A multi-wavelength numerical model in support to quantitative retrievals of aerosol properties from automated-lidar-ceilometers and test applications for AOT and PM10 estimation

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Abstract. Knowledge of the vertical distribution of aerosol particles is a key factor in the study of climate, air pollution, and meteorological processes. The use of automated lidar-ceilometers (ALC) systems for the aerosol vertically-resolved characterization has increased in the recent years thanks to their low construction and operation costs, and to their capability at providing continuous, unattended measurements. At the same time there is a need to convert the ALC signals into usable geophysical quantities. In fact, the quantitative assessment of the aerosol properties from ALC measurements and the relevant assimilation in meteorological forecast models is amongst the main objectives of the EU COST Action TOPROF (Towards Operational ground-based PROFiling with ALCs, doppler lidars and microwave radiometers). Concurrently, the E-PROFILE program of the European Meteorological Services Network (EUMETNET) focuses on the harmonization of ALC measurements and data provision across Europe. Within these frameworks, we implemented a model-assisted methodology to retrieve key aerosol properties (extinction coefficient, surface area and volume) from elastic lidar and/or ALC measurements. The method is based on results from a large set of aerosol scattering simulations (Mie-theory) performed at UV, visible and near IR wavelengths using a "Monte-Carlo" approach to select the input aerosol microphysical properties. An average 'continental aerosol type' (i.e. clean-to-moderately polluted continental aerosol conditions) is addressed in this study. Based on the simulation results, we derive mean functional relationships linking the aerosol backscatter coefficients to the above-mentioned variables. Applied in the data inversion of single wavelength lidars and/or ALCs, these relationships allow quantitative determination of the vertically-resolved aerosols backscatter, extinction, volume and surface area and, in turn, of the extinction-to-backscatter ratio (i.e., the lidar-ratio, LR) and of extinction-to-volume conversion factor (c_v) at 355, 532, 1064 nm. These variables provide valuable information for visibility, radiative transfer and air quality applications. This study also includes 1) validation of the model simulations with real measurements and 2) test applications of the proposed model-based ALC inversion methodology. In particular, our model simulations were compared to backscatter and extinction coefficients independently retrieved by Raman lidar systems operating at different continental sites within the European Aerosol Research LIdar NETwork (EARLINET). This comparison shows good model-measurements agreement, with LR discrepancies below 20%. The model-assisted quantitative retrieval of both aerosol extinction and volume was then tested using raw data from three different ALCs systems (CHM15k-Nimbus), operating within the Italian Automated Lidar-ceilometer Network (ALICENET). To this purpose, a one-year-record of the ALCs-derived aerosol optical thickness (AOT) at each site was compared to direct AOT measurements performed by co-located sun-sky photometers. This comparison shows an overall AOT agreement within 30% at all sites. At one site, the model-assisted ALC estimation of the aerosol volume and mass (i.e., PM10) in the lowermost levels was compared to values measured at the surface-level by co-located in situ instrumentation. Within this exercise, the ALC-derived daily-mean

38 mass concentration was found to well reproduce the corresponding (EU regulated) PM10 values measured by the local
39 Air Quality agency in terms of both temporal variability and absolute values. Although limited in space and time, the
40 good performances of the proposed approach in these preliminary tests suggest it could possibly represent a valid option
41 to extend the capabilities of ALCs at providing quantitative information for operational air quality and meteorological
42 monitoring.

43 **1   Introduction**

44 Due to the impact of atmospheric aerosols on both air quality and climate, substantial efforts have been made to expand
45 our knowledge of their sources, properties and fate. Aerosol particles affect the Earth's radiation budget mainly by two
46 different processes: 1) by scattering and absorbing both solar and terrestrial radiation (aerosol direct effect, Haywood
47 and Boucher, 2000 and aerosol semi-direct effect, Johnson et al., 2004) and 2) by serving as cloud and ice condensation
48 nuclei (aerosol indirect effect, Lohmann and Feichter, 2005, Stevens and Feingold, 2009 and Feingold et al., 2016). The
49 complexity of these processes and the extreme spatial and temporal variability of the aerosol sources, physical and
50 chemical properties and atmospheric processing make the quantification of their impacts very difficult. Aerosols have
51 also proven detrimental effects on human health (e.g., D'Amato et al., 2013, World Health Organization, 2013,
52 Lelieveld et al., 2015). In fact, their concentration (often evaluated in terms of particulate matter mass, or PM) is
53 regulated by specific air quality legislation worldwide. In Europe, the Air Quality Directive 2008/50 defines the
54 'objectives for ambient air quality designed to avoid, prevent or reduce harmful effects on human health and the
55 environment as a whole' (EC, 2008).

56 Among the aerosol observational systems, the LIDAR technique has been proved to be the optimal tool to provide
57 range-resolved, accurate aerosol data necessary in radiative transfer computations (e.g. Koetz et al., 2006, Tosca et al.,
58 2017) and is often usefully employed in supporting air quality studies (e.g. Menut et al., 1997, He et al., 2012). With a
59 spectrum of different system types (elastic backscatter, Raman, High Spectral Resolution, and multi-wavelength lidars),
60 each with specific pro and cons (Lolli et al., 2018), this technique allows retrievals of aerosol and cloud optical
61 properties and relevant distribution within the atmospheric column at several ground-based observational sites (Fernald
62 et al., 1972; Klett, 1981; Shipley et al., 1983, Kovalev and Eichinger, 2004, Heese and Wiegner, 2008; Ansmann et al.,
63 2012). Since 2006, the Cloud Aerosol Lidar and Infrared Pathfinder Satellite Observation (CALIPSO) platform (Winker
64 et al. 2003) also provides a unique, global view of aerosol and cloud vertical distributions through space-based
65 observations (at the operating wavelengths of 532 and 1064 nm). Recently, within the Cloud-Aerosol Transport System
66 (CATS) mission, a lidar was also installed at the International Space Station (ISS, McGill et al., 2015 and York et al.,
67 2016). Space-borne lidar observations are however affected by some drawbacks, as: 1) limited temporal resolution and
68 spatial coverage (the CALIPSO spatial distance between two consecutive ground tracks is about 1000 kilometers and
69 each track has a footprint of 70 m), 2) the contamination of unscreened clouds, and 3) difficulties in quantitatively
70 characterizing the aerosol properties in the lowermost troposphere (Pappalardo et al., 2010). Ground-based lidar
71 networks thus still represent key tools in integrating space-borne observations to study aerosol properties and their 4D
72 distribution. An example of these networks is the European Aerosol Research LIdar NETwork (EARLINET,
73 http://www.earlinet.org/), which, since 2000, provides an extensive collection of ground-based data for the aerosol
74 vertical distribution over Europe (Bösenberg et al., 2003, Pappalardo et al., 2014). The advanced multi-wavelength
75 elastic and Raman lidars employed in this network allows independent retrieval of aerosol extinction (α_a) and
76 backscattering coefficient (β_a) profiles. Yet, despite their unsurpassed potential in data accuracy, advanced lidar
77 networks such as EARLINET have the unsolved problems of the sparse spatial and temporal sampling and of the

78 complexity of operations. In fact, the typical distance between the EARLINET stations is of the order of several
79 hundreds of kilometers and regular measurements of EARLINET are only performed on selected days of the week
80 (Mondays and Thursdays) and for a few hours (mainly at nighttime, due to low signal-to-noise ratio of the Raman
81 signal in daylight). Furthermore, these systems are complicated to be operated, require specific expertise and are
82 therefore unsuitable for operational applications.

83 Nowadays, hundreds of single channel Automated Lidar Ceilometers (ALCs) are in operation over Europe and
84 worldwide. Although such simple lidar-type instruments were originally designed for cloud base detection only, the
85 recent technological advancements make now these systems reliable and affordable for aerosol measurements,
86 increasing the interest in using this technology in different aerosol-related sectors (e.g. air quality, aviation security,
87 meteorology, etc.). Recent studies showed that the ALC technology is now mature enough to be used for a quantitative
88 evaluation of the aerosol physical properties in the lower atmosphere (Wiegner, M. and A. Geiß, 2012, Wiegner et al.,
89 2014) and the exploitation of the full potential of ALCs in the aerosol remote sensing is a current matter of discussion in
90 the lidar community (e.g. Madonna et al., 2015, 2018). The evaluation of ALC capabilities at providing quantitative
91 aerosol information is among the main objectives of the EU COST Action ES1303, TOPROF (Towards Operational
92 ground-based PROFiling with ALCs, doppler lidars and microwave radiometers). An effort in this direction is also
93 underway in the framework of E-PROFILE, one of the observation programs of the EUropean METorological services
94 NETwork (EUMETNET). In fact, several ALC stations are progressively joining E-PROFILE to develop an operational
95 network to produce and exchange ALC-derived profiles of attenuated backscatter. A recent project funded by the EU
96 LIFE+ program (DIAPASON, Desert-dust Impact on Air quality through model-Predictions and Advanced Sensors
97 ObservatioNs, LIFE+2010 ENV/IT/391) also prototyped and tested an ALC system with an additional depolarization
98 channel, capable of discriminating non spherical aerosol types, such as desert dust (Gobbi et al., 2018). Such upgraded
99 ALC systems could further improve the capabilities of the operational aerosol profiling in a near future.

100 Given the necessity to couple advancement in instrumental technology with tools capable of translating raw data into a
101 robust, quantitative and usable information, we propose and characterize here a methodology to be applied to elastic
102 backscatter lidars and/or ALC measurements to retrieve, in a quasi-automatic way, vertically-resolved profiles of some
103 key aerosol optical and microphysical properties. This effort is intended to contribute better exploit these systems
104 potential in integrating data collected by more advanced lidar systems/networks. In particular, the ALC-derived aerosol
105 properties addressed in this study are: backscatter (β_a, $km^{-1} sr^{-1}$), extinction (α_a, km^{-1}), surface area (S_a, $cm^2 cm^{-3}$) and
106 volume (V_a, $cm^3 cm^{-3}$), the latter being convertible into aerosol mass concentration ($\mu g\ m^{-3}$) via assumption on particle
107 density. To this purpose, we developed a numerical aerosol model to perform a large set of aerosol scattering
108 simulations. Based on results from this numerical model, we derive mean functional relationships linking β_a to α_a, S_a
109 and V_a, respectively. These relationships are then applied in the ALC data inversion and analysis. A similar approach
110 was applied in past studies for lidar-based investigations of stratospheric (Gobbi, 1995) and tropospheric aerosols
111 (maritime, desert dust and continental type) at visible and UV lidar wavelengths, (Barnaba and Gobbi, 2001, Barnaba
112 and Gobbi, 2004a, hereafter BG01, BG04a, respectively, Barnaba et al., 2004). Here we extend this approach to all the
113 Nd:YAG laser harmonics commonly used by both advanced lidars and ALC systems (i.e. 355, 532, 1064 nm
114 wavelengths) and specifically address an 'average-continental' aerosol type, intended to represent clean-to-moderately
115 polluted continental aerosol conditions (see Section 2.1). In fact, despite the known differences that can be encountered
116 across the continent both in the short and the long-term (e.g., Putaud et al. 2010), this aerosol type is expected to
117 climatologically dominate over most of Europe.

118 Overall, this investigation is organized as follows: in Section 2 we describe the aerosol model set up to reproduce clean
119 to moderately polluted continental conditions, and the Monte Carlo methodology followed to compute the
120 corresponding bulk optical and physical properties. Section 3 shows and discusses the results of the numerical model,
121 and presents the model-based, mean functional relationships linking the different variables at 355, 532 and 1064 nm. In
122 Section 4 we evaluate both the model simulations capability to reproduce real measurements in continental aerosol
123 conditions, and the capability of the model-based ALC inversion approach to derive quantitative geophysical
124 information. The EARLINET database was used for the first task while tests on the accuracy of the model-based ALC
125 inversion were performed evaluating both the ALC-derived aerosol volume and optical thickness (AOT, i.e. the
126 vertically integrated aerosol extinction). To this purpose we applied the proposed methodology to three ALC systems
127 operating within the Italian Automated LIdar-CEilometer NETwork (ALICE-NET, www.alice-net.eu). In particular, the
128 ALC-derived AOT and aerosol volume (plus mass) were compared, respectively, to reference measurements performed
129 by ground-based sun photometers and in situ aerosol instruments (optical counters and PM10 samplers).

130 Section 5 summarizes the developed approach and main results, critically examining strengths and weaknesses. It also
131 includes discussion on the perspectives of the application of this (or similar) methodology in operational ALC
132 networks.

133

134 **2 The aerosol model**

135 A numerical aerosol model was set up to calculate mean functional relationships between the aerosol backscatter (β_a)
136 and some relevant aerosol properties, as α_a, S_a and V_a. This is done in a two-step procedure (Figure 1), following an
137 approach similar to that developed by BG01 and BG04a:

138 1) Generation of a large set (here 20000) of aerosol optical and physical properties by randomly varying, within
139 appropriate ranges, the microphysical parameters describing the aerosol size distribution and composition (blue box in
140 Fig. 1);

141 2) Based on results at point 1), determination of mean functional relationships linking such key variables (grey box in
142 Fig. 1).

