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- 2 Raindrop Fall Velocities from an Optical Array Probe and 2D-Video
- 3 Disdrometer
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12 Abstract

We report on fall speed measurements of rain drops in light-to-heavy rain events from 13 two climatically different regimes (Greeley, Colorado, and Huntsville, Alabama) using 14 15 the high resolution (50 microns) Meteorological Particle Spectrometer (MPS) and a 3rd generation (170 microns resolution) 2D-video disdrometer (2DVD). To mitigate wind-16 effects, especially for the small drops, both instruments were installed within a 2/3-scale 17 18 Double Fence Intercomparison Reference (DFIR) enclosure. Two cases involved lightto-moderate wind speeds/gusts while the third case was a tornadic supercell that 19 passed over the site with high wind speeds/gusts. As a proxy for turbulent intensity, 20 21 maximum wind speeds from 10-m height at the instrumented site recorded every 3 s were differenced with the 5-min average wind speeds and then squared. The fall speed 22 versus size from 0.1-2 mm were derived from the MPS data and the 2DVD was used for 23 sizes >0.7 mm. Consistency of fall speeds from the two instruments in the overlap 24 region (0.7-2 mm) gave confidence in the data guality and processing methodologies. 25 26 Our results indicate that under light-to-moderate wind gusts, the mean fall speeds agree well with fits to the terminal velocity measured in the laboratory by Gunn and Kinzer 27 from 100 microns up to precipitation sizes. In the supercell case the very strong gusts 28 29 and inferred high turbulence intensity caused a significant broadening of the fall speed 30 distributions with the mean fall speeds about 25-30% less than the terminal velocity of Gunn-Kinzer, i.e. sub-terminal fall speeds. 31

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34 1 Introduction

Knowledge of the terminal fall speed of raindrops as function of size is important in 35 modelling collisional break-up and coalescence processes (e.g., List et al., 1987), in the 36 37 radar-based estimation of rain rate, in retrieval of drop size distribution using Doppler spectra at vertical incidence (e.g., Sekhon and Srivastava, 1971) and in soil erosion 38 studies (e.g., Rosewell 1986). The terminal velocity measurements of Gunn and Kinzer, 39 40 1949) under calm laboratory conditions, and fits to their data (e.g., Atlas et al., 1973; Foote and du Toit, 1969; Beard and Pruppacher, 1969) are still considered the standard 41 against which measurements using more modern optical instruments in natural rain are 42 43 compared (Löffler-Mang and Joss, 2000; Barthazy et al., 2004; Schönhuber et al., 2008; 44 Testik and Rahman, 2016). More recently, the broadening and skewness of the fall speed distributions of a given size (3 mm) in one intense rain event was attributed to 45 mixed-mode amplitude oscillations (Thurai et al., 2013). Super- and sub-terminal fall 46 speeds in intense rain shafts have been detected and attributed, respectively, to drop 47 48 breakup fragments (sizes < 0.5 mm), and high wind/gusts (sizes 1-2 mm) (Montero-Martinez et al., 2009; Larsen et al., 2014; Montero-Martinez and Garcia-Garcia, 2016). 49

The fall speeds and concentration of small drops (< 1 mm) in natural rain are difficult to 50 measure given the poor resolution (>170 microns) of most optical disdrometers and/or 51 sensitivity issues. While cloud imaging probes (with high resolution 25-50 microns) on 52 aircraft have been used for many years they generally cannot measure the fall speeds. 53 A relatively new instrument, the Meteorological Particle Spectrometer (MPS) is a droplet 54 imaging probe that was built by Droplet Measurements Technologies (DMT, Inc.) under 55 contract from the US Weather Service specifically designed for drizzle as small as 50 56 µm and rain drops up to 3 mm. This instrument in conjunction with a lower resolution 57 2D-Video Disdrometer (Schoenhuber et al., 2008) is used in this paper to measure fall 58 speed distributions in natural rain. 59

This paper briefly describes the instruments used, presents fall speed measurements from two sites under relatively low wind conditions, and one case from an unusual tornadic supercell with high winds and gusts and ends with a brief discussion and summary of the results.

