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7 **New and improved infrared absorption cross sections for**
8 **trichlorofluoromethane (CFC-11)**
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11 by

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39 **Abstract**

40 Trichlorofluoromethane (CFC-11), a widely used refrigerant throughout much of the
41 twentieth century and a very potent (stratospheric) ozone depleting substance (ODS), is now
42 banned under the Montreal Protocol. With a long atmospheric lifetime, it will only slowly
43 degrade in the atmosphere, so monitoring its vertical concentration profile using infrared-
44 sounding instruments, thereby validating stratospheric loss rates in atmospheric models, is of
45 great importance; this in turn requires high quality laboratory spectroscopic data.

46 This work describes new high-resolution infrared absorption cross sections of
47 trichlorofluoromethane / dry synthetic air over the spectral range 710 – 1290 cm^{-1} ,
48 determined from spectra recorded using a high-resolution Fourier transform spectrometer
49 (Bruker IFS 125HR) and a 26-cm-pathlength cell. Spectra were recorded at resolutions
50 between 0.01 and 0.03 cm^{-1} (calculated as $0.9/\text{MOPD}$; MOPD = maximum optical path
51 difference) over a range of temperatures and pressures (7.5 – 760 Torr and 192 – 293 K)
52 appropriate for atmospheric conditions. This new cross-section dataset improves upon the
53 one currently available in the HITRAN and GEISA databases through an extension to the
54 range of pressures and temperatures, better signal-to-noise and wavenumber calibrations, the
55 lack of channel fringing, the better consistency in integrated band intensities, and additionally
56 the coverage of the weak combination band $\nu_2 + \nu_5$.

57

58 **1. Introduction**

59 Chlorofluorocarbons (CFCs) were first developed in the 1930s as safe, reliable, and
60 non-toxic refrigerants for domestic use. Trichlorofluoromethane, known as CFC-11 or
61 Freon-11, and dichlorodifluoromethane, known as CFC-12 or Freon-12, were the two most
62 widely used CFCs in applications ranging from refrigerators and air conditioners to
63 propellants in spray cans and blowing agents in foam production.

64 Ultimately, however, CFCs proved too good to be true. The explosion in their use
65 led to a steady increase in their atmospheric abundances. While inert in the troposphere, it
66 was this stability which enabled them to reach the stratosphere where dissociation by
67 ultraviolet radiation released chlorine atoms, which catalyse the destruction of stratospheric
68 ozone (Solomon, 1999). The realisation of this impending environmental disaster prompted
69 international action and in 1987 the Montreal Protocol was ratified; this led to the phasing out
70 of the worldwide production and use of CFCs. CFCs are still released into the atmosphere
71 from “banks”, such as old refrigerators, however these are not regulated by the Protocol
72 (Harris et al., 2014). Banks are the major source of emissions for many ODSs, including
73 CFC-11 which has a long atmospheric lifetime of 52 years (Harris et al., 2014).

74 At present, CFC-11 is the second most abundant CFC in the atmosphere and
75 contributes the second-highest amount of chlorine to the stratosphere, behind CFC-12. In
76 addition to its role in stratospheric ozone destruction – it has the highest ozone depletion
77 potential (1.0) (Harris et al., 2014) of all the CFCs – CFC-11 is a particularly strong
78 greenhouse gas – it has a 100-yr global warming potential of 5160 (Harris et al., 2014).

79 As a key species in stratospheric ozone destruction, CFC-11 atmospheric
80 concentrations are monitored in situ at the surface, e.g. the annual global mean mole fraction
81 of CFC-11 measured by the AGAGE (Advanced Global Atmospheric Gases Experiment)
82 network in 2012 was 235.5 ppt (Carpenter et al., 2014). However, in order to measure
83 concentrations in the stratosphere where ozone destruction occurs, remote-sensing techniques
84 are required. Table 1 contains a listing of limb sounders capable of measuring CFC-11, as
85 described in the literature.

