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2 The instrument constant of sky radiometer (POM-02),

3 Part I: Calibration constant

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12

13 **Abstract**

14 Ground-based networks have been developed to determine the spatiotemporal
15 distribution of aerosols using radiometers. In this study, the accuracy of the calibration

16 constant (V_0) for the sky radiometer (POM-02) which is used by SKYNET was

17 investigated. The temperature dependence of the sensor output was also investigated,

18 and the dependence in the 340, 380, and 2200 nm channels was found to be larger than

19 for other channels, and varied with the instrument. In the summer, the sensor output

20 had to be corrected by a factor of 1.5 to 2% in the 340 and 380 nm channels and by 4%

21 in the 2200 nm channel in the measurements at Tsukuba. In the other channels, the

22 correction factors were less than 0.5%. The accuracy of V_0 from the normal Langley

23 method is between 0.2 and 1.3%, except in the 940 nm channel. The effect of gas

24 absorption was less than 1% in the 1225, 1627, and 2200 nm channels. The

25 degradation of V_0 for shorter wavelengths was larger than that for longer

26 wavelengths. The accuracy of V_0 estimated from the side-by-side measurements was



27 0.1 to 0.5%. The V_0 determined by the improved Langley (IML) method had a
28 seasonal variation of 1 to 3%. The RMS error from the IML method was about 0.6 to
29 2.5%, and in some cases, the maximum difference reached 5%. The trend in V_0 after
30 removing the seasonal variation was almost the same as for the normal Langley
31 method. The calibration method for water vapor in the 940 nm channel was developed
32 using an empirical formula for transmittance. The accuracy of V_0 was better than 1%
33 on relatively stable and fine days. A calibration method for the near-infrared channels,
34 1225, 1627, and 2200 nm, was also developed. The logarithm of the ratio of the sensor
35 output can be written as a linear function of the airmass, by assuming that the ratio of
36 the optical thicknesses between the two channels is constant. The accuracy of V_0 was
37 better than 1% on days with good conditions.

38

39 1. Introduction

40 Atmospheric aerosols are an important constituent of the atmosphere. Aerosols
41 change the radiation budget directly by absorbing and scattering solar radiation and
42 indirectly through their role as cloud condensation nuclei (CCNs), thereby increasing
43 cloud reflectivity and lifetime (e.g., Ramanathan et al. 2001, Lohmann and Feichter
44 2005). Aerosols also affect human health as one of the main components of air
45 pollution.

46 Atmospheric aerosols have a large variability in time and space. Therefore,
47 measurement networks covering an extensive area on the ground and from space have
48 been developed and established to determine the spatiotemporal distribution of
49 aerosols related to climate forcing and air quality on multiple time scales and on
50 regional, hemispheric, and global spatial scales.

51 Ground-based observation systems, such as those using radiometers, are more
52 accurate and easier to install and maintain than space-based systems. Therefore,
53 ground-based observation data are used to validate data obtained from space-based
54 systems (Kahn et al. 2005, Remer et al. 2005, Mélin et al. 2010). Well-known
55 ground-based networks include AERONT (AERosol RObotic NETwork) (Holben et al.
56 1998), SKYNET (Takamura and Nakajima 2004) and PFR-GAW (Precision Filter
57 Radiometer-Global Atmosphere Watch) (Wehrli 2005).

58 In ground-based observation networks, direct solar irradiance and sky radiance are



59 measured, and the aerosol characteristics are retrieved by analyzing these data:
60 optical thickness, single scattering albedo, phase function, complex refractive index,
61 and size distribution. To improve the measurement accuracy, it is important to know
62 the characteristics of the instruments and to calibrate the instruments. Furthermore,
63 from the view point of the validation of optical properties retrieved from the satellite
64 measurement data, it is important to know the accuracy of the ground-based
65 measurements.

66 In SKYNET, the radiometers POM-01 and POM-02, manufactured by Prede Co. Ltd.,
67 Japan, are used. These radiometers are called ‘sky radiometers’, and measure both the
68 solar direct irradiance and sky-radiances. The objectives in this study are to
69 investigate the current status of and problems with the sky radiometer.

70 There are two constants that we must determine to make accurate measurements.
71 One is the calibration constant, and the other is the solid view angle (SVA) of the
72 radiometer. The calibration constant is the output of the radiometer to the
73 extra-terrestrial solar irradiance. The SVA is a constant which relates the sensor
74 output to the sky radiance. The ambient temperature affects the sensor output, and
75 this temperature dependence must be considered when analyzing data from POM-01
76 and POM-02. In this study, the temperature dependence of POM-02 and the calibration
77 of the sensor are described. The SVA is described in detail in Part II.

78 In section 2, we briefly describe the data used in this study. In section 3, firstly, the
79 temperature characteristics of POM-02 are described. Though the majority of POM-01
80 and POM-02 users do not explicitly consider the temperature dependence of the
81 instruments, some channels have a large temperature dependence.

82 Secondly, the accuracy of the calibration constant is described for the sensor output
83 for the extra-terrestrial solar irradiance. Most POM-01 and POM-02 users calibrate
84 the sky radiometers with the Improved Langley (IML) method (Tanaka et al. 1986),
85 because this method only needs on-site measurement data and special measurements
86 for calibration are not required. One of the goals of this paper is to examine the
87 accuracy of the IML method, but before that, in section 4, we consider the accuracy of
88 the normal Langley method using the data obtained at NOAA Mauna Loa observatory
89 (MLO), which is one of most suitable places for sky radiometer calibration by the
90 normal Langley method, and the accuracy of the calibration constant obtained from
91 side-by-side measurement. In section 5, though Campanelli et al. (2004) have already
92 estimated the accuracy of the IML method, we estimate it again and show the time
93 variation and the relation between the calibration constant and temperature
94 dependence, and review the IML method. Then, in section 6, an example of the



95 accuracy of the calibration using a calibrated integrating sphere is shown.

96 In SKYNET, the 940, 1627, and 2200 nm channels were not used. Therefore, the
97 precipitable water vapor (PWV) and the optical thickness at 1627 nm are not
98 estimated. However, these parameters are estimated in AERONET. In sections 7 and 8,
99 calibration methods for these channels are shown using on-site measurement data. In
100 section 9, the results are summarized.

101

102 **2. Data**

103 In this study, the measurements by two POM-02 sky radiometers which are used by
104 the Japan Meteorological Agency/Meteorological Research Institute (JMA/MRI) were
105 conducted. One is used as a calibration reference, POM-02 (Calibration reference), and
106 the other is used for continuous measurement at the Tsukuba MRI observation site,
107 POM-02 (Tsukuba).

108 To calibrate the reference POM-02 by the normal Langley method, the measurements
109 were conducted at the NOAA Mauna Loa Observatory (MLO) for about one month
110 every year, for more than twenty years. The MLO (19.5362°N, 155.5763°W) is located
111 at an elevation of 3397.0 amsl on the northern slope of Mt. Mauna Loa, Island of
112 Hawaii, Hawaii, USA. The atmospheric pressure is about 680 hPa. The MLO is one of
113 the most suitable places to obtain data for a Langley plot and for a solar disk scan.
114 Using these data, the calibration constant is estimated and the SVA is calculated.

115 The continuous observation was performed at the JMA/MRI (36.05°N, 140.13°E) in
116 Tsukuba, which is located about 50 km northeast of Tokyo. Using these continuous
117 measurement data, the calibration constants for the IML method were calculated
118 using the SKYRAD software package (Nakajima et al. 1996, OpenCLASTR,
119 <http://www.ccsr.u-tokyo.ac.jp/~clastr/>). Usually, the calibration of POM-02 for
120 continuous measurement is conducted by comparison with the side-by-side
121 measurement data from the reference POM-02.

122 The temperature dependence of the sensor output was measured using special
123 equipment, which was originally used to measure the temperature dependence of the
124 pyranometer. This equipment is managed and maintained by a branch of the JMA
125 Observation Department.

126

127 **3. Temperature dependence of sensor output**

128 In this section, the temperature characteristics of the POM-02 are described. The
129 POM-02 is temperature-controlled. However, the temperature control is insufficient.
130 Therefore, the sensor output of the POM-02 is dependent on the environmental



131 temperature.

132 The purpose of the temperature control is to keep the temperature inside the
133 instrument from decreasing below levels which will reduce the instrument's accuracy.
134 Instruments are usually designed to operate at temperatures either 20 or 30 °C. When
135 the temperature inside the instrument is below its threshold temperature (20 or 30 °C),
136 the instrument is heated. However, there is no cooling mechanism for when the
137 temperature inside the instrument is higher than its threshold temperature for
138 optimum operation. To monitor the temperature inside the instrument, a temperature
139 sensor is attached near the filter turret. Furthermore, the near-infrared detector,
140 which is thermoelectrically cooled, is equipped with a temperature sensor and
141 temperature data can be recorded.

142 In Fig. 1, an example of the relation between the inside temperature of the
143 instrument and the environmental temperature for POM-02 (Calibration Reference) is
144 shown. The red line is the temperature near the filter turret and the blue line is the
145 temperature of the near-infrared detector. The temperature control setting of this
146 POM-02 is 20 °C. The ambient temperature was varied in the order of 20, 0, -20, 0, 20,
147 40, 20 °C. Since the mounting position of the temperature sensor and the thermal
148 structure of the instrument were different for each product, not every POM-02
149 temperature responds in the same way.

150 In Fig. 2, the relation between the sensor output and the inside temperature near
151 the filter turret for POM-02 (Calibration Reference) is shown. The sensor output is
152 normalized by the sensor output at 20 °C. The ambient environmental temperature
153 was varied from -20 to 40 °C. The detector used for wavelengths less than 1020 nm
154 was a Si photodiode, and the detector for the 1225, 1627, and 2200 nm wavelengths
155 was a thermoelectrically cooled InGaAs photodiode. In this study, the former
156 wavelength region is referred to as the “visible region” and the latter one is the
157 “near-infrared region”.

158 The temperature dependence of the sensor output in the 340 and 2200 nm channels
159 was larger than in the other channels. The range of the atmospheric temperature at
160 Tsukuba was about -5 to 35 °C, and the resulting inside temperatures were between 15
161 and 35 °C (Fig. 1), with the change in the temperature less than 1.5% except for the
162 340 and 2200 nm channels. The temperature dependence of the sensor output varies
163 with the channel.

164 In the 340 nm channel, the sensor output decreased by 7% when the internal
165 temperature increased from 20 to 40 °C. In the 2200 nm channel, the sensor output
166 decreased by a rate of 5 to 6% per 10 degrees of temperature increase. Therefore, the



167 temperature dependence of the sensor output cannot be ignored in these two channels.

