Dear Stephan Bansmer and Co-authors,

thanks for your detailed responses and the revised manuscript which is greatly improved.

There are some further points that need revision before this manuscript can be published, see below:



General comments:

Your manuscript contains an overwhelming number of figures. Several of them could be combined into fewer figures (details below), other figures might not be as relevant that they would need to appear in the main manuscript. Please consider to move some of your figures into a supplement.

Formatting of units: (citing from the manuscript preparation guidelines) "Also units in the denominator should be formatted with negative exponents (e.g. km h-1 instead of km/h)."

(Page and line numbers refer to manuscript version 4 in your MS records, i.e. the version prior to the nomenclature version.)

Specific comments:

Abstract:

- **p1, I5:** "...static temperature..." = static air temperature?
- p1, l8: "... than usually used..." what is usually used?
- p1, I8/9: "Ice water contents up to 20 g=m3 can be simulated." "Simulated" sounds like you talk about numerical model simulation, maybe better "generated"?
- p1, I9: "... measurement techniques for particle sizing have performed..." Either "have been performed" or "are performed".
- p1, l11: "Finally, some applications of the icing wind tunnel are mentioned." Maybe replace "mentioned" with "presented" or "introduced".

Introduction:

- p2, I30: You introduce "static air temperature" but then talk about "static temperature". Please be consistent throughout the manuscript!
- p3, I5: "...for which the model..." Which model? It is the first time a model is mentioned, so you should explain more.
- p3, I7: "super-cooled" throughout the manuscript you use "supercooled" as well as "super-cooled", be consistent!
- p3, l30: "European Aviation Safety Agency (EASA)" already introduced in line 16, so use acronym only.

Section 2:

p4, l27: duplicate "was"

p4, l28/29: Better: "The space available at the installation site suddenly becomes a further design constraint."
p5, l8: "temperature bound" = "temperature threshold"?
p5, l11: "Figure 46" – figures should be numbered in the order they appear in the text!
p7, I9: "PLC" – this acronym has not been introduced!
p7, l16: figure13 – again, figures should be numbered in the order they appear in the text!
p7, I24: Check format of units.
p7, l32: "introduced" – better "installed"
p8, l1 : "where as" – "whereas"
p11, l19: " sealed with another acrylic compound." – specify "another" or use "an" instead of "another"
p11, l20: ", no steps" – do you mean "gaps"
p12, l9: Is there a reference for TAUICE?
p12, l13: duplicate "of"
Section 3
p13, Section 2.9: Can you specify accuracies for the static pressure probes and the Vaisala HMT-337?
p14, l32/Figure 16: In the introduction you introduce Median Volume/Mass Diameters as important parameters, why do you use d _{eq} here?
p14, l31 – p15, l14: I wonder whether this paragraph could be shortened. What value does it add to the manuscript:
p16, l5: "IWT" – the acronym has not been introduced yet, please write out
p16, l32/33: "ice water content IWC" – IWC had been introduced in the introduction, no need to write out here.
Section 4:
p17, l15: "icing-wind-tunnel" – be consistent, either "icing-wind-tunnel" or "icing wind tunnel"
p17, l19: PDI – what does this acronym stand for?
p18, l1: typo: "user?s preference."
p18, l12-32: Is there no reference for the IKP?
p18, l17: "water vapour" – be consistent american english (or british english) throughout the manuscript.

p18, **I31/32**: Why do you turn off the shattering artefact removal feature of the software for the wind tunnel measurements? Don't you expect particle shattering to happen here, too? Especially in high IWC or LWC conditions! Furthermore, it becomes hard then to compare in situ field measurements and wind tunnel measurement when this treatment differs!

p19, **l15/16**, **l24**: "2DS" vs. "2D-S" – be consistent.

p23, 17: "...higher concentrations of small particles..." – could this not be due to particle shattering introducing numerous small particles (as the removal treatment was turned off)?

Section 5

p23, I27: "European Aviation Safety Agency" = EASA (acronym already introduced)

p24, l17: "generic test articles" – do you mean "test objects" or "test models"?

p25, **l10/l22**: "3-D" vs. "3D" – be consistent!

Acknowledgements:

p26, l12: typo "contact" – contract

Figures:

As mentioned before, the figures should be numbered in the order they are mentioned in the text.

Furthermore, I find that you have a huge number of figures. I suggest you have a good revision of which figures are important and less important figures could be moved to a supplement. Also, maybe some figures could be grouped together, or even eliminated. See following suggestions:

- The manuscript contains a number of technical figures, these could potentially be grouped together (and then labeled as subfigures), e.g. Figures 6 and 7; Figures concerning the ice particle generation and dosing, including the ice particle conveyance pipe; maybe you have further ideas where figures would fit together thematically. I also feel that some of these figures might be well suited for a supplement.
- Figures that could be joined together and labeled as subfigure (a) and (b): 14 + 15, 21 + 19, 25 + 26, 27 + 28, 34 + 35, 36 + 37, 39, 41 + 42
- Figure 45 is not needed and can be removed what value does it add to the manuscript?
- Figure 23 + 24: You show images for the IKP and HSI probes, but not for PDI and 2DS probes. Are these figures needed, what value do they add to the manuscript?

Figure 41: Where exactly is the cylinder? not as obvious as in Figure 42.) Why are the side views of figures 41 and 42 so different, is Figure 41 a model simulation? Also, I think these two figures could be joined into one.

Units in figures: "Also units in the denominator should be formatted with negative exponents (e.g. $km h^{-1}$ instead of km/h)."

Design, Construction and Commissioning of the Braunschweig Icing Wind Tunnel

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Abstract. Beyond its physical importance in both fundamental and climate research, atmospheric icing is considered as a severe operational condition in many engineering applications like aviation, electrical power transmission and wind-energy production. To reproduce such icing conditions in a laboratory environment, icing wind tunnels are frequently used. In this paper, a comprehensive overview on the design, construction and commissioning of the Braunschweig icing wind tunnel is given. The tunnel features a test section of $0.5 \text{ m} \times 0.5 \text{ m} \times 0.5 \text{ m}$ with peak velocities of up to $40 \text{ m/s} 40 \text{ ms}^{-1}$. The static air temperature ranges from -25° C to $+30^{\circ}$ C -25° C to 30° C. Supercooled droplet icing with liquid water contents up to $3 \text{ g/m}^3 3 \text{ g m}^{-3}$ can be reproduced. The unique aspect of this facility is the combination of an icing tunnel with a cloud chamber system for making ice particles, which are more realistic than those usually used for mixed phase and ice crystal icing experiments. Ice water contents up to 20 g m^{-3} can be generated. We further show, how current state-of-theart measurement techniques for particle sizing have performed on ice particles. The data is compared to those of in-flight measurements in mesoscale convective cloud systems in tropical regions. Finally, some applications of the icing wind tunnel are mentioned presented.

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1 Introduction

5 Atmospheric icing affects the operational performance of many man-made devices and ground structures. Examples include icing of electrical power networks (Farzaneh (2008)), wind turbines (Hochart et al. (2008)), communication towers (Mulherin (1998)) and aircraft (Gent et al. (2000)).

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Based upon the understanding of the physics and chemistry of clouds (Lamb and Verlinde (2011)), several icing mechanisms are well-known to atmospheric scientists and engineers: supercooled droplet freezing (Jung et al. (2012)) – including freezing drizzle and rain (Zerr (1997)), snow accretion (Makkonen (1989)) and mixed-phase icing (Currie et al. (2012)). Particularly, in the field of aircraft icing, supercooled large droplet (SLD) icing with median droplet diameters above 50μ m 50μ m is considered as an independent category (Politovich (1989)).

Supercooled means that the droplets are in a metastable liquid state although their temperature is below the freezing point. Solidification can take place through homogeneous or heterogeneous nucleation. In absence of any impurities like aerosols, homogeneous nucleation starts with exceeding a barrier of free enthalpy triggering the creation of critical embryos. The change of free enthalpy is a function of latent heat of fusion (promoting phase change) and energy needed to create a new water-solid interface (hindering phase change). With increasing supercooling the latent heat of fusion increases thus making homogeneous nucleation more probable. Droplet impact on a solid substrate promotes *heterogeneous nucleation* initiating a fast development of ice dendrites in the liquid (Schremb and Tropea (2016)).

- 15 The following important parameters are of particular significance in defining the boundary conditions of the icing process:
 - Classical, fluid-mechanical testing parameters like **Reynolds number** $Re = \frac{\varrho}{\mu} U_{\infty} c$, $Re = \frac{\varrho_{\text{air}}}{\mu_{\text{min}}} U_{\infty} c$, **Mach number** $Re = \frac{U_{\infty}}{\mu_{\text{min}}} U_{\infty} c$, $Re = \frac{\varrho_{\text{air}}}{\mu_{\text{min}}} U_{\infty} c$, $Re = \frac{\varrho_{\text{min}}}{\mu_{\text{min}}} U_{\infty} c$, $Re = \frac{$
- Water concentration in the air. On the one hand, this is the Liquid Water Content (LWC), a measure of the mass of water per unit volume of air. Typical values for atmospheric icing conditions range from 0.1 g/m³ to 3 g/m³ 0.1 g m⁻³ to 3 g/m³ 0.1 g m⁻³ and have been measured in the updraft of continental convective systems e.g. above the Amazon basin, see Wendisch et al. (2016) and Braga et al. (2017). On the other hand, in the presence of ice crystals, an Ice Water Content (IWC) has to be specified. Inside continental and oceanic mesoscale convective systems, IWC-values are about 1 g/m³ 1 g m⁻³ (Gayet et al. (2014)). While ice water content in mid-latitude cirrus is generally low (Voigt et al. (2017)), peek values of 6 g/m³ 6 g m⁻³ have been observed in deep convective systems. Commonly, a Total Water Content (TWC) is defined as the sum of LWC and IWC.
 - Median volume diameter (MVD) of the statistical distribution of water droplets in the liquid cloud. MVD attempts to reduce the size distribution to a single, representative scalar diameter. This idea is attributed to LANGMUIR (Langmuir (1961); Finstad et al. (1988)). In the presence of ice crystals, a median mass diameter (MMD) is defined because of their variable density.

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Static air temperature. For water droplet icing, static air temperatures close to ^{0°} € 0°C promote glaze ice formations, whereas rime ice formation becomes predominant at lower static air temperatures (Wagner (1997)).

- **Humidity**, i.e. the amount of water vapor vapour in the air, is an important parameter for mixed-phase ice accretion (Currie et al. (2012)). Jung et al. (2012) showed its relevance for pure droplet icing. A distinction is made between absolute humidity AH and relative humidity RH. Absolute humidity is defined by the mass of water vapor vapour per unit volume of dry air. With the partial pressure of water vapor propor vapour proporty vapour, the static air temperature T of air T_{air} and the specific gas constant of the vapor T_{air} vapour T_{air} and the specific gas constant of the vapor T_{air} vapour T_{air} and the specific gas constant of the partial pressure of water vapor T_{air} vapour T_{air} vapour T_{air} and the specific gas constant of the partial pressure of water T_{air} vapour T_{air} vapour T_{air} vapour T_{air} and the specific gas constant of the partial pressure of water T_{air} vapour T_{air} vapour T_{air} vapour T_{air} to the equilibrium vapor vapour pressure of water T_{air} vapour T_{air} va

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- Accumulation time t_{acc} , for which the <u>test</u> model is exposed to the cloud of <u>super-cooled</u> supercooled droplets and/or ice crystals.
- Among the above enumeration of boundary conditions, temperature is of particular importance, because it governs the freezing dynamics of impacting super-cooled supercooled water droplets. At very low temperatures, the droplets will solidify shortly after their impingement and entrap the surrounding air. The resulting ice accretion is called rime ice, see figure 1. With increasing temperature, the solidification process of impacting droplets is retarded, yielding wind-driven water film dynamics. At locations with increased convective heat transfer, the water film freezes, yielding a nearly transparent glaze ice shape with typical horn formations.

The choice of boundary conditions for conducting a specific icing wind tunnel experiment depends on the application under investigation. When testing to support industrial product developments, certification requirements are often defining the boundary conditions. In Europe, icing of large transport aircraft is addressed in the document CS25 (EASA (2016)) of the European Aviation Safety Agency (EASA). In its Appendix C, information on the standard LWC-MVD envelope can be found, where the liquid water content is in the range between 0.1 and 2.9 g/m³ 2.9 g m⁻³ and MVD between 15 and 50 μ m50 μ m . Icing certification for wind energy converters is governed by the technical standards of the International Electrotechnical Comission Commission (IEC). Among others, the *Det Norske Veritas Germanischer Lloyd* (DNV GL) is an accredited registrar and classification society responsible for certification of wind energy converters, having created calculation rules and detailed specification based on the IEC, see GL (2010). It has to be noted that industrial products are in many cases significantly larger than the dimension of a typical icing wind tunnel. Hence, appropriate scaling laws are necessary to overcome the limited ranges of air speed, test section dimension and icing cloud characteristics in icing wind tunnels, see Bond and Anderson (2004).

Besides supercooled droplet icing, the aircraft industry has become aware of another icing hazard in the recent past. The phenomenon called *ice crystal icing* has been identified to effect ice accretion on heated aircraft assembly such as engine compressor blades or stagnation pressure probes. Thrust losses and engine damages as well as biased flight parameter display and loss of the autopilot can be caused. Ice crystal icing mostly appears in tropical regions in the vicinity of convective cloud systems. A comprehensive treatment on the topic was published by Mason et al. (2006). The Federal Aviation Administration (FAA) of the United States and EASA extended their icing regulations to Appendix D and P for ice crystal icing conditions in

2010 and 2011. Enhanced research activities on the topic have been promoted lately by NASA (Flegel and Oliver (2016)) and the National Research Council of Canada (Currie et al. (2014)).

Icing is a very challenging field of study that incorporates aspects of meteorology, fluid mechanics, thermodynamics, physics and engineering. Despite the fact that many researchers have been involved in the study of ice accretion over numerous decades, the physics of this phenomenon is still not completely understood. Icing wind tunnels are therefore an essential pillar to advance our knowledge in face of this multidisciplinary challenge. Enduring changes in technical standards or certification processes to continuously improve the safety against icing related incidents further emphasize the industrial need for icing wind tunnel testing. The cost estimations of the Ice Protection Harmonization Working Group (IPHWG) for aircraft certification with respect to SLD-icing underlines this requirement, see figure 2.

Against this background, the design and construction of the Braunschweig Icing Wind Tunnel began in 2010. It was the goal to contribute a tunnel with sufficiently large dimensions in the test section to support both fundamental and applied icing research with reasonable low operating costs. During the design, construction and commissioning process, many lessons have been learned, which can not be found in the literature and thus form the outline of this publication. The major components of the Braunschweig Icing Wind Tunnel and their design constraints are presented in chapter 2. The special topic of mixed phase icing, where both supercooled droplets and ice crystals are involved in the icing process, is treated in chapter 3. Together with the international co-authoring partners, the commissioning of the tunnel was realized, which is presented in chapter 4. Finally, some applications and test results that show the tunnel's capabilities are highlighted, with some concluding remarks on health and safety considerations.

2 Design and Construction of the Braunschweig Icing Wind Tunnel

2.1 Overall Design

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The aim to build a low-budget icing wind tunnel for research purposes entailed major choices on the operational envelope. The main targets in icing wind tunnel design are to obtain a homogeneous distribution of flow velocity, temperature, and a uniform icing cloud in the test section. Therefore, it is the test section, which governs the design choices.

The first choice was on the dimension of the test section. Addressing customers in aviation, automotive and energy industries requires reasonable aerodynamic testing where the wind-tunnel wall boundary layers are significantly smaller than the size of the test section. In this regard, investigations with Reynolds numbers up to $2 \cdot 10^6$ are a frequent demand. Moreover, the test section shall allow for mounting significant aircraft subsystems and particle sizing instrumentation. Hence, a cross section of $\frac{0.5 \text{ m} \times 0.5 \text{ m}}{0.5 \text{ m} \times 0.5 \text{ m}} \times 0.5 \text{ m}$ was considered as a lower bound for sizing the icing wind tunnel. Additionally, it has to be taken into account that the test section size is in relation with the overall dimension of the icing wind tunnel. The space available at the installation site suddenly becomes a further design constraint.

