



1                                   **Spectral Intensity Bioaerosol Sensor (SIBS):**  
2                                   **A new Instrument for Spectrally Resolved Fluorescence Detection**  
3                                   **of Single Particles in Real-Time**

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28                                   formance Evaluation, Polystyrene Latex Spheres, Biofluorophores



29 **Abstract**

30 Primary biological aerosol particles (PBAP) in the atmosphere are highly relevant for the Earth sys-  
31 tem, climate, and public health. The analysis of PBAP, however, remains challenging due to their  
32 high diversity and large spatiotemporal variability. For real-time PBAP analysis, light-induced fluo-  
33 rescence (LIF) instruments have been developed and widely used in laboratory and ambient studies.  
34 The interpretation of fluorescence data from these instruments, however, is often limited by a lack of  
35 spectroscopic information. This study introduces a new instrument – the Spectral Intensity Bioaero-  
36 sol Sensor (SIBS) – that resolves fluorescence spectra for single particles and, thus, promises to ex-  
37 pand the scope of fluorescent PBAP quantification and classification.

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39 The SIBS shares key design components with the latest versions of the Wideband Integrated Bioaero-  
40 sol Sensor (WIBS) and the findings presented here are also relevant for the widely deployed WIBS-  
41 4A and WIBS-NEO as well as other LIF instruments. The key features of the SIBS and findings of  
42 this study can be summarized as follows:

- 43 - Particle sizing yields reproducible linear responses for particles in the range of 300 nm to 20  $\mu\text{m}$ .  
44 The lower sizing limit is significantly smaller than for earlier commercial LIF instruments (e.g.,  
45 WIBS-4A and the Ultraviolet Aerodynamic Particle Sizer (UV-APS)), expanding the analytical  
46 scope into the accumulation mode size range.
- 47 - Fluorescence spectra are recorded for two excitation wavelengths ( $\lambda_{\text{ex}} = 285$  and 370 nm) and a  
48 wide range of emission wavelengths ( $\lambda_{\text{mean}} = 302 - 721$  nm) with a resolution of 16 detection  
49 channels, which is higher than for most other commercially available LIF bioaerosol sensors.
- 50 - Fluorescence spectra obtained for 16 reference compounds confirm that the SIBS provides suffi-  
51 cient spectral resolution to distinguish major modes of molecular fluorescence. For example, the  
52 SIBS resolves the spectral difference between bacteriochlorophyll and chlorophyll *a/b*.
- 53 - A spectral correction of the instrument-specific detector response is essential to use the full fluo-  
54 rescence emission range.
- 55 - Asymmetry factor (AF) data were assessed and were found to provide only limited analytical  
56 information.
- 57 - In test measurements with ambient air, the SIBS worked reliably and yielded characteristically  
58 different spectra for single particles in the coarse mode with an overall fluorescent particle frac-  
59 tion of ~4 % ( $3\sigma$  threshold), which is consistent with earlier studies in comparable environments.



## 60 1. Introduction

61 Aerosol particles are omnipresent in the atmosphere, where they are involved in many environmental  
62 and biogeochemical processes (e.g., Baron & Willeke, 2001; Després et al., 2012; Fuzzi et al., 2006;  
63 Hinds, 1999; Pöschl, 2005; Pöschl & Shiraiwa, 2015). Primary biological aerosol particles (PBAP),  
64 also termed bioaerosols, represent a diverse group of airborne particles, consisting of whole or frag-  
65 mented organisms including, e.g., bacteria, viruses, archaea, algae, and reproductive units (pollen  
66 and fungal spores), as well as decaying biomass (e.g., Deepak & Vali, 1991; Després et al., 2012;  
67 Fröhlich-Nowoisky et al., 2016; Jaenicke, 2005; Madelin, 1994; Pöschl, 2005) and can span sizes  
68 from few nanometers up to  $\sim 100 \mu\text{m}$  (Hinds, 1999; Schmauss and Wigand, 1929). The increasing  
69 awareness of the importance of PBAP regarding aerosol-cloud interactions, health aspects, and  
70 spread of organisms on local, continental or even intercontinental scales has led to a growing interest  
71 by scientific researchers and the public (e.g., Després et al., 2012; Fröhlich-Nowoisky et al., 2016).

72 Due to inherent limitations (e.g., poor time resolution and costly laboratory analyses) of tradi-  
73 tional off-line techniques (e.g., light microscopy and cultivation-based methods) for PBAP quantifi-  
74 cation, several types of real-time techniques have been developed within the last several decades to  
75 provide higher time resolution and lower user costs (e.g., Caruana, 2011; Després et al., 2012;  
76 Fennelly et al., 2017; Ho, 2002; Jonsson & Tjærnhage, 2014; Sodeau & O'Connor, 2016). One prom-  
77 ising category of real-time instruments – meaning that particles are sampled and analyzed both in-  
78 stantly and autonomously – involves application of light- induced fluorescence (LIF). The main prin-  
79 ciple of this technique is the detection of intrinsic fluorescence from fluorophores ubiquitous in bio-  
80 logical cells, such as those airborne within PBAP. These fluorophores include a long list of biological  
81 molecules such as aromatic amino acids (e.g., tryptophan and tyrosine), co-enzymes (e.g., reduced  
82 pyridine nucleotides (NAD(P)H)), flavin compounds (e.g., riboflavin), as well as biopolymers (e.g.,  
83 cellulose and chitin) and chlorophyll (e.g., Hill et al., 2009; Li et al., 1991; Pan et al., 2010; Pöhlker  
84 et al., 2012, 2013). Detailed information of biological fluorophores can be found elsewhere (Pöhlker  
85 et al., 2012 and references therein).

86 Today, commercial on-line LIF instruments such as the Ultraviolet Aerodynamic Particle Sizer  
87 (UV-APS, TSI Inc. Shoreview, MN, USA) and the Wideband Integrated Bioaerosol Sensor (WIBS,  
88 developed at the University of Hertfordshire, U.K. and currently licensed and manufactured by Drop-  
89 let Measurement Technologies (DMT, Longmont, CO, USA)) are commonly applied for research  
90 purposes. Detailed descriptions of the UV-APS (e.g., Agranovski et al., 2003; Brosseau et al., 2000;  
91 Hairston et al., 1997) and the WIBS series (e.g., Foot et al., 2008; Kaye et al., 2000, 2005; Stanley et  
92 al., 2011) are given elsewhere. Concisely, the UV-APS uses an  $\lambda_{\text{ex}} = 355 \text{ nm}$  laser excitation source  
93 and spans an emission range between  $\lambda_{\text{em}} = 420\text{--}575 \text{ nm}$ . In contrast, the WIBS applies two pulsed  
94 xenon flash lamps emitting at  $\lambda_{\text{ex}} = 280$  and  $370 \text{ nm}$ , whereas fluorescence emission is detected in



95 three detection channels,  $\lambda_{em}=310-400$  nm (at  $\lambda_{ex}=280$  nm) and  $\lambda_{em}=420-650$  nm (at  $\lambda_{ex}=280$  and  
96 370 nm). Both instruments provide spectrally unresolved fluorescence information. The latest WIBS  
97 model is currently the WIBS-NEO, whose design is based on a WIBS-4A but with an extended par-  
98 ticle size detection range between  $\sim 300$  nm and  $100\ \mu\text{m}$  (nominal). Both UV-APS and WIBS models  
99 have been examined in a variety of laboratory validations (e.g., Agranovski et al., 2003, 2004;  
100 Brosseau et al., 2000; Healy et al., 2012; Hernandez et al., 2016; Kanaani et al., 2007; O'Connor et  
101 al., 2013; Saari et al., 2013, 2014; Savage et al., 2017; Toprak & Schnaiter, 2013) and have been  
102 deployed to investigate both indoor and outdoor atmospheric aerosol via longer-term measurements  
103 (e.g., Bhangar et al., 2014; Crawford et al., 2015; Fernández-Rodríguez et al., 2018; Foot et al., 2008;  
104 Gabey et al., 2010, 2013; Gosselin et al., 2016; Healy et al., 2014; Huffman et al., 2010, 2012, 2013;  
105 Perring et al., 2015; Schumacher et al., 2013; Twohy et al., 2016; Ziemba et al., 2016).

106 Although LIF instruments do not offer the same ability to qualitatively identify sampled particles  
107 as, e.g., off-line microscopy, mass spectrometry, or culture-based methods, they provide size-re-  
108 solved information as well as fast sampling and fine-scale temporal information for single particles  
109 not accessible with off-line techniques. Nevertheless, these instruments present significant chal-  
110 lenges. For example, quantification of PBAP by LIF instruments is hindered by the fact that some  
111 biological materials reveal weak fluorescence characteristics that does not rise above detection  
112 thresholds (Huffman et al., 2012). In addition to this complication, the detection threshold is not a  
113 universally defined parameter and varies for each channel between different units of the same type  
114 of instruments (e.g., Hernandez et al., 2016; Savage et al., 2017). Furthermore, unambiguous spec-  
115 troscopic characterization of bioparticles is fundamentally challenging, because fluorescence spectra  
116 of even individual molecules in condensed matter are relatively broad due to radiative decay path-  
117 ways of excited electrons. Further, bioparticles are chemically complex, each comprised of a mixture  
118 of dozens of types of fluorophores that can each emit a unique emission spectrum that smears together  
119 into an even broader fluorescence spectrum from each particle (Hill et al., 2009, 2015; Pan, 2015).  
120 Another difficulty is that many non-biological particles, such as certain mineral dusts and polycyclic  
121 aromatic hydrocarbons (PAHs), may fluoresce, making it more difficult to distinguish patterns aris-  
122 ing from biological particles (e.g., Pöhlker et al., 2012 and references therein; Savage et al., 2017).  
123 Lastly, most currently available commercial LIF instrumentation are limited to recording data in 1-3  
124 spectrally integrated emission channels, which limits the interpretation of fluorescence information.  
125 Recent efforts to apply more complex clustering algorithms to the spectrally unresolved WIBS-type  
126 data are proving helpful at adding additional discrimination, but aerosol characterization using in-  
127 strumentation with such low spectra resolution is likely to be fundamentally limited (e.g., Robinson  
128 et al., 2013; Ruske et al., 2017; Savage & Huffman, 2018).



129 The evolution of LIF techniques over the last several decades has significantly expanded our  
130 knowledge on spatiotemporal patterns of PBAP abundance in the atmosphere. Nevertheless to further  
131 improve the applicability of LIF instrumentation to widespread PBAP detection, it is necessary both  
132 to design LIF instrumentation with adequate instrumental properties (e.g., high spectral resolution)  
133 and to standardize their operation by characterizing instruments thoroughly with known standards  
134 (Robinson et al., 2017). Working toward this goal, a number of LIF instruments that offer analysis  
135 of single bioparticles by providing resolved fluorescence spectra have been developed (e.g., Hill et  
136 al., 1999; Pan et al., 2003, 2010; Pinnick et al., 2004; Ruske et al., 2017), however relatively little  
137 has been done to offer these commercially. Examples for commercially available instruments provid-  
138 ing resolved fluorescence spectra are the PA-300 ( $\lambda_{\text{ex}}=337$  nm;  $\lambda_{\text{em}}=390-600$  nm, 32 fluorescence  
139 detection channels) (Crouzy et al., 2016; Kiselev et al., 2011, 2013) and the follow-up model Rapid-  
140 E ( $\lambda_{\text{ex}}=337$  nm;  $\lambda_{\text{em}}=350-800$  nm, 32 fluorescence detection channels) (<http://www.plair.ch/>), both  
141 manufactured by Plair SA, Geneva, Switzerland. Beside of resolved fluorescence detection, both  
142 instruments also provide measurements of the decay of fluorescence signals, also referred to as flu-  
143 orescence lifetime.

144 Introduced here is a new instrument for the detection and characterization of individual particles;  
145 the Spectral Intensity Bioaerosol Sensor (SIBS, Droplet Measurement Technologies). Technical  
146 properties of the instrument are described in detail and its performance is validated with sizing and  
147 fluorescence particle standards, as well as with ambient air. Due to the dual excitation and spectrally  
148 resolved fluorescence in combination with a broad size detection range, the SIBS has the potential to  
149 increase the selectivity of fluorescent biological and non-biological particle detection and discrimi-  
150 nation. Because the SIBS uses a comparable optical system as the WIBS-4A and WIBS-NEO, tech-  
151 nical details presented here are broadly important to a growing community of scientists investigating  
152 both indoor and outdoor aerosol. Insights and data presented will thus contribute to ongoing discus-  
153 sions within the community of LIF users and will also stimulate discussions about needs for future  
154 instrument improvements.

155

## 156 2. Materials and methods

### 157 2.1 Chemicals and materials

158 Supplemental table S1 summarizes 19 polystyrene latex spheres (PSLs, 5 doped with fluorescent  
159 dye) and 6 polystyrene divinylbenzene (PS-DVB) particles, which were purchased from Thermo  
160 Fisher (Waltham, MA, USA), Bangs Laboratories Inc. (Fishers, IN, USA), Duke Scientific Corp.  
161 (Palo Alto, CA, USA), and Polysciences Inc. (Warrington, PA, USA). A detailed study regarding  
162 steady-state fluorescence properties of PSLs and PS-DVB particles used within this study can be



163 found in Könemann et al. (2018). Additionally, we analyzed particles comprised separately of seven  
164 pure biofluorophores (tyrosine, tryptophan, NAD, riboflavin, chlorophyll *a* and *b*, and bacteriochloro-  
165 rophyll) (Table S2) as well as one microorganism (*Saccharomyces cerevisiae*; baker's yeast, bought  
166 at a local supermarket). Table S2 also includes reference particles used for asymmetry measurements,  
167 namely iron oxide (Fe<sub>3</sub>O<sub>4</sub>), carbon nanotubes, and ammonium sulfate. Ultrapure water (MilliQ, 18  
168 MΩ) and ≥ 99.8 % ethanol (CAS Nr. 64-17-5, Carl Roth GmbH und Co. KG, Karlsruhe, Germany)  
169 were used as solvents.

170

## 171 2.2 Aerosolization of reference particles

172 PSLs were aerosolized from aqueous suspensions with a portable aerosol generator (AG-100; DMT).  
173 For both fluorescent and non-fluorescent PSLs, one drop of the suspension (or alternatively three  
174 drops for 3 and 4 μm PSLs) was diluted into 10 ml ultrapure water in plastic medical nebulizers  
175 (Allied Healthcare, St. Louis, MO, USA). The majority of water vapor from the aerosolization pro-  
176 cess condenses inside the mixing chamber (~570 cm<sup>3</sup>) of the aerosol generator. By using a tempera-  
177 ture and relative humidity (RH) sensor (MSR 145 data logger, MSR Electronics GmbH, Seuzach,  
178 Switzerland) monitoring the flow directly after the aerosol generator we measured RH values of  
179 ~33% (sample flow: 1.4 l/min, dilution: 5 l/min), ~39% (sample flow: 1.4 l/min, dilution: 4 l/min),  
180 and ~54% (sample flow: 2.3 l/min, dilution: 2 l/min). Because of the low RH measured, we did not  
181 use additional drying (e.g., diffusion dryer) to decrease humidity in the sample flow. Hence, the outlet  
182 of the aerosol generator was directly connected to the SIBS inlet with ~30 cm of conductive tubing  
183 (<sup>1</sup>/<sub>4</sub> inch). PSLs were measured for 1 min. Non-fluorescent 4.52 μm PSLs were measured for 2 min,  
184 because of the low number concentrations due to poor aerosolization efficiency and gravitational  
185 settling of larger particle sizes.

