

# Point-by-point response for referee comments on "Limited angle tomography of mesoscale gravity waves by the infrared limb-sounder GLORIA"

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## 1 Response to Referee Comment of Dr. Ole Christensen

P7I23: How is the accuracy and precision calculated? In particular what errors go into this error analysis? Are things like pointing error of the instrument, thermal noise, calibration error, instrument characterization errors, spectroscopic errors or errors in the assumptions of the background atmosphere included? As short list is required for a reader to be able to judge whether all significant terms are included.

A description of all errors implemented in the diagnostics is included in the manuscript. These errors include pointing errors, misrepresented background gases, uncertainties in spectral line characterization, calibration errors and measurement noise.

P8I12: please clarify what "25 and 60 neighboring points" means in terms of approximate distance or pressure.  
Corresponds to 750km in the horizontal and 3km in the vertical. Has been included in the text.

P14I10: It is a bit unclear to me that "horizontal orientation in this area... is not recovered". For figure 9a, 9e, 9i, it does indeed look like the wave fronts are oriented from north-west to south-east. However, a degradation may be seen, but I have a bit hard time connecting the text with the figure. Perhaps explicitly writing out the directional error in the text (20 degrees?) would make it easier for the reader to follow.

The following text is added to better guide the reader through the Figure: Within the tangent point region (longitude between 1.5° and 2.5°) the wave fronts are oriented almost west to east. Behind the tangent point region (longitude above 2.5°), the wave fronts slowly approach the expected west-north-west to east-south-east orientation. The orientation error of this retrieval ranges between 30° and 10° depending on the area of interest.

P18 figure 11: In panel f) there is a large response in the retrieved values far below the tangent points. This is rather counter-intuitive, and the reason for it (I assume its caused by the horizontal regularization) should be written out in the text.

This information has been included in the text, however at an earlier position (Section 3.1): The retrieval can also reproduce some signal outside the area covered by tangent points. One reason for that is, the path of the LOS which goes through higher altitudes before and after the tangent point, thus collecting information there. Another reason are the horizontal and vertical correlation lengths of 100 km and 1 km, which are used for the retrieval and smear out the signal.

## 5 2 Response to Referee Comment of Dr. Vern Hart

P2L23: You have described the FAT imaging volume as a cylinder, which is to be expected given the circular flight pattern described previously. You then go on to describe the LAT imaging region as simply a ‘3-D volume.’ I suggest this imaging region, its shape, and dimensions require more of an explanation. In the conclusion (P17L11), you refer to the FAT volume is ‘cubic.’ Is this a contradiction?

10 A more detailed explanation of the volume reconstructed with LAT has been added to Section 2.2, as it would require too much space in the introduction (P2L23). Figure 2a has been changed to improve the understanding of the volume recovered with LAT. The term cubic in the conclusion has been replaced by “cylindrical” to avoid confusion.

The LAT flightpath geometry should be described in more detail. Specifically, the vertical curvature of the flight path is  
15 never discussed. Did this cause inconsistencies in image resolution as the altitude between GWs and GLORIA varied? Were edge effects or ‘smearing’ a problem as the altitude increased?

A more detailed explanation of the volume reconstructed with LAT has been added to Section 2.2. Figure 2a has been changed to improve the understanding of the volume recovered with LAT. The flight altitude has been kept constant during the whole study. This information is added to the text. This is also mostly the case during real aircraft measurements of gravity  
20 waves. Thus, smearing effects due to altitude changes do not play a role here.

You describe a 200-km horizontal extent perpendicular to the flight path (P2L25). The resolution of overhead multi-angle images typically degrades moving towards the edge of the images (away from the nadir). Was this effect compensated for in any of your preprocessing steps?

25 For the GLORIA retrieval presented in this paper, a region of interest of 48 x 128 pixels is used, albeit the optical system is designed for a region of interest of 128 x 128 pixels. In this configuration, the imaging quality is at, or close to, the diffraction limit for every pixel used and the edge effects are thus not a dominant error source in the retrieval process. The co-addition of pixels further reduces the relative impact of edge effects and the spectral errors are mostly corrected during level 1 processing (Friedl-Vallon et al., 2014; Kleinert et al., 2014). Furthermore, every single pixel has a point spread function (PSF) which is  
30 about 1.9 arcmin wide. As the tangent points for lower altitudes are further away from the airplane, the projected vertical PSF measured in meters is wider for lower tangent altitudes. This effect is taken into account in the radiative-transfer and instrument model of the retrieval (Ungermann et al., 2015).

The initialization of tomography algorithms can have a profound effect on the resulting reconstructions and I believe your initialization (P6L27) needs to be described in far more detail. This climatological field ac should be discussed further. How was it developed? How do you know it is accurate?

5 The retrieval setup was described in more detail. Especially the construction of the Covariance matrices was given in more detail.

A detailed description of the used climatology and its accuracy is discussed in the given reference Remedios et al. (2007). The influence of the climatology on the result of the sensitivity study is negligible, because we are not interested in absolute temperature values but the perturbations. That is also one of the reasons, why the climatology is directly subtracted from the retrieval result before the analysis. Thus a detailed discussion of the climatology would exceed the scope of this paper.

10

It doesn't appear you have done any synthetic testing of your tomography algorithm. In this process, artificial projection (brightness) data are produced from a known synthetic structure and are then used to produce a 3D reconstruction. This result is compared with the original synthetic object to determine the accuracy of the algorithm. This approach can also be used to identify issues such as edge effects or biasing which results from the initialization model. Could you comment on this or perhaps include some sample synthetic results?

15

The main point of this paper is to test the tomographic algorithm with synthetic data. To clarify this, the terminology has been changed and the before called "true" structures have been renamed to "synthetic". The authors hope that this helps to distinguish synthetic retrieval tests (Sections 3.1 & 3.2) from tests with real measurements (Section 3.3). The effects, which are introduced by the retrieval algorithm are discussed in section 3.2 extensively and include for example the "V-shape".

20

P7L21: You state the horizontal resolution along the flightpath is 30 km. Is this the image resolution? This isn't consistent with a field of view spanning only 100-200 km.

The horizontal field-of view of the detector is around 7km, which is mentioned in Section 2.2. The 100-200km are not a resolution but a coverage which is gained by overlapping forward and rearward looking measurements. We improved the description of the measurement geometry in Section 2.2 to avoid this confusion.

25

You might consider quantifying the spatial agreement between the resulting images in Figure 11 (i.e., (b) and (e), (c) and (f)). Common metrics include the structural similarity index measurement (SSIM) or the Pearson correlation coefficient (PCC).

The following has been included in the text: A more quantitative comparison of the similarities of both retrievals can be given by the Pearson correlation coefficient. Including only areas which are covered by tangent points gives a Pearson correlation coefficient of 0.91. Expanding this area to places with measurement content larger than 0.8 (includes areas crossed by a LOS before or after the tangent point) still leads to a Pearson correlation coefficient of 0.75. Thus, as expected the two retrievals are highly correlated.

30

35 P7L21: Why did you decide to use values 'larger than half the maximum'?

The resolution describes which area around a measurement point which influence its value. However, this influence can in theory reach till infinity. Atmospheric values closer to a measurement point in general have higher influence than values far away. One way to determine the influence of the atmosphere around the measurement point is to look at the averaging kernel matrix. Points with AVK smaller than the fwhm have lower impact and can therefore be neglected. This method is already  
5 used in 1-D and 2-D nadir and limb sounding (Rodgers, 2000; von Clarmann et al., 2009). In this paper we just extended this method to 3-D.

I am a little confused by your statement in the caption of Figure 2. You state the line-of-sight measurements assume a parabolic shape due to the cartesian coordinate system. Are you implying they are straight on a curved surface and then warped  
10 when plotted using straight axes?

Yes, we are implying that the originally straight lines are wrapped, when using straight axes in x, y and z. This has been clarified in the text.

P10L4: In discussing tomographic retrieval of gravity waves from atmospheric data, you have only cited one study. However,  
15 there have been multiple studies on this topic, using data from satellites such as ODIN and AIM. I suggest expanding these references.

Additional citations were added.

You need to be clearer about the differences between the rows in Figure 3. (a) and (c) show true waves which are nearly  
20 identical. Why are the retrieved waves so different?

Figure 3 shows only one wave, but multiple cross sections through the 3-D volume. The first row shows horizontal cross sections, rows 2 and 3 vertical cross sections along  $0^{\circ}\text{N}$  and  $0^{\circ}\text{E}$ , respectively. This has been clarified in the caption.

P19L9: What kind of artifacts are you referring to? Can you describe them?

25 The retrieval creates v-shaped phase fronts. This is added to the manuscript.

It is unclear what 'lz' and 'lh' refer to in Figure 9.

'lz' and 'lh' stand for the vertical and horizontal wavelengths of the waves, respectively. This has been clarified in the caption.

## References

- Friedl-Vallon, F., Gulde, T., Hase, F., Kleinert, A., Kulesa, T., Maucher, G., Neubert, T., Olschewski, F., Piesch, C., Preusse, P., Rongen, H., Sartorius, C., Schneider, H., Schönfeld, A., Tan, V., Bayer, N., Blank, J., Dapp, R., Ebersoldt, A., Fischer, H., Graf, F., Guggenmoser, T., Höpfner, M., Kaufmann, M., Kretschmer, E., Latzko, T., Nordmeyer, H., Oelhaf, H., Orphal, J., Riese, M., Schardt, G., Schillings, J., Sha, M. K., Suminska-Ebersoldt, O., and Ungermann, J.: Instrument concept of the imaging Fourier transform spectrometer GLORIA, *Atmos. Meas. Tech.*, 7, 3565–3577, <https://doi.org/10.5194/amt-7-3565-2014>, 2014.
- Kleinert, A., Friedl-Vallon, F., Guggenmoser, T., Höpfner, M., Neubert, T., Ribalda, R., Sha, M. K., Ungermann, J., Blank, J., Ebersoldt, A., Kretschmer, E., Latzko, T., Oelhaf, H., Olschewski, F., and Preusse, P.: Level 0 to 1 processing of the imaging Fourier transform spectrometer GLORIA: generation of radiometrically and spectrally calibrated spectra, *Atmos. Meas. Tech.*, 7, 4167–4184, <https://doi.org/10.5194/amt-7-4167-2014>, 2014.
- Remedios, J. J., Leigh, R. J., Waterfall, A. M., Moore, D. P., Sembhi, H., Parkes, I., Greenhough, J., Chipperfield, M., and Hauglustaine, D.: MIPAS reference atmospheres and comparisons to V4.61/V4.62 MIPAS level 2 geophysical data sets, *Atmos. Chem. Phys. Discuss.*, 7, 9973–10017, <https://doi.org/10.5194/acpd-7-9973-2007>, 2007.
- Rodgers, C. D.: *Inverse Methods for Atmospheric Sounding: Theory and Practice*, vol. 2 of *Series on Atmospheric, Oceanic and Planetary Physics*, World Scientific, Singapore, 2000.
- Ungermann, J., Blank, J., Dick, M., Ebersoldt, A., Friedl-Vallon, F., Giez, A., Guggenmoser, T., Höpfner, M., Jurkat, T., Kaufmann, M., Kaufmann, S., Kleinert, A., Krämer, M., Latzko, T., Oelhaf, H., Olschewski, F., Preusse, P., Rolf, C., Schillings, J., Suminska-Ebersoldt, O., Tan, V., Thomas, N., Voigt, C., Zahn, A., Zöger, M., and Riese, M.: Level 2 processing for the imaging Fourier transform spectrometer GLORIA: derivation and validation of temperature and trace gas volume mixing ratios from calibrated dynamics mode spectra, *Atmos. Meas. Tech.*, 8, 2473–2489, <https://doi.org/10.5194/amt-8-2473-2015>, 2015.
- von Clarmann, T., De Clercq, C., Ridolfi, M., Höpfner, M., and Lambert, J.-C.: The horizontal resolution of MIPAS, *Atmos. Meas. Tech.*, 2, 47–54, <https://doi.org/10.5194/amt-2-47-2009>, 2009.

# Limited angle tomography of mesoscale gravity waves by the infrared limb-sounder GLORIA

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**Abstract.** Three-dimensional measurements of gravity waves are required, in order to quantify their direction resolved momentum fluxes and get a better understanding of their propagation characteristics. Such 3-D measurements of gravity waves in the lowermost stratosphere ~~can be~~ have been provided by the airborne Gimballed Limb Observer for Radiance Imaging of the Atmosphere (GLORIA) using full angle tomography. Closed flight patterns of sufficient size are needed to acquire the full set of angular measurements for full angle tomography. These take about two hours and are not feasible everywhere due to scientific reasons or air traffic control restrictions. Hence, this paper investigates the usability of limited angle tomography for gravity wave research based on synthetic observations. Limited angle tomography uses only a limited set of angles for tomographic reconstruction and can be applied on linear flight patterns. A synthetic end-to-end simulation has been performed to investigate the sensitivity of limited angle tomography towards gravity waves with different wavelengths and orientations in respect to the flight path. For waves with wave fronts roughly perpendicular to the flight path, limited angle tomography and full angle tomography can derive wave parameters like wavelength, amplitude and wave orientation with similar accuracy. For waves with horizontal wavelength above 200 km and vertical wavelength above 3 km, the wavelengths can be retrieved with less than 10% error, the amplitude with less than 20% error and the horizontal wave direction with an error below 10°. This is confirmed by a comparison of results obtained from full angle tomography and limited angle tomography for a real ~~measurement case~~ measurements taken on 25 January 2016 over Iceland. The reproduction quality of gravity wave parameters with limited angle tomography, however, depends strongly on the orientation of the waves with respect to the flight path. Thus, full angle tomography might be preferable in cases where the orientation of the wave cannot be predicted or waves with different orientations exist in the same volume and thus the flight path cannot be adjusted accordingly. Also for low amplitude waves and short scale waves full angle tomography has advantages due to its slightly higher resolution and accuracy.