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65 2 Instrumentation and Measurements

The principal instruments used in this study are the MPS and 3rd generation 2D-video disdrometer (2DVD), both located within a 2/3-scale Double Fence Intercomparison Reference (DFIR; *Rasmussen et al.*, 2012) wind shield. As reported in (*Notaros et al.*, 2016), the 2/3-scale DFIR was effective in reducing the ambient wind speeds by nearly





a factor of 2-3 based on data from outside and inside the fence. The flow field in and around the DFIR has been simulated by (*Theriault et al.*, 2015) assuming steady ambient winds. They found that depending on the wind direction relative to the octagonal fence, weak up/down drafts could be generated above the sensor areas. For 5 m/s speeds, the up/down drafts could range between -0.4 (down) to 0.2 m/s (up).

The instrument set-up was the same for the two sites (Greeley, Colorado and Huntsville, Alabama). Huntsville has a very different climate from Greeley, and its altitude is 212 m MSL as compared with 1.4 km MSL for Greeley. According to the Köppen–Trewartha climate classification system (*Trewartha and Horn,* 1980), this labels Greeley as a semiarid-type climate, whereas Huntsville is a humid subtropical-type climate (*Belda et al.*, 2014).

The MPS is an optical array probe (OAP) that uses the technique introduced by 81 Knollenberg (1970, 1976, 1980) and measures drop diameter in the range from 0.05-3.1 82 mm. A 64 element photo-diode array is illuminated with a 660 nm collimated laser 83 beam. Droplets passing through the laser cast a shadow on the array and the decrease 84 in light intensity on the diodes is monitored with the signal processing electronics. A two 85 dimensional image is captured by recording the light level of each diode during the 86 period that the array is shadowed. The fall velocity is derived using two methods. One 87 uses the same approach as described in (Montero-Martinez et al., 2009) where the fall 88 velocity is calculated from the product of the true air speed clock and ratio of the image 89 height -to-width. Note that "width" is the horizontal dimension parallel to the array and 90 "height" is along the vertical. The second method computes the fall velocity from the 91 maximum horizontal dimension (spherical drop shape assumption) divided by the 92 amount of time that the image is on the array, a time measured with a 2 MHz clock. In 93 94 order to be comparable to the results of (Montero-Martinez et al., 2009), their approach is implemented here for sizes > 250 µm. The fall velocity of smaller, slower moving 95 droplets, are measured using the second technique. 96

The limitations and uncertainties associated with OAP measurements have been well documented (*Korolev et al.*, 1991; 1998; *Baumgardner et al.*, 2017). All possible corrections have been applied, including the removal of artifacts due to splashing, and oversizing that results from out-of-focus droplets (*Korolev* 2007). The sizing and fall speed errors primarily depend on the digitization error (± 25 microns). The fall speed accuracy according to the manufacturer (DMT) is <10% for 0.25 mm and <1% for sizes greater than 1 mm, limited primarily by the accuracy in droplet sizing.

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105 The 3rd generation 2DVD is described in detail by (*Schoenhuber et al.,* 2007; 2008) and 106 its accuracy of size and fall speed measurement has been well documented (e.g.,





107 Thurai et al., 2007; 2009; Huang et al., 2008; Bernauer et al., 2015). Considering the horizontal pixel resolution of 170 microns and other factors, the effective sizing range is 108 D> 0.7 mm. The fall velocity accuracy is determined primarily by the accuracy of 109 calibrating the distance between the two orthogonal light "sheets" or planes and is < 5% 110 for fall velocity <10 m s⁻¹. In our application, we utilize the MPS for measurement of 111 small drops with D < 1.2 mm and to compare the measurements with the 2DVD in the 112 overlap region of $D \approx 0.7-2.0$ mm to ensure consistency of observations. The only fall 113 velocity threshold used for the 2DVD is the lower limit set at 0.5 m s⁻¹ in accordance 114 with the manufacturer guidelines for rain measurements. 115

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117 2.1 Fall Speeds from Greeley, Colorado

118 We first consider a long duration (around 20 h) rain episode on 17 April 2015 which 119 consisted of a wide variety of rain types/rates (mostly light stratiform < 8 mm h⁻¹) as 120 described in Table 2 of (*Thurai et al.*, 2017). Two wind sensors at height of 1 m were 121 available to measure the winds outside and inside the DFIR. Average wind speeds 122 were, respectively, < 1.5 m s⁻¹ inside the DFIR and < 4 m s⁻¹ outside with light gusts. 123 These wind sensors were specific to the winter experiment described in (*Notaros et al.*, 124 2016) and were unavailable for the rain measurement campaign after May 2015.

