Preparation of primary standard mixtures for atmospheric oxygen measurements with less than 1 μmol mol⁻¹ uncertainty for oxygen molar fractions

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10 Abstract. Precise monitoring of changes in atmospheric O_2 levels was implemented by preparing primary standard mixtures with less than 1 μ mol mol⁻¹ standard uncertainty for O₂ molar fractions. In this study, these mixtures were crafted in 10 L 11 12 high-pressure aluminium alloy cylinders using a gravimetric method in which unknown uncertainty factors were theoretically 13 determined and subsequently reduced. Molar fractions of the constituents (CO₂, Ar, O₂, and N₂) in the primary standard 14 mixtures were mainly resolved using masses of the respective source gases (CO_2 , Ar, O_2 , and N_2) that had been filled into the 15 cylinders. To precisely determine the masses of the source gases, the difference in mass of the cylinder before and after filling 16 the respective source gas was calculated by compared with an almost identical reference cylinder. Although the masses of the 17 cylinders filled with source gas with respect to the reference cylinder tended to deviate in relation to temperature differences 18 between the source gas filled cylinder and surrounding air, the degree of the deviation could be efficiently reduced by 19 measuring the two cylinders at the exact same temperature. The standard uncertainty for the cylinder mass obtained in our 20 weighing system was determined to be 0.82 mg. The standard uncertainties for the O₂ molar fractions in the primary standard 21 mixtures ranged from 0.7 μ mol mol⁻¹ to 0.8 μ mol mol⁻¹. Based on the primary standard mixtures, the annual average molar fractions of atmospheric O₂ and Ar in 2015 at Hateruma Island, Japan, were found to be 209339.1 \pm 1.1 µmol mol⁻¹ and 22 23 9334.4 \pm 0.7 µmol mol⁻¹, respectively. The molar fraction for atmospheric Ar was in agreement with previous reports.

24 **1 Introduction**

25 Observation of atmospheric O₂ molar fractions provides important information about the global carbon cycle (Keeling and 26 Shertz, 1992; Bender et al., 1996; Keeling et al., 1996, 1998a; Stephens et al., 1998; Battle et al., 2000; Manning and Keeling, 27 2006). For example, long-term observation allows the estimation of land biotics and oceanic CO₂ uptake (Manning and 28 Keeling, 2006; Tohjima et al., 2008; Ishidoya et al., 2012a, 2012b). Various measurement techniques have been developed 29 for this purpose, including the utilization of interferometry (Keeling et al., 1998b), mass spectrometry (Bender et al., 1994; 30 Ishidoya et al., 2003; Ishidoya and Murayama, 2014), a paramagnetic technique (Manning et al., 1999; Aoki et al., 2018; 31 Ishidoya et al., 2017), a vacuum-ultraviolet absorption technique (Stephens et al., 2003), gas chromatography (Tohjima, 2000), 32 and a method utilizing fuel cells (Stephens et al., 2007; Goto et al., 2013). In all these cases, calibration using standard mixtures

is required to precisely determine the relationship between the analysis output and O_2 molar fractions obtained.

Molar fractions of O_2 and Ar are commonly expressed as functions of the O_2/N_2 and Ar/N₂ ratios relative to an arbitrary

reference (Keeling and Shertz, 1992; Keeling et al., 2004) in per meg (one per meg is equal to 1×10^{-6}).

2
$$\delta(O_2/N_2) = \frac{[n(O_2)/n(N_2)]_{sam}}{[n(O_2)/n(N_2)]_{ref}} - 1$$
 (1)
3

4
$$\delta(\text{Ar/N}_2) = \frac{[n(\text{Ar})/n(N_2)]_{\text{sam}}}{[n(\text{Ar})/n(N_2)]_{\text{ref}}} - 1$$
 (2)

- 5
- 6
- 7

8 In these equations, n depicts the amount of each substance, and the subscripts "sam" and "ref" refer to sample air and reference 9 air, respectively. As the O_2 and Ar molar fractions in air are 20.946% and 0.943%, respectively, a respective change of 4.8 10 per meg and 107 per meg in $\delta(O_2/N_2)$ and $\delta(Ar/N_2)$ corresponds to a change of 1 µmol mol⁻¹ in molar fractions of O₂ and Ar. 11 Reported peak-to-peak amplitudes of seasonal cycles and trends in atmospheric $\delta(O_2/N_2)$ were within the range of 50 – 150 per meg (10 – 30 μ mol mol⁻¹ for O₂ molar fractions) and -20 per meg yr⁻¹ (-4 μ mol mol⁻¹ yr⁻¹ for O₂ molar fractions) 12 (Keeling et al., 1993; Battle et al., 2000; Van der Laan-Luijkx et al., 2013). To monitor these slight variations, the development 13 14 of primary standard mixtures with standard uncertainty of less than 5 per meg for O_2/N_2 ratios (1 µmol mol⁻¹ for O_2 molar 15 fractions) or less (Keeling et al., 1993; WMO, 2016) is required. In this study, the primary O₂ standard mixture with the 16 recommended uncertainty of 5 per meg (1 μ mol mol⁻¹) or less is hereafter referred to as a "highly precise O₂ standard mixture 17 (HPO)".

18 In general, standard mixtures need to be prepared in which molar fractions of the greenhouse gas species, such as CO_2 , CH_4 , 19 and N_2O , are stable enough during the observation period to enable monitoring of long-term changes in atmospheric molar 20 fractions of their species. For this purpose, it is indispensable to establish methods for determining absolute molar fractions 21 of greenhouse gases in the standard mixtures with required precision. Approved primary standard mixtures exist for CO₂, CH₄, 22 and N₂O, prepared by manometry (Zhao et al., 1997) or gravimetry (Tanaka et al., 1983; Matsueda et al., 2004; Dlugokencky 23 et al., 2005; Hall et al., 2007). However, preparing an HPO is challenging since it is necessary to prepare it with the relative 24 uncertainty of less than one-fifth of that for the CO₂ molar fraction in the CO₂ standard mixture. Since there is no common 25 scale for atmospheric O_2 observation, such as the ratio of O_2/N_2 determined using HPOs, each laboratory has employed 26 reference air determined using its own reference scale instead of a universal scale. This reference scale is determined based 27 on O₂/N₂ ratios in primary standard mixtures filled in high-pressure cylinders and is considered to be sufficiently stable during 28 the observation period (e.g. Keeling et al., 1998b; Tohjima et al., 2008; Ishidoya et al., 2012b). However, there are many 29 deterioration risks of the O₂/N₂ ratio in aluminium cylinders used for reference air and the primary standard mixtures. these 30 include fractionations of O₂ and N₂ induced by pressure, temperature and water vapor gradients (Keeling et al., 1998b), 31 adsorption/desorption of the constituents on the inner surface (Leuenberger et al., 2015), and permeation/leakage of the 32 constituents from/through the valve (Sturm et al., 2004; Keeling et al., 2007). In order to avoid these risks, the cylinders are 33 handled in accordance to certain best practices, including orienting cylinders horizontally to minimize thermal and 34 gravitational fractionation (Keeling, et al. 2007, Leuenberger, et al., 2015).

35 Although causes behind the fractionation should be sufficiently described by now, the effects of permeation and 36 adsorption/desorption have not been completely understood in the long-term scale. To enable comparison of O_2/N_2 values reported based on reference air with the difference scale directly, an independent development of primary standard mixtures with standard uncertainty of 5 per meg (1 μ mol mol⁻¹) or less is needed. In a pioneering study, Tohjima et al. (2005) first

- 3 prepared primary standard mixtures for the atmospheric O₂ measurement based on a gravimetric method. The standard
- 4 uncertainties for the O_2/N_2 ratio (the O_2 molar fraction) were noted at 15.5 per meg (2.9 µmol mol⁻¹), which was larger than
- 5 the required standard uncertainty of 5 per meg (1 μ mol mol⁻¹) or less. Since the 2.9 μ mol mol⁻¹ standard uncertainty recorded
- 6 by Tohjima et al. (2005) was significantly larger than the gravimetrically expected value of 1.6 μ mol mol⁻¹, it was suggested
- 7 that there are unknown factors exerting influence on the mass results of cylinders.
- 8 Taking these facts into consideration, in this study we set out to develop a new gravimetric method to prepare an HPO by 9 scientifically understanding the unknown uncertainty factors in the cylinder mass measurement. Our laboratory has built upon 10 the weighing system proposed by Matsumoto et al. (2004), in which gravimetry was used to prepare standard mixtures. 11 Although this system allows accurate mass measurements with a standard uncertainty of 2.6 mg, this proves insufficient to 12 prepare an HPO. A new mass comparator with better repeatability was recently introduced to the weighing system. In the 13 present study, we theoretically identified the unknown factors and presented an improved means of minimizing them. The 14 standard uncertainties for molar fractions of all constituents in HPOs, prepared using these improved means, are discussed. 15 Moreover, the molar fractions of all constituents in the HPOs were validated by comparing the gravimetric value with the 16 measured values of CO₂ mole fraction, Ar/N₂ ratio, and O₂/N₂ ratio. In order to validate the scale of O₂/N₂ ratio at the National 17 Institute of Advanced Industrial Science and Technology (AIST) determined using the HPOs prepared in this study, the annual 18 average of O₂/N₂ ratios in 2015 at Hateruma Island (24°03'N, 123°49'E, Japan) obtained from our measurements of air samples 19 were preliminarily compared with the annual average of O_2/N_2 ratios in 2015 at Hateruma Island on the NIES scale determined 20 by Tohjima et al.(NIES). Additionally, the molar fractions for atmospheric Ar and O₂ in air samples were determined using
- 21 the HPOs and compared with previously reported values.