143 The following Section describes rationale and set-up of the first step, the second one being thoroughly discussed in
144 Section 3.

145 **2.1 Selection of the aerosol microphysical parameters**

146 As anticipated, an average 'continental' aerosol type (i.e. describing clean to moderately polluted continental
147 conditions, e.g. Hess et al., 1998) was targeted in this study, this being the aerosol type expected to dominate over
148 Europe. Based on a scheme originally proposed by d'Almeida et al. (1991), and on a large set of following
149 observational evidences (e.g. Van Dingenen et al., 2004), in this work its size distribution is described as an external
150 mixture of three size modes. These are (in order of increasing size range): 1) a first ultrafine mode; 2) a second fine
151 mode, mainly composed of water-soluble particles; 3) a third mode of coarse particles.

152 A three-mode lognormal size distribution described by Eq. (1) is employed to this purpose:

153 $$n(r) = \frac{dN}{d\log r} = \sum_{i=1}^{3} \frac{N_i}{\sqrt{2\pi}\log\sigma_i} exp\left[-\frac{(\log r - \log r_i)}{2(\log\sigma_i)^2}\right] \tag{1.}$$

154 In Eq. (1), r_i, σ_i and N_i are respectively the modal radius, the width and the particle number density of the ith aerosol

155 mode (i = 1, 2, 3). At each computation, r_i and σ_i are randomly chosen within a relevant variability range. Values of N_i

156 are conversely obtained by firstly randomly choosing the total number of particles, N_{tot}, to be included in the whole size

157 distribution ($N_{tot} = N_1 + N_2 + N_3$), and then by applying specific rules for the number mixing ratio, $x_i = N_i/N_{tot}$, of each

158 component to this total. To reproduce clean to moderately polluted continental conditions, the value of N_{tot} is made

159 variable between 10^3 and 3×10^4 cm^{-3} (e.g. Hess et al., 1998; Van Dingenen et al., 2004). Being the result of different

160 sources/processes, the three modes are also assumed to have a different composition, this impacting the optical

161 computations through the relevant particle refractive index (m_i), with both its real and imaginary component ($m_i = m_{r\_i}$

162 $- i \times m_{im\_i}$). The Mie theory for spherical particles of radius r_i and refractive index m_i is then used to compute the

163 extinction and backscatter coefficients (see below).

164 A description of the assumptions made for each mode and relevant parameter, mostly based on literature data (Table 1),

165 is given hereafter, the summary of the relevant variability chosen for each parameter being provided in Table 2.

166 1) First Mode

167 This ultrafine mode is the one more directly simulating fresh, anthropogenic emissions. The number mixing ratio $x_{i=1}$

168 ($N_{i=1}/N_{tot}$) of this mode is let variable between 10% (rural conditions, Van Dingenen et al., 2004) and 60% (more

169 polluted conditions, Hess et al., 1998). The variability of its modal radius ($r_1 = 0.005 - 0.03$ µm) is chosen to include

170 from nucleation mode particles to Aitken mode particles. To take into account the wide variability of species within this

171 ultrafine mode, from non-absorbing (e.g., inorganic particles) to highly absorbing materials (e.g. black carbon), wide

172 ranges of variability has been set for its refractive indexes (at λ=355 nm: m_{r_1} in the range 1.40 - 1.8, and m_{im_1} in the

173 range 0.01 – 0.47, see Table 2 for the corresponding values at λ=532 and 1064 nm).

174 2) Second mode

175 The second aerosol mode accounts for 40-90% of N_{tot}, with (dry) r_2 between 0.03 and 0.1 µm. Its composition (m_{r_2},

176 and m_{im_2}) is also made highly variable so to include water soluble inorganic and organic particles (Hess et al., 1998;

177 BG04a; Dinar et al. 2008). In this case, at λ=355 nm, m_{r_2} is in the range 1.40 - 1.7 and m_{im_2} is in the range 0.0001 –

178 0.01 (Table 2).

179 3) Third mode

180 This coarser aerosol mode (modal radius r_3 in the range 0.3 – 0.5 µm) is mainly intended to account for soil derived

181 (dust-like) particles that are a primary continental emission. A quite narrow variability is thus fixed for its m_{r_3} and

182 m_{im_3} values (1.5 – 1.6 and 0.0001 – 0.01, respectively at 355 nm). The relevant number mixing ratio x_3 (N_3/N_{tot}) is set

183 variable between 0.001% and 0.5 %, this mode contributing mostly to the total aerosol volume (thus mass) rather than

184 to the total number of particles.

185 As mentioned, refractive indexes were also made wavelength dependent, as this feature is also typically observed as

186 linked to the different particle composition. In particular, for the second mode (water-soluble particles) we include an

187 increase with the wavelength of the upper boundary values of m_{im_2} and a decrease of m_{r_2} at λ =1064 nm (d'Almeida et

188 al., 1991). For the (dust-like) third-mode particles, the upper boundary values of m_{im_3} are set to decrease with

189 increasing wavelengths (Gasteiger et al., 2011, Wagner et al., 2012).

190 For convenience, the aerosol parameters boundaries summarized in Table 2 refer to dry particles and to ground level.

191 However, the effect of a variable RH, its variability with altitude as well as the generally observed decrease of particle

192   number with altitude is also considered in the model. More specifically, the number of particles in each mode, N_i, and

193   RH are both made altitude-dependent through the following equations (Patterson et al., 1980, BG01):

194    $$N_i(z) = N_i(0) \times \exp\left(\frac{-z}{H_i}\right), \tag{2}$$

195    $$RH(z) = 70 \times \exp\left(\frac{-z}{5.5\ km}\right) \times (1 + dRH), \tag{3}$$

196   the altitude z being variable here between 0 and 5 km. $N_i(0)$ and H_i in eq.2 are the number of particles at the ground and

197   the scale height for each mode, respectively.

198   To describe the altitude effect, in eq. (2) an exponential decrease with height of the particle number density is assumed.

199   To rescale the particle number density of the different modes, $H_{i=1=2}$ is set equal to 5.5 km (Barnaba et al., 2007) while

200    $H_{i=3}$ (coarse particles) is set to 0.8 km (Barnaba et al., 2007). In eq. (3), the additional term (1+dRH) allows a further

201   variability with respect to the mean RH(z) profile assumed, here dRH is randomly chosen between -60 and +60). Values

202   of RH greater than 95% are discarded to avoid divergence.

203   Additionally, while first and third modes are assumed to be water insoluble, the second mode (i=2) is fully hygroscopic.

204   Aerosol humidification is thus considered to act on both particle size and refractive indices of the second aerosol mode

205   (e.g., BG01), as:

206    $$r_{2\_RH} = r_{2\_0}\sqrt{\frac{2-0.01RH}{2(1-0.01RH)}} \quad , \tag{4}$$

207    $$m_{2\_RH} = m_w + (m_{2\_0} - m_w)\left(\frac{r_{2\_0}}{r_{2\_RH}}\right)^3, \tag{5}.$$

208   In eq. 4 and 5, r_{2_RH} and m_{2_RH} are the RH-corrected modal radius and refractive index for the second mode,

209   respectively; r_{2_0} and m_{2_0} are the particle dry modal radius and refractive index, respectively; m_w is the water refractive

210   index (assumed as equal to $1.34 – i7e^{-9}$, $1.33 - i1.3e^{-9}$ $1.33 - i2.9e^{-6}$ at 355, 532 and 1064 nm, respectively).

211   Finally, following Barnaba et al. (2007), an increase of the width of the size distribution with altitude (eq. 6) has been

212   introduced for the first and second aerosol mode:

213    $$\sigma_{1,2}(z) = \sigma_{1,2,z0} \times \exp\left(\frac{z}{30}\right). \tag{6}$$

214   In fact, Barnaba et al., (2007) showed that this was necessary to better reproduce the observed decrease of the Lidar

215   Ratio (LR) with altitude, and likely related to a broadening of the particle size distribution with aging.

216   Once the value of each microphysical parameter is randomly selected within its relevant variability range, and once

217   corrections are applied following eqs. (2) – (6), each resulting aerosol size and composition-resolved distribution is used

218   to compute the aerosol S_a and V_a, as well as to feed a Mie code (assumption of spherical particles, Bohren and

219   Huffman, 1983) to compute β_a, and α_a, (BG01, see also Fig. 1). Overall, the equations used are as follows:

220    $$\beta_a = \int Q_{bsc}(r,\lambda,m)\,\pi r^2 \frac{dN}{d\log r}\frac{1}{r\ln 10}\,dr \tag{7}$$

221    $$\alpha_a = \int Q_{ext}(r,\lambda,m)\,\pi r^2 \frac{dN}{d\log r}\frac{1}{r\ln 10}\,dr \tag{8}$$

222    $$S_a = 4\pi \int r^2 \frac{dN}{d\log r}\frac{1}{r\ln 10}\,dr \tag{9}$$

223 $$V_a = \frac{4}{3}\pi \int r^3 \frac{dN}{d\log r} \frac{1}{r\ln 10} dr,$$ (10)

224 where Q_{bsc} (r_i, λ, m_i) and Q_{ext} (r_i, λ, m_i) are, respectively, the backscatter and the extinction efficiencies. As mentioned,

225 the optical computations are made at the three different wavelengths: 355, 532, 1064 nm (i.e., those of Nd:YAG laser

226 harmonics, the most common wavelengths used by ground-based and space-borne aerosol lidars).

227 Since in our simulations the third aerosol mode is intended to represent dust-like particles, an empirical correction for

228 non-sphericity is also applied to the Mie-derived optical properties of this mode. This procedure is based on BG01,

229 which uses the results of Mishchenko et al. (1997) obtained for surface-equivalent mixtures of prolate and oblate

230 spheroids.

231 **2.2   Model simulation results**

232 In Figure 2 we show the results of 20000 simulations of continental aerosol optical and physical properties derived

233 randomly varying the relevant aerosol size distributions and compositions as described in the previous section. In

234 particular, the results for α_a, S_a and V_a are shown as a function of β_a in Figure 2a, b, c (blue crosses) referring to $\lambda =$

235 1064 nm. For each variable (A), average value per bin of β_a and relevant standard deviations (<A> ± dA) are shown as

236 red dots and vertical bars, respectively. Note that 10 equally spaced bins per decade of β have been considered, and that

237 <A> ± dA are only shown for bins containing at least 1% of the total number of pairs. Corresponding relative errors

238 (dA/<A>) are depicted in Figure 2d, e, f. Some sensitivity tests of these model outputs to the variability of the input

239 microphysical parameters employed are provided in Appendix A.

240 Based on these results, at step-two of the procedure (see scheme in Figure 1), we derive aerosol-specific mean

241 relationships linking aerosol extinction, surface area and volume (α_a, S_a and V_a) to its backscatter (β_a). To this purpose,

242 we used a seventh-order polynomial fit in log-log coordinates. The choice of a seventh-order polynomial fit was made

243 for homogeneity with BG01 and BG04a. These relationships are shown as green lines in Figure 2a, b, c while the

244 relevant fit parameters are reported in Table 3 referring to $\lambda = 1064$ nm (fit parameters related to computations at $\lambda =$

245 355 and 532 nm, are given in Table A1 and Table A2, Appendix B).

246 The red vertical bars of Figure 2 also highlight the ranges of α_a, S_a and V_a which are statistically significant, i.e. those in

247 which, at $\lambda = 1064$ nm, the model provides at least 1% of the total points per corresponding bin of β_a. These are: 10^{-4} -

248 10^{-1} km^{-1}, 10^{-7} - 10^{-5} cm^2/cm^3 and 10^{-13} - 10^{-10} cm^3/cm^3, for α_a, S_a and V_a respectively, corresponding to the backscatter

249 range $9\times10^{-5} \leq \beta_a \leq 4\times10^{-3}$ km^{-1} sr^{-1}. In terms of aerosol properties variability, the relative errors associated to α_a and V_a

250 show almost no dependence on β_a, with values between 30% and 40%. Conversely, the modeled aerosol surface area

251 exhibits a larger dispersion, with relative error values spanning the range 40% - 70%, and decreasing as β_a increases.