Figure 1(a) shows the fall speeds versus D from the 2DVD (shown as contoured 125 frequency of occurrence), along with mean and $\pm 1\sigma$ standard deviation from the MPS. 126 Also shown is the (Atlas et al., 1973) fit to the terminal fall speed measurements of 127 (Gunn and Kinzer, 1949) at sea level and after applying altitude corrections (Beard. 128 1976) for the elevation of 1.4 km MSL for Greeley. Panel (b) shows the histogram of fall 129 speeds for two selected diameter intervals (0.5±0.1 mm) and (1±0.1 mm). Panel (a) 130 demonstrates the excellent "visual" agreement between the two instruments in the 131 132 overlap size range (0.7-2 mm) as well as with the fit to the Gunn-Kinzer laboratory data. Notable is the remarkable agreement in mean fall speeds between the Gunn-Kinzer fit 133 and the MPS for D< 0.5 mm down to near the lower limit of the instrument (0.1 mm). 134 Few measurements have been reported of fall speeds in this size range. 135

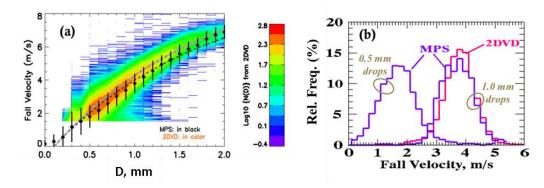
The histograms in Fig. 1(b) show good agreement between 2DVD and MPS for 1 mm 136 drop sizes. The visual agreement between the two instruments is excellent with respect 137 to the mode, symmetry, spectral width and lack of skewness in the distributions. The 138 mean is 3.8 m s⁻¹ while the spectral width or standard deviation from MPS data is 0.6 m 139 s⁻¹. The corresponding coefficient of variation (ratio of standard deviation to mean) is 140 15.7%. The finite bin width (±0.1 mm) used causes a "spread" of around 0.5 m s⁻¹ which 141 142 is clearly a significant contributor to the measured coefficient of variation. The definition 143 of sub- or super-terminal fall speeds by (Montero-Martinez et al., 2009) is based on fall





speeds that are, respectively, less than 0.7 times the mean value or greater than 1.3 times the mean value (i.e., exceeding 30% threshold on either side of the mean terminal fall speed). From examining the 1 mm size fall speed histogram there is negligible evidence of occurrences with fall speeds <2.66 m s⁻¹ (sub) or >4.94 m s⁻¹ (super).

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Figure 1. (a) Fall velocity versus diameter (D). The contoured frequency of occurrence from
 2DVD data is shown in color (log scale). The mean fall velocity and ±1o standard deviation bars
 are from MPS. The dark dashed line is from the fit to the laboratory data of Gunn and Kinzer
 (1949) and the grey dashed line is the same except corrected for the altitude of Greeley, CO
 (1.4 km MSL). (b) Relative frequency histograms of fall velocity for the 0.5±0.1 mm and 1±0.1
 mm bins.

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The histogram from MPS for the 0.5 mm sizes shows positive skewness with mean of 1.8 m s⁻¹, spectral width of 0.65 m s⁻¹ and corresponding coefficient of variation nearly doubling to 35%. The finite bin width (\pm 0.1 mm) causes a "spread" of 0.4 m s⁻¹ which contributes to the measured coefficient of variation. Nevertheless, it is not possible to rule out the occurrence of sub- or super-terminal fall speeds, respectively, less than 1.26 m s⁻¹ or exceeding 2.34 m s⁻¹ (i.e., exceeding 30% of the mean value) based on our data.