22 2 Materials and Methods

23 **2.1 Preparation of the highly precise O₂ standard mixtures**

- Eleven HPOs were prepared in 10 L aluminium alloy cylinders (Luxfer Gas Cylinders, UK), with a diaphragm valve (G-55,
- Hamai Industries Limited, Japan) with poly(chlorotrifluoroethylene) (PCTFE) as sealant in accordance with ISO 6142-1:2015.
- 26 Pure CO₂ (>99.998 %, Nippon Ekitan Corporation, Japan), pure Ar (G1-Grade, 99.9999 %, Japan Fine Products, Japan), pure
- 27 O₂ (G1-Grade, 99.99995 %, Japan Fine Products, Japan), and pure N₂ (G1-Grade, 99.99995 %, Japan Fine Products, Japan)
- were used as source gases to prepare the HPOs. The value of δ^{13} C in pure CO₂ (which was adjusted to the atmospheric level)
- 29 was -8.92% relative to Vienna Pee Dee Belemnite (VPDB). Impurities in the source gases were identified and quantified
- 30 using gas chromatography with a thermal conductivity detector for N_2 , O_2 , CH_4 and H_2 in pure CO_2 , and gas chromatography
- 31 with a mass spectrometer for O_2 and Ar in pure N_2 and N_2 in pure O_2 . A Fourier transform infrared spectrometer was used for
- 32 detection of CO₂, CH₄ and CO in pure N₂, O₂, and Ar. A galvanic cell-type O₂ analyser was employed to quantify O₂ in pure
- 33 Ar. A capacitance-type moisture meter measured H₂O in pure CO₂, and a cavity ring-down-type moisture meter measured
- 34 H_2O in pure N_2 , O_2 and Ar.
- 35 Primarily, standard mixtures of CO_2 in Ar were prepared by combining pure CO_2 and pure Ar using a gravimetric method.
- 36 The molar ratios of CO₂ to Ar were close to the atmospheric molar ratio of CO₂ (400 μ mol mol⁻¹ or 420 μ mol mol⁻¹) to Ar

1 $(9340 \,\mu\text{mol mol}^{-1})$. The 10 L aluminium cylinder was used to prepare the HPO after evacuation by a turbomolecular pump.

- 2 The source gases were filled in the order of the mixtures of CO_2 in Ar, pure O_2 and pure N_2 in a filling room where the
- 3 temperature was controlled at 23 ± 1 °C and humidity was not regulated. The mass of the CO₂ in Ar standard mixture filled
- 4 was determined by the difference in the mass of the cylinder before and after filling with the mixture. The masses of the filled
- 5 pure O_2 and N_2 were treated in the same manner. The final pressure in the cylinder was 12 MPa, and masses of the individual
- gases were approximately 18 g of CO₂ in Ar standard mixture, 300 g of pure O₂, and 1000 g of pure N₂. 6

7 2.2 Weighing procedure for a cylinder

- 8 The masses obtained for the cylinders were determined using the same weighing system as reported by Matsumoto et al. 9 (2004), except for the mass comparator. The mass comparator used in the study of Matsumoto et al. was replaced with a new 10 mass comparator (XP26003L, Mettler Toledo, Switzerland), which had a maximum capacity of 26.1 kg, a sensitivity of 1 mg, 11 and a linearity of 20 mg. In this study, a cylinder whose mass was measured is hereafter referred to as a "sample cylinder". 12 Mass measurement for sample cylinders was performed in a weighing room where temperature and humidity were controlled 13 at 26 ± 0.5 °C and 48 ± 1 %, respectively. The temperature, humidity, and atmospheric pressure surrounding our weighing 14
- system were measured using a USB connectable logger (TR-73, T and D Corporation, Japan).
- 15 Mass measurements of the sample cylinder were conducted with respect to an almost identical reference cylinder aiming to 16 reduce any influence exerted by zero-point drifts, sensitivity issues associated with the mass comparator, changes in buoyancy 17 acting on the cylinder, and/or adsorption effects on the cylinder surface due to the presence of water vapor (Alink et al., 2000; 18 Milton et al., 2011). This is carried out according to several consecutive weighing operations in the ABBA order sequence, where "A" and "B" denote the reference cylinder and the sample cylinder, respectively. The process of loading and unloading 19 20 of the cylinders was automated. One complete cycle of the ABBA sequence required five minutes. The "mass reading" 21 recorded by our weighing system was given as the difference between both cylinders mass readings, which was obtained by 22 subtracting the reference cylinder reading from the sample cylinder reading.
- 23 Because the output of mass comparators is generally known to be nonlinear, there is a tendency to under- or over-estimate 24 mass readings for the sample cylinders in our weighing system. This is because the relation between the output of mass 25 comparators and the mass of artefacts tend to be different among various scale ranges. To reduce the influence of this 26 nonlinearity, sample cylinders were weighed only when the difference in readings between the sample and reference cylinders 27 was less than 500 mg. This was achieved by placing standard weights on a weighing pan alongside the sample or reference 28 cylinder. Any mass differences obtained in our weighing system took into account the masses and buoyancies of the standard 29 weights. The masses of the standard weights were traceable to the International System of Units. The standard uncertainties 30 of the masses were 0.25 mg, 0.045 mg, 0.028 mg, 0.022 mg, 0.018 mg, 0.014 mg, 0.011 mg, and 0.0090 mg for the 500 g,
- 31 100 g, 50 g, 20 g, 10 g, 5 g, 2 g, and 1 g weights, respectively.
- 32 Temperatures of the sample and reference cylinders were alternately measured by a thermocouple-type thermometer with a
- 33 resolution of 0.1 K (TX1001 digital thermometer, probe-90030, Yokogawa Test and Measurement Corporation, Tokyo, Japan)
- 34 before and after weighing of the cylinders.

1 2.3 Determination procedure of isotopic abundances for O and N

2 Each HPO was prepared using pure O₂ from two 48 L cylinders and pure N₂ from three or four 48 L cylinders as source gases. The isotopic abundances (^{16}O , ^{17}O , ^{18}O , ^{14}N , and ^{15}N) for pure O₂ and N₂ may be different between cylinders, resulting in 3 4 abundance differences among each HPO. The averaged values of isotopic abundances in pure O_2 (two cylinders) and pure N_2 5 (three or four cylinders) used for the respective HPOs were calculated based on the ratios of ¹⁸O/¹⁶O, ¹⁷O/¹⁶O, and ¹⁵N/¹⁴N in the HPOs. These were calculated using the equations, ${}^{18}\text{O}/{}^{16}\text{O} = [\delta({}^{18}\text{O}/{}^{16}\text{O}) + 1] \times ({}^{18}\text{O}/{}^{16}\text{O})_{\text{ref}}, {}^{17}\text{O}/{}^{16}\text{O} = [\delta({}^{17}\text{O}/{}^{16}\text{O}) + 1] \times ({}^{18}\text{O}/{}^{16}\text{O})_{\text{ref}}, {}^{18}\text{O}/{}^{16}\text{O})_{\text{ref}}, {}^$ 6 $({}^{17}O/{}^{16}O)_{ref}$, and ${}^{15}N/{}^{14}N = [\delta({}^{15}N/{}^{14}N) + 1] \times ({}^{15}N/{}^{14}N)_{ref}$. The terms $\delta({}^{17}O/{}^{16}O)$, $\delta({}^{18}O/{}^{16}O)$, and $\delta({}^{15}N/{}^{14}N)$ which were 7 determined by a mass spectrometer (Delta-V, Thermo Fisher Scientific Inc., USA) represent the deviation from the 8 corresponding atmospheric value (Ishidoya and Murayama, 2014). The isotopic ratios of $\delta({}^{17}O/{}^{16}O)$, $\delta({}^{18}O/{}^{16}O)$, and 9 $\delta(^{15}N/^{14}N)$ were approximately equal to those of $\delta(^{17}O^{16}O/^{16}O^{16}O)$, $\delta(^{18}O^{16}O/^{16}O^{16}O)$, and $\delta(^{15}N^{14}N/^{14}N)$, since 10 ${}^{17}O^{17}O^{/16}O^{16}O, \ {}^{18}O^{18}O^{/16}O^{16}O \ \text{and} \ {}^{15}N^{15}N^{/14}N^{14}N \ \text{tended to be much less than} \ {}^{17}O^{16}O^{/16}O^{16}O, \ {}^{18}O^{16}O^{/16}O^{16}O \ \text{and} \ {}^{15}N^{15}N^{14}N^{14}N \ \text{tended to be much less than} \ {}^{17}O^{16}O^{16}O^{16}O, \ {}^{18}O^{16}O^{/16}O^{16}O \ \text{and} \ {}^{15}N^{14}N^{14}N^{14}N \ \text{tended to be much less than} \ {}^{17}O^{16}O^{16}O^{16}O, \ {}^{18}O^{16}O^{/16}O^{16}O \ \text{and} \ {}^{15}N^{14}N^{14}N^{14}N \ \text{tended to be much less than} \ {}^{17}O^{16}O^{16}O^{16}O, \ {}^{18}O^{16}O^{/16}O^{16}O \ \text{and} \ {}^{15}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N^{14}N$ 11 $^{15}N^{14}N^{14}N^{14}N$. Values of $(^{18}O/^{16}O)_{ref}$, $(^{17}O/^{16}O)_{ref}$, and $(^{15}N/^{14}N)_{ref}$ refer to ratios of $^{18}O/^{16}O$, $^{17}O/^{16}O$, and $^{15}N/^{14}N$ in reference 12 13 air. In the present study, natural air in a 48 L aluminium cylinder (Cylinder No. CRC00045), equipped with a diaphragm valve 14 (G-55, Hamai Industries Limited, Japan) was used as reference air in the AIST scale (hereafter referred to AIST reference 15 air). The AIST reference air was prepared in 2011 by filling natural air into the cylinder to a pressure of about 13 MPa after drying cryogenically by using a highly efficient water trap at the dew point temperature lower than -80 °C (Ishidoya and 16 17 Murayama, 2014). It has been confirmed that the $\delta(O_2/N_2)$ on the AIST scale is higher by several hundred per meg than that 18 on the Scripps Institution of Oceanography (SIO) scale (e.g. Manning and Keeling, 2006) by reflecting the difference of the 19 filling years of the respective references for AIST and SIO scales (Ishidoya et al., 2016). The corresponding atmospheric values shown in Table 1 were used as the ratios of (¹⁷O/¹⁶O)_{ref}, (¹⁸O/¹⁶O)_{ref}, and (¹⁵N/¹⁴N)_{ref}, as isotopic abundances in the 20 21 troposphere are considered to be constant (Junk and Svec, 1958; Baertschi, 1976; Li et al., 1988; Barkan and Luz, 2005). 22 Because differences between isotopic ratios of N_2 , O_2 , Ar in the AIST reference air and air samples at Hateruma were 23 sufficiently small to be negligible, their fractionations due to preparation of the AIST reference air are ignored.

