The authors have responded comprehensively to the first round of reviews. The manuscript is greatly improved and I believe that it could be published in ACP once the following comments have been addressed.

We are deeply grateful to the reviewer for his/her positive comments to our revised manuscript. The two anonymous reviewers in the first round of review are also greatly appreciated for providing the constructive comments which have helped us to improve the quality of our paper.

The level of English language writing is still an issue. There are a large number of grammatical errors. I recognize that the authors have made efforts to improve this, including the hiring of professional language editing help. However, the level is simply still too low. There are still far too many grammatical errors to list in a scientific review. I recommend that another round of professional copy editing is required.

Response:

Thanks for the comments. Another round of professional copy editing has been performed to our revised manuscript before it is submitted back to the journal. The changes are marked with different colors in the revision.

The authors have scaled their Aqauadag calibrations by a size-independent factor to make them more similar to fullerene calibrations, which are thought to be more representative of ambient BC. The use of a size-independent scaling factor may introduce error since the ratio of Aquadag to fullerene peak heights in the SP2 is size dependent (Laborde et al., 2012). To avoid this problem, Baumgardner et al., (2012) introduced a single-point scaling procedure. I would recommend the authors also use this approach to make their results consistent with other SP2 studies.

Response:

Thanks for the professional comments. We have recalculated the calibration curve for the incandescence signal of our SP2 using the single-point scaling procedure provided by Baumgardner et al. (2012), i.e., scaling the peak high of broadband incandescence signals for Aquadag with 300 nm mobility diameter (i.e., 8.9 fg rBC mass) by a factor

of 0.75, and axis intercept at zero. The derived calibration factor is 0.0039, very close to the value (0.0040) we used in the current manuscript. This deviation (~2.5%) is even lower than that (~3%) during the calibrations performed pre- and post-campaign. The negligible difference in the estimated calibration factor should influence few on the results presented in the current manuscript. If the calibration factor is changed, there should be a massive data processing but has few impacts on the final results. Thus, we retain the original method in our current paper and add a sentence to express the similar calibration factor derived from the two method (Lines 284-287).

The authors have done a good job of including extra information in the revised supplementary information that increases confidence in their measurements.

The campaign was conducted over 19 days but only very limited time-resolved measurements are presented (one comparison between a clean period and a polluted period, which is very interesting). This seems like a missed opportunity. Were there any interesting variations in Dfm or rho_eff over time? Even if not this would still be interesting to know. Is it possible to add a Figure that shows the time series of Dfm and rho_eff at given mobility diameters? Perhaps also with the PM2.5 time series shown in Fig. S7 to see if there is any correlation. One of the advantages of the tandem DMA-SP2 technique is that it can perform measurements at relatively high time resolution, it seems a shame not to use this advantage.

Response:

Thanks for providing the useful comments. Actually, the initial purpose of this experimental setup was to investigate the mixing state of size-resolved BC-containing particles at a high time resolution, as well as the morphology and effective density of the uncoated BC aggregates. The results presented in this paper is one part of the whole study on the properties of size-resolved BC particles. We had tried to investigate the temporal variations in the morphology and effective density of the uncoated BC aggregates during the data processing. Unfortunately, there have no sufficient data volume to provide the stability and reliable size distribution from which the typical rBC

mass for a given mobility size was derived at a high time resolution, even on the daily scale. As shown in the attached figure, even for the two compared periods presented in our paper (i.e., a clean period and a polluted episode) with relatively adequate time, the data volume is still not sufficient enough as expected to obtain a wonderful number size distribution of the mass-equivalent diameter of the rBC core of *ext*BC. There are obvious fluctuations in the size distribution especially at larger mobility sizes (Fig. R1). Thus, we roughly discussed the differences in the mass-mobility relationship of extBC between the clean period and polluted episode in our current manuscript, although we know the time-resolved mass-mobility relationship should be more interesting. The low data volume is related to the low particle numbers at a certain mobility size and the low $PM_{2.5}$ concentration in this campaign (23 µg m⁻³ on average).



Fig. R1 Number size distributions of the mass-equivalent diameter of the rBC core of *ext*BC normalized by the peak value at five represented mobility diameters (140, 225, 350, 500 and 750 nm) during a clean period (left) and a polluted episode (right). Lognormal fitting is performed for the major peak of each distribution.

The focus of this manuscript is on externally mixed BC particles (the revised manuscript now contains good, solid arguments for how the authors have isolated these particles). I think it would also be interesting to present the campaign-averaged Dfm and rho_eff of internally mixed BC particles (i.e. those displaying delay time greater than 2 microseconds). This will require extra calculations and I know it is outside the main focus of the manuscript. But I think its important to put the extBC results in

context, which is the main focus of the manuscript. For example, is it actually true that the extBC particles display low Dfm because they are uncoated and more aggregatelike? This could be partially answered by checking if the coated particles that were present at the same time displayed higher Dfm.

Response:

The morphology and effective density of internally mixed BC were studied in the previous literature by using a VTDMA-SP2 system (Zhang et al., 2016). In our study, the DMA-SP2 is not likely to be used to study the morphology and effective density of internally mixed BC particles since the mobility diameter of these internally mixed BC particles were selected while only the mass of rBC core were determined by the SP2. Even the mass of the coating material can be estimated using the LEO fitting method and assumed density, larger uncertainties will be induced. Therefore, the DMA-SP2 system provide an advantage to study the mass-mobility relationship of externally mixed BC particles. For internally mixed BC particles, a VTDMA-SP2 system is required. As presented in Zhang et al. (2016), the average effective density of internally mixed BC particles measured at a suburban site nearby Beijing was 1.2 g/cm³, significantly higher than the values of externally mixed BC particles. Correspondingly, the internally mixed BC particles had a lower shape factor than the externally mixed ones. It means the internally mixed BC particles have a more compact structure. Zhang et al. (2016) also showed that the effective density of internally mixed BC particles was increased with the relative coating-thickness (Dp/Dc) perhaps due to the reconstruction of BC aggregates during their aging processing in the atmosphere.

Zhang, Y. X., Zhang, Q., Cheng, Y. F., Su, H., Kecorius, S., Wang, Z. B., Wu, Z. J., Hu, M., Zhu, T., Wiedensohler, A., and He, K. B.: Measuring the morphology and density of internally mixed black carbon with SP2 and VTDMA: new insight into the absorption enhancement of black carbon in the atmosphere, Atmos. Meas. Tech., 9, 1833–1843, doi:10.5194/amt-9-1833-2016, 2016.

Specific comments:

L 46: Higher rho_eff values than what? This is an ambiguous statement. I guess it is meant higher than one might expect based on the trend observed at lower mobility diameters. But this is a poorly defined reference point.

Response:

We have changed the expression to make the meaning there clearer and more readable. (Lines 46–48)

L 104: What is meant by the term 'virtual Df'? Please clarify.

Response:

The 'virtual $D_{\rm f}$ ' means the 'fractal dimension' defined as the scaling exponent between the radius of gyration ($R_{\rm g}$) which is a root mean square radius that quantified the overall size of the aggregate, and the radius of primary spherules (*a*) composing the aggregates, expressed as:

$$\mathbf{N} = k_0 (\frac{R_g}{a})^{D_f},$$

where *N* is the number of primary paritcles and k_0 is the scaling prefactor. We have added similar expression in the revision (Lines 106–108).

L 108: This sentence needs rewriting to clarify that it is the parameter effective density that is difficult to characterize by TEM, not the tandem measurements. (At least I think this is the intended meaning of the sentence.)

Response:

Modified (Lines 113–114).

L 113: Change to 'system detection limit' Response: Corrected (Line 117). L 121: 'Extrapolation' is a better and more specific word to use than 'applicability' here. Response:

Modified (Lines 125–127).

L 421: Are these extBC masses averaged over the full campaign? If so this needs to be explicitly stated.

Response:

Yes, the extBC masses presented here are the campaign average values. We have explicated it in the revision (Line 437).