186 *S. cerevisiae* was analyzed using a method similar to the one stated above, with the exceptions  
187 that the suspension was prepared with a spatula tip of material mixed into ultrapure water and that a  
188 diffusion dryer (20 cm, 200 g silica) was added to remove excess water vapor. *S. cerevisiae* was  
189 measured for 5 min. Chlorophyll *a*, *b*, and bacteriochlorophyll samples were diluted in 10 ml ethanol.  
190 Between each measurement, the setup was cleaned by aerosolizing ultrapure water for 5 min.

191 PS-DVB particles and biofluorophores (Table S1 and S2) were aerosolized in a dry state. For this  
192 purpose, air at a flowrate of ~0.6 l/min was sent through a HEPA filter into a 10 ml glass vial. A  
193 small amount of each solid powder sample (~1 g) was placed inside the vial and entrained into the  
194 particle-free airstream. Additionally, the sample was physically agitated by tapping the vial. The  
195 outlet was connected with ~20 cm conductive tubing into the inlet of the SIBS. The tubing and glass



196 vial were cleaned after each measurement to prevent particle contaminations from previous measure-  
197 ments. Each powder was sampled until cumulative number concentrations  $> 5000$  particles were  
198 reached.

199 In contrast to the monodisperse and spherical PSL standards, the biofluorophore aerosolization  
200 process provided a polydisperse and morphologically heterogeneous particle distribution with sig-  
201 nificant particle fractions at sizes  $< 1 \mu\text{m}$ . Therefore, we only used particles in a size range between  
202 1 and  $2 \mu\text{m}$  with sufficient fluorescence intensity values for subsequent data analysis. The only ex-  
203 ceptions are the chlorophyll types, where a size range between  $0.5$  and  $2 \mu\text{m}$  (chlorophyll *a* and *b*)  
204 and  $0.5$  and  $1 \mu\text{m}$  (bacteriochlorophyll) were used due to a less efficient particle aerosolization.

205 The fluorescent background of the SIBS was measured daily by firing the xenon lamps into the  
206 optical chamber in the absence of particles (forced trigger mode). In this case, the diaphragm pump  
207 was turned off and the inlet blocked to prevent particles reaching the optical chamber. One forced  
208 trigger mode was performed per day with  $100$  xenon shots per min over a duration of  $5$  min. The  
209 background signal ( $+ 1\sigma$  standard deviation (SD)) was subtracted from derived fluorescence emission  
210 of each sample. Additionally, the background signal was reviewed periodically between each bio-  
211 fluorophore measurement to verify that, e.g., optical components are not coated with residues from  
212 previous measurements. A significant change in background signal was not observed between indi-  
213 vidual measurements.

214 For particle asymmetry measurements, iron (II, III) oxide ( $\text{Fe}_3\text{O}_4$ ), carbon nanotubes, and ammo-  
215 nium sulfate were aerosolized in dry state, and  $2 \mu\text{m}$  non-fluorescent PSLs and ultrapure water were  
216 aerosolized with the aerosol generator method outlined above with SIBS integration times of  $3$  min  
217 in all cases. Due to the broad distribution of asymmetry factor (AF) values for particles below  $1 \mu\text{m}$ ,  
218 only the size fraction  $\geq 1 \mu\text{m}$  was used for subsequent analyses. Furthermore, we observed that AF  
219 bins between  $0$  and  $1$ , and AF bin  $100$  tend to produce increased signal responses, especially for high  
220 particle concentrations, for which they were discarded within the analyses. The origin of this effect  
221 is unknown, but most likely related to detector noise.

222 For collection of particles for microscopy measurements, the sample flow was bypassed and led  
223 through a custom-made particle impactor, which was connected to a mass flow controller (D-6321-  
224 DR, Bronkhorst High-Tech B.V., Ruurlo, Netherlands) and a membrane pump (N816.1.2KN.18,  
225 KNF, Freiburg, Germany). Particles were collected out of the sample flow onto glass cover slips ( $15$   
226 mm diameter) at a flow rate of  $2$  l/min over a duration of  $1$  min.

227

### 228 **2.3 Reference fluorescence spectra**

229 A Dual-FL fluorescence spectrometer (Horiba Instruments Incorporated, Kyoto, Japan) was used as  
230 an offline reference instrument to validate the SIBS spectra. Aqualog V3.6 (Horiba) software was





231 used for data acquisition. The spectrometer was manufacturer-calibrated with NIST Fluorescence  
232 Standard Reference Materials (SRMs 2940, 2941, 2942, and 2943). Aforementioned standard fluor-  
233 ophores were analyzed using the SIBS excitation wavelengths at  $\lambda_{\text{ex}} = 285$  and 370 nm. The Dual-  
234 FL<sup>1</sup> spectrometer uses a xenon arc lamp as excitation source and a CCD (charge-coupled device) as  
235 emission detector, capable of detecting fluorescence emission between 250 and 800 nm. Unless oth-  
236 erwise stated, a low detector gain setting (2.25 e<sup>-</sup> per count) and an emission resolution of 0.58 nm  
237 was used for all measurements with the Dual-FL. Subsequently, we use the term “reference spectra”  
238 for all measurements performed with the Dual-FL. In total, 100 individual spectra were recorded for  
239 each sample and averaged spectra were analyzed in Igor Pro (Wavemetrics, Lake Oswego, Oregon  
240 USA). Background measurements (solvent<sup>2</sup> in the absence of particles) were taken under the same  
241 conditions as for sample measurements and subtracted from the emission signal. For direct compar-  
242 ison to spectra recorded by the SIBS, reference spectra were re-binned by taking the sum of the  
243 fluorescence intensity within the spectral bin width of each SIBS detection channel (Table 1).

244 For PSL measurements, 1.5  $\mu\text{l}$  of each PSL stock solution was diluted in 3.5 ml ultrapure water  
245 in a 10 x 10 x 40 mm UV quartz cuvette (Hellma Analytics, Müllheim, Germany) and constantly  
246 stirred with a magnetic stirrer to avoid particle sedimentation during measurements. Chlorophyll *a*  
247 and *b* and bacteriochlorophyll were handled equally, however concentrations were individually ad-  
248 justed to prevent the detector from being saturated and to avoid self-quenching or inner filter effects  
249 (Sinski and Exner, 2007). Concentrations were used as follows: chlorophyll *a*: 300 nmol/l, chloro-  
250 phyll *b*: 1  $\mu\text{mol/l}$ , and bacteriochlorophyll: 3  $\mu\text{mol/l}$ . PSLs, chlorophyll *b*, and bacteriochlorophyll  
251 measurements were performed with an integration time of 2 s. For chlorophyll *a* an integration time  
252 of 1 s was used.

253 All other biofluorophores, *S. cerevisiae*, and PS-DVB particles were measured in dry state using  
254 a front surface accessory (Horiba). The sample was placed into the surface holder and covered with  
255 a synthetic fused silica window. To limit detector saturation from more highly fluorescent particle  
256 types, the surface holder was placed at a 70° angle to the fluorescence detector for NAD and ribofla-  
257 vin, 75° for tyrosine, 80° for *S. cerevisiae*, and 85° for tryptophan and PS-DVB particles and subse-  
258 quently excited at  $\lambda_{\text{ex}}=285$  and 370 nm. Emission resolution and detector gain settings were used as  
259 for measurements of samples in solution, except for an integration time of 1 s for all dry samples.  
260 Background measurements were performed as described above and subtracted from each sample.  
261 Excitation-emission matrices (EEMs) were measured with the same samples as for single wavelength  
262 measurements. EEMs were recorded at excitation wavelengths between  $\lambda_{\text{ex}} = 240$  and 800 nm (1 nm  
263 increments) and an emission range between  $\lambda_{\text{em}} = 247$  and 829 nm (0.58 nm increments). Exposure

<sup>1</sup> Technical information taken from Dual-FL operation manual, rev. A, 30 NOV 2012; Horiba.

<sup>2</sup> Note that  $\geq 99.8$  % ethanol was used as solvent for chlorophyll *a*, *b*, and bacteriochlorophyll instead of ultrapure water.





264 times of 1 s were used, except for 2  $\mu\text{m}$  green, 3  $\mu\text{m}$  non-fluorescent PSLs (2 s), and NAD (0.5 s).  
265 EEMs were analyzed using Igor Pro.

266

#### 267 **2.4 Calibration lamps and spectral correction**

268 The relative responsivity of a fluorescence detector can vary substantially across its emission range  
269 and, therefore, must be spectrally corrected as a function of emission wavelength (e.g., DeRose, 2007;  
270 Lakowicz, 2004). For spectral correction it was important to choose: (i) light sources covering the  
271 full spectral emission range of the SIBS, with temporal stability on the timescale of many months  
272 and (ii) a calibrated and independent spectrometer to serve as spectral reference.

273 A deuterium-halogen lamp (DH-Mini; Ocean Optics, Largo, FL, USA) and a halogen projector  
274 lamp (EHJ 24 V, 250 W; Ushio Inc., Tokyo, Japan) were used as calibration light sources. Both  
275 lamps were connected to a 50 cm optical fiber (FT030, Thorlabs, Newton, NJ, USA) and vertically  
276 fixed inside the optical chamber of the Dual-FL spectrometer. An aluminum mirror was attached to  
277 the end fitting of the optical fiber, reflecting light in a 90° angle into the detector opening. The pro-  
278 jector halogen lamp was allowed to warm up for 30 s before each measurement. For all power levels  
279 (100, 150, 200, and 250 W), an integration time of 3 s was used. The DH-Mini was operational for  
280 30 min before each measurement. Settings were used as for the projector halogen lamp, however,  
281 due to the low emission a high detector gain setting (9 e<sup>-</sup> per count) was used with an integration time  
282 of 25 s. As described in Sect. 2.3, 100 single measurements were taken and averaged (Fig. S1). For  
283 the SIBS, both light sources were measured in the same way as for the reference spectra. Measure-  
284 ments were performed with a detector amplification at 610 V (see Sect. 4.2). Background measure-  
285 ments were taken as described in Sect. 2.2. Projector halogen lamp spectra (at all power levels) were  
286 recorded for 3 min, the DH-Mini, due to its low emission intensity, for a duration of 5 min.

287 For the halogen projector lamp, averaged intensity values in each spectral bin were acquired at  
288 each power level (150, 200, and 250 W). Spectra measured at 100 W were discarded due to the low  
289 and unstable emission at wavelengths shorter than ~500 nm (Fig. S1). Reference spectra and spectra  
290 recorded by the SIBS were normalized onto the SIBS detection channel 9 ( $\lambda_{\text{mean}} = 528.0$  nm), which  
291 is, theoretically, the detection channel with the highest responsivity (see Sect. 4.3). The individual  
292 spectral correction factors were calculated by dividing the reference spectra by the spectra derived  
293 from the SIBS. The final correction factors are a combination out of both light sources where the  
294 detection channels 1-5 ( $\lambda_{\text{mean}} = 302.2 - 415.6$  nm) include the correction factors for the DH-Mini and  
295 the detection channels 6-16 ( $\lambda_{\text{mean}} = 443.8 - 721.1$  nm) the correction factors for the halogen projector  
296 lamp. At the intersection between channel 5 and 6, both corrections (DH-Mini, halogen) are in good  
297 agreement ( $\Delta_{\text{correction}} = 0.6$  in channel 6). For all particle measurements described in the following



298 sections, the background signal and raw sample spectra recorded by the SIBS were multiplied by  
299 those correction factors.

300

### 301 **2.5 Microscopy of selected reference particles**

302 Bright field microscopy was conducted using an Eclipse Ti2 (Nikon, Tokyo, Japan) with a 60x im-  
303 mersion oil objective lens and an additional optical zoom factor of 1.5, resulting in a 90x magnifica-  
304 tion. Glass cover slips, used as collection substrates in the particle impactor (Sect. 2.2), were put onto  
305 a specimen holder and fixed with tape. Images were recorded using a DS Qi2 monochrome micro-  
306 scope camera with 16.25 megapixels and z-stacks of related images were created using the software  
307 NIS-Elements AR (both Nikon).

308

### 309 **2.6 Ambient measurement setup and data analysis**

310 The SIBS was operated between the 5<sup>th</sup> of April to the 7<sup>th</sup> of May 2018 on the roof (fourth floor inside  
311 a roof laboratory) of the Max Planck Institute for Chemistry in Mainz, Germany (49°59'28.2"N,  
312 8°13'44.5"E) similar to measurements as described in Huffman et al. (2010) using a UV-APS. The  
313 period between the 12<sup>th</sup> and 18<sup>th</sup> of April 2018 is described here to highlight the capability of the  
314 SIBS to monitor ambient aerosol. Beside of the SIBS, four additional instruments (data not shown  
315 within this study) were connected with ~20 cm conductive tubing ( $\frac{1}{4}$  inch) to a sample airflow  
316 splitter (Grimm Aerosol Technik GmbH & Co. KG, Ainring, Germany). The splitter was connected  
317 to 1.5 m conductive tubing ( $\frac{5}{8}$  inch), bent out of the window, and connected to 2.4 m stainless steel  
318 tubing ( $\frac{5}{8}$  inch, Dockweiler AG, Neustadt-Glewe, Germany) vertically installed. Between a TSP  
319 head (total suspended particles, custom-made) and the stainless steel tubing, a diffusion dryer (1 m,  
320 1 kg silica) was installed. Silica was exchanged every third to fourth day and periodic forced trigger  
321 measurements were performed. The total flow was ~8.4 l/min.

322 For measurements presented here, only particles were included if they showed fluorescence emis-  
323 sion in at least two consecutive spectral channels. This filter was applied to limit noise introduced  
324 from measurement artifacts from a variety of sources and will need to be investigated in more detail.  
325 The conservative analysis approach here suggests that the values reported are likely to be a lower  
326 limit for fluorescent particle number and fraction. The observations are in line with previous meas-  
327 urements, however, giving general support that the SIBS measurements are reasonable. Note that the  
328 maximum repetition rate of the xenon lamps is 125 Hz, corresponding to maximum concentrations  
329 of 20 particles per  $\text{cm}^{-3}$  (see Sect. 3.3). Because ~50% of the total coarse particle number were excited  
330 by xenon 1 and xenon 2, the fluorescent particle concentrations and fluorescent fractions are cor-  
331 rected accordingly.



### 332 3. Design and components of the SIBS

333 The SIBS is based on the general optical design of the WIBS-4A (e.g., Foot et al., 2008; Healy et al.,  
334 2012; Hernandez et al., 2016; Kaye et al., 2005; Perring et al., 2015; Robinson et al., 2017; Savage  
335 et al., 2017; Stanley et al., 2011) with improvements based on a lower particle sizing limit, resolved  
336 fluorescence detection, and a broader emission range. The instrument provides information about  
337 size, particle asymmetry, and fluorescence properties for individual particles in real-time. The exci-  
338 tation wavelengths are optimized for the detection of the biological fluorophores tryptophan,  
339 NAD(P)H, and riboflavin. However, other fluorophores in PBAP will certainly fluoresce at these  
340 excitation wavelengths as many of them cluster in two spectral fluorescence “hotspots” as summa-  
341 rized in Pöhlker et al. (2012 and references therein) and as shown for WIBS-4A measurements by  
342 Savage et al. (2017). Figure 1 shows an overview of excitation wavelengths and emission ranges of  
343 the UV-APS, WIBS-4A, WIBS-NEO, and SIBS for bioaerosol detection in relation to the spectral  
344 location of selected biofluorophores, such as tyrosine, tryptophan, NAD(P)H, riboflavin, and chloro-  
345 phyll *b*. At  $\lambda_{\text{ex}} = 285$  nm, the SIBS excites fluorophores in the “protein hotspot”, at  $\lambda_{\text{ex}} = 370$  nm  
346 fluorophores in the “flavin/coenzyme hotspot” (Pöhlker et al., 2012). In contrast to the UV-APS, the  
347 SIBS is able to detect fluorescence signals from chlorophyll due to the extended upper spectral range  
348 of detection (up to  $\lambda_{\text{em}} = 721$  nm). Both the WIBS-4A and WIBS-NEO cover the spectral emission  
349 range for chlorophyll *b*, however, cannot provide resolved spectral information to separate it from  
350 other fluorophores. Table 2 summarizes and compares parameters and technical components of the  
351 SIBS, WIBS-4A, and WIBS-NEO. The individual components are described in detail in the subse-  
352 quent sections.