The maximum velocity is primarily determined by the power of the wind tunnel drive, see section 2.2. However, when considering the heat balance in icing wind tunnels, see section 2.3, the importance of the chilling device becomes evident. Figure 3 shows cost estimations for both drive and cooling power, which is based on price lists of leading manufactures in their respective field. Realizing that the cost for tunnel cooling is about 10 to 20 times more expensive than for the tunnel drive, the maximum velocity in the test section is in direct relation to the available investment costs for building the tunnel. Given an upper bound for the tunnel drive power, a compromise between test section size and maximum velocity has to be made. Consistent with the above requirements of test section size, the maximum velocity was estimated to be $\frac{40 \text{ m}}{\text{s}} \cdot \frac{40 \text{ m} \text{ s}^{-1}}{\text{s}}$. A larger test section size would decrease the maximum velocity given a constant drive power. Since many aeronautical applications even demand for higher speeds, the decisions of a maximum velocity of 40 m s^{-1} and a test section of $0.5 \text{ m} \times 0.5 \text{ m}$ were taken.

The lower temperature bound limit inside the test section was set to $-20 \,^{\circ}\text{C} - 20 \,^{\circ}\text{C}$. Below this bound, one usually observes rime ice accretions on the test models. In contrast, glaze ice formation with temperatures closer to $0 \,^{\circ}\text{C} - 0 \,^{\circ}\text{C}$ is significantly more complex from the perspective of fundamental research that shall be conducted in the present icing wind tunnel. An overview of the entire tunnel and its adjacent components is given in Figure 36. A comparison with the specifications of selected other icing wind tunnels can be found in table 1.

2.2 Tunnel Drive

Given an air temperature of $-20^{\circ}C$ $-20^{\circ}C$, a maximum speed U_{∞} and a cross-sectional area $AA_{\rm lsec}$ of the tunnel test-section, the required jet power $P_{\rm jet}$ can be calculated to

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$$P_{\text{jet}} = \frac{1}{2} \varrho_{\text{air}} U_{\infty}^{2} \cdot \dot{V}_{\text{air}}$$

$$= \frac{1}{2} \varrho_{\text{air}} U_{\infty}^{2} \cdot \left(U_{\infty} \cdot A_{\text{tsec}} \right) = \frac{1}{2} \varrho_{\text{air}} U_{\infty}^{3} \cdot A_{\text{tsec}}$$

$$= \frac{1}{2} \cdot 1.38 \frac{kg}{m^{3}} 1.38 \text{ kg m}^{-3} \cdot \left(40 \frac{m}{s} 40 \text{ ms}^{-1} \right)^{3} \cdot 0.25 m^{2} 0.25 m^{2}$$

$$\approx 11 kW 11 kW ...$$

In wind-tunnel design, a power factor λ is introduced, which indicates the ratio of fan power $P_{\rm fan}$ and jet power $P_{\rm jet}$. For a conventional, closed-loop, low-speed wind tunnel, a value of $\lambda \approx 1.5$ can be assumed. The fan power has to compensate pressure and skin friction losses of the individual tunnel components. High losses are caused by the wall boundary layers of the test section and the tunnel diffuser. Furthermore, the first and the second corner after the diffuser create significant losses. Since the vanes of the first corner are subject to ice accretion, their losses are tremendously higher compared to conventional closed-loop tunnels. Consequently, the power factor was dimensioned to a large value of $\lambda = 2.3$, yielding a fan power of $\frac{P_{\rm tan}}{P_{\rm tan}} \approx 25 \, {\rm kW}$. An axial fan with an efficiency $\eta_{\rm fan} \approx 67\%$ was selected for the tunnel drive, the electrical power

input is therefore $\frac{37kW}{37kW}$. Using a frequency converter, the speed of the three-phase asynchronous motor of the fan drive can be controlled, allowing variable tunnel speeds from 5 to $\frac{40m}{s}$ 40 ms⁻¹.

2.3 Chilling Device

To adjust and maintain a constant temperature of the air inside the icing wind tunnel, which is well below the freezing point of water, a chilling device is necessary, whose cooling capacity has to exceed the power of the wind tunnel drive.

2.3.1 Power Estimation

There are three major heat sources that have to be compensated by the chilling device: the power of the tunnel drive, the heat input through the wind-tunnel walls, and the heat transfer of the water spray. The latter is composed of the *sensible heat* $\dot{Q}_{sensible}$ to super-cool supercool the water spray from the temperature at which it leaves the pneumatic atomizer (about $20^{\circ}C$ - $20^{\circ}C$) down to the air temperature inside the tunnel test-section and the *latent heat* \dot{Q}_{latent} . As soon as the supercooled water droplet impacts on a solid substrate, the drop solidifies due to the heterogeneous nucleation and releases the latent heat. Both sensible and latent heat are a function of the liquid water content (LWC)

$$\dot{m}_{
m water} =
m LWC \cdot \dot{V}_{
m air}$$

$$=
m LWC \cdot \left(U_{\infty} \cdot A_{
m tsec} \right)$$
5 $\dot{Q}_{
m sensible} = \dot{m}_{
m water} \cdot c_{
m sensible}$, water $\cdot \Delta T$

$$\dot{Q}_{
m latent} = \dot{m}_{
m water} \cdot c_{
m latent latent, water} \quad ,$$

where $\dot{m}_{\rm water}$ is the water mass flow rate of the spray atomizers and ΔT the temperature difference, with the values for demineralized water

$$c_{\underline{\text{sensiblep,water}}} = \underbrace{4.183 \frac{kJ}{kg \cdot K}}_{4.183 \text{ kJ kg}^{-1} \text{ K}^{-1}}$$

$$c_{\underline{\text{latent latent,water}}} = \underbrace{334 \frac{kJ}{kg}}_{334 \text{ kJ kg}^{-1}} .$$

For typical operational conditions of the icing wind tunnel, i.e. $U_{\infty} = 40m/s$, $\Delta T = 30K$. $U_{\infty} = 40 \,\mathrm{ms}^{-1}$, $\Delta T = 30 \,\mathrm{K}$, and approximations for the fan power and the heat input through the wind-tunnel walls, the necessary cooling capacity can be estimated, see figure 4. When LWC is in the range of $\frac{1g/m^3}{1 \,\mathrm{g\,m}^{-3}}$, a continuous cooling power of $\frac{30kW}{30 \,\mathrm{kW}}$ is necessary. To allow for a five-minute peak load of $\frac{1WC = 3g/m^3}{1 \,\mathrm{m}^{-3}}$ and static tunnel LWC = $3 \,\mathrm{g\,m}^{-3}$ and static air temperatures down to $\frac{20^{\circ}C}{1 \,\mathrm{m}^{-20}}$. In the tunnel, the system was dimensioned for a maximum cooling power of $\frac{80kW}{1 \,\mathrm{m}^{-20}}$.

2.3.2 Refrigeration Unit

To provide the icing wind tunnel with 30 kW 30 kW continuous cooling capacity and 80 kW maximum cooling capacity, a customized cooling unit was built, see the upper right part of figure 36. Core of the system is a refrigeration unit 4 that cools the heat transfer fluid (Therminol® D12) down to a temperature of -32°C -32°C. To compensate for the volume change of the fluid, a surge tank 4 filled with nitrogen is installed. The cold fluid is then stored in a 4m³ 4m³-sized buffer tank 5, which acts as a hydraulic switch; on its right side, it is connected to the refrigeration unit. On its left side, it is connected to the heat exchanger in the wind tunnel. As a result, the operation of the refrigeration unit is decoupled from the heat exchanger in the wind tunnel. Using a pump and electrically operated valves 6, the volume flow from the buffer tank to the heat exchanger in the wind tunnel is controlled. Depending on this flow rate, the continuous cooling power of 30 kW 30 kW of the refrigeration unit can be well exceeded. For a 80kW-80 kW peak load operation, the complete buffer tank is first cooled to -32°C -32°C, and afterwards the chilled heat transfer fluid is pumped through the heat exchanger at the wind tunnel within a short period of 7 minutes. The recovered waste heat from the cooling process is fed into the regenerative heating system of the institute building. Because of its cold-resistant properties, the piping of the cooling system and the storage tanks are made of stainless steel 1.4541, and then insulated with a 4 cm thick insulation layer made of Kaiflex®. The entire system is controlled by PLCprogrammable logic control.

2.3.3 Heat Exchanger

The heat transfer between the Therminol® D12 fluid and the circulating wind tunnel air takes place through two consecutive heat exchangers, see figure 5. While flowing from the bottom to the top of the first heat exchanger, the Therminol® D12 fluid increases in temperature. With only one heat exchanger, this would yield an undesired vertical temperature gradient in the passing air flow. The second compartment, in which the Therminol® D12 fluid flows the opposite way, thus compensates the vertical gradient in air temperature. The entire assembly is installed in front of the third corner in the wind tunnel, see also figure 12. The large distance to the test section promotes a homogenization of the temperature distribution over the tunnel cross-section through turbulent effects.

Mechanically, the heat exchangers have a cross-sectional area of $\frac{1.6 \times 1.6 \,\mathrm{m}^2}{1.6 \,\mathrm{m} \times 1.6 \,\mathrm{m}}$ and are composed of elliptical tubes made of stainless steel and rectangular, smooth steel fins, which are firmly connected to each other by hot galvanizing. The fin pitch is $\frac{3.5 \,\mathrm{mm}}{3.5 \,\mathrm{mm}}$. To protect the galvanized steel surface against corrosion, it is coated with ZACOSIN® 2000Q, an epoxy resin-based protective coating that provides high thermal conductivity by embedded aluminum aluminium particles, see also section 2.8. Alternatively, the heat exchanger could also have been made entirely of stainless steel. Compared to standard steel, however, stainless steel has a significantly lower thermal conductivity ($\frac{50 \,\mathrm{W}}{\,\mathrm{m}}$ K compared to $\frac{15 \,\mathrm{W}}{\,\mathrm{m}}$ K which is why the surface of the heat exchanger would have to be much larger.

The heat exchangers also include a condensation drainage. Condensation occurs predominantly in the initial cooling phase

from room temperature to cold operational temperatures. When continuously running at low temperatures, the effect of condensation is only minor.

2.4 Settling Chamber

To obtain a low turbulence level inside the test section of the tunnel, a settling chamber with honeycomb and screens is introducedinstalled. An influence on the longitudinal (parallel to the main flow direction) and lateral flow structures (perpendicular to the main flow direction) has to be distinguished (Scheiman and Brooks (1981)). The honeycomb reduces both the lateral components of mean velocity and of the larger turbulent eddies, where as whereas the screens reduce the longitudinal components of turbulence or mean-velocity variations across the sectional area. For the Braunschweig Icing Wind Tunnel, a combination of screen - honeycomb - screen - screen was selected. The screens have a mesh size of 6mm 6 mm, the honeycomb diameter is 10mm 10 mm, its length is 100mm 100 mm.

2.5 Spray System

A spray system is required to produce a uniform drop distribution in the tunnel test section. In the past five years, two different spray systems have been developed; one system suited for large droplets ($\frac{\text{MVD} \approx 80 \,\mu\text{m}}{\text{MVD}} \approx 80 \,\mu\text{m}$) and a high liquid water content ($\frac{1.1 \, \text{g/m}^3 < \text{LWC} < 3 \, \text{g/m}^3 \cdot 1.1 \, \text{gm}^{-3} < \text{LWC} < 3 \, \text{gm}^{-3}$) and a second spray system for low MVD (8 μ m < MVD < 48 μ m) and low LWC (0.1 $\, \text{gm}^{-3} < \text{LWC} < 2 \, \text{gm}^{-3}$).

The spray system for large droplets is composed by a grid of 5×5 air-assisted atomizers. Inside one atomizer, a thin jet of water is destabilized by the shear forces of a transversely directed air stream and finally breaks down into small droplets. The diameter of the water jet, which is in the range of $100\mu m$ 100 μm , determines the water volume flux as well as the drop size. The higher the applied air pressure, the higher the aerodynamic shear forces, and the smaller the drop size. In order to enable a high variability of the drop size distribution for basic tests, the atomizers are designed in a modular manner, that is, the water jet and the air atomizer cap can be exchanged separately.

Stainless steel has been chosen as a material for the atomizers, since otherwise conventional atomizers made of nickel-plated brass are prone to wear due to frequent temperature changes inside the icing tunnel (from $+20^{\circ}$ C to -20° C \times C \times

The high liquid water content $(1.1 \,\mathrm{g\,m^{-3}} < \mathrm{LWC} < 3 \,\mathrm{g\,m^{-3}})$ provided by the spray system is a technical constraint of airassisted atomizers. Below a certain water flow rate that is needed for low LWC, their operation becomes unstable. This is because the low water pressure inside the atomizer cap prevents the formation of a stable water jet, which is needed prior to atomisation. The second spray system circumvents that problem and was specially designed for low liquid water contents, see

figure 6. To achieve a LWC of $\frac{0.1 \text{g/m}}{0.1 \text{ g m}^{-3}}$, a water volume flux can be estimated to

$$\dot{V}_{\underline{\text{H}}_{2}\underline{\text{O}}_{\text{water}}} = \frac{\text{LWC} \cdot \dot{V}_{\text{air}}}{\varrho_{\text{H}_{2}\underline{\text{O}}}} \frac{\text{LWC} \cdot \dot{V}_{\text{air}}}{\varrho_{\text{water}}} = \frac{\text{LWC} \cdot A_{\text{test section}} \cdot U_{\infty}}{\varrho_{\text{H}_{2}\underline{\text{O}}}} \frac{\text{LWC} \cdot A_{\text{tsec}} \cdot U_{\infty}}{\varrho_{\text{water}}}$$

$$= \frac{0.1 \frac{\text{g}}{\text{m}^{3}} \cdot 0.25 \,\text{m}^{2} \cdot 40 \,\frac{\text{m}}{\text{s}}}{1000 \,\frac{\text{kg}}{\text{m}^{3}}} \frac{0.1 \,\text{g} \,\text{m}^{-3} \cdot 0.25 \,\text{m}^{2} \cdot 40 \,\text{m} \,\text{s}^{-1}}{1000 \,\text{kg} \,\text{m}^{-3}}$$

$$= 3.6 \frac{1}{\text{h}} 3.61 \,\text{h}^{-1} \quad . \tag{1}$$

- 5 Using a spray matrix of 5×5 atomizers, this yields a volume flux of $0.14\frac{1}{h}$ $0.141h^{-1}$ per atomizer, which is ten times lower than the typical minimum flow rate of air-assisted atomizers in industrial environments. In consequence, electrically-actuated air-assisted atomizers have been chosen for the spray system. The atomizers are pulsating up to 10000 times a minute, making the spray appear to be constant. By changing the duty cycle, the LWC can be adjusted without modifying the pressure of water and air supply, which is a major advantage compared to classical air-assisted atomizers.
- Water and air are conditioned for the operation of the air-assisted atomizers, see figure ??again figure 6. A pressurized vessel of stainless steel serves as a reservoir of water, whose temperature can be regulated between 2° C and 80° C and 80° C. Water pressure and air pressure can be adjusted separately between 0.5 and 9 bar. Water flow rate and air flow rate are measured and adjusted by flow regulators. Possible pollutant particles, which could block the nozzles, are separated with water filters. In larger icing test facilities, where tunnel air temperatures below -40° C can be adjusted, the air duct of the spray atomizers is heated in order to prevent droplet freeze out. In our setup, droplet freeze out is less likely due to the minimum temperatures of around -20° C, which are well above the homogeneous nucleation limit of water. Therefore, the atomizer's air duct is not heated, and air at room temperature is used to operate the atomizers.

2.6 Tunnel Nozzle Design

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Having passed the settling chamber, the flow is accelerated in the tunnel nozzle that contracts towards the dimensions of the test section. Beyond altering the turbulence structure (Uberoi (1956)), the shape of the tunnel nozzle also influences the trajectories of the water droplets that are injected by the spray system. Especially large droplets and ice particles are affected.