## 20 1 Introduction

Gravity waves (GWs) couple the atmosphere vertically by transporting energy and momentum from the surface to altitudes as far as the mesosphere. On this way through the atmosphere, GWs interact with the mean flow and, thus, are responsible for the wind reversal in the mesosphere and influence prominent circulation patterns such as the quasi-biennial oscillation (QBO) of stratospheric tropical winds, and the meridional Brewer-Dobson circulation (BDC) in the stratosphere. These circulation

patterns can then affect surface temperature and pressure patterns (Sigmond and Scinocca, 2010; Kidston et al., 2015; Sandu et al., 2016; Scaife et al., 2016).

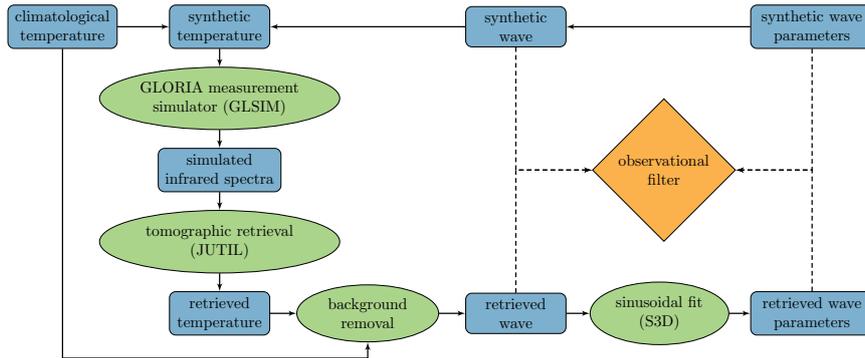
Due to their small scales, GWs are implemented into many climate projection and weather prediction models in the form of simplified sub-models, so-called GW parameterisations. For practical reasons, these parameterisations assume solely vertical propagation of GWs. However, multiple studies highlight the importance of 3-D propagation (Preusse et al., 2009; Sato et al., 2009; McLandress et al., 2012; Kalisch et al., 2014; Ribstein and Achatz, 2016). To gain a better understanding of 3-D propagation and improve the GW parameterisations in models, 3-D measurements of GWs are required (Geller et al., 2013).

The Gimballed Limb Observer for Radiance Imaging of the Atmosphere (GLORIA) can provide such 3-D measurements of gravity waves (Krisch et al., 2017). GLORIA is an airborne infrared limb sounder, which can change its horizontal viewing direction from  $45^\circ$  (right forward) to  $135^\circ$  (right backward) with respect to the aircraft's heading (Friedl-Vallon et al., 2014; Riese et al., 2014). If a volume of air is measured from all surrounding angles, it may be reconstructed using tomographic methods (Natterer, 2001). Measuring emitted radiation from all  $360^\circ$  around the volume, for instance, by flying a circle, is called full angle tomography (FAT). This technique improves the horizontal resolution of limb sounders by an order of magnitude (Ungermaun et al., 2010b). In contrast to FAT, limited angle tomography (LAT) does not measure the volume from all sides but only from a limited set of angles. Due to the horizontal scanning capabilities of the GLORIA instrument, this is already possible on a linear flight path. However, LAT inversion problems are in general seriously ill-posed (Natterer, 2001).

Since flying in a circular pattern of sufficient size can take more than two hours, FAT is only suitable for measurements in steady atmospheric states, where the conditions do not change during the acquisition time. Accounting for the change of the atmosphere during acquisition is possible for trace gas retrievals by including advection (Ungermaun et al., 2011). The temperature structure, however, is governed by a multitude of waves with different spatial and temporal scales. Our apriori knowledge of the temporal development of these waves is not sufficient to retrieve a fast changing temperature structure using FAT. Furthermore, air traffic control restrictions and trade-offs with other instruments can force linear flight patterns.

Using FAT, a cylindrical 3-D volume with about 400 km ~~horizontal~~horizontal diameter and several kilometres in altitude can be reconstructed (Ungermaun et al., 2011; Kaufmann et al., 2015; Krisch et al., 2017). In contrast, the 3-D volume, which is reconstructed using LAT, has a ~~horizontal extent of only 100–200~~more complex structure with horizontal extent in the viewing direction of about 150 km ~~perpendicular to the flight path depending on the altitude~~ (Ungermaun et al., 2011). Is this smaller horizontal extent sufficient to derive the 3-D orientation of GWs? In this paper we will investigate, how well the wave structures are retrieved using LAT and to which accuracy wave parameters can be derived from LAT temperature retrievals.

In order to assess the differences between FAT and LAT using the GLORIA instrument quantitatively, we will derive the observational filter of both methods. The observational filter is a measure for the sensitivity of an instrument towards measuring different GWs and should be taken into account when comparing measurements from different instruments or measurements with model results (Alexander, 1998; Preusse et al., 2002; Ern et al., 2006; Trinh et al., 2016). In general, an observational filter tells by how much the amplitude or the momentum flux is underestimated by the measurement technique. As it is possible to reconstruct 3-D volumes with GLORIA and, thus, derive a 3-D wave vector, we extend the concept of the observational filter. Besides the usual observational filter regarding the wave amplitude, we introduce an observational filter for the wave



**Figure 1.** Methodological concept of the performed simulation study. A detailed description can be found in Sec. 2.1.

orientation. This observational filter for the wave orientation with respect to the flight direction is important for the planning of research flights: The wave orientation can often be predicted and the flight path adjusted respectively.

This paper is structured as follows: In Sec. 2 a description of the methodological concept on how to derive an observational filter (Sec. 2.1) is followed by a short introduction to the concept of limb sounding and the GLORIA instrument (Sec. 2.2).

- 5 Afterwards, the methods used for the estimation of the observational filter are described in detail: the GLORIA measurement simulator (Sec. 2.3), the tomographic retrieval concept (Sec. 2.4), the background removal algorithm (Sec. 2.5) and the three-dimensional wave fitting routine (Sec. 2.6). Finally, a definition of the observational filter is given in Sec. 2.7. Section 3 presents and discusses the results of the simulation study for FAT (Sec. 3.1) and LAT (Sec. 3.2) and compares the measurement quality of both methods for a real GLORIA measurement case on 25 January 2016 over Iceland (Sec. 3.3).

## 10 2 Methods

### 2.1 Methodological concept

- The goal of this simulation study is to determine the capability of the GLORIA infrared limb imager to measure mesoscale GWs with limited angle tomography (LAT). The accuracy of reconstructing GW parameters, such as horizontal and vertical wavelength, amplitude and wave orientation is studied. For this purpose, an end-to-end simulation was performed which is
- 15 described in this section. Figure 1 shows the concept of this end-to-end simulation.

To simulate a realistic atmosphere, a complete climatological field  $\mathbf{a}_c \in \mathbb{R}^n$  from Remedios et al. (2007) was used. The climatological temperature field  $T_c \in \mathbb{R}^m$  was perturbed at each point  $\mathbf{x}_i \in \mathbb{R}^3$  in space by a simulated, so-called true synthetically generated wave field

$$w_{ts}^i = \hat{T} \cdot \sin(\mathbf{k}\mathbf{x}_i + \phi), \quad (1)$$

where  $\hat{T}$  is the temperature amplitude,  $\mathbf{k} \in \mathbb{R}^3$  the 3-D wave vector, and  $\phi$  the phase of the wave. In the present simulation study the temperature amplitude  $\hat{T}$  is arbitrarily chosen to be 3 K, the wave phase  $\phi$  to be  $0^\circ$ . For an improved conception, from now on, the 3-D wave vector  $\mathbf{k} = (k_0, k_1, k_2)$  will mainly be expressed in terms of vertical wavelength  $\lambda_z = \frac{2\pi}{k_2}$ , horizontal wavelength  $\lambda_h = \frac{2\pi}{\sqrt{k_1^2 + k_0^2}}$ , and horizontal wave direction  $\varphi = \arctan2(k_1, k_0)$ . Thus, if a downward pointing wave vector

5 ( $\lambda_z > 0$ , upward propagating wave) is assumed, a wave with wave direction of  $0^\circ$  has east-west oriented wavefronts and is tilted towards the north, a wave with wave direction of  $180^\circ$  is tilted towards the south. These ~~true synthetic~~ waves are added to the climatological temperature field to gain the ~~so-called true temperature field~~  $\mathbf{T}_r \in \mathbb{R}^m$  ~~a synthetic temperature field~~  $\mathbf{T}_s \in \mathbb{R}^m$ .

From this predetermined atmospheric state and with a given flightpath, the GLORIA measurement simulator (Sec. 2.3) calculates a set of infrared spectra, as would be measured by the GLORIA instrument. A tomographic retrieval (Sec. 2.4) is

10 then performed using these simulated infrared spectra. This retrieval uses only a well-defined set of infrared radiances (Tab. 3) and can reconstruct the atmosphere only in a reduced area, limited by the measurement geometry.

A background removal algorithm then subtracts the climatological temperature field  $\mathbf{T}_c$  from the retrieved temperature field  $\mathbf{T}_r \in \mathbb{R}^m$  to obtain the retrieved wave structure  $\mathbf{w}_r \in \mathbb{R}^m$ . In a real measurement case (Sec. 3.3) the background field is unknown and has to be identified by mathematical filtering methods (Sec. 2.5). To solely investigate the sensitivity of the

15 measurement concept and exclude any additional effects, these filtering methods are not used for the simulation study.

Finally, the retrieved wave structure is compared to the ~~original true synthetic~~ wave structure. By repeating this process for different horizontal and vertical wavelengths, the observational filter of LAT is established (Sec. 2.7). To interpret the retrieved wave structure with respect to gravity waves, the wave parameters amplitude, phase and wave vector have to be derived, using the small-volume few-wave decomposition method S3D (Sec. 2.6 and Lehmann et al., 2012). Comparing these retrieved wave

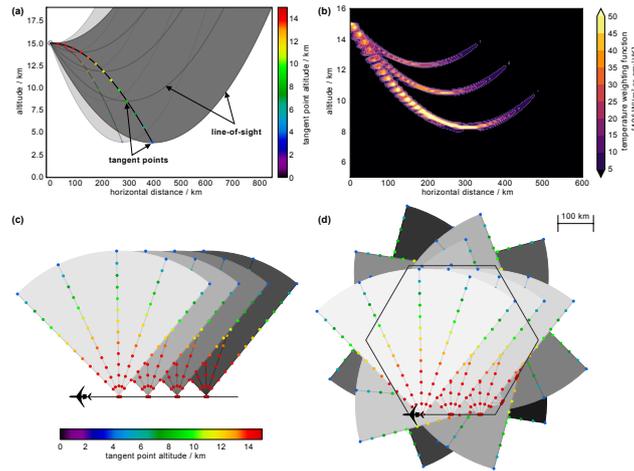
20 parameters to the prescribed ~~true synthetic~~ wave parameters gives detailed information on the usability of the different retrievals for GW research.

## 2.2 The GLORIA infrared limb imager

The Gimballed Limb Observer for Radiance Imaging of the Atmosphere (GLORIA) is an airborne Fourier Transform Spectrometer (FTS) which combines a classical Michelson interferometer with a 2-D detector array (Friedl-Vallon et al., 2014).

25 GLORIA measures the infrared radiation in the spectral range from 780 to 1400  $\text{cm}^{-1}$ , which is emitted by molecules in the atmosphere along the line-of-sight (LOS). The interferometer spectrally resolves this radiation to reveal characteristic molecular emissions. ~~A schematic of the limb sounding geometry is depicted in Fig. 2a.~~ Due to the exponentially declining density of the atmosphere with altitude, most radiation along the LOS is emitted at lower altitudes and, thus, around the tangent point. Moreover, for geometrical reasons, a comparatively long part of the LOS samples altitudes close to the tangent point, while

30 higher atmospheric layers are passed only briefly (Fig. 2a). As a consequence, conventional limb sounders are more sensitive to changes in the atmosphere around the tangent point (Fig. 2b). The horizontal resolution of conventional limb sounders along the line-of-sight is roughly 200–300 km (von Clarmann et al., 2009; Ungermann et al., 2012). In flight direction the horizontal resolution of 1D retrievals, which mainly depends on the horizontal field of view, is on the order of several kilometres for the airborne limb imager GLORIA.