353

#### 354 3.1 Aerosol inlet and flow diagram

355 The design for the aerosol inlet of the SIBS is identical to the inlet of the WIBS-4A and WIBS-NEO.  
356 A detailed flow diagram is shown in Figure S2. Aerosol is drawn in via an internal pump as laminar  
357 air flow through a tapered delivery nozzle (Fig. S2a) where sheath ( $\sim 2.2$  l/min) and sample flow ( $\sim 0.3$   
358 l/min) are separated.

359

#### 360 3.2 Size and shape analysis

361 After passing the delivery nozzle, entrained particles traverses a 55 mW continuous-wave diode laser  
362 at  $\lambda = 785$  nm (#2 in Fig. 2 and position #1 in Fig. S3). Unlike in the WIBS-4A and WIBS-NEO (635  
363 nm diode laser), the triggering laser in the SIBS is in the near-infrared (IR) region ( $> 700$  nm) and,  
364 therefore, outside the detectable emission range of the 16-channel photomultiplier tube (PMT) to  
365 avoid scattered light from the particle trigger laser being detected (see Fig. 1). The side and forward



366 scattered light is collected and used for subsequent measurements. Side scattered light is collected  
367 by two concave mirrors, which are directed at  $90^\circ$  from the laser beam axis, and reflect the collected  
368 light onto a dichroic beam splitter (#7 in Fig. 2). A PMT (H10720-20, Hamamatsu Photonics K.K.,  
369 Japan) converts incoming light signals into electrical pulses, which are used for particle triggering  
370 and sizing (#6 in Fig. 2). For the determination of the optical particle size, the SIBS uses a calculated  
371 calibration curve according to the Lorenz-Mie Theory, assuming spherical and monodisperse PSLs  
372 with a refractive index of 1.59 (Brandrup et al., 1989; Lorenz, 1890; Mie, 1908). Compared to aero-  
373 dynamic sizing, which depends on particle morphology and density (e.g., Reid et al., 2003; Reponen  
374 et al., 2001), the calculated optical diameter can vary significantly if the assumption of sphericity is  
375 not fulfilled. In contrast, optical sizing is not as affected by differences in material density. The in-  
376 strument operator must thus be aware of uncertainties in measured particle size due to, e.g., particle  
377 morphology, spatial orientation of a particle when traversing the trigger laser or changing refractive  
378 indices. In contrast to the WIBS-4A, the SIBS and WIBS-NEO detect a range of particle sizes be-  
379 tween  $\sim 0.3$  and  $100 \mu\text{m}$  (nominal), achievable by using one PMT gain setting instead of switching  
380 between a “Low Gain” and “High Gain” setting. Physical and technical details of this Gain-switching  
381 method are patent pending and are not publicly available.

382 The forward-scattered light is measured by a quadrant PMT (#5 in Fig. 2) to detect the scatter  
383 asymmetry for each particle (Kaye et al., 1991, 1996). A OG-515 long pass filter (Schott AG, Mainz,  
384 Germany) prevents incoming light from the xenon flash lamps in a spectral range below  $515 \pm 6 \text{ nm}$   
385 from reaching the Quadrant PMT. To calculate the AF, the root-mean-square variations for each  
386 quadrant of the PMT of the forward-scattered light intensities are used (Gabey et al., 2010). The AF  
387 broadly relates whether a particle is more spherical or fibril. Theoretically, for a perfectly spherical  
388 particle, the AF would be 0, whereas an elongated particle would correspond to an AF of 100 (Kaye  
389 et al., 1991). However, due to electrical and optical noise of the Quadrant PMT, the AF value of a  
390 sphere is usually between ca. 2 and 6 (according to WIBS-4A service manual (DOC-0345 Rev A)).  
391 Because the AF value depends on physical properties of optical components, the baseline for spher-  
392 ical particles may shift even within identical instruments (Savage et al., 2017). For example, the study  
393 by Toprak & Schnaiter (2013) reported an average AF value for spherical particles of 8 using a  
394 WIBS-4A. In contrast, AF values shown by Foot et al. (2008) were, on average, below  $\sim 5$  for  
395 spherical particles measured with a WIBS-2s prototype.

396

### 397 **3.3 Fluorescence excitation**

398 Two xenon flash lamps (L9455-41, Hamamatsu) (#3 and #4 in Fig. 2) are used to induce fluores-  
399 cence. They emit light pulses, which exhibit a broad excitation wavelength range of 185 to 2000 nm.  
400 The light is optically filtered to obtain a defined excitation wavelength. Further information about



401 spectral properties of the xenon flash lamps can be found elsewhere (Specification sheet  
402 TLSZ1006E04, Hamamatsu, May 2015). Figure 3 displays relevant optical properties of the lamps  
403 and filters used within the SIBS, WIBS-4A, and WIBS-NEO. For the SIBS, a BrightLine® FF01-  
404 285/14-25 (Semrock Inc., Rochester, NY, USA) single-band bandpass filter is used with  $\lambda_{\text{mean}} = 285$   
405 nm and an effective excitation band<sup>3</sup> of 14 nm width is used for xenon 1. For xenon 2, the single-  
406 band bandpass filter BrightLine® FF01-370/36-25 (Semrock) is used with  $\lambda_{\text{mean}} = 370$  nm and with  
407 an effective excitation band of 36 nm width. The only difference between all three instruments is that  
408 the WIBS-4A and WIBS-NEO use a different single-band bandpass filter for xenon 1 (Semrock,  
409 BrightLine® FF01-280/20-25;  $\lambda_{\text{mean}} = 280$  nm; effective excitation band of 20 nm). The excitation  
410 light beam is focused on the sample flow within the optical cavity, resulting in a rectangular beam  
411 shape of ~5 mm by 2 mm. Xenon 1 is triggered when particles pass position 2 in Figure S3 and  
412 approximately 10  $\mu\text{s}$  later xenon 2 is triggered as the particles move further to position 3 in Figure  
413 S3. After firing, the flash lamps need ~5 ms to recharge. During the recharge period, particles are  
414 counted and sized but no fluorescence information is recorded. The maximum repetition rate of the  
415 xenon lamps yields a measurable particle number concentration of  $\sim 2 \times 10^4 \text{ l}^{-1}$  (corresponding to  
416  $20 \text{ cm}^{-3}$ ).

417 Irradiance values from light sources becomes a crucial factor when interpreting derived fluores-  
418 cence data of LIF instruments because the fluorescence intensity is directly proportional to the inten-  
419 sity of incident radiant power, described by the relationship:

420

$$421 \quad F = \phi I_0 (1 - e^{-\epsilon bc}) \quad (1)$$

422

423  $\phi$ : quantum efficiency,  $I_0$ : intensity of incident light,  $\epsilon$ : molar absorptivity,  $b$ : path length (cell),  $c$ :  
424 molar concentration (Guilbault, 1990).

425 To measure the irradiance of each xenon lamp after optical filtering, we used a thermal power  
426 head (S425C, Thorlabs), which was placed at a distance of 11.3 cm (focus length from xenon arc  
427 bow to sample flow intersection) from the xenon lamp measuring over a duration of 1 min at 10  
428 xenon shots per s. By measuring new xenon lamps, we observed an average irradiance of 14.8  
429  $\text{mW}/\text{cm}^2$  for xenon 1 and 9.6  $\text{mW}/\text{cm}^2$  for xenon 2, corresponding to ~154 % higher irradiance (spec-  
430 trally integrated) from xenon 1. A second set of lamps, used intermittently for three years including  
431 several months of continuous ambient measurements and a lab study with high particle concentra-  
432 tions, exhibited average irradiance values of 10.8  $\text{mW}/\text{cm}^2$  ( $1\sigma$  SD 1.8  $\text{mW}/\text{cm}^2$ ) for xenon 1 and 4.9

<sup>3</sup> The effective excitation band is defined as “guaranteed minimum bandwidth” (GMBW), describing the spectral region a bandpass filter transmits light relative from the mean wavelength. For example, a GMBW of 14 nm means that light is transmitted in a 7 nm spectral range above and below the mean wavelength.



433  $\text{mW}/\text{cm}^2$  ( $1\sigma$  SD  $1.9 \text{ mW}/\text{cm}^2$ ) for xenon 2, corresponding to  $\sim 220\%$  higher irradiance from xenon  
434 1. Comparing the nominal, transmission-corrected irradiance data from the two xenon lamps pro-  
435 vided by the lamp supplier (Fig. 3a and 3b, red dashed lines), an irradiance imbalance between xenon  
436 1 and xenon 2 can be assumed for all three LIF instruments discussed here (SIBS, WIBS-4A, and  
437 WIBS-Neo).

438 Results shown here are comparable to multiple WIBS studies (e.g., Hernandez et al., 2016;  
439 Perring et al., 2015; Savage et al., 2017), where fluorescence emission intensities at  $\lambda_{\text{ex}} = 280 \text{ nm}$   
440 (xenon 1) also show a tendency to be higher than those at  $\lambda_{\text{ex}} = 370 \text{ nm}$  (xenon 2).

441

#### 442 **3.4 Spectrally resolved fluorescence detection**

443 Fluorescence emission from excited particles is collected by two parabolic mirrors in the optical  
444 cavity and delivered onto a custom-made dichroic beam splitter (Semrock, #7 in Fig. 2). The beam  
445 splitter allows transmission of incoming light between  $\sim 300$  and  $710 \text{ nm}$ , with an average transmis-  
446 sion efficiency of  $96\%$ . At wavelengths shorter than  $300 \text{ nm}$ , the transmission decreases rapidly to  $<$   
447  $20\%$  at  $275 \text{ nm}$ . At the upper detection end of the SIBS ( $\lambda_{\text{mean}} = 721 \text{ nm}$ ), the transmission efficiency  
448 decreases to  $\sim 89\%$ . The scattering light from the diode laser is reflected at a  $90^\circ$  angle onto the PMT  
449 used for particle detection and sizing. At the excitation wavelength of  $785 \text{ nm}$ , the reflection effi-  
450 ciency is stated at  $\sim 95\%$  (Fig. S4).

451 After passing the dichroic beamsplitter, the photons are led into a grating polychromator (A  
452 10766, Hamamatsu) (#8 in Fig. 2). A custom-made transmission grating (Hamamatsu) is used to  
453 diffract incoming light within a nominal spectral range between  $290.8 - 732.0 \text{ nm}$ . In case of the  
454 SIBS, a grating with  $300 \text{ g}/\text{mm}$  groove density and  $400 \text{ nm}$  blaze wavelength is used, resulting in a  
455 nominal spectral width of  $441.2 \text{ nm}$  and a resolution of  $28.03 \text{ nm}/\text{mm}$ . After passing the transmission  
456 grating, the diffracted light hits a 16-channel linear array multi-anode PMT (H12310-40, Hamama-  
457 tsu) (#9 in Fig. 2) with defined mean wavelengths for each channel as shown in Table 1.

458 For each single particle detected, two spectra are recorded, at  $\lambda_{\text{ex}} = 285$  and  $370 \text{ nm}$ . The detect-  
459 able band range of the PMT overlaps the excitation wavelength of xenon 2. Therefore, a notch optical  
460 filter (Semrock) is placed between the optical chamber and the grating polychromator to prevent the  
461 detector from being saturated. Incoming light at wavelengths shorter than  $300 \text{ nm}$  and from  $362$   
462  $377 \text{ nm}$  is blocked from reaching the PMT resulting in a reduced spectral bin width for detection  
463 channels 1, 3 and 4. The first three detection channels are omitted because their mean wavelengths  
464 are below  $\lambda_{\text{ex}} = 370 \text{ nm}$  (see also Fig. 1). Accordingly, the emission spectra for xenon 2 excitation  
465 begin at channel 4 ( $\lambda_{\text{mean}} = 387.3 \text{ nm}$ ).

466 Technical data (xenon flash lamps, filters, dichroic beam splitter, PMT responsivity, and trans-  
467 mission grating) described in the previous sections (3.3 and 3.4) were provided by Hamamatsu and



468 Semrock. Note that transmission/reflection efficiencies of the dichroic beam splitter, cathode radiant  
469 sensitivity of the PMT, and diffraction efficiency data are modeled. Thus, individual components  
470 may differ slightly from modeled values, even within the same production batch. Neither company  
471 assumes data accuracy or provides warranty, either expressed or implied.

472 The SIBS was originally designed and marketed to record time-resolved fluorescence lifetime.  
473 The fluorescence lifetime of most biofluorophores, serving as targets for bioaerosol detection, are  
474 usually below 10 ns (e.g., Chorvat & Chorvatova, 2009; Herbrich, et al., 2012; O'Connor et al., 2014;  
475 Richards-Kortum & Sevick-Muraca, 1996). However, by choosing xenon lamps as excitation source,  
476 recording relevant fluorescence lifetimes in this ns range is hampered by the relatively long decay  
477 time of the xenon lamp excitation pulse ( $\sim 1.5 \mu\text{s}$ ). In principle, fluorescence lifetime measurements  
478 would be possible if the xenon lamps were replaced by appropriate laser excitation sources in the  
479 SIBS optical design.

480

### 481 **3.5 Software components and data output**

482 The SIBS uses an internal computer (#10 in Fig. 2) with embedded LabView-based data acquisition  
483 software allowing the user to control functions in real time and change multiple measurement param-  
484 eters. As an example, the “Single Particle” tab out of the SIBS interface is shown in Figure S5. Here,  
485 the user can define, e.g., the sizing limits of the SIBS (upper and lower threshold) and the minimum  
486 size of a particle being excited by the xenon flash lamps. Furthermore, forced trigger measurements  
487 can be performed while on this particular tab. Subsequently, the term “forced trigger measurement”  
488 will be replaced by “background signal measurement”. A local Wi-Fi network is installed so that the  
489 SIBS can be monitored and controlled remotely. A removable hard drive is used for data storage.  
490 Data is stored in a HDF5 format to minimize storage space and optimize data write speed. Resulting  
491 raw data are processed in Igor Pro. As an example: by using a minimum sizing threshold of 500 nm,  
492 the SIBS data output per day, operating in a relatively clean environment ( $\sim 40$  particles per  $\text{cm}^{-3}$ ),  
493 can span several hundreds of MB. In contrast, the data output can increase up to  $\sim 3$  GB daily in  
494 polluted areas ( $\sim 680$  particles per  $\text{cm}^{-3}$ ). By lowering the minimum sizing threshold to 300 nm, the  
495 data volume can exceed 10 GB per day when sampling in a moderately polluted environment ( $\sim 180$   
496 particles per  $\text{cm}^{-3}$ ).