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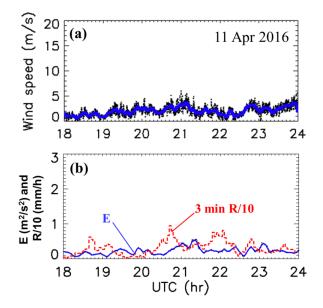
165 2.2 Fall Speeds from Huntsville, Alabama

The first Huntsville event occurred on 11 April 2016 and consisted of precipitation associated with the mesoscale vortex of a developing squall line that moved across northern Alabama between 1800 and 2300 UTC and produced over 25 mm of rainfall in the Huntsville area. Figure 2(a) shows the ambient 10-m height wind speeds (3 s and 5-





min averaged) recorded at the site. Maximum speeds were less than 5 m s⁻¹ and wind 170 gusts were light. As no direct in situ measurement of turbulence was available we use 171 the approach by (Garrett and Yuter, 2014) who estimate the difference between the 172 maximum wind speed, or gust, that was sampled every 3 s, and the average wind 173 speed from successive 5 min intervals. The estimated turbulent intensity is proportional 174 to $E = (Gusts - AverageWind)^2/2$. Figure 2(b) shows the E values which were small 175 (maximum $E < 0.4 \text{ m}^2 \text{ s}^{-2}$) indicative of low turbulence. Also, shown in Fig. 2(b) is the 176 2DVD-based time series of rainfall rate (R) averaged over 3 mins; the maximum R is 177 around 10 mm h^{-1} . 178



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180 <u>Figure 2</u>: (a) 3-s raw and 5-min averaged wind speeds at 10-m height. (b) turbulent 181 intensity estimates E, and 3-min averaged R (note: plot is R/10).

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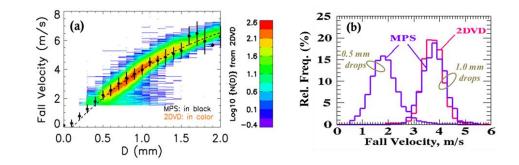
Figure 3(a) shows the fall velocity versus D comparison between the two instruments while panel (b) shows the histograms for the 0.5 and 1 mm bins. Similar to the Greeley event, the mean fall speed agreement between both instruments in the overlap region is excellent and consistent with the fit to the Gunn-Kinzer laboratory data. As in Fig. 1(a), the MPS data in Fig. 3(a) is in excellent agreement with Gunn-Kinzer fit for sizes < 0.5 mm.

The 0.5 and 1 mm histogram shapes in Fig. 3(b) are quite similar to the Greeley case shown in Fig. 1(b). The mean and standard deviations from the MPS data for the 0.5 and 1 mm bins are, respectively, $[2 \ 0.62]$ and $[3.88 \ 0.44]$ m s⁻¹. The comments made





192 earlier with respect to Fig. 1(b) of the Greeley event are also applicable here. In
193 particular, the fall speed histogram for the 0.5 mm sizes cannot rule out the occurrence
194 of sub- or super-terminal fall speeds based on our data.



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196 <u>Figure 3</u>. (a) as in Fig. 1(a) except for 11 April 2016 event. The dashed line is fit to 197 Gunn-Kinzer at sea level. (b) as in Fig. 1(b) except for 11 April 2016 event.