**Figure 2.** In Panel (a) a simple schematic of the limb sounding geometry is given. The lines-of-sight (airplane is flying into the paper plane. Images taken under  $90^\circ$  azimuth cover the dark grey area with LOS. The respective tangent points (bright coloured dots) increase in distance with decreasing altitude. The tangent points of forward and rearward looking images (light grey and pale coloured dots) are closer to the flight path. The LOS, which are straight lines in reality, have a parabolic shape in this plot due to the transformation into a cartesian coordinate system used for this plot with x-axis following the earth surface. Panel (b) shows the weighting function along three different LOS indicating the contribution of the respective part of the atmosphere to the observed signal. In panels (c) and (d) the principles of limited angle tomography (LAT) and full angle tomography (FAT) with the GLORIA instrument are depicted, respectively. Shown are top views in bird perspective onto the flight path. The dots again indicate the tangent points and are coloured according to their altitude. Each grey sector indicates one horizontal scan from  $45^\circ$  to  $135^\circ$ . The lighter the grey, the later in time are these measurements taken.

GLORIA operates in two different modes: The chemistry mode, which has a high spectral sampling of  $0.0625 \text{ cm}^{-1}$ , and the dynamics mode with a coarser spectral sampling of only  $0.625 \text{ cm}^{-1}$ . However, the coarser spectral sampling leads to a faster interferogram acquisition and accordingly an improved spatial sampling ( $0.45 \text{ km}$  instead of  $2.25 \text{ km}$ ). This improved spatial sampling is used to scan the atmosphere horizontally in steps of  $4^\circ$  from  $45^\circ$  (right forward) to  $135^\circ$  (right backward) with respect to the aircraft's heading. In this way, the same volume of air is measured under different angles, which allows for tomographic retrievals (Sec. 2.4). The tangent points of forward (azimuth of  $45^\circ$ ) and rearward (azimuth of  $135^\circ$ ) looking images are closer to the flight path than for images with  $90^\circ$  azimuth (Fig. 2). Thus, due to the panning ability, the tangent point covered area of GLORIA is a banana-shaped curtain with horizontal extent across flight track around  $100\text{--}200 \text{ km}$ . Inside this banana-shaped curtain, the horizontal resolution across flight track is improved by a factor of 2 compared to conventional limb sounders. Thus, the panning ability of GLORIA improves both the horizontal sampling and the resolution across flight track.

Figure 2c and d show the measurement concept a top view of the measurement concepts for LAT and FAT used in this paper, respectively. In FAT (Fig. 2d), the full volume of the hexagon is covered by measurements from all  $360^\circ$  around. The tangent points cover a cylindrical 3-D volume inside the hexagonal flight path. In LAT however (Fig. 2c), the air is sampled by measurements from only one side, thus,  $90^\circ$  instead of  $360^\circ$ . Thus the volume covered by tangent points has a banana shape in the

vertical with horizontal extent across flight track between 100–200 km depending on altitude. In the along flight track direction the shape is extended as far as the aircraft flies. A detailed description of the different GLORIA measurement concepts can be found in Riese et al. (2014).

### 2.3 GLORIA measurement simulator

- 5 The GLORIA measurement simulator can replicate an infrared spectrum which GLORIA would measure on a flight through a given atmospheric state. In the present paper, for the LAT cases, a linear flight along the zero meridian from 5° S to 5° N at a constant flight altitude of 15 km was arbitrarily chosen. For the FAT cases, a hexagonal flight pattern with 400 km diameter around 0° N and 0° E again at constant flight altitude of 15 km is used. The ~~true atmospheric state~~  $\mathbf{a}_t \in \mathbb{R}^m$  ~~synthetic atmospheric state~~  $\mathbf{a}_s \in \mathbb{R}^m$  is composed of a temperature field, a pressure field, and the distributions of several trace gases. The climatological atmospheric state  $\mathbf{a}_c$  from Remedios et al. (2007) is transformed into the ~~true atmospheric state~~  $\mathbf{a}_t$  ~~synthetic atmospheric state~~  $\mathbf{a}_s$  by inserting the ~~true temperature field~~  $T_t$  ~~synthetic temperature field~~  $T_s$  (see Sec 2.1). A radiative transfer model  $\mathbf{F} : \mathbb{R}^n \rightarrow \mathbb{R}^o$  maps this ~~true synthetic~~ atmospheric state onto a set of radiances  $\mathbf{F}(\mathbf{a}_t)\mathbf{F}(\mathbf{a}_s)$ . This radiative transfer model also includes instrument and geometry effects such as increase of the field-of-view of single detector pixels along LOS and the instrument line shape (Ungermann et al., 2015). As GLORIA can measure radiances only with finite precision, a measurement error  $\epsilon \in \mathbb{R}^o$  is added to get the simulated radiances  $\mathbf{y} \in \mathbb{R}^o$ :

$$\mathbf{y} = \mathbf{F}(\mathbf{a}_{ts}) + \epsilon \quad (2)$$

This straight forward calculation is done with the Juelich Rapid Spectral Simulation Code V2 (JURASSIC2) for each pixel of the GLORIA detector along the chosen flight path leading to a set of thousands of simulated infrared spectra.

- This model was developed to efficiently handle imager instruments and tomographic retrievals. It is based on JURASSIC (Hoffmann, 2006), which was previously used as forward model for the evaluation of several satellite- and air-borne remote sensing experiments (e.g. Eckermann et al., 2006; Hoffmann et al., 2008; Weigel et al., 2010; Kalicinsky et al., 2013; Ungermann et al., 2016). It contains several approaches of varying computational complexity and accuracy for computing radiances, but we employ in this work the fast table based approach based on the emissivity growth approximation (EGA; e.g. Weinreb and Neuendorffer, 1973; Gordley and Russell, 1981).

### 2.4 Tomographic temperature retrieval

The Juelich Tomographic Inversion Library (JUTIL) software package is used for mapping the simulated infrared spectra back to the geophysical quantities, in our case the retrieved atmospheric state  $\mathbf{a}_r \in \mathbb{R}^n$  including the retrieved temperature  $T_r$ . This retrieval represents an ill-posed problem, which is solved by approximating it with a well-posed one using a Tikhonov regularisation scheme (Tikhonov and Arsenin, 1977). This leads to the minimisation problem:

$$J(\mathbf{a}) = (\mathbf{F}(\mathbf{a}) - \mathbf{y})^T \mathbf{S}_\epsilon^{-1} (\mathbf{F}(\mathbf{a}) - \mathbf{y}) + (\mathbf{a} - \mathbf{a}_a)^T \mathbf{S}_a^{-1} (\mathbf{a} - \mathbf{a}_a) \rightarrow \min \quad (3)$$

**Table 1.** Standard deviations, vertical ( $c_z$ ) and horizontal ( $c_h$ ) correlation lengths used for the tomographic retrieval.

<u>atmospheric quantity</u>	<u><math>\sigma</math></u>	<u><math>w_0</math></u>	<u><math>w_1</math></u>	<u><math>c_z</math></u>	<u><math>c_h</math></u>
<u>temperature</u>	<u>1 K</u>	<u><math>10^{-3}</math></u>	<u><math>10^{-2}</math></u>	<u>1.0 km</u>	<u>100 km</u>
<u><math>O_3</math></u>	<u>141 ppbV</u>	<u><math>10^{-6}</math></u>	<u><math>10^{-5}</math></u>	<u>8.0 km</u>	<u>6400 km</u>
<u><math>CCl_4</math></u>	<u>13 pptV</u>	<u><math>10^{-5}</math></u>	<u><math>10^{-4}</math></u>	<u>2.0 km</u>	<u>800 km</u>
<u><math>HNO_3</math></u>	<u>987 pptV</u>	<u><math>10^{-7}</math></u>	<u><math>10^{-4}</math></u>	<u>3.2 km</u>	<u>1280 km</u>

with  $\mathbf{S}_\epsilon \in \mathbb{R}^{o \times o}$  the measurement error covariance matrix and  $\mathbf{S}_a \in \mathbb{R}^{n \times n}$  the covariance matrix of the atmospheric state vector. As an apriori state  $\mathbf{a}_a$  the climatological field  $\mathbf{a}_c$  is used.

The first term of the costfunction represents the inversion of the forward model, which can have many different mathematical solutions. To choose a physically meaningful solution, a regularisation term using a covariance matrix  $\mathbf{S}_a$  is added. This covariance matrix is constructed as follows:

$$\mathbf{S}_a^{-1} = \frac{w_0}{\sigma^2} \|\mathbf{a}\|^2 + \frac{w_1}{\sigma^2} \left( \left\| c_z \frac{\partial}{\partial z} \mathbf{a} \right\|^2 + \left\| c_h \frac{\partial}{\partial x} \mathbf{a} \right\|^2 + \left\| c_h \frac{\partial}{\partial y} \mathbf{a} \right\|^2 \right). \quad (4)$$

The standard deviations  $\sigma$ , weighting factors  $w_0$  and  $w_1$ , and the correlations lengths  $c_z$  and  $c_h$  used for the retrieval are given in Tab. 1. These factors are chosen ad-hoc and cannot be interpreted directly as physically meaningful correlation lengths. A more physical regularisation scheme is currently under development and will be described by Krasauskas et al. (2018).

15 This minimization problem is solved with a truncated conjugate gradient based trust region scheme. More details on the retrieval algorithms used for GLORIA Level 2 processing are described by Ungermann et al. (2015).

Since the temperature perturbations due to the wave are small compared to the background temperature  $T_c$ , the retrieval can be linearized around this background temperature (Rodgers, 2000; Ungermann et al., 2010a):

$$\mathbf{y} - \mathbf{y}_a = \mathbf{F}'(\mathbf{a}_a)(\mathbf{a} - \mathbf{a}_a) + \epsilon. \quad (5)$$

15  $\mathbf{F}'(\mathbf{a}_a) = \left. \frac{\partial \mathbf{F}}{\partial \mathbf{a}} \right|_{\mathbf{a}_a}$  is the Jacobian matrix of the forward model evaluated at  $\mathbf{a}_a$  and  $\mathbf{y}_a = F(\mathbf{a}_a)$  are the simulated radiances of the background state. With the retrieval gain matrix  $\mathbf{G}(\mathbf{a}_a) = (\mathbf{F}'(\mathbf{a}_a)^T \mathbf{S}_\epsilon^{-1} \mathbf{F}'(\mathbf{a}_a) + \mathbf{S}_a^{-1})^{-1} \mathbf{F}'(\mathbf{a}_a)^T \mathbf{S}_\epsilon^{-1}$  and the Jacobian matrix  $\mathbf{F}'(\mathbf{a}_a)$  the averaging kernel matrix  $\mathbf{A}(\mathbf{a}_a) = \mathbf{G}(\mathbf{a}_a) \mathbf{F}'(\mathbf{a}_a)$  can be calculated, which converts the **true temperature perturbation  $\mathbf{w}_t = \mathbf{a}_t - \mathbf{a}_a$**  synthetic temperature perturbation  $\mathbf{w}_s = \mathbf{a}_s - \mathbf{a}_a$  into the retrieved temperature perturbation  $\mathbf{w}_r = \mathbf{a}_r - \mathbf{a}_a$ :

$$20 \quad \mathbf{G}(\mathbf{a}_a)(\mathbf{y} - \mathbf{y}_a) = \mathbf{G}(\mathbf{a}_a)(\mathbf{F}'(\mathbf{a}_a)(\mathbf{a} - \mathbf{a}_a) + \epsilon) \quad (6)$$

$$\mathbf{a}_r - \mathbf{a}_a = \mathbf{A}(\mathbf{a}_a)(\mathbf{a}_{ts} - \mathbf{a}_a) + \mathbf{G}(\mathbf{a}_a)\epsilon \quad (7)$$

$$\mathbf{w}_r = \mathbf{A}(\mathbf{a}_a)\mathbf{w}_{ts} + \mathbf{G}(\mathbf{a}_a)\epsilon. \quad (8)$$

**Table 2.** Systematic error sources included in the retrieval diagnostics with respective standard deviations and correlation lengths.

Error source	standard deviation	correlation lengths
pointing errors	0.05°	vertical: $\infty$ , temporal: 0
misrepresented background gases	taken from Remedios et al. (2007)	taken from Remedios et al. (2007)
uncertainties in spectral line characterization	5%	temporal: $\infty$
calibration errors gain	1%	temporal: $\infty$
calibration errors offset	5 nW	temporal: 0

The noise term  $\mathbf{G}(\mathbf{a}_a)\boldsymbol{\epsilon}$  is negligible with regard to the other terms and is disregarded in the present study. For selected cases, the linear approximation has been validated by a comparison of linear and non-linear retrieval results. The retrievals of the real measurements in Sec. 3.3 are non-linear.

Further, the gain matrix  $\mathbf{G}(\mathbf{a}_a)$  is used to calculate the influence of an arbitrary error source described by a covariance matrix  $\mathbf{S}_\epsilon$  on the retrieval result  $\mathbf{G}^{-1}\mathbf{S}_\epsilon\mathbf{G}$ . Covariance matrices describing different systematic error sources are assembled using an auto-regressive approach with reasonable standard deviations and correlation lengths (Tarantola, 2004). The standard deviations and correlation lengths used for the different systematic errors are summarized in Tab. 2. The effects of instrument errors on the GLORIA retrieval are discussed in detail by Kleinert et al. (2018). The covariance matrix for measurement noise is taken from theoretical estimates given by Friedl-Vallon et al. (2014) that agree well with estimates derived from real measurements (Kleinert et al., 2014).

As most systematic errors are fully correlated over all measurements of one flight and the noise terms are negligible in magnitude compared to the wave perturbation terms, the error term  $\mathbf{G}(\mathbf{a}_a)\boldsymbol{\epsilon}$  is disregarded in the simulation study.