252 A key parameter for the inversion of lidar signals is the so-called Lidar Ratio (*LR*), i.e. the ratio between α_a and β_a

253 (Ansmann et al., 1992). In Figure 3 we thus show the results of our simulations in terms of LR vs β_a at the three λ (355,

254 532 and 1064 nm, Figure 3a, b, c, respectively) and relevant dLR/LR values (Figure 3d, e, f, respectively). The color

255 code is the same of Fig. 2. Additional horizontal black lines have been inserted representing values (solid central lines)

256 of the 'weighted-LR' ± 1 standard deviation (dotted side lines), i.e. the LR weighted by the number of simulated points

257 in each considered backscatter bin. The 'weighted-LR' values derived at 355, 532 and 1064 nm, are 50.1 ± 17.9 sr, 49.6

258 ± 16.0 sr and 37.7 ± 12.6 sr, respectively. Figure 3 also allows showing that the statistically significant regions of

259 simulated backscatter values shifts towards smaller values with increasing λ (e.g. at $\lambda = 355$, the β_a extending regions is

260 4×10^{-5} - 2×10^{-2} km^{-1} sr^{-1}, whereas, at 532 nm, it ranges between 2×10^{-5} - 1×10^{-2} km^{-1} sr^{-1}). Furthermore, Figure 3

reveals a quite different shape of the LR vs β_a functional relationships (green curves) at different wavelengths. At 355 and 532 nm the curve is concave, with quite similar LR maxima of the fitting curve (54.3 and 53.8 sr at approximately $\beta_a = 4\times10^{-4}$ km^{-1} sr^{-1} and 2×10^{-3} km^{-1} sr^{-1}, respectively). At 1064 nm the curve is conversely monotonic, with a flex point at $\beta_a = 3$-4×10^{-4} km^{-1} sr^{-1}. A larger data dispersion also characterizes the results at $\lambda = 355$ and 532 nm (LR values from 10 to 90 sr) in comparison to $\lambda = 1064$ nm (LR in the range 18 – 80 sr, except for a minor number of outliers). This translates into different LR relative errors at UV, VIS and infrared (IR) wavelengths. At 1064, dLR/LR slightly decreases for increasing backscatter, with values around 35%. At the shorter wavelengths, it increases as a function of β_a, with a large (>40%) relative error for values of $\beta_a > 2\times10^{-3}$ km^{-1} sr^{-1}.

To insert our results into a more general context, we compared the derived, model-based weighted-LR values to some LR data reported in the literature (Table 4). In particular, we selected some of the works using the aerosol model developed to invert the Calipso lidar data (Omar et al., 2009). This latter considers six different aerosol sub-types: clean continental (CC), clean marine (CM), dust (D), polluted continental (PC), polluted dust (PD), and smoke (S). Our model-derived LR at 532 nm falls in the middle of the range (35-70 sr) fixed by the Calipso CC and PC aerosol classes. The work by Papaggianopoulos et al. (2016), in which the LR values are adjusted accordingly to EARLINET observations, reports a LR range at 532 nm of 47-62 sr. At the same wavelength, the aerosol range defined by the LIVAS climatology (LIdar climatology of Vertical Aerosol Structure for space-based lidar simulation studies, Amiridis et al., 2015), is 54-64 sr. In both cases, our model seems to be closer to the LR values of CC aerosol type, which is compatible to our intention to simulate clean-to-moderately polluted continental aerosol type. At 532 nm, our LR value is also reasonably in between the CC and PC LR values derived by Omar et al. (2009), but again closer to the CC LR value. The very small decrease of LR values between 532 and 355 nm estimated by LIVAS for the CC aerosol is also consistent with our results. Similarly, our model predicts a lower mean LR in the near IR with respect to the green, in agreement with results of Amiridis et al. (2015) in CC conditions and not to those in polluted conditions. Table 4 also includes the continental aerosol LR values estimated in the work of Düsing et al. (2018) through comparison between airborne in situ and ground-based lidar measurements. Our model is in good agreement with their LR values at 355 and 532 nm. At 1064 nm, the algorithm developed by Düsing et al. (2018) provided a value of LR around 15 sr. On the other hand, in the same study the authors found that, rather, a value of LR = 30 sr gives the better accord between their Mie and lidar-based α_a, this value being closer to our model-derived one at 1064 nm (LR =37.7). The difference between these two values is explained by the authors to be probably due to the estimation of the aerosol particle number size distribution, a critical parameter for a reliable modeling of aerosol particle backscattering.

As a last added value of the outcome from our model-based results, we derive here and provide in Table 5 extinction-to-volume conversion factors, $c_v = V_a/\alpha_a$ (e.g., Ansmann et al., 2010) at three different wavelengths (355, 532 1064 nm), and compare these to similar outcomes from other studies. To our knowledge, values of continental particles c_v at three wavelengths are only available in Mamouri and Ansmann (2017). Note that c_v, is also proportional, through the particle density ρ_a, to the inverse of the so-called 'mass-to-extinctions efficiency' (MEE, i.e. $\alpha_a/(V_a*\rho_a)$) a parameter important in several aerosol-related applications (e.g. the estimation of particulate matter mass from satellite AOT or in modules of global circulation and chemical transport models to compute aerosol radiative forcing effects, Hand and Malm, 2007). For convenience, model-derived MEE values are also included in Table 5.

3 Evaluation of the model performances and potential of its application

In this section, we evaluate the capability of the model results to reproduce 'real' aerosol conditions and explore the potential of the proposed model-based ALC inversion in producing quantitative geophysical information. In particular:

- In Section 3.1 we compare our simulations to real observations of independent backscatter and extinction coefficients made by different EARLINET Raman lidars (Bösenberg et al., 2001, Pappalardo et al., 2014).

- In Section 3.2, our model results are used to invert measurements acquired by some ALCs systems operating within ALICE-NET, which networks several ALCs systems (Nimbus CHM15k by Lufft) located across Italy and run by Italian research institutions and environmental agencies. Here we use data from some of these systems to derive the aerosol optical and physical properties (e.g. the aerosol optical thickness, AOT, and the aerosol volume and mass).

3.1 Comparison of the modelled aerosol optical properties to EARLINET measurements

As mentioned EARLINET Raman stations perform coordinated measurements two days per week following a schedule established in 2000 (Bösenberg et al., 2003). Overall, the EARLINET database includes the following categories: 'climatology', 'CALIPSO', 'Saharan dust', 'volcanic eruptions', 'diurnal cycles', 'cirrus', and 'others' (forest fires, photo smog, rural or urban, and stratosphere). To be comparable to our results, we used EARLNET β_a and α_a coefficients at 355 nm and at 532 nm within the quality assured (QA) 'climatology' category (Pappalardo et al., 2014). However, note that additional data filtering was necessary to screen out residual, likely unreliable values within this QA-'climatology' category. In particular, we only selected those EARLINET QA data further satisfying the following criteria:

- β_a and α_a coefficients evaluated independently, i.e. only obtained using the Raman method (Ansmann et al., 1992);
- β_a and $\alpha_a > 0$;
- $LR < 100$;
- Relative errors on β_a and $\alpha_a < 30\%$.

Then, we selected those sites in Europe expected to be mostly impacted by 'continental' aerosols and having the largest datasets (e.g., at least 100 points) at 355 and 532 nm. Overall, 5 sites satisfied these conditions (Table 6), and namely Madrid (Spain), Potenza and Lecce (Italy), Leipzig and Hamburg (Germany). Finally, being interested in continental conditions here, we filtered out those measurements dates affected by desert dust at the measuring sites, i.e. we removed from our 'model-measurement comparison data set' all the dates within the EARLINET 'climatology' category also belonging to the EARLINET 'Saharan dust' category.

Figure 4 depicts the results of the model-measurements comparison at the sites fulfilling our requirements in terms of LR vs β_a at λ=355 nm (the corresponding results at λ=532 nm, including Madrid in place of Hamburg, are given in Appendix C, Figure C1). The colored area represents the model-simulated data range, while the color code indicates the absolute number of simulated values (i.e. counts) in each β_a - LR pair. The EARLINET-measured values are reported as black open circles. Note that, being the model simulations performed over an altitude range 0-5 km (see Section 2.1) only those simulations corresponding to the altitude range (Δz) covered by the measurements at each EARLINET station was taken into account here. Figure 4 shows the model results to well encompass the measured LR vs β_a data, with few measurements outside the modeled range (most of the exceptions are found for Potenza). Statistically, the

9

337 highest number density of simulated data well fits the observations, with the exception of Hamburg (Figure 4a), which
338 however has the lowest number of measured data (it is not an EARLINET station any longer, see Table 6).

339 In Figure 5 the previous results at λ=355 nm are converted in terms of 'mean' LR per bin of β_a for both model (blue)
340 and observations (red, again, only β_a bins containing at least 1% of the total modeled data were considered). This view
341 shows that there is a general good agreement between the modeled and the measured LR values, and in their variation
342 with β_a. Some major deviations are found for Potenza and are further discussed in the following. The model-
343 measurements accordance shown in Figure 5 was evaluated in quantitative terms by computing mean LR relative
344 differences at both λ = 355 and 532 nm, i.e., we derived ($[(LR_{mod}- LR_{meas})/ LR_{meas}]*100$) values, where LR_{mod} and
345 LR_{meas} are the lidar ratio values computed by model and derived by lidar measurements, respectively. These values are
346 reported in Table 7 for each considered EARLINET station, together with the measurements-based mean LR in each
347 observational site (computed weighting the number of observations per β_a-spaced bins).

348 Results in Figure 5 and Table 7 also give some hints on the capability of the aerosol type assumed (and its admitted
349 ranges of variability) to reproduce 'real' continental aerosol conditions in different sites across Europe. In fact, the four
350 continental sites selected with our criteria are still expected to be partially impacted by different aerosol types.

351 - A good agreement between the model and the observations in terms of LR mean values is found for Hamburg (Figure
352 5a), with mean LR differences of the order of 5% (Table 7). Still, the measured LR values have a high variability and
353 their distribution is positioned towards high values of β_a (1×10^{-3} to 4×10^{-3} km^{-1} sr^{-1}). This could be due to the presence
354 of different aerosols types as slightly polluted marine and polluted aerosol (Matthias and Bösenberg, 2002).

355 - A good accord for Leipzig (Fig. 5c) also indicates that this site is mostly dominated by 'pure' continental particles. In
356 fact, the distribution of observed LR points in Fig. 4, which covers β_a values ranging from 2×10^{-4} to 3×10^{-3} km^{-1} sr^{-1}, is
357 well centered to the modeled simulations highest density (counts > 40). Table 7 shows that at both wavelengths mean
358 discrepancies with LR measurements keep well below 10%.

359 The highest differences in Fig. 5 are found in some southern Europe EARLINET sites:

360 - In Lecce (Fig. 5b), the best agreement between model and observations is found for the lowest values of β_a (between
361 9×10^{-4} to 1×10^{-3} km^{-1} sr^{-1}, see Table 7). Also, the increase from 10% to 18% in the discrepancies at 355 and 532 nm
362 indicates some model problems in correctly reproducing the spectral variability of the optical properties, suggesting
363 some mismatch between modeled and real aerosol sizes in this site (see discussion below).

364 - In Potenza (Fig. 5d), a significant difference between the mean LR curves emerges for β_a values > 6×10^{-4} km^{-1} sr^{-1} ,
365 with observed LR values lower than those simulated here.

366 These discrepancies could be due to the influence of marine aerosols at both stations (De Tomasi et al., 2006, Mona et
367 al., 2006, Madonna et al., 2011), which is expected to produce lower LR values for high values of β_a (e.g. BG01). In
368 fact, Madrid shows better performances, with dLR/LR comparable to those in Leipzig.

369 To provide some insight into the reasons of the model-measurements differences at LC and PO sites, some specific
370 model sensitivity tests have been performed and are reported in Appendix D. In particular, for Lecce, we found that
371 better agreement between the observed and simulated LR vs β_a behavior at 355 nm is obtained by reducing the
372 variability range of N_{tot} (from 500 - 10000 cm^{-3} to 500 - 5000 cm^{-3} at ground). This indicates that LC is likely affected
373 by cleaner continental aerosol type conditions. The sensitivity simulations done for understanding the mismatches with
374 Potenza measurements show that an extension of the variability range of the coarse mode radius is needed to reproduce

375 the observed decrease of LR for increasing backscatter (Figure 5d). This suggests the presence of coarse particles larger
376 than those assumed in such clean continental environment (Appendix D). This is compatible with the suspect of marine
377 air contamination, although at this stage we are not able to exclude additional contamination of coarser particle of soil
378 origin.

379 Overall, mean LR differences between our 'average-continental' model and data at selected continental sites in Europe
380 keep lower than 20% (Table 7), this indicating it reasonably well reproduces the clean-to-moderately polluted
381 continental aerosol conditions we intended to simulate.