168 In Fig. 3, the temperature dependence of the sensor output for POM-02 (Tsukuba) is
169 shown. The temperature dependence of the sensor output in the 340, 380, and 2200 nm
170 channels is large for this POM-02. In the 340 and 380 nm channels, the rate of sensor
171 output decrease was about 1.5% per 10 degrees, and in the 2200 nm channel, the rate
172 of sensor output decrease was about 3% per 10 degrees. In the other channels, the
173 temperature dependence of the sensor output was less than 1% for temperatures
174 between 0 and 40 °C.

175 Though only two examples were shown here, the temperature dependence of the
176 sensor output differed between instruments. If we want to accurately consider the
177 temperature dependence of the sensor output, we need to measure it for each
178 instrument or only use channels with a small temperature dependence.

179

180 4. Normal Langley method

181 To determine the accuracy of the IML method, first the accuracies of the normal
182 Langley method and the transfer of the calibration constant are investigated. The
183 transferred calibration constant can be obtained by comparing the side-by-side
184 measurement of the direct solar irradiance.

185

186 4.1 Accuracy of normal Langley method

187 Figure 4 shows an example of a Langley plot using the data obtained at MLO. In
188 these Langley plots, the data in both the morning and afternoon are plotted. The linear
189 regression lines were determined using the data with airmasses from 2 to 6 AM. From
190 these examples, by using data taken at a location with suitable conditions, it is
191 possible to determine an accurate calibration line.

192 At MLO, ten to twenty measurements for the Langley calibration can usually be
193 taken over a period of 30 to 40 continuous observation days depending on the weather
194 conditions. Near-infrared channels are more sensitive to weather conditions than
195 channels in the visible range, because there are water vapor absorption bands in the
196 near-infrared channel and the water vapor in the atmosphere tends to fluctuate.

197 Table 1 shows the calibration constants (V_0) determined using the data taken from
198 October 2015 to November 2015 at MLO. The calibration constants were calculated for
199 the following four cases.

200 Case 1, no gas absorption, and no temperature correction (NGABS, NTPC)

201 Case 2, no gas absorption, and temperature correction (NGABS, TPC)



202 Case 3, gas absorption, and no temperature correction (GABS, NTPC)
203 Case 4, gas absorption, and temperature correction (GABS, TPC)
204 The error of the calibration constants ($(SD)/V_0$, SD: standard deviation, V_0 is the

205 mean value) were 0.2 to 1.3% except in the 940 nm channel, where the mean V_0 and
206 standard deviation were calculated from all data without weighting. If the weighted
207 mean is used, a smaller error is expected: less than 1%. From these results, it can be
208 seen that the calibration constant can be reliably determined by the normal Langley
209 method using the data taken at MLO.

210 The effect of the temperature dependence is large in the 340 and 2200 nm channels,
211 based on the ratio of (GABS,NTPC)/(GABS,TPC) (= (Case 3)/(Case 4)). In the other
212 channels, the effect of the temperature dependence is less than 0.9%. The range of the
213 atmospheric temperature was about 5 to 15 °C when the measurements for the
214 calibration at MLO were being conducted. Therefore, the effect of the temperature
215 dependence on the sensor output is small.

216 From the ratio of (NGABS,TPC)/(GABS,TPC) (= (Case 2)/(Case 4)), the effect of gas
217 absorption is more than 10% in the 940 nm channel, less than 0.4% in the 1225 and
218 1627 nm channels, and about 1% in the 2200 nm channel. These channels have weak
219 gas absorption by water vapor, CO₂, and CO.

220 As seen from the ratio of (NGABS,NTPC)/(GABS,TPC) (= (Case 1)/(Case 4)), the
221 calibration constants, except in the 340, 940, and 2200 nm channels, can be
222 determined with an error of less than 1% without consideration of the temperature
223 effect and gas absorption by using the data taken at MLO.

224 To calibrate the 940 nm channel, the vertical distribution of water vapor is necessary.
225 The vertical distribution of water vapor is constructed with radiosonde data from the
226 nearest site, precipitable water vapor (PWV) by the Global Positioning System (GPS),
227 and relative humidity measured at MLO, and the transmittance is calculated
228 (Uchiyama et al. 2014). The radiosonde measurements were taken twice a day, and the
229 PWV by GPS were the 30-minute averages. The temporal resolution of these data is
230 not high enough to accurately determine the vertical distribution of the water vapor,
231 resulting in a large error in the calibration constant in the 940 nm channel.

232 Figure 5 shows the annual variation of the calibration constants (V_0) for POM-02
233 (Calibration Reference). The lens in the visible region (Si photodiode region) was
234 replaced in 2013 and the interference filter in the 1225 nm channel was replaced in



235 2014. Since insufficient data were taken due to bad weather conditions in 2007 and
236 2008, the calibration could not be performed with sufficient accuracy. Therefore, the
237 degradation is not smooth in some channels.

238 In general, the degradation at shorter wavelengths is larger than at longer
239 wavelengths in the Si photodiode region. During the period from 2006 to 2013, the
240 changes of V_0 in the 340, 380, and 400 nm channels were -10% per year, -7% per year,

241 and -4% per year, respectively. The changes of V_0 in the 500, 675, and 870 nm
242 channels were about -1% per year, and that in the 1020 nm channel was almost zero.
243 These results indicate that calibration is necessary at least once a year to monitor the
244 degradation of V_0 .

245 The accuracy of the calibration in the near-infrared region is sensitive to weather
246 conditions. Therefore, the annual variation of the calibration constant is not always
247 smooth. However, from 2009 to 2013, the annual change of the calibration constant
248 was less than 1%.

249

250 **4.2 Accuracy of transfer by direct solar measurement**

251 The calibration constant for one instrument can be used to estimate the calibration
252 constant for another instrument by comparison with the simultaneous measurements
253 of the solar direct irradiance.

254 Table 2 shows the results of the calibration constant transferred from POM-02
255 (Calibration Reference) to POM-02 (Tsukuba) and POM-02 (Fukuoka) in December
256 2014. The comparison measurements were conducted over 11 days for POM-02
257 (Tsukuba) and 8 days for POM-02 (Fukuoka). The error (SD/V_0) is 0.1 to 0.5%

258 depending on the wavelength, where mean V_0 is the arithmetic mean. The error is
259 0.5% even for water vapor in the 940 nm channel; usually the fluctuation of the sensor
260 output is large due to fluctuations in the water vapor amount. If the weighted mean is
261 used as the expected value, a smaller error than that of the arithmetic mean is
262 expected.

263 The observations for the comparison depend on the weather conditions, but if
264 instruments are calibrated, it is the most straightforward and accurate way to transfer
265 and determine the calibration constant for different instruments.



266

267 **5. Improved Langley method**

268 **5.1 Comparison between Improved Langley and normal Langley method**

269 In the Improved Langley (IML) method, the temperature dependence of the sensor
270 output is not usually explicitly considered. This means that the calibration constant
271 determined by the IML method implicitly includes the temperature dependence of the
272 sensor output. Before comparing the calibration constant determined by the IML
273 method and that transferred from POM-02 (Calibration Reference), we examined how
274 much the sensor output changes with ambient temperature change.

275 In Fig. 6, the monthly mean values of the inside temperature of POM-02 (Tsukuba)
276 and the temperature of the near-infrared detector are shown. As seen from the figure,
277 these temperatures were controlled in the period from November to April. In Fig. 7, the
278 temperature correction factors are shown, where the reference temperature is 20 °C. In
279 the summer, the sensor output must be corrected by 1.5 to 2% in the 340 and 380 nm
280 channels and by 4% in the 2200 nm channel. In the other channels, the corrections
281 were less than 0.5%: the temperature effect on these channels was small.

282 In Fig. 8, the calibration constants determined by the IML method from January
283 2014 to December 2015 are shown. To compare between the IML and normal Langley
284 methods, the calibration constants transferred from POM-02 (Calibration Reference)
285 are also shown. The observations for the calibration transfer were conducted in
286 December 2013, December 2014, and December 2015, and the calibration constants for
287 POM-02 (Tsukuba) were determined. The calibration constants in other months were
288 obtained by linear interpolation and the temperature correction factor was also taken
289 into consideration. In Fig. 8, the running means of the monthly IML value are also
290 shown.

291 For every channel, the calibration constants determined by the IML method have a
292 seasonal variation: they are larger in the winter and smaller in the summer. The
293 amplitude of the seasonal variation is larger than that of the temperature correction
294 factor. Furthermore, the annual trend of the calibration constant, after removing the
295 seasonal variation, is almost the same as by the normal Langley method.

296 In the 380 nm channel, the calibration constant changes due to the temperature
297 dependence of the sensor output: in the summer the calibration constant decreases by
298 about 2%. The calibration constant (V_0) determined by the IML method changes by up
299 to 6%. Even if the effect of the temperature change is subtracted from the seasonal
300 variation, there is a difference of about 4% between the calibration coefficients. In the



301 400, 500, 675, and 870 nm channels, there is a difference of 1 to 2% between the
302 calibration coefficients, and in the 340 nm channel, there is a difference of 3% between
303 the calibration coefficients. In the 1020 nm channel, since the interference filter was
304 changed in September 2015, a direct comparison is difficult. In Table 3, the statistics of
305 the difference between both calibration coefficients are shown. The RMS error is about
306 0.6 to 2.5% depending on the wavelength. This result is almost the same as in
307 Campanelli et al. (2004). However, the maximum difference between both calibration
308 coefficients was about 1.3 to 4.7%, and these differences are rather large. The statistics
309 of the 3-point running mean for the IML method are also shown in Table 3. The errors
310 are a bit smaller than for the non-smoothed values: the RMS error is about 0.5 to 1.7%.

311 Though the period of comparison is only two years, the calibration constant by the
312 IML method represents the annual trend and implicitly includes the temperature
313 dependence of the sensor output. However, the calibration constant has a seasonal
314 variation of 1 to 3%, and in some cases, the maximum difference reaches about 5%. The
315 2% error in the calibration constant is not significant in a turbid atmosphere, but it is
316 significant in a clear atmosphere, such as in polar and ocean regions. Furthermore,
317 there is a possibility that the seasonal variation of the calibration constant causes an
318 artificial seasonal variation in the retrieved parameters. The seasonal variation can be
319 reduced by smoothing, such as with a running mean. However, over-smoothing
320 dampens the temperature effect of the sensor output.

321

322 5.2 Review of Improved Langley method

323 In this section, the Improved Langley method is briefly reviewed and we examine
324 possible issues with the IML method.