Conventional infrared temperature retrievals for limb sounding instruments are based on optically thin spectral lines for example in the CO<sub>2</sub> Q-branch region at 790.75 cm<sup>-1</sup> (12.6 μm) (Riese et al., 1997; Ungermann et al., 2010a). Nadir sounders in contrast use spectral lines with different opacity to improve the vertical resolution (Hoffmann and Alexander, 2009). Transferring this concept to limb sounding and including additional lines with high opacity into limb sounding retrievals increases the resolution along LOS (Ungermann et al., 2011). Thus, including some ozone emission lines between 980 cm<sup>-1</sup> and 1014 cm<sup>-1</sup> with different optical depths in our retrieval and applying LAT improved the horizontal resolution in viewing direction at 10.5 km altitude to 70 km. The resolution is derived through fitting a 3-D ellipsoid around all points of the averaging kernel matrix  $\mathbf{A}$  larger than half the maximum. The horizontal resolution along the flight path is 30 km and the vertical resolution 400 m. The FAT retrievals have a horizontal resolution of 20 km in both directions and a vertical resolution of 200 m. The precision (random error) of both methods within the tangent point area is below 0.05 K, the accuracy (sum over all systematic errors) of FAT below 0.5 K, and the accuracy of LAT below 0.7 K.

For the GW sensitivity study a retrieval setup with the spectral ranges 1 to 7 in Tab. 3 is used. For the real measurement retrievals in Sec. 3.3, the spectral ranges 8 and 9 in Tab. 3 are included in the retrieval to improve the knowledge about the CCl<sub>4</sub>

**Table 3.** Spectral ranges used for the retrievals presented in this paper. The last column indicates the retrieved quantity for each spectral range. For the simulation study (Sec. 3.1 and 3.2), the spectral ranges 1 to 7 are used. For the real measurement retrievals (Sec. 3.3), the spectral ranges 8 to 13 are added.

	Spectral range / $\text{cm}^{-1}$			Used for
1	790.625	–	791.250	temperature
2	791.875	–	792.500	temperature
3	956.875	–	962.500	temperature
4	980.000	–	984.375	temperature, $\text{O}_3$
5	992.500	–	997.500	temperature, $\text{O}_3$
6	1000.625	–	1006.250	temperature, $\text{O}_3$
7	1010.000	–	1014.375	temperature, $\text{O}_3$
8	793.125	–	795.000	$\text{CCl}_4$
9	796.875	–	799.375	$\text{CCl}_4$
10	883.750	–	888.125	$\text{HNO}_3$
11	892.500	–	896.250	$\text{HNO}_3$
12	900.000	–	903.125	$\text{HNO}_3$
13	918.750	–	923.125	$\text{HNO}_3$

background radiation in the  $\text{CO}_2$  Q-branch region. Furthermore, the spectral ranges 10 to 13 are used additionally to retrieve the trace gas  $\text{HNO}_3$ .

To improve the convergence speed and the quality of the real measurement retrievals, apriori fields are taken from different models. The temperature apriori was constructed from European Centre for Medium-Range Weather Forecasts (ECMWF) operational analyses at resolution T1279/L137 by applying the background removal described in Sec. 2.5. ~~The background removal ensures that any GW signature~~ This process is designed to remove all GW signatures from the apriori field. Thus, GW signatures in the retrieval result does not originate from the apriori field originate almost entirely from measurement content. The pressure field was taken directly from ECMWF. The apriori fields of several trace gases ( $\text{CH}_4$ ,  $\text{CO}_2$ ,  $\text{H}_2\text{O}$ ,  $\text{O}_3$ , ...) are taken from the Whole Atmosphere Community Climate Model version 4 (WACCM4).

## 10 2.5 Background removal

The atmospheric temperature distribution is mainly determined by the stratification and the balanced flow. However, GWs cause small-scale perturbations to this background temperature structure. Before identifying GWs with wave-fitting algorithms, a scale separation of large-scale background and small-scale perturbations has to be performed. This scale separation is called background removal.

For the simulation study in Sec. 3.1 and Sec. 3.2 the background temperature is known to be the climatological temperature field  $T_c$ . The background removal subtracts this temperature field  $T_c$  from the retrieved temperature field  $T_r$  and the temperature residual or so-called retrieved wave field  $w_r$  remains.

For the retrievals of real measurements presented in Sec. 3.3, one-dimensional Savitzky-Golay filters (Savitzky and Golay, 1964) are applied in all three spatial directions with third-order polynomials over 25 and 60 neighbouring points in vertical and both horizontal directions, respectively. This corresponds to 750 km in both horizontal directions and 3 km in the vertical. In this way, a spatial separation in into large scale background and small scale temperature residuals is achieved.

## 2.6 Three-dimensional sinusoidal fitting routine (S3D)

To compare the retrieval results with the original waves and interpret the structures with regard to gravity waves, wave parameters (horizontal and vertical wavelengths, wave amplitude and wave direction) have to be derived. This is done in overlapping sub-volumes of 5 km vertical and 400 km x 400 km horizontal extent. In these sub-volumes a sinusoid is fitted to the retrieved wave field  $w_r$  using a least square method (Lehmann et al., 2012). In this process, the following equation is minimized:

$$\chi^2 = \sum_i \frac{(f(\mathbf{x}_i) - w_r(\mathbf{x}_i))^2}{\sigma^2(\mathbf{x}_i)} \quad (9)$$

with the sinusoidal function

$$f(\mathbf{x}) = \hat{T} \cdot \sin(\mathbf{k}\mathbf{x} + \phi) = A \cdot \sin(\mathbf{k}\mathbf{x}) + B \cdot \cos(\mathbf{k}\mathbf{x}), \quad (10)$$

and a weighting function  $\sigma^2(\mathbf{x})$ , which is chosen to be 1 if a tangent point exists in this grid cell and  $10^5$  if not. Systematic tests of this algorithm by superposition of two sinusoids show that even for wavelengths up to 2.5 times the cube size, the wave parameters of both waves are fitted with errors below 1%.

Due to the fact that the wave structures of the real measurements vary strongly in space, a smaller cube size of 3.6 km x 160 km x 160 km is chosen for the S3D fits presented in Sec. 3.3. This cube size is sufficient to derive vertical wavelengths up to 9 km and horizontal wavelengths up to 400 km.

## 2.7 Observational filter

The relative error of the retrieval within an area A can be calculated as follows:

$$S = \sum_{\mathbf{x}_i \in A} \frac{\underbrace{w_t(\mathbf{x}_i) - w_r(\mathbf{x}_i)}_{\text{red}}}{\underbrace{w_t(\mathbf{x}_i)}_{\text{red}}} \frac{\underbrace{w_s(\mathbf{x}_i) - w_r(\mathbf{x}_i)}_{\text{blue}}}{\underbrace{w_s(\mathbf{x}_i)}_{\text{blue}}} \quad (11)$$

For a fair comparison, the area A must be chosen in a way, that it covers the measurement region, meaning a region covered with tangent points. In our case, the area A was chosen to be between 9.5–11.5 km altitude,  $1.75^\circ - 2.25^\circ$  longitude, and  $-1^\circ - 1^\circ$  latitude for LAT and 9.5–11.5 km altitude,  $-1^\circ - 1^\circ$  longitude, and  $-1^\circ - 1^\circ$  latitude for FAT.

This relative error is a helpful measure for the reproducibility of GWs by the measurement setup and the retrieval concept. However, it does not give detailed information upon which wave parameters can be derived. Thus, we further define more

specific relative errors for the important gravity wave parameters, horizontal  $\lambda_h$  and vertical  $\lambda_z$  wavelengths, amplitude  $\hat{T}$  and horizontal wave orientation  $\varphi$ :

$$S_\xi = \sum_{\mathbf{x}_i \in A} \frac{\xi_t(\mathbf{x}_i) - \xi_r(\mathbf{x}_i)}{\xi_t(\mathbf{x}_i)} \frac{\xi_s(\mathbf{x}_i) - \xi_r(\mathbf{x}_i)}{\xi_s(\mathbf{x}_i)} \text{ for } \xi \in \lambda_h, \lambda_z, \hat{T}, \varphi \quad (12)$$

These more specific relative errors help to define the quality of our measurement for gravity wave analysis.

- 5 The observational filter  $O = 1 - S$  is a measure of the sensitivity of an instrument and defines which gravity waves can be detected. The knowledge of the observational filter is necessary for meaningful comparisons of measurements from different instruments or measurement and model results (Alexander, 1998; Preusse et al., 2002; Ern et al., 2006; Ern and Preusse, 2012; Trinh et al., 2016).

### 3 Results and discussion

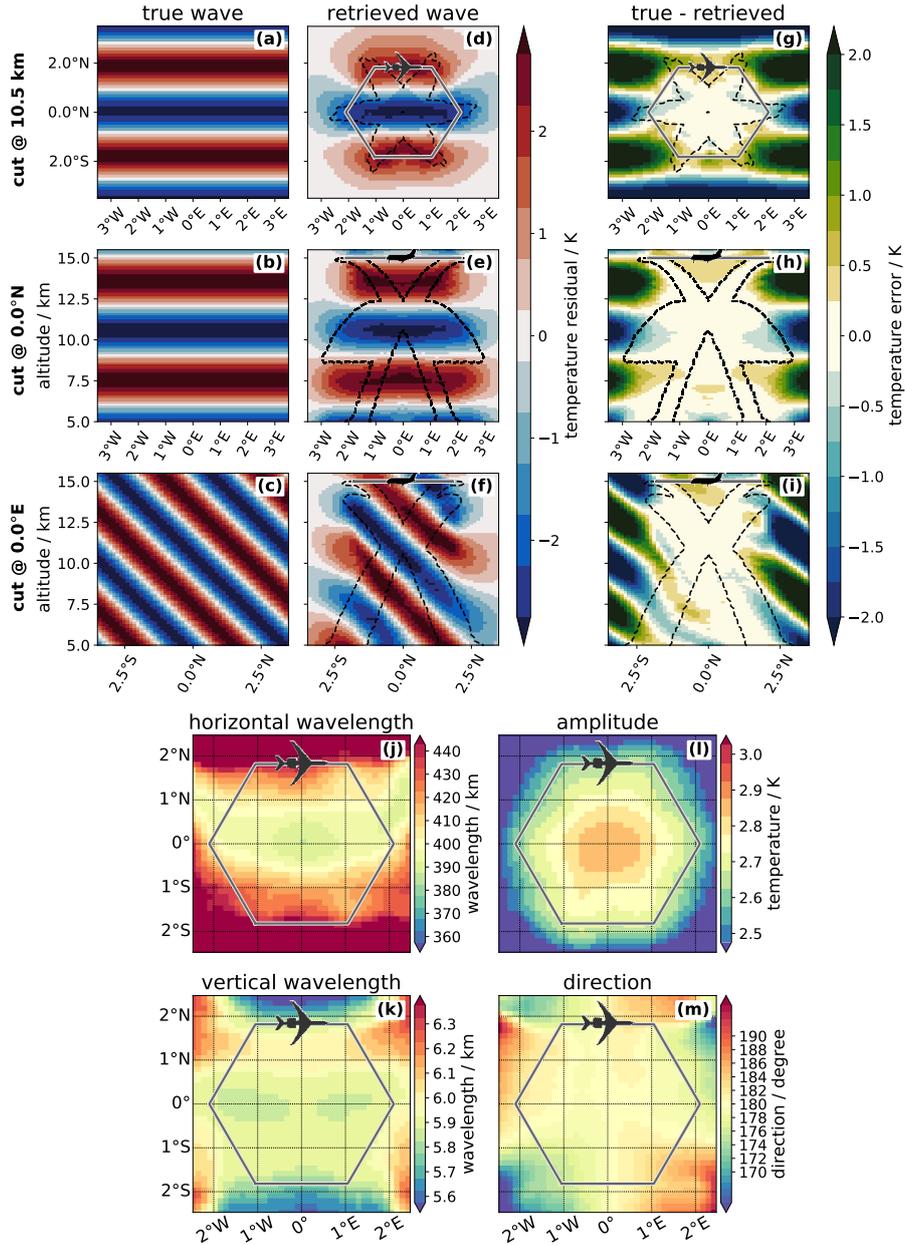
#### 10 3.1 Full angle tomography (FAT)

Former studies demonstrated the feasibility of 2-D tomography of rearward looking satellite instruments for the retrieval of 3-D atmospheric structures such as gravity waves (Degenstein et al., 2004; Ungermann et al., 2010a; Hultgren et al., 2013). However, the concept of 3-D tomography with sideways looking airborne instruments (Riese et al., 2014) or sub-limb instruments (Song et al., 2017; Hart et al., 2018) is quite different and, thus, may exhibit different characteristics. In 2-D tomography a volume is reconstructed from rearward looking measurements on a moving platform which slice the volume into multiple 2-D images. In 3-D FAT, a volume is reconstructed from measurements at different sides all around the volume. Thus, with FAT the problem of wave orientation with respect to instrument and flight path (Ungermann et al., 2010a) does not appear.