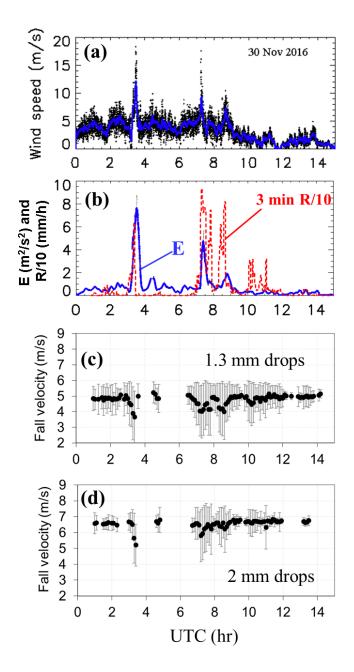
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The second case considered is from 30 November 2016 wherein a supercell passed 199 over the instrumented site from 0300-0330 UTC producing about 15 mins later a long-200 lived EF-2 tornado. Strong winds were recorded at the site with 5-min averaged speeds 201 reaching 10-12 m s⁻¹ between 0320-0330 and E values in the range to 7-8 m² s⁻² 202 indicating strong turbulence (Fig. 4a,b). The rain rates peaked at 70 mm h⁻¹ during this 203 time (Fig. 4b). About 3 h later several squall-line type storm cells passed over the site 204 from 0700-0900 UTC again with strong winds but considerably lower E values 2-4 m² s⁻² 205 and maximum R of 80 mm h⁻¹. After 1000 UTC the E values were much smaller (< 0.5 206 $m^{2} s^{-2}$) indicating calm conditions. The peak R is also smaller at 30 mm h⁻¹ at 1000 UTC. 207

Figure 4 panels (c) and (d) show the mean and $\pm 1\sigma$ of the fall speeds from the 2DVD for 208 the 1.3 and 2 mm drop sizes, respectively. The MPS data are not shown here since 209 during this event it was located outside the DFIR on its turntable and we did not want to 210 confuse the wind-effects between the two instruments. It is clear from Fig. 4(c) that 211 during the supercell passage (0300-0330 UTC) the mean fall speed for 1.3 mm drops 212 decreases (from 5 to 3.5 m s⁻¹) and the standard deviation increases (from 0.5 to 1.5 m 213 s⁻¹). The same trend can be seen for the subsequent squall-line rain cell passage from 214 0700-0900 UTC. 215







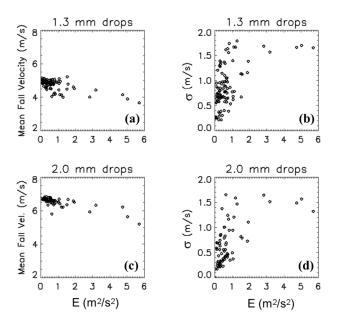
218 <u>Figure 4</u>. (a) as in Fig. 2(a) except for 30 Nov 2016 event. (b) as in Fig. 2(b). (c) mean 219 and $\pm 1\sigma$ standard deviation of fall speeds from 2DVD for 1.3 ± 0.1 mm sizes. (d) as in (c) 220 except for 2 ± 0.1 mm sizes.





222 To expand on this observed correlation, Fig. 5 (a,b) show scatterplots of the mean fall speed and standard deviation versus E for the 1.3 mm drops (panels c and d show the 223 same but for the 2 mm drops). The mean fall speed decreases with increasing E nearly 224 225 linearly for E>1 m² s⁻². This decrease relative to Gunn-Kinzer terminal fall speeds is termed as "sub-terminal" and our data is in general agreement with (Montero-Martinez 226 and Garcia-Garcia 2016) who found an increase in the numbers of sub-terminal drops 227 228 with sizes between 1-2 mm under windy conditions using a 2D-Precipitation probe with resolution of 200 microns (similar to 2DVD) but without wind fence. The standard 229 deviation of fall speeds (σ_f) versus *E* is shown in panels 6 (b,d). When *E*>1 m² s⁻², the σ_f 230 is nearly constant at 1.5 m s⁻¹ for both drop sizes. For E<1, the σ_f is more variable and 231 essentially uncorrelated with E. From the discussion related to Fig. 1(b) and 3(b), $\sigma_{\rm f}$ 232 values exceeding 0.5 m s⁻¹ can be attributed to physical as opposed to instrumental and 233 finite bin width effects. Thus, the fall speed distributions are considerably broadened 234 when $E>1 \text{ m}^2 \text{ s}^{-2}$ due to increasing turbulence levels which is again consistent with the 235 findings of (Montero-Martinez and Garcia-Garcia, 2016) as well as (Garett and Yuter, 236 2014). The latter observations, however, were of graupel fall speeds in winter 237 precipitation using a multi-angle snowflake camera (Garrett et al., 2012). 238

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241 <u>Figure 5.</u> (a,b) mean fall speed and standard deviation, respectively, versus E for 1.3
 242 mm sizes. (c,d) same but for 2 mm sizes.