325 The solar direct irradiance at the surface based on the Beer-Bouguer-Lambert Law
326 is written as follows:

$$327 \quad F = F_0 \exp(-m\tau) \quad (1)$$

328 where F and F_0 are the solar irradiance at the surface and the top of the
329 atmosphere, respectively, $m = 1/\mu_0$ is the airmass, μ_0 is the cosine of the solar
330 zenith angle, and τ is the layer optical thickness.

331 The single scattering by aerosol in the almucantar of the sun is given by the following
332 equation:



$$\begin{aligned}
 333 \quad I_1(\mu_0, \phi) &= m\tau\omega_0 P(\cos\Theta)F_0 \exp(-m\tau) \\
 &= m\tau_{sca} P(\cos\Theta)F_0 \exp(-m\tau)
 \end{aligned} \tag{2}$$

334 where a one-layer plane-parallel atmosphere is assumed, $\tau_{sca} = \tau\omega_0$ is the layer
 335 scattering optical thickness, ϕ is the azimuthal angle measured from the solar
 336 principal plane, ω_0 is the single scattering albedo, and $P(\cos\Theta)$ is the normalized
 337 phase function at the scattering angle Θ . The Improved Langley method is based on
 338 these equations.

339 If the sensor output is proportional to the input energy, the following equation can be
 340 written for the direct solar measurement:

$$341 \quad V = V_0 \exp(-m\tau) \tag{3}$$

342 where $V = CF$, $V_0 = CF_0$, and C is the proportional constant (sensitivity). The
 343 contribution of scattered light in the field of view is ignored.

344 The sensor output for the measured single scattering V_1 can be written as follows:

$$\begin{aligned}
 345 \quad V_1 &= CI_1(\mu_0, \phi)\Delta\Omega \\
 &= Cm\tau\omega_0 P(\cos\Theta)F_0 \exp(-m\tau)\Delta\Omega \\
 &= m\tau\omega_0 P(\cos\Theta)V_0 \exp(-m\tau)\Delta\Omega
 \end{aligned} \tag{4}$$

346 where $\Delta\Omega$ is the SVA.

347 From these equations, the following equations can be obtained:

$$348 \quad m\tau = \frac{V_1}{\omega_0 P(\cos\Theta)V_0 \exp(-m\tau)\Delta\Omega} \tag{5}$$

$$349 \quad m\tau_{sca} = \frac{V_1}{P(\cos\Theta)V_0 \exp(-m\tau)\Delta\Omega} \tag{6}$$

350 Then from eq. (3), we can get the following equations:

$$351 \quad \ln V = \ln V_0 - m\tau \tag{7}$$

$$352 \quad \ln V = \ln V_0 - m\tau_{sca} / \omega_0 \tag{8}$$

353 If m , $m\tau$, and $m\tau_{sca}$ can be obtained, the logarithm of the sensor output can be

354 linearly fitted with m , $m\tau$, and $m\tau_{sca}$. The case when the x-axis is m and the



355 y-axis is $\ln V$ corresponds to the normal Langley method, and the case when the
 356 x-axis is $m\tau$ or $m\tau_{sca}$ and the y-axis is $\ln V$ is the Improved Langley method. In
 357 the normal Langley method, the intersection of the y-axis and the regression line is
 358 $\ln V_0$ and the slope of the regression line is $-\tau$. There are two IML methods. If the
 359 x-axis is $m\tau$, the intersection of the y-axis and the regression line is $\ln V_0$ and the
 360 slope is -1 . Otherwise, if the x-axis is $m\tau_{sca}$, the intersection of the y-axis and the
 361 regression line is $\ln V_0$ and the slope is $-1/\omega_0$. The SKYRAD package adopts the
 362 latter method.

363 In the SKYRAD package, two observable quantities are analyzed. One is the direct
 364 solar irradiance (eq. (3)), and the other is defined as

$$365 \quad R(\lambda, \Theta) = \frac{V(\lambda, \Theta)}{V(\lambda, 0)m\Delta\Omega} \quad (9)$$

366 where $V(\lambda, \Theta)$ is the sensor output of the sky radiance measurement for the
 367 scattering angle Θ , $\cos\Theta = \mu_0^2 + (1 - \mu_0^2)\cos\phi$, $\Delta\Omega$ is the SVA of the sky
 368 radiometer, and $V(\lambda, 0)$ is the direct solar irradiance. This is the sky radiance
 369 normalized by the direct solar irradiance.

370 $V(\lambda, \Theta)$ is composed of the single scattering and multiple scattering radiances.
 371 Therefore, eq. (9) can be expressed as follows:

$$372 \quad R(\lambda, \Theta) = \frac{V_1(\lambda, \Theta)}{V(\lambda, 0)m\Delta\Omega} + R_m(\lambda, \Theta) \quad (10)$$

$$= \tau\omega_0 P(\cos\Theta) + R_m(\lambda, \Theta)$$

373 In the SKYRAD package, in the retrieval process for the Improved Langley method the
 374 single scattering and multiple scattering components are estimated by solving the
 375 radiative transfer equation. Once the single scattering component is retrieved, $m\tau$
 376 and $m\tau_{sca}$ are estimated. In the SKYRAD package, the complex refractive indexes for
 377 each channel are given and the measurement data with a scattering angle of less than
 378 30 degrees are used.

379 Once $m\tau$ is obtained, the calibration constants can be estimated from
 380 $\ln V_0 = \ln V + m\tau$. However, in the SKYRAD package, $\ln V_0$ and W_0 are determined



381 by fitting to $\ln V = \ln V_0 - m\tau_{sca}/W_0$. In eq. (8), W_0 is the single scattering albedo.
382 The single scattering albedo is defined as the ratio of the scattering coefficient to the
383 extinction coefficient. Therefore, the single scattering albedo must be a value between
384 zero and one. However, W_0 is frequently greater than 1. The fitted error, number of
385 measurements, and the transmittance are checked. Then, the data passing the check
386 criterion are chosen as the calibration constants.

387 For the 500 nm channel, Figs. 9 shows a scatter plot of V_0 and the optical thickness
388 at 500 nm, a scatter plot of W_0 and V_0 , and (c) a time series of V_0 from January
389 2014 to December 2015. In this case, the V_0 values with errors less than 0.01 were
390 chosen. Figure 9 (a) shows that there is no dependence of V_0 on the optical thickness.
391 Also, there is no dependence of V_0 on the Ångström exponent (not shown). Figure 9
392 (b) shows that V_0 and W_0 are negatively correlated, and that even if the correct W_0
393 is determined, the V_0 are scattered with a width of about 4%. Figure 9 (c) also shows
394 that V_0 is determined to within about a 2% uncertainty and that V_0 has a seasonal
395 variation. These figures show that the V_0 determined by the IML method in the
396 SKYRAD package has a 2% uncertainty.

397 As stated above, V_0 can be determined from $\ln V_0 = \ln V + m\tau$. However, the V_0
398 determined from this equation are systematically overestimated by several percent
399 and have a larger seasonal variation than by linear fitting to $\ln V = \ln V_0 - m\tau_{sca}/W_0$.

400 In the Improved Langley method, the refractive index is fixed. This causes the
401 estimation error in the single and multiple scattering in eq. (10). This may also be the
402 cause of the 2% uncertainty. Some iterative process may be necessary to reduce this
403 uncertainty.



404

405 **6. Calibration using the calibrated light source**

406 In this section, the accuracy of the calibration using the calibrated integrating
407 sphere is described. If POM-02 can be calibrated using the calibrated light source, then
408 POM-02 can be calibrated quickly without being influenced by the weather.

409 In this study, the integrating sphere, which is calibrated and maintained by the
410 Japanese Aerospace Exploration Agency (JAXA), was used (Yamamoto et al. 2002).
411 This integrating sphere is used to calibrate the radiometers which are used to validate
412 satellite remote sensing products.

413 To use the light source, the extra-terrestrial solar irradiance and the SVA of the sky
414 radiometer are necessary, as well as the radiance emitted by the light source. The
415 extra-terrestrial solar irradiance by Gueymard (2004) was used here, along with the
416 SVA obtained by processing the solar disk scan data.

417 When the integrating sphere is measured by POM-02, the sensor output is written
418 as follows:

$$419 \quad V_{sph}(\lambda_0) = \int_{\Delta\lambda} C(\lambda)\varphi(\lambda)I_{sph}(\lambda)d\lambda \cdot \Delta\Omega \bigg/ \int_{\Delta\lambda} \varphi(\lambda)d\lambda \quad (11)$$

420 where $V_{sph}(\lambda_0)$ is the sensor output in channel λ_0 , $C(\lambda)$ is the sensitivity at
421 wavelength λ , $\varphi(\lambda)$ is the spectral response function of the interference filter,
422 $I_{sph}(\lambda)$ is the spectral radiance from the integrating sphere at wavelength λ , and
423 the emitted radiance from the integrating sphere is assumed to be homogeneous. This
424 equation is approximated as follows:

$$425 \quad V_{sph}(\lambda_0) \cong C(\lambda_0)\bar{I}_{sph}(\lambda_0) \cdot \Delta\Omega \quad (12)$$

426 where

$$427 \quad \bar{I}_{sph}(\lambda_0) = \int_{\Delta\lambda} \varphi(\lambda)I_{sph}(\lambda)d\lambda \bigg/ \int_{\Delta\lambda} \varphi(\lambda)d\lambda \quad (13)$$

428 When the extra-terrestrial solar irradiance is measured, the sensor output is written
429 as follows:

$$430 \quad V_{sun}(\lambda_0) = \int_{\Delta\lambda} C(\lambda)\varphi(\lambda)F_0(\lambda)d\lambda \bigg/ \int_{\Delta\lambda} \varphi(\lambda)d\lambda \quad (14)$$

431 where $V_{sun}(\lambda_0)$ is the sensor output in channel λ_0 , and $F_0(\lambda)$ is the
432 extra-terrestrial solar spectral irradiance. This equation is approximated as follows:



433 $V_{sun}(\lambda_0) \cong C(\lambda_0)\bar{F}_0(\lambda_0)$ (15)

434 where

435
$$\bar{F}_0(\lambda_0) = \int_{\Delta\lambda} \varphi(\lambda)F_0(\lambda)d\lambda / \int_{\Delta\lambda} \varphi(\lambda)d\lambda$$
 (16)

436 From eqs. (12) and (15), $V_{sun}(\lambda_0)$ is written as follows:

437
$$V_{sun}(\lambda_0) \cong V_{sph}(\lambda_0) \frac{\bar{F}_0(\lambda_0)}{\bar{I}_{sph}(\lambda_0) \cdot \Delta\Omega}$$
 (17)

438 In Table 4, the calibration constants for POM-02 (Calibration Reference) determined
439 from the integrating sphere measurement are compared with the results of the
440 Langley method. At POM-02 (Calibration Reference), the relative difference was 0.7 to
441 7.6% in channels 2 to 8, and 0.5 to 1.8% in channels 9 to 11. The integrating sphere
442 used in channels 2 to 8 is different from that in channels 9 to 11.