L 535: Did any of the DMA-(APM, ELPI) studies use a thermodenuder to remove volatile components? This would not necessarily successfully remove all volatile material (as indicated on L 564 for the Rissler study). But it would be good to indicate if one was used or not here (as is done for the Rissler study), to indicate how much volatile material one might expect to be present in these previous measurements. Response:

In the previous literature for the study on the mass-mobility relationship of BC aggregates using the tandem method, the thermodenuder was seldom employed to remove volatile components because previous studies mostly focused on the BC aggregates freshly emitted from the diesel exhaust. The mass fraction of volatile components in these freshly emitted BC aggregates was considered generally low. The thermodenuder might have been used in some other DMA-(APM, ELPI) studies we missed to mention. However, the thermodenuder was not used in any of the studies cited in our current manuscript, except for Risslar et al. (2014).

L 567: My interpretation of this observation (that the discrepancy between DMA-APM and DMA-SP2 measurements is smaller at larger diameters) is that the particles at larger diameters are actually less coated (hence SP2 mass would agree more closely with APM mass). The ~300 nm diameter at which the discrepancy starts to decrease is consistent with the shift from a line with Dfm of 2.51 to one with Dfm of 2.07 shown in Fig. 3

(the kinks in the curves in Figs. 3 and 4 both occur ~300 nm). I'm not sure which interpretation is correct but perhaps this alternative one should at least be mentioned. Response:

We think carefully of this interpretation and found that it is with lower likelihood. Results shown in this study focus on the externally mixed BC particles, i.e., BC particles without coatings or with thin coatings. Thus, the coating effect should be negligible for the SP2 measured extBC particles regardless of the mobility size. We ever suspected that whether the smaller diesel exhaust particles (e.g., <300 nm) measured were more likely to be coated by other components resulting in a higher mass or effective density determined by the DMA-APM system. However, there is no sufficient evidence to support this hypothesis. Moreover, as shown in Fig. S9, the number fraction of ambient extBC in the entire SP2 detectable BC-containing particles showed a minimum at ~300 nm mobility diameter. It means that BC particles in this size range are more likely to be affected and coated by other components than the smaller or larger BC particles. Even the extBC was discussed in this study, the extBC particles in this mobility size range should also be more likely to be affected by other components resulting in a more compact structure than those with smaller or larger mobility sizes. The more compact BC aggregates also resulted in a relatively constant dynamic shape factor in the 200–350 nm mobility diameter range (Fig. 5).

1	A study of the morphology and effective density of externally
2	mixed black carbon aerosols in ambient air using a size-
3	resolved single-particle soot photometer (SP2)
4	Yunfei Wu ^{1*} , Yunjie Xia ^{1, 2} , Rujin Huang ³ , Zhaoze Deng ⁴ , Ping Tian ⁵ , Xiangao Xia ⁴ ,
5	Renjian Zhang ^{1*}
6	1 Key Laboratory of Regional Climate-Environment for Temperate East Asia (RCE-TEA), Institute of
7	Atmospheric Physics, Chinese Academy of Sciences, Beijing 100029, China
8	2 University of Chinese Academy of Sciences, Beijing, 100049, China
9	3 Key Laboratory of Aerosol Chemistry and Physics, State Key Laboratory of Loess and Quaternary Geology,
10	Institute of Earth and Environment, Chinese Academy of Sciences, Xi'an 710061, China
11	4 Key Laboratory of Middle Atmosphere and Global Environment Observation (LAGEO), Institute of
12	Atmospheric Physics, Chinese Academy of Sciences, Beijing 100029, China
13	5 Beijing Weather Modification Office, Beijing 100089, China
14	Correspondence to: Yunfei Wu (wuyf@mail.iap.ac.cn) and Renjian Zhang (zrj@mail.iap.ac.cn)
15	

16 Abstract

17 The morphology and effective density of externally mixed black carbon (extBC) aerosols, important factors affecting the radiative forcing of black carbon, were studied 18 by using a tandem technique coupling a differential mobility analyzer (DMA) with a 19 single-particle soot photometer (SP2). The study extended the mass-mobility 20 relationship to large *ext*BC particles with a mobility diameter (d_{mob}) larger than 350 nm, 21 22 a size range was seldom included in previous tandem measurements of BC aggregates in the atmosphere. The experiment was conducted at an urban site in Beijing during a 23 24 19-day winter period from 23 January to 10 February 2018. Ambient dry particles were selected by the DMA, and the size-resolved extBC particles were distinguished from 25 particles with a thick coating (internally mixed) according to the time delay between 26 the incandescence signal peak and the scattering peak detected by the SP2. The masses 27 of the extBC particles were then quantified. The time differences between the DMA 28

29 size selection and the SP2 measurement were processed previously. The normalized 30 number size distributions were investigated at the prescribed d_{mob} sizes in the range of 140–750 nm to provide the typical mass of *ext*BC at each d_{mob} . On this basis, the mass-31 mobility relationship of the ambient extBC was established, inferring a mass-mobility 32 scaling exponent $(D_{\rm fm})$ (an important quantity for characterizing the morphology of 33 fractal-like BC aggregates) with a value of 2.34±0.03 in the mobility range investigated 34 35 in this study. This value is comparable with those of diesel exhaust particles, implying a predominant contribution of vehicle emissions to the ambient extBC in urban Beijing. 36 Compared to the clean period, a higher $D_{\rm fm}$ value was observed in the polluted episode, 37 indicating a more compact BC aggregate structure than that in the clean period. The 38 effective densities (ρ_{eff}) of the *ext*BC in the same d_{mob} range were also derived, with 39 40 values gradually decreasing from 0.46 g cm⁻³ at 140 nm mobility to 0.14 g cm⁻³ at 750 nm mobility. The ρ_{eff} values were slightly lower than those measured using the DMA-41 aerosol particle mass analyzer (APM) system. The difference in ρ_{eff} values was likely 42 due to the lower BC masses determined by the SP2 compared to those measured by the 43 APM at the same mobility, since the SP2 measured the refractory BC (rBC) mass 44 instead of the total mass of the BC aggregate, which consists of both rBC and a possible 45 46 fraction of nonrefractory components measured by the APM. Higher-The ρ_{eff} values were observed in the 280–350 nm $d_{\rm mob}$ range, and were much closer to the values for 47 soot aggregates reported in the literature. The higher ρ_{eff} values It might be related to the 48 more compact structure of BC aggregates in this range, resulting from the 49 reconstruction effect by volatile and/or semivolatile components in the atmosphere. The 50 51 reconstruction effect might also result in a hiatus in the increased dynamic shape factor 52 in the range of 200-350 nm, which generally increased presented an overall increase 53 from 2.16 to 2.93 in the 140–750 nm d_{mob} range.

54

55 1 Introduction

Black carbon (BC), a byproduct of incomplete combustion, is the main light-absorbing
 component in atmospheric aerosols. BC can lead to positive radiative forcing second in
 magnitude only second to CO₂ and thus warming of the earth's atmosphere (IPCC,

59 2013). However, there remains a large amount of uncertainty regarding the radiative 60 forcing induced by BC due to its complexity and variability in morphology, mixing 61 state and hygroscopicity. Freshly emitted BC particles usually exhibit as-fractal-like aggregates made upcomposed of a number of primary carbon spherules (Park et al., 62 2004; Sorensen, 2011), which are generally hydrophobic. The condensation of organic 63 and/or inorganic components leads to the collapse of the fractal-like aggregates and, in 64 65 turn, a compact structure of BC particles (Slowik et al., 2007; Zhang et al., 2008). Changes in the morphology of BC particles affect their optical properties. Encasement 66 67 by organic and/or inorganic coatings also increases the absorption of BC particles through the lensing effect (Shiraiwa et al., 2010; Peng et al., 2016). In addition, water-68 soluble coatings increase the hydrophilic ability of the BC particles (Zhang et al., 2008; 69 70 McMeeking et al., 2011), indirectly affecting the radiative forcing by affecting the cloud 71 processes.