24 2.4 Analytical methods

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In this study, a mass spectrometer was used to determine O_2/N_2 and Ar/N_2 ratios in the HPOs. A cavity ring-down spectrometer was used to examine consistency among molar fractions of CO_2 in the HPOs. In this section, we describe the analytical

methods and relationships between the absolute O_2/N_2 (Ar/N₂) ratios and the mass-spectrometry based isotopic ratios.

28 2.4.1 Evaluations of O₂/N₂ and Ar/N₂ ratios in highly precise O₂ standard mixtures and natural air

Ratios of O_2/N_2 and Ar/N_2 in the HPOs were validated by comparison of gravimetrically calculated values with the measured values obtained by the mass spectrometer (Delta-V, Thermo Fisher Scientific Inc., USA). The mass spectrometer was adjusted to measure ion beam currents for masses 28 ($^{14}N^{14}N$), 29 ($^{15}N^{14}N$), 32 ($^{16}O^{16}O$), 33 ($^{17}O^{16}O$), 34 ($^{18}O^{16}O$), 36 (^{36}Ar), 40 (^{40}Ar), and 44 ($^{12}C^{16}O^{16}O$) simultaneously. Isotopic ratios of $\delta(^{15}N^{14}N)^{14}N$), $\delta(^{17}O^{16}O^{16}O)^{16}O)$, $\delta(^{18}O^{16}O/^{16}O^{16}O)$, $\delta(^{16}O^{16}O/^{14}N^{14}N)$, $\delta(^{36}Ar/^{40}Ar)$, and $\delta(^{40}Ar/^{14}N^{14}N)$ were determined against the AIST reference air using the mass spectrometer. In our prepared HPOs, the ratios of $\delta(O_2/N_2)_{HPO_{grav}}$ and $\delta(Ar/N_2)_{HPO_{grav}}$, comprised of all isotopes of O_2 , N_2 and Ar and gravimetrically calculated, are not equal to the isotopic ratios of $\delta(^{16}O^{16}O/^{14}N^{14}N)_{HPO_{meas}}$ and $\delta(^{40}\text{Ar}/^{14}\text{N})_{\text{HPO}_{meas}} \text{ measured by the mass spectrometer. This is because the isotopic ratios in source gases are different from the corresponding atmospheric values. The subscripts "HPO_grav" and "HPO_meas" hereafter refer to the gravimetric value and the measured value in the HPO. Thus, mass-spectrometry based isotopic ratios need to be converted to values equivalent to the <math>\delta(O_2/N_2)_{\text{HPO}_{grav}}$ ratio and the $\delta(Ar/N_2)_{\text{HPO}_{grav}}$ ratio. The values of $\delta(O_2/N_2)_{\text{HPO}_{meas}}$ and $\delta(Ar/N_2)_{\text{HPO}_{meas}}$ were calculated using mass-spectrometry based on isotopic ratios $^{15}\text{N}^{14}\text{N}/^{14}\text{N}^{14}\text{N}$, $^{17}\text{O}^{16}\text{O}/^{16}\text{O}$, $^{18}\text{O}^{16}\text{O}/^{16}\text{O}$, $^{36}\text{Ar}/^{40}\text{Ar}$, and $^{38}\text{Ar}/^{40}\text{Ar}$ as depicted in equations (3) and (4). Isotopic species of $^{17}\text{O}^{17}\text{O}$, $^{18}\text{O}^{17}\text{O}$, $^{18}\text{O}^{18}\text{O}$, $^{15}\text{N}^{15}\text{N}$, were negligible because the abundance of these species was very small.

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9
$$\delta(0_2/N_2)_{\text{HPO}_{meas}} = \left[\delta(^{16}O^{16}O/^{14}N^{14}N)_{\text{HPO}_{meas}} + 1 \right] \times \\ 10 \qquad \left[\frac{1 + {}^{17}O^{16}O/^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O}{1 + {}^{15}N^{14}N/^{14}N^{14}N} \right]_{\text{HPO}} / \left[\frac{1 + {}^{17}O^{16}O/^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16}O^{16$$

$$12 \qquad \delta(\text{Ar/N}_2)_{\text{HPO}_{\text{meas}}} = \left[\delta(^{40}\text{Ar}/^{14}\text{N}^{14}\text{N})_{\text{HPO}_{\text{meas}}} + 1\right] \times \left[\frac{1+^{36}\text{Ar}/^{40}\text{Ar}+^{38}\text{Ar}/^{40}\text{Ar}}{1+^{15}\text{N}^{14}\text{N}/^{14}\text{N}^{14}\text{N}}\right]_{\text{HPO}} / \left[\frac{1+^{36}\text{Ar}/^{40}\text{Ar}+^{38}\text{Ar}/^{40}\text{Ar}}{1+^{15}\text{N}^{14}\text{N}/^{14}\text{N}^{14}\text{N}}\right]_{\text{ref}} - 1$$

$$13 \qquad (4)$$

14

The values of ¹⁵N¹⁴N/¹⁴N, ¹⁷O¹⁶O/¹⁶O, and ¹⁸O¹⁶O/¹⁶O¹⁶O in the HPOs and the AIST reference air were calculated using 15 isotope abundances of O and N determined by the procedure described in section 2.3 (Table 1). The ³⁶Ar/⁴⁰Ar ratio of pure 16 Ar filled in the HPOs was calculated using equation ${}^{36}\text{Ar}/{}^{40}\text{Ar} = [\delta({}^{36}\text{Ar}/{}^{40}\text{Ar})_{\text{HPO meas}} + 1] \times ({}^{36}\text{Ar}/{}^{40}\text{Ar})_{\text{ref.}}$ The 17 $\delta(^{36}\text{Ar}^{/40}\text{Ar})_{\text{HPO meas}}$ value was determined by mass spectrometry of the HPOs. The $(^{36}\text{Ar}^{/40}\text{Ar})_{\text{ref}}$ value obtained was the 18 19 atmospheric value (36 Ar/ 40 Ar = 0.003349 ± 0.000004), because isotopic abundances of Ar in the AIST reference air were equal to that of the atmospheric value. The value of ³⁸Ar/⁴⁰Ar in the HPOs and the AIST reference air, which could not be 20 21 measured, was assumed to be ${}^{38}\text{Ar}/{}^{40}\text{Ar} = 0.000631 \pm 0.000004$ taken from previous reports as the atmospheric values. 22 Deviations of respective abundances of ³⁸Ar from the atmospheric value were considered to be less than the uncertainty of the 23 atmospheric value for ³⁸Ar. The atmospheric values of isotopic abundances for Ar were reported in an IUPAC technical report 24 (Böhlk, 2014).

- On the other hand, the absolute O_2/N_2 ratio in the AIST reference air was calculated by substituting the $(O_2/N_2)_{HPO_grav}$ in the
- HPOs and the $\delta(O_2/N_2)_{HPO_meas}$ for $(O_2/N_2)_{sam}$ and for $\delta(O_2/N_2)$ in equation (1). The absolute Ar/N₂ ratio in the AIST reference
- air was calculated in same manner (see the section 5.3).

28 2.4.2 Measurements of CO₂ in highly precise O₂ standard mixtures

- 29 Molar fractions of CO₂ in HPOs were verified using a cavity ring-down spectrometer (G2301, Picarro Inc., USA) equipped
- 30 with a multi-port valve (Valco Instruments Co. Inc., USA) for gas introduction and a mass flow controller (SEC-N112,
- 31 100SCCM, Horiba STEC, CO., Ltd, Japan). The cavity ring-down spectrometer was calibrated by three primary standard
- 32 gases $(364.50 \pm 0.14 \,\mu\text{mol mol}^{-1}, 494.04 \pm 0.14 \,\mu\text{mol mol}^{-1}, \text{and } 500.32 \pm 0.14 \,\mu\text{mol mol}^{-1})$ that had been prepared from pure
- 33 CO₂ and purified Air (G1 grade, Japan Fine Products, Japan) in accordance with ISO 6142-1:2015. The individual standard
- 34 gases were continuously measured for 30 minutes, respectively and the data last ten minutes were used. The pure CO₂ was
- 35 the same as the source gas used for preparation of the HPOs.