86 The infrared (IR) spectra for large molecules like trichlorofluoromethane are highly
87 complex, consisting of very many closely spaced spectroscopic lines, making the task of
88 generating line parameters from measurements an almost impossible one. For the purposes
89 of atmospheric remote sensing, it is possible to use absorption cross sections in forward
90 models instead of line parameters, however this requires laboratory measurements of air-
91 broadened spectra over a range of temperatures and pressures. The accuracy of retrievals of

92 CFC-11 abundances for the limb sounders in Table 1 is very much dependent on the quality
93 of the underlying spectroscopy; ideally absorption cross-section datasets should cover a range
94 of atmospherically relevant pressure-temperature (PT) combinations, with accurate
95 wavenumber scales and band intensities, and properly resolved spectral features. This work
96 presents new spectroscopic data, optimised for limb sounding instruments, which improve
97 upon those currently available in the HITRAN (Gordon et al., 2017) and GEISA (Jacquinet-
98 Husson et al., 2016) databases.

99

100 **2. Infrared spectroscopy of trichlorofluoromethane**

101 **2.1. Spectroscopic background**

102 There are two stable isotopes of carbon and chlorine, and one of fluorine, resulting
103 in eight stable isotopologues of trichlorofluoromethane, namely $^{12/13}\text{C}^{35}\text{Cl}_3\text{F}$,
104 $^{12/13}\text{C}^{35}\text{Cl}_2^{37}\text{ClF}$, $^{12/13}\text{C}^{35}\text{Cl}^{37}\text{Cl}_2\text{F}$, and $^{12/13}\text{C}^{37}\text{Cl}_3\text{F}$; these belong to the point groups C_{3v} , C_s ,
105 C_s and C_{3v} , respectively. Taking into account the natural abundances of $^{12}\text{C} / ^{13}\text{C}$ (~ 99%
106 and ~1%), and $^{35}\text{Cl} / ^{37}\text{Cl}$ (~ 76% and ~24%), the most abundant isotopologues are therefore
107 $^{12}\text{C}^{35}\text{Cl}_3\text{F}$, $^{12}\text{C}^{35}\text{Cl}_2^{37}\text{ClF}$, and $^{12}\text{C}^{35}\text{Cl}^{37}\text{Cl}_2\text{F}$, with abundances of 43%, 41%, and 13%,
108 respectively.

109 As a non-linear molecule with five atoms, trichlorofluoromethane possesses nine
110 normal vibrational modes; in the C_{3v} point group there are three non-degenerate
111 fundamentals of A_1 symmetry (ν_1 , ν_2 , and ν_3), and three doubly-degenerate fundamentals of
112 E symmetry (ν_4 , ν_5 , and ν_6). For the C_s point group, the ν_1 , ν_2 , and ν_3 modes possess A'
113 symmetry, with the doubly-degenerate ν_4 , ν_5 , and ν_6 modes each splitting into one A' and
114 one A'' mode (Snels et al., 2001). Since the splittings in the ν_4 , ν_5 , and ν_6 levels are small, it
115 is normal to label these bands assuming C_{3v} symmetry. The 710 – 1290 cm^{-1} spectral range
116 covered in the present work contains two strong fundamental bands, $\nu_1 \sim 1081.28 \text{ cm}^{-1}$ and
117 $\nu_4 \sim 849.5 \text{ cm}^{-1}$, and a weaker combination band, $\nu_2 + \nu_5 \sim 936.5 \text{ cm}^{-1}$; reported frequencies
118 are those for the most abundant isotopologue, $^{12}\text{C}^{35}\text{Cl}_3\text{F}$ (von Lilienfeld et al., 2007 ; Snels et
119 al., 2001). Isotopologues complicate the already dense CCl_3F rotation-vibration spectrum;
120 each has slightly different molecular parameters, with bands shifted by small amounts
121 relative to each other. These main band systems are shown in Figure 1 in the plot of the new
122 absorption cross section at 191.7 K and 7.535 Torr. Details on the measurement conditions
123 and derivation of this cross section are given in Section 3.

124

125 **2.2. A brief history of trichlorofluoromethane absorption cross sections**

126 High resolution (0.03 cm^{-1}) absorption cross sections of pure trichlorofluoromethane
127 at 296 K were first included in HITRAN as part of the 1986 compilation (Massie et al., 1985;
128 Rothman et al., 1987). The HITRAN 1991/1992 compilation saw the first introduction of
129 temperature-dependent cross sections (203 – 293 K) for CFC-11 (McDaniel et al., 1991;
130 Rothman et al., 1992; Massie and Goldman, 1992); as before these were derived from
131 measurements of pure CCl_3F at 0.03 cm^{-1} resolution.