72 Laboratory studies indicate that freshly emitted BC particles can become thickly coated within a few hours in the atmosphere (Pagels et al., 2009; Peng et al., 2016). Thus, 73 74 many studies have focused on the optical properties and radiative forcing of thickly-75 coated BC particles (Jacobson, 2001; Khalizov et al., 2009; Liu et al., 2017). However, 76 in situ measurements have shown that a great number of uncoated and/or thinly-coated 77 BC particles exist in the ambient atmosphere, with a fraction even higher than that of the-aged BC particles (Schwarz et al., 2008). In general, thickly-_coated BC particles 78 account for <50% of the BC-containing particles in urban areas based on single-particle 79 soot photometer (SP2) measurements (Huang et al., 2012; Wang et al., 2014; Wu et al., 80 81 2016). The existence of a large fraction of uncoated and/or thinly-coated BC particles 82 is likely due to continuous emission from combustion processes such as vehicle exhaust 83 (Wang et al., 2017). Therefore, studies on the radiative forcing of BC particles without thick coatings are also essential, especially in urban areas. First, the morphologies and 84 sizes of these quasi-bare BC particles, which are the essential quantities for calculating 85 the optical properties of BC particles in numerical models, should be investigated 86 (Scarnato et al., 2013; Bi et al., 2013). 87

The morphology of fractal-like BC aggregates is generally characterized by a quantity $\frac{3}{3}$

89 called the $\frac{1}{2}$ fractal dimension² (D_f), which has been well documented in the review 90 literature (Sorensen, 2011). The ideal diffusion-limited cluster aggregation (DLCA), to which the BC aggregates belong, has a $D_{\rm f}$ value of 1.78 \pm 0.1. Recent studies have also 91 reported a similar $D_{\rm f}$ value of ~1.82 for bare-like soot particles by using transmission 92 electron microscopy (TEM) analysis of aerosol samples collected in four different 93 environments (Wang et al., 2017). A significant increase in the $D_{\rm f}$ was observed when 94 95 the soot particles were partly coated or embedded. In the past two decades, the morphologies of the BC aggregates were have also been widely studied by using 96 tandem mobility techniques (Park et al., 2008). Measurements obtained by-using an 97 impactor (e.g., the an electrical low-pressure impactor (ELPI)) or a particle mass 98 analyzer (e.g., the an aerosol particle mass analyzer (APM), the or a centrifugal particle 99 100 mass analyzer (CPMA)) connected in tandem with a differential mobility analyzer 101 (DMA) have revealed the relationship between particle mass and mobility (Park et al., 102 2003; Maricq and Xu, 2004; Olfert et al., 2007; Rissler et al., 2014; Sorensen, 2011; and associated references therein). The derived mass-mobility scaling exponents $(D_{\rm fm})$ 103 104 which were have also been called 'fractal dimensions' in some of these references, 105 varied in-over a wide range of 2.2-2.8 for diesel exhaust particles. These values were 106 inherently higher than the virtual $D_{f_{a}}$ which is defined as the scaling exponents between 107 the radius of gyration of an aggregate and the radius of primary spherules composing the aggregate, due to the improper interpretation of mobility measurements, which 108 wasas demonstrated in detail in Sorensen (2011). The $D_{\rm f}$ of diesel particles obtained 109 using TEM is ~1.75, corresponding to a large $D_{\rm fm}$ value of ~2.35 obtained from based 110 111 on the mass-mobility relationship (Park et al., 2004). The mobility size-dependent 112 effective densities (ρ_{eff}) of BC aggregates were also determined from the DMA-ELPI 113 or DMA-APM (or CPMA) measurements, which were difficult to characterize using 114 the TEM techniques. The previous tandem measurements generally provided the mass-mobility relationship 115 116 of particles with a mobility diameter (d_{mob}) not exceeding 350 nm due to the system

- 117 detection limit (Park et al., 2003; Maricq and Xu, 2004; Olfert et al., 2007; Rissler et
- 118 al., 2014). A condensation particle counter (CPC) is connected next to the DMA-APM

119 system to measure the number concentrations of mobility size-selected particles at 120 various APM voltages. The voltage is proportional to the particle mass, and the voltage 121 resulting in the maximum concentration is in turn considered the typical voltage, in turn, of the mass of particles with a prescribed mobility size. Because the larger particles 122 (e.g., $d_{mob}>350$ nm) are less abundant in the atmosphere than the-smaller particles, 123 124 larger uncertainties exist in the DMA-APM-CPC measurements for these larger 125 particles (Geller et al., 2006). Hence, T the applicability extrapolation of the massmobility relationship established on the basis of tandem measurements of smaller 126 127 mobility diameters (e.g., $d_{mob} < 350$ nm) to larger particles (e.g., $d_{mob} > 350$ nm) is insufficient. 128

The SP2 was developed on the basis of the laser-induced incandescence technique and 129 130 provides-an advantages to in the study of individual BC particle properties, including 131 mass, size and mixing state. The SP2 determines the refractory BC (rBC) mass from 132 particle to particle, thus provideing the masses of BC aggregates throughout a wide size 133 range (70–500 nm mass-equivalent diameter-claimed by according to the manufacturer) 134 with high sensitivity and accuracy (Schwarz et al., 2006). Recently, a tandem system 135 consisting of an SP2 connected next-to a DMA was developed to study the properties 136 of size-resolved BC aerosols in the atmosphere. The mass distributions and mixing 137 states of the size-selected BC were investigated in northern Indiaby using a DMA-SP2 tandem system (Raatikainen et al., 2017). Coupling an SP2 with a volatility tandem 138 139 DMA (VTDMA), the rBC core size distributions of internally mixed BC and those measured by the VTDMA were compared at the prescribed mobility size ranges. 140 141 Subsequently, the morphology and effective density of the internally mixed BC particles were studied (Zhang et al., 2016). The hygroscopic properties of BC particles 142 143 were studied by-using a hygroscopicity tandem DMA (HTDMA)-SP2 coupling system (McMeeking et al., 2011; Liu et al., 2013). Few studies have been performed on the 144 145 morphology and effective density of fractal-like BC aggregates uncoated by that are not coated with other components, especially those in the ambient atmosphere, using the 146 147 DMA-SP2 measurements, especially in the ambient atmosphere.

148 Using the DMA-SP2/CPC system, Gysel et al. (2012) revealed that the SP2 was unable

149 to reliably detect BC particles from a PALAS spark discharge soot generator due to the 150 lower detection limit of the SP2 for loosely packed agglomerates made up of small 151 primary spherules (~5-10 nm in diameter). However, they also claimed that a well-152 aligned SP2 was expected to have a detection efficiency adequate to measure the BC aggregates (e.g., diesel exhaust soot) in the atmosphere because these BC aggregates 153 have larger primary spherules and substantially higher effective densities than the 154 155 agglomerates made up of small primary spherules. Therefore, in this study, a DMA-SP2 tandem system was built to examine the mass-mobility relationship (from which 156 157 the morphology and effective density were further derived) of uncoated BC aggregates, 158 especially in the large particle size range (e.g., $d_{\text{mob}}>350$ nm), which was has seldom been included in previous tandem measurements. Moreover, the uncoated BC 159 160 aggregates were distinguished from the thickly_-coated BC particles by using the SP2, 161 thus allowing the study on of the mass-mobility relationship of ambient BC aggregates in different atmospheric environments. Previous DMA-ELPI or APM tandem 162 measurements were mainly conducted in the laboratory or in the source environments 163 164 (e.g., in the a tunnel) where fresh BC aggregates were predominant.