497

## 498 **4. Results and data validation**

### 499 **4.1 Validation of SIBS sizing**

500 To validate the optical sizing of the SIBS, twenty particle size standards were analyzed, covering a  
501 broad size range from 0.3 to 20  $\mu\text{m}$  in particle diameter. Overall, the particle size measurements from





502 the SIBS (optical diameter) show good agreement with the corresponding measurements of physical  
503 diameter reported by PSL and PS-DVB manufacturers (Fig. 4). For the SIBS and WIBS-NEO, the  
504 manufacturer states a nominal minimum size detection threshold of 300 nm. Figure 4 shows that a  
505 linear response between optical particle size and physical particle size extends to at least 300 nm.  
506 Smaller particles were not investigated. The upper size detection threshold is reported by the manu-  
507 facturer to be nominally 100  $\mu\text{m}$ . However, the upper limit was not investigated here due to the dif-  
508 ficulty in aerosolizing particles larger than this. In most field applications, the upper particle size cut  
509 is often far below this value due to unavoidable sedimentation losses of large particles in the inlet  
510 system (e.g., Moran-Zuloaga et al., 2018.; Von der Weiden et al., 2009). Note that the size distribu-  
511 tions of physical diameter for PS-DVB standards are broader compared to the PSL standards, as  
512 reported by the manufacturer (Table S1). This also translates to broader distributions of optical di-  
513 ameter measured by the SIBS than for the PSL particles. The 0.356  $\mu\text{m}$  PSL sample was an outlier  
514 with respect to the overall trend, showing an optical diameter of 0.54  $\mu\text{m}$ . We suspect that this devi-  
515 ation between physical and optical size can be explained by a poor quality of this particular PSL  
516 sample lot rather than an instrumental issue, and so it was not included in the calculation of the trend  
517 line (Fig. 4). Furthermore, the SIBS was shown to undersize the PSLs between 0.6 and 0.8  $\mu\text{m}$ ,  
518 however, the overall trend exhibits a coefficient of determination of  $r^2 > 0.99$ .

519 As mentioned in Sect. 3.2, an important point regarding the SIBS and WIBS-NEO is that the size  
520 calibration within the unit cannot be changed by the user, meaning that the PMT output voltages are  
521 transformed directly to outputted physical diameter within the internal computer using a proprietary  
522 calculation. It is still important, however, for the user to perform sizing calibration checks frequently  
523 to verify and potentially post-correct particle sizing of all particle sizing instruments, including the  
524 SIBS and WIBS-NEO.

525

#### 526 4.2 Amplification of fluorescence detector

527 As with all optical detection techniques, adequate understanding of detection thresholds is an essen-  
528 tial aspect of instrument characterization and use (e.g., Jeys et al., 2007; Savage et al., 2017). Ap-  
529 plication of appropriate voltage gain settings must be applied to the physical detection process so as  
530 not to lose information about particles that cannot be recovered by post-processing of data. Yet par-  
531 ticles in the natural atmosphere exhibit an extremely broad range of fluorescence intensities (many  
532 orders of magnitude), arising from the breadth of quantum yields for fluorophores occurring in aer-  
533 osols and from the steep increase of fluorescence emission intensity with particle size (2<sup>nd</sup> to 3<sup>rd</sup>  
534 power) (e.g., Hill et al., 2015; Könemann et al., 2018; Sivaprakasam et al., 2011; Swanson &  
535 Huffman, 2018). This range of fluorescence properties is generally broader than the dynamic range  
536 of any single instrument, and so a UV-LIF instrument can be operated, e.g., to either: (i) apply a



537 higher detector gain to allow high sensitivity toward detecting weakly fluorescing particles, often  
538 from rather small particles ( $< 1 \mu\text{m}$ ), at the risk of losing fluorescence information for large or  
539 strongly fluorescent particles due to detector saturation or (ii) apply a lower detector gain to prefer-  
540 entially detect a wide range of more highly fluorescent particles, but at the risk of not detecting  
541 weakly fluorescent or small particles.

542 Amplification voltage of the 16-channel PMT used in the SIBS can be adjusted between 500 and  
543 1200 V. Each of the 16 detection channels can also be individually adjusted using digital gain settings  
544 within the SIBS acquisition software. This channel-specific gain does not affect the amplification  
545 process (e.g., the dynode cascade), but rather modifies the output signal of single detection channels  
546 digitally. The digital gain is applied only after the signal collection process, and so cannot compensate  
547 for a signal that is below the noise threshold or that saturates the detector. The digital gain was thus  
548 left at the maximum gain level (255 arbitrary units (a.u.)) for all channels during particle measure-  
549 ments discussed here.

550 To explore the influence of amplification voltage on particle detectability,  $0.53 \mu\text{m}$  purple PSLs  
551 were chosen to arbitrarily represent the lower limit of detectable fluorescence intensity. Using larger  
552 ( $0.96 \mu\text{m}$ ) particles comprised of the same purple fluorophore, Könemann et al. (2018) showed that  
553 the particles were only narrowly detectable above the fluorescence threshold in each of the three  
554 channels of a WIBS-4A (same unit as used in Savage et al., 2017) and so the smaller,  $0.53 \mu\text{m}$  PSLs  
555 were chosen here as a first proxy for the most weakly fluorescing particles we would expect to detect.  
556 To improve the signal to noise ratio (SNR) for the lower fluorescence detection limit, the PMT am-  
557 plification voltage was varied in seven steps between 500 and 1000 V (corresponding to a gain from  
558  $10^3$  to  $10^6$ , specification sheet TPMP01060E02, Hamamatsu, June 2016) for purple PSLs and back-  
559 ground signals (Fig. 5a). Whereas PSL spectra at a PMT amplification of 500 V were indistinguish-  
560 able from the background signal ( $+ 1\sigma$  SD), spectra show a discernable peak at  $\lambda_{\text{mean}} = 415.6 \text{ nm}$   
561 above 600 V. Subsequently, the SIBS was operated with a PMT amplification voltage of 610 V  
562 corresponding to the lowest SNR threshold accepted (Fig. 5a, b). The detection of small biological  
563 particles was tested by measuring the emission spectrum of *S. cerevisiae* as an example of a PBAP  
564 (see also Pöhlker et al., 2012). On average, the size of intact *S. cerevisiae* particles range between  $\sim 2$   
565  $- 10 \mu\text{m}$  (e.g., Pelling et al., 2004; Shaw et al., 1997). To test the ability of the SIBS to detect low  
566 intensity emissions, we separately analyzed *S. cerevisiae* particles between  $0.5$  and  $1 \mu\text{m}$ , which most  
567 likely includes cell fragments caused by the aerosolization process (Fig. 5c). The tryptophan-like  
568 emission, peaking in detection channel 2 ( $\lambda_{\text{mean}} = 330.6 \text{ nm}$ ) for  $\lambda_{\text{ex}} = 285 \text{ nm}$ , reveals intensity values  
569 below 100 a.u., which are comparable to fluorescence intensity values derived from  $0.53 \mu\text{m}$  purple  
570 PSLs (detection channel 5,  $\lambda_{\text{mean}} = 415.6 \text{ nm}$ , Fig. 5d). These two tests for *S. cerevisiae* and  $0.53 \mu\text{m}$   
571 purple PSLs confirmed the instruments ability to detect emission spectra from particles at least as



572 strongly fluorescent as these two test cases, leaving a wide range to detect larger and more intensely  
573 fluorescing particles. By using a  $3\sigma$  SD threshold, the fluorescence peak at  $\lambda_{\text{mean}} = 415.6$  nm of 0.53  
574  $\mu\text{m}$  purple PSLs is still detectable but cannot be distinguished from the background signal at a  $6\sigma$  SD  
575 threshold anymore. Therefore, fluorescence intensity values at the lower detection limit should be  
576 treated with care. Corrected spectra of both *S. cerevisiae* and 0.53  $\mu\text{m}$  purple PSLs can be found in  
577 the supplement (Fig. S6). By operating the SIBS at relatively low detector amplification, very weak  
578 fluorescence, especially from small particles ( $< 1 \mu\text{m}$ ) might not exceed the detection threshold dur-  
579 ing field applications and would be missed. Further investigation will be necessary to choose ampli-  
580 fication voltages appropriate for individual applications where smaller or otherwise weakly fluores-  
581 cent particles might be particularly important. For all subsequent measurements discussed here, a  
582 PMT amplification voltage of 610 V was used.

583 Saturation only occurred for 15 and 20  $\mu\text{m}$  non-fluorescent PS-DVB particles. As highlighted in  
584 Figure S7, the polystyrene/detergent signal (Könemann et al., 2018) at  $\lambda_{\text{ex}} = 285$  nm for 10  $\mu\text{m}$  PS-  
585 DVB particles can be spectrally resolved (Fig. S7b), whereas the spectrum for 15  $\mu\text{m}$  PS-DVB par-  
586 ticles (Fig. S7e) is altered due to single particles ( $\sim 10\%$  out of 400 particles) saturating the detector  
587 (at 62383 a.u.). By comparing the defined lower detection end (Fig. 5) to the upper end (Fig. S7), a  
588 quantitative difference of approximately three orders of magnitude can be estimated, indicating a  
589 wide detectable range at the chosen amplification voltage setting.

590

### 591 4.3 Wavelength-dependent spectral correction of detector

592 The 16 cathodes of the PMT should be considered as independent detectors with wavelength depend-  
593 ent, individual responsivity and amplification characteristics. In combination with physical properties  
594 of technical components (e.g., excitation sources, optical filters, gratings), an instrumental-specific  
595 spectral bias might result in incorrect or misleading spectral patterns if not corrected (e.g., DeRose,  
596 2007; DeRose et al., 2007; Holbrook et al., 2006). To compensate for such potential instrumental  
597 biases, we used a spectral correction approach as described in Sect. 2.4. The spectral correction fac-  
598 tors are comparable to the theoretical responsivity of the PMT with the highest correction for chan-  
599 nels 1-4 ( $\lambda_{\text{mean}} = 302.2 - 387.3$  nm) and 14-16 ( $\lambda_{\text{mean}} = 666.5 - 721.1$  nm) (Fig. 6). Channel 8 ( $\lambda_{\text{mean}} =$   
600 500.0 nm) shows the highest responsivity and channels 6 and 7 ( $\lambda_{\text{mean}} = 443.8$  and 471.9 nm) exhibit  
601 a noticeable lower responsivity than their adjacent channels (see also Sect. 4.4.1). The spectral cor-  
602 rection shows several peaks (e.g., detector channels 3, 5, and 8) and dips (e.g., detector channels 4,  
603 6, and 7) (Fig. 6), however, this pattern is due to gain variations for different channels and is not  
604 noise.

605 It is important to note that the detector settings and spectral correction uniquely refer to the SIBS  
606 unit as it was used for the current study. Due to technical and physical variability as stated above, it



607 is likely that the spectral correction required for other SIBS units would be somewhat different. Fur-  
608 thermore, the wavelength-dependent detector correction may change over time due to material fa-  
609 tigue or contaminations in the optical chamber affecting background signal measurements. Periodic  
610 surveillance and adjustments are therefore required, especially after measurements where the instru-  
611 ment was exposed to high particle concentrations or was operated during extreme weather or envi-  
612 ronmental conditions (e.g., temperature, humidity, vibration). For particle sizing verification, we rec-  
613 ommend the use of 0.5, 1, and 3  $\mu\text{m}$  non-fluorescent PSLs. Regarding a fluorescence response check,  
614 we recommend 2  $\mu\text{m}$  green and 2  $\mu\text{m}$  red PSLs for the validation of the spectral responsivity maxi-  
615 mum and the upper (near-IR) detection range. To our knowledge, no fluorescent dyed PSLs are avail-  
616 able to verify the response within the lower spectral detection range (UV) of the SIBS. However, the  
617 polystyrene signal of 3  $\mu\text{m}$  non-fluorescent PSLs (Fig S7g, h, i, see also Könemann et al., 2018)  
618 represents a compromise between signal strength at  $\lambda_{\text{ex}} = 285 \text{ nm}$  and aerosolization efficiency (com-  
619 pared to PSLs with larger sizes) for a spectral responsivity validation.

620

#### 621 **4.4 Fluorescence spectra of standards**

##### 622 **4.4.1 PSL standards**

623 The SIBS spectra for the four different PSL standards, covering an emission range from UV to near-  
624 IR, generally agree well with the corresponding reference spectra (Fig. 7). Each of the two excitation  
625 wavelengths probe separate fluorescent modes, which appear at approximately the same emission  
626 wavelength for a given PSL type (e.g.,  $\lambda_{\text{em}} = \sim 580 \text{ nm}$  for red PSLs, Fig. 7j), as discussed by  
627 Könemann et al. (2018). Moreover, even the rather weak polystyrene and detergent fluorescence,  
628 systematically associated with PSL suspensions (Könemann et al., 2018), is resolved by the SIBS at  
629  $\lambda_{\text{ex}} = 285 \text{ nm}$  and  $\lambda_{\text{em}} = \sim 300 \text{ nm}$  (Fig. 7b, e, h, k). It is further noteworthy that emissions at  $\lambda_{\text{ex}} = 285$   
630 nm are generally higher than derived emissions at  $\lambda_{\text{ex}} = 370 \text{ nm}$  (Fig. 7c, f, i, l), supporting the finding  
631 that a particle receives higher irradiance values from xenon 1 than from xenon 2 (see also Sect. 3.3).

632 As mentioned in Sect. 4.3, detection channels 6 and 7 require relatively large correction factors.  
633 For 2.07  $\mu\text{m}$  purple PSLs (Fig. 7b, c), the SIBS spectra closely match the references spectra after  
634 correction. For the 2.1  $\mu\text{m}$  blue PSLs (Fig. 7e, f), however, the corrected spectrum matches the ref-  
635 erence spectrum well, except at detection channel 6 ( $\lambda_{\text{mean}} = 443.8 \text{ nm}$ ), where the SIBS spectrum is  
636 lower than the reference spectrum by approximately 50%. This effect was also observed for 1  $\mu\text{m}$   
637 blue PSLs (Thermo Fisher, B0100), doped with the same fluorophore (data not included in this  
638 study). The reason for this malfunction is unknown. Nevertheless, because this effect only occurs  
639 noticeably for highly fluorescent blue PSLs and NAD (see also Sect. 4.4.2), one explanation could  
640 be that the instrument-dependent dynode cascade (the electronic amplification stages) for this partic-



641 ular detection channel is suppressed, resulting in a lower amplification efficiency. In this case, rela-  
642 tively low signals could be amplified correctly, whereas medium or high intensity emission could  
643 only be amplified up to a certain level. The amplification threshold for detection channel 6 is, how-  
644 ever, unknown and needs further verification.

645

#### 646 **4.4.2 Biofluorophore standards**

647 Figure 8 and 9 highlight fluorescence spectra of different biofluorophores measured by the SIBS,  
648 which correspond to related reference spectra (compare also Pöhlker et al., 2012), showing that amino  
649 acids (fluorescence emission only at  $\lambda_{\text{ex}} = 285$  nm), co-enzymes and flavin compounds (fluorescence  
650 emission at  $\lambda_{\text{ex}} = 285$  and 370 nm), and chlorophylls (fluorescence emission only at  $\lambda_{\text{ex}} = 370$  nm)  
651 can be spectrally distinguished.

652 The uncorrected spectrum of tryptophan (Fig. S9) highlights the necessity of a spectral correction  
653 to compensate for the low detector responsivity within the UV and near-IR bins. If the fluorescence  
654 signal of tryptophan remains uncorrected, the spectra is shifted slightly to longer wavelengths (red-  
655 shifted) due to the low responsivity of channel 2 in comparison to channel 3, resulting in misleading  
656 spectral information. For NAD (Fig. 8h, i), fluorescence intensity values of channel 6 are lowered  
657 due the suppressed amplification efficiency in this particular channel as described for blue PSLs  
658 (Sect. 4.4.1).