132 While the two previous HITRAN editions (1986 and 1991/1992) neglected pressure-
133 broadening effects on the CCl_3F spectra, cross sections for 33 distinct PT combinations
134 (201–296 K and 40–760 Torr N_2 -broadened) over two wavenumber ranges, $810\text{--}880\text{ cm}^{-1}$
135 and $1050\text{--}1120\text{ cm}^{-1}$, were introduced into HITRAN 1996 (Li and Varanasi, 1994; Rothman
136 et al., 1998). Another 22 PT combinations covering lower pressures and temperatures over
137 the same wavenumber ranges were added to HITRAN 2000 (provided by Varanasi, cited
138 within Rothman et al., 2003), bringing the overall PT coverage to 190–296 K and 8–760
139 Torr. Out of these 55 PT combinations, four pairs possess both temperature and pressure
140 within 1 K and 5 Torr, respectively. This dataset, henceforth referred to as the Varanasi
141 dataset, has been used widely for remote-sensing applications since it was first introduced; it
142 is still the dataset included in the most recent GEISA 2015 (Jacquinet-Husson et al., 2016)
143 and HITRAN 2016 (Gordon et al., 2017) spectroscopic databases. Despite its widespread
144 use, the Varanasi dataset has some deficiencies which will be discussed in Section 4,
145 alongside a comparison with the new spectroscopic data taken as part of the present work.

146

147 **3. New absorption cross sections of air-broadened trichlorofluoromethane**

148 **3.1. Experimental**

149 The experimental setup at the Molecular Spectroscopy Facility (MSF), Rutherford
150 Appleton Laboratory (RAL) and the experimental procedures have been described previously
151 for related measurements (e.g. Harrison et al., 2010; Harrison, 2015b; Harrison, 2016); the
152 reader is referred to one of these previous studies for more information. Instrumental
153 parameters associated with the Fourier Transform Spectrometer (FTS) used for the
154 measurements, sample details, and the cell configuration are summarised in Table 2. The
155 sample pressures and temperatures for each air-broadened spectrum, along with their
156 experimental uncertainties and associated spectral resolutions, are listed in Table 3.

157

158 **3.2. Generation of absorption cross sections**

159 The procedure used to generate absorption cross sections from measured spectra has
 160 been reported previously (e.g. Harrison et al., 2010; Harrison, 2015b; Harrison, 2016), so the
 161 full details are not provided here. The wavenumber scale of the cross sections is calibrated
 162 against the positions of isolated N₂O absorption lines taken from the HITRAN 2012 database
 163 (Rothman et al., 2013). The absorption cross sections, $\sigma(\nu, P_{air}, T)$ in units of cm² molecule⁻¹,
 164 at wavenumber ν (cm⁻¹), temperature T (K) and synthetic air pressure P_{air} , are normalised
 165 according to

$$167 \int_{710 \text{ cm}^{-1}}^{1290 \text{ cm}^{-1}} \sigma(\nu, P_{air}, T) d\nu = 9.9515 \times 10^{-17} \text{ cm molecule}^{-1}, \quad (1)$$

168
 169 where the value on the right hand side is the average integrated band intensity over
 170 the spectral range 710 – 1290 cm⁻¹ for three 760-Torr-N₂-broadened trichlorofluoromethane
 171 spectra (at 278, 298, and 323 K) from the Pacific Northwest National Laboratory (PNNL) IR
 172 database (Sharpe et al., 2004). This intensity calibration procedure counters problems with
 173 trichlorofluoromethane adsorption in the vacuum line and on the cell walls, and furthermore
 174 assumes that the integrated intensity over each band system is independent of temperature.
 175 The reader is referred to Harrison et al. (2010) for a more complete explanation of the
 176 underlying assumption, and references cited within Harrison (2015a, 2015b, and 2016) for
 177 details on previous successful uses of this approach.

178 A selection of the derived absorption cross sections is presented in Figure 2,
 179 showing the expected behaviour with temperature at a total pressure of ~ 200 Torr; the
 180 wavenumber range covers the microwindow for the ACE-FTS v3.6 retrieval scheme.

