about 200 km. In addition, some variability in the MLS dataset can be attributed to the slightly different geolocation of two consecutive overpass profiles. For example, the average horizontal distance between the locations of Summit and the MLS overpass is 160 km. Further analysis revealed that the difference between the MLS 1 and MLS 2 datasets depends also on the time when the daily MLS observation takes place. For example, when MLS 2 data are from the observation period close to local solar noon (14:11 to 15:10 UTC) and MLS 1 data are measured close to sunrise (5:28 to 6:26 UTC), MLS 2 data for Layers 7 – 9 are biased high by 3-6% relative to the MLS 1 dataset, while MLS 2 data for Layer 10 are biased low by 8%. This time-of-day dependency and its variation with altitude is by and large consistent with diurnal variations of the ozone profile measured by various instruments at Mauna Loa, Hawaii (Parrish et al., 2014), and by a microwave radiometer at Bern, Switzerland (Studer et al., 2014). This suggests that the time-of-day effect observed at Summit is caused by actual diurnal changes of the ozone profile rather than potential time-dependent systematic errors in the MLS dataset.

Another source of variability in the retrieval/MLS and retrieval/sonde comparisons is the different vertical resolutions of MLS (about 2-3 km), sondes (0.1 km) and our Umkehr retrievals (about 10 km for $\sigma_a = 0.4$ and about 25 km for $\sigma_a = 0.1$).

If measurement and forward model were without error, an Umkehr profile would resemble the actual profile smoothed by the

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AKs. To reduce the effect of the differing resolution, the higher-resolution MLS profiles could be convolved with the AKs of the Umkehr profile prior to comparing the two profiles. This technique has for example been applied by Nair et al. (2011) when comparing lidar and SBUV profiles. We did not use this method because it artificially reduces the true difference that is observed when comparing a high-resolution (sonde, MLS) with a low-resolution (Umkehr) profile. Of note, Nair et al. (2011) found that the smoothing technique does not make a significant difference in seasonally averaged data such as those presented in Fig. 5 and Table 2.

The bias between Umkehr retrievals and MLS or sonde data is reduced when correcting the forward model for the systematic error presented by the pseudospherical approximation. It is interesting to note that the correction is only in the -0.5 to 2.0 % range (Fig. 1b) but reduces the retrieval bias by up to 4 % (Layer 6 in spring, see Table 2). Considering that the uncertainty of our measurements is 3 % (1σ), systematic errors in the measurement vector in the 2 - 3 % range could conceivably be responsible for the remaining bias of Umkehr and MLS profiles indicated in Fig. 5 and Table 2. To test this hypothesis, we modified the measurement vector within reasonable limits and recalculated the profiles. We found that the bias between Umkehr and MLS profiles cannot be significantly reduced further, suggesting that the bias cannot be attributed to measurement errors alone.

The difference between results corrected for cloud effects and uncorrected results is very small, implying that clouds affect the accuracy of the retrievals only marginally. However, this conclusion may not apply to locations with thicker clouds or locations affected by aerosols and should be tested if the method is used at other sites.

If S_e is well defined, the most important parameter to optimize the results is σ_a . The objective is to find the right balance between sensitivity to the *a priori* profile on one hand and sensitivity to (unavoidable) errors in the measurement vector or forward model on the other. We chose $\sigma_a = 0.1$ and 0.4 . The smaller value quantifies the standard deviation of the actual variability of the ozone profile at Summit. While calculations with this value lead to good results, the solution may not be optimal for profiles at the fringe of the distribution (e.g., result for Layer 3 in Fig. 3). A small σ_a also results in a small value of d_s . However, statistics for results calculated with $\sigma_a = 0.1$ and 0.4 are quite similar (Table 2), suggesting that any value for σ_a between 0.1 and 0.4 leads to acceptable profiles. Determining the best value for sites other than Summit requires consideration of the measurement system and variability of the ozone profile at this site.

There are various ways to optimize the Global-Umkehr method for specific applications or locations. For example, if two instruments were to measure side by side, the uncertainty used to set up S_e could be better estimated by comparing the measurements of the two systems. Furthermore, the method to set up S_a could be modified to take into account that the

variability of the ozone profile depends on altitude (Eq. (4) uses the same standard deviation σ_a for all layers). [More current ozone absorption cross section data could be used \(e.g., Orphal et al., 2016\) than the Bass and Paur \(1985\) data implemented in this work and by OMI. If temperature profile data are available, these could be utilized to account for the temperature dependence of the ozone absorption cross section. Wavelengths, bandpass, and SZAs used for the measurement](#)

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vector could be further optimized to reduce uncertainties related to the Ring effect or the temperature dependence of the ozone absorption cross section. For example, by degrading the spectral resolution (currently set to 2 nm), the impact of the Ring effect could be diminished. Finally, the MYSTIC Monte Carlo model, which was used to calculate the correction function $R(\theta)$ (see Fig. 1b), was run with a scalar radiative transfer solver, which did not take polarization into account.