659 All biofluorophores (except chlorophyll types) were aerosolized as dry powders (see Sect. 2.2)  
660 to avoid fluorescence solvatochromism effects, means solvent-dependent spectral shifts relative to  
661 the dry fluorophore state, which serves as a reference case here (e.g., Johnson et al., 1985). Solvato-  
662 chromism of fluorophores in aqueous solution – the only atmospherically relevant case – typically  
663 shifts fluorescence emissions to longer wavelengths due to the stabilized excited state caused by polar  
664 solvents (Lakowicz, 2004). This spectral red-shift can be seen in Figure S10, where the peak maxi-  
665 mum for NAD shows a difference of  $\sim 15$  nm between a dry and water-solved state, whereas ribofla-  
666 vin reveals an even higher shift of  $\sim 37$  nm. Here, solvatochromism serves as an example for fluores-  
667 cence spectra that vary substantially as a function of the fluorophore's microenvironments (e.g., sol-  
668 vent polarity, pH, temperature).

669 Each of the three types of chlorophyll exhibit the weakest emission of all biofluorophores meas-  
670 ured within this study, however the SIBS was able to detect the fluorescence signal at  $\lambda_{\text{ex}} = 370$  nm  
671 for all three (Fig. 9). The spectral difference between chlorophyll *a* and *b* is only minor at  $\lambda_{\text{ex}} = 370$   
672 nm ( $\Delta\lambda = 8.3$  nm) for which the spectral resolution of the SIBS is not capable of distinguishing be-  
673 tween both types (Fig. 9a, b, c, d and Fig. S11) (e.g., French et al., 1956; Welschmeyer, 1994). Nev-  
674 ertheless, the SIBS shows the ability to distinguish between chlorophyll *a* and *b*, and bacteriochloro-  
675 phyll due to the red-shift in the bacteriochlorophyll spectrum ( $\Delta\lambda = 28.5$  nm at  $\lambda_{\text{ex}} = 370$  nm, between



676 chlorophyll *a* and bacteriochlorophyll). This may provide a further discrimination level regarding  
677 algae, plant residues, and cyanobacteria. Bacteriochlorophyll also shows a second and even stronger  
678 emission peak at  $\lambda_{\text{ex}} = 370$  nm ( $\lambda_{\text{em}} = \sim 800$  nm) that could help further distinguish it from chlorophyll  
679 *a* and *b*, but the SIBS spectrometer cannot currently detect this far into the IR (e.g., Rijgersberg et  
680 al., 1980; Van Grondelle et al., 1983).

681 Overall, fluorescence emissions recorded by the SIBS are in good agreement with measured ref-  
682 erence spectra. However, care must be taken as to the interpretation of fluorescence emissions cov-  
683 ering broad spectral ranges, which span regimes with large differences between individual correction  
684 factors (e.g., channel 15 ( $\lambda_{\text{mean}} = 693.9$  nm, Fig.7l) and channel 2 ( $\lambda_{\text{mean}} = 330.6$  nm, Fig.8k). For the  
685 SIBS, namely the first two UV detection channels and the last two near-IR channels have to be treated  
686 with care. Further investigations are required for a careful assessment of how the spectral correction  
687 can be applied properly onto fluorescent and non-fluorescent atmospheric particles.

688

#### 689 **4.5 Particle asymmetry measurements**

690 The AF of spherical particles such as PSLs (Fig. 10a, b) and ultrapure water droplets is approximately  
691 10 (Table 3), which is slightly higher than reported values for spherical particles by, e.g., Savage et  
692 al. (2017) (AF= ~5) or Toprak & Schnaiter (2013) (AF= ~8) using a WIBS. It is noteworthy that the  
693 AF of water droplets increases slightly with increasing droplet size and, therefore, contributes to the  
694 mean value (Fig. S13). This effect is most likely based on a decreasing surface tension with increas-  
695 ing droplet size for which the droplet morphology is changed to a more oval shape within the sample  
696 flow. A similar effect regarding a potential droplet deformation using an Airborne Particle Classifier  
697 (APC) was observed by Kaye et al., (1991). Even if the morphology of ammonium sulfate (crystal-  
698 line, Fig. 10d) and Fe<sub>3</sub>O<sub>4</sub> (irregular clusters, Fig. 10f) is diverse, their differences in AFs is only  
699 minor (~13 and 14, Table 3), indicating that most naturally occurring aerosols (e.g., sea salt, soot,  
700 various bacterial and fungal clusters) will occur in a AF regime between ~10 and 20. Only rod-shaped  
701 carbon nanotubes (110-170 nm diameter, 5-9  $\mu\text{m}$  length) show increased AFs with a mean value at  
702 ~22 (Table 3) at which also, e.g., bacteria would occur (Fig. 10h). No particles observed exhibited  
703 average AF values >25, as would have been expected for, e.g., carbon nanotubes. Because the range  
704 of AF values for homogenous particles is relatively broad and the differences between morphologi-  
705 cally diverse particle types is only minor (Table 3), the question can be raised to what extent particles  
706 could be distinguished based on the AF under ambient conditions. As also discussed by Savage et al.  
707 (2017), the AF values reported by SIBS and WIBS units should be treated with extreme care.

708 The validation of asymmetry measurements is challenging due to unavoidable particle and aero-  
709 solization effects (e.g., particle agglomeration and spatial orientation within the sample flow) and the





710 lack of standardized procedures for AF calibrations. Measurements performed in this study do, there-  
711 fore, only serve as a rough AF assignment. Moreover, even if both the SIBS and WIBS use the same  
712 technical components for defining AFs, a direct intercomparison cannot be applied due to technical  
713 variability (e.g., PMT related signal-to-noise ratio or the alignment of optical components). Addi-  
714 tionally, it is currently unknown in how far the 785 nm diode laser of the SIBS affect asymmetry  
715 measurements compared to the WIBS using a 635 nm diode laser.

716

#### 717 4.6 Initial ambient measurements

718 Several weeks of initial ambient SIBS measurements have been conducted on the roof of the Max  
719 Planck Institute for Chemistry in Mainz, Germany. At the same location, Huffman et al. (2010) con-  
720 ducted one of the first ambient UV-APS studies. Moreover, Toprak & Schnaiter (2013) conducted a  
721 WIBS-4A study at a comparable site in central Germany. The aim of this brief ambient section is to  
722 validate that the SIBS-derived key aerosol and fluorescence data are consistent with the aforemen-  
723 tioned studies. We found a good agreement between the coarse mode ( $\geq 1 \mu\text{m}$ ) number concentrations  
724 ( $N_{T,c}$ ) of the SIBS ( $N_{T,c}$  ranging from 0.19 to 1.24  $\text{cm}^{-3}$ , with a mean of 0.59  $\text{cm}^{-3}$ ), the UV-APS (mean  
725  $N_{T,c}$ : 1.05  $\text{cm}^{-3}$  (Huffman et al., 2010)), and the WIBS-4A (mean  $N_{T,c}$ : 0.58  $\text{cm}^{-3}$  (Toprak and  
726 Schnaiter, 2013)) (Fig.11a). Furthermore, good agreement was found between coarse mode fluores-  
727 cent number concentrations ( $N_{F,c}$ ) of the SIBS with a  $3\sigma$  SD threshold (mean  $N_{F,c(3\sigma)}$ : 0.019  $\text{cm}^{-3}$ ),  
728 the UV-APS (mean  $N_{F,c}$ : 0.027  $\text{cm}^{-3}$  (Huffman et al., 2010)), and the WIBS-4A with a  $3\sigma$  SD thresh-  
729 old (mean  $N_{F,c(3\sigma)}$ : 0.031  $\text{cm}^{-3}$  (Toprak and Schnaiter, 2013)) (Fig.11a). Similarly, the fraction of  
730 fluorescent particles in the coarse mode ( $N_{F,c}/N_{T,c}$ ) compares well between SIBS with a  $3\sigma$  SD thresh-  
731 old (mean  $N_{F,c(3\sigma)}/N_{T,c}$ : 4.2 %), the UV-APS (mean  $N_{F,c}/N_{T,c}$ : 3.9 % (Huffman et al., 2010)), and the  
732 WIBS-4A with a  $3\sigma$  SD threshold (mean  $N_{F,c(3\sigma)}/N_{T,c}$ : 7.3 % (Toprak and Schnaiter, 2013)) (Fig.11b).  
733 Expectedly, a  $1\sigma$  SD threshold gives much higher fluorescent fractions of 39.2 %, whereas a  $6\sigma$  SD  
734 threshold corresponds with much lower fluorescent fractions of 1% (Fig.11b). Note that no perfect  
735 match between our results and the studies by Huffman et al. (2010), and Toprak & Schnaiter (2013)  
736 can be expected, since the measurements took place with different sampling setups and during dif-  
737 ferent seasons. Furthermore, the spectrally resolved SIBS data makes the definition of fluorescent  
738 fraction more complex than for UV-APS and WIBS data (see Sect. 2.6). However, the overall good  
739 agreement confirms that the SIBS produces reasonable results in an ambient setting. Further, the  
740 single particle fluorescence spectra are reasonable with respect to typical biofluorophore emissions  
741 (Pöhlker et al., 2012). Exemplary spectra ( $\lambda_{\text{ex}} = 285$  and 370 nm) of ambient single particles can be  
742 found in the supplement (Fig.S14). An in-depth analysis of extended SIBS ambient datasets is subject  
743 of ongoing work.

744





## 745 5. Summary and conclusions

746 Real-time analysis of atmospheric bioaerosols using commercial LIF instruments has largely been  
747 restricted to data recorded in only 1-3 spectrally integrated emission channels, limiting the interpre-  
748 tation of fluorescence information. Instruments that can record resolved fluorescence spectra over a  
749 broad range of emission wavelengths may thus be required to further improve the applicability of  
750 LIF instrumentation to ambient PBAP detection. Introduced here is the SIBS as a new aerosol fluo-  
751 rescence detector, which provides resolved fluorescence spectra ( $\lambda_{\text{mean}} = 302 - 721 \text{ nm}$ ) from each  
752 of two excitation wavelengths ( $\lambda_{\text{ex}} = 285 \text{ and } 370 \text{ nm}$ ) for single particles. The current study intro-  
753 duces the SIBS by presenting and experimentally validating its key functionalities. This work criti-  
754 cally assesses the strengths and limitations of the SIBS with respect to the growing interest in real-  
755 time bioaerosol quantification and classification. It should be noted that the study is an independent  
756 evaluation that was not conducted or endorsed by the manufacturer. Overall, this work confirms a  
757 precise particle sizing between 300 nm and 20  $\mu\text{m}$  and the particle discrimination ability based on  
758 spectrally resolved fluorescence information of several standard compounds.

759 The SIBS was operated at a low PMT detector amplification setting (610 V) to retain capacity to  
760 detect large or brightly fluorescent particles. It was confirmed, however, that even weak fluorescence  
761 signals from 0.53  $\mu\text{m}$  purple PSLs and from small *S. cerevisiae* fragments (0.5 - 1  $\mu\text{m}$ ) can be clearly  
762 distinguished from the background signal. Saturation events were only observed for the polysty-  
763 rene/detergent signal from relatively large 15 and 20  $\mu\text{m}$  PS-DVB particles. Nevertheless, the fluo-  
764 rescence intensity detection threshold is highly instrument-dependent due to the complex interaction  
765 of single technical components across individual instruments. For example, xenon 1 exhibited ~154  
766 % higher irradiance than xenon 2 (both new lamps) due to differences in the properties of xenon  
767 emission and the optical filters used. For used xenon lamps (> 4000 hours of use), an even higher  
768 difference of ~220 % was observed. Thus, a defined fluorescence detection threshold will most likely  
769 change over time due to, e.g., material fatigue. These observations are valid not only for the SIBS,  
770 but also for the WIBS-4A and WIBS-NEO and lead to important implications for interpretation of  
771 particle data. In particular, a particle that exhibits measurable fluorescence in WIBS channel FL1,  
772 but only weak fluorescence in channel FL3 could be assigned as an “A-type” particle in one instru-  
773 ment but an “AC-type” particle in an instrument with slightly stronger xenon 2 irradiance. These  
774 differences in classification can be extremely important to interpretation of ambient data (e.g., Savage  
775 et al., 2017; Perring et al. 2015).

776 The PMT used in the SIBS shows a wavelength-dependent sensitivity distribution along all 16  
777 detection channels. To compensate for this characteristic and to be able to use the broadest possible  
778 fluorescence emission range, the measured emission spectra were corrected with respect to reference  
779 spectra acquired from deuterium and halogen lamps. A spectral correction over a broad emission



780 range also introduces drawbacks, however, that LIF-instrument users should keep in mind while in-  
781 terpreting derived fluorescence information. In particular, the first two (UV) and the last two (near-  
782 IR) detection channels should be treated with care, because they have larger correction factors com-  
783 pared to adjacent channels. Ultimately, the correction factor and amplification voltages applied to the  
784 detector will be experiment-specific and will need to be investigated with respect to individual ex-  
785 perimental aims.

786 Fluorescence spectra of fluorescent PSLs, amino acids, co-enzymes, and flavins measured by the  
787 SIBS agree well with corresponding spectra recorded with an offline reference spectrometer. Thus,  
788 the SIBS was shown to be capable of clearly distinguishing between different particle types based on  
789 resolved fluorescence information. Furthermore, the extended fluorescence emission range ( $\lambda_{em} = >$   
790 700 nm) enables the SIBS also to distinguish chlorophyll *a* and *b* from bacteriochlorophyll, poten-  
791 tially opening new possibilities for the detection of, e.g., algae, plant residues, and cyanobacteria.

792 Particle asymmetry measurements revealed that spherical PSLs have an AF of  $9.9 (\pm 3.6)$ ,  
793 whereas other materials (ammonium sulfate,  $Fe_3O_4$ , and carbon nanotubes) show AF values of  $13.1$   
794  $(\pm 8.1)$ ,  $14.4 (\pm 7.4)$ , and  $21.6 (\pm 12.7)$ , respectively. Because differences of measured AF value  
795 between morphologically diverse particle types are small and within the ranges of uncertainty for the  
796 measurement of a given set of particles, it is questionable how well particles can be distinguished  
797 based on the AF by the quadrant PMT as presently measured. Users of SIBS and WIBS instruments  
798 should apply extreme care if using AF data. It is also likely that different instrument units may have  
799 very different AF responses with respect to this measurement. At a minimum, each individual unit  
800 needs to be rigorously calibrated to known particle types to determine if AF values are sufficiently  
801 different (e.g., separated by several standard deviations) to justify scientific conclusions based on the  
802 metric.

803 Exemplary ambient data, measured between the 12<sup>th</sup> and 18<sup>th</sup> of April 2018 on the roof of the  
804 Max Planck Institute for Chemistry in Mainz (Germany), are consistent with LIF measurement data  
805 using a UV-APS (Huffman et al., 2010) and a WIBS-4A (Toprak and Schnaiter, 2013). Total coarse  
806 particle number concentrations revealed a mean value of  $0.59 \text{ cm}^{-3}$  ( $1.05 \text{ cm}^{-3}$  (Huffman et al., 2010);  
807  $0.58 \text{ cm}^{-3}$  (Toprak and Schnaiter, 2013)) of which  $\sim 4.2\%$  are considered to be fluorescent using a  $3\sigma$   
808 SD threshold ( $3.9\%$  (Huffman et al., 2010);  $7.3\%$  (Toprak and Schnaiter, 2013)), including only par-  
809 ticles that show fluorescence emission in, at least, two adjoining detection channels. Using a  $1\sigma$  and  
810  $6\sigma$  SD threshold results in fluorescent fractions of  $39.2\%$  and  $1\%$  respectively. However, the applica-  
811 bility of different threshold strategies for the SIBS is currently under investigation and needs further  
812 verifications.