181

182 **3.3. Absorption cross section uncertainties**

183 The accuracy of the wavenumber scale for the new absorption cross sections is
 184 comparable to the accuracy of the N₂O lines used in the calibration; according to the
 185 HITRAN error codes, this is between 0.001 and 0.0001 cm⁻¹. The uncertainty in the intensity
 186 is dominated by systematic errors. A true measure of the random errors as a function of
 187 wavenumber would ideally require multiple concentration-pathlength burdens at each PT
 188 combination, however only one is available for each; however, as indicated in Section 4.4,
 189 these are small and make minimal contribution to the overall error budget. The maximum
 190 systematic uncertainties in the sample temperatures (μ_T) and total pressures (μ_P) are 0.4 %

191 and 0.7 %, respectively (see Table 3). The photometric uncertainty (μ_{phot}), associated with
192 the detection of radiation by the MCT detector and systematic error arising from the use of
193 Bruker's non-linearity correction for MCT detectors, is estimated to be ~2 %. The
194 pathlength error (μ_{path}) is estimated to be negligibly small, lower than 0.1 %. According to
195 the PNNL metadata, the systematic error in the PNNL CCl₃F spectra used for the intensity
196 calibration is estimated to be less than 3 % (2σ). Equating the error, μ_{PNNL} , with the 1σ
197 value, i.e. 1.5 %, and assuming that the systematic errors for all the quantities are
198 uncorrelated, the overall systematic error in the dataset can be given by:

199

$$200 \quad \mu_{\text{systematic}}^2 = \mu_{\text{PNNL}}^2 + \mu_{\text{T}}^2 + \mu_{\text{P}}^2 + \mu_{\text{phot}}^2. \quad (2)$$

201

202 Note that using PNNL spectra for intensity calibration effectively nullifies the errors in the
203 trichlorofluoromethane partial pressures and cell pathlength, so these do not have to be
204 included in Eq. 2. According to Eq. 2, the systematic error contribution, $\mu_{\text{systematic}}$, to the new
205 cross sections is ~3% (1σ).

206

207

208 **4. Comparison between absorption cross-section datasets**

209 In this section the new dataset presented in this work is compared with the older
210 Varanasi dataset, which has a stated uncertainty of 2 % (Li and Varanasi, 1994). The
211 comparison focuses on their wavenumber scales, integrated band strengths, artefacts such as
212 channel fringing, signal-to-noise ratios, spectral resolution, and PT coverage. Given the
213 various problems identified in sections below, the 2 % uncertainty is a significant
214 underestimate. In addition, the new dataset includes the weak combination band, $\nu_2 + \nu_5$, not
215 present in the Varanasi measurements. These new data will provide a more accurate basis for
216 retrieving CFC-11 from atmospheric spectra recorded in the limb.

217

218

219 **4.1. Wavenumber scale**

220 It is likely that the wavenumber scale for the Varanasi dataset was never calibrated;
221 this has been observed in a number of recent studies for other halogenated species in which
222 new datasets have been compared with older Varanasi datasets, e.g. HFC-134a (Harrison,
223 2015a), CFC-12 (Harrison, 2015b), and HCFC-22 (Harrison, 2016). As explained earlier, the

224 absolute accuracy of the wavenumber scale for the new dataset lies between 0.001 and
225 0.0001 cm^{-1} . In comparison, the ν_4 band in the Varanasi cross sections is shifted too low in
226 wavenumber; this shift varies between cross sections, e.g. by $\sim 0.002 \text{ cm}^{-1}$ (a correction
227 factor of ~ 1.000002) for the 190 K / 7.5 Torr ν_1 Varanasi measurement and by $\sim 0.007 \text{ cm}^{-1}$
228 (a correction factor of ~ 1.000007) for 216.1 K / 100.0 Torr ν_1 .