382 **3.2   Model results application to Nimbus CHM15-k ALC measurements**

383 To test and validate the model-based inversion methodology, we used the derived functional relationships (Section 2.2)
384 to invert and analyze the measurements of some ALICENET ALCs (Lufft CHM15k systems). These instruments are
385 biaxial ceilometers that emit laser pulses at 1064 nm (Nd:YAG-laser, class M1) with a typical pulse energy of 8 μJ and
386 a pulse repetition rate of about 6500 Hz. The instruments have a specified range of 15 km and full overlap at around
387 1500 m (Heese et al., 2010). The manufacturer provides the overlap correction functions ($O(z)$) for each system. As
388 shown recently by Wiegner and Geiß (2012) and Wiegner et al. (2014), a promising strategy to retrieve the aerosol
389 backscatter coefficient from ALC measurement is adopting the forward solution of the Klett inversion algorithm (Klett
390 1985). This solution requires a known calibration constant of the system (i.e. absolute calibration, c_L) and an
391 assumption on the LR. The advantage with respect to the backward solution is that calibration is not affected by the low
392 SNR in the upper troposphere and it is needed occasionally. Furthermore, starting close to the surface, the data retrieval
393 allows resolving aerosol layers in the boundary layer even if their optical depth is high. The forward solution of the
394 Klett inversion algorithm is thus adopted here. For convenience, we report here the equations used within our procedure
395 to obtain β_a from ALC measurements, which are also described in Wiegner and Geiß (2012, equations 1 – 3):

396 $$\beta_a(z) = \frac{Z(z)}{LR\ N(z)} - \beta_m(z)$$ (11)

397 with

398 $$Z(z) = LR\ z^2 P(z) exp\left[-2 \int_0^z (LR\ \beta_m - \alpha_m)dz'\right]$$ (12)

399 and

400 $$N(z) = c_L - 2 \int_0^z Z(z')dz'.$$ (13)

401 Here, β_m and α_m are the molecular backscatter and extinction coefficients calculated from climatological, monthly air
402 density profiles and $z^2 P(z)$ is the ALC range (z) corrected signal (P) (also referred to as RCS), that is the raw data
403 obtained by the considered ALCs. As anticipated, knowledge of the calibration constant c_L is needed to solve eq. 13
404 (and thus 11, forward solution). In our analysis of ALC daily records, the constant c_L has been obtained by the
405 "backward approach" (Rayleigh calibration) applied to night-time, cloud-free ALC signal averaged over 1 or 2 hours at
406 75 m height resolution. This allows for using the best c_L retrieval (that is the night-time, lowest noise one), in the
407 forward solution of the lidar equation, which guarantees operating over the best signal to noise range of the ALC signal.

3.2.1 Model-based retrieval of aerosol optical properties

Operatively, inversion of the aerosol properties, $\alpha_a(z)$ and $\beta_a(z)$, is performed using an iterative technique, since we need to correct the backscatter signal at each altitude z for extinction losses. The iterative procedure is stopped when convergence in the integrated aerosol backscatter (IAB=$\Sigma_0^{zcal}\beta_a(z)$) is reached (e.g. BG01). At each step, aerosol extinction is derived using the functional relationship $\alpha_a = \alpha_a (\beta_a)$ of Table 3.

An example of the outcome of this retrieval methodology is depicted in Figure 6. It shows the time-height (24h, 0 - 6 km) contour plot of α_a retrieved at 1064 nm during a whole day of measurements (June 26, 2016) performed by the ALICENET system of Aosta San Christophe (ASC, 45.8°N, 7.4°E 570 m a.s.l., Northern Italy, Figure 7a). Time and altitude resolutions are 1 min and 15 m, respectively. Note that ALC data are cloud-screened using the cloud mask of the Lufft firmware.

The aerosol optical thickness (AOT) is obtained vertically integrating the ALC-derived $\alpha_a(z)$ from the surface up to a fixed height z_{AOT}, above which the aerosol contribution is assumed to be negligible. In Figure 6, the ALC-derived AOT values at 1064 nm (pink curve, with a temporal resolution of 5 min) is superimposed to the extinction contour. Reference AOT values from a co-located sun-sky radiometer (a Prede POM-02 system) are shown by orange circles. These were extrapolated at 1064 nm from the instrument 1020 nm-channel using the Angström exponent derived fitting AOT values at all the radiometer wavelengths. This example illustrates the very good performances of our model-assisted inversion scheme, and the capability of this approach to extend to nighttime the (daylight-only) radiometer observations.

To evaluate the performances of our model-assisted retrieval of $\alpha_a(z)$ over a more statistically significant dataset, the same approach illustrated in Figure 6 was applied to a longer record in the ASC site, plus Nimbus CHM-15k ALC datasets from two additional ALICENET sites: San Pietro Capofiume (SPC, 44°39N, 11°37E, 10 m a.s.l.) and Rome Tor Vergata (RTV, 41.88°N, 12.68°E, 100 m a.s.l.). The location of the instruments is shown in Figure 7a (red circles), while some information on system types and site characteristics is given in Table 8. The data analyzed here were collected during the following periods: April 2015 – June 2017, June 2012 – June 2013 and February 2014 – September 2015, for ASC, SPC and RTV, respectively.

In those sites, reference AOTs were collected by three co-located sun-sky radiometer, and namely using two SKYNET Prede sun-sky radiometers at ASC and SPC (POM-02L and POM-02, respectively, www.euroskyrad.net) and an AERONET Cimel CE 318-2 instrument operational at RTV (https://aeronet.gsfc.nasa.gov, Rome Tor Vergata station, data level 2.0). Only AOT values between 0.01 and 0.2 at 1064 nm were considered. This range allows for excluding the data points with 1064 nm-AOT lower than the sunphotometer expected accuracy (dAOT=0.01) and those where we found aerosol extinction to cause significant deterioration of our ALC signal. Overall a total of 1237, 268, 850 AOT pairs were analyzed at ASC, SPC and RTV, respectively.

Also note that, although CHM-15k data are already corrected for the $O(z)$ function provided by the manufacturer, the variation of the ALC internal temperature was shown to lead to $O(z)$ differences up to 45% in the first 300 m above ground (Hervo et al., 2016). For this reason, in our analyses the lowest valid altitude of the CHM-15k for both the SPC and RTV systems was fixed to be about 400 m. A linear fit of the first two valid ALC points is then used to extrapolate $\alpha_a(z)$ down to the ground (z_0). Conversely, due to the optimal characterization down to the ground of $O(z)$ provided by Lufft for the CHM-15k system installed at ASC, values at z_0 at this site are not those extrapolated but actually those

12

446　measured. The maximum altitude of aerosol extinction vertical integration to derive the AOT, z_{AOT}, was selected as the
447　first height above 4000 m where the range corrected signal (RCS) has a SNR < 1.

448　Results of the long-term AOT comparison are summarized in Figure 7 and Table 9. For each site under investigation,
449　Figure 7 shows the histograms of the AOT differences between the hourly-mean coincident AOTs as derived by the
450　ALCs and measured by the sun-photometers (red curve, corresponding AOT vs AOT scatter plots at the three
451　considered sites are given in Appendix E). To evaluate the advantage of our approach with respect to more standard
452　lidar inversions, we also computed AOT differences using two fixed-LR values. In particular, we used LR = 52 sr (i.e.
453　the value suggested by the E-Profile network, black lines) and LR = 38 sr (i.e. the weighted mean LR value derived
454　from our model, see Section 3, blue lines). Figure 7 shows that the best agreement is found at ASC. The distribution of
455　AOT difference has a maximum around 0 for each of the three inversions schemes, with very low dispersion. The full
456　width at half maximum, FWHM, is in fact around 0.015, and approximately 55% of the data are included in the interval
457　-0.01 – 0.01, which is even within the expected error of photometric measurement. For SPC and RTV, the red and blue
458　histograms are peaked around 0, whereas the black ones are shifted, with maxima around 0.01-0.02 and 0.02-0.03 for
459　SPC and RTV, respectively. These two sites have higher dispersion (FWHM = 0.03), and approximately 30% of the
460　data are included in the interval -0.01 – 0.01 for the red and blue histograms at both sites, which is probably due to the
461　different aerosol load affecting the different ALICENET stations. As pointed out by the low value of the average AOT
462　computed at ASC for the analyzed dataset (<AOT> = 0.027), low pollution levels generally characterize this site, with
463　some exceptions due to wind-driven aerosol transport from the nearby Po valley (Diémoz et al,2018a, 2018b this issue).
464　On the contrary, RTV (<AOT> = 0.044) and, especially, SPC in the Po Valley (<AOT> = 0.076) are characterized by
465　higher aerosol content and pollution levels, which explain the larger histogram dispersions. Note that the high frequency
466　of fog events in winter markedly reduces the number of analyzed AOT pairs at SPC site, while some desert dust
467　affected days at both SPC (e.g., Bucci et al., 2018) and RTV (e.g., Barnaba et al., 2017) were removed from our
468　datasets (no desert-dust affected dates in ASC).

469　Table 9 summarizes the long-term performances of the model-based procedure in deriving quantitative AOT from the
470　ALC systems at the three investigated sites. It includes values of the average differences between the ALC-derived and
471　sunphotometers-measured AOT (both bias, < dAOT >, and absolute difference <|dAOT|>, with associated standard
472　deviations) obtained using both the proposed model-based approach and the fixed-LR inversions. For SPC and RTV
473　sites, these numbers show that the best ALC–photometer accordance is reached when employing either the model-based
474　or the fixed LR=38 sr inversion scheme. In fact, these two approaches have similar performances in terms of mean
475　dAOT values (<|dAOT|> = 0.011, 0.013 and 0.013, 0.014 for SPC and RTV, respectively), mean percent error
476　(<|dAOT|> /<AOT> = 0.16, 0.19 and 0.31, 0.33) and a very low mean relative bias (<dAOT>/<AOT> = -0.043, 0.005
477　and 0.088, 0.11). On the other hand, the fixed LR=52 sr retrieval produces an overestimation of AOT in both SPC and
478　RTV (<dAOT>/<AOT> = 0.33 and 0.44) with larger discrepancies between retrieved and observed AOTs (<|dAOT|>
479　= 0.021 and 0.026, <|dAOT|>/<AOT> = 0.38 and 0.49). For the ASC site, due to the low aerosol content, the
480　differences among the inversion schemes are almost negligible.

481　Overall, for the three sites, the statistics over the long-term datasets employed showed good results of the model-based
482　approach with similar behavior of the retrievals with a fixed LR of 38 sr, while a fixed LR value of 52 sr produces an
483　overestimation of the AOT at SPC and RTV. As different sites have different (and not known a-priori) characteristic LR
484　values, these results highlight the potential of the model-based approach to derive quite accurate β_a and α_a coefficients
485　without the need to choose and fix an arbitrary LR value.

3.2.2 Model-based retrieval of aerosol volume (and mass)

In this section we provide examples of the applicability of the proposed approach to derive air-quality relevant parameters. In particular, we use the ALC, β_a-retrieved data and the 7th-order polynomial fit linking β_a (at $\lambda = 1064$ nm) to V_a (see also Table 3 and Figure 2c) to derive the aerosol volume (and mass).

The ALC-estimates were firstly compared to aerosol volume derived in situ at the ASC site by two different optical particle counters (OPCs) on 29th December 2016 and 5th September 2017. For the case of the 29th December 2016, a TSI Optical Particle Sizer (OPS) 3330 was employed. This instrument has 16 channels that can be programmed to provide the number concentration at different (and logarithmically spaced) diameter size ranges within the interval 0.3 - 10 μm. Further details can be found in the TSI manual (2011). For the case of the 5th September 2017, the Fidas®200s OPC was used. This spectrometer is able to retrieve high-resolution particle spectra (size measurements between 0.15 and 27 μm, with 32 channels/decade, Pletscher et al., 2016). For both dates, Figure 8 shows the time (x-axis, 24h) vs. height (left y-axis) contour plots of the ALC-based retrieval of the aerosol volume concentration (cm^3/cm^3). The OPC-derived aerosol volume concentration measured at ground-level is reported as a function of time (x-axis) on the right y-axis (grey curve). The corresponding ALC-derived volume concentration (integrating the ALC data between 0 and 75 m) is shown by a pink curve (same right y-axis). Daily mean volume concentration values derived by OPCs and by ALC are also plotted (grey cross and pink triangle, respectively). The horizontal bar in the upper part of the figure indicates the ranges of RH measured in-situ during the analyzed cases.

The OPC-to-ALC comparison is certainly affected by intrinsic factors, as differences on the atmospheric layer sampled (at ground and integrated between 0 and 75 m, for OPC and ALC, respectively) and on the probing methods (in-situ and remote sensing, dried air sampled by OPC and ambient conditions sampled by the ALC). Furthermore, as mentioned in Section 4.2.1, a major critical issue of ALC retrievals at low levels is the correction for the overlap function, which needs to be experimentally characterized and verified for each instrument.

These issues are visible in the given example of Figure 8. In fact, in the upper panel, the agreement between the ALC-derived and the TSI-OPC aerosol V_a values is good between 0 and 7 UTC. In the following hours both instruments register an increase of the aerosol volume, although with some discrepancies in absolute values. Starting from 18 UTC, the ALC derives an aerosol volume concentration higher than the OPC one by a factor of 3-3.5. This disagreement could be related to both the presence/arrival of fine particles (<0.3 μm) not measured by the optical counter (see for example Diémoz et al., 2018a), or to aerosol hygroscopic effects (increase of volume associated to hygroscopic growth seen by ALC but not by the OPC which dries the air samples). This latter effect is confirmed by the large RH values (RH > 90%) measured after 18 UTC. The lower panel shows a good agreement between the ALC-derived and the Fidas OPC V_a values, in particular until 4 UTC and after 16 UTC. Some differences emerge around 7 UTC and between 11 and 15 UTC, where the ALC volume is lower by a factor of 2 compared to the in situ Fidas V_a values. The smaller minimum detectable size of the Fidas OPC instrument with respect to the OPS is likely the reason for the better accord between ALC and OPC V_a values in this test date. In this case, the effect of RH seems to be less important, and indeed RH values keep lower than 90%.