443 The value of the extra-terrestrial solar spectrum is dependent on the database. In
444 Fig. 10, three data sets are shown (Thuillier et al. 2003, Gueymard 2004, and 1985
445 Wehrli Standard Extraterrestrial Solar Irradiance Spectrum (Wehrli 1985, Neckel and
446 Labs 1981)). The ratios of the solar spectrum to Gueymard (2004) are also shown.
447 These figures show that there is a several percent difference in the values depending
448 on the wavelength. The SVA error is 1% (see Part II); the disk scan data were taken at
449 MLO, where measurement conditions were good for the solar disk scan. The accuracy
450 of the integrating sphere was 1.7% (Yamamoto et al. 2002). Considering the magnitude
451 of these errors, the above differences in the calibration constants seem reasonable.
452 However, to reduce the optical thickness error below 0.01, a calibration coefficient
453 error of several percent is too large. For estimating the optical thickness from
454 measurements of the direct solar irradiance, the calibration coefficient determined by
455 the Langley method is better.

456

457 7. Calibration of 940 nm channel

458 The calibration constant depends on the extra-terrestrial solar irradiance in the 940
459 nm band, the spectral response function of the interference filter, the spectral
460 sensitivity of the detector, and the transmittance of radiometer optics. Calibration
461 methods for the 940 nm channel, which is in the water vapor absorption band, have
462 been considered extensively in previous studies (Reagan et al. 1987a, 1987b, 1995;
463 Bruegge et al. 1992; Thome et al. 1992, 1994; Michalsky et al. 1995, 2001; Schmid et al.



464 1996, 2001; Shiobara et al. 1996; Halthore et al. 1997; Cachorro et al. 1998;
465 Plana-Fattori et al. 1998, 2004; Ingold et al. 2000; Kiedron et al. 2001, 2003). For
466 example, Uchiyama et al. (2014) developed the Langley method which takes account of
467 the gas absorption, and the empirical relationship between the transmittance and
468 precipitable water vapor (PWV) was determined from the theoretical calculation by
469 giving the spectral response function and the model atmosphere. The PWV is
470 estimated from the transmittance for the 940 nm channel. The empirical formula is
471 usually used for the transmittance of the 940 nm channel by water vapor.

472 Most POM-02 users have taken measurements without calibrating the 940 nm
473 channel for a long time. To make use of these accumulated data, it is necessary to
474 develop a calibration method using data at the observation site. Campanelli et al.
475 (2014) developed a method to determine the calibration constant and parameters for
476 the empirical formula of the transmittance using the on-site surface meteorological
477 data and simultaneous POM-02 data. However, it is difficult to obtain the empirical
478 formula for transmittance by the column water vapor from the surface measurement
479 data.

480 In this study, given the spectral response function, the empirical transmittance
481 formula is produced by the method shown in Uchiyama et al. (2014). Then, the
482 modified Langley method shown below is performed using the empirical formula and
483 the observation data.

484 The water vapor transmittance is approximated as follows:

$$485 \quad Tr(\text{H}_2\text{O}) = \exp(-a(pwv)^b) \quad (18)$$

486 where a and b are fitting coefficients (see Appendix), and pwv is PWV.

487 The sensor output V is written as follows (Uchiyama et al. 2014):

$$488 \quad V = \frac{V_0}{R^2} \exp(-m(\tau + \tau_R)) Tr(\text{H}_2\text{O}) \\ = \frac{V_0}{R^2} \exp(-m(\tau + \tau_R)) \exp(-a(m \cdot pwv)^b) \quad (19)$$

489 where V_0 is the calibration coefficient, R is the distance between the earth and the
490 sun, τ is the aerosol optical thickness at 940 nm, and τ_R is the optical thickness of
491 the molecular scattering (Rayleigh scattering). The aerosol optical thickness τ at 940
492 nm is interpolated from the optical thicknesses at 870 and 1020 nm.

493 The above equation can be rewritten as follows:



$$\ln VR^2 + m(\tau - \tau_R) = \ln V_0 - a(pwv)^b m^b \quad (20)$$

495 The parameters on the left-hand side are known: V is the measurement value, R
496 and m can be calculated from the solar zenith angle, and τ_R is estimated from the
497 surface pressure. If pwv is constant, then the right-hand side of the equation is a
498 linear function of m^b . Therefore, the values on the left-hand side can be fitted by a
499 linear function of m^b , and the intersection of the y-axis and the fitted line is $\ln V_0$.

500 Before the above-mentioned method was applied to the MRI data, it was first applied
501 to the data taken at MLO, which has more stable weather conditions than Tsukuba,
502 MRI. The results applied to the data taken at MLO in October and November 2014 and
503 in October and November 2015 are shown in Table 5.

504 The calibration coefficients determined in 2014 and 2015 were 2.2973×10^{-4}
505 ($SD/V_0=0.052$) and 2.2954×10^{-4} ($SD/V_0=0.047$), respectively.

506 The calibration coefficients determined by the Langley method with consideration of
507 gas absorption in 2014 and 2015 were 2.3364×10^{-4} ($SD/V_0=0.093$) and 2.3157×10^{-4}
508 ($SD/V_0=0.097$), respectively. Though there are differences of 1.7% and 0.9%, these
509 values are very similar. The error of the new method is smaller than the method which
510 takes account of gas absorption more accurately than the new method. This may be
511 due to errors in the estimates of the water vapor amount and distribution: the PWV is
512 obtained from the GPS PWV, which has a low time resolution (30-minute average) and
513 some data are missing, and the vertical distribution is estimated from only two
514 radiosonde measurements per day near MLO.

515 The water vapor amount tends to fluctuate. Though the restriction that the PWV be
516 constant is severe, the above method is applied to the data taken at Tsukuba, MRI and
517 the calibration constants are compared with the calibration constant for POM-02
518 (Calibration Reference), which was calibrated by the Langley method with
519 consideration of the gas absorption using the data taken at MLO and interpolated to
520 the observation day (see Table 6).

521 The ratio of the calibration coefficients in the period from December 14, 2014 to
522 January 5, 2015 (10 cases) was 1.0094, and in the period December 1, 2015 to
523 December 30, 2015 (17 cases) it was 0.99818. Thus, the difference between the two



524 methods is less than 1%.

525 Although it seems that the above-mentioned modified Langley method does not work
526 well at all locations and under all weather conditions, the calibration constant of the
527 940 nm channel can be determined by applying the above-mentioned method on a
528 relatively stable and fine day at the observation site.

529

530 **8. Calibration coefficients of near-infrared channels**

531 The measurements for the near-infrared channels, 1225, 1627, and 2200 nm, of
532 POM-02 have been performed at many SKYNET sites, but the data have not been
533 analyzed, because most POM-02 users cannot calibrate these channels by themselves.

534 These channels can be calibrated with the Langley method with a reasonable
535 accuracy by taking into account the gas absorption. However, many users cannot make
536 these measurements for the Langley method. Furthermore, the scattering of light in
537 these channels is small and the IML method cannot be applied.

538 For some observation days, data with a very high correlation between channels may
539 be obtained. In this case, if the calibration constant of one channel is known, then the
540 calibration constants of the other channels can be inferred. The general method for the
541 case when the ratio of the optical thicknesses is constant was shown by Forgan (1994).

542 In this study, by assuming that the channels in the visible region including the 940
543 nm channel are calibrated, a similar method was applied to the near-infrared channels
544 to determine the calibration constant and the accuracy was investigated.

545 The sensor output of POM-02 is written as follows:

$$546 \quad V = \frac{V_0}{R^2} \exp(-m(\tau + \tau_R)) Tr(gas) \quad (21)$$

547 where V is the sensor output, V_0 is the calibration constant, R is the distance
548 between the earth and the sun, m is the airmass, τ is the aerosol optical thickness,
549 τ_R is the optical thickness of the molecular scattering (Rayleigh scattering), and
550 $Tr(gas)$ is the transmittance of the gas absorption.

551 The sensor output for channels 1 and 2 are as follows:

$$552 \quad V_1 = \frac{V_{01}}{R^2} \exp(-m(\tau_1 + \tau_{R1})) Tr_1(gas) \quad (22)$$

$$553 \quad V_2 = \frac{V_{02}}{R^2} \exp(-m(\tau_2 + \tau_{R2})) Tr_2(gas). \quad (23)$$

554 The calibration constant of channel 1 is assumed to be known and that of channel 2 is
555 determined.



556 From eqs. (22) and (23), the following equation is obtained:

$$557 \frac{V_2}{V_1} = \frac{V_{02} \exp(-m(\tau_2 + \tau_{R2})) Tr_2(gas)}{V_{01} \exp(-m(\tau_1 + \tau_{R1})) Tr_1(gas)}$$

558 Therefore,

$$559 \ln \frac{V_2}{V_1} = \ln \frac{V_{02}}{V_{01}} - m(\tau_2 - \tau_1) - m(\tau_{R2} - \tau_{R1}) + \ln \frac{Tr_2(gas)}{Tr_1(gas)} \quad (24)$$

$$560 \ln \frac{V_2}{V_1} + m(\tau_{R2} - \tau_{R1}) - \ln \frac{Tr_2(gas)}{Tr_1(gas)} = \ln \frac{V_{02}}{V_{01}} - m(\tau_2 - \tau_1) \quad (25)$$

$$= \ln \frac{V_{02}}{V_{01}} - \left(\frac{\tau_2}{\tau_1} - 1\right) \tau_1 m$$

561 If the water vapor amount is estimated from the 940 nm channel, and the mixing
 562 ratio of CO₂ and CO is given, then the transmittance of gas can be estimated. Given
 563 the observation time and the latitude and longitude of the observation site, the
 564 air mass is calculated, and τ_{R1} and τ_{R2} are calculated from the surface pressure.
 565 Therefore, the left-hand side of eq. (25) is known. Furthermore, if the ratio of the
 566 optical thickness τ_2/τ_1 is constant, then this equation is a linear function of $m\tau_1$.

567 Therefore, the intersection of the y-axis and the linearly fitted line is $\ln V_{02}/V_{01}$, and if

568 V_{01} is known, then V_{02} is also known. Although this condition is not always satisfied,
 569 sometimes a linear fit will provide sufficient accuracy.