229

230 **4.2. Integrated band strengths**

231 Integrated band strengths for the Varanasi cross sections have been calculated over
232 the spectral ranges of the cross-section files, $810 - 880$ and $1050 - 1120 \text{ cm}^{-1}$, covering the ν_4
233 and ν_1 bands respectively, and compared with those for the new absorption cross sections
234 calculated over the same ranges; plots of integrated band strength against temperature for
235 each dataset, including the PNNL spectra, and wavenumber range can be found in Figure 3.
236 At each temperature the Varanasi integrated band strengths display a small spread in values,
237 most notably for the ν_4 band, however there is no evidence for any temperature dependence,
238 backing up the assumption in Section 3.2 that the integrated intensity over each band system
239 is independent of temperature. The small spread in values is likely due to inconsistencies in
240 the baselines for the Varanasi cross sections, which are larger for the ν_4 band. Additionally,
241 according to the PNNL spectra and the new measurements, the ν_4 cross section at 810 cm^{-1} is
242 non-zero due to the presence of a weak hot band. Therefore, calculating integrated band
243 strengths for the new dataset over the $810 - 880 \text{ cm}^{-1}$ range creates a very small temperature
244 dependence in the ν_4 integrated band strengths. Unfortunately, the wavenumber ranges do
245 not extend far enough to obtain an unambiguous measure of the baseline position for the
246 Varanasi data, and the cross sections in the HITRAN and GEISA databases have had all
247 negative cross section values set to zero, which has the effect of adjusting the baseline
248 positions by a small amount near the band wings.

249

250 **4.3. Channel fringes**

251 Most of the absorption cross sections in the Varanasi CFC-11 dataset contain
252 noticeable channel fringes above the noise level (refer to Figure 4 for an example of this); in
253 transmittance these would equate to peak-to-peak amplitudes as high as $\sim 2-3 \%$. For the
254 measurements described in the present work, wedged cell windows were used to avoid
255 channel fringes by preventing reflections from components in the optical path of the
256 spectrometer.

257

258 **4.4. Signal-to-noise ratios (SNRs)**

259 The SNRs of the transmittance spectra measured in the present work have been
260 calculated using Bruker's OPUS software at $\sim 990 \text{ cm}^{-1}$ where the transmittance is close to 1;
261 the values range from 2600 to 4700 (rms), equivalent to percentage transmittances between
262 0.04 and 0.02 %. A direct comparison with the Varanasi dataset, however, is not possible
263 without the original transmittance spectra or, at the very least, information on the
264 experimental mixing ratios. Further complicating issues, the Varanasi cross sections are
265 missing negative values near the baselines (refer to Section 4.2) and many have channel
266 fringes superimposed. However, it is apparent from a direct inspection that the new cross
267 sections have improved SNR, in some cases substantially so, such as shown in Figure 5.

269 **4.5. Spectral resolution**

270 All spectra used to create the Varanasi cross-section dataset were either recorded at
271 0.01 (for sample mixtures of 75 Torr and below) or 0.03 cm^{-1} spectral resolution (defined as
272 0.9/MOPD). In the present work 0.01 cm^{-1} resolution was used for mixtures below 10 Torr,
273 0.03 cm^{-1} for 300 Torr and above, and 0.015 and 0.0225 cm^{-1} for intermediate pressures. The
274 spectra recorded at 191.6 K and 98.14 / 200.0 Torr were mistakenly recorded at spectral
275 resolutions of 0.0225 / 0.0300 cm^{-1} instead of the planned 0.015 / 0.0225 cm^{-1} . However,
276 careful inspection indicated that there was no under-resolving of spectral features for these
277 two measurements. Overall, the dataset comparison indicates that the spectral resolutions
278 chosen for the Varanasi measurements were suitable.

280 **4.6. Pressure-temperature coverage**

281 An absorption cross-section dataset used in remote sensing should cover all possible
282 combinations of pressure and temperature appropriate for the region of the atmosphere being
283 observed; in this case the focus is on the mid-troposphere ($\sim 5 \text{ km}$) up to the stratosphere.
284 Extrapolating beyond the temperatures and pressures represented within the dataset is
285 generally unreliable, so forward model calculations should at the very least use a four-point
286 interpolation scheme. With this in mind, the P and T of the laboratory measurements were
287 chosen to cover the range of P and T from ACE-FTS v3.0 data. The additional cross sections
288 outside the range of the Varanasi P and T will ensure a better coverage for analysing
289 atmospheric limb spectra. Figure 6 provides a graphical representation of the PT
290 combinations for both datasets, illustrating the improved PT coverage (30 PT combinations
291 in total) relative to the Varanasi dataset. The sampling density in PT space is lower than for

292 the Varanasi dataset; However, due to the congestion and lack of any strong sharp features in
293 the spectra, it is not anticipated that this will have any noticeable effect for remote-sensing
294 applications.