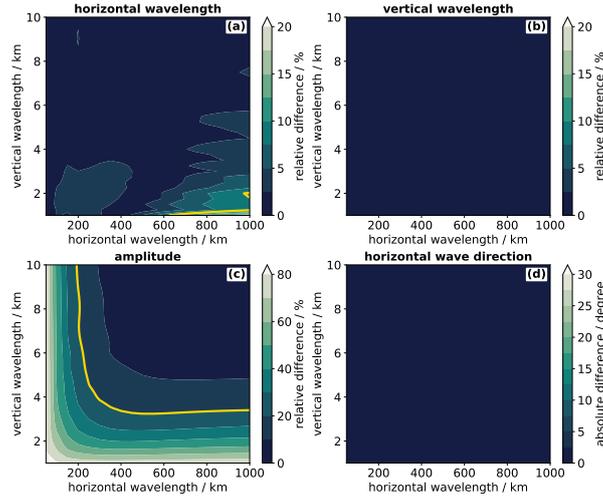
Figure 3 shows a 3-D FAT of a GW with horizontal wavelength of 400 km, vertical wavelength of 6 km and horizontal wave direction  $\varphi_r = 180^\circ$   $\varphi_s = 180^\circ$ . Pictured are three cross sections through the 3-D volume: at 10.5 km altitude (first row), at 20  $0^\circ$  N (second row), and at  $0^\circ$  E (third row). The first column shows the true synthetic wave, the second column the retrieved wave and the third column the difference of both. This synthetic wave has east-west oriented phase fronts (Fig. 3a) and is tilted to the south (Fig. 3c).

In general, the signal is well reproduced by the retrieval. Within the tangent point area (dotted lines) the temperature error is below 0.5 K. This is in good agreement with the determined accuracy of 0.5 K (Sec. 2.4). The retrieval can also reproduce some signal outside the area covered by tangent points. One reason is, the path of the LOS which goes through higher altitudes before and after the tangent point, thus collecting information there. Another reason are the horizontal and vertical correlation lengths of 100 km and 1 km, which are used for the retrieval and smear out the signal.

A S3D fit (Sec. 2.6) was performed for this retrieval at 10.5 km altitude with a cube size of 5 km x 400 km x 400 km. The results of this fit can be seen in Fig. 3j-m. Within the hexagonal flight pattern the horizontal and vertical wavelength, and the horizontal wave direction are well reproduced. The original amplitude of 3 K is underestimated by 0.1 K.



**Figure 3.** FAT retrieval and S3D results for a wave with 400 km horizontal, 6 km vertical wavelength, and horizontal wave direction  $\varphi_r = 180^\circ$ ,  $\varphi_s = 180^\circ$ . Panels (a-c) show the true synthetic wave, panels (d-e) the retrieved wave, and panels (g-i) the difference between both. In the first row (a, d, g) horizontal cross sections at 10.5 km altitude are depicted. The rows 2 (b, e, h) and 3 (c, f, i) show vertical cross sections along 0° N and 0° E, respectively. The black dashed lines mark the area covered by tangent points. In panels (j-m) the S3D results horizontal wavelength  $\lambda_h$ , vertical wavelength  $\lambda_z$ , horizontal wave direction  $\varphi$ , and temperature amplitude  $\hat{T}$  are pictured.



**Figure 4.** Specific observational filters for (a) horizontal wavelength  $\lambda_h$ , (b) vertical wavelength  $\lambda_z$ , (c) temperature amplitude  $\hat{T}$ , and (d) horizontal wave direction  $\varphi$  for the FAT retrieval. The yellow lines mark errors of 10% for the horizontal wavelength (Panel (a)) and 20% for the amplitude (Panel (c)).

These S3D results are used to construct the specific observational filters in Fig. 4. A mean value of the S3D fit results between  $1^\circ$  S,  $1^\circ$  N,  $1^\circ$  W, and  $1^\circ$  E gives the specific observational filter of the respective wavelength pair. The horizontal wavelength, the vertical wavelength and the horizontal wave direction are well reproduced for all tested waves. Further, there appears no phase shift in the FAT retrieval in contrast to conventional 1-D retrievals (Ungermaun et al., 2010a). However, the amplitude of the waves is slightly reduced for waves with horizontal wavelength below 200 km or vertical wavelength below 3 km.

This simulation study shows, that FAT is able to properly reconstruct the wave vectors of mesoscale GWs. However, the observational filter of the temperature amplitude has to be taken into account, when comparing these measurements to different data sets.

### 3.2 Limited angle tomography (LAT)

#### 10 3.2.1 Dependence of the retrieval results on horizontal and vertical wavelengths

Figure 5 shows a comparison of the LAT retrieval results for different wavelengths. The waves in column 1 and 4 have a larger horizontal wavelength of 600 km compared to the waves in columns 2 and 3 with 200 km horizontal wavelength. The vertical wavelength of 6 km of the waves in columns 1 and 3 is longer than the vertical wavelength of 2 km of the waves in columns 2 and 4. The waves with large vertical wavelength in columns 1 and 3 are well reproduced by the LAT retrieval within the tangent point covered area with errors below 0.5 K. The waves with short vertical wavelengths show larger temperature errors of up to 1.5 K within the tangent point area. This difference comes from the curved LOS through the straight wavefronts, which leads to an averaging over different wave phases. For the waves with short vertical wavelength the LOS crosses multiple opposite

wave phases, which decreases the measurement signal. A similar dependence of the sensitivity on the alignment of phasefronts with LOS was observed for sub-limb viewers (Wu and Waters, 1996; McLandress et al., 2000).

All retrieved waves show a slight V shape pattern, which is more emphasized for the waves with short vertical wavelength. This V shape is probably caused by the parabola shape of the LOS. The retrieval does not know, where along the line-of-sight  
5 how much of the measured radiation was emitted, unless crossing measurements give sufficient information. As the LAT has fewer measurements at different angles, the temperature signal is redistributed according to the weighting function (Fig. 2b) along the LOS. This can be nicely seen in the vertical cross sections in Fig. 5(g, o), where the warm temperature follows the LOS upwards behind the tangent point. This vertical shift of temperature also causes the northward oriented V shape in the horizontal cross sections.

10 As already for the FAT case, the specific observational filters were calculated using the S3D fits of the LAT retrievals (Fig. 6). The deviations of horizontal and vertical wavelengths are mainly below 10%. Only for very short vertical and very long horizontal wavelengths errors of above 20% appear. This is probably due to the above mentioned V shape deformation of the wave, which is more difficult to fit with one single sinusoidal wave. The same problem appears for the horizontal wave direction. For waves with short vertical and long horizontal wavelengths and, thus a strong V shape, the direction cannot be  
15 derived properly anymore. For the rest of the waves the direction error stays everywhere below 10°. The observational filter for the amplitude shows a similar pattern as for the FAT case.

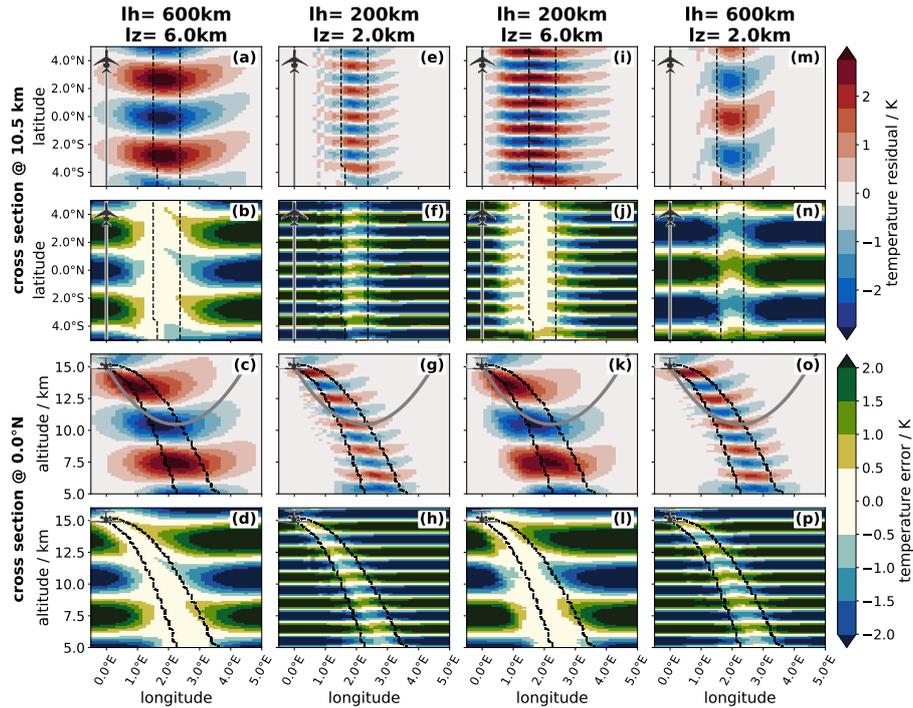
### 3.2.2 Dependence of retrieval results on the wave orientation

Due to the limited measurement sector, the orientation of the wave with respect to the instrument position might be important for LAT. Figure 7 depicts the retrieval results for waves with horizontal wave directions turned by 30° ( $\varphi_r = 210^\circ$ ,  $\varphi_s = 210^\circ$ )  
20 compared to those in Fig. 5. ~~In the vertical, the~~ The wavefronts are tilted southward and westward and, thus, the vertical tilt is towards the instrument. They decrease in height with increasing distance to the flight path.

Overall, the structures are reproduced reasonably well. As for the perfectly perpendicularly-aligned waves already, waves with long vertical wavelengths (Fig. 7a–d and Fig. 7i–l) are reproduced better than waves with short vertical wavelengths (Fig. 7e–h and Fig. 7m–p).

25 Due to the tilt of the waves towards the aircraft, the LOS is partly aligned with the wavefronts before the tangent point. This effect is stronger for steep waves such as in Fig. 7k than for relatively flat waves such as in Fig. 8c, g and o. Due to this alignment the area of best sensitivity is shifted towards the aircraft for the steep wave. Spreading the signal now around this shifted sensitivity maximum, just spreads the signal along the same wave phase, as the LOS has little curvature in this region. Therefore, no strong shape deviation is observed. For the flat waves a similar V shape can be observed as for the waves in  
30 Fig. 5, due to a spreading of signal along LOS around the tangent point.

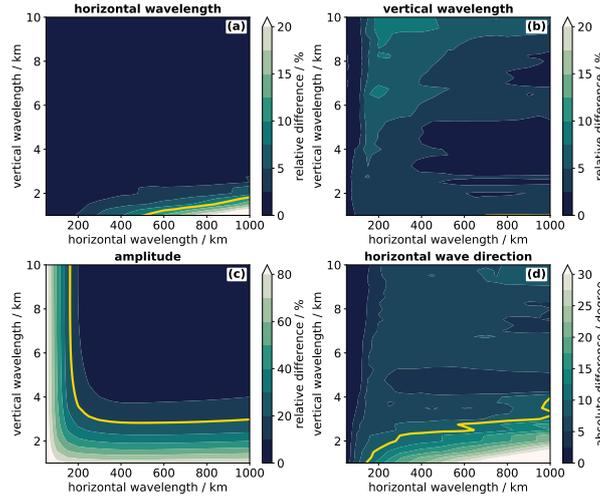
In the observational filter (Fig. 8) a small decrease in the quality of amplitude reproduction can be seen compared to the observational filter of perfectly east-west aligned waves (Fig. 6). However, the wavelengths and wave direction are barely influenced and reproduced at a similar high quality. The V shape of the waves only occurs outside the tangent point region, thus proper horizontal wave directions can be observed.



**Figure 5.** Cross sections of retrieved waves (first and third row) and differences between true-synthetic and retrieved waves (second and fourth row) of the LAT retrieval. The different columns show waves with different horizontal ( $lh$ ) and vertical ( $lz$ ) wavelengths. The true-synthetic horizontal wave orientation of all waves is  $\varphi_T = 180^\circ - \varphi_S = 180^\circ$  and, thus, these waves have wavefronts perpendicular to the flight path. The black dashed lines mark the area covered by tangent points. The grey line in the vertical cross sections indicates a LOS for a measurement with  $90^\circ$  azimuth angle and tangent point altitude of 10.5 km.

Fig. 9 shows the retrieval results for waves turned by  $-210^\circ$  ( $\varphi_T = 30^\circ - \varphi_S = 30^\circ$ ) compared to Fig. 5. These waves are tilted northward and eastward and, thus, the vertical tilt is away from the flight path. Only for the wave with large horizontal and large vertical wavelength (Fig. 9a–d) the temperature amplitude is reproduced well within the tangent point region. However, the horizontal orientation in this area, which should be similar to Fig. 7a from north-west to south-east is not recovered. Within the tangent point region (longitude between  $1.5^\circ$  and  $2.5^\circ$ ) the horizontal wave fronts are oriented almost west to east. Behind the tangent point region (longitude above  $2.5^\circ$ ), the wave fronts slowly approach the expected west-north-west to east-south-east orientation. Thus, the orientation error ranges between  $10^\circ$  and  $30^\circ$  depending on the area of interest. The same happens for waves with short vertical wavelengths (Fig. 9e–h and m–p): The information about the horizontal wave direction is lost within the retrieval. Again a V shape appears for all these waves. Due to the inverse vertical tilt compared to Fig. 7, the opening of the V shape is this time to the south.