3 Identifying and minimizing unknown factors of uncertainty

As previously mentioned,, there are several unknown factors that influence mass readings obtained for sample cylinders. Identifying and minimizing these unknown factors is discussed in this section.

4 **3.1 Factors to cause deviations of mass readings**

5 Generally, mass readings of a sample cylinder obtained from a mass comparator tend to vary due to numerous factors such as 6 buoyancy, adsorption/desorption and thermal effects. The buoyancy effect is caused by changes in the density of the 7 surrounding air due to variations in ambient temperature, humidity, and pressure. Adsorption effect can greatly influence mass 8 readings of a sample cylinder by adsorption and desorption of water vapor from surrounding ambient air on the external 9 surface of a sample cylinder (Alink et al., 2000; Mizushima, 2004, 2007; Milton et al., 2011). The thermal effect is related to 10 temperature gradients between a sample cylinder and the surrounding ambient air (Gläser, 1990, 1999; Mana et al., 2002; 11 Gläser and Borys, 2009; Schreiber et al., 2015), which is able to change the weight force of the sample cylinder through 12 frictional forces exerted on the vertical surface of a sample cylinder and pressure forces on the horizontal surface. Both the 13 friction and pressure forces are caused by the upward or downward flow of air that is heated or cooled, respectively, by the 14 sample cylinder.

15 When the ABBA technique is employed for mass measurements under identical experimental conditions, the deviations of 16 the mass readings due to the factors described above become negligible because they are equally exerted on both the sample 17 and the reference cylinder. In fact, the buoyancy effect could be cancelled by adopting the ABBA technique in our mass 18 measurements (see Section 4.3.1). On the other hand, the identical experimental conditions tend to be disturbed by the 19 temperature change on the sample cylinder surface by adiabatic compression of the source gases and by the temperature 20 difference between the filling room and the weighing room. Mass readings of the sample cylinder deviate from true values 21 when thermal effects due to a change in the sample cylinder surface temperature are exerted independently and at varying 22 degrees on the sample and reference cylinders. Moreover, the amount of water adsorbed on the sample cylinder surface can 23 also be influenced by humidity if the level in the filling room is different from that in the weighing room. This non-uniformity 24 of temperatures and the water amount between the sample cylinder surface and the surrounding ambient air is assumed to be 25 the main contributor of uncertainties in obtained mass readings of the sample cylinder (Matsumoto et al., 2008). In order to 26 identify and minimize the contribution to the non-uniformity, we examined the equilibrium of both humidity and temperature 27 for the surface of the sample cylinder used in this study before carrying out any measurement.

28 **3.2 Identifying and minimizing unknown uncertainty factors**

Equilibrium in the temperature and the water amount between the sample cylinder surface and its surrounding ambient air is considered to be achieved by placing the sample cylinder on our weighing system for an appropriate duration of time before the mass reading. Here, the equilibrium between the reference cylinder surface and its surrounding ambient air is always achieved as the reference cylinder is permanently left on the weighing system. The equilibrium for the sample cylinder is easily disturbed by the processes of its evacuation and filling of the source gases. To quantify the appropriate time interval needed to restore equilibrium, the mass readings of the sample cylinder were recorded after the evacuation and after the filling. These values were plotted against the time elapsed after evacuation and filling (Figure 1). The surface temperature of the 1 sample cylinder recorded after the evacuation was 2 K lower, while the temperatures recorded after the filling for CO_2 in Ar 2 standard mixture, pure O_2 , and pure N_2 were -0.7 K, 1 K, and 6 K higher than that of the reference cylinder, respectively. In 3 this experiment, the equilibria were considered to be achieved when the standard deviation of the mass readings remained 4 constant for two or more hours with the repeatability value of < 0.82 mg (see in Section 4.3.1.). Interestingly, the mass readings 5 recorded after the evacuation and filling for CO_2 in Ar mixture tended to decrease as time elapsed, while those after filling 6 with pure O₂ and N₂ gases tended to increase. Deviations in mass readings had some connection with the temperature

7 difference between the reference and sample cylinders. The results imply that warmer cylinders appear lighter.

8 Appropriate time intervals were defined as the time elapsed from the evacuation or the filling time to the point of re-achieving 9 equilibrium. This time interval was noted as 5 h after complete evacuation. Time intervals required after filling with relevant 10 gases were different depending on the filled gas species to some extent. For the CO_2 in Ar mixture, equilibria were achieved 11 in 3 h to 5 h, while 4 h to 5 h were required for O_2 equilibration and 7 h to 9 h for N_2 . These intervals indicate that preparation

12 of a single HPO requires several days. To determine the mass of the sample cylinder in as short time as possible, a clear 13 indicator for carrying out mass measurement is needed.

14 As described above, the deviations in mass readings are considered to have some relation with the temperature differences in 15 the sample and reference cylinders. Therefore, we proceeded to examine this relationship to understand whether the 16 temperature difference can be the indicator. The closed squares shown in Figure 2 indicate that the deviation of mass readings 17 was proportional to temperature differences. This deviation rate was determined to be -14.3 mg K^{-1} . Although the results 18 indicate that a temperature difference of 0.1 K causes a deviation of 1.4 mg, the deviation in the recorded mass readings 19 ensures the repeatability value of 0.82 mg that is achieved by reducing the temperature difference to below 0.06 K. By 20 conducting measurements of the cylinder temperature using the thermocouple-type thermometer and ensuring that the mass 21 readings were taken when the temperature of both cylinders were the same, we were able to reduce the deviation contributing 22 to the mass readings.

23 To validate the proposed weighing procedure, the reproducibility of mass readings obtained after disturbing the equilibrium 24 was measured. Hence, the mass reading sequence after a cooling or heating cycle of the cylinders was examined. Figure 3 25 illustrates the results in which four heating cycles (number 1 to 4) and four cooling cycles (number 5 to 8) were conducted. 26 In this experiment, temperatures of the cooled or heated cylinder were 1 K to 3 K lower or 10 K to 20 K higher than that of 27 the reference cylinder, respectively. When mass readings were recorded after ensuring equal temperatures of both the sample 28 and reference cylinders, no difference in mass readings recorded after the cooling and heating cycles was detected. The 29 reproducibility of mass readings was estimated to be 0.44 mg with regards to the standard deviation of the mass readings 30 shown in Figure 3. The fact that reproducibility was lower than the repeatability value validated the weighing procedure. The

31 contributions to mass readings by non-equilibrium conditions were negligible using the established weighing procedure.

32 It is difficult to determine whether deviations in mass readings recorded for sample cylinders were caused by thermal or 33 adsorption effects simply by analysing these results. This is because both effects are related to temperature fluctuations.

34 However, the thermal effect influenced the slope of the calibration function solely through temperature differences, whereas

35 the adsorption/desorption effect influenced the slope of the calibration function via a combination of both ambient temperature

- 36
- and humidity. This is because the adsorbed or desorbed amounts of water on the surface of both cylinders are highly dependent
- 37 on the cylinder temperature and humidity of the surrounding ambient air. To determine which of these effects contributed the

most to the deviations, the relationship between the deviations and temperature differences was investigated under various conditions in the weighing room. Humidity was stringently controlled at 30%, 50%, 65%, and 80%, whereas temperature levels were maintained at 22 °C, 26 °C, and 29 °C. As shown in Figure 2, the slope did not depend on the humidity nor temperature. These results indicate that the dominant factor of deviations in mass readings was rather an effect of thermal gradients than adsorption, because the deviations depended on the temperature difference only. Therefore, we focused on minimizing the impact of any thermal gradient in further experiments.

7 4 Evaluation of uncertainty factors for the O₂ Standard Mixtures

In this section, we discuss any uncertainty factors associated with molar fractions of constituents in the HPOs. The gravimetric molar fraction (y_k) of the constituent k (CO₂, Ar, O₂, and N₂) was calculated using the molar mass (M_i) and the molar fraction $(x_{i,j})$ of the constituent i (CO₂, Ar, O₂, N₂ and impurities) in the filled source gas j (CO₂ in Ar standard mixture, pure O₂, and pure N₂). Additionally, mass (m_j) of the source gases filled into the sample cylinder were incorporated into equation (5) in accordance with ISO 6142-1:2015.

13

14
$$y_k = \frac{\sum_{j=1}^{r} \left(\frac{x_{k,j} \times m_j}{\sum_{i=1}^{q} x_{i,j} \times M_i} \right)}{\sum_{j=1}^{r} \left(\frac{m_j}{\sum_{i=1}^{q} x_{i,j} \times M_i} \right)}$$
 (5)

1	5
т	J

In this equation, *r* and *q* represent the number of source gases *j* and constituents *i*, respectively, while $x_{k,j}$ is the molar fraction of the constituent *k* in the source gas *j*. Uncertainties $(u(y_k))$ associated with the gravimetric molar fraction were calculated according to the law of propagation.