244 3 Discussion and Conclusions

We have reported on raindrop fall speed distributions using a high resolution (50 245 microns) droplet spectrometer (MPS) collocated with moderate resolution (170 microns) 246 247 2DVD to cover the entire size range (0.1 mm onwards) expected in natural rain. The only comparable earlier study is by (Montero-Martinez et al., 2009) who used collocated 248 2D-cloud and precipitation probes (2D-C, 2D-P) but restricted their data to calm wind 249 250 conditions. Their main conclusion was that the distribution of the ratio of the measured fall speed to the terminal fall speed for 0.44 mm size, while having a mode at 1 was 251 strongly positively skewed with tails extending to 5 especially at high rain rates. In our 252 253 data shown in Fig. 1(b) and 3(b), there is no such strong positive skewness, and the corresponding ratio does not exceed 2. 254

255 Larsen et al., (2014) appear to confirm the ubiquitous existence of super-terminal fall 256 speeds for sizes < 1 mm using different instruments one of them being a 2DVD. It is well-known that "mis-matched" drops cause erroneous fall speed estimates from 2DVD 257 for tiny drops. To clarify the "mis-matched" drop problem: it is very difficult to match a 258 drop detected in the top light-beam plane of the 2DVD to the corresponding drop in the 259 bottom plane for sizes < 0.5 mm (Schoenhuber et al., 2008; Appendix in Huang et al., 260 2010; Bernauer et al., 2015). It is not clear if (Larsen et al., 2014) accounted for this 261 problem in their analysis. 262

Our histograms of fall speeds for 1 mm sizes (Fig. 1b and 3b) under calm wind conditions from both MPS and 2DVD did not show any evidence of either sub- or superterminal speeds, rather the histograms were symmetric with mean close to the Gunn-Kinzer terminal velocity value. However, for the 0.5 mm sizes, our histogram of fall speeds using the MPS under calm conditions cannot rule out the occurrence of both sub- and super-terminal fall speeds, after accounting for instrumental and finite bin width effects.

In a later study using only the 2D-P probe, (Montero-Martinez and Garcia-Garcia, 2016) 270 found sub-terminal fall speeds and broadened distributions under windy conditions for 271 1-2 mm sizes in general agreement with our results using 2DVD. Stout et al., (1995) 272 simulated the motion of drops in isotropic turbulence and determined that there would 273 be a significant reduction of the average drop settling velocity (relative to terminal 274 velocity) of greater that 35% for drops around 2 mm size when the ratio of rms velocity 275 fluctuations (due to turbulence) relative to drop terminal velocity is around 0.8. Whereas 276 we did not have a direct measure of the rms velocity fluctuations, the proxy for 277 turbulence intensity (E) related to wind gusts during supercell passage (very large E 278 279 around 7 m² s⁻²) clearly shows a significant reduction in mean fall speeds of 25-30% relative to terminal speed for 1.3 and 2 mm sizes with significant broadening of the fall 280 speed distributions relative to calm conditions. 281





When $E<0.5-1 \text{ m}^2 \text{ s}^{-2}$, our data show that the mean fall speeds are within a few per cent (<5%) of the *Gunn-Kinzer* terminal velocity over the entire range from 100 microns and larger to precipitation sizes. For sizes > 1 mm, no significant broadening of the fall speed distribution over that ascribed to instrument and/or finite bin widths effects were observed. While our dataset is limited to three events they cover a wind range of rain rates, wind conditions and two different climatologies. Analysis of further events with direct measurement of turbulent intensity would be needed to generalize our findings.

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- 290 Data Availability
- 291 Data used in this paper can be accessed from:
- 292 ftp://lab.chill.colostate.edu/pub/kennedy/merhala/Bringi_et_al_2017_GRL_datasets/
- 293

294 Competing interests

VNB and MT declare they have no conflict of interest. DB is employed by Droplet
Measurements Technologies, Inc. located in Longmont, Colorado, USA who make the
Meteorological Particle Spectrometer used in this study.

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