For steep waves (Fig. 9k) the main signal is again shifted, this time behind the tangent point area, where the LOS and the wavefronts are well aligned. Thus the spreading of the signal does not influence these waves as strongly as the flat waves and



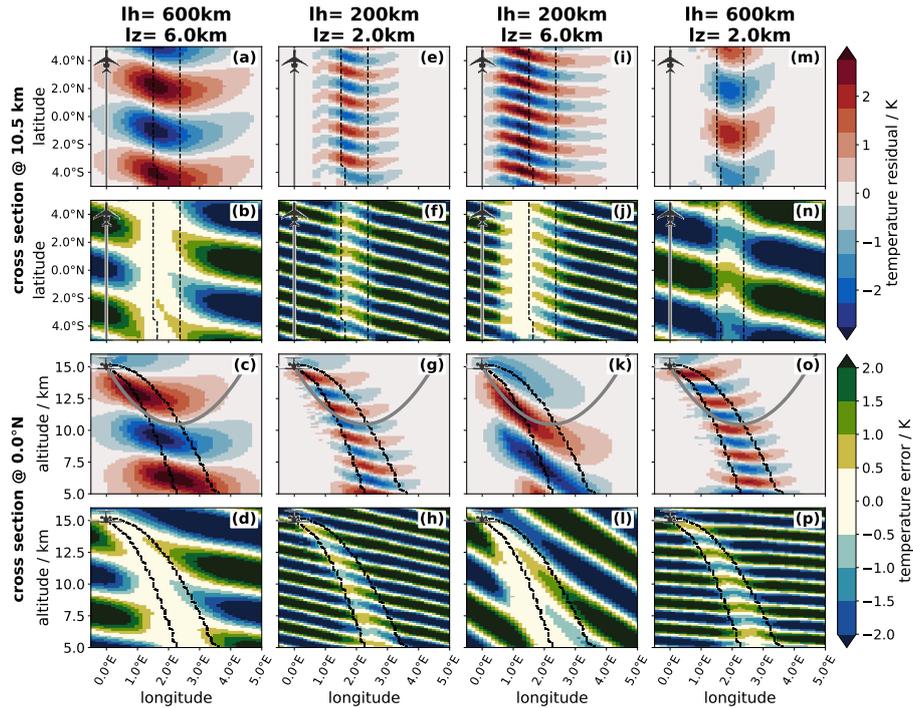
**Figure 6.** Specific observational filters for (a) horizontal wavelength  $\lambda_h$ , (b) vertical wavelength  $\lambda_z$ , (c) temperature amplitude  $\hat{T}$ , and (d) horizontal wave direction  $\varphi$  for the LAT retrieval and true-synthetic horizontal wave orientation  $\varphi_t = 180^\circ$   $\varphi_s = 180^\circ$  and, thus, these waves have wavefronts perpendicular to the flight path. The yellow lines mark errors of 10%, 10%, 20% and  $10^\circ$  in Panels (a-d), respectively.

the horizontal orientation does not get lost in the retrieval. The decreased amplitude compared to Fig. 7 can be explained by the fact that the maximum of the weighting function along LOS is located slightly before the tangent point (Fig. 2b).

A similar picture is given from the observational filter in Fig. 10. Even though the amplitude is underestimated for very steep waves, the horizontal wave orientation can be derived accurately. However, the flatter the wave gets, the worse the derived horizontal wave direction. For waves with horizontal to vertical wavelengths ratio of above 200, the direction error exceeds  $30^\circ$ . Also the horizontal wavelengths reproduction is decreased somewhat compared to the two cases before (Fig. 6 and Fig. 8).

Further tests with horizontal wave direction  $30^\circ < \varphi_t < 90^\circ$  and  $210^\circ < \varphi_t < 270^\circ$   $30^\circ < \varphi_s < 90^\circ$  and  $210^\circ < \varphi_s < 270^\circ$  show a drastic decline in the amplitude sensitivity towards waves with short horizontal wavelengths. For waves tilted away from the flight path ( $\varphi_t > 30^\circ$   $\varphi_s > 30^\circ$ ) the fit quality of the horizontal wave direction and the horizontal wavelength decreases drastically already at  $\varphi_t = 40^\circ$   $\varphi_s = 40^\circ$ .

These studies show that LAT applied to gravity waves gives best results for waves with wavefronts perpendicular to the flight path and, thus, horizontal wave vector  $\varphi_t = 180^\circ$   $\varphi_s = 180^\circ$  or  $\varphi_s = 0^\circ$ . However, if the wave is slightly turned, the quality of the derived wave parameters is not affected strongly as long as the wave is tilted towards the instrument ( $180^\circ \leq \varphi_t \leq 210^\circ$   $180^\circ \leq \varphi_s \leq 210^\circ$ ). In general, waves are best retrieved when their aspect ratio of horizontal and vertical wavelengths, i.e. their steepness, is favourable for an alignment with the LOS. In these cases, tilts towards and away from the instrument may give reasonable results.

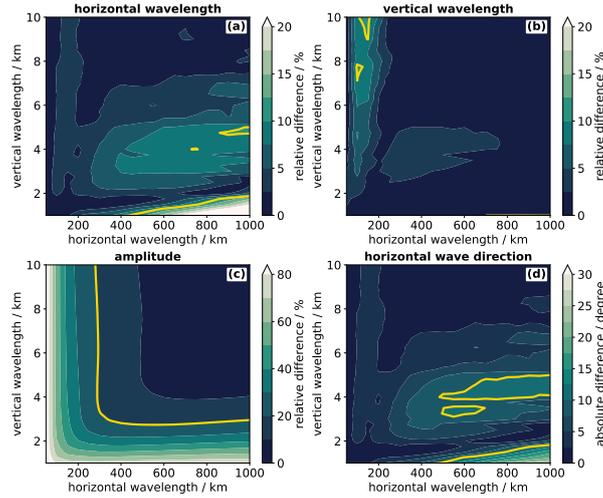


**Figure 7.** Cross sections of retrieved waves (first and third row) and differences between true-synthetic and retrieved waves (second and fourth row) of the LAT retrieval. The different columns show waves with different horizontal (lh) and vertical (lz) wavelengths. The true-synthetic horizontal wave orientation of all waves is  $\varphi_r = -210^\circ$   $\varphi_s = 210^\circ$  and, thus, these waves are tilted towards the aircraft.

### 3.3 Comparison of LAT and FAT results for a real measurement case on 25 January 2016 over Iceland

From December 2015 to March 2016, GLORIA was deployed on board of the German research aircraft HALO for a research campaign covering several scientific targets such as demonstrating the use of infrared limb imaging for gravity wave studies (GWEX), studying the full life cycle of a gravity wave (GW-LCYCLE), investigating the Seasonality of Air mass transport and origin in the Lowermost Stratosphere (SALSA), and observing the Polar Stratosphere in a Changing Climate (POLSTRACC). On 25 January 2016, a research flight over Iceland investigated a GW excited at the Icelandic Mountains (Krisch et al., 2017). A linear flight leg of 500 km length crossing the wavefronts, was followed by a hexagonal flight pattern with 460 km diameter around the wave structure.

Krisch et al. (2017) present the retrieval results of FAT using only measurements taken during the hexagonal flight. Fig. 11 compares these FAT results with LAT results using only measurements taken on the 500 km linear flight leg through the middle of the volume. In general the LAT results (Fig. 11 d-e) agree very well with the FAT results (Fig. 11 a-c) within the volume covered by both. FAT as well as the LAT retrieval, show a superposition of waves with longer and shorter horizontal wavelengths. Differences in strength and scale of the waves, for example in cross section #2, can be explained due to the

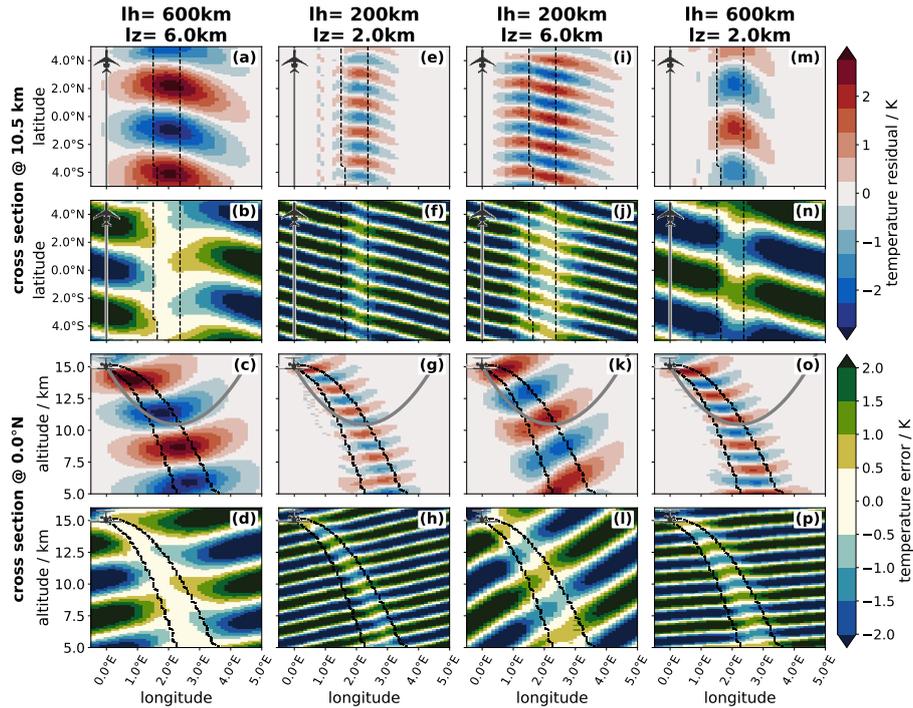


**Figure 8.** Specific observational filters for (a) horizontal wavelength  $\lambda_h$ , (b) vertical wavelength  $\lambda_z$ , (c) temperature amplitude  $\hat{T}$ , and (d) horizontal wave direction  $\varphi$  for the LAT retrieval with true-synthetic horizontal wave orientation  $\varphi_r = 210^\circ$   $\varphi_s = 210^\circ$  and, thus, waves, which are tilted towards the aircraft. The yellow lines mark errors of 10%, 10%, 20% and 10° in Panels (a-d), respectively.

different tangent point coverage of both methods. Especially higher altitudes in cross section #2 are not well covered with tangent points in the FAT retrieval (Fig. 11c). This is probably the reason why the temperature residual is slightly lower in the FAT retrieval compared to LAT. Also the smaller scale waves in this region (Fig. 11f) are less prominent in the FAT retrieval (Fig. 11c).

- 5 A more quantitative comparison of the similarities of both retrievals can be given by the Pearson correlation coefficient. Including only areas which are covered by tangent points gives a Pearson correlation coefficient of 0.91. Expanding this area to places with measurement content larger than 0.8 (includes areas crossed by a LOS before or after the tangent point) still leads to a Pearson correlation coefficient of 0.75. Thus, as expected the two retrievals are highly correlated.

A main advantage of 3-D tomographic measurement of GWs over conventional limb measurements is the ability to derive  
 10 the horizontal wave direction. This is done by applying the S3D fitting routine. Figure 12 shows the wave parameters obtained from these fits for both cases. Within the confidence area of our fits, all wave parameters agree very well for both methods. The observed GW has a horizontal wavelength of around 200 km, a vertical wavelength around 5.5 km, amplitudes up to 2 K and a horizontal wave direction of 160°. Thus, the wave vector is turned by 20° compared to the flight direction. The wavefronts are tilted southward and, thus, away from the flight path. Figure 10 predicted for such waves a wavelength reproduction of  
 15 more than 90% and an error in the estimation of the horizontal wave direction below 7.5°. This can be confirmed with the real measurement case.

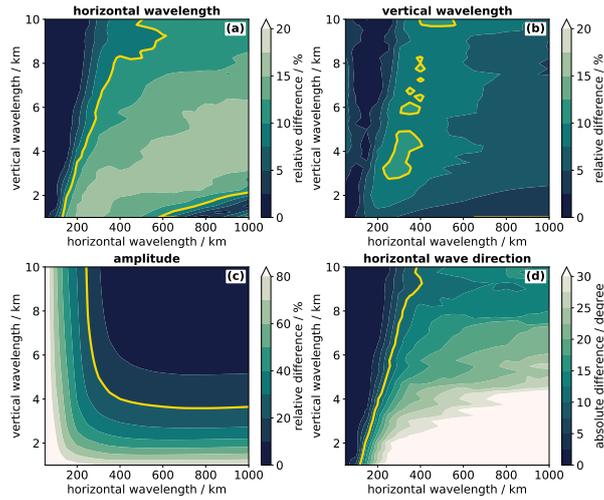


**Figure 9.** Cross sections of retrieved waves (first and third row) and differences between true-synthetic and retrieved waves (second and fourth row) of the LAT retrieval. The different columns show waves with different horizontal ( $l_h$ ) and vertical ( $l_z$ ) wavelengths. The true-synthetic horizontal wave orientation of all waves is  $\varphi_r = 30^\circ$   $\varphi_s = 30^\circ$  and, thus, these waves are tilted away from aircraft.