165 Beijing, the capital of China, has suffered from severe air pollution issues in recent 166 years. Studies have revealed that emissions from coal combustion and/or biomass burning for industry activities and residential heating have played a predominate 167 predominant role in the particulate pollution in Beijing, especially in during the polluted 168 169 episodes (Zhang et al., 2013; Huang et al., 2014; Wu et al., 2017; Ma et al., 2017a, b). 170 Thus, the variation in the mass-mobility relationship of uncoated BC aggregates was 171 also compared for a polluted episode and a clean episode to examine the possible 172 influence of a source change on the morphology of these BC aggregates. In addition, a 173 better mobility size resolution (33 logarithmic size bins from 20 to 750 nm) was set for 174 our DMA-SP2 system than that-was used in previous similar studies, in which only a 175 few mobility diameters in the range of ~150–350 nm were selected (Zhang et al., 2016; 176 Liu et al., 2013; McMeeking et al., 2011). Similar to the study presented in-by 177 Raatikainen et al. (2017), the high size resolution provides an advantage is advantageous 178 for calculating the BC mass and number size distribution in the polluted region in our 179 future studies.

180

208

181 2 Measurements

182 2.1 Experimental setup

183 A tandem system comprisinges a size selection unit and a measurement section was built and deployed in an ambient experiment that was conducted on the roof of a 184 185 building (approximately 8 m above the ground) on the campus of the Institute of Atmospheric Physics, Chinese Academy of Sciences (IAP, CAS) during the winter 186 187 from 23 January to 10 February 2018 (19 days in total). The site is 1L ocated in an urban 188 area of Beijing;, it the site is a few hundred meters away from two main roads and thus may be significantly affected by vehicle emissions. More information about on the 189 190 measurement site is described in previous studies (e.g., Wu et al., 2016, 2017).

191 As shown in Fig. 1, polydisperse aerosols in the sample air were drawn through the size selection unit (a model 3087 neutralizer, a model 3080 classifier and a model 3081 192 DMA, TSI Inc., Shoreview, MN, USA) to generate quasi-monodisperse particles with 193 194 a certain electrical d_{mob} . Before entering the system, the ambient air was dried prior 195 toby passing through a 12-inch-long Nafion dryer (model MD-700-12F-3, Perma Pure 196 LLC, Toms River, NJ, USA). A vacuum pump was used to draw the dry sheath air (e.g., 197 particle-free indoor air) opposite to the flow direction of the sample air to provide thean appropriate vacuum degree required for the dryer. The size-selected particles were 198 delivered to the measurement section for analysis by various methods, including an SP2 199 (Droplet Measurement Technologies, Boulder, CO, USA), a CPC (model 3776, TSI 200 Inc., Shoreview, MN, USA) and two microaethalometers (model AE51, AethLabs, San 201 Francisco, CA, USA). The operational flow rates were set to 0.1, 0.3 and 0.15 LPM 202 203 (STP) for the SP2, CPC and two AE51s, respectively. The sheath flow rate was set to 3 LPM, resulting in a ratio of sheath-to-sample flow rate of 4.3:1 for the DMA. Particles 204 in the range of 15-750 nm in mobility diameter could be selected. The flow rate for 205 each instrument was calibrated using a soap film flowmeter (model Gilian Gilibrator-206 2, Sensidyne, Petersburg, FL, USA) before the experiment to ensure the accuracy of the 207

selected particle sizes and measurements. The scientific purpose of this experimental

setup was to study the mixing states of size-selected BC particles, the mass and number
size distribution of BC, as well as the morphology and effective density of the uncoated
BC aggregates that are discussed in the current study. <u>Since Because</u> only the DMA and
SP2 were involved in the measurements presented in this study, the setting and
operation of the two instruments were described and discussed in detail.

214



215

Fig. 1 Schematic of the experimental setup for size-resolved measurements of blackcarbon.

218

219 2.2 Particle size selection

The DMA was connected to an external computer on which a program was run to 220 221 control the voltage of the DMA, i.e., the particle mobility diameter (d_{mob}) . Thirty-three 222 <u>*d*</u>_{mob}mobility diameters values were set in the program to cyclically control the particles 223 selected by the DMA and gradually increased from 20 nm to 750 nm onat the 224 logarithmic scale. Stepwise size selection was repeated until the operator stopped the 225 program. A short cycle lasting for 18 s with for each of the 33 diameters lasting for 18 s-and a long cycle lasting for 36 s with for each size lasting for 36 s were set to 226 227 alternately operate in this experiment (Fig. S1 in the supplemental file). The purpose of 228 this these settings was to identify the time difference between the size selection and the 229 subsequent measurement, as described in the following sections.-

230

231 **2.3 Black carbon measurement**

232 The individual particulate rBC mass was measured by the SP2 according to the laser-

233 induced incandescence signal when the particle passed through the intense Nd:YAG 234 intracavity continuous laser beam (Schwardz et al., 2006) with a Gaussian distribution. 235 The rBC mass in the SP2 detection range (~0.3-250 fg in this study, dependent on the laser intensity of a specific instrument) is proportional to the peak of the incandescence 236 signal independent of the mixing state of the BC particles. If a BC particle is coated 237 with nonrefractory components, the coating will evaporate before the rBC core 238 239 incandesces, leading to a time lag between the peaks of incandescence and scattering signals that are synchronously detected by the SP2 (Moteki and Kondo, 2007). 240 According to the frequency distribution of the time lag, there was a significant 241 242 distinction between thickly-coated (i.e., internally mixed) BC particles (intBC) and 243 thinly-coated or uncoated (i.e., externally mixed) BC particles (extBC) (Fig. S2) with 244 a minimum frequency at $\sim 2 \mu s$. BC-containing particles with delay times shorter than 245 2 µs were identified as extBC. The delay time threshold might vary slightly from one 246 SP2 to another, ; e.g. for example, Zhang et al. (2016) reported a short time lag of 1.6 µs. However, the delay time threshold should be constant for a given instrument. In 247 248 previous measurements using the same SP2 employed in this study, the critical delay time was maintained at 2 µs regardless of the ambient conditions, such as the pollution 249 250 level (Wu et al., 2016, 2017). A fraction of BC-containing particles with thin or even 251 moderate coatings might also be recognized as extBC using the time-delay approach (Laborde et al., 2012). The effects of these thinly or even moderately coated BC 252 253 particles will beare discussed in Section 3.2 by reducing the delay time threshold from 254 2 µs to 1.2 µs and 0.4 µs, respectively.-255 The scattering signal of a single particle synchronously detected by the SP2 can be used

- to estimate the optical size of the particle. The mixing state of a BC-containing particle can be deduced by comparing the optical size of the particle and the mass-equivalent size of the rBC core. <u>Since-Because</u> the nonrefractory coating of a BC-containing particle is evaporated due to the light absorption and heating of the rBC core when it passes through the laser beam, the scattering cross-section of this particle, which is
- 261 proportional to the scattering intensity at a given incident light intensity, is gradually
- decreased. To estimate the initial optical size of this particle, an approach called

263 leading-edge-only (LEO) fitting was developed (Gao et al., 2007). A small fraction of 264 the measured scattering signal at its in the initial stage before the particle is perturbed 265 by the laser is employed in the LEO fitting to reconstruct the expected scattering distribution of the initial particle. In this method, the location of the leading edge in the 266 beam is also required, which is determined from a two-element avalanche photodiode 267 (APD) signal. Unfortunately, the notch in the two-element APD of our SP2 failed to fix 268 269 in an adequate position (e.g., before the peak location of the laser beam) in this experiment. Thus, the optical size and the consequent coating thickness of the BC-270 271 containing particle cannot be estimated. However, the coating thickness is not a crucial quantity in our current study on the morphology and density of uncoated BC aggregates. 272 The coating thickness can provide a validation of our discrimination of extBC but 273 274 should have little influence on our final analysis and the discussion presented in the 275 following sections.