570 This method was applied to the data of POM-02 (Calibration Reference) from
 571 December 2014 to December 2015. The 500 nm was chosen as channel 1 in eq. (25).
 572 The data used here had an RMS error of 0.005. In Fig. 11(a), the monthly mean of

573 V_{02}/V_{01} and the standard deviation are shown. The lines of the ratio, which are

574 interpolated from the calibration constant determined using the data taken in October
 575 and November of 2014 and 2015 at MLO, are also shown. In Fig. 11 (b), the ratio of the
 576 calibration constant by the above method and the interpolated value of the calibration
 577 constant determined from the MLO data are shown. In the 1627 nm channel, the
 578 differences are less than 2% throughout the year and the differences in December and
 579 January are less than 1%. In the 1225 nm channel, the differences are less than 2%
 580 except in April 2015. In the 2200 nm channel, the differences in some months are more
 581 than 3%. However, in December 2015, the differences in all channels are less than 1%.

582 This shows that the calibration constant can be determined by the method shown here
 583 to within an accuracy of 1% under suitable conditions. Currently, there is no method to



584 calibrate the near-infrared channel from on-site observation data. The method shown
585 here is an alternative to the Langley method.

586

587 **9. Summary and conclusion**

588 Atmospheric aerosols are an important constituent of the atmosphere. Measurement
589 networks covering an extensive area from ground and space have been developed to
590 determine the spatiotemporal distribution of aerosols. SKYNET is a ground-based
591 monitoring system using sky radiometers POM-01 and POM-02 manufactured by
592 Prede Co. Ltd., Japan. To improve their measurement accuracy, it is important to know
593 the characteristics of the instruments and accurately calibrate them accordingly.

594 There are two constants that we must determine to make accurate measurements.
595 One is the calibration constant, and the other is the SVA of the radiometer.
596 Additionally, the temperature dependence of the sensor output is another important
597 characteristic.

598 In this study, the data obtained by two sky radiometers POM-02 of the JMA/MRI are
599 considered. One of the sky radiometers is used as a calibration reference, and the other
600 is used for continuous measurement at the Tsukuba MRI observation site.

601 The sensor output of POM-02 is dependent on the environmental temperature. The
602 temperature dependence of the sensor output in the 340, 380, and 2200 nm channels
603 was larger than in other channels. For example, the sensor output in the 340 and 380
604 nm channels of POM-02 (Tsukuba) increased at a rate of about 1.5% per 10 degrees,
605 and that in the 2200 nm channel increased at a rate of about 3% per 10 degrees. In the
606 other channels, the sensor output increased at a rate of less than 1% when the sensor's
607 internal temperature was 0 to 40 °C. The temperature dependence of the two POM-02
608 examined here was different for each instrument. If we want to make accurate
609 measurements, we need to measure the temperature dependence for each instrument
610 or use the channels with a small temperature dependence.

611 For the measurement at Tsukuba, the temperature inside the POM-02 (Tsukuba)
612 was controlled during the winter and spring seasons from November to April, but was
613 not regulated, and thus was high during the summer. In the summer, sensor output
614 must be corrected by 1.5 to 2% in the 340 and 380 nm channels and by 4% in the 2200
615 nm channel. In the other channels, the corrections were less than 0.5%.

616 As well as determining the accuracy of the IML method, this study investigated the
617 accuracy of the normal Langley method and of the calibration transfer. From the data
618 taken at MLO, the error in the calibration constants determined by the normal
619 Langley method (SD/V_0) was 0.2 to 1.3%, except in the 940 nm channel. The effect of



620 gas absorption was more than 10% in the 940 nm channel, but was less than 0.4% in
621 the 1225 and 1627 nm channels and less than 1% in the 2200 nm channel, which all
622 have weak gas absorption.

623 The comparison measurements for transferring the calibration constant were
624 conducted in December at Tsukuba over about ten days. The error (SD/V_0) for the
625 transfer method was 0.1 to 0.5% depending on the wavelength. Though the
626 measurements for the comparison depend on the weather conditions, if there are
627 calibrated instruments, then it is a straightforward and accurate way to determine the
628 calibration constant.

629 The annual variation of the calibration constants (V_0) for POM-02 (Calibration
630 Reference) was also investigated. Roughly speaking, the degradation in the shorter
631 wavelengths was larger than that in the longer wavelengths in the Si photodiode
632 region. The changes in the 340 nm channel were -10% per year from 2006 to 2013. The
633 change in the near-infrared region (thermoelectrically cooled InGaAs photodiode) was
634 less than 1% from 2009 to 2013. These results indicate that calibration of the
635 instruments is necessary at least once a year to monitor the degradation of V_0 .

636 The calibration constant determined by the IML method and that transferred from
637 the POM-02 (Calibration Reference) were compared using the data taken at Tsukuba
638 from December 2013 to December 2015.

639 For every channel, the calibration constants determined by the IML method had a
640 seasonal variation of 1 to 3%. The calibration constants determined by the IML method
641 implicitly include the temperature dependence of the sensor output. However, even if
642 the change due to the temperature variation is subtracted from the seasonal variation,
643 there is a difference of 1 to 4% between the two calibration coefficients. The RMS
644 errors of the differences between the two calibration coefficients were about 0.6 to 2.5%.
645 This result is almost the same as that of Campanelli et al. (2004). However, in some
646 cases, the maximum difference reached up to 5%. Furthermore, the annual trend of the
647 calibration constant excluding the seasonal variation was almost the same as for the
648 normal Langley method.

649 The IML method was reviewed and its error characteristics were investigated. There
650 was no dependence of V_0 on the optical thickness and Ångström exponent, and V_0
651 and W_0 were negatively correlated. Furthermore, even if the optical thickness and



652 W_0 are determined, the V_0 had an uncertainty of $\pm 2\%$. The time series of V_0 showed
653 a seasonal variation with an uncertainty of $\pm 2\%$. These results show that the V_0
654 determined by the IML method using the SKYRAD package has about a $\pm 2\%$
655 uncertainty.

656 We also tried to determine V_0 using the calibrated integrating sphere as the light
657 source. The relative differences of V_0 were about 1 to 8% depending on the
658 wavelength. Considering the magnitude of the errors in the extra-terrestrial solar
659 spectrum, SVA, and the integrating sphere, the above differences in the calibration
660 constants seem reasonable. However, to reduce the optical thickness error below 0.01,
661 an error of several percent in the calibration coefficient is too large.

662 The calibration method for water vapor channel in the 940 nm was developed using
663 the on-site measurement data. V_0 was determined by the modified Langley method
664 using a pre-determined empirical transmittance equation. The differences in the
665 calibration coefficients between the normal Langley method and the above method
666 were less than 1% on relatively stable and fine days.

667 The calibration method for the near-infrared 1225, 1627, and 2200 nm channels was
668 also developed using the on-site measurement data. It is assumed that channels in the
669 visible wavelength region and the 940 nm channel are calibrated. Then, if the ratio of
670 the optical thicknesses between two channels is constant, the logarithm of the ratio of
671 the sensor output can be written as a linear function of the airmass. Here, the
672 calibration constant for one of the two channels is known and the transmittance of
673 water vapor is calculated using the PWV estimated from the 940 nm channel. By
674 fitting the logarithm of the ratio of sensor output to a linear function of the airmass,
675 the ratio of the calibration constants is determined. By this method, the calibration
676 constants could be determined to an accuracy of within 1% on the days with good
677 weather conditions.

678 In this study, it is shown that some channels have a non-negligible temperature
679 dependence in the sensor output and that the calibration constants determined by the
680 IML method showed a seasonal variation and had an uncertainty of $\pm 2\%$. In one
681 channel, the maximum error reached about 5%. Reducing the uncertainty of the IML
682 method is a task for future work, along with the problems related to the determination



683 of calibration constants. In particular, the calibration constants for the 940 nm channel
684 and the near-infrared channels must be determined using on-site measurement data.

685

686 **Appendix**

687 Coefficients of water vapor transmittance.

688 Details of the method for determining the coefficients a and b are described in
689 Uchiyama et al. (2014). The coefficients a and b depend on the vertical structure of
690 the atmospheric temperature and humidity. Therefore, it is difficult to choose suitable
691 values that can be applied under all atmospheric conditions. The range of variability of
692 transmittance for an atmospheric profile is limited. Atmospheric transmittance is
693 computed for a broad range of atmospheric conditions, and values for a and b were
694 chosen that best fit the ensemble conditions.

695 The value of coefficients determined by our method for POM-02 (Calibration
696 Reference) are $a = 0.139186$, $b = 0.631$. The values of the coefficients for the
697 trapezoidal spectral response function, which has full width at a half maximum of 10
698 nm and central wavelength of 940 nm, are $a = 0.147101$, $b = 0.625$.

699

700 **Acknowledgements**

701 This work was supported by the NIES GOSAT-2 project, Japan. This work was
702 partially supported by JSPS KAKENHI Grant Number JP17K00531.

703

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831 Table titles

832 Table 1 Example of calibration constants (V_0) determined from the data taken at MLO.

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834 Table 2 Results of the calibration constant transferred from POM-02 (Calibration
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848 Figure captions

849 Fig. 1 Relation between the inside temperature of the instrument and the ambient
850 environmental temperature for POM-02 (Calibration Reference).

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852 Fig. 2 Relation between the sensor output and the inside temperature near the filter
853 turret for POM-02 (Calibration Reference). The sensor output is normalized by that at
854 20°C. The error bars are the standard deviation. (a) 340, 380, 400, and 500 nm. (b) 675,
855 870, 940, and 1020 nm. (c) 1225, 1627, and 2200 nm.

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865 Fig. 6 Monthly mean values and standard deviation of the inside temperature of
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870 POM-02 (Tsukuba) from December 2013 to December 2016.

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872 Fig. 8 Time series of the calibration constant for POM-02 (Tsukuba) from January 2014
873 to December 2015. Blue open squares with error bars denote the calibration constants
874 determined by the IML method. The green line shows the 3-point running mean of IML,
875 and the red line is the calibration constant transferred from POM-02 (Calibration
876 Reference).

877

878 Fig. 9 (a) Scatter plot of the optical thickness at 500 nm and V_0 by the IML method

879 for the 500 nm channel. (b) Scatter plot of W_0 and V_0 by the IML method for the 500

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882

883 Fig. 10 (a) Extra-terrestrial solar spectra. The red line is Gueymard (2004), the blue
884 line is Thuillier et al. (2003), and the green line is Wehrli (1985). (b) Ratios of the solar
885 spectrum to Gueymard (2003).

886

887 Fig. 11 (a) Monthly mean of V_{02}/V_{01} and the standard deviation, here

888 $V_{01} = V_0(500\text{nm})$. (b) Ratio of V_{02} to the interpolated value of the calibration constant

889 determined by the Langley method. The red symbols are 1225 nm, blue ones are 1627
890 nm, and green ones are 2200 nm.

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Table 1 Example of calibration constants (V_0) determined from the data taken at MLO.