In general, high RH values (RH >= 90%) are known to markedly affect the aerosol mass estimation from remote sensing techniques and its relationship with 'reference' PM2.5 or PM10 measurements methods, usually performed in dried conditions (e. g. Barnaba et al., 2010; Adam et al., 2012, Li et al., 2016, Li et al., 2017). This theme is also discussed in Diemoz et al. 2018a for the ALC measurement site of Figure 8. Nevertheless, even with the mentioned

525 limitations, results in Fig. 8 well show the potential of the developed method in providing sound values of aerosol

526 volume, and hence, mass, in average-RH regimes, giving support to more standard PM10 air quality monitoring.

527 To give a further example in this direction, the model–assisted retrievals of aerosol mass over a longer time period were

528 used to derive daily-mean aerosol mass concentrations (PM10), a measurement typical of air quality stations. To this

529 purpose, for the two-months period June-July 2012, we derived daily mean values of aerosol volume at the SPC site

530 using the functional relationships $V_a = V_a(\beta_a)$, and then converted these into mass (PM10) using typical values of

531 aerosol densities (ρ_a). Results are shown in Figure 9. It compares the daily average PM10 concentration measured in

532 situ at SPC by the Italian Regional Environmental Protection Agency (ARPA, red solid curve) and the model-assisted,

533 ALC-derived daily mass concentration obtained assuming both a fixed particle density $\rho_a = 2$ g/cm^3 (blue dotted curve),

534 and a range of it between 1.5-2.5 g/cm^3 (shaded area), this range covering approximately the typical ρ_a values at the

535 SPC site. Yellow shaded areas indicate the presence of dust events (e.g. Bucci et al., 2018) that are excluded from the

536 results reported in the next paragraph.

537 More in detail, the daily-mean, ALC-derived mass concentrations were estimated in two steps: 1) estimation of hourly

538 mass values for the selected height; 2) computation of the daily values through the median of the hourly values. To

539 guarantee a good daily representativeness, the second step is applied only to those days in which at least 50% of the

540 hourly values is available in all the following temporal ranges: 00 - 05 UTC, 06 - 11 UTC, 12 - 17 UTC, 18 - 23 UTC.

541 Note that, due to the uncertainties associated to the *O(z)* in the first hundreds of meters (as previously mentioned, the

542 ALC system at SPC has an old firmware, and its overlap function is not optimally characterized), we used the 225 m

543 level as more trustworthy to estimate ALC mass concentration. On the other hand, during the considered period of the

544 year (i.e. June and July), the comparison to ground-level PM10 at SPC is expected to be only slightly affected by this

545 height difference, particularly in daytime, due to the strong convection within the mixing layer. Possible exception

546 could be in nocturnal conditions when vertical gradients in the lowermost hundreds of meters can occur. However, our

547 statistical (3-year) ALC records show the mixing layer height at SPC to descend below 250 m only 4-5 hours per day in

548 July (usually between 22 and 3 UTC, i.e., when emissions are at a minimum). Overall, Figure 9 confirms a good

549 agreement between the ALC-derived and the ARPA reference PM10 values, with a correlation coefficient (R) of 0.64.

550 In fact, mean, absolute mean and relative differences, between the two series are: $<dPM10> = 2.3 \pm 6.0$ g/cm^3,

551 $<|dPM10|> = 4.8 \pm 4.3$ g/cm^3 and $<(dPM10/PM10)> = 0.14 \pm 0.27$. This agreement attests that SPC site can indeed be

552 considered an 'average' continental site and suggests the potential of this approach to derive information on aerosol

553 volume and mass. Still, due to the specificity of each site and to the limited period considered here, these results cannot

554 be taken as representative of all continental sites at all times. Further studies at different places and over longer time

555 periods would be necessary to better assess the uncertainty of the proposed retrieval, including uncertainties due to the

556 variability of 'continental' conditions (in terms of particle size distribution, compositions, hygroscopic effects, etc…),

557 but also of the instrument-dependent performances (e.g. overlap corrections, etc…).

558 **4   Summary and Discussion**

559 Thanks to their low construction/operation costs and to their capability at providing continuous, unattended

560 measurements, the use of automated-lidar-ceilometers (ALCs) for aerosol characterization has increased in the recent

561 years. Several numerical approaches were recently proposed to estimate the aerosol vertical profile either using

562 ceilometer measurement only, or coupling these with ancillary measurements (e.g., Stachlewska et al., 2010; Flentje et

563 al., 2010; Wiegner et al., 2012; Wiegner et al., 2014; Cazorla et al., 2017, Román et al., 2018).

564   This work proposes a methodology to retrieve key aerosol properties (as extinction coefficient, surface area and

565   volume, thus mass) from lidar/ALC measurements using in support the results from a specifically developed aerosol

566   numerical model to drive the retrievals. In particular, the numerical model uses a "Monte-Carlo" approach to simulate a

567   large set (20000) of aerosol microphysical properties intended to reproduce the variability of 'average' (clean-to-

568   moderately polluted) continental conditions, i.e., those expected to dominate over Europe. Based on the assumption of

569   particle sphericity, relevant computations of aerosol physical (surface area and volume, S_a and V_a) and optical

570   (backscattering and extinction coefficients, $ß_a$ and α_a through Mie scattering theory) properties were performed at three

571   commonly used lidar wavelengths (i.e., at the Nd:YAG laser harmonics 355, 532, 1064 nm). Fitting procedures of this

572   large set (20,000) of $ß_a$ vs. α_a, S_a and V_a data-pairs were then used to derive mean functional relationships linking $ß_a$ to

573   α_a, S_a and V_a, respectively. The model's statistical uncertainties (i. e., those related to the variability of the

574   microphysical parameters used in input to the computations of the bulk physical/optical properties) associated to these

575   so-derived mean relationships were found to be within 30% and 40% for $ß_a$ vs α_a and $ß_a$ vs V_a, respectively, while $ß_a$ vs

576   S_a exhibits a larger dispersion (relative standard uncertainty of 40%-70%, depending on $ß_a$). It is worth mentioning that

577   these are higher than those associated to the retrievals of aerosol bulk parameters using the complete set of Raman

578   lidars observations (three aerosol backscattering and two extinction coefficients, i.e., the so called 3+2 approach),

579   assuming, as in our case, no random uncertainty in the lidar input data. For example, Veseloski et al. (2012) found a

580   maximum uncertainty of 15% for particle volume and surface area estimation, in the case of 0% random uncertainty in

581   the lidar input data. Note however, that such multi-wavelength lidar systems are 10 to 20 times more expensive than

582   ALC systems, need to be operated by highly trained operators, and are rarely run all day round.

583   The model results also allowed exploring the expected dependence of the (continental aerosol) lidar ratio (LR) on $ß_a$ at

584   355, 532 and 1064 nm, and in turn, the mean, 'weighted'-LR value at each wavelength (found to be 50.1 ± 17.9 sr, 49.6

585   ± 16.0 sr and 37.7 ± 12.6 sr, at 355, 532 and 1064 nm respectively). Availability in literature of LR values at 1064 nm

586   are scarce and its monotonic increase with $ß_a$ found in this work (Figure 3) suggests that the use of a fixed LR value for

587   the inversion of ALC signals should be done with caution and carefully evaluated case by case. A similar, non-

588   monotonic behavior characterizes the shapes of LR vs $ß_a$ curve at 355 and 532 nm.

589   We tested the reliability of our model results in two ways: 1) the model numerical computations were compared to

590   'real' lidar measurements (specifically selected within the EARLINET database), and 2) the model-assisted retrievals of

591   aerosol optical (AOT) and physical (V_a, PM10) properties by real, operational ALC systems were compared to

592   corresponding 'reference' measurements performed by co-located, independent instrumentation.

593   In particular, in task 1) our simulations were compared to backscatter and extinction coefficients at 532 and 355 nm

594   independently retrieved by advanced Raman lidar systems operating at different EARLINET sites in Europe (namely

595   Hamburg and Leipzig in Germany, Madrid in Spain, Lecce and Potenza in Italy). The model simulations were found to

596   statistically well match the observations (Figures 4, 5 and C1). Mean discrepancies between model and measurement-

597   based LR were found to be lower than 20%, suggesting a good capability of the assumed aerosol model (and admitted

598   range of variability) to represent 'real', 'average continental' aerosol conditions in different sites across Europe. Some

599   differences emerged for so southern Italy EARLINET sites, possibly affected by the influence of marine aerosols,

600   leading to lower LR values for high values of $ß_a$.

601   For task 2) we applied the proposed model-based inversion to different ALC systems (Lufft CHM-15k), part of the

602   Italian ALICENET network. We firstly tested the ability of the proposed approach to derive aerosol extinction by

603   comparing hourly-mean, vertically-integrated α_a (i.e., hourly mean AOT) derived by three ALC systems to

604   corresponding AOT measurements from co-located sun-photometers (ALICENET sites of Aosta San Cristophe (ASC),

605   San Pietro Capofiume (SPC) and Rome Tor Vergata (RTV), Figure 7). ALC-sun photometer agreement was found to be
606   within 30%. Tests on the use of fixed LR were also performed to investigate the advantage of the proposed approach
607   with respect to more standard ones. To this purpose, we used the (1064 nm) fixed-LR value suggested by the E-Profile
608   EUMETSAT Program and the 'weighted mean' derived from our model (52 sr and 38 sr, respectively). While for the
609   ASC site negligible differences were found among the three retrieval schemes, for both SPC and RTV sites the best
610   ALC – sun photometer accordance in AOT is reached when employing the model-based or the fixed LR=38 sr
611   inversion schemes, with a mean error around 16-19 % and 31-33 % for SPC and RTV, respectively. Applying the fixed
612   LR value of 52 sr produces an overestimation of the AOTs, with a mean relative bias equal to 33 % and 44 % at SPC
613   and RTV, respectively. This suggests that, at 1064 nm, the LR value for continental aerosol is lower than the one
614   assumed by the E-Profile procedure and, more in general, this highlights the advantage of a procedure not requiring an
615   a-priori, and to some extent arbitrary, choice of the LR value.

616   As a second test in task 2, values of aerosol volume (and mass) derived using the model-assisted ALC retrieval were
617   compared to in situ aerosol measurements performed by OPCs and PM10 analyzers. A continuous, two-months
618   comparison (June – July 2012) between daily average aerosol mass concentration as measured in situ and derived by
619   ALC (in the lowest altitudes) at SPC, showed a mean relative difference of around 15% (Figure 9).

620   Overall, the good results obtained in our validation efforts are encouraging but necessarily related to the specific
621   conditions at the measuring sites considered and to the characteristics of the instruments employed. They are therefore
622   not necessarily representative of results obtainable in all European continental sites, and at all times. Further tests using
623   wider datasets covering a variety of sites and ALC instrumentation would be desirable to better understand potential
624   and limits of the applicability of the proposed method over the larger scale. An obvious intrinsic limitation is that the
625   method is dependent on the considered aerosol type which in this study was tuned to reproduce average continental
626   aerosol conditions. Errors associated to the application of the derived functional relationship might be larger if more
627   'specific' aerosol conditions (e.g. contamination by sea salt or desert dust particles) affect a given site. In the future, the
628   information coming from ALC systems with an additional depolarization channel (as tested in the DIAPASON Project,
629   Gobbi et al., 2018) could be used to force the retrieval to different model schemes (e.g. switching from 'no dust' to
630   'dust' schemes conditions) in the same vertical profile. This will enhance the capabilities of ALCs to operatively
631   estimate and characterize the aerosol optical properties (e.g. Gasteiger and Freudenthaler, 2014).

632   Additionally, although our validation exercises returned results well within the uncertainties related to the model
633   statistical variability alone (i.e., the relative errors associated to the mean functional relationships), the expected total
634   uncertainty to be associated to the method should include terms that have not been specifically addressed in this work,
635   as for example the instrumental error itself.
636   On the other hand, the proposed approach has the main advantage of allowing the operational (i.e. 24/7) retrieval of
637   fairly reliable, remote sensing profiles of aerosol optical (β_a, α_a) and physical (S_a, V_a) properties (with associated
638   uncertainties and limitations) by means of relatively simple and robust instruments. This could temporally and spatially
639   complement the information coming from more advanced lidar networks (for example, the Raman channel of multi-
640   wavelength system cannot be used in daylight conditions) and, more in general, could represent a valid option to
641   deliver, in quasi real time, the 3D aerosol fields useful for operational air quality (e.g. integration of the in situ surface
642   measurements) and for meteorological and climate monitoring (e.g. aerosol-cloud interaction and aerosol transport and
643   dispersion processes).
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Data availability

AERONET Rome-Tor Vergata sun photometer AOT data were downloaded from the AERONET web page (AERONET, 2018). SKYNET sun photometer AOT data were downloaded from the SKYNET webpage (SKYNET, 2018). EARLINET backscattering and extinction coefficients were downloaded from the EARLINET webpage (EARLINET, 2018). ALICENET ALC raw data are available upon request at alicenet@isac.cnr.it.