Before the experiment, the incandescence signal was calibrated using DMA-selected 276 277 monodisperse Aquadag particles. The effective densities of the mobility size-selected 278 Aquadag particles were referred todetermined based on the polynomial equation as a 279 function of the d_{mob} reported in Gysel et al. (2011). The incandescence signal is more 280 sensitive to the Aquadag particles than to the ambient BC particles, i.e., because the 281 Aquadag particle induces a higher incandescence signal peak (by a factor of ~25%) than fullerene soot or an ambient BC particle with the same mass (Laborde et al., 2012). 282 Thus, the peak intensity of the incandescence signal was reduced by a factor of 25% 283 when calculating the calibration coefficient. The calculated calibration factor, 284 285 determined as the slope of the linear regression of rBC masses against the scaled peak 286 heights of SP2's broadband incandescence signal, is consistent well with the factor 287 estimated using a single-point scaling procedure (Baumgardner et al., 2012). The same calibration was performed again after the experiment. The calibration factors varied 288 little (<3%), indicating the stability of the SP2 measurement during the entire 289 experiment (Fig. S3). The uncertainty in the individual rBC mass determination is 290 estimated to be ~10% due to the uncertainties in the rBC mass calibration and the 291 effective density of the calibration material. An additional uncertainty may also arise in 292 10

293 the determination of extBC masses when the time-delay approach is used to distinguish 294 the mixing state of BC particles. The uncertainty will be further discussed in Section 3.2. 295

296

297 **3 Data processing**

3.1 Identifying the time difference between the size selection and the SP2 298 299 measurement

There exists a considerable difference between the time recorded by the size selection 300 program and that recorded by the SP2, due to the time cost of the particles transmitting 301 from the DMA to the SP2, as well as the system clock difference between the computer 302 on which the size selection program runs and that for the SP2 data acquisition. As 303 304 shown in Fig. S1, the SP2 measurement occurs significantly later than the size selection. We have developed two methods to identify the time difference. The first method is 305 306 toinvolves finding the time difference between the local peak in the particle number concentration (including both scattering and incandescence) detected by the SP2 and 307 308 the beginning of the corresponding size selection cycle. During the experiment, stepwise size selection was cyclically performed to produce quasi-monodisperse 309 310 particles with sizes gradually increasing from 20 nm to 750 nm. Thus, at the beginning of each new cycle, the voltage of the DMA should first drop drastically from a high 311 value to a low one to make the particle size decrease from 750 nm to 20 nm. As a result, 312 313 some particles with sizes in the SP2 efficiently detectable range of the SP2 (~100-500 nm) are measured during the descent period, producing a local peak in the number 314 315 concentration. Since Because it takes only a few seconds for the descent, identifying the occurrence time of the local peak position in-based on the SP2 clock and the 316 317 beginning time of the size selection in-based on the external computer clock will 318 provides the time difference of for each cycle.-The other method is to involves checking the consistency of the number and/or mass 319

size distributions between the short-duration cycle and long-duration cycle. Although 320

321 the durations of each size in the short cycle and long cycle are different (18 s vs. 36 s),

322 the time difference between the size selection and the measurement should be uniform

for adjacent short and long cycles. Setting an initial time difference and calculating the mean number and/or mass concentration of each particle size, the number and/or mass size distributions are obtained. Then, the correlation coefficients between the size distributions during short and long cycles are calculated. Changing the time difference gradually, we can obtain a set of correlation coefficients as functions of the time differences. The time difference resulting in the maximum correlation coefficient is considered the difference between the size selection and the measurement.

Since the detection efficiency of the SP2 decreases dramatically in the small particle range (Fig. S4), the size distributions of the SP2-detected particles are inadequate for further calculation of the correlation coefficients. Therefore, the former method was employed in <u>the</u> current study to identify the time difference between the size selection and the SP2 measurement. The latter method will be used to examine the time difference between the size selection and the AE51/CPC measurements in our future study on the number and mass size distributions of BC.

337

338 **3.2** Determination of the typical masses of extBC at prescribed mobility sizes

Particles in a certain size range are selected by the DMA instead of absolutely 339 monodisperse particles at-in a given mobility size due to the effect of the transfer 340 341 function. In addition, larger particles with multiple charges are also selected. The 342 frequency and number size distributions of extBC as a function of the mass-equivalent 343 diameter of rBC (d_{me}) at different mobility sizes are presented in Figs. S5 and S6, respectively. Note that the number size distribution has been normalized by the peak 344 value of the corresponding distribution. Since the frequency and number size 345 distributions of *ext*BC are quite insufficient at small particle sizes ($d_{me} < 70$ nm) due to 346 347 the low detection efficiency of the SP2 (Fig. S4), only the distributions with a $d_{\rm mob}$ 348 larger than 140 nm are presented. In our-the following study, we mainly address the morphology and effective density of extBC in the 140-750 nm d_{mob} range. The 349 350 normalized number size distributions at five representative d_{mob} values (e.g., 140, 351 225, 350, 500, and 750 nm) are also shown in Fig. 2. The eExtBC particles having with 352 a considerable $d_{\rm me}$ range were observed for a certain $d_{\rm mob}$, indicating a wide transfer

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function of the DMA due to the relatively low ratio of sheath-to-sample flow (4.3:1). Multicharged particles also affected the size distribution, especially in the $d_{\rm mob}$ range of 100–400 nm (Ning et al., 2013). As shown in Fig. S6 and Fig. 2, a minor peak is obviously observed at in the right tail of the major peak at each size distribution with afor $d_{\rm mob}$ values of <350 nm.

As mentioned above, a fraction of thinly and/or moderately coated BC particles might 358 359 also be recognized as extBC according to the time delay between the SP2 incandescence and scattering signal peaks. These particles also have impacts on the size distribution 360 361 of extBC at for a given mobility size. A thinly-coated BC particle can be expected to have a larger mass than a bare BC with the same mobility due to the restructuring of 362 the thinly coated BC particle by coating materials. These thinly coated BC particles will 363 364 increase the size distribution at-in_their right tail when mixed with the effect of 365 multicharged particles. It is currently difficult or even impossible to separate the effects of the thinly coated and multicharged particles based on the size distribution of extBC 366 367 at the current stage. To examine the possible effect of these thinly coated particles, we 368 tightened the criterion of the delay time for the discrimination of extBC, gradually 369 decreasing from <2.0 µs to <1.2 µs and <0.4 µs. As shown in Figs. S5 and S6, a decrease 370 in the delay time threshold results in a significant reduction in the data volume used in 371 the analysis but has few effects on the major peak location of the distribution, which is 372 used as the typical $d_{\rm me}$ of *ext*BC at for a given mobility size. The typical $d_{\rm me}$ values, determined as the mode values of the lognormal function that are employed to fit the 373 major peak of the size distribution at a certain mobility size, vary little with the delay 374 375 time thresholds (Table S1). The maximum discrepancy in the $d_{\rm me}$ is <3% throughout the prescribed mobility size range in this study (140-750 nm). The delay time 376 377 threshold-caused change mainly appears inat the right tail of the normalized number size distribution. Reducing the delay time threshold to 0.4 µs results in a significant 378 decrease in the fraction of particles with a large $d_{\rm me}$ compared to the 2.0 µs and 1.2 µs 379 thresholds (Fig. S6). These large particles are likely attributed to thinly and/or even 380 moderately coated BC particles whose structures are relatively more compact than the 381 absolutely bare BC particles. Therefore, we propose that thinly and/or even moderately 382 13

383 coated BC and multicharged particles should both have effects on the size distribution 384 of extBC, mainly atin its right tail, while havingbut have little influence on the typical $d_{\text{me}_{a}}$ that which is considered as the peak d_{me} of the distribution at for a given mobility 385 size. The uncertainty in the typical $d_{\rm me}$ due to the time-delay approach that was utilized 386 to distinguish the *ext*BC is approximately 3% at a given d_{mob} , which is in turn ~10% in 387 of the corresponding mass of extBC. Combining the uncertainty in the rBC mass 388 389 determined by the SP2 (~10%), the total uncertainty in the determined mass of extBC 390 should be ~20% in the studied mobility range of 140-700 nm. For-To achieve anthe adequate data volume used infor the analysis, the results and discussion presented in 391 392 the following sections are based on the database of extBC discriminated as BC-393 containing particles with delay times lower of less than 2.0 µs, unless otherwise 394 specified.

395



396

Fig. 2 Campaign average number size distribution of the mass-equivalent diameter of
the rBC core of *ext*BC normalized by the peak value at five represented representative
mobility diameters (140, 225, 350, 500 and 750 nm) selected by the DMA. Lognormal
fitting is performed for the major peak of each distribution.