Wavelength (nm)	340	380	400	500	675	870	940	1020	1225	1627	2200
Case 1 V_0 (NGABS,NTPC)($\times 10^{-4}$)	0.19881	0.39320	1.6407	2.7731	3.2849	2.4785	1.9679	1.5521	0.88035	1.4385	0.72407
SD($\times 10^{-4}$)	0.00117	0.00175	0.0062	0.0082	0.0054	0.0145	0.1769	0.0078	0.01116	0.0123	0.00449
ERR(=SD/ V_0)	0.00588	0.00446	0.0038	0.0030	0.0016	0.0058	0.0899	0.0050	0.01267	0.0085	0.00621
Case 2 V_0 (NGABS,TPC)($\times 10^{-4}$)	0.20507	0.39505	1.6470	2.7624	3.2572	2.4820	1.9540	1.5658	0.87322	1.4258	0.68699
SD($\times 10^{-4}$)	0.00179	0.00168	0.0062	0.0078	0.0071	0.0143	0.1759	0.0095	0.01155	0.0132	0.00792
ERR(=SD/ V_0)	0.00874	0.00425	0.0038	0.0028	0.0022	0.0058	0.0900	0.0061	0.01323	0.0093	0.01153
Case 3 V_0 (GABS,NTPC) ($\times 10^{-4}$)	0.19881	0.39319	1.6407	2.7732	3.2850	2.4785	2.3426	1.5524	0.88353	1.4397	0.73073
SD($\times 10^{-4}$)	0.00117	0.00175	0.0062	0.0081	0.0054	0.0144	0.2356	0.0078	0.01080	0.0122	0.00442
ERR(=SD/ V_0)	0.00588	0.00445	0.0038	0.0029	0.0016	0.0058	0.1006	0.0050	0.01222	0.0085	0.00604
Case 4 V_0 (GABS,TPC) ($\times 10^{-4}$)	0.20506	0.39504	1.6469	2.7626	3.2573	2.4820	2.3261	1.5664	0.87646	1.4268	0.69471
SD($\times 10^{-4}$)	0.00179	0.00168	0.0062	0.0078	0.0072	0.0143	0.2340	0.0098	0.01133	0.0132	0.00837
ERR(=SD/ V_0)	0.00874	0.00425	0.0038	0.0028	0.0022	0.0057	0.1006	0.0062	0.01292	0.0093	0.01205
(Case 1)/(Case 4) - 1.0	-0.0305	-0.0047	-0.0038	0.0038	0.0085	-0.0014	-0.1540	-0.0091	0.0044	0.0082	0.0423
(Case 2)/(Case 4) - 1.0	0.0000	0.0000	0.0001	-0.0001	0.0000	0.0000	-0.1600	-0.0004	-0.0037	-0.0007	-0.0111
(Case 3)/(Case 4) - 1.0	-0.0305	-0.0047	-0.0038	0.0038	0.0085	-0.0014	0.0071	-0.0089	0.0081	0.0090	0.0518
				: ABS(ERR)> 0.03							: ABS(ERR)<0.01

V_0 : mean value of calibration constant in 2015 MLO observation

SD: standard deviation

ERR=SD/ V_0

GABS: consideration of gas absorption

NGABS: no consideration of gas absorption



TPC: consideration of temperature correction

NTPC: no consideration of temperature correction

Table 2 Results of the calibration constant transferred from POM-02 (Calibration Reference) to POM-02(Tsukuba) and POM-02 (Fukuoka) in December 2014.

Site: Tsukuba

Period: 2015/12/01 to 2016/01/01

No. of days: 11

SN: PS1202091 Calibrated by Sky radiometer
PS1207831

Wavelength (nm)	340	380	400	500	675	870	940	1020	1225	1627	2200
$V_0 (\times 10^{-4})$	0.17469	0.25711	1.1621	2.9248	3.4792	2.2969	1.9900	0.79227	0.87065	1.4074	0.76879
$SD (\times 10^{-4})$	0.00050	0.00065	0.0021	0.0039	0.0045	0.0085	0.0087	0.00427	0.00321	0.0072	0.00402
ERR (=SD/ V_0)	0.00284	0.00253	0.0018	0.0013	0.0013	0.0037	0.0044	0.00539	0.00369	0.0051	0.00523

Site: Fukuoka

Period: 2015/12/4 to 2015/12/20

No. of days: 8

SN: PS1202071 Calibrated by Skyradiometer PS1207831

Wavelength (nm)	340	380	400	500	675	870	940	1020	1225	1627	2200
$V_0 (\times 10^{-4})$	0.18374	0.23346	1.2332	2.9179	3.5176	2.3021	1.9827	1.8899	0.84113	1.2783	0.60461
$SD (\times 10^{-4})$	0.00028	0.00025	0.0014	0.0025	0.0041	0.0044	0.0106	0.0031	0.00279	0.0027	0.00113
ERR (=SD/ V_0)	0.00155	0.00107	0.0011	0.0008	0.0012	0.0019	0.0053	0.0016	0.00331	0.0021	0.00186

V_0 : mean value,

SD: standard deviation



Table 3 Statistics of difference between IML method and normal Langley method.

Wavelength (nm)	340	380	400	500	675	870	1020
V0 ($\times 10^{-4}$)	0.17600	0.26022	1.1840	2.9161	3.4681	2.2863	1.2487
BIAS ($\times 10^{-4}$)	-0.00136	-0.00428	-0.0042	0.0083	-0.0125	0.0017	0.0048
RMS ($\times 10^{-4}$)	0.00325	0.00649	0.0198	0.0225	0.0209	0.0197	0.0309
DFMAX ($\times 10^{-4}$)	0.00725	0.01218	0.0368	0.0475	0.0445	0.0489	0.0558
DFMIN ($\times 10^{-4}$)	0.00006	0.00059	0.0004	0.0011	0.0028	0.0002	0.0006
BIAS/V0	-0.0077	-0.0164	-0.0036	0.0028	-0.0036	0.0008	0.0039
RMS/V0	0.0184	0.0249	0.0167	0.0077	0.0060	0.0086	0.0247
DFMAX/V0	0.0412	0.0468	0.0311	0.0163	0.0128	0.0214	0.0447
DFMIN/V0	0.0003	0.0023	0.0004	0.0004	0.0008	0.0001	0.0005
V0_3RM ($\times 10^{-4}$)	0.17604	0.26037	1.1848	2.9150	3.4678	2.2858	1.2794
BIAS ($\times 10^{-4}$)	-0.00114	-0.00389	-0.0030	0.0093	-0.0124	0.0022	0.0030
RMS ($\times 10^{-4}$)	0.00303	0.00594	0.0178	0.0181	0.0171	0.0163	0.1014
DFMAX ($\times 10^{-4}$)	0.00495	0.01065	0.0321	0.0398	0.0352	0.0338	0.3663
DFMIN ($\times 10^{-4}$)	0.00022	0.00006	0.0007	0.0032	0.0013	0.0006	0.0004
BIAS/ V0_3RM	-0.0065	-0.0149	-0.0025	0.0032	-0.0036	0.0010	0.0023
RMS/ V0_3RM	0.0172	0.0228	0.0150	0.0062	0.0049	0.0071	0.0793
DFMAX/ V0_3RM	0.0281	0.0409	0.0271	0.0137	0.0102	0.0148	0.2863
DFMIN/ V0_3RM	0.0013	0.0002	0.0006	0.0011	0.0004	0.0002	0.0003

V0: mean calibration constant (IML method) during Jan. 2014 to Dec. 2015.

V0_3RM: mean calibration constant (IML method, 3-point running mean) during Jan. 2014 to Dec. 2015.

BIAS: bias (mean of differences between IML and normal Langley methods)

RMS: root mean squares of differences between IML and normal Langley methods

DFMAX: maximum difference between IML and normal Langley methods



DFMIN: minimum difference between IML and normal Langley methods



Table 4 Calibration constants for POM-02 determined by using the calibrated integrating sphere measurement

λ_0 (nm)	\bar{F}_0 (mW/m ² /nm)	\bar{I}_{sph} (mW/m ² /sr/nm)	$\Delta\Omega$ ($\times 10^{-4}$) (sr)	V_{sph} ($\times 10^{-10}$)	V_{sun} ($\times 10^{-4}$)	V_0 ($\times 10^{-4}$)	$(V_{sun} - V_0)/V_0$ (%)	I
340	1036.2	-	2.3970	-	-	0.19884	-	-
380	1210.6	24.5	2.4370	1.9699	0.39941	0.39280	1.68	PTFE(4.17A(50W)x4)
400	1523.3	49.6	2.4190	13.376	1.6982	1.6434	3.34	PTFE(4.17A(50W)x4)
500	1964.6	238.1	2.4170	87.342	2.9817	2.7703	7.63	PTFE(4.17A(50W)x4)
675	1496.5	764.1	2.4220	409.02	3.3075	3.2850	0.69	PTFE(4.17A(50W)x4)
870	958.1	1171.1	2.4310	772.57	2.6000	2.4708	5.23	PTFE(4.17A(50W)x4)
940	822.0	1218.8	2.4520	878.84	2.4173	2.3364	3.46	PTFE(4.17A(50W)x4)
1020	698.1	1236.8	2.4520	682.48	1.5710	1.5559	0.97	PTFE(4.17A(50W)x4)
1225	466.5	537.3	1.9800	204.73	0.89767	0.88715	1.19	PTFE(3.30A(50W)x4)
1627	236.0	377.2	2.0000	459.62	1.4378	1.4456	-0.54	PTFE(3.30A(50W)x4)
2200	82.0	128.2	2.0570	237.19	0.73756	0.72472	1.77	PTFE(3.30A(50W)x4)

V_0 : calibration constant by normal Langley method



Table 5 Calibration constant at 940 nm by modified Langley method using the data taken at MLO.

		Langley	modified Langley	Ratio
2014	V0 ($\times 10^{-4}$)	2.3364	2.2973	0.9833
	SD ($\times 10^{-4}$)	0.2183	0.1195	
	SD/V0	0.0934	0.0520	
	No. of data	19	19	
2015	V0 ($\times 10^{-4}$)	2.3157	2.2954	0.9912
	SD ($\times 10^{-4}$)	0.2236	0.1077	
	SD/V0	0.0966	0.0469	
	No. of data	30	20	

The data taken at MLO in 2014 and 2015 were used.

V0: mean value

SD: standard deviation

Ratio = (modified Langley V0)/(Langley V0)



Table 6 Same as Table 5 but using the data taken at Tsukuba, MRI.