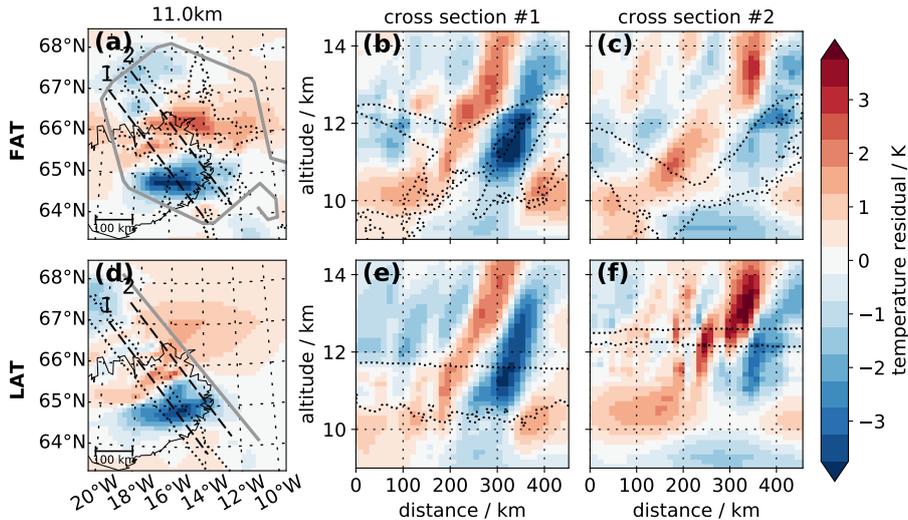
#### 4 Conclusions

This paper investigates the use of LAT applied to airborne limb imaging for gravity wave research. In contrast to FAT, which allows for the reconstruction of a large, eubicylindric, 3-D volume, LAT can only reconstruct a band of 200 km around a banana-shaped vertical curtain parallel to the flight path. The horizontal resolution is 30 km in flight direction and 70 km perpendicular to flight direction. The vertical resolution is on the order of 400 m. This volume and resolution are sufficient to properly derive all important wave parameters such as the horizontal and vertical wavelengths, the amplitude, and the wave direction for waves with wavefronts perpendicular to the flight path. This is feasible due to the perfect alignment of wave phases and LOS and agrees well with earlier studies for other limb sounding concepts (Wu and Waters, 1996; McLandress et al., 2000; Ungermann et al., 2010a).

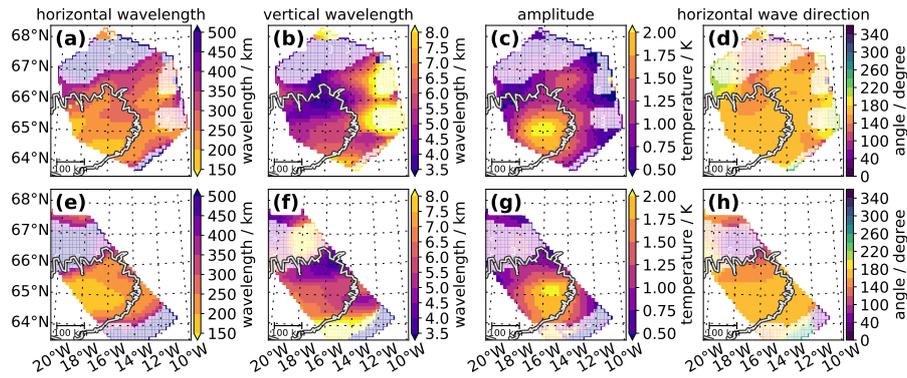
10 The quality of the 3-D reconstruction strongly depends on the orientation of the wave with respect to the instrument. If the waves are slightly turned away from the perfect orientation the quality of the derived wave parameters is not strongly affected as long as the wavefront is tilted towards the instrument. If the wavefronts are tilted away from the instrument, the retrieval will create artefacts in form of v-shaped phase fronts, which reduce the quality of the derived horizontal wave directions and



**Figure 10.** Specific observational filters for (a) horizontal wavelength  $\lambda_h$ , (b) vertical wavelength  $\lambda_z$ , (c) temperature amplitude  $\hat{T}$ , (d) and horizontal wave direction  $\varphi$  for the LAT retrieval with true synthetic horizontal wave orientation  $\varphi_r = 30^\circ$   $\varphi_s = 30^\circ$  and, thus, waves, which are tilted away from aircraft. The yellow lines mark errors of 10%, 10%, 20% and  $10^\circ$  in Panels (a-d), respectively.



**Figure 11.** Comparison of FAT (a-c) and LAT (d-f) retrieval results for a research flight over Iceland on 25 January 2016. Shown are the temperature residuals after background removal as described in Sec. 2.5. The grey line indicates the part of the flight path, from which the measurements are used for the retrieval. The left column depicts horizontal crosssections at 11 km altitude, the two right columns present vertical cross sections along the dashed lines of the left column. The dotted lines mark the area covered by tangent points. Panels (a) to (c) are adapted from Krisch et al. (2017) Fig. 3.



**Figure 12.** Wave parameters as obtained from the S3D fit with fitting cubes of 160 km x 160 km x 3.6 km at center height of 11.5 km for FAT (a-d) and LAT (e-h). Nonsignificant fitting results with wavelengths above 2.5 times the cube size are shaded. Panel (a) to (d) are adapted from Krisch et al. (2017) Fig. 5.

wavelengths. For waves with horizontal wavelength under 300 km, the amplitude error is larger for waves with wavefronts tilted away from the instrument than for waves with wavefronts tilted towards the instrument. In general, the better the alignment of the wave phases and the LOS is, the more information is attained by the tomographic retrieval. Thus, steeper waves can be derived with better accuracy than flatter waves. For steep waves with a horizontal to vertical wavelength ratio below 200 correct wave directions can be derived independently of the tilt. However, for waves turned by more than  $40^\circ$  compared to the perfect, perpendicular case, the reconstruction quality decreases drastically for all tested waves.

The capacity of LAT for GW research was demonstrated by comparing LAT and FAT for a real measurement case on 25 January 2016 above Iceland. The temperature residuals agree very well with each other. The wave parameters derived with a sinusoidal fitting routine yield similar results.

In summary, for many GW cases the observation in LAT mode can be recommended. However, for short scale waves FAT is preferable due to the higher spatial resolution of 20 km x 20 km x 200 m. The slightly better accuracy of 0.5 K for FAT compared to 0.7 K for LAT also makes FAT favourable for low amplitude waves. Furthermore, when the precise orientation of the wave cannot be predicted before the flight, FAT should be the method of choice. Nevertheless, for many other cases, LAT might be preferred due to its shorter acquisition time.

*Data availability.* The tomographic retrieval data used in Sec. 3.3 is available from the HALO database (Krisch, 2017, 2018).

*Competing interests.* The authors declare that they have no conflict of interest.

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## References

- Alexander, M. J.: Interpretations of observed climatological patterns in stratospheric gravity wave variance, *Geophys. Res. Lett.*, 103, 8627–8640, <https://doi.org/10.1029/97JD03325>, 1998.
- Degenstein, D. A., Llewellyn, E. J., and Lloyd, N. D.: Tomographic retrieval of the oxygen infrared atmospheric band with the OSIRIS  
5 infrared imager, *Can. J. Phys.*, 82, 501–515, 2004.
- Eckermann, S. D., Wu, D. L., Doyle, J. D., Burris, J. F., McGee, T. J., Hostetler, C. A., Coy, L., Lawrence, B. N., Stephens, A., McCormack, J. P., and Hogan, T. F.: Imaging gravity waves in lower stratospheric AMSU-A radiances, Part 2: validation case study, *Atmos. Chem. Phys.*, 6, 3343–3362, <https://doi.org/10.5194/acp-6-3343-2006>, 2006.
- Ern, M. and Preusse, P.: Gravity wave momentum flux spectra observed from satellite in the summertime subtropics: Implications for global  
10 modeling, *Geophys. Res. Lett.*, 39, <https://doi.org/10.1029/2012GL052659>, 2012.
- Ern, M., Preusse, P., and Warner, C. D.: Some experimental constraints for spectral parameters used in the Warner and McIntyre gravity wave parameterization scheme, *Atmos. Chem. Phys.*, 6, 4361–4381, <https://doi.org/10.5194/acp-6-4361-2006>, 2006.
- Friedl-Vallon, F., Gulde, T., Hase, F., Kleinert, A., Kulesa, T., Maucher, G., Neubert, T., Olschewski, F., Piesch, C., Preusse, P., Rongen, H., Sartorius, C., Schneider, H., Schönfeld, A., Tan, V., Bayer, N., Blank, J., Dapp, R., Ebersoldt, A., Fischer, H., Graf, F., Guggenmoser, T.,  
15 Höpfner, M., Kaufmann, M., Kretschmer, E., Latzko, T., Nordmeyer, H., Oelhaf, H., Orphal, J., Riese, M., Schardt, G., Schillings, J., Sha, M. K., Suminska-Ebersoldt, O., and Ungermann, J.: Instrument concept of the imaging Fourier transform spectrometer GLORIA, *Atmos. Meas. Tech.*, 7, 3565–3577, <https://doi.org/10.5194/amt-7-3565-2014>, 2014.
- Geller, M. A., Alexander, M. J., Love, P. T., Bacmeister, J., Ern, M., Hertzog, A., Manzini, E., Preusse, P., Scaife, A. A., and Zhou, T.: A comparison between gravity wave momentum fluxes in observations and climate models, *J. Clim.*, 26, 6383–6405,  
20 <https://doi.org/10.1175/JCLI-D-12-00545.1>, 2013.
- Gordley, L. L. and Russell, J. M.: Rapid inversion of limb radiance data using an emissivity growth approximation, *Appl. Optics*, 20, 807–813, <https://doi.org/10.1364/AO.20.000807>, 1981.
- Hart, V. P., Taylor, M. J., Doyle, T. E., Zhao, Y., Pautet, P.-D., Carruth, B. L., Rusch, D. W., and Russell, J. M.: Investigating Gravity Waves in Polar Mesospheric Clouds Using Tomographic Reconstructions of AIM Satellite Imagery, *Journal of Geophysical Research: Space  
25 Physics*, 123, 955–973, <https://doi.org/10.1002/2017JA024481>, <https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1002/2017JA024481>, 2018.
- Hoffmann, L.: Schnelle Spurengasretrieval für das Satellitenexperiment Envisat MIPAS, Tech. Rep. JUEL-4207, Forschungszentrum Jülich, Jülich, Germany, ISSN 0944-2952, 2006.
- Hoffmann, L. and Alexander, M. J.: Retrieval of stratospheric temperatures from Atmospheric Infrared Sounder radiance measurements for  
30 gravity wave studies, *J. Geophys. Res.*, 114, D07 105, <https://doi.org/10.1029/2008JD011241>, 2009.
- Hoffmann, L., Kaufmann, M., Spang, R., Müller, R., Remedios, J. J., Moore, D. P., Volk, C. M., von Clarmann, T., and Riese, M.: Envisat MIPAS measurements of CFC-11: retrieval, validation, and climatology, *Atmos. Chem. Phys.*, 8, 3671–3688, <https://doi.org/10.1016/j.asr.2005.03.112>, 2008.
- Hultgren, K., Gumbel, J., Degenstein, D., Bourassa, A., Lloyd, N., and Stegman, J.: First simultaneous retrievals of horizontal  
35 and vertical structures of Polar Mesospheric Clouds from Odin/OSIRIS tomography, *J. Atm. Sol.-Terr. Phys.*, 104, 213 – 223, <https://doi.org/https://doi.org/10.1016/j.jastp.2013.06.013>, <http://www.sciencedirect.com/science/article/pii/S1364682613001892>, 2013.