295
296

297 **5. Conclusions**

298 New high-resolution IR absorption cross sections for air-broadened
299 trichlorofluoromethane (CFC-11) have been determined over the spectral range 710 – 1290
300 cm^{-1} , with an estimated systematic uncertainty of $\sim 3\%$. Spectra were recorded at
301 resolutions between 0.01 and 0.03 cm^{-1} (calculated as 0.9/MOPD) over a range of
302 atmospherically relevant temperatures and pressures (7.5 – 760 Torr and 192 – 293 K).
303 These new absorption cross sections improve upon those currently available in the HITRAN
304 and GEISA databases. In particular, they cover a wider range of pressures and temperatures,
305 they have a more accurately calibrated wavenumber scale, they have more consistent
306 integrated band intensities, they do not display any channel fringing, they have improved
307 SNR, and additionally they cover the weak combination band, $\nu_2 + \nu_5$.

308

309 **Data availability**

310 The new CFC-11 absorption cross section dataset will be made available to the
311 community via the HITRAN and GEISA databases, but in the meantime is available
312 electronically from the author.

313

314

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319 measurements.

320

321 **Figure Captions**

322 Figure 1. The absorption cross section of trichlorofluoromethane / dry synthetic air at 191.7
323 K and 7.535 Torr (this work), with vibrational band assignments for the main band systems
324 in the 710 – 1290 cm^{-1} spectral region.

325

326 Figure 2. The new absorption cross sections of trichlorofluoromethane / dry synthetic air at a
327 total pressure of ~ 200.0 Torr over a range of temperatures (191.6, 202.4, 216.6, 232.6,
328 252.5, and 273.8 K). The observed narrowing of the ν_4 band as the temperature decreases is
329 due to the decline in Boltzmann populations of the upper rovibrational levels of the ground
330 state.

331

332 Figure 3. Integrated band strength as a function of temperature for the new, Varanasi, and
333 PNNL cross-section datasets over the wavenumber ranges 810 – 880 and 1050 – 1120 cm^{-1} .

334

335 Figure 4. The Varanasi absorption cross section of trichlorofluoromethane / dry synthetic air
336 at 232.7 K and 250.0 Torr (black), with the new cross section at 232.6 K and 201.0 Torr
337 overlaid (red; this work). Channel fringes in the Varanasi cross section are clearly visible.

338

339 Figure 5. The Varanasi absorption cross section of trichlorofluoromethane / dry synthetic air
340 at 216.5 K and 7.50 Torr (black), with the new cross section at 216.7 K and 7.50 Torr
341 overlaid (red; this work). Additional noise in the Varanasi cross section is clearly visible.

342

343 Figure 6. A graphical representation of the PT coverage for both the new and Varanasi
344 datasets.

345

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529 **Tables**

530

531 Table 1: Summary of limb sounders past and present capable of measuring CFC-11.

Instrument	Platform	Years	
ATMOS (Atmospheric Trace MOlecule Spectroscopy)	Space shuttle	1985, 1992, 1993, 1994	Chang et al., 1996; Irion et al., 2002
CIRIS 1A (Cryogenic InfraRed Radiance Instrumentation for Shuttle)	Space shuttle	1991	Bingham et al., 1997
CRISTA (CRyogenic Infrared Spectrometers and Telescopes for the Atmosphere)	Space shuttle	1994, 1997	Offermann et al., 1999
CLAES (Cryogenic Limb Array Etalon Spectrometer)	UARS (Upper Atmosphere Research Satellite)	1991 – 1993	Roche, et al., 1993
ILAS (Improved Limb Atmospheric Spectrometer)	ADEOS (ADvanced Earth Observing Satellite)	1996 – 1997	Yokota, et al., 2002
ILAS II	ADEOS II	2003	Wetzel et al., 2006,
HIRDLS (High Resolution Dynamics Limb Sounder)	Aura	2004 – 2008	Hoffmann et al., 2014
MIPAS (Michelson Interferometer for Passive Atmospheric Sounding)	ENVISAT (ENVironmental SATellite)	2002 – 2012	e.g. Hoffmann et al., 2005; Dinelli et al., 2010; Kellmann et al., 2012
ACE-FTS (Atmospheric Chemistry Experiment – Fourier transform spectrometer)	SCISAT	2004 –	Brown et al., 2011