		Langley	modified Langley	Ratio
2014	V0 ($\times 10^{-4}$)	2.3343	2.3562	1.0094
	SD ($\times 10^{-4}$)	0.0002	0.1429	
	SD/V0	0.0001	0.0598	
	No. of data	10	10	
2015	V0 ($\times 10^{-4}$)	2.3132	2.3090	0.9982
	SD ($\times 10^{-4}$)	0.0006	0.1043	
	SD/V0	0.0003	0.0452	
	No. of data	17	17	

The data taken at MRI, Tsukuba in December 2014 and December 2015 were used.

V0: mean value

SD: standard deviation

Ratio = (modified Langley V0)/(Langley V0)



Fig.1

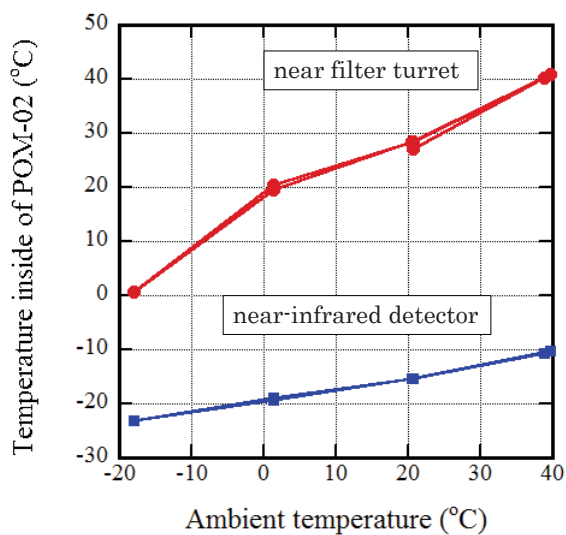


Fig. 1 Relation between the inside temperature of instrument and the ambient environmental one for POM-02 (Calibration Reference).



Fig. 2

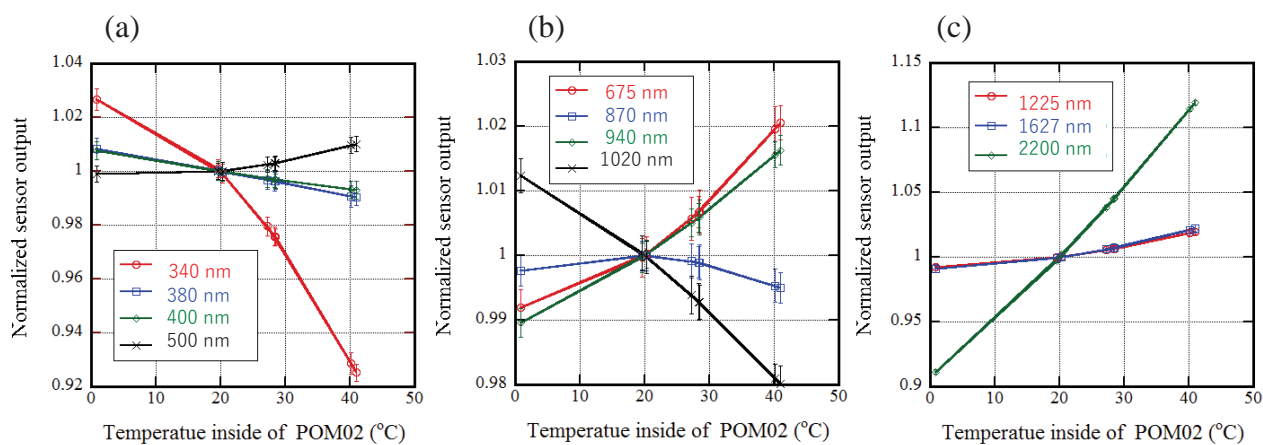


Fig. 2 Relation between sensor output and the inside temperature near filter turret for POM-02 (Calibration Reference). The sensor output is normalized by that at 20° C, (a) 340, 380, 400 and 500nm. (b) 675, 870, 940 and 1020nm. (c) 1225, 1627 and 2200nm.



Fig. 3

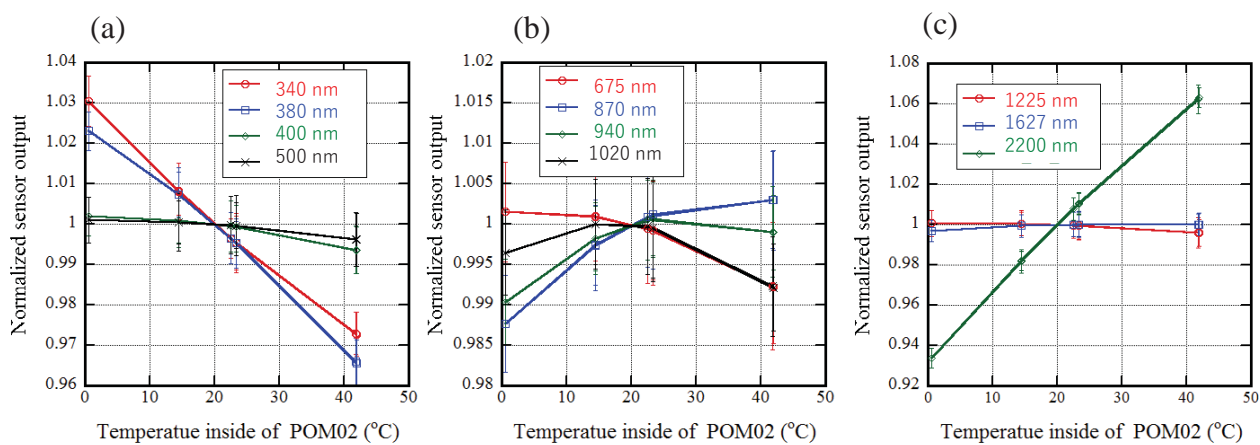


Fig. 3 Same as Fig. 2 but for POM-02 (Tsukuba).



Fig. 4

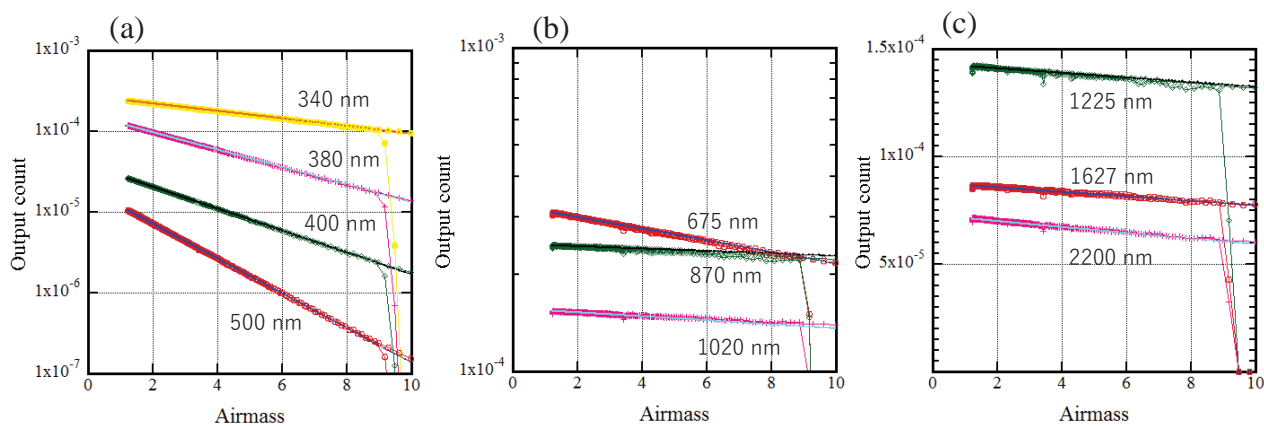


Fig. 4 Examples of Langley plots using the data obtained at MLO, on November 3, 2015.



Fig. 5

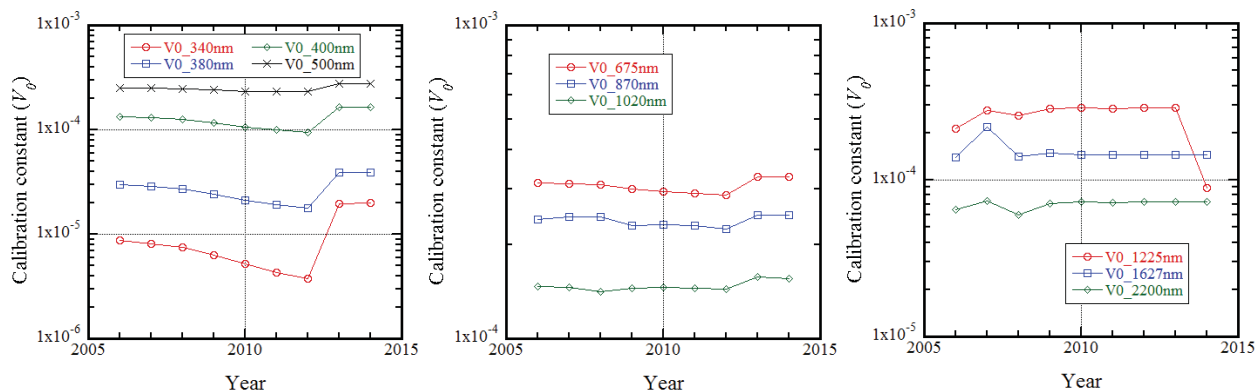


Fig. 5 The annual variation of the calibration constants (V_ϕ) for POM-02 (Calibration Reference).



Fig. 6

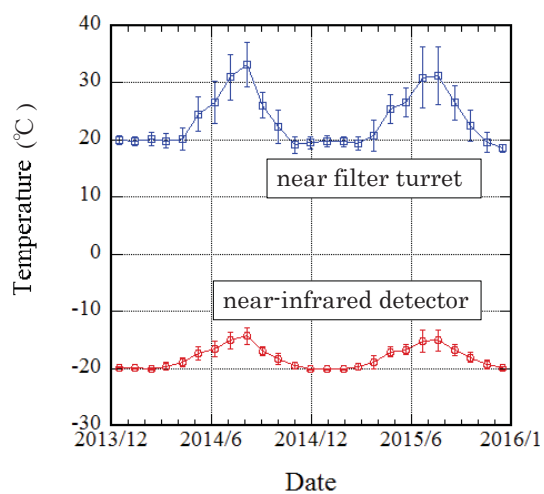


Fig. 6 Monthly mean values and standard deviation of the inside temperature of POM-02 (Tsukuba) (blue line) and the temperature of near-infrared detector (red line) from December, 2013 to December, 2016.