813 The results suggest that the SIBS has the ability to increase the selectivity of detection of fluo-  
814 rescent biological and non-biological particles by use of two excitation wavelengths and 16-channel



815 resolved fluorescence information in combination with a broad detectable emission range. The ap-  
816 plicability of described methods onto ambient datasets is currently under investigation. Data shown  
817 here and the detailed insights of technical components used in the SIBS will be broadly beneficial  
818 for users of LIF instruments providing resolved fluorescence information, but also for users of vari-  
819 ous generations of WIBS and other LIF instruments widely applied within the bioaerosol community.  
820

## 821 **6. Data availability**

822 The data of the key results presented here can be provided upon request. For specific data requests,  
823 please refer to the corresponding authors.

824

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- 1068

1069 **Appendix A:** List of acronyms and symbols.

Acronym/Symbol	Description
<b>AF</b>	Asymmetry factor
<b>APC</b>	Airborne Particle Classifier
<b>CCD</b>	Charge-coupled device
<b>DMT</b>	Droplet Measurement Technologies
<b>EEM</b>	Excitation-emission matrix
<b>EM</b>	Emission
<b>EX</b>	Excitation
<b>IR</b>	Infrared
<b>LIF</b>	Light-induced fluorescence
$N$	Particle number concentration ( $\text{cm}^{-3}$ )
$N_{T,c}$	$N$ of total coarse particles (1-20 $\mu\text{m}$ )
$N_{F,(no)}$	$N$ of fluorescent coarse particles (1-20 $\mu\text{m}$ ) at 1, 3, or 6 $\sigma$
<b>NAD</b>	Nicotinamide adenine dinucleotide
<b>NAD(P)H</b>	Nicotinamide adenine dinucleotide and nicotinamide adenine dinucleotide phosphate
<b>NIST</b>	National institute of standards and technology
<b>PBAP</b>	Primary biological aerosol particles
<b>PMT</b>	Photomultiplier tube
<b>PAH</b>	Polycyclic aromatic hydrocarbons
<b>PSL</b>	Polystyrene latex sphere
<b>PS-DVB</b>	Polystyrene-divinylbenzene
<b>SD</b>	Standard deviation
<b>SIBS</b>	Spectral intensity bioaerosol sensor
<b>SNR</b>	Signal to noise ratio
<b>TSP</b>	Total suspended particles
<b>UV</b>	Ultraviolet
<b>UV-APS</b>	Ultraviolet aerodynamic particle sizer
<b>Vis</b>	Visible light
<b>WIBS</b>	Wideband integrated bioaerosol sensor

1070



1071 **Table 1.** Lower, mean, and upper wavelength at each PMT detection channel. Nominal data accord-  
1072 ing to manufacturer Hamamatsu.

Channel	$\lambda_{\text{lower}}$ (nm)	$\lambda_{\text{mean}}$ (nm)	$\lambda_{\text{upper}}$ (nm)
1	298.2	302.2	316.2
2	316.6	330.6	344.6
3	345.0	359.0	362.5
4	377.5	387.3	401.3
5	401.5	415.6	429.7
6	429.8	443.8	457.8
7	457.9	471.9	485.9
8	486.0	500.0	514.0
9	514.0	528.0	542.0
10	541.9	555.9	569.9
11	569.7	583.7	597.7
12	597.4	611.4	625.4
13	625.0	639.0	653.0
14	652.8	666.5	680.2
15	679.9	693.9	707.9
16	707.1	721.1	735.1



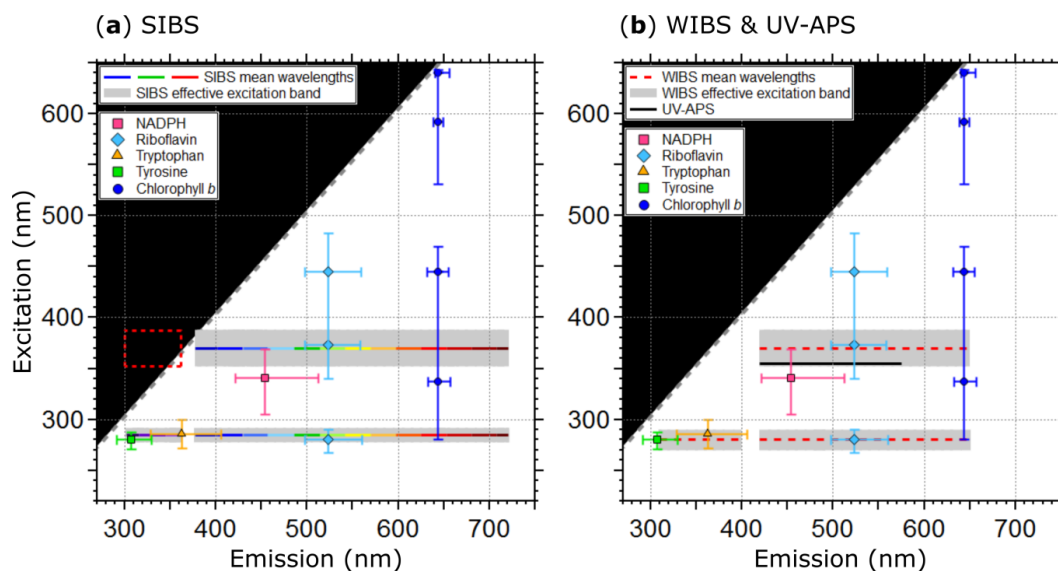
1073 **Table 2.** Parameters and technical components of the SIBS in comparison to the WIBS-NEO and  
 1074 WIBS-4A. Data are taken from manufacturer information.

	SIBS	WIBS-NEO	WIBS-4A
<b>Measured parameters</b>	Particle size Asymmetry Factor Fluorescence spectra	Particle size Asymmetry Factor Integrated fluorescence in 3 channels	Particle size Asymmetry Factor Integrated fluorescence in 3 channels
<b>Particle size range</b>	~0.3 – 100 $\mu\text{m}$	~0.3 – 100 $\mu\text{m}$	~0.5 – 20 $\mu\text{m}$
<b>Maximum concentration</b>	$\sim 2 \times 10^4$ particles/L	$\sim 2 \times 10^4$ particles/L	$\sim 2 \times 10^4$ particles/L
<b>Fluorescence excitation</b>	$\lambda_{\text{ex}} = 285$ and $\lambda_{\text{ex}} = 370$ nm	$\lambda_{\text{ex}} = 280$ and $\lambda_{\text{ex}} = 370$ nm	$\lambda_{\text{ex}} = 280$ and $\lambda_{\text{ex}} = 370$ nm
<b>Fluorescence emission</b>	$\lambda_{\text{em}} = 302 - 721$ nm (16-channel PMT)	$\lambda_{\text{em}} = 310-400$ nm and $\lambda_{\text{em}} = 420-650$ nm	$\lambda_{\text{em}} = 310-400$ nm and $\lambda_{\text{em}} = 420-650$ nm
<b>Flow rate</b>	Sample flow: ~0.3 l/min Sheath flow: ~2.2 l/min (re-circulating)	Sample flow: ~0.3 l/min Sheath flow: ~2.2 l/min (re-circulating)	Sample flow: ~0.3 l/min Sheath flow: ~2.2 l/min (re-circulating)
<b>Laser</b>	785 nm diode laser, 55 mW	635 nm diode laser, 15 mW	635 nm diode laser, 12 mW
<b>Pump</b>	Diaphragm pump	Diaphragm pump	Diaphragm pump
<b>Power requirements</b>	200 W, 90 - 230 VAC	150 W, 90 - 230 VAC	150 W, 90 - 230 VAC
<b>Weight (kg)</b>	20.1	12.5	13.6
<b>Dimension W x L x H (cm)</b>	42.5 x 61.5 x 23.5	45.1 x 36.2 x 24.1	30.4 x 38.2 x 17.1



1075 **Table 3.** Asymmetry factor (AF) values for reference particles. Values are based on the mean of a  
1076 Gaussian fit applied onto each particle histogram (see also Fig. 10), including  $1\sigma$  SD.

	<b>AF</b>
<b>2 <math>\mu\text{m}</math> non-fluorescent PSLs</b>	$9.9 \pm 3.6$
<b>Ultrapure water</b>	$11.9 \pm 2.9$
<b>Ammonium sulfate</b>	$13.1 \pm 8.1$
<b>Fe<sub>3</sub>O<sub>4</sub></b>	$14.4 \pm 7.4$
<b>Carbon nanotubes</b>	$21.6 \pm 12.7$

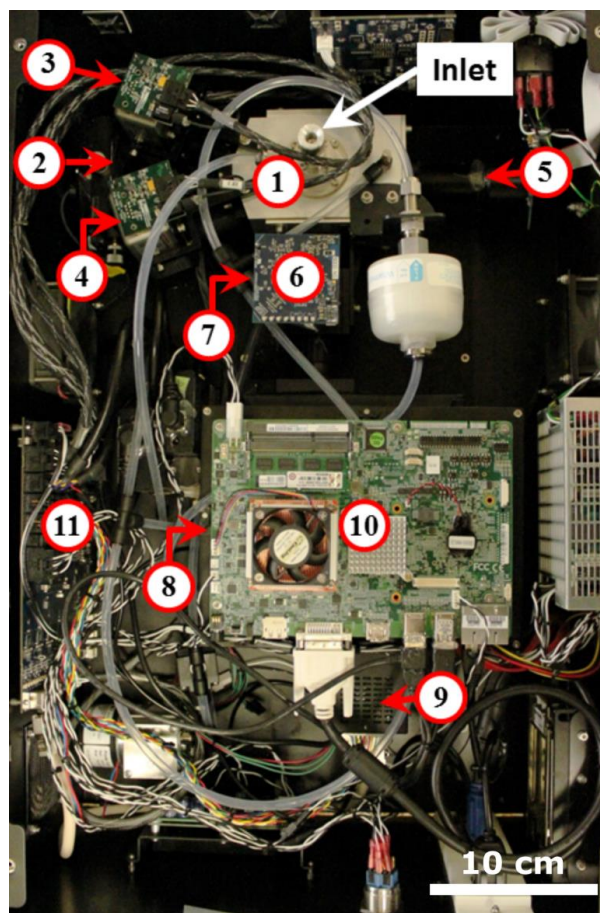


1077 **Figure 1.** Optical design and overview of excitation and emission specifications of the LIF instru-  
 1078 ments UV-APS, WIBS, and the SIBS with spectral locations of the autofluorescence modes of the  
 1079 biofluorophores tyrosine, tryptophan, NAD(P)H, riboflavin, and chlorophyll *b* (as examples). Here  
 1080 the term WIBS includes the WIBS-4A and WIBS-NEO, because both instruments use the same op-  
 1081 tical components. Spectral properties of the emission bands of LIF instruments are illustrated as hor-  
 1082 izontal lines. The color-coded bars in (a) illustrate the spectrally resolved fluorescence detection of  
 1083 the two excitation wavelengths ( $\lambda_{\text{ex}} = 285$  and  $370$  nm) by the SIBS. The “blind spot” (white notch)  
 1084 at  $\lambda_{\text{ex}} = 285$  nm between  $\lambda_{\text{em}} = 362 - 377$  nm (a) originates from a notch optical filter, used to block  
 1085 incident light from the excitation sources. Grey dashed lines show the 1<sup>st</sup> order elastic scattering. At  
 1086  $\lambda_{\text{ex}} = 370$  nm, the detection range of the SIBS includes the spectral range where  $\lambda_{\text{em}} < \lambda_{\text{ex}}$ , for which  
 1087 fluorescence is not defined and so data within the red dashed rectangle is omitted (a). Grey bars  
 1088 indicate the effective excitation bands of optical filters used for the WIBS and SIBS (see also Sect.  
 1089 3.3 and Fig. 3). The effective excitation bands in the WIBS and SIBS occur in a spectral range span-  
 1090 ning several nanometers (up to 36 nm), in contrast to the UV-APS (black line, b), which uses a laser  
 1091 source with a defined excitation (Figure adapted from Pöhlker et al., 2012).