401

402 3.3 Theoretical calculation of the morphology and effective density-

403 The structure of *ext*BC, agglomerated by primary spherules with diameters of 20-60 nm

404 (Alexander et al., 2008), can be characterized by its mass-mobility relationship, which

405 is approximately expressed as a power-law relationship between the mass of the 406 agglomerate particle (m) and its mobility diameter (d_{mob}), expressed as

$$407 mtextbf{m} = k \cdot d_{mob}^{D_{fm}} (1)$$

where the prefactor *k* is a constant and $D_{\rm fm}$ is the mass-mobility scaling exponent, which was sometimes erroneously called the ⁴fractal dimension² in previous studies (e.g., Park et al., 2003). This quantity corresponds well to the virtual $D_{\rm f}$ and represents the morphology of the BC aggregates (Sorensen, 2011). The $D_{\rm fm}$ value of a sphere is 3. Thus, the morphology of a particle becomes increasingly closer to that of a sphere as the $D_{\rm fm}$ increases gradually to 3.

The effective density (ρ_{eff}) of the *ext*BC particles is calculated as the ratio of the BC mass (*m*) measured using the SP2 and the BC volume, which is based on the d_{mob} selected by the DMA, expressed as

417
$$\rho_{eff} = \frac{6m}{\pi d_{mob}^3} \tag{2}$$

418 Combining Eqs. 1 and 2, ρ_{eff} can also be expressed as a function of d_{mob} ,

419
$$\rho_{eff} = K \cdot d_{mob}^{D_{fm}-3}$$
 (3)

420 where *K* is a constant, corresponding to the prefactor *k* in the mass-mobility relationship. 421 The dynamic shape factor is also calculated to indicate the morphology of the *ext*BC 422 particles. It is derived from the ratio of the slip-corrected mass-equivalent diameter (d_{me}) 423 and d_{mob} , expressed as

424
$$\chi = \frac{d_{mob} \cdot C_c(d_{me})}{d_{me} \cdot C_c(d_{mod})}$$
(4)

where $d_{\rm me}$ is calculated from the BC mass (*m*) by assuming the BC particle to be a compact sphere with a density of 1.8 g cm⁻³ (Taylor et al., 2015), and $C_{\rm c}$ is the Cunningham slip correction factor parameterized by particle diameter (*d*)

428
$$C_c(d) = 1 + \frac{2\lambda}{a} \left[\alpha + \beta \exp(-\frac{\gamma \cdot d}{2\lambda}) \right]$$
(5)

where λ is the mean free path of the gas molecules, which is set to 65 nm in this study according to Zhang et al. (2016). The values of the three empirical parameters α , β and γ are 1.257, 0.4 and 1.1, respectively (Eq. 9.34 on page 407 in Seinfeld and Pandis, 2006).

434 4 Results and discussion

433

435 **4.1 Mass-mobility relationship of the ambient** *ext*BC

436 A power-law relationship was applied to the d_{mob} -determined extBC mass values, delivering the <u>a campaign average</u> mass-mobility scaling exponent $(D_{\rm fm})$ of the ambient 437 extBC (Fig. 3). In the d_{mob} range of 140–750 nm, the fitted D_{fm} is 2.34, with one standard 438 439 deviation of 0.03. The fitted $D_{\rm fm}$ is close to the lower limit of the $D_{\rm fm}$ values of diesel 440 exhaust particles presented in the literature, indicating the dominant contribution of 441 diesel exhaust to the extBC in our measurement site in urban Beijing. Depending on the fuel type, engine type and load, the $D_{\rm fm}$ of diesel exhaust particles measured by the 442 DMA-APM or DMA-ELPI systems ranged between 2.22 and 2.84 (Olfert et al., 2007; 443 444 Maricq and Xu, 2004; Park et al., 2003 and references therein). The higher D_{fm} values 445 in the literature are likely attributed to the higher fraction of volatile and/or semivolatile 446 components (e.g., sulfate) in the diesel exhaust (Park et al., 2003; Olfert et al., 2007). 447 Accompanied by The presence of these volatile and/or semi-volatile components would 448 result in a more compact structure of the particle, leading to a higher $D_{\rm fm}$ value 449 compared tofor coated particles than for the bare BC aggregate. Since Because the rBC 450 mass instead of the whole particle mass of extBC was measured by the SP2, the a 451 relatively low D_{fm} value was expected and reasonable in this study. In addition, the 452 relatively low D_{fm} value observed in urban Beijing is also likely implies to imply high fuel quality (e.g., low sulfur content) and efficient combustion in vehicle engines, which 453 decrease the organic and/or inorganic fractions in diesel exhaust particles. The $D_{\rm fm}$ 454 value for the ambient soot agglomerates measured with a DMA-APM system near a 455 diesel truck-dominated highway was 2.41 (Geller et al., 2006), slightly higher than the 456 457 value in our study. According to Sorensen (2011), the ideal fractal-like DLCA with a virtual $D_{\rm f}$ of 458 approximately 1.78 should have an expected $D_{\rm fm} \approx 2.2$ in the slip flow regime in which 459 the BC aggregates are generally observed. The slightly larger $D_{\rm fm}$ value of ambient 460

extBC (2.34) in the current study might indicate a more compact structure than the ideal

462 fractal-like DLCA due to the reconstruction effect by other components in the

463 atmosphere. The reconstruction effect appears to be more significant in the smaller 464 particle range than in the larger particle range. The smaller BC particles are more likely to be coated by volatile and/or semivolatile materials, which will be discussed in detail 465 in the next section. We piecewise fitted the mass-mobility relationship using the power 466 law function in the mobility ranges of 140-350 nm and 350-750 nm. A $D_{\rm fm}$ of 467 2.51±0.04 that was obtained in the smaller mobility range (140–350 nm) was obviously 468 469 larger than the fitted value in the whole size range (140-750 nm). In contrast, a much 470 lower $D_{\rm fm}$ with <u>a</u> value of 2.07±0.02 was observed in the larger mobility range (350-750 nm). These results indicate that the ambient extBC particles with larger mobility 471 472 diameters were likely less influenced by the reconstruction effect than those with smaller mobility diameters. 473





Fig. 3 The mass of extBC particles as a function of the mobility diameter in the range 476 477 of 140-750 nm (black circles), fitted by a power-law relationship (red line). The power law functions piecewise fitted in the 140-350 nm mobility range (green line) range and 478 479 in the 350-750 nm mobility range (blue line) are overlaid. The dashed lines represent 480 in-the uncertainties in the determined extBC masses.

481

482 Table 1 The typical mass-equivalent diameters (dme) and corresponding masses of extBC 483 at for different mobility sizes (d_{mob}) selected by the DMA in the whole campaign, in a polluted episode and in a clean period. The effective densities (ρ_{eff}) and dynamic shape 484 17

485	factors (χ) at t	he d_{mob} selected	by the DMA	in throughout the	whole campaign	are also
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486 presented	d.
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$d_{\rm mob}~({\rm nm})$	d _{me} (nm)				mass (fg)			χ
	total	polluted	clean	total	polluted	clean		
140	88.8	87.2	88.5	0.66	0.63	0.65	0.46	2.16
160	97.5	96.9	98.1	0.87	0.86	0.89	0.41	2.27
180	106.2	106.1	107.0	1.13	1.13	1.15	0.37	2.35
200	115.6	116.1	115.5	1.46	1.48	1.45	0.35	2.40
225	127.9	128.6	128.4	1.97	2.01	1.99	0.33	2.41
250	140.5	142.2	141.0	2.62	2.71	2.64	0.32	2.41
280	155.8	158.0	154.4	3.56	3.72	3.47	0.31	2.41
315	172.6	174.8	170.6	4.85	5.04	4.68	0.30	2.40
350	188.2	191.8	185.9	6.28	6.65	6.05	0.28	2.41
400	207.4	213.7	207.4	8.41	9.20	8.41	0.25	2.43
450	226.4	232.3	225.9	10.94	11.81	10.87	0.23	2.50
500	243.8	251.4	242.2	13.65	14.98	13.39	0.21	2.62
560	262.6	271.1	260.1	17.06	18.77	16.58	0.19	2.71
630	283.2	293.5	282.5	21.42	23.83	21.25	0.16	2.81
700	305.1	312.7	305.0	26.76	28.83	26.73	0.15	2.89
750	319.6	328.8	323.5	30.76	33.49	31.92	0.14	2.93