- Kalicinsky, C., Grooss, J. U., Guenther, G., Ungermann, J., Blank, J., Hofer, S., Hoffmann, L., Knieling, P., Olschewski, F., Spang, R., Stroh, F., and Riese, M.: Observations of filamentary structures near the vortex edge in the Arctic winter lower stratosphere, *Atmos. Chem. Phys.*, 13, 10 859–10 871, <https://doi.org/10.5194/acp-13-10859-2013>, 2013.
- Kalisch, S., Preusse, P., Ern, M., Eckermann, S. D., and Riese, M.: Differences in gravity wave drag between realistic oblique and assumed vertical propagation, *J. Geophys. Res. Atmos.*, 119, 10,081–10,099, <https://doi.org/10.1002/2014JD021779>, 2014.
- Kaufmann, M., Blank, J., Guggenmoser, T., Ungermann, J., Engel, A., Ern, M., Friedl-Vallon, F., Gerber, D., Grooß, J. U., Guenther, G., Höpfner, M., Kleinert, A., Kretschmer, E., Latzko, T., Maucher, G., Neubert, T., Nordmeyer, H., Oelhaf, H., Olschewski, F., Orphal, J., Preusse, P., Schlager, H., Schneider, H., Schuettemeyer, D., Stroh, F., Suminska-Ebersoldt, O., Vogel, B., M. Volk, C., Woiwode, W., and Riese, M.: Retrieval of three-dimensional small-scale structures in upper-tropospheric/lower-stratospheric composition as measured by GLORIA, *Atmos. Meas. Tech.*, 8, 81–95, <https://doi.org/10.5194/amt-8-81-2015>, <http://www.atmos-meas-tech.net/8/81/2015/>, 2015.
- Kidston, J., Scaife, A. A., Hardiman, S. C., Mitchell, D. M., Butchart, N., Baldwin, M. P., and Gray, L. J.: Stratospheric influence on tropospheric jet streams, storm tracks and surface weather, *Nature Geosci.*, 8, 433–440, <https://doi.org/10.1038/ngeo2424>, 2015.
- Kleinert, A., Friedl-Vallon, F., Guggenmoser, T., Höpfner, M., Neubert, T., Ribalda, R., Sha, M. K., Ungermann, J., Blank, J., Ebersoldt, A., Kretschmer, E., Latzko, T., Oelhaf, H., Olschewski, F., and Preusse, P.: Level 0 to 1 processing of the imaging Fourier transform spectrometer GLORIA: generation of radiometrically and spectrally calibrated spectra, *Atmos. Meas. Tech.*, 7, 4167–4184, <https://doi.org/10.5194/amt-7-4167-2014>, 2014.
- Kleinert, A., Krisch, I., Ungermann, J., Adibekyan, A., Gutschwager, B., and Monte, C.: Characterization of blackbody inhomogeneity and its effect on the retrieval results of the GLORIA instrument, *Atmos. Meas. Tech.*, p. submitted, 2018.
- Krasauskas, L., Ungermann, J., Ensmann, S., Krisch, I., Kretschmer, E., Preuße, P., and Riese, M.: Tomographic airborne limb sounder retrievals on irregular grid with second order regularisation, *Atmos. Meas. Tech. Discuss.*, to be submitted, 2018.
- Krisch, I.: F10\_20160125\_GLORIA\_dynamics\_mode\_tomography\_FZJ.nc, <https://halo-db.pa.op.dlr.de/dataset/5079>, last access: 10 July, 2017.
- Krisch, I.: F10\_20160125\_GLORIA\_dynamics\_mode\_linear\_flight\_tomography\_FZJ.nc, <https://halo-db.pa.op.dlr.de/dataset/5445>, last access: 27 February, 2018.
- Krisch, I., Preusse, P., Ungermann, J., Dörnbrack, A., Eckermann, S. D., Ern, M., Friedl-Vallon, F., Kaufmann, M., Oelhaf, H., Rapp, M., Strube, C., and Riese, M.: First tomographic observations of gravity waves by the infrared limb imager GLORIA, *Atmos. Chem. Phys.*, 17, 14 937–14 953, <https://doi.org/10.5194/acp-17-14937-2017>, 2017.
- Lehmann, C. I., Kim, Y.-H., Preusse, P., Chun, H.-Y., Ern, M., and Kim, S.-Y.: Consistency between Fourier transform and small-volume few-wave decomposition for spectral and spatial variability of gravity waves above a typhoon, *Atmos. Meas. Tech.*, 5, 1637–1651, <https://doi.org/10.5194/amt-5-1637-2012>, 2012.
- McLandress, C., Alexander, M. J., and Wu, D. L.: Microwave Limb Sounder observations of gravity waves in the stratosphere: A climatology and interpretation, *J. Geophys. Res.*, 105, 11 947–11 967, 2000.
- McLandress, C., Shepherd, T. G., Polavarapu, S., and Beagley, S. R.: Is Missing Orographic Gravity Wave Drag near 60 degrees S the Cause of the Stratospheric Zonal Wind Biases in Chemistry Climate Models?, *J. Atmos. Sci.*, 69, 802–818, <https://doi.org/10.1175/JAS-D-11-0159.1>, 2012.
- Natterer, F.: *The mathematics of computerized tomography*, Society for Industrial and Applied Mathematics, Philadelphia, PA, USA, 2001.

- Preusse, P., Dörnbrack, A., Eckermann, S. D., Riese, M., Schaeler, B., Bacmeister, J. T., Broutman, D., and Grossmann, K. U.: Space-based measurements of stratospheric mountain waves by CRISTA, 1. Sensitivity, analysis method, and a case study, *J. Geophys. Res.*, 107(D23), <https://doi.org/10.1029/2001JD000699>, 2002.
- Preusse, P., Eckermann, S. D., Ern, M., Oberheide, J., Picard, R. H., Roble, R. G., Riese, M., Russell III, J. M., and Mlynczak, M. G.: Global ray tracing simulations of the SABER gravity wave climatology, *J. Geophys. Res. Atmos.*, 114, <https://doi.org/10.1029/2008JD011214>, 2009.
- Remedios, J. J., Leigh, R. J., Waterfall, A. M., Moore, D. P., Sembhi, H., Parkes, I., Greenhough, J., Chipperfield, M., and Hauglustaine, D.: MIPAS reference atmospheres and comparisons to V4.61/V4.62 MIPAS level 2 geophysical data sets, *Atmos. Chem. Phys. Discuss.*, 7, 9973–10017, <https://doi.org/10.5194/acpd-7-9973-2007>, 2007.
- 10 Ribstein, B. and Achatz, U.: The interaction between gravity waves and solar tides in a linear tidal model with a 4-D ray-tracing gravity-wave parameterization, *Journal of Geophysical Research: Space Physics*, 121, 8936–8950, <https://doi.org/10.1002/2016JA022478>, <http://dx.doi.org/10.1002/2016JA022478>, 2016JA022478, 2016.
- Riese, M., Preusse, P., Spang, R., Ern, M., Jarisch, M., Grossmann, K.-U., and Offermann, D.: Measurements of trace gases by the Cryogenic Infrared Spectrometers and Telescopes for the Atmosphere (CRISTA) experiment, *Adv. Space Res.*, 19, 563–566, [https://doi.org/10.1016/S0273-1177\(97\)00172-5](https://doi.org/10.1016/S0273-1177(97)00172-5), 1997.
- 15 Riese, M., Oelhaf, H., Preusse, P., Blank, J., Ern, M., Friedl-Vallon, F., Fischer, H., Guggenmoser, T., Hoepfner, M., Hoor, P., Kaufmann, M., Orphal, J., Ploeger, F., Spang, R., Suminska-Ebersoldt, O., Ungermann, J., Vogel, B., and Woiwode, W.: Gimballed Limb Observer for Radiance Imaging of the Atmosphere (GLORIA) scientific objectives, *Atmos. Meas. Tech.*, 7, 1915–1928, <https://doi.org/10.5194/amt-7-1915-2014>, 2014.
- 20 Rodgers, C. D.: *Inverse Methods for Atmospheric Sounding: Theory and Practice*, vol. 2 of *Series on Atmospheric, Oceanic and Planetary Physics*, World Scientific, Singapore, 2000.
- Sandu, I., Bechtold, P., Beljaars, A., Bozzo, A., Pithan, F., Shepherd, T. G., and Zadra, A.: Impacts of parameterized orographic drag on the Northern Hemisphere winter circulation, *J. Adv. Model. Earth Syst.*, 8, 196–211, <https://doi.org/10.1002/2015MS000564>, 2016.
- Sato, K., Watanabe, S., Kawatani, Y., Tomikawa, Y., Miyazaki, K., and Takahashi, M.: On the origins of mesospheric gravity waves, *Geophys. Res. Lett.*, 36, <https://doi.org/10.1029/2009GL039908>, 2009.
- 25 Savitzky, A. and Golay, M. J. E.: Smoothing and Differentiation of Data by Simplified Least Squares Procedures., *Analytical Chemistry*, 36, 1627–1639, <https://doi.org/10.1021/ac60214a047>, <http://dx.doi.org/10.1021/ac60214a047>, 1964.
- Scaife, A. A., Karpechko, A. Y., Baldwin, M. P., Brookshaw, A., Butler, A. H., Eade, R., Gordon, M., MacLachlan, C., Martin, N., Dunstone, N., and Smith, D.: Seasonal winter forecasts and the stratosphere, *Atmospheric Science Letters*, 17, 51–56, <https://doi.org/10.1002/asl.598>, <http://dx.doi.org/10.1002/asl.598>, 2016.
- 30 Sigmund, M. and Scinocca, J. F.: The Influence of the Basic State on the Northern Hemisphere Circulation Response to Climate Change, *Journal of Climate*, 23, 1434–1446, <https://doi.org/10.1175/2009JCLI3167.1>, 2010.
- Song, R., Kaufmann, M., Ungermann, J., Ern, M., Liu, G., and Riese, M.: Tomographic reconstruction of atmospheric gravity wave parameters from airglow observations, *Atmospheric Measurement Techniques*, 10, 4601–4612, <https://doi.org/10.5194/amt-10-4601-2017>, <https://www.atmos-meas-tech.net/10/4601/2017/>, 2017.
- 35 Tarantola, A.: *Inverse Problem Theory*, Society for Industrial and Applied Mathematics, Philadelphia, 2004.
- Tikhonov, A. N. and Arsenin, V. Y.: *Solutions of ill-posed problems*, Winston, Washington D.C., USA, 1977.

- Trinh, Q. T., Kalisch, S., Preusse, P., Ern, M., Chun, H.-Y., Eckermann, S. D., Kang, M.-J., and Riese, M.: Tuning of a convective gravity wave source scheme based on HIRDLS observations, *Atmos. Chem. Phys.*, 16, 7335–7356, <https://doi.org/10.5194/acp-16-7335-2016>, 2016.
- Ungermann, J., Hoffmann, L., Preusse, P., Kaufmann, M., and Riese, M.: Tomographic retrieval approach for mesoscale gravity wave observations by the PREMIER Infrared Limb-Sounder, *Atmos. Meas. Tech.*, 3, 339–354, <https://doi.org/10.5194/amt-3-339-2010>, 2010a.
- Ungermann, J., Kaufmann, M., Hoffmann, L., Preusse, P., Oelhaf, H., Friedl-Vallon, F., and Riese, M.: Towards a 3-D tomographic retrieval for the air-borne limb-imager GLORIA, *Atmos. Meas. Tech.*, 3, 1647–1665, <https://doi.org/10.5194/amt-3-1647-2010>, 2010b.
- Ungermann, J., Blank, J., Lotz, J., Leppkes, K., Hoffmann, L., Guggenmoser, T., Kaufmann, M., Preusse, P., Naumann, U., and Riese, M.: A 3-D tomographic retrieval approach with advection compensation for the air-borne limb-imager GLORIA, *Atmos. Meas. Tech.*, 4, 2509–2529, <https://doi.org/10.5194/amt-4-2509-2011>, <http://www.atmos-meas-tech.net/4/2509/2011/>, 2011.
- Ungermann, J., Kalicinsky, C., Olschewski, F., Knieling, P., Hoffmann, L., Blank, J., Woiwode, W., Oelhaf, H., Hösen, E., Volk, C. M., Ulanovsky, A., Ravegnani, F., Weigel, K., Stroh, F., and Riese, M.: CRISTA-NF measurements with unprecedented vertical resolution during the RECONCILE aircraft campaign, *Atmos. Meas. Tech.*, 5, 1173–1191, <https://doi.org/10.5194/amt-5-1173-2012>, 2012.
- Ungermann, J., Blank, J., Dick, M., Ebersoldt, A., Friedl-Vallon, F., Giez, A., Guggenmoser, T., Höpfner, M., Jurkat, T., Kaufmann, M., Kaufmann, S., Kleinert, A., Krämer, M., Latzko, T., Oelhaf, H., Olschewski, F., Preusse, P., Rolf, C., Schillings, J., Suminska-Ebersoldt, O., Tan, V., Thomas, N., Voigt, C., Zahn, A., Zöger, M., and Riese, M.: Level 2 processing for the imaging Fourier transform spectrometer GLORIA: derivation and validation of temperature and trace gas volume mixing ratios from calibrated dynamics mode spectra, *Atmos. Meas. Tech.*, 8, 2473–2489, <https://doi.org/10.5194/amt-8-2473-2015>, 2015.
- Ungermann, J., Ern, M., Kaufmann, M., Müller, R., Spang, R., Ploeger, F., Vogel, B., and Riese, M.: Observations of PAN and its confinement in the Asian summer monsoon anticyclone in high spatial resolution, *Atmos. Chem. Phys.*, 16, 8389–8403, <https://doi.org/10.5194/acp-16-8389-2016>, 2016.
- von Clarmann, T., De Clercq, C., Ridolfi, M., Höpfner, M., and Lambert, J.-C.: The horizontal resolution of MIPAS, *Atmos. Meas. Tech.*, 2, 47–54, <https://doi.org/10.5194/amt-2-47-2009>, 2009.
- Weigel, K., Riese, M., Hoffmann, L., Hofer, S., Kalicinsky, C., Knieling, P., Olschewski, F., Preusse, P., Stroh, F., Spang, R., and Volk, C. M.: CRISTA-NF measurements during the AMMA-SCOUT-O<sub>3</sub> aircraft campaign, *Atmos. Meas. Tech.*, 3, 1437–1455, <https://doi.org/10.5194/amt-3-1437-2010>, 2010.
- Weinreb, M. P. and Neuendorffer, A. C.: Method to Apply Homogeneous-path Transmittance Models to Inhomogeneous Atmospheres, *J. Atmos. Sci.*, 30, 662–666, [https://doi.org/10.1175/1520-0469\(1973\)030<0662:MTAHPT>2.0.CO;2](https://doi.org/10.1175/1520-0469(1973)030<0662:MTAHPT>2.0.CO;2), 1973.
- Wu, D. L. and Waters, J. W.: Gravity-wave-scale temperature fluctuations seen by the UARS MLS, *Geophys. Res. Lett.*, 23, 3289–3292, 1996.