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948 **Table 1. Aerosol parameter values as reported in literature for continental-type aerosols.**

Reference	r_1 (µm) σ_1	r_2 (µm) σ_2	r_3 (µm) σ_3	N_1/N_{tot} (%)	N_2/N_{tot} (%)	N_3/N_{tot} (%)	m_{r_1}, m_{im_1}	m_{r_2} m_{im_2}	m_{r_3} m_{im_3}	N_{tot} (cm^{-3})	Aerosol type
Whitby (1978)[1]	0.008 1.6	0.034 2.1	0.46 2.2	0.56	0.44	4×10^{-4}	-	-	-	1800	Clean continental
D'Almeida et al. (1991)[2]	0.012 2.0	0.029 2.24	0.471 2.51	0.06	0.94	2×10^{-6}	1.75 0.44	1.53 0.012	1.53 0.008	20000	Average continental
Hess et al. (1998)[2]	0.012 2.0	0.021 2.24	0.471 2.51	0.56	0.44	0.3×10^{-4}	1.75 0.44	1.53 0.012	1.53 0.008	15300	Average continental
Barnaba and Gobbi (2004a)[1]	0.007- 0.012 1.7 – 2.0	0.021 – 0.077 2.03 – 2.24	0.403 – 0.5 2.11 – 2.24	6.1– 54.2	45.8 – 93.9	$(2 -26.1) \times 10^{-4}$	1.25–2.00 0.07–1.00	1.53 6×10^{-3}	1.53 8×10^{-3}	$10^3 - 10^4$	
Omar et al. (2009)[1]	-	0.093-0.10 1.53-1.61	0.68-0.76 1.9-2.1	-	0.999-1	$(0.02-3) \times 10^{-4}$	-	1.38-1.40 $(0.1 -6.3) \times 10^{-3}$	1.40-1.46 $(3.4-6.3) \times 10^{-3}$	-	Clean and polluted continental
Levy et al. (2007)[2]	0.018 2.0	0.005 2.97	0.5 2.97	1	1×10^{-7}	1×10^{-13}	1.75 0.44	1.53 6×10^{-3}	1.53 8×10^{-3}	-	
Barnaba et al. (2007)[1]	-	0.05-0.1 1.35-1.70	0.4-0.5 1.5-2.0	-	0.98-0.99	0.01-0.02	-	1.35-1.55 $(2.5 -20) \times 10^{-3}$	1.53-1.6 $(1.0 -80) \times 10^{-4}$	$1\text{-}3 \times 10^3$	Continental - coastal
Amiridis et al. (2015)[1]	-	0.03-0.9 1.6-2.2	0.47-0.69 1.9-2.5	-	1	$(4 – 8) \times 10^{-7}$	-	1.42-1.45 $(2.3 –6) \times 10^{-3}$	1.45-1.53 $(2.3 –6) \times 10^{-3}$	-	Clean and polluted continental

949 [1]The refractive index is at λ=532 nm.

950 [2]The refractive index is at λ=550 nm.

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953 **Table 2. Variability ranges used in this study. Values refer to ground and dry conditions (see text for details).**

Parameter	Mode I	Mode II	Mode III
r_i (μm)	0.005 - 0.03	0.03 - 0.1	0.3 - 0.5
σ	1.35 – 1.7	1.35 – 1.7	1.5 – 2.4
N_i/N_{tot} (%)	10 - 60	40 - 90	0.01 – 0.5
m_{r_i} (355 nm)	1.40 – 1.80	1.40 – 1.70	1.50 – 1.60
(532 nm)	1.40 – 1.80	1.40 – 1.70	1.50 – 1.60
(1064 nm)	1.42 – 1.82	1.37 – 1.66	1.50 – 1.60
m_{im_i} (355 nm)	$1\times10^{-2} - 0.47$	$1\times10^{-4} - 0.010$	$1\times10^{-4} - 0.02$
(532 nm)	$9\times10^{-3} - 0.44$	$1.2\times10^{-4} - 0.012$	$1\times10^{-4} - 0.01$
(1064 nm)	$9\times10^{-3} - 0.44$	$1.5\times10^{-4} - 0.015$	$1\times10^{-4} - 0.005$
N_{tot} (cm^{-3})	500 – 10000		

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956   **Table 3. Parameters of the Seventh-Order Polynomial Fits ($y=a_0+a_1x+a_2x^2+a_3x^3+a_4x^4+a_5x^5+a_6x^6+a_7x^7$) for λ = 1064 nm, with**

957   **$x=\log(\beta_a)$ (in km^{-1} sr^{-1}) and $y=\log(\alpha_a, S_a, or V_a)$ in (km^{-1}, cm^2/cm^3 and cm^3/cm^3, respectively).**

Functional relantionship at 1064 nm	Extinction coefficient	Surface area	Volume
a_0	3.797837507651898	12.019452592845141	-5.314834128998254
a_1	3.294032541389781	30.825966279368547	2.500484347793244
a_2	0.962603336867675	24.518531616019207	-1.196109537503000
a_3	0.241796629870675	10.625241994796593	-1.583236058579546
a_4	0.064609145804688	2.634051072085453	-0.681801883947768
a_5	0.017721752150233	0.373150843707711	-0.145232662646142
a_6	0.002722551625862	0.027971628176431	-0.015471229968392
a_7	0.000157245409783	0.000854381337164	-0.000658925756875

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960 **Table 4. Mean weighted LR at 355, 532 and 532 nm derived in this work and comparison to the corresponding aerosol**
961 **subtypes (clean continental, CC, and polluted continental, PC) from relevant literature.**

LR (sr)	*λ=355 nm*	*λ=532 nm*	*λ=1064 nm*
Omar et al., (2009) (Calipso aerosol model)	-	70 ± 25 (PC)	30 (PC)
	-	35 ± 16 (CC)	30 (CC)
Amiridis et al. (2015) (LIVAS database)	59.5[*] (PC)	64 (PC)	-
	56.5[*] (CC)	54 (CC)	-
Papagiannopoulos et al.. (2016) (EARLINET measurements)	-	62 ±10 (PC)	-
	-	47 ± 4 (CC)	-
Düsing et al.. (2018) (in-situ and lidar measurements)	55	55	30; 15[**]
This work	50.1 ± 17.9	49.6 ± 16.0	37.7 ± 12.6

962 * derived using the extinction-related and backscatter-related Ångström exponents given by Amiridis et al. (2013)

963 ** see the explanation in the text for the two different values

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965 **Table 5. Extinction-to-volume conversion factors, $c_v = V_a/\alpha_a$ (and corresponding 'mass-to-extinctions efficiency' values, MEE**
966 **$= \alpha_a/(V_a * \rho_a)$, given assuming $\rho_a = 2$ g/cm^3) of continental particles as derived from our model at different wavelengths**

Reference	c_v [10^{-6}m] (corresponding MEE, [m^2g^{-1}])			Notes
Wavelength [nm]	355	532	1064	
Hess et al. (1998)	-	0.35 (1.43)	-	OPAC, clean continental model
Hess et al. (1998)	-	0.28 (1.79)	-	Opac, polluted continental model
Barnaba and Gobbi (2004b)	-	0.18 (2.78)	-	Continental model
Ansmann et al. (2011b)	-	0.18 (2.78)	-	Germany, fine aerosol fraction
Lewandosky et al. (2010)	-	-	0.77 – 2 (0.25-0.65)	Mexico city basin
Sicard et al. (2012)	-	0.26 (1.92)	-	AERONET, Spain
Mamouri and Ansmann (2017)	0.17 (2.94)	0.30 (1.67)	0.96 (0.52)	Germany, continental anthropogenic pollution
Mamouri and Ansmann (2017)	0.23 (2.17)	0.41 (1.22)	1.41 (0.35)	Cyprus, continental anthropogenic pollution
Mamali et al. (2018)		0.14, 0.24 (3.57, 2.03)		Cyprus, fine non-dust aerosol fraction
This work	0.12 (4.17)	0.19 (2.63)	0.60 (0.83)	Continental (clean to moderately polluted)

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969 **Table 6. Main characteristics of the dataset of the EARLINET continental sites considered in this study. The listed dataset**
970 **refers to the data downloaded from the EARLINET site (last access on the 11th of January 2018).**

Station	Number of points at 355 and at 532 nm)	Altitude range (Δz, in km)	Period
Lecce (LC) 40.33 N, 18.10 E, 30 m a.s.l.	1012 – 109	1 – 4	Aug2007 – Oct2013
Leipzig (LE) 51.35 N, 12.43 E, 90 m a.s.l.,	5186 – 4549	1.5 – 4	Aug2008 – Sept2016
Potenza (PO) 40.6 N, 15.72 E, 760 m a.s.l.	1244 – 219	1.5 – 4	May2000 – Aug2009
Hamburg (HH) 53.57 N, 9.97 E, 25 m a.s.l.	243 – n.a.	0.5 – 4	Apr2001 – Oct2002
Madrid (MA) 40.45 N, 3.73 E, 669 m a.s.l.	n.a. – 492	0.5 – 4	Jun2006 – Jun2008

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973   **Table 7 Mean LR discrepancies between our model results and EARLINET measurements and weighted LR at 355 and 532**
974   **nm for the considered EARLINET stations.**

	$[(LR_{mod} - LR_{meas})/LR_{meas}]*100$		EARLINET weighted LR (sr)	
Station	$\lambda=355\ nm$	$\lambda=532\ nm$	$\lambda=355\ nm$	$\lambda=532\ nm$
LC	10	18	51.8	44.5
LE	6	9	52.6	51.0
PO	17	7	44.9	57.2
HH	5	-	53.3	-
MA	-	6	-	54.2

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977 **Table 8. Main characteristics of the ALC and co-located sun-sky radiometer equipment located at the considered**
978 **ALICENET sites.**

	Site type	ALC model	ALC firmware	Sun photometer model
ASC	alpine	Nimbus CHM150104	0.743	POM-02
SPC	rural	Nimbus CHM110115	0.556	POM-02L
RTV	semi-rural	Nimbus CHM070052	0.720	CIMEL CE-318

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981 **Table 9. Results of the comparison between the AOT measured by sun-photometers and the one derived by ALCs (model-**
982 **based and fixed LR inversion schemes) at three ALICENET stations. Mean differences (expressed in terms of <dAOT> =**
983 **<(AOT$_{ceil}$-AOT$_{phot}$)>, <|dAOT|> (module), <dAOT/AOT> and <|dAOTI/AOT|>) are reported with their standard deviations.**

ALICENET sites	<dAOT>	<|dAOT|>	<dAOT/AOT>	<|dAOTI/AOT|>
ASC				
Variable LR from our model	-0.004 ± 0.015	0.010 ± 0.013	-0.25 ± 0.57	0.31 ± 0.35
	0.002 ± 0.021	0.009 ± 0.015	0.31 ± 0.58	0.33 ± 0.35
LR = 52 sr				
LR = 38 sr	-0.004 ± 0.014	0.009 ± 0.012	-0.23 ± 0.43	0.30 ± 0.32
SPC				
Variable LR from our model	-0.001 ± 0.020	0.013 ± 0.016	-0.005 ± 0.28	0.19 ± 0.20
	0.021 ± 0.026	0.026 ± 0.02	0.33 ± 0.35	0.38 ± 0.26
LR= 52 sr				
LR= 38 sr	-0.003 ± 0.019	0.011 ± 0.014	-0.043 ± 0.24	0.16 ± 0.18
RTV				
Variable LR from our model	0.004 ± 0.020	0.014 ± 0.014	0.11 ± 0.49	0.33 ± 0.30
	0.016 ± 0.023	0.021 ± 0.018	0.44 ± 0.59	0.49 ± 0.45
LR= 52 sr				
LR= 38 sr	0.003 ± 0.019	0.013 ± 0.013	0.088 ± 0.460	0.31 ± 0.27