19

$$20 u^2(y_k) = \sum_{j=1}^r \sum_{i=1}^q \left(\frac{\partial y_k}{\partial x_{i,j}}\right)^2 \times u^2(x_{i,j}) + \sum_{i=1}^q \left(\frac{\partial y_k}{\partial M_i}\right)^2 \times u^2(M_i) + \sum_{j=1}^r \left(\frac{\partial y_k}{\partial m_j}\right)^2 \times u^2(m_j) (6)$$

21

22 In this equation, u(A) depicts the standard uncertainty for A. Gravimetric molar fractions of the constituent k and its associated 23 uncertainty in the molar fractions for the HPOs prepared in this study were calculated using equation (5) and equation (6), 24 and they are listed in Table 2. The standard uncertainties for the constituents N₂, O₂, Ar, and CO₂ were $0.8 - 1.0 \,\mu$ mol mol⁻¹, 25 $0.7 - 0.8 \ \mu mol \ mol^{-1}$, $0.6 - 0.7 \ \mu mol \ mol^{-1}$, and $0.03 \ \mu mol \ mol^{-1}$, respectively. Table 3 lists the contribution of the purity of 26 the source gases, molar masses of the constituents, and masses of the source gases to the gravimetric molar fraction. These 27 correspond to the square root of the first, second, and third terms found in equation (6), respectively. Uncertainty factors in 28 the gravimetric molar fractions in the HPOs were mainly those of the mass for the source gases filled into the sample cylinder. 29 Contributions from other sources of uncertainty were negligible. Purity of the source gases and molar masses of the 30 constituents i, as well as the masses of the source gases and their associated standard uncertainties, are described in Sections 31 4.1, 4.2, and 4.3.

1 **4.1 Purity of source gas**

2 Pure O₂, N₂, Ar, and CO₂ were used as source gases to prepare the HPOs. Molar fractions of impurities presented in source

3 gases and their associated standard uncertainties were determined based on the primary standard gases prepared in accordance 4 with ISO 6142-1:2015. When the molar fraction of impurity *h* was under its detection limit (L_h), the molar fraction (x_h) and

5 standard uncertainty $(u(x_{h,j}))$ of h in the source gas j was calculated using the equations $x_{h,j} = L_{h,j}/2$ and $(x_{h,j}) = L_{h,j}/2$

6 $L_{h,i}/2\sqrt{3}$. The calculated values for the impurities and purities of the source gases are listed in Table 4.

7 **4.2 Molar masses of constituents**

- 8 Molar masses (M_i) of the constituents were calculated using most recent atomic masses and isotopic abundances reported by 9 the IUPAC. However, IUPAC values for atomic masses of O and N have large standard uncertainties because they reflect 10 variability present in individual isotopic abundances of natural terrestrial matter. Using IUPAC values, standard uncertainties for N₂ and O₂ molar fractions in the HPOs were calculated to be 4 μ mol mol⁻¹. In addition, atmospheric values of their isotopic 11 12 abundances could not be used for calculating molar masses of the source gases even though pure O_2 and N_2 were produced 13 from air. This is because isotopically abundant O and N in pure O_2 and N_2 tends to deviate from the corresponding atmospheric 14 value during the production process. Therefore, the isotopic abundances of O and N in HPOs were precisely determined using 15 mass spectrometry. Their isotopic abundance and associated standards of uncertainty are shown in Table 1. The difference of
- 16 isotopic abundances of O and N between respective HPOs were negligible.
- 17 Based on the isotopic abundances obtained in this study and the atomic masses of O and N in the filled source gases, pure O₂ 18 and N2 were determined with relative standard uncertainties of 0.000029 % and 0.000006 %, respectively. It was shown that 19 the uncertainty in molar masses is negligible (Table 3). Although the grade and supplier of the pure O_2 and N_2 used in this study were the same as those of the source gases used by Tohiima et al. (2005), the atomic masses (15.999366 (1) for O and 20 21 14.006717 (4) for N) obtained for the two elements were different from Tohjima et al. reported values (15.999481 (8) for O and 14.006677 (4) for N). These differences resulted in a deviation of 0.4 μ mol mol⁻¹ and 1.2 μ mol mol⁻¹ for O₂ and N₂, 22 23 respectively. Since these results infer that the ratios of O and N isotopes change with production time, the isotopic abundances 24 of O and N in the source gases have to be precisely determined whenever HPOs are prepared using different pure O_2 and N_2 . 25 On the other hand, standard uncertainties in the atomic mass presented in an IUPAC technical report by De Laeter et al. (2003) 26 were sufficient for further use in the case of Ar and CO₂ as source gases.

27 **4.3 Determining the masses of the filled gases**

Masses of individual gases that were filled into the sample cylinders were calculated using the mass difference before and after the filling. The standard uncertainty of the obtained mass was calculated by combining standard uncertainties of mass readings of the sample cylinder before and after filling each gas. To determine uncertainty in mass reading of the sample cylinder, three factors were evaluated i.e., the repeatability $u(m_{rep})$ of the mass readings, permeation $u(m_{gas \ permeate})$ of the source gases during weighing, and buoyancy change $u(m_{buoyancy})$ due to the expansion of the cylinder. The standard uncertainties $(u(m_{cyl}))$ were defined according to equation (7).

1
$$u^{2}(m_{cyl}) = u^{2}(m_{rep}) + u^{2}(m_{gas \, permeate}) + u^{2}(m_{buoyancy}).$$
 (7)

2

These factors are discussed in detail in sections 4.3.1, 4.3.2, and 4.3.3. The standard uncertainties of masses for the individual
filled gases were 1.2 mg regardless of the gas species.

5 4.3.1 Repeatability of mass readings

6 The repeatability of mass readings was evaluated by continuous mass measurement of sample cylinders using the ABBA 7 technique over three days. This is because preparation of a single HPO requires three days. Mass readings were recorded after 8 the sample cylinder was left on the weighing system for at least a week. Air density was likewise measured for three days by 9 carefully monitoring temperature, humidity, and pressure changes in ambient air (Figure 4). Our findings indicate that the 10 mass readings remain stable during the three-day experiment. The standard deviation of mass readings (0.82 mg) is 11 represented as repeatability $u(m_{rep})$. The fact that the mass readings were not affected by changes in the air density also 12 indicates that buoyancy issues influencing the sample cylinder were cancelled out by changes simultaneously affecting the 13 reference cylinder.

14 **4.3.2 Permeation of source gases during weighing**

15 All of the cylinders used in this study have diaphragm valves, which were joined to the cylinders via pipe fittings and sealed 16 with Teflon tape. The seal of diaphragm valves was made from PCTFE, through which gases tend to permeate quite slowly 17 (Sturm, 2004). Since permeation of the source gases during weighing the sample cylinders resulted in evaluation error of the 18 masses for the source gases, we examined the permeability of purified air by monitoring mass of the sample cylinder filled 19 with purified air at a pressure of 8 MPa. The changes in mass readings were measured for over four months. From these results, 20 it was determined that the permeability was 0.013 mg day⁻¹. This effect was considered to be negligible because it is much lower than the repeatability. Therefore, the contribution of permeability $(u(m_{gas \ permeate}))$ to the standard uncertainty 21 22 calculations $(u(m_{cyl}))$ was ignored. On the other hand, the amount of air permeating from the sample cylinder during the 23 course of a year was calculated to be about 4.7 mg. This quantity may cause changes in the composition of the HPO if the 24 mixture is kept for extended periods of time, since the gas permeability depends on the gas species (Sturm, 2004).

25 4.3.3 Buoyancy effect of cylinder expansion

Oh et al. (2013) reported that volume in the 10 L aluminium cylinders linearly increases with increase in internal pressure, and volume expansion was determined to be 24 ± 2 ml when the pressure difference in the sample cylinders was 12 MPa. Tohjima et al. (2005) likewise reported a volume expansion of 22 ± 4 ml when the pressure difference was 10 MPa. In this study, we adopted a volume expansion of the sample cylinders as 55 ± 5 ml, measured by a cylinder supplier, when the pressure difference was 25 MPa. Compared to the expansion rate with respect to pressure variations reported by Oh (2.0 ± 0.2 ml MPa⁻¹) (2013) and Tohjima (2.2 ± 0.4 ml MPa⁻¹) (2005), expansion rate of the sample cylinders used in this study was

- determined to be 2.2 ± 0.2 ml MPa⁻¹. The pressure difference recorded before and after filling with source gases was 0.12
- 33 MPa, 2.5 MPa, and 9.4 MPa for CO₂ in Ar standard mixture, pure O₂, and pure N₂, respectively. These pressure differences
- 34 were subsequently used to calculate buoyancy effects, which were reported as 0.3 mg, 6.4 mg, and 23.9 mg for CO_2 in Ar

1 standard mixture, pure O₂, and pure N₂, respectively. In turn, these buoyancy effects caused changes in the gravimetric molar

2 fraction of $+0.5 \,\mu\text{mol mol}^{-1}$ and $-0.5 \,\mu\text{mol mol}^{-1}$ for O₂ and N₂, respectively. The final mass readings were corrected to take

3 these changes into account. The standard uncertainties $u(m_{buoyancy})$ in linear expansion were considered to be negligible.