5 Lacis et al. (1998) calculated that modelling errors for irradiance resulting from the omission of polarization in these calculations can be as large as 1.3% for a Rayleigh atmosphere. However, errors for 310 and 337 nm (i.e., the wavelengths used in the Global-Umkehr method) agree to within 0.1%. We therefore conclude that the omission of polarization is not an import error source in our calculations.

10 We used *a priori* profiles that are independent of the total ozone column. Zenith-sky Umkehr retrievals from Dobson instruments that have historically been processed with the algorithm developed by Mateer and DeLuisi (1992) used TOC-dependent *a priori* profiles to constrain the retrieval. While this practice can lead to artefacts when calculating trends (Petropavlovskikh et al., 2005; Stone et al., 2015), the approach may be the best choice if a profile with the smallest uncertainty possible is sought for a specific purpose.

15 The Global-Umkehr method was tested with spectroradiometric measurements from a polar location because we only operate instruments at high-latitude sites. Inversions using high-latitude data are more challenging compared to retrievals for lower latitudes because of the limited range of SZAs at polar regions, the long time that is required to scan the range of SZAs necessary for the retrieval, and the high short- and long-term variability of the ozone profile. We have therefore confidence that the method would work well for mid- and low-latitude locations. Confirmation of this assertion is subject of future tests.

20 5 Conclusions

An optimal estimation method has been developed to retrieve vertical ozone profiles from measurements of global spectral irradiance in the UV. The method is similar to the widely used zenith-sky Umkehr technique, which inverts measurements of zenith sky radiance. To our knowledge, this is the first time that the Umkehr technique was applied to measurements of global irradiance. High-quality measurements of global spectral irradiance are now available for more than 25 years at

25 several NDACC locations (De Mazière et al., 2017), and the Global-Umkehr method has the potential to make these long-term datasets available for studying changes in the vertical distribution of ozone.

Compared to the standard zenith-sky Umkehr method, multiple scattering effects become more important when exploiting global irradiance measurements, which also include contributions from photons received from directions close to the horizon. Therefore, the sphericity of the viewing geometry needs to be taken into account. We have shown that this

30 challenge can be overcome by using a forward model with pseudospherical approximation plus additional corrections.

The method was evaluated with spectroradiometric measurements from Summit, Greenland, and validated with balloon sonde and MLS observations. For calculations using the corrected forward model, the bias between our retrieved profiles and

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MLS observations ranges between -5% (Layers 5 and 6) and $+3\%$ (Layer 2&3). The magnitude of this bias is comparable, if not smaller, to values reported in the literature for the standard [zenith-sky](#) Umkehr method. The distribution of the difference between retrieval and MLS observations was quantified with the interquartile and 10th – 90th percentile ranges. Depending on altitude, the interquartile ranges vary between 7 % and 13 % and the 10th – 90th percentile ranges run between

14 % and 24 %. Of note, interquartile ranges calculated from the differences of two MLS profiles that were measured on consecutive days vary between 5 % and 10 %, suggesting that a considerable portion of the retrieval/MLS differences can be attributed to real changes in the ozone profile. For Umkehr Layer 2 and higher, retrieval/MLS and retrieval/sonde differences are by and large consistent. The poor sensitivity of the Umkehr method to the altitude range of Layer 0&1 leads to are relatively large scatter (e.g., the interquartile range is 25 %) of the retrieval/sonde differences for this layer.

10 The effect of the parameter σ_a , which controls the sensitivity of the solution on the *a priori* profile, was extensively assessed. It was found that results calculated with a small value of $\sigma_a = 0.1$ (emphasis on *a priori*) generally agree to within

2 – 3 % of those calculated with a large value of $\sigma_a = 0.4$ (emphasis on measurements). By setting σ_a to a large value, retrieval errors may occasionally become large if the measurement vector is affected by unforeseen conditions (e.g., changing ozone, variable clouds). For example, the maximum retrieval/MLS difference was 50 % for $\sigma_a = 0.4$ but only

15 32 % for $\sigma_a = 0.1$.

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The retrieved ozone profiles were integrated over altitude. The resulting TOCs agreed almost ideally with TOCs measured by OMI: depending on the correction method, the retrieval/OMI bias ranged between -0.2% and 0.7% with a standard deviation of less than 2.0 %.

While the Global-Umkehr method was only tested for a high-latitude site, we are confident that it will also work at lower latitudes, but this assertion requires confirmation by future tests.

Data availability

“Version 2” spectra from the SUV-150B spectroradiometer at Summit are available from the Arctic Data Center at <https://arcticdata.io/>.

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30 [comments.](#)

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