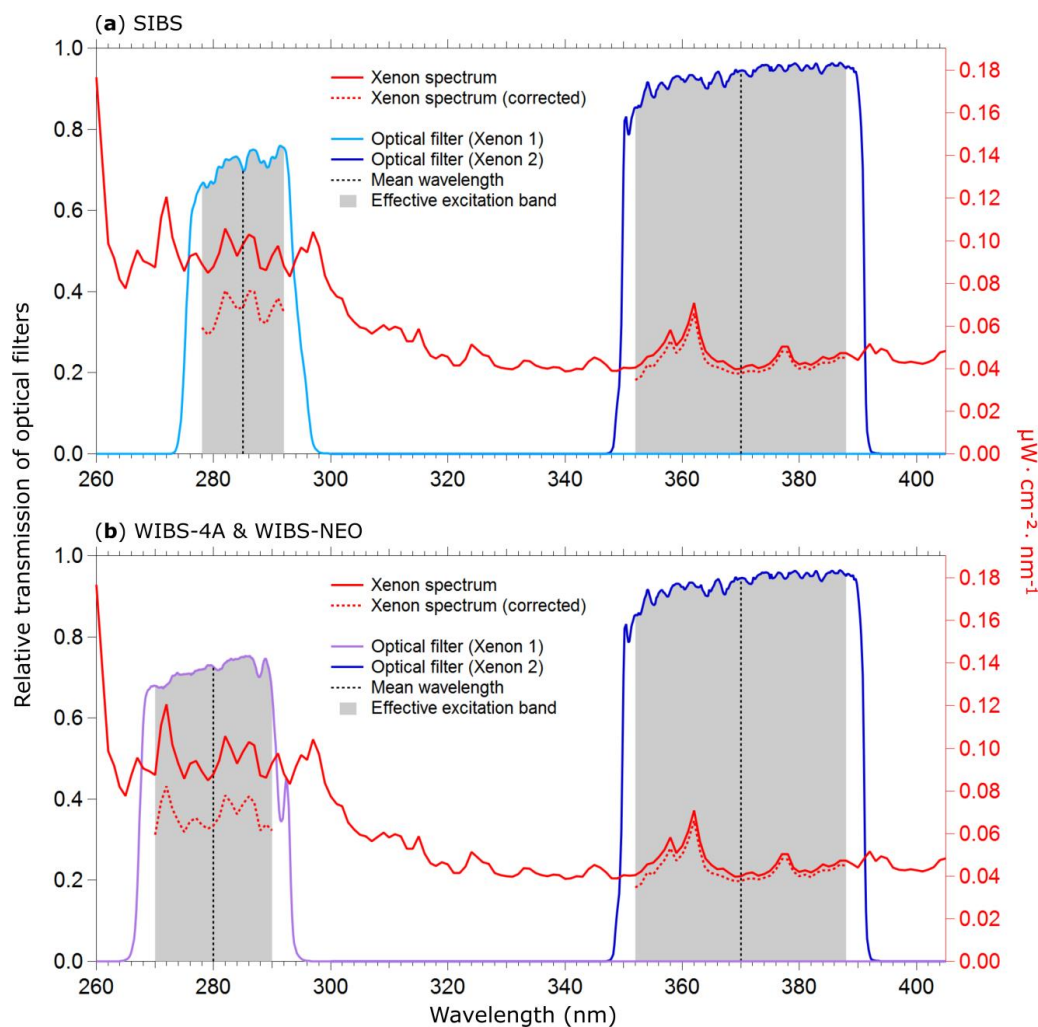




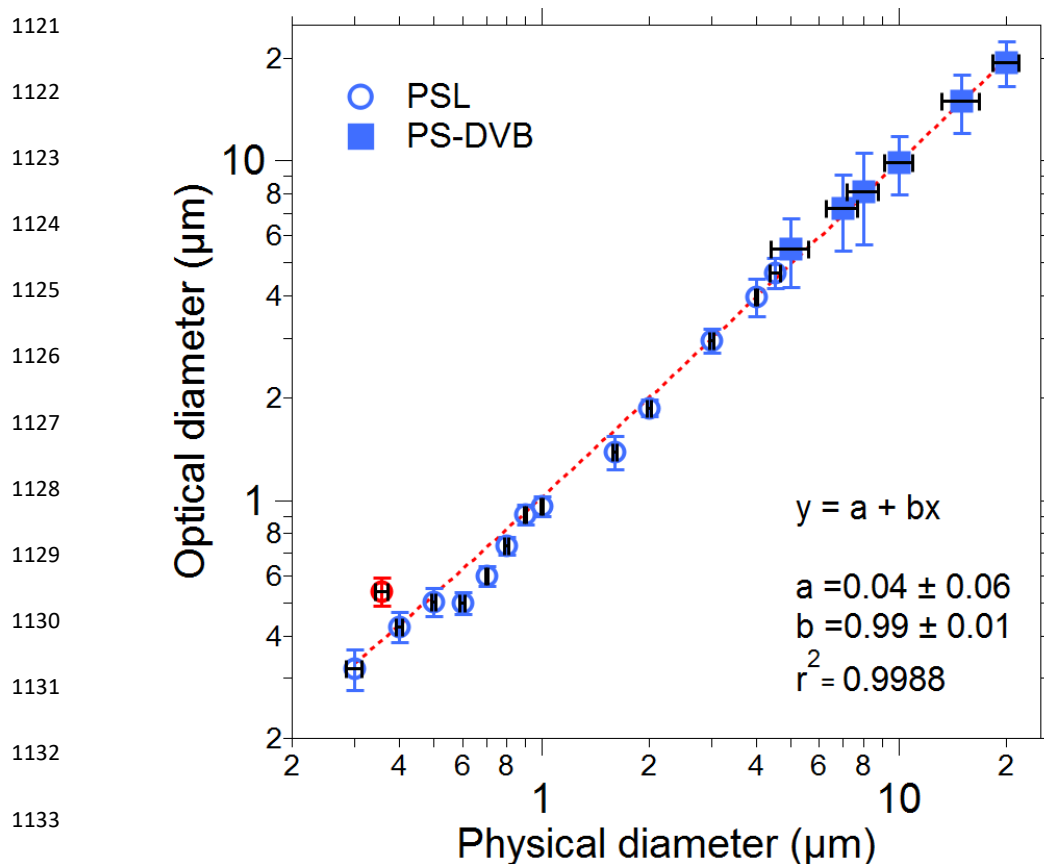
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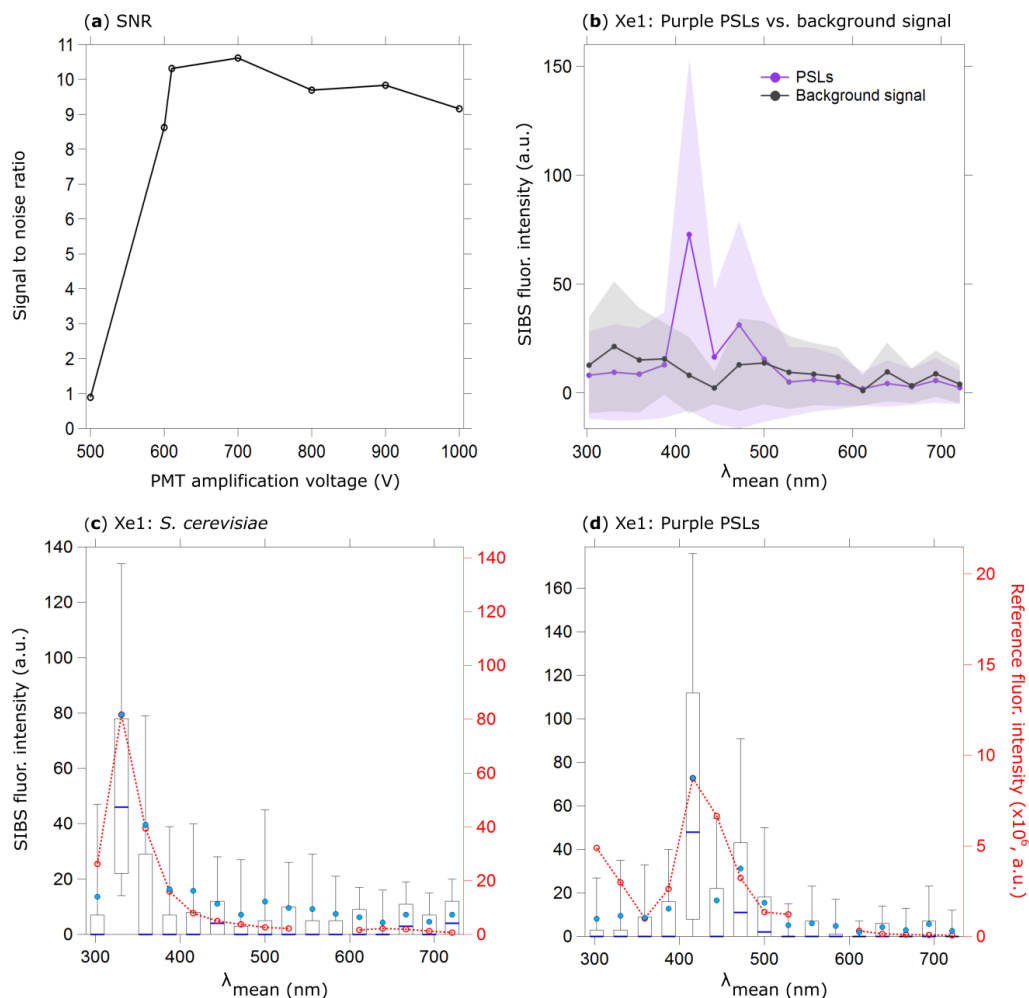
1106 **Figure 2.** Technical components within SIBS body. (1) Optical cavity. (2) Continuous wave diode  
1107 laser used for particle detection and sizing. (3) and (4) Xenon light sources. (5) Quadrant PMT used  
1108 for the determination of particle asymmetry. (6) PMT used for particle detection and sizing. (7) Di-  
1109 chroic beamsplitter separates side-scattered light (particle sizing) and fluorescence emission (not vis-  
1110 ible; below component (6)). (8) Grating polychromator (below component (10)). (9) 16-channel PMT  
1111 used for detection of fluorescence. (10) Embedded computer unit. (11) Control-board.



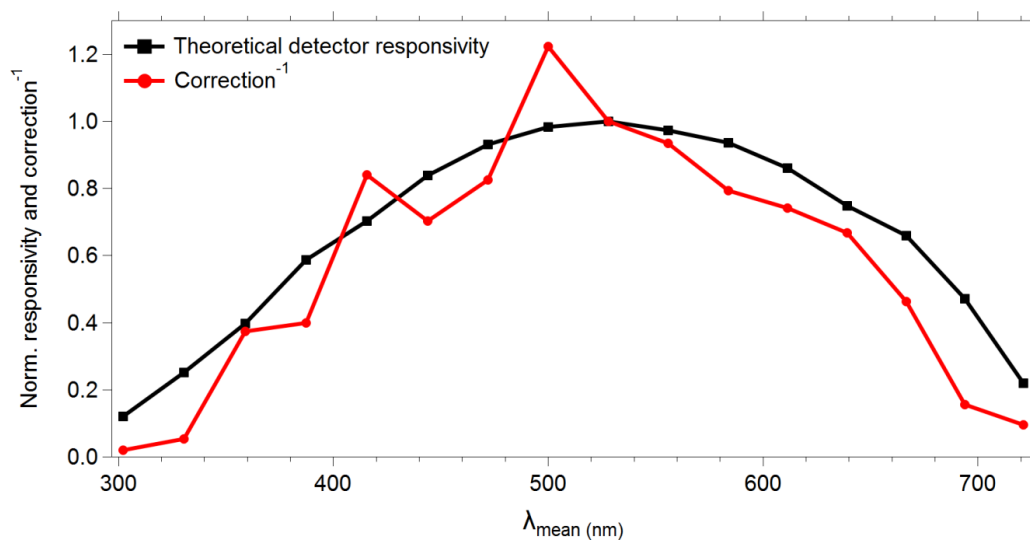
1112 **Figure 3.** Irradiance from xenon flash lamps based on specifications of lamps and optical filters.  
 1113 Purple and blue lines show optical transmission of filters (left axes) applied to select excitation wave-  
 1114 length. Gray bands indicate where filter transmit light relative from the mean wavelength. Red lines  
 1115 show theoretical irradiance values of the xenon flash lamp (right axes): solid line (raw output), dashed  
 1116 line (relative output after filtering). Relative output shown as raw output multiplied by effective ex-  
 1117 citation band of the bandpass filters used in the: (a) SIBS ( $\Delta\lambda_{\text{ex}}(\text{Xenon1}) = \sim 14$  nm;  $\Delta\lambda_{\text{ex}}(\text{Xenon2}) = \sim 36$   
 1118 nm), and (b) WIBS-4A and WIBS-NEO ( $\Delta\lambda_{\text{ex}}(\text{Xenon1}) = \sim 20$  nm;  $\Delta\lambda_{\text{ex}}(\text{Xenon2}) = \sim 36$  nm). Xenon lamp  
 1119 operating conditions: 600 V main voltage, 0.22  $\mu\text{F}$  main capacitance, 126 Hz repetition rate, 500 mm  
 1120 distance. (Data courtesy: Xenon flash lamps / Hamamatsu; Single-band bandpass filters / Semrock).



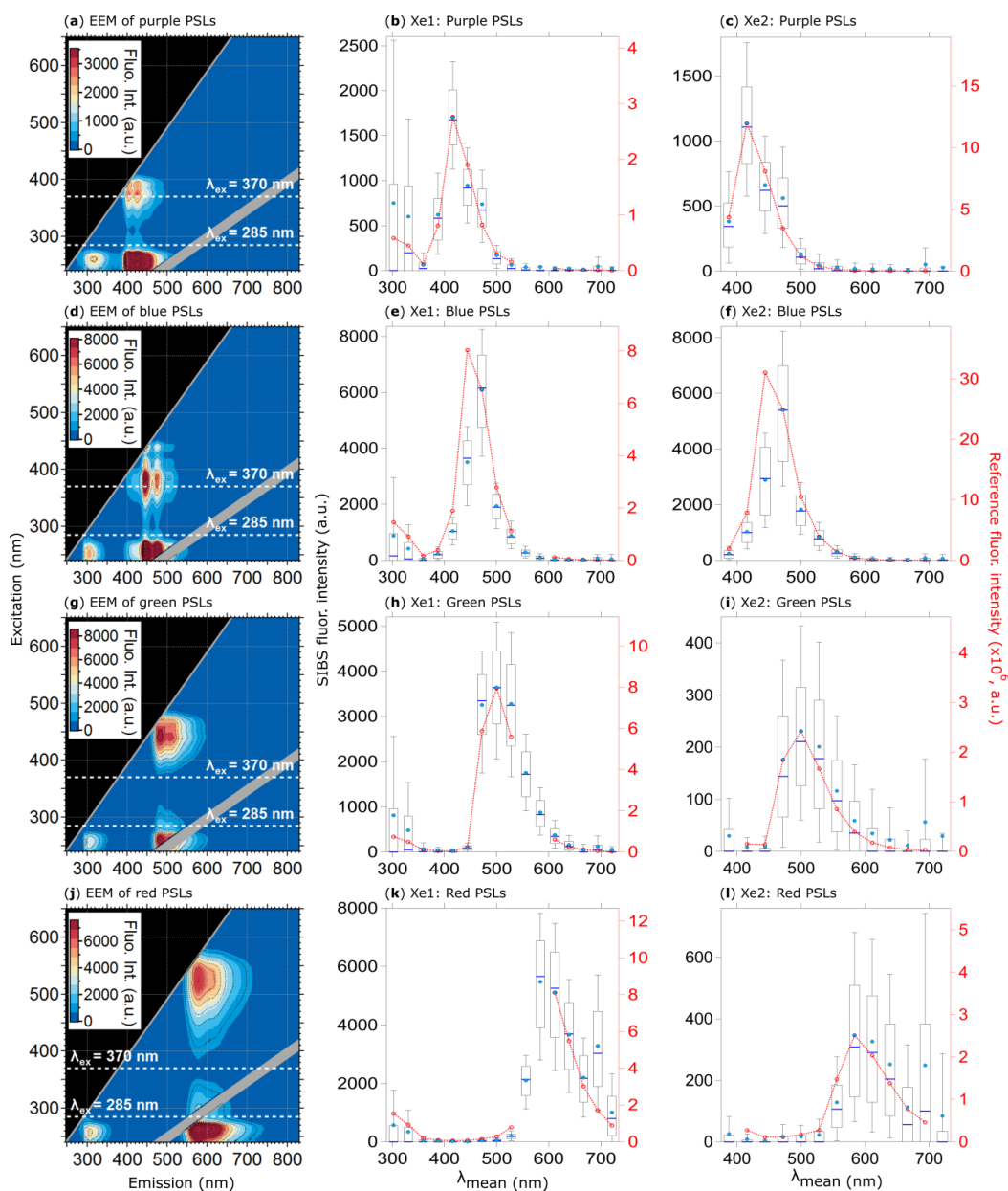
1134 **Figure 4.** Size calibration of SIBS. Black horizontal bars indicate  $1\sigma$  SD as stated by each manufac-  
1135 turer (Table S1). Optical diameter values and related  $1\sigma$  SD are based on a Gaussian fit, which was  
1136 used to average size distributions of several thousand homogeneous particles for each measurement.  
1137 The linear fit (red dashed line) excludes the  $0.356\ \mu\text{m}$  PSL sample (red marker), an outlier potentially  
1138 caused by a poor quality PSL batch. Only non-fluorescent particle standards were used for determin-  
1139 ing the sizing accuracy.



1140 **Figure 5.** SIBS signal to noise ratio (SNR) in (a): emission of 0.53  $\mu\text{m}$  purple PSLs (5260 particles,  
 1141 background signal + 1 $\sigma$  SD subtraction) divided by background signal at different PMT amplification  
 1142 voltages (both at Xe1, channel 5, averaged, and uncorrected). Background signal measured over 5  
 1143 min. In (b), fluorescence emission in contrast to background signal at a PMT amplification voltage  
 1144 of 610 V are shown (same parameters as in (a)). Shaded area: 1 $\sigma$  SD. Fluorescence intensity values  
 1145 are shown in arbitrary units. Fluorescence emission spectra of (c) *S. cerevisiae* (yeast; 2048 particles,  
 1146 0.5 – 1  $\mu\text{m}$ ) and (d) PSLs (as in (b)). Red dashed lines and markers (right axes) show averaged and  
 1147 re-binned reference spectra. Box and whisker plots (left axes) show SIBS spectra: median (blue line),  
 1148 mean (circle), boxes 75 and 25 percentile, whiskers 90 and 10 percentile. Data coinciding with 1<sup>st</sup> or  
 1149 2<sup>nd</sup> order elastic scattering were removed from reference spectra.