4 5 Validation of the constituents in the highly precise O₂ standard mixtures

5 The O₂ molar fractions in the HPOs deviate from gravimetric values with deviation of the molar fractions of other constituents. 6 In this section, the molar fractions of all constituents were validated. The gravimetric and measured values for the CO₂ molar 7 fractions were compared, along with the values of $\delta(Ar/N_2)_{HPO_grav}$, $\delta(Ar/N_2)_{HPO_meas}$, $\delta(O_2/N_2)_{HPO_grav}$ and $\delta(O_2/N_2)_{HPO_meas}$ to 8 validate the molar fractions of Ar, O₂ and N₂ in the HPOs. Table 5 shows the $\delta(O_2/N_2)_{HPO_meas}$ and $\delta(Ar/N_2)_{HPO_meas}$ values 9 calculated using equation (3) and (4), as well as the values for $\delta(^{15}N^{14}N/^{14}N^{14}N)_{HPO_meas}$, $\delta(^{17}O^{16}O/^{16}O)_{HPO_meas}$, 10 $\delta(^{18}O^{16}O/^{16}O^{16}O)_{HPO_meas}$, $\delta(^{16}O^{16}O/^{14}N^{14}N)_{HPO_meas}$, $\delta(^{36}Ar/^{40}Ar)_{HPO_meas}$, and $\delta(^{38}Ar/^{40}Ar)_{HPO_meas}$.

11 5.1 Determining the absolute (O₂/N₂) and (Ar/N₂) ratios in AIST reference air

12 The absolute $(O_2/N_2)_{HPO_grav}$ and $(Ar/N_2)_{HPO_grav}$ ratios, as well as the $\delta(O_2/N_2)_{HPO_meas}$ and $\delta(Ar/N_2)_{HPO_meas}$ values of the HPOs 13 are listed in Table 2 and Table 5. Using these values, the absolute ratios for O_2/N_2 and Ar/N_2 in AIST reference air were 14 0.2680929 ± 0.0000016 and 0.0119542 ± 0.0000009 , respectively. On the AIST scale, these values corresponded to $\delta(O_2/N_2)$ 15 = 0 and $\delta(Ar/N_2) = 0$. Associated standard uncertainties were determined with regards to the law of propagation of uncertainty. 16 The HPOs prepared in this study did not include minor components of Ne, He, Kr, CH₄, H₂ and N₂O which are present in air samples. Therefore, the variation of the $\delta({}^{15}N{}^{14}N{}^{14}N{}^{14}N)$, $\delta({}^{17}O{}^{16}O{}^{16}O{}^{16}O)$, $\delta({}^{18}O{}^{16}O{}^{16}O{}^{16}O)$, $\delta({}^{16}O{}^{16}O{}^{14}N{}^{14}N)$, $\delta({}^{36}Ar{}^{40}Ar)$, 17 and $\delta({}^{40}Ar/{}^{14}N^{14}N)$ values obtained by the measurement of mass spectrometry was examined when molar fractions of Ne 18 19 change from 0 μ mol mol⁻¹ to 500 μ mol mol⁻¹. Consequently, the isotopic ratios did not change significantly depending on 20 Ne. Since abundance of Ne is highest in minor components, the changes of the molar fractions of other minor components 21 also might fail to affect the isotopic ratios measured in this study.

22 5.2 CO₂ molar fractions and Ar/N₂ ratio

- Three primary standard gases were used to measure CO_2 molar fractions in HPOs. Table 2 illustrates the gravimetric and measured values and associated standard uncertainties. The CO_2 molar fractions in the cylinder labelled CPB28679, which had been prepared on 29 March 2017, were not measured. Differences between the gravimetric and measured values (obtained
- by subtracting the measured value from gravimetric value) were found to range from $-0.17 \,\mu\text{mol mol}^{-1}$ to 0.03 $\mu\text{mol mol}^{-1}$.
- 27 The gravimetric values were in line with the measured values, both being within the accepted levels of uncertainty.
- From these results, mass of the CO₂ in Ar standard mixture which we evaluated were considered validated. Figure 5 shows
- 29 the plot of the $\delta(Ar/N_2)_{HPO_meas}$ values relative to the $\delta(Ar/N_2)_{HPO_grav}$ values, as well as the residuals of the $\delta(Ar/N_2)_{HPO_meas}$
- 30 values that had been estimated using the best fitted line by the least squares method. The standard deviation of the residuals
- 31 was 78 per meg. This standard deviation represents a scatter in the $(Ar/N_2)_{HPO_{grav}}$ values, since the standard uncertainty for
- 32 $\delta(Ar/N_2)_{HPO_meas}$ was much smaller than the obtained standard deviation (Ishidoya and Murayama, 2014). The standard
- 33 uncertainties for $\delta(Ar/N_2)_{HPO_{grav}}$ values ranged from 65 per meg to 77 per meg. Standard uncertainties were comparable to

1 the standard deviations obtained for the residuals, supporting the validity of uncertainty calculations for the constituents, Ar 2 and N₂.

5.3 O₂/N₂ ratio 3

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Figure 5 shows a plot of the $\delta(O_2/N_2)_{HPO_meas}$ values listed in Table 5 against the $\delta(O_2/N_2)_{HPO_grav}$ values listed in Table 2, as well as residuals from the fitted line by the least squares method. The slope of the fitted line was determined to be $1.00162 \pm$ 0.00029, which indicated that the discrepancy between change rates of the $\delta(O_2/N_2)_{HPO}$ meas values and the $\delta(O_2/N_2)_{HPO}$ grav values was within 0.16%. The standard deviation of the residuals was 3.6 per meg, which proved to be in line with the standard uncertainties for the corresponding gravimetric values since the standard uncertainties for the values of $\delta(O_2/N_2)_{HPO}$ grav ranged from 3.2 per meg to 4.0 per meg. The agreement with the gravimetric and measured values reinforced the idea that the method 10 for calculating the uncertainties of the constituents, O_2 and N_2 , was proper and accurate. On the other hand, the $\delta(O_2/N_2)_{HPO_meas}$ values were lower than their $\delta({}^{16}O{}^{16}O{}^{14}N{}^{14}N)_{HPO}$ meas counterparts by 18.2 per meg to 27.1 per meg (Table 5) and differences were larger than the standard uncertainties obtained for both values. This means that the $\delta(O_2/N_2)$ values in the HPOs should

be expressed by $\delta(O_2/N_2)_{HPO}$ meas rather than $\delta({}^{16}O^{16}O/{}^{14}N^{14}N)_{HPO}$ meas. 13

14 The O_2/N_2 ratio of gases delivered from the cylinders may differ from the gravimetric O_2/N_2 ratio by either homogeneous or

15 inhomogeneous fractionation (Leuenberger et al., 2015 Langenfelds et al., 2005 Keeling et al., 2004). In this study, we used

16 the same type of values and cylinders as in the study of Tohjima et al. (2005). Tohjima et al. examined changes in the O_2/N_2

ratio of the HPOs by releasing the inner air into a room at a flow rate of 8 mL min⁻¹ and found that the fractionation of O_2 and 17

18 N_2 during air release was negligible. Therefore, we chose not to evaluate fractionation in this study.

19 **6** Comparison with Previous Values

20 To confirm the consistency of the results obtained using the HPOs, we preliminarily compared O_2/N_2 ratios on both the AIST

21 and NIES scale using annual average of $\delta(O_2/N_2)$ values in the air sample from Hateruma Island collected from January to

22 December in 2015. Additionally, the molar fraction of atmospheric O₂ and Ar were determined based on the HPOs and then

23 compared with previously reported values to confirm consistency of the results.

24 6.1 Comparison between O₂/N₂ ratios on the AIST and NIES scales

25 We observed the atmospheric $\delta(O_2/N_2)$ by analysing air samples collected at Hateruma Island from January to December in 2015. For the air samples, we confirmed that the isotopic ratios of N_2 and O_2 did not differ significantly from the atmospheric 26 27 values, such that we regard the $\delta({}^{16}O/{}^{16}N)$ value measured by the mass spectrometer is equivalent to $\delta(O_2/N_2)$ in 28 equation (1). Twice a month, the air samples were collected in two Pyrex glasses arranged in series (one for AIST and the 29 other for NIES). Using these air samples, we determined that the annual average of $\delta(O_2/N_2)$ in 2015 on the AIST scale was 30 -63 ± 3 per meg. The number following the symbol \pm denotes the standard uncertainty of the measurement. Using equation 31 (1), the $\delta(O_2/N_2)$ value on the AIST scale was then converted to the absolute O_2/N_2 ratio by utilizing the absolute ratio of 32 O_2/N_2 in AIST reference air determined in section 5.1. In 2015, the absolute O_2/N_2 ratio on Hateruma Island was 0.2680761 33 ± 0.0000018 . This value can be converted to the corresponding $\delta(O_2/N_2)$ value on the NIES scale using equation (1), since the absolute O_2/N_2 ratio in NIES reference air was reported to be 0.2681708 by Tohjima et al. (2005). The converted $\delta(O_2/N_2)$ 34

- 1 value, which is expressed as $\delta(O_2/N_2)_{\text{NIES on AIST}}$ in the following descriptions, was found to be -353 ± 6 per meg on the NIES
- 2 scale. The uncertainty expresses the 95% confidence interval.
- 3 On the other hand, the annual average of $\delta(O_2/N_2)$ in Hateruma island in 2015 measured by NIES on the NIES scale was -395
- \pm 5 per meg (Tohjima, Y., personal communications). The number following the symbol \pm denotes the standard uncertainty
- 5 of the measurement (Tohjima et al., 2008). The $\delta(O_2/N_2)$ value is expressed as $\delta(O_2/N_2)_{\text{NIES on NIES}}$. There was a difference of
- 6 41 per meg between both values of $\delta(O_2/N_2)_{\text{NIES on AIST}}$ and $\delta(O_2/N_2)_{\text{NIES on NIES}}$. This difference falls outside of the 95%
- 7 confidence interval. The disagreement between $\delta(O_2/N_2)_{\text{NIES on AIST}}$ and $\delta(O_2/N_2)_{\text{NIES on NIES}}$ suggests that there are some
- 8 inconsistencies between the gravimetric methods developed by Tohjima et al. (2005) and this study. Additionally, other
- 9 sources of error may exist (e.g. difference between instruments, sampling method, introduction method). Therefore, a direct
- 10 comparison of the O_2/N_2 ratio or the O_2 molar fraction between the AIST and NIES scales, by a round-robin experiment of
- 11 the HPOs developed in this study, is required some time in the near future to quantify the differences in absolute values and
- 12 span of each gravimetric scale.