Two different nozzle geometries have been considered. These are

- the AVA nozzle, with its contour given by

$$\frac{z\left(x\right)}{z_{\text{out}}}\frac{z\left(x\right)}{z_{\text{tsec}}} = \left(\sqrt{K} - 1\right)\left(1 - \frac{x}{L}\right)^{3}\left[2 - \left(1 - \frac{x}{L}\right)^{3}\right] + 1, \text{ and}$$

- the Witoszinski nozzle, with its contour given by

$$\frac{\frac{z\left(x\right)}{z_{\text{out}}}\frac{z\left(x\right)}{z_{\text{tsec}}} = \frac{1}{\sqrt{1 + \left(\frac{1}{K} - 1\right)\frac{\left[1 - \left(\frac{x}{L}\right)^{2}\right]^{2}}{\left[1 + \frac{1}{3}\left(\frac{x}{L}\right)^{2}\right]^{3}}}} \quad ,$$

where x, z are spatial coordinates, L is the nozzle length, $\frac{z_{\text{in}}, z_{\text{out}}}{z_{\text{sc}}, z_{\text{ISC}}}$ are half of the inlet and outlet diameter of the nozzle, and K is given by $\frac{K = \left(\frac{z_{\text{in}}}{z_{\text{out}}}\right)^2 K = \left(\frac{z_{\text{sc}}}{z_{\text{out}}}\right)^2}{z_{\text{coc}}}$. To provide enough time for the super-cooling supercooling process of the droplets on their way to the tunnel test section, a large nozzle length of 3.5 m 3.5 m is foreseen.

The droplet trajectories are mainly governed by inertial forces, drag and gravity. A water droplet with diameter d, velocity vector $\begin{pmatrix} u_d \\ w_{ol} \end{pmatrix}$ is transported by the air velocity $\begin{pmatrix} u_{air} \\ w_{ol} \end{pmatrix}$ resulting in a drag force

$$\boldsymbol{D} = \frac{\varrho_{\rm air}}{2} \left| \Delta \boldsymbol{u} \right|^2 d^2 c_{\underline{\boldsymbol{d}} \, \underline{\mathcal{D}}} \cdot \frac{\Delta \boldsymbol{u}}{\left| \Delta \boldsymbol{u} \right|} \quad , \label{eq:delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_delta_de$$

where $\Delta u = \binom{u_{\rm air}}{w_{\rm air}} - \binom{u_d}{w_d}$ represents the slip velocity. Assuming laminar flow around the spherical droplets with a mass of m_d at low Reynolds numbers $\frac{e}{Re} = \frac{\varrho_{\rm air} \cdot |\Delta u| \cdot d}{\mu_{\rm air}}$ $\frac{e}{Re} = \frac{\varrho_{\rm air} \cdot |\Delta u| \cdot d}{\mu_{\rm air}}$, the drag coefficient $\frac{e}{d} \cdot c_D$ can be approximated to $\frac{e}{d} = \frac{6\pi}{Re} \cdot c_D = \frac{6\pi}{Re}$. Applying the principle of d'Alembert, this yields a system of differential equations, given by

$$m_d rac{\partial inom{u_d}{w_d}}{\partial t} = m{D} + m_d m{g}$$
 ,

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where t represents the time and q is the acceleration due to gravity.

The effect of gravity on the droplet trajectories is depicted in figure 7, assuming an air velocity of $\frac{40 \text{ m}}{s} 40 \text{ ms}^{-1}$ at the nozzle exit. Small droplets with a size of $\frac{60 \mu \text{ m}}{60 \text{ µm}}$ follow the streamlines of air flow inside the nozzle. With increasing droplet size, the gravitational deflection becomes predominant, particularly at the low speed regions of the nozzle. For $\frac{300 \mu \text{ m}}{200 \text{ µm}}$ droplets, their trajectory from the lowest spray-bar even collides with the nozzle wall. Such undesired collisions alter the size distribution of the droplet cloud and have to be avoided.

Figure 8 demonstrates the effect of the nozzle geometry on the droplet slip velocity. Note that after the nozzle length of L = 3.5 m 3.5 m, a straight segment of 0.5 m in length is attached. In the low speed region x < 1.5 m 1.5 m, the AVA nozzle accelerates the flow less than the Witoszinski nozzle. Consequently, the slip in velocity between droplet and air is larger in that region. Downstream, for $\frac{1}{8} = 3 \text{ m} x > 3 \text{ m}$, the slip is caught up. While the small droplet of $\frac{60 \mu \text{ m}}{8} = \frac{60 \mu \text{m}}{8} = \frac{60 \mu \text{m}}{8} = \frac{300 \mu \text{m}}{8} = \frac$

we observed that the surrounding low speed air-flow in the settling chamber rapidly decelerates the droplet. This effect on the droplet trajectory is thus negligible in our tunnel.

Furthermore, the cooling process of the droplet on its trajectory has to be considered. Cooling is promoted by a convective heat transfer at the interface between droplet and air. When the droplet is smaller than $\frac{100 \, \mu \, m}{100 \, \mu m}$, the surface temperature penetrates the entire droplet volume in about $\frac{10 \, \text{ten}}{100 \, \text{ten}}$ milliseconds, assuming a Fourier number of 1 for the heat conduction problem, see also Yao and Schrock (1976). Hence, isothermal cooling can be assumed as an approximation for the thermal energy balance, yielding a differential equation for the droplet temperature T_d

$$\frac{\pi}{6} d^3 \varrho_d c_{\rm p,\,d} \frac{\partial T_d}{\partial t} = {\rm Nu} \cdot \pi d^2 \left(T_{\rm air} - T_d \right) \cdot \frac{k_{\rm air}}{d} \quad ,$$

where ϱ_d and $c_{p,d}$ represent the density and the specific heat capacity of the droplet, k_{air} is the heat conductivity of air, and Nu is the Nusselt number, estimated by $\frac{Nu = 2 + 0.6 \, Pr^{\frac{1}{3}} \, Re^{\frac{1}{2}}}{10.6 \, Pr^{\frac{1}{3}} \, Re^{\frac{1}{2}}} = \frac{1}{10.6 \, Pr^{\frac{1}{3}} \, Re^{\frac{1}{3}}} = \frac{1}{10.6 \, Pr^{$

Using the above equation, the temperature evolution of the droplet inside the tunnel nozzle is plotted in figure 8 for a given static air temperature of $T_{\rm air} = 263.15$ K. After one meter traveling metre travelling distance, the droplet of $\frac{60 \,\mu}{\rm m} \, \frac{1}{\rm m} \, \frac{1}{\rm$

In summary, droplet conditions up to a maximum size of $\frac{150 \,\mu\text{ m}}{150 \,\mu\text{m}}$ can be well provided in the present tunnel environment, tanking into account a reasonably low deflection due to gravity, a low slip velocity and sufficient droplet cooling. Furthermore, the AVA geometry was selected for the wind tunnel nozzle, because it enables a slightly higher heat transfer to supercool the droplets.

2.7 Tunnel Walls

2.7.1 Mechanical Design

To minimize the heat transport through the wind tunnel walls, they consist of three different layers with a total thickness of 80 mm 80 mm, see figure 9. Between two layers of aluminum aluminium, an insulating foam with low heat conductivity is introduced. The inner aluminum aluminium shell which is potentially exposed to the icing cloud is anodized for a protection against corrosion. The flange of each wall segment is made of purenit®, a highly compressed material based on PUR/PIR rigid foam providing a high thermal insulation. A rubber seal prevents the water from penetrating into the wall. The flanges were fastened with many screws, thus avoiding water leakage. In order to make the joints fully watertight, the inner groove was

sealed with an acrylic compound. Due to an excellent manufacturing accuracy, no steps between adjacent wall segments are observed.

2.7.2 Aerodynamic Wall Effects

Icing wind tunnel tests usually involve models with large dimensions to minimize uncertainties when applying scaling laws for ice accretion. The flow around these models thus interferes with the walls of the test section. Two different wall effects have to be considered. On the one hand this is the viscous effect of junction flow (Simpson (2001)), which occurs when the boundary layer of the wind tunnel wall encounters the test model attached to the same wall. The adverse pressure gradient forces the wall boundary layer to separate with horseshoe vortices that wrap around the obstacle. These unsteady vortices result in high turbulence intensities, high surface pressure fluctuations and heat transfer rates, which finally alter the ice accretion in that area. This effect was computationally studied for a NACA0012 airfoil with the icing code FENSAP (Beaugendre et al. (2003)), see figure 10. FENSAP solves the Reynolds-averaged Navier-Stokes equations using a finite element approach, calculates droplet trajectories with an Eulerian scheme and finally determines the ice accretion using a model of Messinger (1953). The upper part of figure 10 clearly identifies the horseshoe vortex while plotting the spanwise velocity in multiple slices and an isosurface along the NACA0012 airfoil. Especially in the vicinity of the leading edge at the junction between airfoil and wind-tunnel wall, the two-dimensional flow behaviour is significantly deteriorated. Because a glaze ice case was simulated $(T_{\infty} = -10\,^{\circ}\text{C}, \text{LWC=1}\,\text{g}\,\text{m}^{-3}$, ice accumulation time $t_{acc} = 120\,\text{s}$), the velocity field strongly influences the heat transfer. At the junction, the decreased heat transfer attenuates the ice accretion, which is shown in the lower part of figure 10. However, the two-dimensional ice shape distant to the tunnel wall, which covers 80% of the wing span, remains unaffected.

The second case of wind-tunnel wall interaction is an inviscid effect. The distant walls that are not connected to the test model will influence the streamlines around the model. Thereby, both pressure and shear stress distribution along the model surface are altered. Interestingly, at low speeds and sufficiently large droplet sizes (MVD=15 µm), the ice accretion is nearly unaffected, when tunnel walls are introduced, see the computational results of the icing code TAUICE in figure 11.

2.8 Corrosion

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Since the operation of the icing wind tunnel entails the usage of water droplets or ice particles, corrosion has to be considered during the design process. The most vulnerable part of the tunnel in this regard is the heat exchanger, as it is made from a mix of steels, which is galvanized together, see section 2.3.3. Without any proper protection, many factors decide on the rate of corrosion of a galvanized surface, the most important for this application being the chemical composition of the water and the type of surface contact.

Water with a high content of carbonates supports the creation of a protective layer of alkaline zinc carbonates $ZnCO_3$ on the galvanized surface. However, with no or only few carbonates and a relatively high content of oxygen, the water reacts with the zinc to form zinc hydroxide $2 Zn + 2 H_2O + O_2 = 2 Zn(OH)_2$. This further forms a compound $2 ZnCO_3 \cdot 3 Zn(OH)_2$,

also known as *white rust*, which does not have the protective function of zinc-carbonates, but only connects loosely to the surface. Similar to normal rust this porous layer keeps water near the surface and delays drying, which further increases the rate of corrosion. Over time, particles flake off, until the zinc is used up.

- 5 The type of surface contact between water and the galvanized surface is important for the rate of corrosion. If the water is sprayed onto the surface and remains as small droplets, each droplet forms a small corrosion element with a large surface that promotes the supply of oxygen. This increases the corrosion rate in comparison to the case, where a galvanized surface is completely submerged in water with only a small contact area to the surrounding air.
- 10 To protect the galvanized steel surface of the heat exchanger against corrosion, it is coated with ZACOSIN® 2000Q, an epoxy resin-based protective coating that provides high thermal conductivity by embedded aluminum aluminium particles.

Another metal in contact with the deionized water is brass, which is used in several valves, pipe connections, etc. Brass forms tarnish, a thin dark layer of stable metal oxides, that protects the base metal. Unlike rust, the oxidation reaction is self limiting. As soon as the layer is formed, no more metal is oxidized. The tarnish did not pose any problem in all parts.

The models for the tunnel are made from aluminium or glass-fibre reinforced plastics (GRP). None of these materials showed any signs of reaction with the water so far. The same holds for the inner wall of the wind tunnel which is made of anodized aluminium. aluminium Even with scratches from installing models no signs of corrosion showed up.

20 2.9 Basic Instrumentation

To monitor the icing wind tunnel operation and thus the boundary conditions of an experiment, several probes are necessary, see figure 12. The velocity U_{∞} in the test section is determined by two static pressure probes p_1 and $p_2 \cdot p_{sc}$ and p_{tsec} . Their typical accuracy for long term stability is $\pm 0.3\%$ of the full scale output. Using the Bernoulli-equation

$$p_{\underline{1}\underline{\text{sc}}} + \frac{\varrho}{\underline{2}} v_1 \frac{\varrho_{\text{air}}}{2} u_{\text{sc}}^2 = p_{\underline{2}\underline{\text{tsec}}} + \frac{\varrho}{\underline{2}} \frac{\varrho_{\text{air}}}{2} U_\infty^2 + \zeta_{\text{nozzle}} \frac{\varrho}{\underline{2}} \frac{\varrho_{\text{air}}}{2} U_\infty^2 \quad ,$$

where $v_{\text{T}} u_{\text{SC}}$ is the velocity in the settling chamber at p_{T} , ϱp_{SC} , ϱ_{air} is the density of air, ζ_{nozzle} is the specific loss coefficient of the wind tunnel nozzle, and a continuity equation

$$v_1 u_{\rm sc} \cdot A_{1 \rm sc} = U_{\infty} \cdot A_{\infty \rm tsec},$$

where A_1 and A_{∞} A_{sc} and A_{tsc} are the cross sectional areas of settling chamber and test section, one can estimate

$$U_{\infty} = \frac{p_2 - p_1}{\frac{\varrho}{2} \left[\left(\frac{A_{\infty}}{A_1} \right)^2 - 1 - \zeta_{\text{nozzle}} \right]} \frac{p_{\text{tsec}} - p_{\text{sc}}}{\frac{\varrho_{\text{air}}}{2} \left[\left(\frac{A_{\text{tsec}}}{A_{\text{sc}}} \right)^2 - 1 - \zeta_{\text{nozzle}} \right]} . \tag{2}$$

Note that the effect of pressure losses in the wind tunnel nozzle is usually neglected, since $\left(\frac{A_{\infty}}{A_{1}}\right)^{2} \gg 1 \left(\frac{A_{\text{lsec}}}{A_{\infty}}\right)^{2} \gg 1$ and $\zeta_{\text{nozzle}} \ll 1$.

A VAISALA HMT-337 monitors both the relative humidity RH and the total temperature T_{tot} in the settling chamber. The uncertainty for the humidity measurement is $\pm 1.8\%$ RH, the accuracy for the temperature measurement is around ± 0.3 °C. T_{tot} is used to calculate the static air temperature T_{∞} in the test section

$$T_{\infty} = T_{\rm tot} - \frac{U_{\infty}^2}{2c_p} \frac{U_{\infty}^2}{2c_{\rm p,air}} \quad ,$$

5

where e_p $c_{p,air}$ is the specific heat capacity of air. Moreover, the temperature information is necessary to determine the density e $e_{p,air}$ of air in equation (2). The power of the tunnel drive P_{fan} and the pressure loss over the heat exchanger Δp_{HE} indicate the performance degradation of the icing wind tunnel due to ice accretion. Furthermore, the state of the spray system S_{spray} is controlled, see section 2.5.

3 Ice Crystals and Mixed Phase Capability

- Ice crystals in the atmosphere have to be considered for aircraft safety, since they partially melt in warm environments and develop a "sticky" character. In particular, heated stagnation pressure probes and engine compressor stator blades can be maleffected by ice crystal icing (Mason et al. (2006)). Icing of aircraft probes can cause false flight parameters displayed inside the cockpit. Ice accretion inside the compressor causes flow blockage, forcing the compressor to operate towards stall conditions. The compressor encounters a decay in rotational speed resulting in significant thrust losses (rollback event) (Oliver (2014)).
- 20 Moreover, total engine flame out may appear if a huge mass of accumulated ice is shed into the combustor.

To provide experimental capability on ice crystal icing, the Braunschweig icing wind tunnel was upgraded with an ice crystal generation and conveyance system, which is presented in this section, see also figure 36.

3.1 Morphology of Ice Crystals

Ice crystal icing conditions are typically encountered in wide ranging convective cloud systems of high ice water content at flight altitude. Such conditions can especially be found in the vicinity of mesoscale convective cloud systems in tropical re-

gions, see Grzych and Mason (2010) and Leroy et al. (2015). In order to better document ice crystal icing conditions, two flight campaigns have been conducted in the course of the HAIC- and HIWC project (Dezitter et al. (2013), Strapp et al. (2016)). The first campaign took place in Darwin, Australia in 2014 during the monsoon period, the second campaign in Cayenne, French Guiana, in 2015 during the rainy season. The flight measurements have been conducted in high ice water content cloud areas in large tropical mesoscale convective systems, mostly over the oceans. Details about the campaigns and about data treatment can be found in Leroy et al. (2016b), Leroy et al. (2016a) Leroy et al. (2016a) and Leroy et al. (2017).