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535 Table 2: FTS parameters, sample conditions, and cell configuration for all measurements

Spectrometer	Bruker Optics IFS 125HR
Mid-IR source	Globar
Detector	Mercury cadmium telluride (MCT) D313 ^a
Beam splitter	Potassium bromide (KBr)
Optical filter	~700–1400 cm ⁻¹ bandpass
Spectral resolution	0.01 to 0.03 cm ⁻¹
Aperture size	3.15 mm
Apodisation function	Boxcar
Phase correction	Mertz
CCl ₃ F (Supelco)	99.9% purity, natural-abundance isotopic mixture; freeze-pump-thaw purified multiple times prior to use
Air zero (BOC Gases)	total hydrocarbons < 3 ppm, H ₂ O < 2 ppm, CO ₂ < 1 ppm, CO < 1 ppm; used ‘as is’
Cell pathlength	26 cm
Cell windows	Potassium bromide (KBr) (wedged)
Pressure gauges	3 MKS-690A Baratrons (1, 10 & 1000 Torr) (±0.05% accuracy)
Refrigeration	Julabo F95-SL Ultra-Low Refrigerated Circulator (with ethanol)
Thermometry	4 PRTs, Labfacility IEC 751 Class A
Wavenumber calibration	N ₂ O

536 “Due to the non-linear response of MCT detectors to the detected radiation, all interferograms
 537 were Fourier transformed using Bruker’s OPUS software with a non-linearity correction
 538 applied.
 539

540

541 Table 3: Summary of the sample conditions for all measurements.

Temperature (K)	Initial CCl ₃ F Pressure (Torr) ^a	Total Pressure (Torr)	Spectral resolution (cm ⁻¹) ^b
191.7 ± 0.8	0.266	7.535 ± 0.035	0.0100
191.5 ± 0.8	0.302	49.83 ± 0.13	0.0150
191.6 ± 0.8	0.302	98.14 ± 0.68	0.0225
191.6 ± 0.8	0.266	200.0 ± 0.3	0.0300
202.3 ± 0.5	0.319	7.508 ± 0.006	0.0100
202.4 ± 0.5	0.309	50.28 ± 0.13	0.0150
202.3 ± 0.5	0.318	99.85 ± 0.30	0.0150
202.4 ± 0.5	0.309	200.4 ± 0.2	0.0225
202.3 ± 0.5	0.318	301.6 ± 0.3	0.0300
216.7 ± 0.5	0.347	7.496 ± 0.018	0.0100
216.7 ± 0.5	0.358	49.93 ± 0.09	0.0150
216.7 ± 0.5	0.357	99.94 ± 0.07	0.0150
216.6 ± 0.5	0.375	201.0 ± 0.2	0.0225
216.7 ± 0.5	0.383	360.4 ± 0.3	0.0300
232.6 ± 0.4	0.407	7.500 ± 0.020	0.0100
232.6 ± 0.4	0.395	49.80 ± 0.15	0.0150
232.6 ± 0.4	0.544	99.67 ± 0.16	0.0150
232.6 ± 0.4	0.417	201.0 ± 0.1	0.0225
232.6 ± 0.4	0.413	399.8 ± 0.3	0.0300
252.5 ± 0.2	0.503	7.477 ± 0.003	0.0100
252.5 ± 0.2	0.486	50.06 ± 0.05	0.0150
252.5 ± 0.2	0.516	200.9 ± 0.1	0.0225
252.5 ± 0.2	0.544	399.9 ± 0.2	0.0300
252.5 ± 0.2	0.607	600.2 ± 0.3	0.0300
273.9 ± 0.2	0.475	7.501 ± 0.001	0.0100
273.8 ± 0.2	0.613	201.6 ± 0.1	0.0225
273.8 ± 0.2	0.598	355.8 ± 0.1	0.0300
273.8 ± 0.2	0.607	760.1 ± 0.2	0.0300
293.1 ± 0.1	0.548	355.8 ± 0.1	0.0300
293.0 ± 0.1	0.566	760.0 ± 0.1	0.0300

542 ^aMKS-690A Baratron readings are accurate to ± 0.05%.543 ^bUsing the Bruker definition of 0.9/MOPD.

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