13 6.2 Determination of atmospheric O₂ and Ar molar fractions and comparison with previous data

The molar fractions for atmospheric O₂ and Ar were determined based on $\delta(O_2/N_2)$ and $\delta(Ar/N_2)$ values for air samples taken at Hateruma Island in 2015. The $\delta(O_2/N_2)$ and $\delta(Ar/N_2)$ values were -62.8 per meg and -62.8 per meg, respectively. Using the $(O_2/N_2)_{ref}$ and $(Ar/N_2)_{ref}$ ratios for AIST reference air, the $\delta(O_2/N_2)$ and $\delta(Ar/N_2)$ values were used to calculate the O_2/N_2 and Ar/N_2 ratios using equation (1) and (2). In 2015, the calculated O_2/N_2 and Ar/N_2 ratios for samples from Hateruma Island were 0.2680761 ± 0.0000018 and 0.0119534 ± 0.0000009, respectively. The molar fractions of O₂ and Ar (x_{O_2} and x_{Ar}) were calculated using the aforementioned O_2/N_2 and Ar/N_2 ratios by using the equations below.

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21
$$x_{0_2} = K \times \frac{O_2/N_2}{(1+O_2/N_2 + Ar/N_2)}$$
 (8)

22
$$x_{Ar} = K \times \frac{Ar/N_2}{(1+O_2/N_2 + Ar/N_2)}$$
 (9)

24 In these two equations, K is the sum of N₂, O₂, and Ar molar fractions in the air samples and was estimated to be 999567.8 \pm 25 0.1 μ mol mol⁻¹. To obtain this value, the molar fractions of Ne (18.18 μ mol mol⁻¹), He (5.24 μ mol mol⁻¹), CH₄ (1.82 μ mol 26 mol⁻¹), Kr (1.14 μ mol mol⁻¹), H₂ (0.52 μ mol mol⁻¹), N₂O (0.32 μ mol mol⁻¹), CO (0.15 μ mol mol⁻¹) and Xe (0.09 μ mol mol⁻¹) 27 reported by Tohjima et al. (2005) and CO₂ (404.7 μ mol mol⁻¹) in 2015 were used. The CO₂ molar fraction was the average 28 CO_2 molar fraction measured using a mass spectrometer. The calculated O_2 and Ar molar fractions were 209339.1 ± 1.1 µmol 29 mol^{-1} and 9334.4 \pm 0.7 $\mu mol mol^{-1}$, respectively. Standard uncertainties were estimated in accordance with the law of 30 propagation of uncertainties. In 2000, Tohjima et al. (2005) reported the atmospheric Ar molar fraction of $9333.2 \pm 2.1 \mu mol$ 31 mol⁻¹ (2005), whereas the value reported for air samples collected on Korea's Anmyeon Island in 2002 and at Niwot Ridge 32 in 2001 was $9332 \pm 3 \mu$ mol mol⁻¹ (Park et al., 2004). Hence, our values for atmospheric Ar were in line previous reports.

²³

1 7 Conclusion

In this study, we demonstrated that the deviation of mass readings of the sample cylinders is susceptible to temperature differences between the sample and reference cylinders. The contribution degree of the temperature difference was -14.3 mg K^{-1} . Our results also indicate that variations in mass readings due to temperature difference could be reduced to negligible levels by obtaining mass readings at thermal equilibrium. A long time is required to reach thermal equilibrium. Since the variations mainly depend on temperature differences rather than factors related to the adsorption phenomena (e.g., the temperature of the cylinder and/or the humidity of ambient air), it was thus concluded that the variations in mass readings were influenced solely by thermal effects.

- We developed a preparation technique for the production of HPOs with atmospheric levels of CO₂, Ar, O₂, and N₂. To determine the O₂ molar fractions with standard uncertainties of less than 1 μ mol mol⁻¹, repeatability in measuring the mass readings of the sample cylinders was determined to be 0.82 mg. The impact of leakage or permeation of the source gases through the cylinder valve, as well as change in buoyancy such as the expansion of the cylinder as a factor of the cylinder inner pressure was evaluated. Additionally, the molar masses of the source gases, pure O₂ and pure N₂, were determined based
- on the abundance of their isotopes. The gravimetrically calculated standard uncertainties were in good agreement with the standard deviation for the corresponding measured values. This indicates that the uncertainty calculations of the gravimetric values for constituents performed in this study were accurate and valid.
- Based on the HPOs prepared in this study, we determined molar fractions of atmospheric Ar and O_2 at Hateruma Island in 2015. These values were 9334.4 \pm 0.7 μ mol mol⁻¹ and 209339.1 \pm 1.1 μ mol mol⁻¹, for Ar and O₂, respectively. The
- 19 atmospheric Ar molar fraction was in line with the values reported by Park (9332 \pm 3 µmol mol⁻¹) and Tohjima (9333.2 \pm 2.1

20 μ mol mol⁻¹) (Park et al., 2004; Tohjima et al., 2005).

21

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- 9

1 Table 1. Isotopic composition and atomic masses of pure oxygen and pure nitrogen used to prepare highly precise O₂ standard

2 mixtures (HPOs).

3

Isotone	Atomic mass ^{a,b}	Isotope abundance		
isotope	Atomic mass	Atmosphere ^a	Source gas ^a	Isotopic ratio of source gas ^e
^{14}N	14.0030740074(18)	0.996337(4)°	0.996346(4)	
¹⁵ N	15.000108973(12)	0.003663(4) ^c	0.003654(4)	$\delta^{15}N = (-2.397 \pm 0.001) \ \text{\%}$
¹⁶ O	15.9949146223(25)	0.9975684(9) ^d	0.9975887(9)	
¹⁷ O	16.99913150(22)	0.0003836(8) ^d	0.0003818(8)	$\delta^{17}O = (-4.66 \pm 0.05) \ \text{\%}$
¹⁸ O	17.9991604(9)	0.0020481(5) ^d	0.0020295(5)	$\delta^{18}O = (-9.075 \pm 0.003) \ \text{\%}$
Sources	Atomi	c mass of nitrogen	a Atomic	e mass of oxygen ^a
Atmosphere	14.006	6726(4)	15.999405(1)	
Source gases	14.006	5717(4)	15.999	366(1)

4

5 ^a Numbers in the parentheses represent the standard uncertainty in the last digits.

^b Atomic mass and the standard uncertainty as determined by De Laeter et al. (2003).

^c Abundance of the isotope and the standard uncertainty as determined using calculations for the absolute ¹⁵N/¹⁴N ratio
 obtained by Junk and Svec (1958).

^d Abundance of the isotope and the standard uncertainty were calculated using ${}^{17}O/{}^{16}O = 12.08$ ‰ and ${}^{18}O/{}^{16}O = 23.88$ ‰ vs.

10 the VSMOW as determined by Barkan and Luz (2005). The absolute isotopic ratio for VSMOW and the standard uncertainty

11 were determined by Li et al. (1988) for ${}^{17}\text{O}/{}^{16}\text{O}$ and Baertschi (1976) for ${}^{18}\text{O}/{}^{16}\text{O}$.

¹² ^e Isotopic ratio is defined as the difference in the corresponding atmospheric value (AIST reference air) measured using a

13 mass spectrometer. Numbers following the symbol \pm denote the standard uncertainty.





Figure 1 Changes in mass readings of sample cylinders plotted against the time elapsed after evacuation of the cylinder and filling of source gases. The mass readings were obtained using the weighing system. Deviation of the mass reading is expressed as the change in amount from the equilibrium value, which was defined as the mass reading when the standard deviation of the values remained constant for two or more hours.





Figure 2 Changes in the mass readings observed for sample cylinders plotted against temperature differences obtained under various conditions (a temperature range from 22 °C to 29 °C, a humidity range from 30% to 80%.)







Figure 3 Reproducibility of mass readings obtained for the sample cylinder after cylinders had been heated at 40 °C (numbers 1 to 4) or cooled at 23 °C (numbers 5 to 8). The error bars represent the standard uncertainty.

Table 2. Gravimetric values of N_2 , O_2 , and CO_2 molar fractions alongside $(O_2/N_2)_{HPO_grav}$, $(Ar/N_2)_{HPO_grav}$, $\delta(O_2/N_2)_{HPO_grav}$, and $\delta(Ar/N_2)_{HPO_grav}$, as well as the measured values of CO₂ molar fractions in HPOs.