Atmospheric ice particles feature a broad diversity of sizes and shapes which depend on the individual ice particle growth history affected by ambient temperature and super-saturation. Figure 13 and ?? show (top) shows examples of ice particle images captured close to $-10\,^{\circ}C$ during the Darwin campaign with the 2D-Stereo probe (see section 4.5 for a description of the instrument). Images from figure 13 correspond to a stratiform part of the cloud where rather constant ice water contents close to $\frac{1.0 \text{ g/m}^3}{1.0 \text{ gm}^{-3}}$ were sampled. Images from figure ?? 13 (bottom) were recorded in convective cores with IWC peak values exceeding $\frac{3 \text{ g/m}^3}{3 \text{ gm}^{-3}}$. In the convective part, ice particles are more numerous and smaller. Close to $-10\,^{\circ}C$ column and capped column type crystals have been found. Larger ice particles ($\frac{5000 \text{ µm}}{600 \text{ µm}} > 600 \text{ µm}$) are rare and resemble graupel (dense and roundish particles). On the contrary, in the stratiform part of the cloud, ice particles larger than $\frac{600 \text{ µm}}{600 \text{ µm}} = \frac{600 \text{ µm}}{600 \text{ µm}$

Figure 14 shows ice particle size and mass distributions (PSD and MSD) depending on ambient temperature using the equivalent diameter d_{eq} for size definition. The size equivalent diameter intends to provide a size information on a non-spherical particle. It is defined as the diameter of a circle of the same area as the shaded pixelnumber for each particle. The size and mass distributions have been averaged for the selected temperature regimes, only cloud areas with total water contents above $\frac{1.0 \text{ g/m}^3}{1.0 \text{ g m}^{-3}}$ have been taken into account. The concentrations of small ($\frac{200 \mu m}{1.0 \text{ g m}^{-3}}$) and large (>1 mm>1 mm) ice particles vary in opposite ways with temperature: For colder temperatures, concentrations of small ice particles increase whereas the number of large particles decreases (cf. figure 14). This temperature dependency might be a consequence of several cloud processes. Nucleation of new ice particles is favored favoured at low temperatures, creating new small ice crystals. On the opposite, growing of ice crystals by collection processes requires larger particle sizes and might be more efficient at higher temperature. Regarding the dynamics of the cloud, small ice particles can also easily be carried aloft by updraft winds. Therefore, there is still no clear and unique scenario describing the formation of high IWC areas in clouds. Cumulative ice particle mass distributions are plotted in the bottom part of figure 14. PSDs have been converted in to MSDs following the work of Fontaine et al. (2014) and Leroy et al. (2016b). Bin masses are linked to the particle size using a power-law relationship $m(D) = \alpha D_{eq}^{\beta} m(d) = \beta d_{eq}^{\gamma}$. However, α and β and γ are not constant; for each time step, β is deduced from the analysis of particle images and $\alpha \beta$ is constrained by additional TWC measurements thus ensuring that the total mass from the MSDs equals the measured TWC. With decreasing atmospheric temperature the cumulative particle mass is carried more and more by small ice particles. The median mass diameter (MMD), reduces from roughly $750 \,\mu m$ at $-10 \,^{\circ}C$ to $320 \,\mu m$ at $-50 \,^{\circ}C$. More details about MMDs in high IWC cloud regions can be found in Leroy et al. (2017).

As mentioned above, ice crystal icing can cause malfunctions of aircraft engines due to inner ice accretions. Inside-engine conditions are characterized by ice particle sizes of about 20 µm 20 µm as ice particles fragment in the fan stage. Due to centrifugal forces, high ice particle concentrations can be found in the casing region of the core engine. The local TWC can exceed the atmospheric TWC by a factor of 4 or even more, see Feulner et al. (2015).

3.2 Ice Crystal Production in a Cloud Chamber

It has been aimed to reproduce closest possible replicates of natural ice crystals. Cloud chamber technology has been identified to be appropriate, because natural ice crystal growth is simulated in an artificial cloud. Usually, cloud chamber technology is applied in meteorological science to investigate ice crystal formation and growth mechanisms and to study the interaction of individual ice crystals, see Connolly et al. (2011). Thus, the productivity of cloud chambers is accepted to be rather low. For icing wind tunnel studies, huge amounts of ice particles are required. In collaboration with the Austrian Neuschnee GmbH, two highly productive cloud chambers have been developed and installed inside a cooling room. Figure 15 and 36 show both cloud chambers and auxiliary equipment. Basically, atomized droplets are forced to freeze out inside the chamber by the expansion of cool pressurized air. Strong circulation of air keeps the particles in suspense until they grow to certain size and settle down to the bottom of the chamber. The production rate of the cloud chamber is not sufficient to enable a direct supply of ice crystals into the wind tunnel. Thus, the particles are collected inside a chest freezer, which is directly connected to the cloud chambers.

The chest freezers operate at 70 °C 70 °C to minimize degeneration and sintering of the particles. Currently, the production rate for both cloud chambers is limited to about 1kg/h1 kg h⁻¹. This production rate is sufficient to perform between 10 to 15 wind tunnel tests a day.

Microscopic images of ice crystals grown inside the cloud chamber are shown in figure ?? the upper right part of figure 15. The primary ice particle habit can be adjusted by variation of ambient temperature inside the cooling chamber. After storing in the chest freezer, the ice particles typically feature aggregates of individual crystals as further illustrated in the following section.

3.3 Ice Crystal Conveyance

To establish defined cloud conditions inside the IWT test section test section of the icing wind tunnel, ice particles are fluidized into an airflow and guided into the wind tunnel, see figure 36. The conveying airflow is extracted from the icing wind tunnel by a bypass construction including a radial fan. An external heat exchanger (aftercooler) cools the air to compensate heat input of the fan and environment. The aftercooler is connected to the refrigeration system of the cooling chamber. Downstream of the heat exchanger, the piping system enters the cooling room where ice particles are supplied to the airflow. Further downstream, the particle-laden flow is discharged into the IWT icing wind tunnel.

To adjust the ice water content inside the test section a defined massflow of ice particles $\dot{m}_{ice,inj}$ has to be supplied at a constant feed rate. Particle dosing is realized by a volumetric dosing machine, shown in figure 16. The dosing machine has a very linear operating behaviorbehaviour: the higher the frequency of the dosing machine motor, the higher the supplied ice particle massflow, see figure 17 (left). After storage and dosing, the particles are partially interlocked and exist as accumulated ice clumps, which are by far larger than the desired particle sizes of several tens of microns. Thus, the particles have to be sieved to desired sizes, which is realized by a custom made sieving machine, see again figure 16. The sieving machine is driven by an electric motor that forces a 800×800 mm 800 mm $\times800$ mm sieve to oscillate at a frequency of 8 Hz. The total mass of ice provided by the dosing machine does not pass the sieving machine. Greater clumps of ice are not fully broken by the sieving procedure and depose on the sieve. The average deposit is about $\frac{15}{15}$ to $\frac{20}{15}$ of the ice mass provided by the sieving machine.

The sieving machine is connected to the conveyance piping system by a conical flexible bag. Ice particles passing the sieve directly fall into the piping system and get dragged by the pipe airflow. To adapt the local pressure inside the piping system to the ambient pressure inside the cooling room, the piping system implies an injector nozzle that induces a local jet inside the pipe, see the lower part of figure 16. The nozzle diameter is carefully adjusted based on preliminary calculations of the pressure distribution inside the whole piping system. The temporal stability of the mass flow is depicted in figure 17 (right). The dashed line shows the ice particle mass supplied by the dosing machine at a frequency of 21 Hz. It is a linear fit, a constant feedrate has been proven by weightingweighing, as indicated by the red data points. The blue dots and the solid blue line represents masses of ice weighted weighted at the outlet of the sieving machine. This curve shows a very linear behavior behaviour as well. One can conclude that after about 20.30 seconds the ice particle massflow injected into the wind tunnel is temporally constant. Weighing of ice mass passed by the sieving machine has been repeated several times, also for lower feedrates and it turned out that after about 30 seconds the provided ice massflow is stable. Correlations between the injected ice particle massflow $m_{ice,inj}$ and the ice water content IWC IWC inside the test section are discussed in section 4.6. Figure ?? The lower left part of figure 16 shows a microscopic image of two ice particles which have been captured at the exit of the sieving machine. As one can see, the particles feature aggregates of tiny ice particles and can thus be considered as close replicates to the natural ice particles shown in figure 13and ??.

The ice particles are injected into the settling chamber of the icing wind tunnel, upstream of the spraybar system, see figure ??.

The velocity of the particle-laden jet is about 20 m/s20 ms⁻¹, which is 5 times higher than the local wind tunnel airspeed. Consequently, the jet mixes and expands by turbulent mixing with the ambient airflow so that the particles are spread among the flow field and mix with the droplets atomized by the spraybar system. The particle trajectories get contracted inside the wind tunnel nozzle. Due to the circular pipe exit the ice particles also cover a circular cross sectional area inside the test section. It has been tried to extend this area by using adapters at the pipe exit and by the use of several injection pipes distributed along the settling chamber. These efforts did not work out due to ice deposits inside the additional assembly, see Baumert et al.

(2016). Consequently, the simplest approach of the single pipe outlet has been chosen. For studies at glaciated and mixed phase conditions, screens and honeycombs in the settling chamber of the tunnel have to be unmounted to avoid ice accretion on these elements.

4 Commissioning of the Braunschweig Icing Tunnel

5 4.1 General Approach

The <u>icing-wind-tunnel icing wind tunnel</u> calibration has been performed with respect to the requirements specified in SAE ARP 5905. Accordingly, both an aero-thermal calibration of the airflow and a calibration of the icing cloud <u>have had</u> to be performed. After presenting the deployed measurement techniques, selected calibration results are presented. Finally, an instrumentation inter-comparison exercise was carried out.

10 **4.2 PDI Probe**

The measurement principle of PDI Phase Doppler Interferometry (PDI) is based on light scattering interferometry (Rudoff et al. (1990)) that uses as measurement scale the wavelength of light and as such its calibration is not as easily degraded. The scattering by spheres much larger than the instrument wavelength is approximated by geometrical optics. Size and velocity are determined by measuring sinusoidal scattering signals on adjacent photo-detectors as particles move through an interference fringe pattern formed in the intersection of two laser beams of same wavelength.

The method does not require frequent calibration since the light wavelength does not change and the detector separations that affect the size measurement are fixed. The sinusoidal nature of the signals detected may be used with the Fourier analysis approach to detect signals reliably even in low signal-to-noise (SNR) environments. Off-axis light scatter detection is used. The advantage of this approach is that a very small, well-defined sample volume can be formed which minimizes uncertainty when computing absolute concentration and reduces the possibility of coincident events (more than one particle residing in the sample volume at one time). The sample volume can be adjusted to balance count rate with coincidence rate to suit the user? s preference. The PDI only needs an initial factory calibration since parameters affecting the measurements as the laser wavelength, beam intersection angle, transmitter and receiver focal lengths, and the detector separation do not change with age or use of the instrument.

The method has been demonstrated to be capable of measuring drop size distributions, droplet velocity distributions, size-velocity correlation, droplet time of arrival, droplet spacing, droplet number density, liquid volume flux, and liquid water content. At the Braunschweig icing wind tunnel, the PDI has been used for size and concentration measurements of cold clouds and for liquid / ice discrimination in mixed phase conditions. Although the PDI method cannot provide a size measurement for ice particles, it is able to extract a velocity information due to the particle's reflectivity. The resulting velocity histogram

was used in an instrumentation inter-comparison campaign (see section 4.6.4) in order to check differences between the mean velocity of ice particles and airspeed upstream of the 2D-S probe volume.

4.3 Cranfield Isokinetic Probe (IKP)

The probe consists of two tubes, where one pipe end is faced towards the upstream direction, thereby collecting a representative sample of air, droplets and ice particles suspended within the flow, see also figure 18. The goal is to measure the total water content TWC of the flow. The volume flux $\dot{Q}_{\rm front}$ has to be adjusted to obtain isokineticity at the probe head. The second tube is faced downstream with its pipe end. Thus, a second sample of air, without any particles or droplets in it, can be collected to determine the water vapour content so that the quantity of water in condensed form (solid and liquid) can be determined.

A high electric current is passing through the metallic structure of the probe head thus enabling its resistive heating. Consequently, the probe stays free of an ice build-up so that the flow conditions at the entrance to the probe can be carefully matched to oncoming flow in order to achieve the correct sampling, not too rich or depleted with respect to the condensed water. The heating of the inside of the probe drives all sampled water droplets and ice particles into the vapor vapour phase. A subsequent cooling system composed of copper tubes plunged into a water tank has been added in order to bring the air at a reasonable temperature (about 30-35 °C30 °C to 35 °C) before entering the measurement system.

The measurement system itself is mainly comprised of two parts. The first part is the mass flow meter measurement and valve to automatically control the mass flow at the probe inlet to maintain iso-kinetic conditions. The second part is the water vapor vapour concentration measurement which is done using a LICOR 7000 system.

The probe features a double wall construction making it possible to use it as a pitot static probe and also providing a means to heat it. A large current is passed through the outer wall, via the tip and back along the inner tube. Tube wall thicknesses are chosen to get the appropriate split of heat between the different parts of the probe.

4.4 HSI Probe

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The high speed imaging (HSI) probe (Baumert et al. (2016)) illuminates a volume using six laser beams and takes shadow images of traversing cloud particles. Thereby, depth of field errors and sampling bias due to particle obscuration are minimized. The shadow images are sampled at a frequency of 300 fps by a CMOS-chip of 640 by 480 pixels. Images of particles in the range of 5 to 1200 μ m 1200 μ m can be recorded, the resolution is 3.795 μ m/pixel3.795 μ mpixel $^{-1}$. The entire device including lasers and camera is remotely controlled. A trigger beam of different wavelength is coaxially aligned to the laser beams, see figure ?? (left). A receiver including an aperture of appropriate size and shape is elevated to 40 degrees to the transmitted beam and detects particles passing the object plane. It will trigger the lasers (laser pulse duration 12 ns) when particles are within the desired measurement volume. The HSI hardware is integrated in a cylindrical canister. Three arms including optical components are mounted on the front end of the canister that allow to take intrusive particle images, see also

figure ??... The image processing software is able to detect and to evaluate irregular shaped particles. Some examples of 2D particle images that were detected by the HSI probe, are shown in figure ?? (right). Particles that are rejected by and includes several particle validation criteria, e.g. out-of-focus or other validation criteriaare marked in red, valid particles are outlined in green colour. detection.

5 **4.5 2DS-2D-S Probe**

The two-dimensional stereo (2D-S) probe (Lawson et al. (2006), Lawson and Baker (2006)) detects the size and concentration of cloud droplets and ice particles in the size range of 10 to $\frac{1280 \ \mu m}{1280 \ \mu m}$ using shadow images of the cloud particles. Two orthogonal diode laser beams illuminate two linear diode arrays consisting of 128 photodiodes with $\frac{10 \ \mu m}{100 \ \mu m}$ pixel $\frac{10 \ \mu m}{100 \ m}$ pixel $\frac{10 \ \mu m}{100 \ m}$ resolution. When a particle crosses the laser beam in the sampling volume, its shadow image on the photodiode array is recorded by high-speed electronics. The diode lasers operate at 45 W and are single-mode and temperature-stabilized. This design with two lasers better defines the sampling volume boundaries and thus minimizes errors associated with the depths of field and the sizing of small particles.

The LaMP 2D-S data processing has been described in detailed in Leroy et al. (2016b) as it has previously been used for the treatment of the data sets collected during the HAIC field campaigns. The software is able to extract various size parameters form the particle images (area equivalent diameter, maximum dimension, etc.) and treats all major artefacts related to optical array probe's measurements, e.g. shattering, splashing, multiple particles in a single image, out of focus images. An area equivalent diameter was used in the present manuscript. For truncated images, the computation of the particle's size is based on the work of Korolev and Sussman (2000): all valid images are considered for the retrieval of the particle size distribution. The size of out-of-focus particles is corrected according to the work of Korolev (2007). The only difference between the analysis of the field campaign and the wind tunnel datasets lies in the shattering treatment as this artefact removal feature has been turned off for the wind tunnel measurements. Shattering of large particles is a function of IWC, see Field et al. (2006). The particle sizes in the wind tunnel were rather small (MMD around 80 µm) thus the fraction of shattered particles is expected to be less than 1%. Additionally, the high concentrations in the wind tunnel would have rendered an inter-particle time analysis difficult because the inter-particle time of shattered and non-shattered particles were similar.