Fig. 7

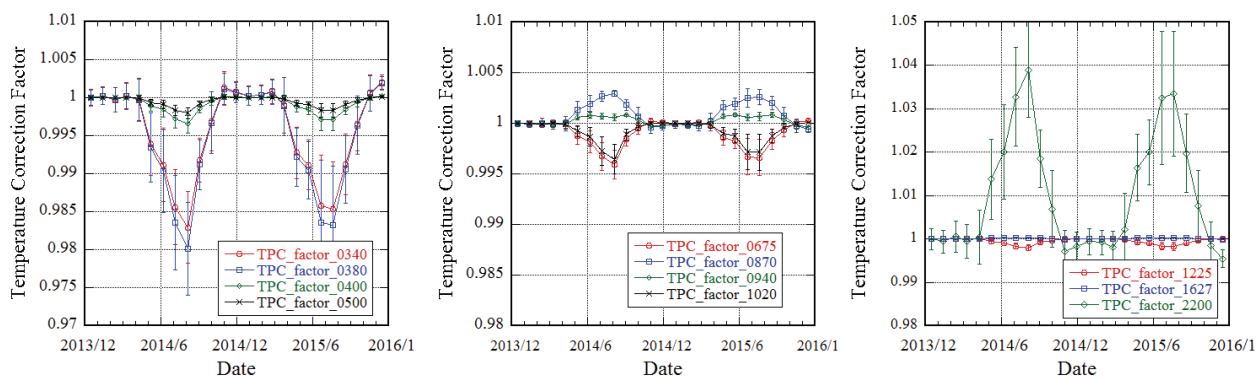


Fig. 7 The monthly means of temperature correction factors and standard deviation for POM-02 (Tsukuba) from December, 2013 to December, 2016.

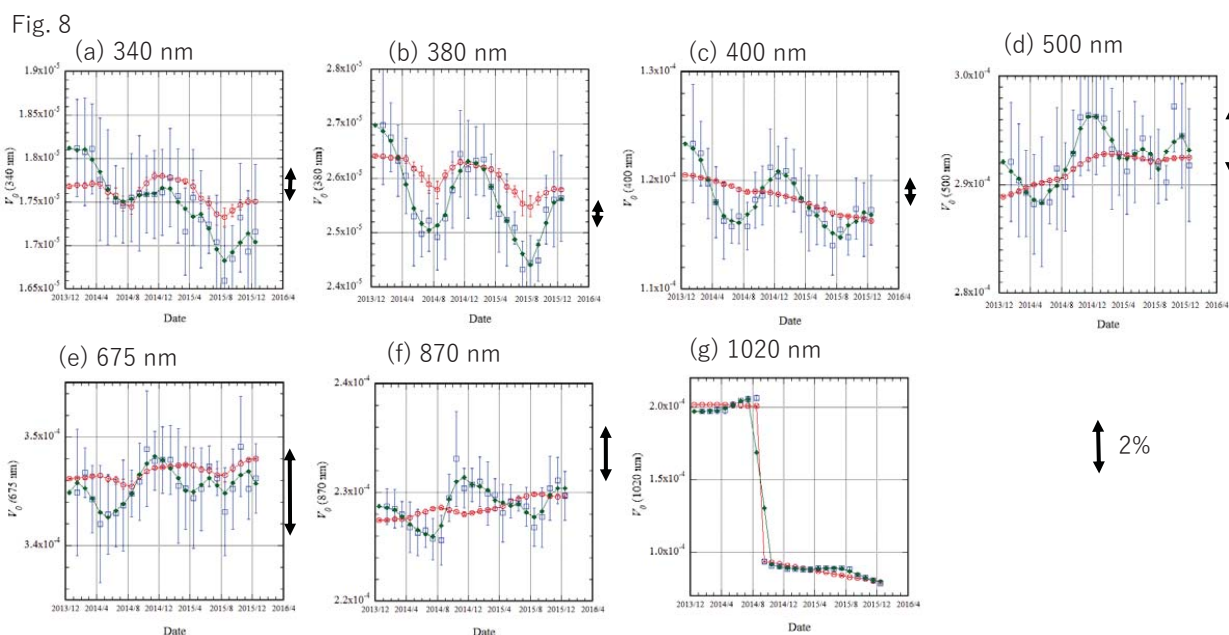


Fig. 8 Time series of calibration constant for POM-02 (Tsukuba) from January, 2014 to December, 2015. Blue open squares with error bar is calibration constant determined by IML method, Green line is 3 points running mean of IML, and Red line is calibration constant transferred from POM-02 (Calibration Reference).



Fig. 9

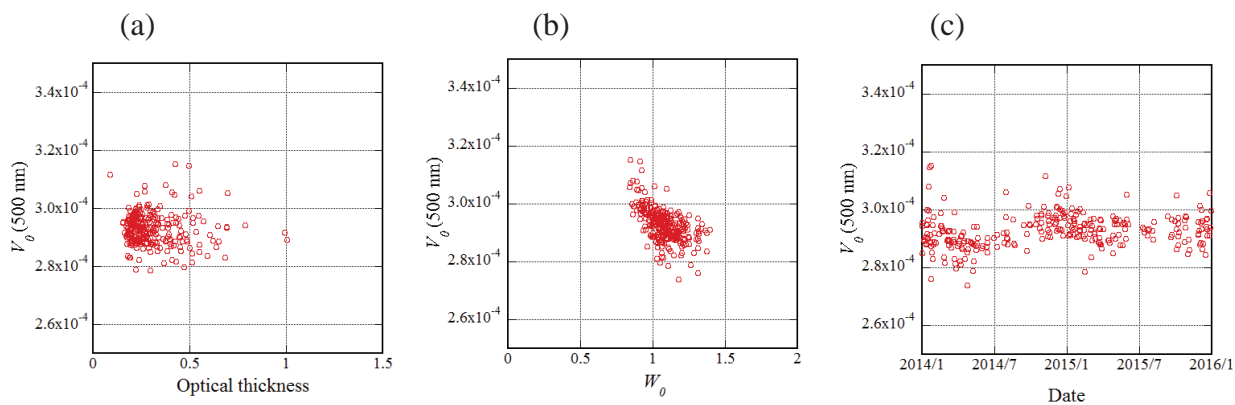


Fig. 9 (a) scatter plot of optical thickness at 500nm and V_{θ} by the IML method for 500nm channel, (b) scatter plot of W_{θ} and V_{θ} by the IML method for 500nm channel, (c) time series of V_{θ} by the IML method for 500nm channel in the period from January 2014 to December 2015 are shown.



Fig. 10

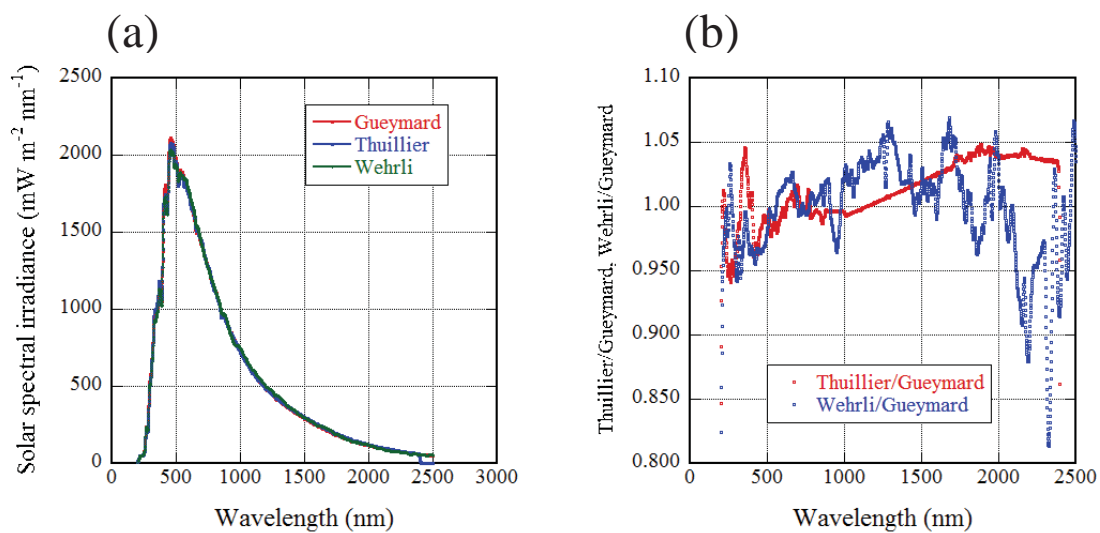


Fig. 10 (a) extra-terrestrial solar spectra, red line is Gueymard (2003), blue line is Thuillier et al (2003), and green line is Wehrli (1985), (b) the ratios of the solar spectrum to Gueymard (2003).



Fig. 11

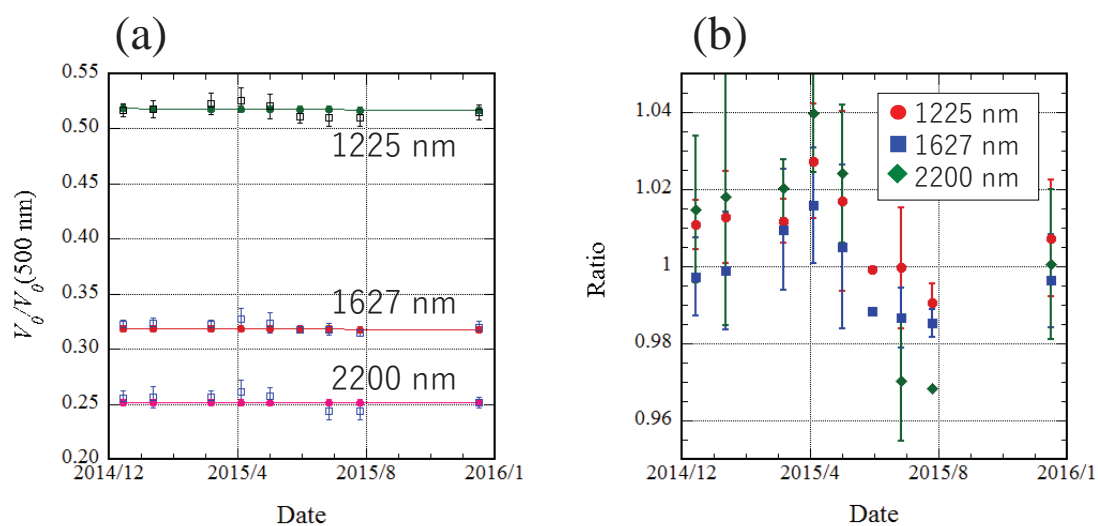


Fig. 11 (a) monthly mean of V_{02}/V_{01} and standard deviation, here $V_{01}=V_{\theta}(500 \text{ nm})$, (b) Ratio of V_{02} to the interpolated value of the calibration constant determined by Langley method, red symbols are 1225nm, blue ones are 1627nm, and green ones are 2200nm