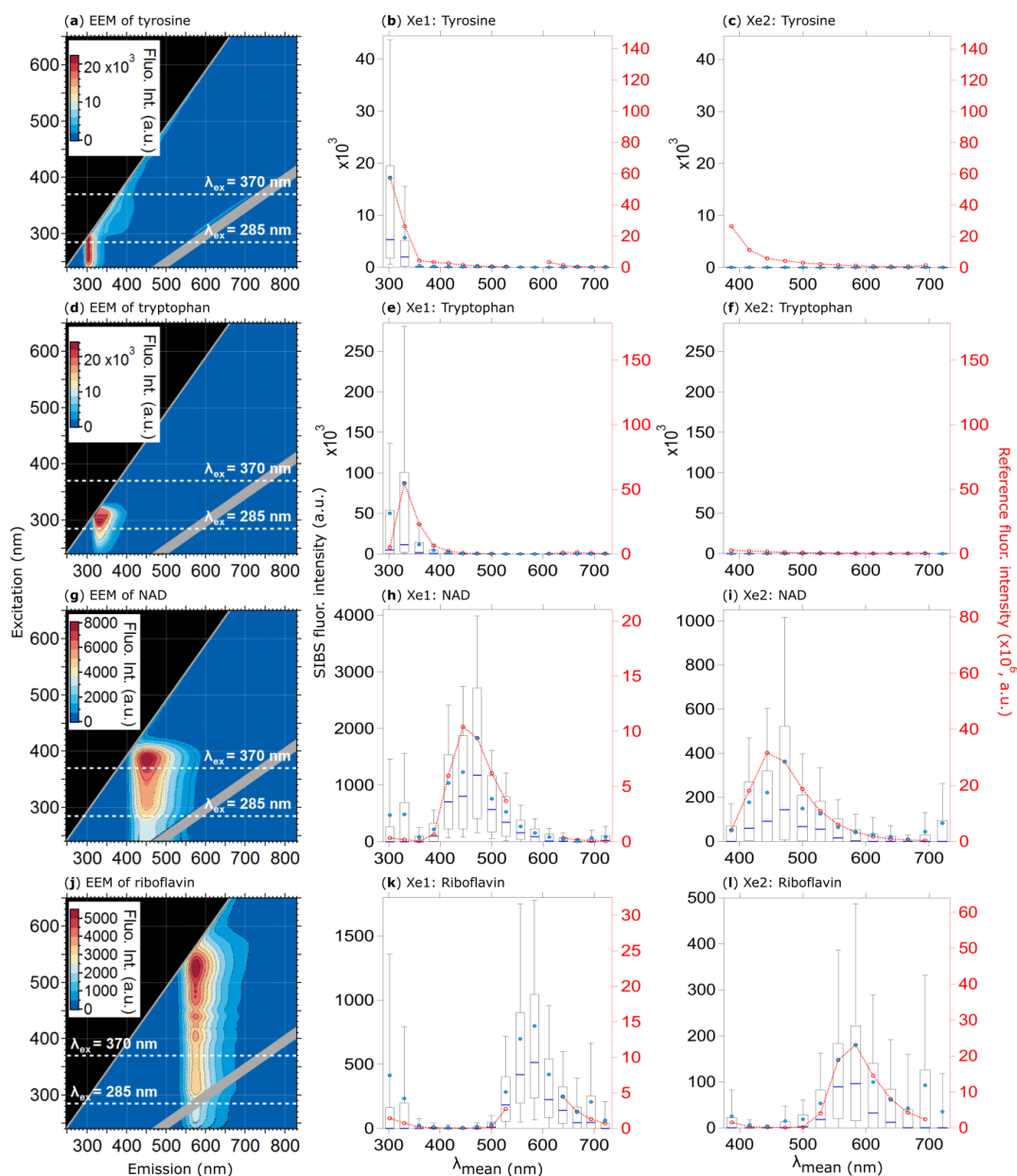


1150 **Figure 6.** Normalized theoretical detector responsivity and spectral correction. Theoretical detector  
1151 responsivity derived from measured cathode radiant sensitivity multiplied by the diffraction effi-  
1152 ciency (as shown in Figure S8). Note that red line shows inverse of spectral correction to match  
1153 detector response.



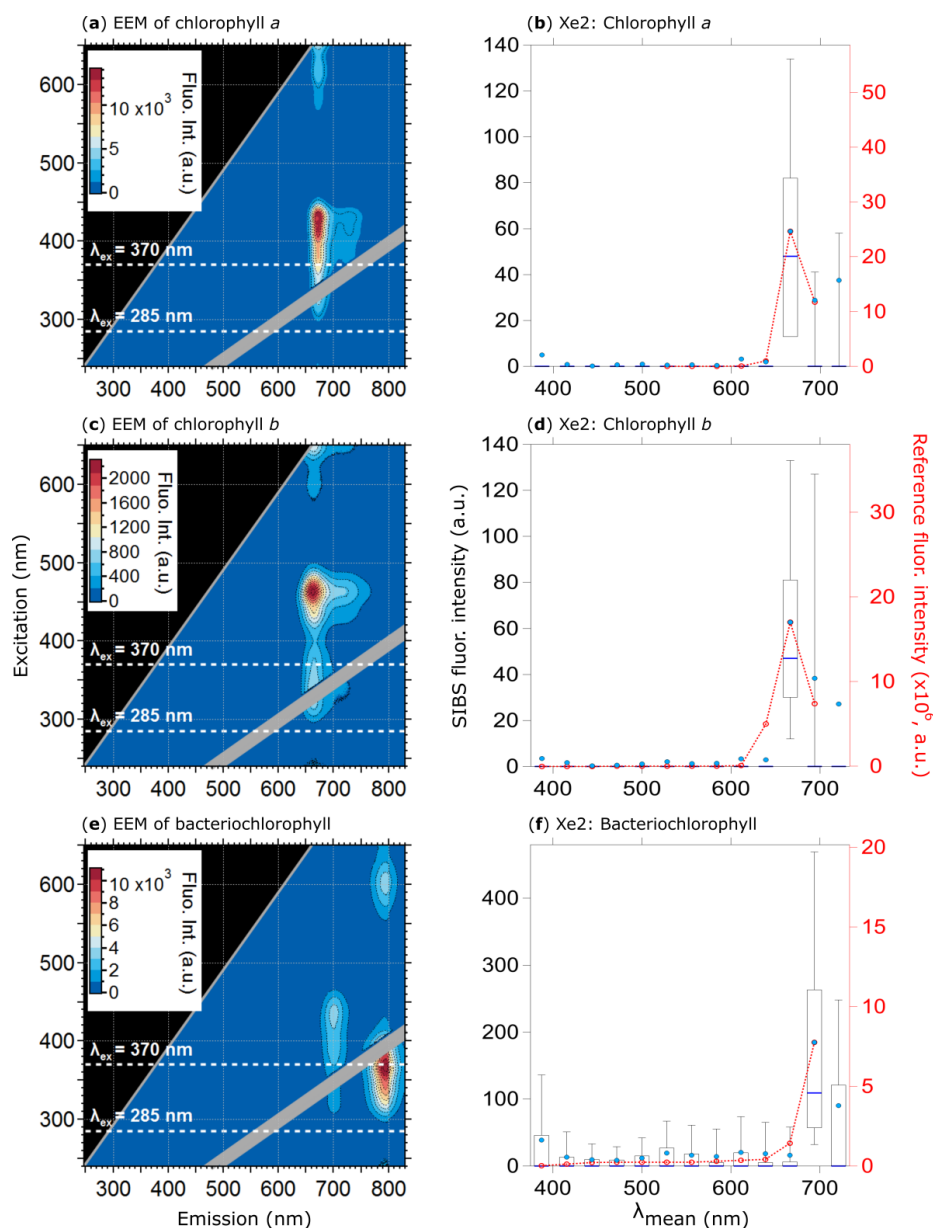
1154 **Figure 7.** Fluorescence emission spectra of PSLs. Steady-state fluorescence signatures displayed as  
 1155 EEMs (left column) and spectra at Xe1 and Xe2 (middle, right columns) for: 2.07  $\mu\text{m}$  purple (**a**, **b**  
 1156 and **c**, 1082 particles), 2.1  $\mu\text{m}$  blue (**d**, **e** and **f**, 1557 particles), 2  $\mu\text{m}$  green (**g**, **h**, and **i**, 1174 particles),  
 1157 and 2  $\mu\text{m}$  red PSLs (**j**, **k**, and **l**, 1474 particles). Within EEMs: white dashed lines show SIBS excita-  
 1158 tion wavelengths ( $\lambda_{\text{ex}} = 285$  and  $370$  nm), grey diagonal lines indicate 1<sup>st</sup> and 2<sup>nd</sup> order elastic scat-  
 1159 tering bands (both bands were subtracted automatically by the Aqualog V3.6 software).





1160 **Figure 8.** Fluorescence emission spectra of biofluorophores. EEMs (left column) and spectra at Xe1  
 1161 and Xe2 wavelengths (middle and right columns) shown for: tyrosine (a, b, and c, 209 particles),  
 1162 tryptophan (d, e, and f, 193 particles), NAD (g, h, and i, 376 particles), and riboflavin (j, k, and l,  
 1163 205 particles). All biofluorophores were size-selected between 1 and 2  $\mu\text{m}$ .

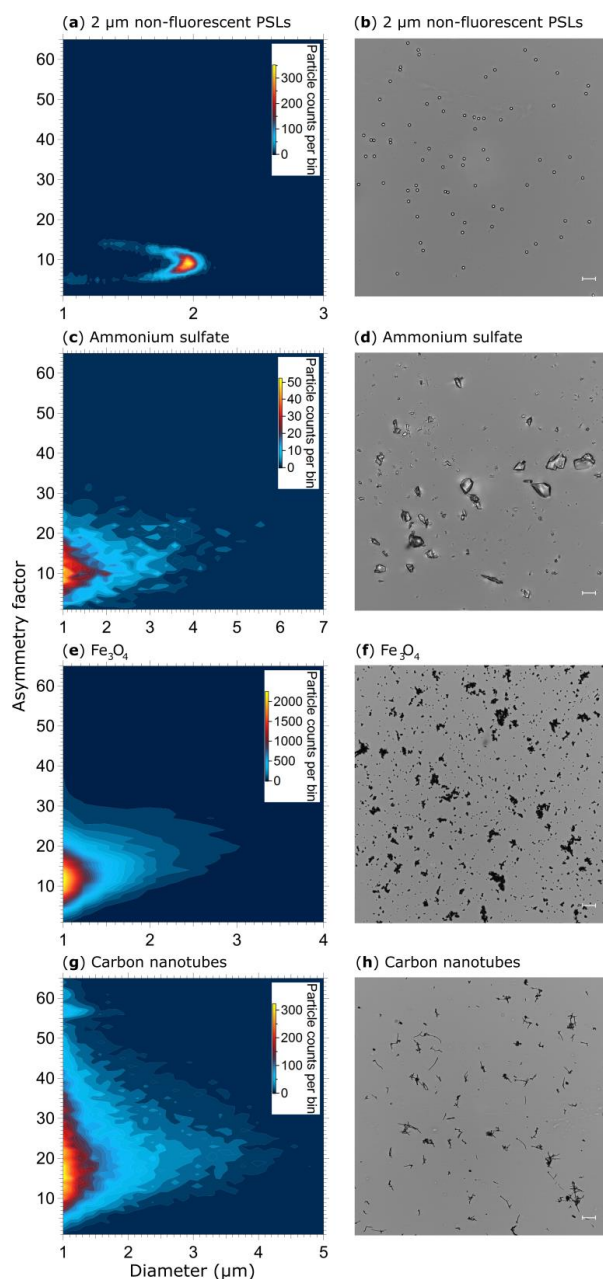




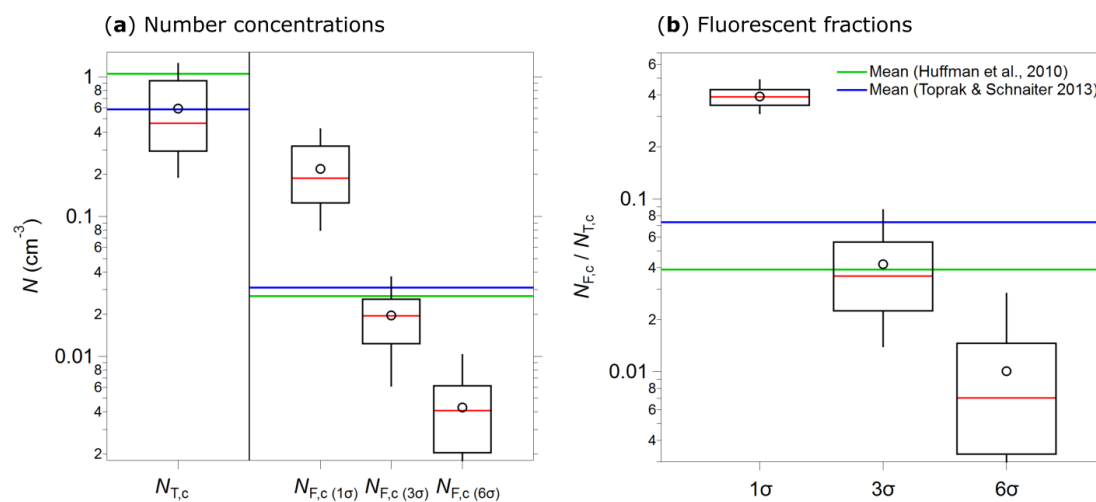
1164 **Figure 9.** Fluorescence emission spectra of three chlorophyll types. Highlighted are EEMs (left col-  
 1165 umn) and spectra at Xe2 (right columns) for: chlorophyll *a* (**a** and **b**, 370 particles), chlorophyll *b* (**c**  
 1166 and **d**, 585 particles), and bacteriochlorophyll (**e** and **f**, 633 particles). Size range chlorophyll *a* and  
 1167 *b*: 0.5 - 2  $\mu\text{m}$ , size range bacteriochlorophyll: 0.5 - 1  $\mu\text{m}$ . Emission spectra at Xe1 are excluded due  
 1168 to a fluorescence artifact caused by solved components from the polymer of the aerosolization bottles  
 1169 (Fig. S12).



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1199 **Figure 10.** Particle asymmetry. Shown are particle density histograms (left column) and microscopy  
1200 images (right column) for: 2  $\mu\text{m}$  non-fluorescent PSLs (**a** and **b**, 17836 particles), ammonium sulfate  
1201 (**c** and **d**, 3496 particles),  $\text{Fe}_3\text{O}_4$  (**e** and **f**, 65097 particles), and carbon nanotubes (56949 particles, **g**).  
1202 Scale bar (right column) indicates a length of 10  $\mu\text{m}$ .



1203 **Figure 11.** Integrated coarse particle (1-20  $\mu\text{m}$ ) number concentrations, measured between the 12<sup>th</sup>  
 1204 and 18<sup>th</sup> of April 2018 (5 min average), for total particles ( $N_{T,c}$ , fluorescent and non-fluorescent) and  
 1205 coarse fluorescent particles ( $N_{F,c}$ ) after 1, 3, and 6 $\sigma$  SD background signal subtraction **(a)**. The fluo-  
 1206 rescent fractions of integrated coarse particle number concentrations ( $N_{F,c} / N_{T,c}$ ) at 1, 3, and 6 $\sigma$  SD  
 1207 are shown in **(b)**. Median (red line), mean (black circles), boxes 75 and 25 percentile, whiskers 95  
 1208 and 5 percentile **(a and b)**. Data from Huffman et al. (2010) (green lines) and Toprak & Schnaiter,  
 1209 (2013) (blue lines) were taken for comparison **(a and b)**.