4.6 Calibration Results

4.6.1 Aerothermal Calibration

An aerothermal calibration has been conducted to ensure and prove adequate airflow quality inside the test section. Investigations on airflow uniformity have been made by means of a five-hole-probe. The probe has been calibrated *a-priori* based on the procedure described by Treaster and Yocum (1978). The measurements have been performed for the IWT-standard setup standard setup of the icing wind tunnel with screens and honeycombs installed in the settling chamber and also for the ice-crystal-icing setup, where these elements have been unmounted.

Figure 19 (top) shows the contour of axial velocity inside the test section for the ice-crystal-icing setup. The pneumatic system for ice-particle conveyance was activated, the exit velocity of the pipe depicted in figure ?? was about 20 m/sconveyance pipe inside the settling chamber was about $20 \,\mathrm{ms}^{-1}$, however, no ice particles were fed into the system. The plot covers a cross-sectional area of $240 \times 400 \,\mathrm{mm}^2$ around the center-line $240 \,\mathrm{mm} \times 400 \,\mathrm{mm}$ around the centre-line position. Measurements have been taken in vertical and horizontal steps of $40 \,\mathrm{mm} 40 \,\mathrm{mm}$. One can clearly observe a footprint of the pneumatic-conveyance jet inside the flow-field. Velocity deviations relative to the center-line centre-line velocity are about $0.4 \,\mathrm{m/s} \cdot 0.4 \,\mathrm{ms}^{-1}$, which complies with ARP 5905 requirements of $+/-1 \,\% \pm 1\%$ spatial flow uniformity. Moreover, vectors of orthogonal velocity components are plotted, indicating a vortex inside the flow field. Since no screens and no honeycombs are installed, the vortex most likely originates from the turbulent flow interaction between conveyance piping, jet flow and surrounding tunnel air flow. The vortex magnitude is very low, local flow angularities comply with ARP 5905 specifications. For the FWT standard setup standard setup of the icing wind tunnel, no vortex can be observed and the flow field is very uniform as shown in figure ??19 (bottom).

Turbulence intensity has been measured by means of hot-wire anemometry. Figure 20 shows the contour plot of turbulence intensity for the ice-crystal-icing setup. The peak value of 1.4 percent 1.4% complies with ARP specifications. Again, the footprint of the pneumatic-conveyance jet is clearly visible.

Besides flow velocity characterization, measurements of total temperature have been carried out. A very uniform temperature distribution could be verified for the IWT-standard setup standard setup of the icing wind tunnel with flow straightener and screens installed. Local temperature deviations during the measurements correspond to temporal variation of tunnel temperature, which can be adjusted with an accuracy of $+/-0.5 \degree C \pm 0.5 \degree C$ around the target value. For the ice-crystal-icing setup, a footprint of the jet can be found in the temperature field, see figure 21. For a static wind tunnel temperature of $0 \degree C$ air temperature in the wind tunnel of $0 \degree C$, the local temperature at center-line position is about $-0.8 \degree C - 0.8 \degree C$. At lower temperatures, the discrepancy diminishes since the jet temperature and the airflow temperature converge. In case of rather low temperatures of $-15 \degree C - 15 \degree C$, the center-line temperature is about 0.4 K higher than the ambient temperature.

4.6.2 Calibration of the droplet cloud

Currently, two spray systems can be mounted inside the IWTicing wind tunnel, see section 2.5. The spray system with high LWC and high MVD has been calibrated by means of the Cranfield Isokinetic Probe (IKP), the canister PDI probe, and the modular HSI probe. IKP measurements have been conducted to correlate the liquid water content (LWC) in the test section with the water flow-rate that is supplied to the spray atomizers. The left image of figure 22 shows temporal plots of LWC obtained for various water flow-rates at $U_{\infty} = 40 \text{ m/s} \cdot U_{\infty} = 40 \text{ m/s}^{-1}$. As illustrated in the right image of figure 22, the dependency between LWC and flow-rate is rather linear. Test points corresponding to the plots in the left image are highlighted by the same colorcolour. For lower flow velocities U_{∞} , higher LWCs values are obtained as anticipated in equation 1. Nevertheless, the

LWC increase is not proportional to $1/U_{\infty}$, because the cross-sectional area of the icing cloud expands at lower wind-tunnel speeds.

Comparisons of liquid water contents derived from canister-HSI and canister-PDI measurements are plotted in the left image of figure 23. The dashed lines stress deviations of +/-10 percent $\pm 10\%$. The liquid water contents assigned to C-HSI and C-PDI have been derived based on volume to mass correlations of detected droplets. It can be observed that there is very good agreement between the measurements of all three probes except for two test points of the PDI at $U_{\infty} = 20 \text{ m/s} \cdot U_{\infty} = 20 \text{ ms}^{-1}$. These measurements had been affected by water condensation inside the C-PDI. The MVD varies in the range of 80 μ m to 95 μ m 80 μ m to 95 μ m, see the right image of figure 23. Higher MVDs have been observed with the C-PDI probe at $U_{\infty} = 40 \text{ m/s} \cdot U_{\infty} = 40 \text{ ms}^{-1}$. Presumably, droplet coalescence is promoted for higher air velocities. Yet the medium droplet size is rather constant for various liquid water contents at the same wind tunnel speed. All these results refer to the ice-crystal-icing setup with the jet switched on. No measurements have been taken away from center-line centre-line position as neither the canister probes nor the IKP could be traversed adequately. Ice accretion studies of selected test articles models have proved a very uniform droplet cloud.

15 4.6.3 Calibration of the ice particle cloud

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Ice-particle cloud calibration has been performed by means of the IKP and the canister-HSI. Figure 24 shows time histories of the ice water content (IWC) determined by the IKP for dosing frequencies $f_{\rm dosing}$ of 5 and 30 Hz at wind tunnel speeds U_{∞} of 20 and $\frac{40 \text{ m/s}}{40 \text{ ms}^{-1}}$. In the left image absolute values are plotted while the right image shows plots of IWC normalized by the target IWC. According to the IKP measurements, stable cloud conditions are established after 20 to 60 seconds, which is surprising since steady state particle supply is established after 20 to 30 seconds. There are two possible explanations. On the one hand, there might be a slight recirculation of ice particles inside the closed-loop wind tunnel – not every injected ice particle might settle down at the forth corner of the tunnel. On the other hand, the background humidity is slowly increasing due to sublimation of the ice particle cloud, which might influence the IKP measurement. Increasing the particle supply of the dosing machine to frequencies higher than 30 Hz does not result in higher IWC, because the permeability of the sieving mesh is limited. According to figure 24, a maximum IWC of about $\frac{19 \text{ g/m}^3}{19 \text{ g m}^{-3}}$ can be achieved at the center-line centre-line position. At f_{dosing} =5 Hz, a minimum IWC of $\frac{3 \text{ g/m}^3}{3 \text{ g m}^{-3}}$ can be adjusted. Lower feed-rates can be obtained by use of a smaller feed pipe of the dosing machine but have not been calibrated with respect to test section IWC. Figure 24 shows oscillations in IWC at low feed-rates because of an unsteady discharge of the dosing machine. This issue has been solved by the use of a rotating rod which supports particle discharge.

Based on the IKP measurements, a linear correlation between IWC and dosing frequency has been determined. Corresponding plots are shown in the left image of figure 25 for wind tunnel speeds of $U_{\infty} = 20 \text{ m/s}$ and $U_{\infty} = 40 \text{ m/s}$ $U_{\infty} = 20 \text{ ms}^{-1}$ and $U_{\infty} = 40 \text{ ms}^{-1}$. Due to the linear operational behavior behaviour of the sieving and dosing machine, a strong linear dependence between IWC and injected ice particle mass $\dot{m}_{ice,inj}$ could be demonstrated as shown in the right image of figure 25.

This dependency allows to assess the accuracy of IWC set for a test run. For wind tunnel tests which include ice particle supply, the IWC is adjusted based on the linear fits of figure 25, left. The frequency of the dosing machine is set according to the demanded IWC. After a test run, the ice particle deposit on the sieving machine is weightedweighed. Based on the correlation of dosing frequency and supplied massflow, see figure 17 (left), the ice particle massflow $\dot{m}_{ice,inj}$ can be determined respecting ice deposit on the sieve. Figure 25 then allows to deduce the real IWC which has been established inside the test section. Additional test campaigns where this procedure has been applied, showed an accuracy of IWC adjustments of about $\frac{15.96}{10.0000} \pm 1.5\%$.

Investigations on spatial IWC uniformity have been performed by means of a custom-made particle-collecting tube system. The system allows to measure a relative distribution of the ice particle concentration inside the test section. A contour plot for $\frac{U_{\infty} = 40 \text{ m/s}}{U_{\infty}} = 40 \text{ ms}^{-1}$ is shown in figure 26. Acceptable uniformity of $\frac{1}{20}$ percent $\pm 20\%$ according to ARP 5905 is given for a circular area of about $\frac{150 \text{ mm}}{150 \text{ mm}}$ in diameter. This area has proved to be sufficient and appropriate for valid icing experiments. In agreement with the turbulent footprint of the pneumatic conveyance jet, see figure 20, the peak IWC is located slightly above the center-line centre-line position, which has to be respected for the mounting of aerodynamic test models.

4.6.4 Probe inter-comparison

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Based on the IKP calibration, a cross comparison of various optical array probes has been performed in the Braunschweig Icing Wind Tunnel. The CNRS and DLR 2DS-probes2D-S probes, as well as the C-HSI have been mounted consecutively at the same position inside the test section. Additionally, modular HSI and PDI probes have been installed externally, to assess the repeatability of the test conditions. Figure 27 gives an overview of the test setup. M-PDI and M-HSI laser beams have been adjusted close to the sampling volume to perform non intrusive measurements in parallel.

Figure 28 presents a comparison of optical array probe measurements at seven glaciated cloud conditions named I01, I02, I03, I05, I06, I09 and I10. Ice water contents have been derived based on the mass-size relation described by Brown and Francis (1995). Up to a diameter of $97 \,\mu m$, the particles have been treated as spherical for the determination of the IWC. Above that threshold, the assumption of spherical particles would lead to an overestimation of the IWC. Ice water contents derived from measurements of both $\frac{2DS-probes-2D-S}{2D-S}$ probes and the C-HSI probe have been related to the estimated wind tunnel IWC. A rather good agreement between the measurements was proven as most of the measured values agreed with the wind tunnel prescribed IWCs with an accuracy of $\pm 20\%$, which is well within the limits of expected instrument accuracy.

Wind tunnel testing of the 2DS-probes 2D-S probes and the C-HSI have generated a great variety of ice particle images. Figure 29 shows examples of ice particles captured by the C-HSI. The aggregate structure which has been observed for ice particles before fluidization and conveyance, see figure 13, could be maintained. Figure 30 shows a series of ice particle images obtained from 2D-S measurements inside the Braunschweig WFicing wind tunnel. Note that the ice particle images were

oriented vertically in a post-processing step. It can be observed that most of the particles feature an irregular, elongated shape. Based on the analyses of ice particle images one can conclude that the particles partially break up during conveyance and are reduced to smaller sizes.

A characteristic size distribution of artificial ice particles inside the IWT icing wind tunnel is shown in figure 31 (top). Furthermore, particle size distributions of atmospheric ice particles (Atm) are included according to figure 14. The artificial ice particle cloud includes higher concentrations of small ice particles with equivalent diameters below 200 μm s200 μm. The amount of particles smaller than 200 μm 200 μm yielded from the wind tunnel measurements has not been identified during the airborne measurements, even for the coldest temperatures. The high amount of small ice particles is reflected in the cumulative mass distribution shown in figure 2231 (bottom). An MMD of about 79 μ m 79 μm was determined for the wind tunnel icing cloud. This value agrees with the initial target of an MMD in the range of 50 to 200 μm 200 μm. The HAIC flight campaigns have shown that against previous assumptions natural ice crystal clouds are characterized by higher MMDs in the range of 300 to 800 μm 800 μm depending on temperature, see figure 2231 (bottom). Yet the high number of small ice particles with MMDs below 80 μm 80 μm may be of relevance for particle size distributions in young and aged contrails (Voigt et al. (2010), Voigt et al. (2011)).

It can be summarized that the initial design targets of the ice-crystal generation system have been achieved. Irregular ice particles with close agreement to natural ice particles can be injected into the wind tunnel. Even hexagonal ice-crystals can be grown by cloud chamber technology and might be used for single particle impact studies. The ice particle size distribution is characterized by an MMD of about $\frac{80 \ \mu m}{80 \ \mu m}$ which is between atmospheric and inside-engine particle sizes. Because the ice-crystal generation process is fully separated from icing wind tunnel operation, the MMD does not change with operating conditions. IWC values of 3 to $\frac{20 \ g/m^3}{20 \ g m^{-3}}$ correspond to total water contents expected for inside-engine conditions and air data probes. Lower ice water contents can be adjusted by modifications of the dosing machine but require further wind tunnel calibration measurements.

25 5 Some applications and test results

5.1 Super-cooled Supercooled Large Droplets Icing

The consideration of super-cooled supercooled large droplets (SLD) icing is part of the certification for large transport air-craft (Politovich (1989)). Appendix O of the document EASA CS25 (EASA (2016)) of the European Aviation Safety Agency provides detailed information on the MVD and LWC envelope. Outstanding requirement is the simulation of liquid clouds with median volume diameters larger than 50μ m. The trajectory 50μ m. The trajectories of such large droplets show a large deviation from the air streamlines. With their increased inertia, the droplets might be impacting or running back in regions on the aircraft wing, which are not protected with de- or anti-icing systems.

Such SLD ice shapes and their surface roughness where investigated in the Braunschweig Icing Wind Tunnel (Steiner and Bansmer (2016)). In order to determine which wind tunnel conditions to use, they had to be related to the relevant dimensionless parameters influencing ice roughness. These are the dimensionless accumulation parameter A_c , given in Bond and Anderson (2004) as

$$5 \quad A_c = \frac{LWC \cdot U_{\infty} \cdot t_{acc}}{\rho_{\text{ice}} \cdot \delta} \frac{LWC \cdot U_{\infty} \cdot t_{acc}}{\varrho_{\text{ice}} \cdot \delta}$$

$$(3)$$

and the stagnation freezing fraction

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$$f_0 = f_0 \left(p, T_{\infty}, U_{\infty}, MVD, LWC, c, \delta \right) \tag{4}$$

where p_{ice} give represents the density of ice, δ the diameter of the inscribed circle of the airfoil nose as a reference length, free-stream pressure p and static free-stream temperature T air temperature T_{∞} .

Within this two-dimensional (A_c and f_0) parameter space, a Latin hypercube sampling was performed to define a set of experiments. The resulting ice shapes have been captured with a mould-and-cast method (Reehorst and Richter (1987)) and are shown in figure 32. They can be categorized in two groups. For shapes in the first group, which are marked in blue, the ice thickness was almost constant over some distance from the stagnation point and then a ridgeline appeared. Shapes in the second group were marked in red, and exhibited an intense yet smooth ice accretion at the stagnation line followed by a much rougher zone with a feather structure.

5.2 Mixed-Phase-Icing of a Cylinder

Comprehensive investigations on ice accretion of generic test articles models at mixed phase conditions have been performed in the recent past (Bansmer and Baumert (2017), Baumert et al. (2018)). The test articles models have been equipped with heat foils to investigate the effect of internal heat conduction on the accretion process.

Figure 33 shows ice shapes on a cylinder model without heat foil operation. A mixed phase cloud with a total water content of $\frac{12g/m^3}{12 \text{ g m}^{-3}}$ and a melting ratio m_r (defined as LWC/TWC in this setup) of 0.12 has been adjusted. The static air temperature has been varied between 0 and $\frac{-15 \, ^{\circ}\text{C}}{-15 \, ^{\circ}\text{C}}$. In contrast to supercooled droplet icing, the ice accretion for mixed-phase icing exhibits a conical shape. The ice accretion process appears to be strongly dependent on air temperature. At constant melting ratio and constant total water content, a decrease in icing severity can be observed for lower temperatures. This effect can be addressed to supercooling of the liquid droplets inside the mixed phase clouds. At lower temperatures enhanced supercooling allows a large amount of liquid to freeze with only a short delay upon impact. Therefore, less liquid is locally available to promote ice particle sticking resulting in a significantly reduced ice accretion. Water imbibition into the supposed porous ice layer seems therefore a necessary condition for increased icing severity.