⁴⁸⁷

The variation in the morphology of extBC was further examined by comparing the 488 mass-mobility relationship in a polluted episode and with that in a consecutively 489 490 subsequent clean period. As shown in Fig. S7, a polluted episode rapidly formed at 491 14:00 (local time, if not specified) on 26 January and lasted one and a half days to 0:00 492 on 28 January 2018. The mean PM_{2.5} mass concentration was $72.1\pm23.1 \ \mu g \ m^{-3}$ in 493 this polluted episode, three times the campaign average value ($23.0\pm26.7 \ \mu g \ m^{-3}$). The 494 $D_{\rm fm}$ value is was 2.42±0.09 in the polluted episode, higher than that (2.33±0.06) 495 observed in the subsequent consecutively clean period from 1:00 on 28 January to 18:00 on 31 January 2018, during which the average PM2.5 concentration was merely 8.9±2.7 496 497 μ g m⁻³ (Fig. S8). The higher $D_{\rm fm}$ in the polluted episode is mainly due to the increase 498 in the masses of *ext*BC at large mobility sizes (e.g., $d_{\text{mob}} > 250$). As shown in Table 1, the typical masses of *ext*BC in the 280–700 nm d_{mob} range in the polluted episode are 499 ~7-13% larger than those in the clean period. Although the differences might result 500 from the uncertainty (~20%) in the mass determination of extBC, the commonly larger 501

502 extBC masses (in the 280–700 nm d_{mob} range) to some degree still imply a possibly 503 more compact structure of extBC aggregates in the polluted episode, which might relate 504 to changes in the dominant sources and the ambient environment. Previous studies have revealed that regionally transported pollutants emitted from coal combustion and/or 505 biomass burning played an important or even predominant role in polluted episodes in 506 Beijing (Wu et al., 2017; Ma et al., 2017a). Thus, a considerable fraction of extBC 507 508 aggregates from these sources is likely to coexist with the local vehicle-emitted BC aggregates in the polluted episode, even though the proportion of extBC in the total BC-509 containing particles decreased (Fig. S9). These transported BC aggregates originating 510 from coal combustion and/or biomass burning might have a more compact structure 511 than those from vehicle exhaust due to the differences in the combustion environments 512 513 and efficiencies. In addition, the BC aggregates might also become more compact due 514 to the reconstruction effect by the volatile and/or semivolatile components, which are 515 generally abundant in polluted episodes. Both possible factors are likely to result in the larger $D_{\rm fm}$ values in the polluted episode. 516

517

518 **4.2 Size-resolved effective densities of the ambient** *ext*BC

In contrast to the mass of *ext*BC (*m*), the effective density of the *ext*BC particles (ρ_{eff}) showed a significant decreasing trend as the d_{mob} increased from 140 nm to 750 nm (Fig. 4 and Table 1). The highest ρ_{eff} of 0.46 g cm⁻³ was observed in the 140 nm d_{mob} , likely because the BC aggregates at the smallest size are made up of the fewest primary spherules. When the d_{mob} increased to 750 nm, ρ_{eff} decreased to as low as 0.14 g cm⁻³, approximately one-third of that at 140 nm. The very low ρ_{eff} values agree well with the fractal-like nature of the *ext*BC particles.–



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Fig. 4 The effective density (ρ_{eff}) of the *ext*BC particles as a function of the mobility diameter (d_{mob}) (black circles). The red line represents the power-law fitting of ρ_{eff} against versus d_{mob} . The variations of ρ_{eff} with d_{mob} measured for the soot agglomerates from diesel exhausts (Park et al., 2003) and near-traffic urban environments (Rissler et al., 2014) are also presented as blue triangles and red squares, respectively. The dashed lines represent-in the uncertainties in the determined ρ_{eff} .

The ρ_{eff} values obtained by the DMA-SP2 measurements are close to those of the lower 535 limits of diesel exhaust particles measured by the DMA-APM (or CPMA) or DMA-536 ELPI systems. Park et al. (2003) reported a decrease in the ρ_{eff} of diesel exhaust particles 537 under a moderate (50%) engine load from 0.95 g cm⁻³ to 0.32 g cm⁻³ as the mobility 538 539 diameter increased from 50 nm to 300 nm (Fig. 4). The ρ_{eff} values presented in Park et al. (2003) are approximately 1.25, 1.18 and 1.05 times those in our study at ~150 nm, 540 220 and 300 nm in mobility diameter, respectively. The differences in $\rho_{\rm eff}$ values 541 between our study and the literature are generally within the uncertainty (~20%) in the 542 mass determination of extBC at prescribed mobility sizes. However, the commonly 543 lower $\rho_{\rm eff}$ values are also likely due to the techniques used to determine the mass of BC 544 545 aggregates. Some previous studies on the ρ_{eff} of diesel exhaust particles using the DMA-APM or DMA-ELPI tandem measurements also showed a slightly larger ρ_{eff} throughout 546 the comparable mobility ranges (e.g., ~150-350 nm) than that measured in this study 547 (Maricq and Xu, 2004; Olfert et al., 2007). The masses of the bare BC particles were 548

determined by the laser-induced incandescence technique of the SP2. In a previous 549 550 tandem system, the APM (or CPMA) or ELPI was utilized to determine the typical mass of BC aggregates at a given mobility, and the BC aggregates are likely composed 551 552 of a fraction of volatile and/or semivolatile components in addition to the bare primary 553 particles. These volatile and/or semivolatile components increase the mass of the whole 554 particle, resulting in a larger ρ_{eff} value for at a certain mobility and cause causing a 555 compact structure of the BC aggregate. For example, Olfert et al. (2007) found that the $\rho_{\rm eff}$ of diesel exhaust particles coated with minor sulfate and water contents (~2% of the 556 total particle mass) was ~0.4 g cm⁻³ at 299 nm, only slightly larger than the value of 557 diesel exhaust particles (0.32 g cm⁻³) measured in Park et al. (2003) and that of *ext*BC 558 in the urban atmosphere (0.31 g cm⁻³) in our study at the same mobility size. However, 559 the $\rho_{\rm eff}$ value increased significantly to ~0.71 g cm⁻³ at a relatively high engine load of 560 561 40% due to the high sulfate levels (~30% of the total particle mass) in the diesel exhaust particles (Olfert et al., 2007). 562

563 The ρ_{eff} values of ambient soot aggregates also showed a similar decreasing trend with increasing d_{mob} based on the DMA-APM system (Geller et al., 2006; Rissler et al., 564 2014). Rissler et al. (2014) showed a decrease in the average ρ_{eff} of BC aggregates from 565 566 0.94 g cm⁻³ to 0.31 g cm⁻³ in the near-traffic urban environment as the $d_{\rm mob}$ increased from 50 nm to 350 nm (Fig. 4), similar to that of the freshly emitted diesel exhaust 567 particles presented in Park et al. (2003). However, based on the same method, the ρ_{eff} 568 569 values of the ambient BC aggregates that mostly originated from diesel exhausts (Geller et al., 2006) are substantially different from those presented in Rissler et al. (2014), 570 especially in the large particle size range. The ρ_{eff} at ~350 nm was 0.17 g cm⁻³ in Geller 571 572 et al. (2006), approximately half of the value presented in Rissler et al. (2014). The 573 reason for the discrepancy might be related to the large measurement uncertainties of 574 the DMA-APM system for large particles, e.g., with <u>dmob</u> sizes greater than 300 nm-in 575 the d_{mob} , since these large particles are less abundant in the ambient atmosphere (Geller et al., 2006). Compared to the results presented in Rissler et al. (2014), the ρ_{eff} values 576 577 of ambient extBC aggregates in our study are slightly lower, e.g., by ~17%, ~18% and 578 ~6% forin d_{mob} values of 150 nm, 250 nm and 350 nm-d_{mob}, respectively. The relatively