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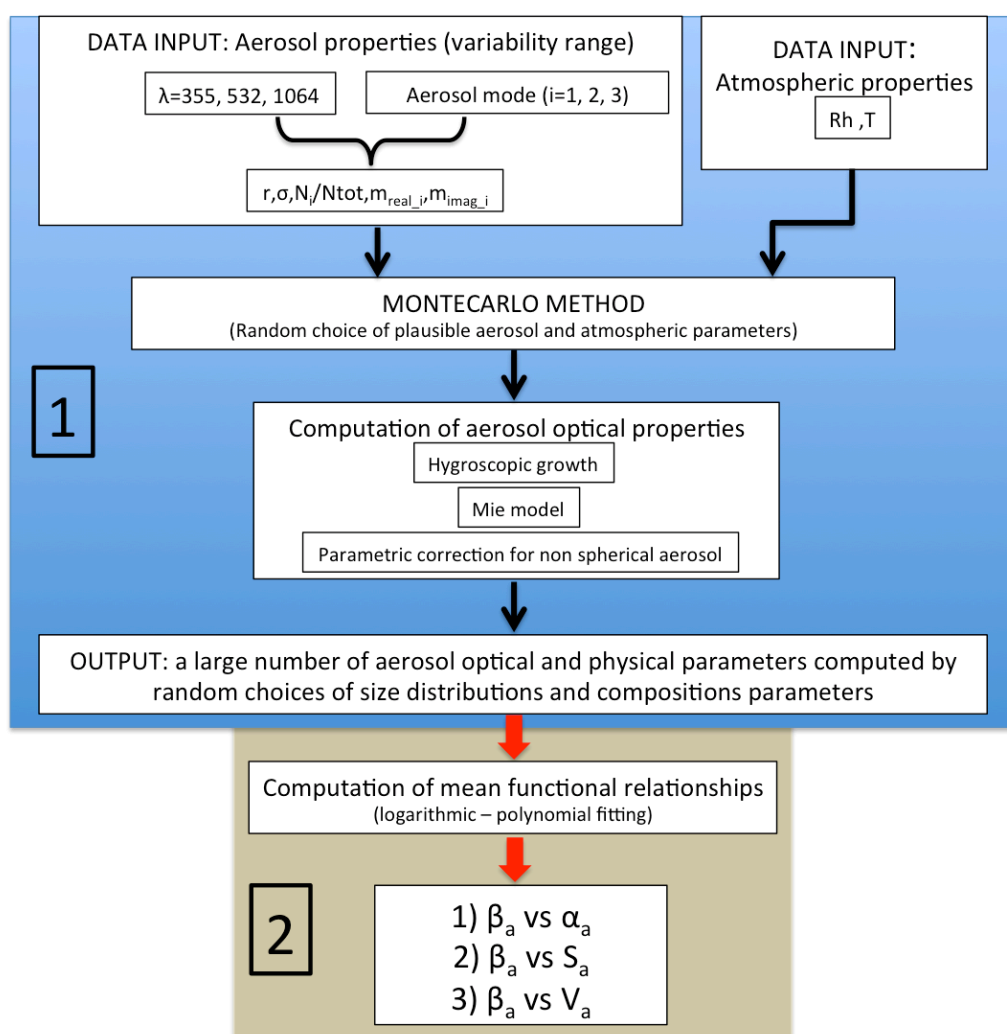
Figure 1. Schematic of the two-step model structure developed to obtain, as a result, functional relationships between the aerosol backscatter (β_a) and the aerosol extinction, surface area and volume (α_a, S_a and V_a, respectively).

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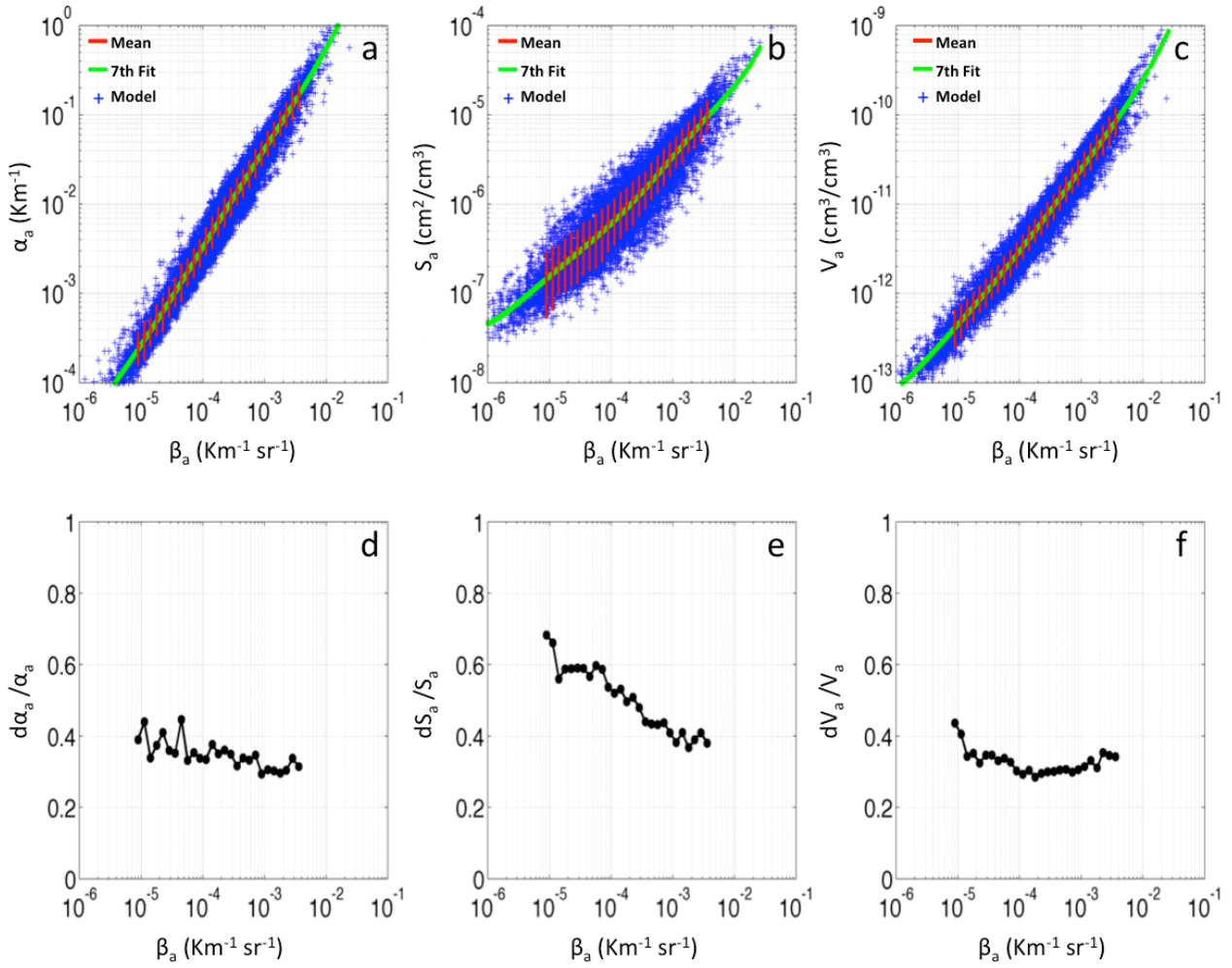
Figure 2. Scatterplots of a) α_a (km^{-1}), b) S_a (cm^2/cm^3) and c) V_a (cm^3/cm^3) vs backscatter β_a (km^{-1} sr^{-1}) and relevant relative errors (panels d, e, f, respectively) as derived from 20000 model computations (blue points) at λ = 1064 nm. Red dots and error bars are the average values per decade of β and their standard deviations, green lines are the 7th-order polynomial fit curve of the 20000 points.

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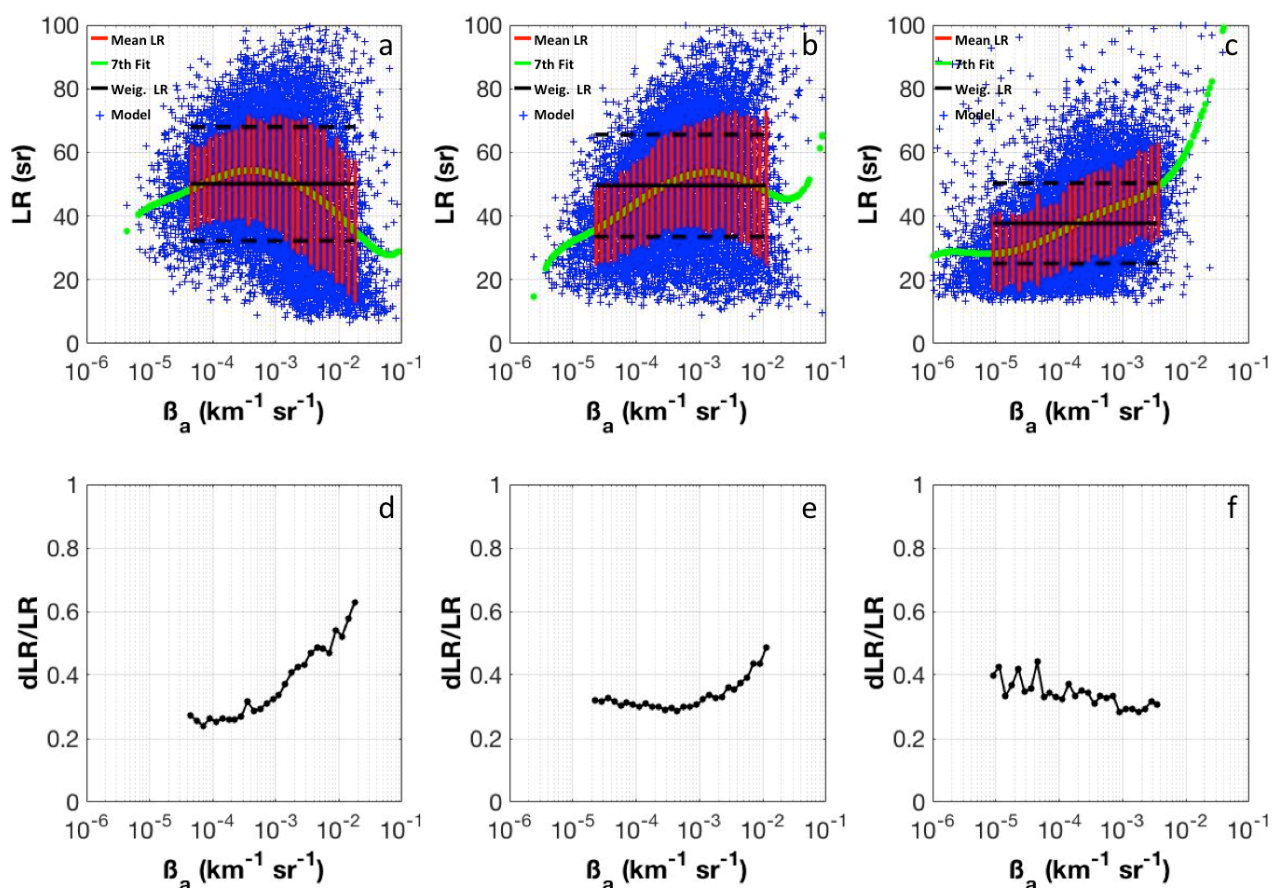
Figure 3. Upper plots: scatterplots of LR (sr) versus β_a (km^{-1} sr^{-1}) at: a) 355 nm; b) 532 nm; c) 1064 nm (blue points). The 7th-order polynomial fit curve (green lines) and the average values per decade of ß together with their standard deviations (red points and red vertical bars, respectively) are also reported. Horizontal black lines are mean values of the 'weighted-LR' and ± 1 s. d. (solid and dotted lines, respectively). Lower plots: relative errors associated with the model-derived LR at d) 355 nm; e) 532 nm; f) 1064 nm.

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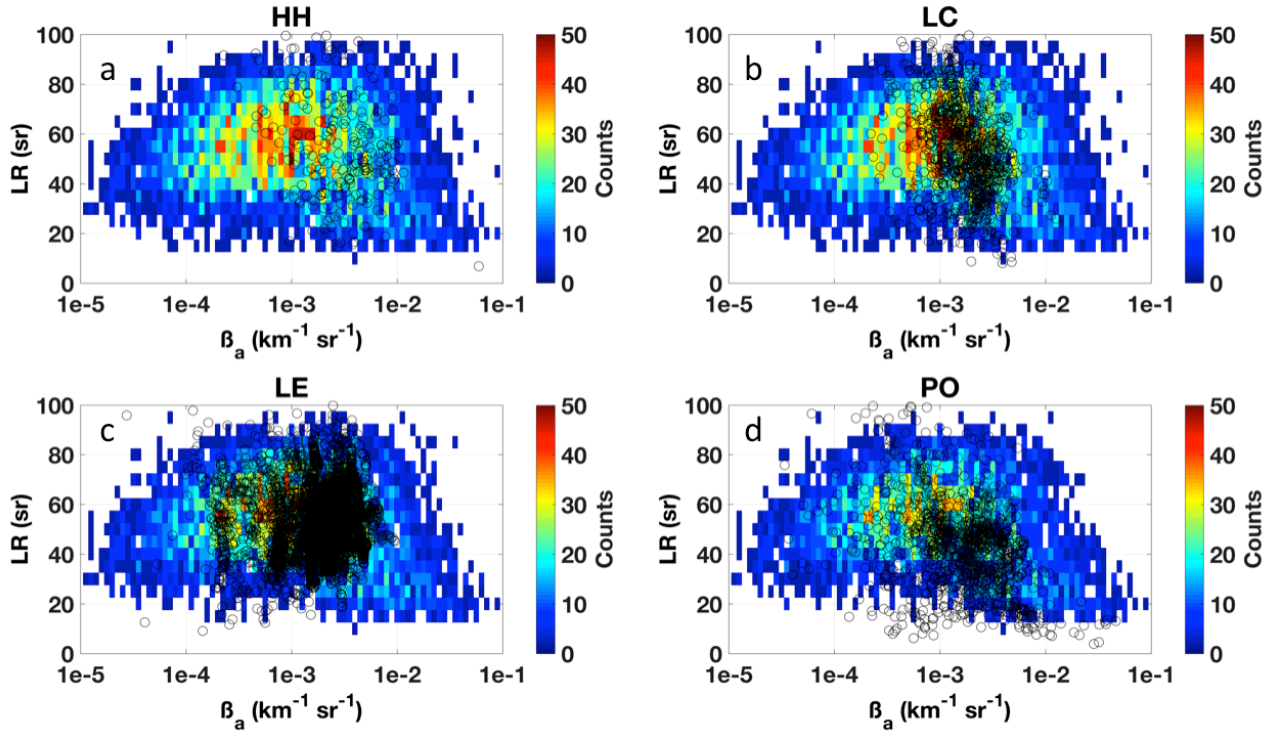
Figure 4. Scatterplots of LR (sr) versus β_a (km^{-1} sr^{-1}) at 355 nm as simulated by our model (colored region) and measured by EARLINET lidars (black open circles) in Hamburg (Germany) (a), Lecce (Italy) (b), Leipzig (Germany) c) and Potenza (Italy) (d). The color area is the region of simulated values, the color code indicating the number of simulated values in each β_a -LR pair (see legend). In particular, the color-2D histogram is computed using a semi-logarithmic box consisting of 10 equally spaced bins per decade of β_a in the x-axis and 5 spaced LR values in the y-axis.