Figure 34 gives an example of mixed phase icing of the cylinder with activated heat foils. In contrast to figure 33, no cone shaped ice accretion develops as the accretion layer constantly melts due to internal heating. Consequently, the accretion does not adhere to the cylinder surface and gets dragged towards the upper side of the cylinder by shear forces induced by the ambient airflow. Capillary forces prevent the accretion layer from instant shedding. As the accretion layer convects downstream, further ice accumulates in the flow stagnation area of the cylinder. In consequence, a cohesive ice accretion structure detaches from the cylinder and growths further downstream.

5.3 Development of Measurement Techniques

The experimental characterization of multiphase flows in icing wind tunnels can be very challenging and is a fruitful topic for developing measurement techniques as already shown in section 4. Here, we want to shortly summarize an optical technique to measure the 3-D-3D geometry of ice accretion.

Two properties of ice accretion are particularly demanding for optical measurement techniques. On the one hand, the ice surface geometry is very complex and includes both large roughness elements and deep chasms, see again figure 32. On the other hand, glaze ice formations appear transparent in the visible light range, which makes them hard to detect with conventional camera systems (Gong and Bansmer (2015b)). Therefore, we developed a general approach based on mid-infrared (MIR) laser scanning (Gong and Bansmer (2015a)), because MIR radiation penetrates ice and water only within a depth of less than 10 micrometers 10 µm.

The MIR-line scanning method uses a CO₂-laser light sheet, which is first created through proper lenses, and then projected onto the scene, see figure 35. The ice surface will absorb this light and emit an infrared radiation signature, which is visible as a deformed line. Using a properly calibrated camera, the deformed line can be transformed using a triangulation function yielding 3D-coordinates. While consecutively scanning with the laser light sheet over the object, the entire 3D ice geometry can be reconstructed. The scanning accuracy was validated with 3D ice objects of known dimensions; a pattern of frustums was used.

6 Health and Safety

Operating the icing wind tunnel with prototype test models requires intervention and human surveillance. In particular, activities inside the icing wind tunnel or in the cooling chamber (see 8 in Figure 36) can be very exhausting and have to be considered for health and safety. As depicted in Figure ??, unsafe behavior According to Ramsey et al. (1983), unsafe behaviour significantly increases when the ambient wet-bulb temperature is lower than 10 °C 10 °C, although workers are appropriately clothed. Frequent crew change is thus a necessary measure to perform these exhausting activities.

Another safety aspect is water handling with respect to legionella, see figure ??. These bacteria multiply with ideal growth rates between 32 to $42 \,^{\circ}\text{C} \, 42 \,^{\circ}\text{C}$ and can infect alveolar macrophages, which may result in several lung diseases. The water, which is to be sprayed inside the icing wind tunnel is increasing this threat since the contaminated droplets smaller than $10 \,\mu$ m $10 \,\mu$ m can be deeply inhaled. To minimize this risk, demineralized water is used that is further sterilized with ultraviolet radiation before entering the pneumatic atomizers.

Competing interests. The authors declare that they have no conflict of interest.

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References

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- Al-Khalil, K., Salamon, L., Tenison, G., et al.: Development of the Cox icing research facility, in: 36th Aerospace Sciences Meeting & Exhibit, AIAA, Citeseer, 1998.
- Bansmer, S. E. and Baumert, A.: From high altitude clouds to an icing wind tunnel: en route to understand ice crystal icing, in: Proceedings of the EUCASS Conference held in Milano, Italy, 2017.
 - Baumert, A., Bansmer, S., Sattler, S., Pervier, H., and Esposito, B.: Simulating natural ice crystal cloud conditions for icing wind tunnel experiments A review on the design, commissioning and calibration of the TU Braunschweig ice crystal generation system, in: AIAA AVIATION Forum, pp. –, American Institute of Aeronautics and Astronautics, https://doi.org/10.2514/6.2016-4053, 2016.
- Baumert, A., Bansmer, S., Trontin, P., and Villedieu, P.: Experimental and numerical investigations on aircraft icing at mixed phase conditions, International Journal of Heat and Mass Transfer, 123, 957–978, 2018.
 - Beaugendre, H., Morency, F., and Habashi, W. G.: FENSAP-ICE's Three-Dimensional In-Flight Ice Accretion Module: ICE3D, Journal of Aircraft, 40, 239–247, https://doi.org/10.2514/2.3113, 2003.
 - Bond, T. H. and Anderson, D. N.: Manual of Scaling Methods, NASA/CR-2004-212875, 2004.
- Braga, R. C., Rosenfeld, D., Weigel, R., Jurkat, T., Andreae, M. O., Wendisch, M., Pöhlker, M. L., Klimach, T., Pöschl, U., Pöhlker, C., et al.:

 Comparing parameterized versus measured microphysical properties of tropical convective cloud bases during the ACRIDICON–CHUVA campaign, Atmospheric Chemistry and Physics, 17, 7365, 2017.
 - Brown, P. and Francis, P.: Improved Measurements of the Ice Water Content in Cirrus Using a Total-Water Probe, Journal of Atmospheric and Oceanic Technology, 12, 410–414, 1995.
 - Connolly, P., Emersic, C., and Field, P.: A Laboratory Investigation into the Aggregation Efficiency of small Ice Crystals, Journal of Atmospheric Chemistry and Physics, 11, 2055 2076, 2011.
 - Currie, T., Struk, P., Tsao, J.-C., Fuleki, D., and Knezevici, D.: Fundamental Study of Mixed-Phase Icing with Application to Ice Crystal Accretion in Aircraft Jet Engines, in: Fluid Dynamics and Co-located Conferences, pp. –, American Institute of Aeronautics and Astronautics, https://doi.org/10.2514/6.2012-3035, 2012.
 - Currie, T. C., Fuleki, D., and Mahallati, A.: Experimental studies of mixed-phase sticking efficiency for ice crystal accretion in jet engines, in: 6th AIAA Atmospheric and Space Environments Conference, p. 3049, 2014.
 - Dezitter, F., Grandin, A., Brenguier, J., Hervy, F., Schlager, H., Villedieu, P., and Zalamansky, G.: HAIC (High Altitude Ice Crystals), 5th AIAA Atmospheric and Space Environments Conference, AIAA-2013-2674, 2013.
 - EASA: Certification Specifications and Acceptable Means of Compliance for Large Aeroplanes CS-25, Amendment 18, https://www.easa.europa.eu/system/files/dfu/CS-25%20Amendment%2018_0.pdf, 2016.
- Farzaneh, M.: Atmospheric Icing of Power Networks, Springer Netherlands, Dordrecht, https://doi.org/10.1007/978-1-4020-8531-4, 2008.
 - Feulner, M., Liao, S., Rose, B., and Liu, X.: Ice Crystal Ingestion in a Turbofan Engine, SAE 2015 International Conference on Icing of Aircraft, Engines and Structures, 2015.
 - Field, P., Heymsfield, A., and Bansemer, A.: Shattering and particle interarrival times measured by optical array probes in ice clouds, Journal of Atmospheric and Oceanic Technology, 23, 1357–1371, 2006.
- Finstad, K. J., Lozowski, E. P., and Makkonen, L.: On the median volume diameter approximation for droplet collision efficiency, Journal of the atmospheric sciences, 45, 4008–4012, 1988.

- Flegel, A. B. and Oliver, M. J.: Preliminary Results from a Heavily Instrumented Engine Ice Crystal Icing Test in a Ground Based Altitude Test Facility, in: 8th AIAA Atmospheric and Space Environments Conference, p. 3894, 2016.
- Fontaine, E., Schwarzenboeck, A., Delanoe, J., Wobrock, W., Leroy, D., Dupuy, R., Gourbeyre, C., and Protat, A.: Constraining mass-diameter relations from hydrometeor images and cloud radar reflectivities in tropical continental and oceanic convective anvils, Atmospheric Chemistry and Physics, 14, 11 367–11 392, 2014.
- Gayet, J.-F., Shcherbakov, V., Bugliaro, L., Protat, A., Delanoë, J., Pelon, J., and Garnier, A.: Microphysical properties and high ice water content in continental and oceanic mesoscale convective systems and potential implications for commercial aircraft at flight altitude, Atmos. Chem. Phys., 14, 899–912, https://doi.org/doi:10.5194/acp-14-899-2014, 2014.
- Gent, R. W., Dart, N. P., and Cansdale, J. T.: Aircraft icing, Philosophical Transactions of the Royal Society A: Mathematical, Physical and Engineering Sciences, 358, 2873–2911, https://doi.org/10.1098/rsta.2000.0689, 2000.
- GL: Rules and Guidelines Industrial Services Guideline for the Certification of Wind Turbines, 2010.
- Gong, X. and Bansmer, S.: 3-D ice shape measurements using mid-infrared laser scanning, Optics express, 23, 4908-4926, 2015a.
- Gong, X. and Bansmer, S.: Laser scanning applied for ice shape measurements, Cold Regions Science and Technology, 115, 64-76, 2015b.
- Grzych, M. and Mason, J.: Weather Conditions Accordated with Jet Engine Power-loss and Damage due to Ingestion of Ice Particles: What

 We've Learned Through 2009, 14th Conference of Aviation, Range and Aerospace Meteorology, 2010.
 - Hochart, C., Fortin, G., Perron, J., and Ilinca, A.: Wind turbine performance under icing conditions, Wind Energy, 11, 319–333, https://doi.org/10.1002/we.258, 2008.
 - IPHWG: Working Group Report on Supercooled Large Droplet Rulemaking, 2005.

5

10

20

- Jung, S., Tiwari, M. K., Doan, N. V., and Poulikakos, D.: Mechanism of supercooled droplet freezing on surfaces, Nat Commun, 3, 615–, 2012.
- Knudsen, J. G. and Katz, D. L.: Fluid dynamics and heat transfer, McGraw-Hill Book Company, 1958.
- Korolev, A.: Reconstruction of the sizes of spherical particles from their shadow images. Part I: Theoretical considerations, Journal of Atmospheric and Oceanic Technology, 24, 376–389, 2007.
- Korolev, A. and Sussman, B.: A technique for habit classification of cloud particles, Journal of Atmospheric and Oceanic Technology, 17, 1048–1057, 2000.
 - Lamb, D. and Verlinde, J.: Physics and Chemistry of Clouds, Cambridge University Press, cambridge Books Online, 2011.
 - Langmuir, I.: The collected works of Irving Langmuir, vol. 10 Atmospheric phenomena, chap. Mathematical investigation of water droplet trajectories, pp. 335–393, Pergamon press, 1961.
 - Lawson, R. and Baker, B.: Improvement in determination of ice water content from two-dimensional particle imagery. Part II: Applications to collected data, Journal of Applied Meteorology and Climatology, 45, 1292–1303, 2006.
 - Lawson, R., O'Connor, D., Zmarzly, P., Weaver, K., Baker, B., Mo, Q., and Jonsson, H.: The 2D-S (Stereo) probe: design and preliminary tests of a new airborne, high speed, high-resolution particle imaging probe, Journal of Atmospheric and Oceanic Technology, 23, 1462–1477, 2006.
- Leroy, D., Fontaine, E., Schwarzenboeck, A., Strapp, J., Lilie, L., Delanoe, J., Protat, A., Dezitter, F., and Grandin, A.: HAIC/HIWC Field
 Campaign Specific Findings on PSD Microphysics in High IWC Regions from In Situ Measurements: Median Mass Diameters, Particle
 Size Distribution Characteristics and Ice Crystal Shapes, SAE Technical Paper 2015-01-2087, SAE 2015 International Conference on
 Icing of Aircraft, Engines and Structures, 2015.

- Leroy, D., Coutris, P., Fontaine, E., Schwarzenboeck, A., Strapp, J., Lilie, L., Korolev, A., McFarqhuar, G., Dezitter, F., and Calmes, A.: HAIC/HIWC field campaigns specific findings on ice crystal characteristics in high ice water content cloud regions, 8th AIAA Atmospheric and Space Environments Conference, 2016a.
- Leroy, D., Fonaine, E., Schwarzenboeck, A., and Strapp, J.: Ice Crystal Sizes in High Ice Water Content Clouds. Part I: On the Computation of Median Mass Diameter from In Situ Measurements. Part I: On the Computation of Median Mass Diameter from In Situ Measurements, Journal of Atmospheric and Oceanic Technology, 33, 2461–2476, 2016b.
 - Leroy, D., Fontaine, E., Schwarzenboeck, A., Strapp, J., Korolev, A., McFarquhar, G., Dupuy, R., Gourbeyre, C., Lilie, L., Protat, A., Delanoue, J., Dezitter, F., and Grandin, a.: Ice Crystal Sizes in High Ice Water Content Clouds. Part II: Statistics of Mass Diameter Percentiles in Tropical Convection Observed during the HAIC/HIWC Project, Journal of Atmospheric and Oceanic Technology, 34, 117–136, 2017.
 - Makkonen, L.: Estimation of wet snow accretion on structures, Cold Regions Science and Technology, 17, 83 88, https://doi.org/http://dx.doi.org/10.1016/S0165-232X(89)80018-7, 1989.

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- Mason, J., Strapp, J., and Chow, P.: The Ice Particle Threat to Engines in Flight, 44th AIAA Aerospace Sciences Meeting and Exhibit, AIAA-2006-206, 2006.
- 15 Messinger, B. L.: Equilibrium temperature of an unheated icing surface as a function of air speed, Journal of the Aeronautical Sciences, 20, 29–42, 1953.
 - Mulherin, N.: Atmospheric icing and communication tower failure in the United States, Cold Regions Science and Technology, 27, 91 104, https://doi.org/http://dx.doi.org/10.1016/S0165-232X(97)00025-6, 1998.
 - Oleskiw, M. M., Hyde, F. H., and Penna, P. J.: In-Flight Icing Simulation Capabilities of NRC's Altitude Icing Wind Tunnel, in: 39th AIAA Aerospace Sciences Meeting & Exhibit, AIAA-2001-0094, https://doi.org/10.2514/6.2001-94, 2001.
 - Oliver, M.: Validation Ice Crystal Icing Engine Test in the Propulsion System Laboratory at NASA Glenn Research Center, 6th AIAA Atmospheric and Space Environments Conference, AIAA-2014-2898, 2014.
 - Pastor-Barsi, C., Arrington, A., and VanZante, J. F.: Aero-Thermal Calibration of the NASA Glenn Icing Research Tunnel (2012 Test), NASA Technical Report E-18454-1, 2012.
- 25 Politovich, M. K.: Aircraft Icing Caused by Large Supercooled Droplets, J. Appl. Meteor., 28, 856–868, https://doi.org/10.1175/1520-0450(1989)028<0856:AICBLS>2.0.CO;2, 1989.
 - Ramsey, J. D., Burford, C. L., Beshir, M. Y., and Jensen, R. C.: Effects of workplace thermal conditions on safe work behavior, Journal of safety Research, 14, 105–114, 1983.
 - Reehorst, A. L. and Richter, G. P.: New Methods and Materials for Molding and Casting Ice Formations, Tech. Rep. NASA-TM-100126, NASA Lewis Research Center, 1987.
 - Rudoff, R., Smith, J., and Bachalo, W.: Development of a Phase Doppler Based Probe for Icing Cloud Droplet Characterization, 28th Aerospace Sciences Meeting, Reno Nevada,, 1990.
 - Scheiman, J. and Brooks, J.: Comparison of Experimental and Theoretical Turbulence Reduction from Screens, Honeycomb, and Honeycomb-Screen Combinations, Journal of Aircraft, 18, 638–643, 1981.
- Schremb, M. and Tropea, C.: Solidification of supercooled water in the vicinity of a solid wall, Physical Review E, 94, 052 804, 2016. Simpson, R. L.: Junction flows, Annual Review of Fluid Mechanics, 33, 415–443, 2001.
 - Steiner, J. and Bansmer, S.: Ice Roughness and its Impact on the Ice Accretion Process, in: 8th AIAA Atmospheric and Space Environments Conference, p. 3591, 2016.