579 higher ρ_{eff} values are also likely attributed to the effects of volatile and/or semivolatile 580 components in the soot aggregates. Rissler et al. (2014) found that the residual mass fraction of volatile and/or semivolatile materials in the soot aggregates was ~10%, even 581 when the sample air was heated to 300 °C before entering the system for measurement. 582 583 It is interesting to note that the ρ_{eff} values appear to be closer to the values presented in the literature using the DMA-APM measurements in the 280–350 nm $d_{\rm mob}$ range (Fig. 584 585 4). As shown in Fig. 3, larger typical masses of extBC in this range are also observed beyond the logarithmic scaled linear curve that is fitted to the mass-mobility 586 587 relationship. The relatively larger masses and ρ_{eff} values might imply a more compact 588 structure of extBC aggregates in this range, which is likely resulted results from the reconstruction effect by the ambient volatile and/or semivolatile components. As shown 589 590 in Fig. S9, the size-resolved number fractions of extBC exhibit a minimum in the 280-591 350 nm d_{mob} range, regardless of whether they are associated with the polluted episode or the clean period. This finding indicates that particles in this mobility range are more 592 593 likely to be thickly coated by other components compared than are to those particles in 594 the smaller or larger mobility ranges. Zhang et al. (2016) also observed an increased coating thickness of the BC-containing particles in the mobility range of 200-350 nm 595 596 (Table 1 in the literature) using the VTDMA-SP2 measurement at a suburban site ~70 597 km away from our observation site, although the variation in the coating thickness in the larger mobility range was not investigated. It should be noted Notably, that the 598 number fraction of extBC at each mobility size presented in Fig. S9 is roughly 599 calculated as the ratio of the extBC number concentration to the sum of extBC and 600 601 intBC, in which the multiply charged effects were not corrected. Although the extBC 602 particles without coatings and/or with thin coatings are the focus of the current study, 603 the higher fraction of thickly coated BC particles in the 280–350 nm d_{mob} range implies a higher possibility that these extBC particles in the same range were affected by 604 volatile and/or semivolatile materials in the atmosphere, in turn resulting in a more 605 compact structure of these BC aggregates. Further detailed studies of the size 606 distribution of BC (including extBC, intBC and both) and non-BC particles based on 607 the combined measurements of SP2 and CPC are needed in our further work to reveal 608 22

609 the potential mechanism for this phenomenon.

Although the ρ_{eff} of *ext*BC at small sizes (d_{mob} <140 nm) cannot be determined due to the lower limit of the DMA-SP2 system, we extended the ρ_{eff} of *ext*BC to a large size range ($350 < d_{mob} < 750$ nm), which was barely investigated in previous studies using tandem measurements. A continuous decrease in ρ_{eff} with increasing d_{mob} was observed even in the large size range between 350 nm and 750 nm (Fig. 4). It is reasonable to infer that the structure of the *ext*BC particles becomes looser when the fractal-like aggregates built up by the primary spherules increase.

617

618 **4.3 Dynamic shape factors of the ambient** *ext*BC

Due to their fractal-like structures, the extBC particles generally have large dynamic 619 620 shape factors (χ) with values in the range of 2.16 to 2.93 (Table 1), much larger than those of *int*BC with an average value of ~1.2 (Zhang et al., 2016). The χ value declined 621 exponentially as a function of coating thickness of BC-containing particles (Zhang et 622 al., 2016). In contrast to the decrease in ρ_{eff} , the χ values of *ext*BC generally increase as 623 the-d_{mob} increases from 140 nm to 750 nm (Fig. 5). The extBC particles 750 nm in 624 625 mobility diameter have a mean χ value of 2.93, approximately 1.36 times that for 140 626 nm $d_{\rm mob}$ particles (Table 1). The larger particles have looser structures, resulting in higher χ values. However, the χ values appear to vary slightly in a narrow range between 627 628 2.40-to and 2.41 in the size range of 200 nm to 350 nm (Fig. 5). The hiatus in the gradual 629 increase in χ should-is also <u>likely</u> related to the more compact structure of *ext*BC particles in the 280-350 nm mobility range, which has been discussed in detail in the 630 631 previous sections .-.



633

Fig. 5 The dynamic shape factor of the *ext*BC particles as a function of the mobilitydiameter in the range of 140–750 nm.

636

637 5 Conclusions

638 The DMA-SP2 system was established to study the morphology and effective density 639 of the ambient extBC particles, especially in the larger mobility size range, i.e., $350 < d_{mob} < 750$ nm, which was seldom investigated in previous tandem measurements. 640 Quasi-monodisperse particles in the d_{mob} range of 20–750 nm were stepwise selected 641 using the DMA and then delivered to the SP2 for rBC mass measurement and mixing 642 state discrimination. The time difference between the size selection and the SP2 643 measurement was previously processed using the local peak approach. The normalized 644 645 number size distribution of extBC, distinguished as having a delay time between the incandescence signal peak and the scattering peak detected by the SP2 of less than 2 µs, 646 as a function of d_{me} was investigated at each prescribed mobility size in the range of 647 140-750 nm. The size distributions at smaller mobility sizes were not presented due to 648 649 the lower limit of the rBC mass determined by using the SP2. The peak $d_{\rm me}$, calculated 650 as the mode value of a lognormal function fitted to the major peak of the size 651 distribution, was determined as the typical d_{me} value at each mobility size. Consequently, 652 the typical mass of extBC at each mobility size was identified. Reducing the time delay 653 time-threshold employed to discriminate the extBC had few effects on the determined masses of extBC, implying the reliability of our study for the extBC particles. The 654

655 uncertainty in the determined extBC masses was ~20%, combining based on a 656 combination of the uncertainty in the SP2-measured rBC mass with that and the 657 uncertainty related to the time- delay approach. On this basis, the mass-mobility relationship of ambient extBC in urban Beijing was investigated. The campaign average 658 $D_{\rm fm}$ value was estimated to be 2.34±0.03 by fitting the derived *ext*BC masses as a 659 660 power-law function of $d_{\rm mob}$ in the range of 140–750 nm, close to the lower-limit $D_{\rm fm}$ 661 value of diesel exhaust particles. A relatively larger $D_{\rm f}$ value was observed in the polluted episode than in the clean period (2.42±0.09 vs. 2.33±0.06), implying a more 662 compact structure of BC aggregates in the polluted episode. 663

A decrease in $\rho_{\rm eff}$ with increasing $d_{\rm mob}$ was observed, with the $\rho_{\rm eff}$ value decreasing from 664 0.46 g cm⁻³ at <u>a d_{mob} value of</u> 140 nm mobility to 0.14 g cm⁻³ at 700 nm. The ρ_{eff} values 665 666 derived using the DMA-SP2 measurement were slightly lower than those based on the 667 DMA-APM measurement. This difference was most likely due to the bias in the extBC mass determination using the SP2 and APM techniques. The pure rBC mass determined 668 using the SP2 in this study was generally lower than the total mass of the BC aggregate, 669 which comprises both rBC and a possible fraction of nonrefractory components. The 670 $\rho_{\rm eff}$ values in the 280–350 nm mobility range appeared to be much closer to the values 671 672 for soot aggregates reported in the literature by using the DMA-APM tandem measurement. This finding might be related to the more compact structure of BC 673 aggregates in this range, which was likely influenced by the reconstruction effect of 674 volatile and/or semivolatile components in the atmosphere. The reconstruction effect 675 might also result in a hiatus in the gradually increased χ value in the range of 200–350 676 nm. Large χ values generally increased from 2.16 to 2.93 with increasing d_{mob} , further 677 implying the high fractal structure of extBC particles. 678

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