Cylinder number	Preparation date	Gravimetric values ^a , µmol mol ⁻¹							Measured values, μmol mol ⁻¹	
		N2	O ₂	Ar	CO ₂	(O ₂ /N ₂) _{HPO_grav}	(Ar/N ₂) _{HPO_grav}	$\delta(O_2/N_2)_{HPO_grav} \ ^b$	$\delta(Ar/N_2)_{HPO_grav} \ ^b$	CO ₂
CPC00556	15 March 2017	780094.1 ± 1.0	210068.3 ± 0.8	9415.2 ± 0.7	422.30 ± 0.03	0.2692858 ± 0.0000011	0.0120693 ± 0.0000009	4471.8 ± 4.0	9619 ± 77	422.37 ± 0.14
CPB28679	29 March 2017	782593.9 ± 0.8	207770.2 ± 0.7	9222.1 ± 0.6	413.64 ± 0.03	0.2654892 ± 0.0000009	0.0117841 ± 0.0000008	-9689.9 ± 3.4	-14244 ± 67	-
CPB16178	5 April 2017	779014.8 ± 1.0	211348.4 ± 0.8	9241.0 ± 0.7	395.78 ± 0.03	0.2713021 ± 0.0000010	0.0118624 ± 0.0000009	11993.0 ± 4.0	-7694 ± 77	395.96 ± 0.14
CPB16345	7 April 2017	781499.3 ± 1.0	208750.7 ± 0.8	9349.6 ± 0.7	400.43 ± 0.03	0.2671156 ± 0.0000011	0.0119636 ± 0.0000009	-3623.2 ± 4.0	777 ± 75	400.40 ± 0.14
CPB16315	12 April 2017	781264.1 ± 0.9	209040.6 ± 0.7	9297.0 ± 0.7	398.18 ± 0.03	0.2675671 ± 0.0000010	0.0118999 ± 0.0000009	-2595.1 ± 3.6	-5191 ± 79	398.21 ± 0.14
CPB16379	17 April 2017	781059.5 ± 0.8	209233.2 ± 0.7	9308.6 ± 0.6	398.68 ± 0.03	0.2678838 ± 0.0000009	0.0119179 ± 0.0000008	-757.9 ± 3.3	-3050 ± 65	398.68 ± 0.14
CPB16349	13 June 2017	780424.7 ± 0.8	209813.5 ± 0.7	9342.7 ± 0.6	419.06 ± 0.03	0.2688452 ± 0.0000009	0.0119713 ± 0.0000008	2828.5 ± 3.4	1419 ± 66	419.22 ± 0.14
CPB28912	15 June 2017	780792.3 ± 0.8	209437.0 ± 0.7	9351.1 ± 0.6	419.44 ± 0.03	0.2682366 ± 0.0000009	0.0119765 ± 0.0000008	558.1 ± 3.4	1851 ± 66	419.54 ± 0.14
CPB28679	22 June 2017	780869.0 ± 0.8	209383.9 ± 0.7	9328.6 ± 0.6	418.44 ± 0.03	0.2681421 ± 0.0000009	0.0119464 ± 0.0000008	205.8 ± 3.4	-664 ± 65	418.54 ± 0.14

^a Numbers following the symbol \pm denote the standard uncertainty.

^b Values were calculated using the absolute O₂/N₂ and Ar/N₂ in AIST reference air as described in section 5.1.

	Uncertainty so	Combined standard		
	Mass of	uncertainty,		
Constituent	source gas ^a			µmol mol ⁻¹
N ₂	0.77	0.11	0.05	0.77
O ₂	0.63	0.03	0.03	0.63
Ar	0.56	0.13	0.02	0.58
CO_2	0.025	0.006	0.011	0.028

Table 3. Typical contribution of each source of uncertainty (including the mass of the source gas, molar mass, and purity) to the standard uncertainties obtained for the molar fractions of N_2 , O_2 , Ar, and CO_2 in the HPO.

^a The values were calculated in the procedure described in section 4.3.

^b The values were calculated in the procedure described in section 4.2.

^c The values were calculated in the procedure described in section 4.1.

Impurity	Source gases, μ mol mol ⁻¹					
Impully	CO ₂	Ar	O ₂	N_2		
N ₂	0.9 ± 0.5	0.12 ± 0.07	0.12 ± 0.07	-		
O_2	0.3 ± 0.1	0.5 ± 0.3	-	0.05 ± 0.03		
				0.05 ± 0.03		
Ar	-	-	0.089 ± 0.052	0.28 ± 0.01		
				0.32 ± 0.03		
CO_2	-	0.002 ± 0.001	0.124 ± 0.004	0.002 ± 0.001		
H_2O	4.8 ± 2.7	0.05 ± 0.03	0.05 ± 0.03	0.05 ± 0.03		
CH_4	0.6 ± 0.3	0.005 ± 0.003	0.005 ± 0.003	0.005 ± 0.003		
CO	-	0.04 ± 0.02	0.04 ± 0.02	0.04 ± 0.02		
H_2	2.2 ± 1.3	-	-	-		
				999999.8 ± 0.1		
Purity	999991.3 ±3.1	999999.3 ± 0.3	999999.6 ± 0.1	999999.6 ± 0.1		
				999999.5 ± 0.1		

Table 4. Impurities in source gases for preparation of HPOs

Numbers following the symbol \pm denote the standard uncertainty.

"-" represents the constituents which were not measured.



Figure 4 Repeatability of mass readings obtained for the sample cylinders and ambient air density for three days. Solid and dashed lines represent mass readings and ambient air density, respectively

Table 5. $\delta(^{15}N^{14}N/^{14}N^{14}N)_{HPO_meas}$, $\delta(^{17}O^{16}O/^{16}O)_{HPO_meas}$, $\delta(^{18}O^{16}O/^{16}O)_{HPO_meas}$, $\delta(^{16}O^{16}O/^{14}N^{14}N)_{HPO_meas}$, $\delta(^{36}Ar/^{40}Ar)_{HPO_meas}$ and $\delta(^{40}Ar/^{14}N^{14}N)_{HPO_meas}$ and $\delta(^{40}Ar/^{14}N^{14}N^{14}N)_{HPO_meas}$ and $\delta(^{40}Ar/^{14}N^{14}N^{14}N)$

Cylinder number	$\delta(^{15}N^{14}N/^{14}N^{14}N)_{HPO_meas}$	$\delta(^{17}O^{16}O/^{16}O^{16}O)_{HPO_meas}$	$\delta(^{18}O^{16}O/^{16}O^{16}O)_{HPO_meas}$	$\delta({}^{16}O{}^{16}O{}^{/14}N{}^{14}N{})_{HPO_meas}$	$\delta(O_2/N_2)_{HPO_meas}$	$\begin{array}{l} \delta(O_2/N_2)_{HPO_meas} \\ - \ \delta(^{16}O^{16}O/^{14}N^{14}N) \end{array}$	$\delta(^{36}Ar/^{40}Ar)_{HPO_meas}$	$\delta(^{40}Ar/^{14}N_2)_{HPO_meas}$	$\delta(Ar/N_2)_{HPO_meas}$
CPC00556	-2365.0 ± 1.2	-4032 ± 50	-7907.8 ± 2.6	4477.5 ± 3.2	4459.2 ± 3.2	-18.2	-2465 ± 50	9649.0 ± 6.5	9658.1 ± 6.5
CPB28679	-2343.5 ± 1.2	-4032 ± 50	-8298.0 ± 2.6	-9704.7 ± 3.2	-9724.4 ± 3.2	-19.7	-1969 ± 50	-14102.6 ± 6.5	-14092.2 ± 6.5
CPB16178	-2372.5 ± 1.2	-4219 ± 50	-8279.7 ± 2.6	12011.7 ± 3.2	11991.7 ± 3.2	-20.0	-2197 ± 50	-7828.0 ± 6.5	-7818.1 ± 6.5
CPB16345	-2351.5 ± 1.2	-4676 ± 50	-9087.6 ± 2.6	-3624.2 ± 3.2	-3647.7 ± 3.2	-23.5	-2311 ± 50	712.0 ± 6.5	721.5 ± 6.5
CPB16315	-2356.2 ± 1.2	-4665 ± 50	-9069.6 ± 2.6	-1946.8 ± 3.2	-1970.2 ± 3.2	-23.4	-2228 ± 50	-4538.2 ± 6.5	-4528.5 ± 6.5
CPB16379	-2416.8 ± 1.2	-4655 ± 50	-9062.8 ± 2.6	-763.6 ± 3.2	-786.6 ± 3.2	-22.9	-2261 ± 50	-3074.4 ± 6.5	-3064.3 ± 6.5
CPB16349	-2407.9 ± 1.2	-4630 ± 50	-9036.0 ± 2.6	2833.1 ± 3.2	2810.2 ± 3.2	-23.0	-2360 ± 50	1485.7 ± 6.5	1495.4 ± 6.5
CPB28912	-2397.2 ± 1.2	-4656 ± 50	-9075.3 ± 2.6	554.6 ± 3.2	531.5 ± 3.2	-23.2	-2348 ± 50	1812.2 ± 6.5	1821.9 ± 6.5
CPB28679	-2390.8 ± 1.2	-5109 ± 50	-9941.2 ± 2.6	212.5 ± 3.2	185.4 ± 3.2	-27.1	-2338 ± 50	-642.8 ± 6.5	-633.2 ± 6.5

These values are on the AIST scale, i.e., determined against AIST reference air and are given in per meg.

Numbers following the symbol \pm denote the standard uncertainty.



Figure 5 Relationship between $\delta(O_2/N_2)_{HPO_grav}$ and $\delta(Ar/N_2)_{HPO_meas}$ on the AIST scale (upper). Fitting residuals $\delta(O_2/N_2)_{HPO_meas}$ and $\delta(Ar/N_2)_{HPO_meas}$ are likewise shown (lower).