- Strapp, J., Korolev, A., Ratvasky, T., Potts, R., Protat, A., May, P., Ackerman, A., Fridlind, A., Minnis, P., Haggerty, J., Riley, J., Lilie, L., and Isaac, G.: The High Ice Water Content Study of Deep Convective Clouds: Report on Science and Technical Plan, techreport, Federal Aviation Administration, https://doi.org/DOT/FAA/TC-14/31, 2016.
- Treaster, A. L. and Yocum, A. M.: The calibration and application of five-hole probes, Tech. Rep. TM 78-10, Pennsylvania State University, institute for science and engineering, applied research laboratory, 1978.
- Uberoi, M. S.: Effect of Wind-Tunnel Contraction on Free-Stream Turbulence, Journal of the Aeronautical Sciences, 23, 754-764, 1956.
- Vecchione, L. and De Matteis, P.: An overview of the CIRA icing wind tunnel, in: 41st Aerospace Sciences Meeting and Exhibit, p. 900, 2003.
- Voigt, C., Schumann, U., Jurkat, T., Schäuble, D., Schlager, H., Petzold, A., Gayet, J.-F., Krämer, M., Schneider, J., Borrmann, S., et al.: Insitu observations of young contrails—overview and selected results from the CONCERT campaign, Atmospheric Chemistry and Physics, 10, 9039–9056, 2010.
 - Voigt, C., Schumann, U., Jessberger, P., Jurkat, T., Petzold, A., Gayet, J.-F., Krämer, M., Thornberry, T., and Fahey, D.: Extinction and optical depth of contrails, Geophysical Research Letters, 38, 2011.
- Voigt, C., Schumann, U., Minikin, A., Abdelmonem, A., Afchine, A., Borrmann, S., Boettcher, M., Buchholz, B., Bugliaro, L., Costa,
 A., et al.: ML-CIRRUS: The Airborne Experiment on Natural Cirrus and Contrail Cirrus with the High-Altitude Long-Range Research
 Aircraft HALO, Bulletin of the American Meteorological Society, 98, 271–288, 2017.
 - Wagner, B.: Ice Accretion Simulation, Tech. Rep. AGARD report no. 344, North Atlantic Treaty Organization, 1997.
 - Wendisch, M., Pöschl, U., Andreae, M. O., Machado, L. A., Albrecht, R., Schlager, H., Rosenfeld, D., Martin, S. T., Abdelmonem, A., Afchine, A., et al.: ACRIDICON–CHUVA Campaign: Studying Tropical Deep Convective Clouds and Precipitation over Amazonia Using the New German Research Aircraft HALO, Bulletin of the American Meteorological Society, 97, 1885–1908, 2016.
 - Yao, S.-C. and Schrock, V.: Heat and mass transfer from freely falling drops, Journal of Heat Transfer, 98, 120-126, 1976.
 - Zerr, R. J.: Freezing Rain: An Observational and Theoretical Study, J. Appl. Meteor., 36, 1647–1661, https://doi.org/10.1175/1520-0450(1997)036<1647:FRAOAT>2.0.CO;2, 1997.

Nomenclature

25 Roman symbols

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- \underline{a} speed of sound, ms^{-1}
- A area, m²
- A_c accumulation parameter
- ΔH absolute humidity, g m⁻³
- c chord length of an airfoil, m
- c_D drag coefficient of a droplet
- c_{latent} specific latent heat of fusion, kJkg⁻¹
- c_{p} specific heat capacity at constant pressure, kJ kg⁻¹ K⁻¹
- d diameter of a droplet, μ m
- D drag force vector, $1 \cdot N$

 f_0 stagnation line freezing fraction

 f_{dosing} dosing frequency, Hz

g gravitational acceleration vector, $\mathbf{1} \cdot \text{m s}^{-2}$

IWC ice water content, gm^{-3}

k heat conductivity, $W m^{-1} K^{-1}$

K contraction of the wind-tunnel nozzle

L length of the wind-tunnel nozzle, m

LWC liquid water content, g m⁻³

m = mass, kg

 $\underline{\dot{m}}$ mass flow rate, kg s⁻¹

 $\underbrace{m_{\mathcal{T}}}$ melting ratio

Mach number

MVD median volume diameter, µm

MMD median mass diameter, µm

Nu Nusselt number

p pressure, Pa

 \underbrace{P} power, W

Pr Prandtl number

 \dot{Q} heat, \dot{W}

 \underbrace{R} specific gas constant, $\operatorname{Jkg^{-1}K^{-1}}$

Re Reynolds number

RH relative humidity

 \mathcal{S}_{spray} operational state of the spray system

t time, s

 t_{acc} accumulation time, s

T static temperature, K

Tu Turbulence level

TWC total water content, $g m^{-3}$

u, w velocity components in x and z direction, ms^{-1}

 U_{∞} freestream velocity of air, ms⁻¹

 \dot{V} volumetric flow rate, $m^3 s^{-1}$

x, z Cartesian coordinates, m

Greek symbols

 α Angle of attack, deg

 β , γ fitting constants for particle sizing

 δ twice the nose radius of an airfoil, m

 Δ difference, e.g. in temperature or in velocity

 ζ specific loss coefficient

 $\underline{\eta}$ efficiency

 λ power factor

 μ dynamic viscosity, $kg m^{-1} s^{-1}$

 ϱ density, kg m⁻³

Subscripts

air air

chill chilling

d droplet

eq equivalent

fan fan

HE heat exchanger

HSI High-speed imaging probe

ice ice

inj injected

IKP Isokinetic probe

jet jet

latent latent

max maximum

nozzle wind tunnel nozzle

PDI Phase Doppler interferometry

sc settling chamber

sensible sensible

target target value

tsec test section

tot total

vapour vapour

water water

wb wet bulb

 \underbrace{x} $\underbrace{\text{in } x \text{-direction}}$

 ∞ freestream condition

Table 1. Performance of different icing wind tunnels compared to the Braunschweig Icing Wind Tunnel. $v_{max} U_{\infty,max}$ represents the maximum air speed in the test section, P_{chill} is the installed cooling power. Data extracted from Pastor-Barsi et al. (2012), Vecchione and De Matteis (2003), Al-Khalil et al. (1998) and Oleskiw et al. (2001).

Tunnel Name	Test Section Size	$v_{ ext{max}}$ $U_{\infty, ext{max}_{\infty}}$	P_{fan}	P_{chill}	λ	$\frac{P_{\mathrm{chill}}}{P_{\mathrm{fan}}}$	Altitude
	$\mathbf{m}\times\mathbf{m}$	$\frac{\text{m/s}}{\infty} \frac{\text{ms}^{-1}}{\infty}$	kW	kW	-	-	capability
NASA IRT	1.83×2.75	174.3	3730		0.20	-	no
CIRA	2.35×1.15	225.0	4000	6400	0.19	1.60	yes
Cox	0.71×1.17	98.3	149	-	0.27	-	no
NRC AIWT	0.57×0.57	100.0	450	420	1.99	0.93	yes
	0.52×0.33	165.0					
Braunschweig	0.50×0.50	40.0	37	80	3.35	2.16	no

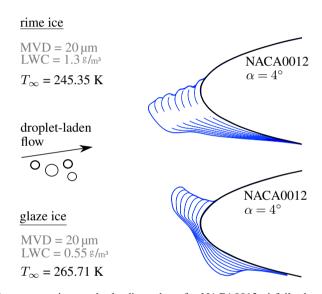
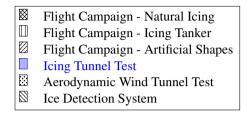


Figure 1. Rime ice and glaze ice shapes over time at the leading edge of a NACA0012 airfoil, chord length 0.53 m. One blue line represents one minute of ice accretion. Further boundary conditions (rime): $\frac{U_{\infty} = 58 \text{ m/s}U_{\infty} = 58 \text{ ms}^{-1}}{U_{\infty} = 58 \text{ ms}^{-1}}$, $\frac{t_{\text{acc}} = 8 \text{ min}}{t_{\text{acc}} = 8 \text{ min}}$; (glaze): $\frac{U_{\infty} = 102.8 \text{ m/s}U_{\infty} = 102.8 \text{ m/s}U_{\infty} = 102.8 \text{ m/s}^{-1}$, $\frac{t_{\text{acc}} = 7 \text{ min}}{t_{\text{acc}} = 7 \text{ min}}$. Data from computations with TAUICE by Jan Steiner.

Microscopic images of ice particles inside the cloud chambers. *Left:* plate crystals. *Right:* Needles. Ice particle structure after storage, dosing and sieving procedure Ice particle conveyance pipe inside the settling chamber of the Braunschweig Icing Wind Tunnel Cumulative particle size distribution: Comparison of artificial ice particle cloud to atmospheric cloud conditions Even for appropriate clothing, exposure to cold temperatures yields an increase in unsafe behavior. Data reproduced from Ramsey et al. (1983). *Legionella pneumophila*, strain Corby. Transmission electron microscopy, negative contrast, scale marker = 500 nm. Picture by Norbert Bannert, Gudrun Holland / Robert Koch Institute.



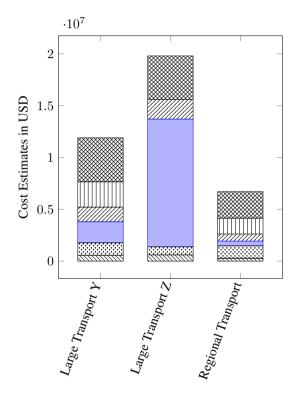


Figure 2. Cost estimation for aircraft icing certification according to FAR Part 25 Appendix O based on the working group report on supercooled large droplet rulemaking IPHWG (2005).

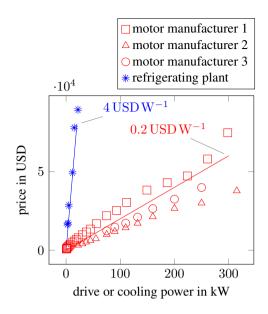


Figure 3. Pricing as of 2017 for asynchronous motors from different manufactures and for CO₂ refrigerating plants before tax and tolls.

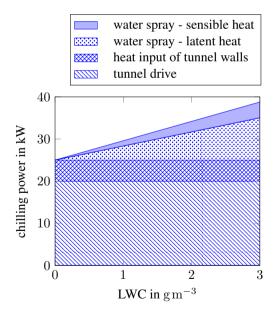


Figure 4. Required chilling power decomposition as a function of the liquid water content for typical operational conditions of the Braunschweig Icing Wind Tunnel.

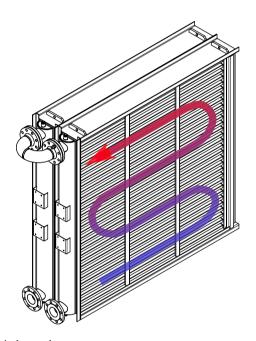


Figure 5. Heat exchanger inside the icing wind tunnel.

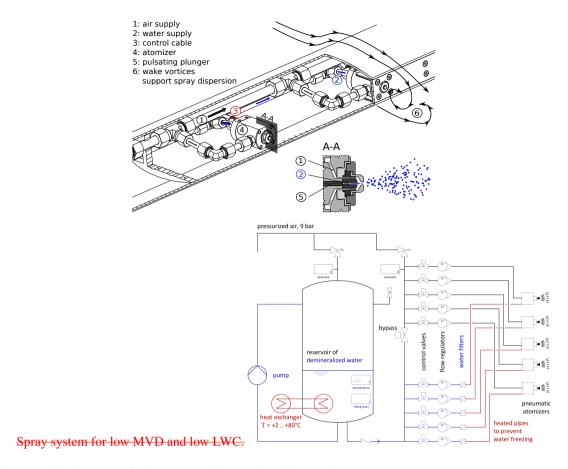


Figure 6. Spray conditioning system for low MVD and low LWC, and its supply system.

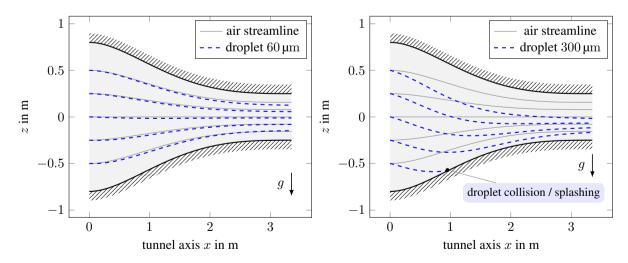


Figure 7. Droplet trajectories in the wind tunnel nozzle and their deviation from the air streamlines under the influence of gravity. Large droplets may collide with the tunnel walls thereby altering the droplet size distribution.

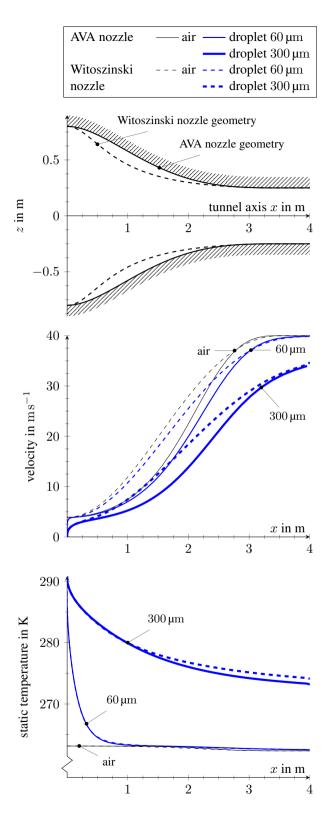


Figure 8. Velocity and temperature slope for air and two different sized droplets inside two tunnel nozzle configurations.

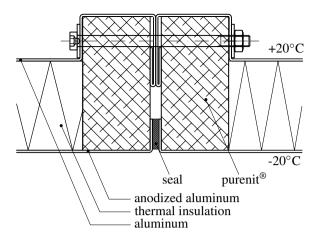


Figure 9. Mechanical design of the icing wind tunnel walls.

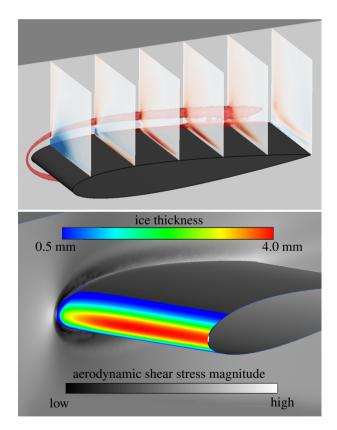


Figure 10. The junction flow between a NACA0012 airfoil (e = 75 emc = 0.75 m, α = 0°) and wind-tunnel side wall creates a horse-shoe vortex (upper image), which is altering the ice accretion (lower image). $\frac{U_{\infty} = 40 \frac{m}{s} U_{\infty} = 40 \text{ ms}^{-1}}{s}$, $\frac{T_{\infty} = -10 \degree \text{C}T_{\infty} = -10 \degree \text{C}}{s}$, LWC=1 $\frac{g}{m^3}$ 1 g m $\frac{-3}{s}$, IWC=0.

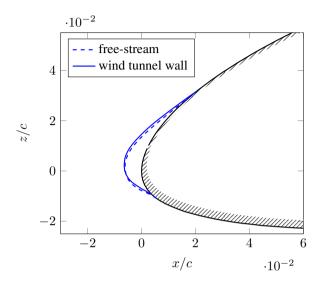


Figure 11. Inviscid wall effect on ice accretion for a AH-94-W-145 airfoil, $\frac{\text{c=0.75 mc}}{\text{c=0.75 m}}$, $U_{\infty} = \frac{40 \text{ m/s}}{40 \text{ ms}} = \frac{10 \text{ °C}}{10 \text{ °C}}$, $U_{\infty} = \frac{10 \text{ °C}$

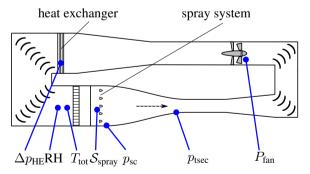


Figure 12. Basic instrumentation of the Braunschweig Icing Wind Tunnel.

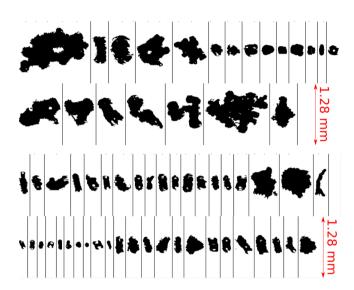


Figure 13. Ice particle images captured by 2DS-2D-S probe in stratiform cloud region during the Darwin Campaign. Top: Stratiform cloud region , $\frac{MMD \approx 150 \mu m}{MMD} \approx 150 \mu m$, $\frac{IWC \approx 1.2 g/m^3}{IWC} \approx 1.2 g m^{-3}$, $\frac{T_{\infty} \approx -10 \, ^{\circ}C}{I} \approx -10 \, ^{\circ}C$. Bottom: Convective cloud region, $\frac{MMD \approx 80 \, \mu m}{I} = 100 \, ^{\circ}C$. Bottom: Convective cloud region, $\frac{MMD \approx 80 \, \mu m}{I} = 100 \, ^{\circ}C$.

Ice particle images captured by 2DS probe in convective cloud region during the Darwin Campaign, $MMD \approx 80 \mu m$, $IWC \approx 3.0 q/m^3$, $T_{\infty} \approx -10 \, ^{\circ}C$

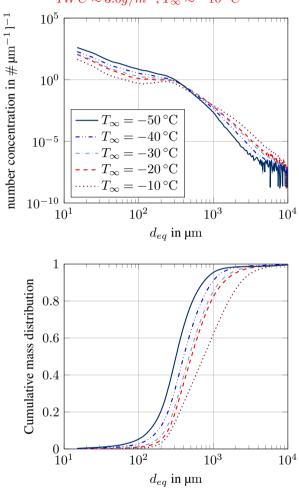
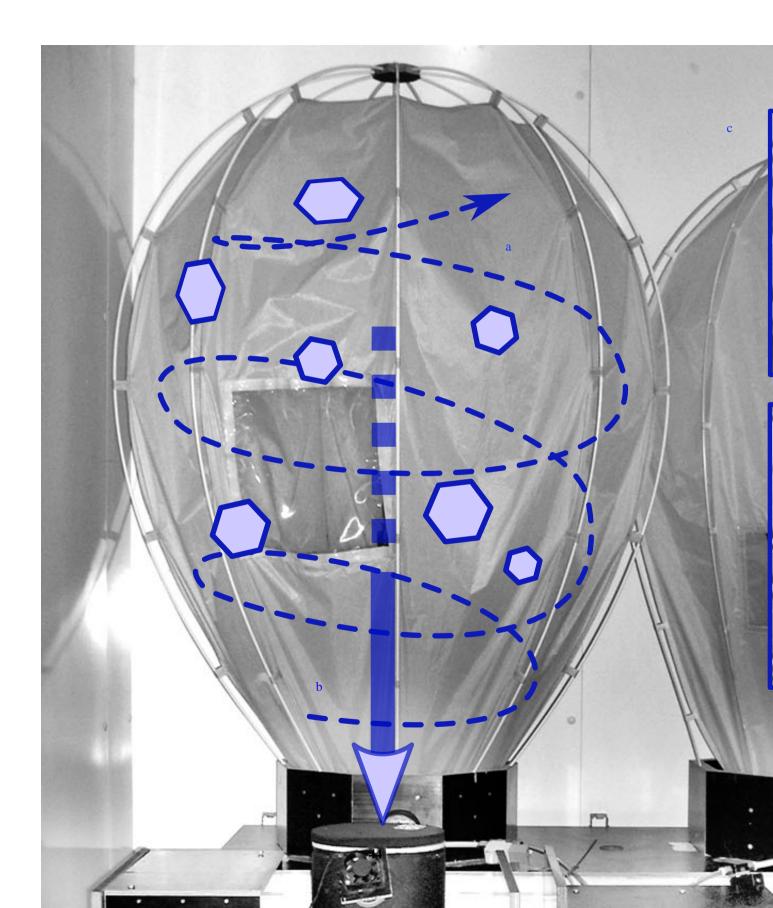
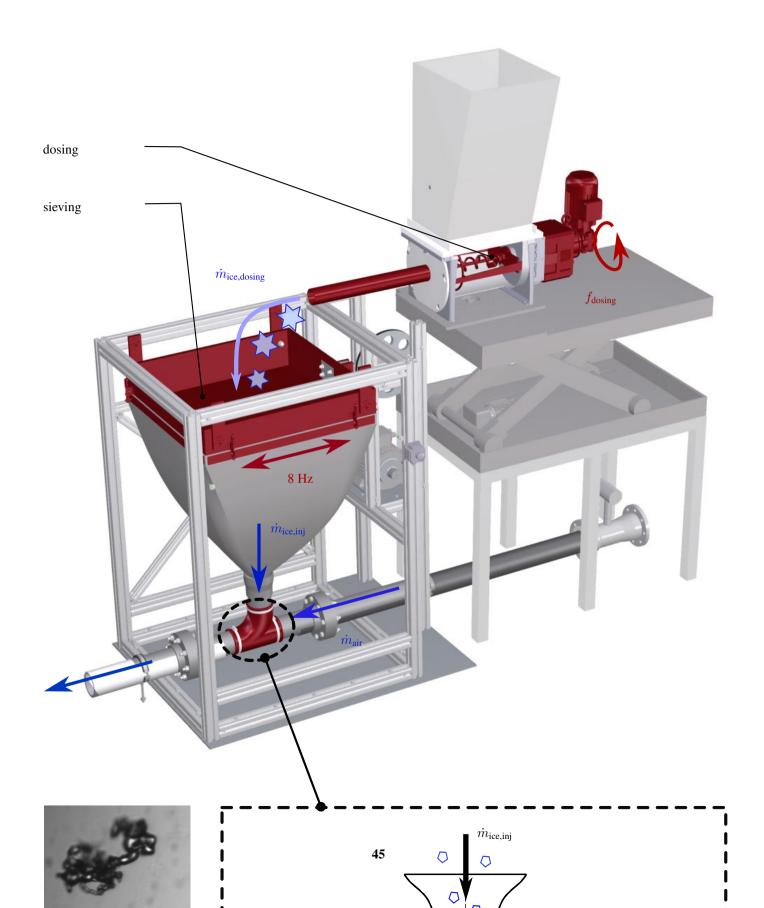
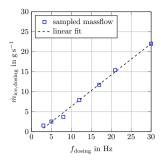


Figure 14. Ice particle size distribution (MMD) based on HAIC Falcon 20 flight measurements from Darwin and Cayenne, dependence on ambient air temperature







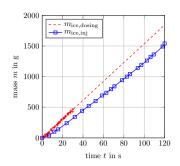


Figure 17. Ice particle massflow. *Left:* Dependence on dosing machine frequency. *Right:* Ice particle supply over time at a dosing frequency of $f_{\text{dosing}} = 21 \text{ Hz}$.

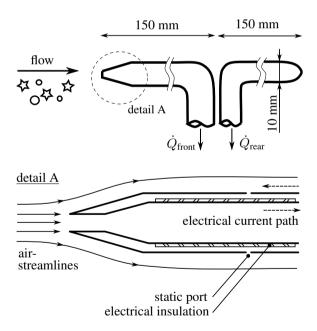


Figure 18. Isokinetic probe to measure the total water content TWC, developed at Cranfield university.

Ice particle detection by means of C-HSI probe. Left: Trigger beam induced by HSI. Right:

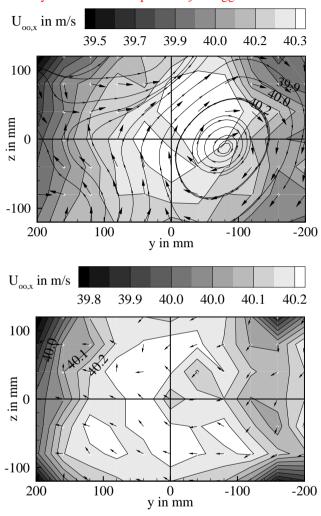


Figure 19. Spatial distribution of axial flow velocity $U_{\infty,x}$, for the . Top: ice-crystal-icing setup. <u>Bottom:</u> standard setup of the icing wind tunnel.

Spatial distribution of axial flow velocity $U_{\infty,x}$, for the standard setup.

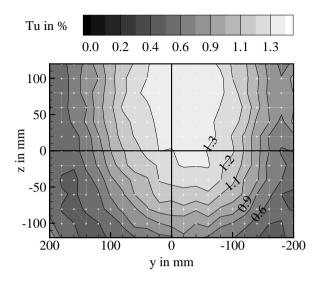


Figure 20. Spatial distribution of turbulence intensity for the mixed phase configuration

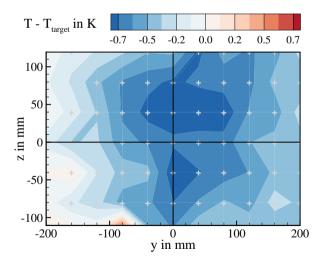
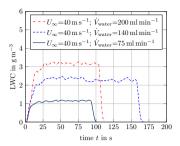


Figure 21. Spatial distribution of temperature variation for the mixed phase configuration



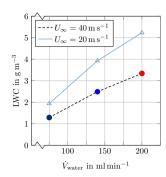
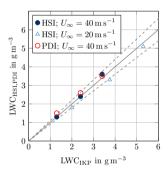


Figure 22. Liquid water content calibration at the tunnel centreline based on IKP measurements. *Left:* LWC history. *Right:* LWC depending on supplied water flow-rate \dot{V}_{water} .



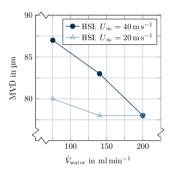
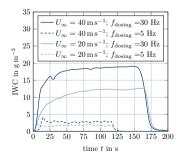


Figure 23. Comparison of HSI, PDI and IKP measurements for supercooled droplet conditions. *Left:* Optical array probes vs. IKP. *Right:* MVD vs. LWC.



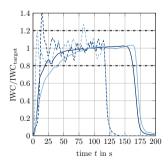
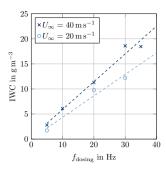


Figure 24. IWC time history based on IKP measurements. Left: IWC history. Right: Normalized IWC history.



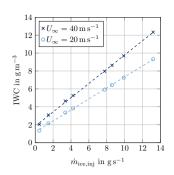


Figure 25. Ice water content depending on dosing frequency f_{dosing} and injected ice particle flow $\dot{m}_{ice,inj}$

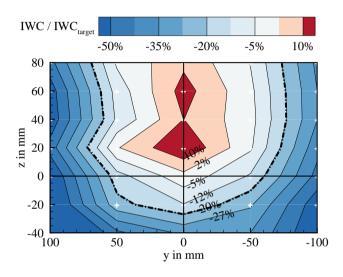


Figure 26. Percentaged deviation of IWC from centerline value, $U_{\infty} = \frac{40 \text{ m/s}}{40 \text{ m/s}} = \frac{20 \text{ Hz}}{100 \text{ m/s}}$.

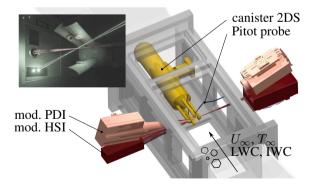


Figure 27. Setup of the probe inter-comparison exercise. Alternatively to the canister 2DS-2D-S probe, which is shown here, the canister HSI probe was mounted. *Upper left:* C-HSI probe mounted inside the 1WTicing wind tunnel, modular PDI and HSI laser beams are adjusted close to the C-HSI sampling volume

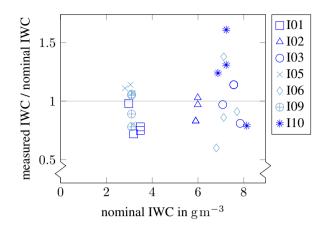


Figure 28. Results of the probe intercomparison between 2DS-LAMP2D-S (LAMP), 2DS-DLR2D-S (DLR), M-HSI, C-HSI for seven different test points at glaciated icing conditions

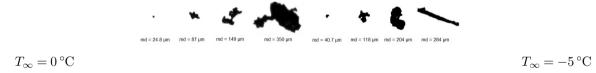


Figure 29. Ice particle images captured by the C-HSI probe inside the Braunschweig Icing Wind Tunnel test section

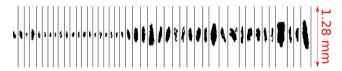


Figure 30. Ice particle images captured by the 2D-S probe inside the Braunschweig Icing Wind Tunnel test section, $\frac{MMD = 79 \mu m}{MMD} = 79 \mu m$, $\frac{TWC = 3.2 g \, \text{m}^{-3}}{TWC} = \frac{15 \, \text{°C}}{TW} = -15 \, \text{°C}$

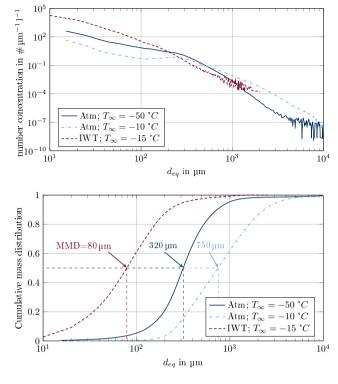


Figure 31. Particle size distribution: Comparison of artificial ice particle cloud to atmospheric cloud conditions. *Top:* Particle size distribution. *Bottom:* Cumulative particle size distribution.

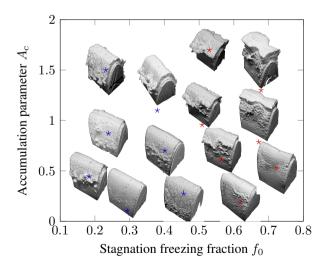
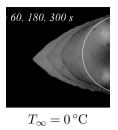
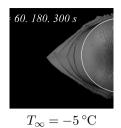
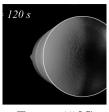


Figure 32. Experimentally observed SLD ice shapes on a NACA0012 leading edge in the parameter space of f_0 and A_c , Re= $2 \cdot 10^6$, MVD= 80μ m80 μ m, AoA= 0° . Results reproduced from Steiner and Bansmer (2016).







 $T_{\infty} = -15\,^{\circ}\mathrm{C}$

Figure 33. Side view of ice accretion shapes on a cylinder for $\overline{TWC} = 12g/m^3$ $\overline{TWC} = 12 g m^{-3}$, $m_r = 0.12$ and $\overline{U_{\infty}} = 40 \text{ m/s} U_{\infty} = 40 \text{ m/s}^{-1}$, with a variation of temperature

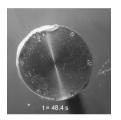






Figure 34. Side view of ice accretion when heating the cylinder, $T_{\infty} = -5$ °C, $T_{\infty} =$

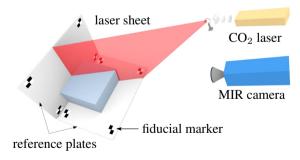


Figure 35. Infrared laser line scanning technique. The mid-infrared camera records the distorted CO_2 -laser line on the ice surface. With a prior calibration, the 3D geometry of the object can be reconstructed. Image reproduced from Gong and Bansmer (2015a).

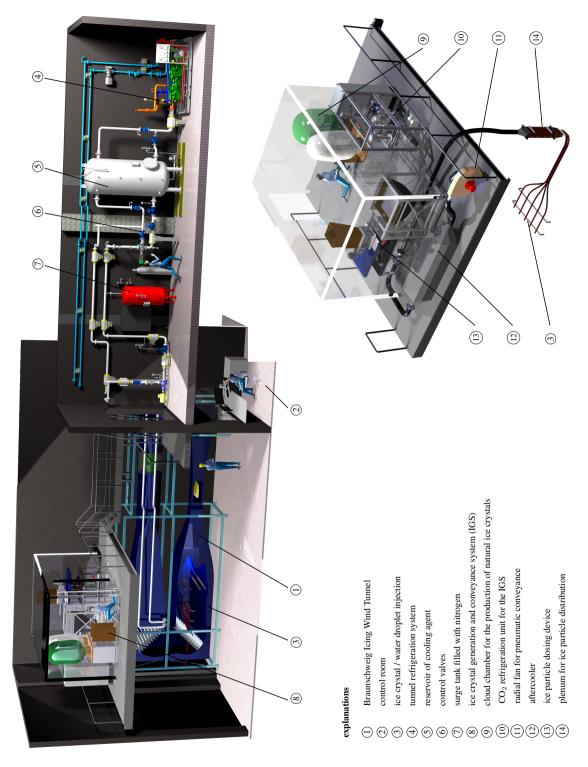


Figure 36. Overview of the Braunschweig Icing Wind Tunnel and its adjacent components.