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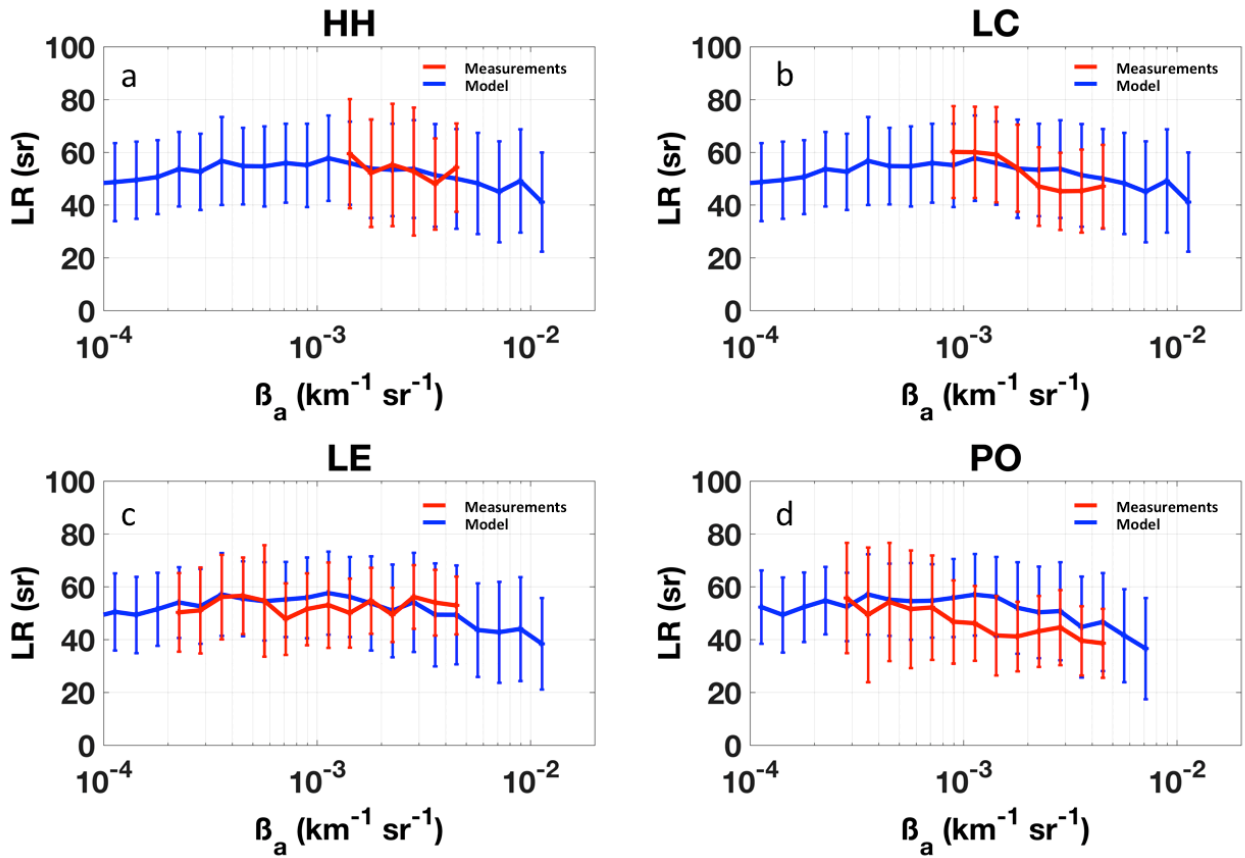
Figure 5. Model-simulated (blue) and lidar measured (red) LR vs $ß_a$ mean curves at 355 nm calculated per 10 equally spaced bins per decade of $ß_a$ in a) Hamburg, b) Lecce, c) Leipzig, and d) Potenza EARLINET lidar station. Vertical bars are the associated standard deviations.
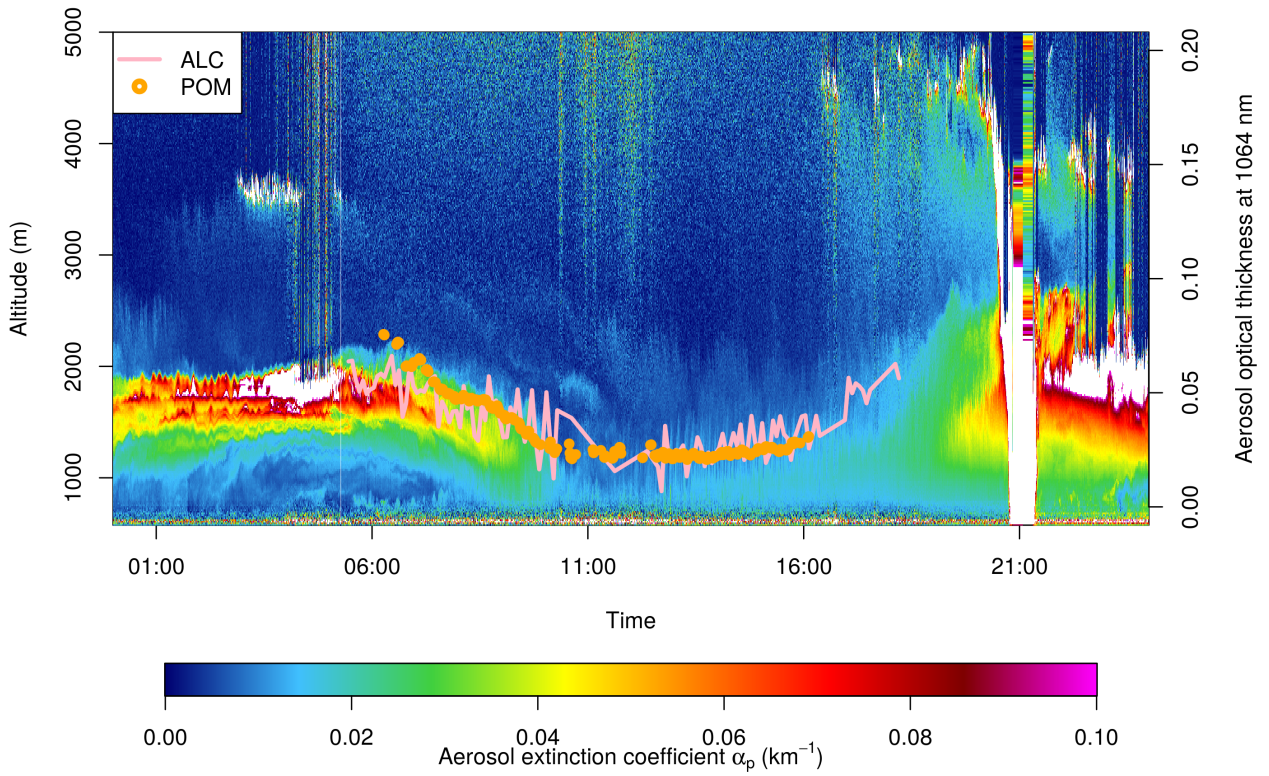
Figure 6. Time-height cross-section of the aerosol extinction coefficients α_a [km^{-1}], as derived at 1064 nm on 26 June 2016 by the ALICENET ALC of Aosta San Christophe (Northern Italy). The orange circle points and the pink line are the AOT values (right y-axis, panel b) measured by a co-located POM-02L radiometer and estimated from the ALC following our approach.

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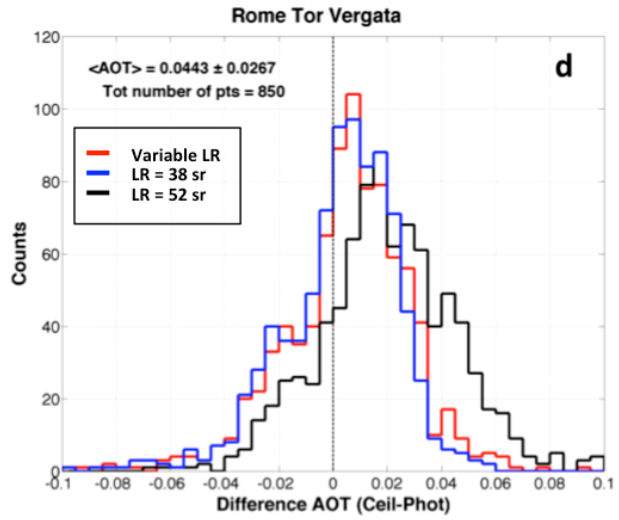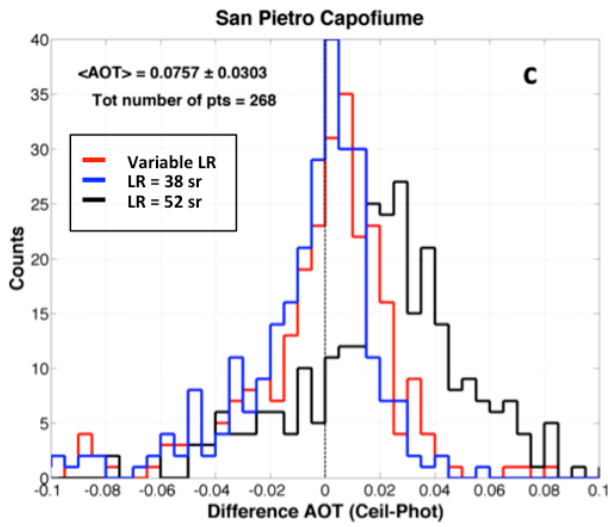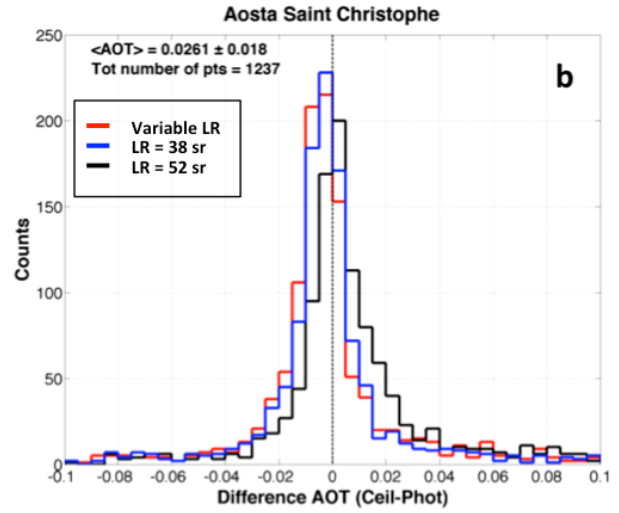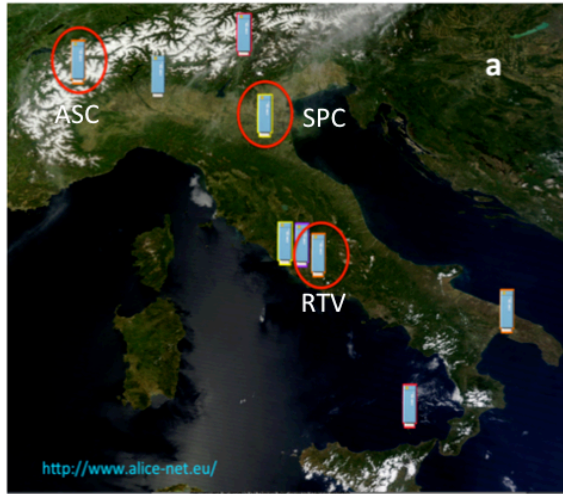
Figure 7. a) geographical map of the ALC network ALICENET. The red circles highlight the selected sites for this study: Aosta San Christophe (ASC), San Pietro Capofiume (SPC) and Rome Tor Vergata (ASC). b-d) histograms of the differences between the hourly-mean coincident AOTs at 1064 nm as derived by ALCs and measured by photometers, at ASC, SPC and RTV, respectively. The different colors (red, blue and black) depict the different inversion schemes: model-based inversion scheme, LR = 38 sr and LR = 52 sr, respectively. In each panel the values of the average measured AOT (and its associated standard deviation) and of the number of considered pairs are also reported.

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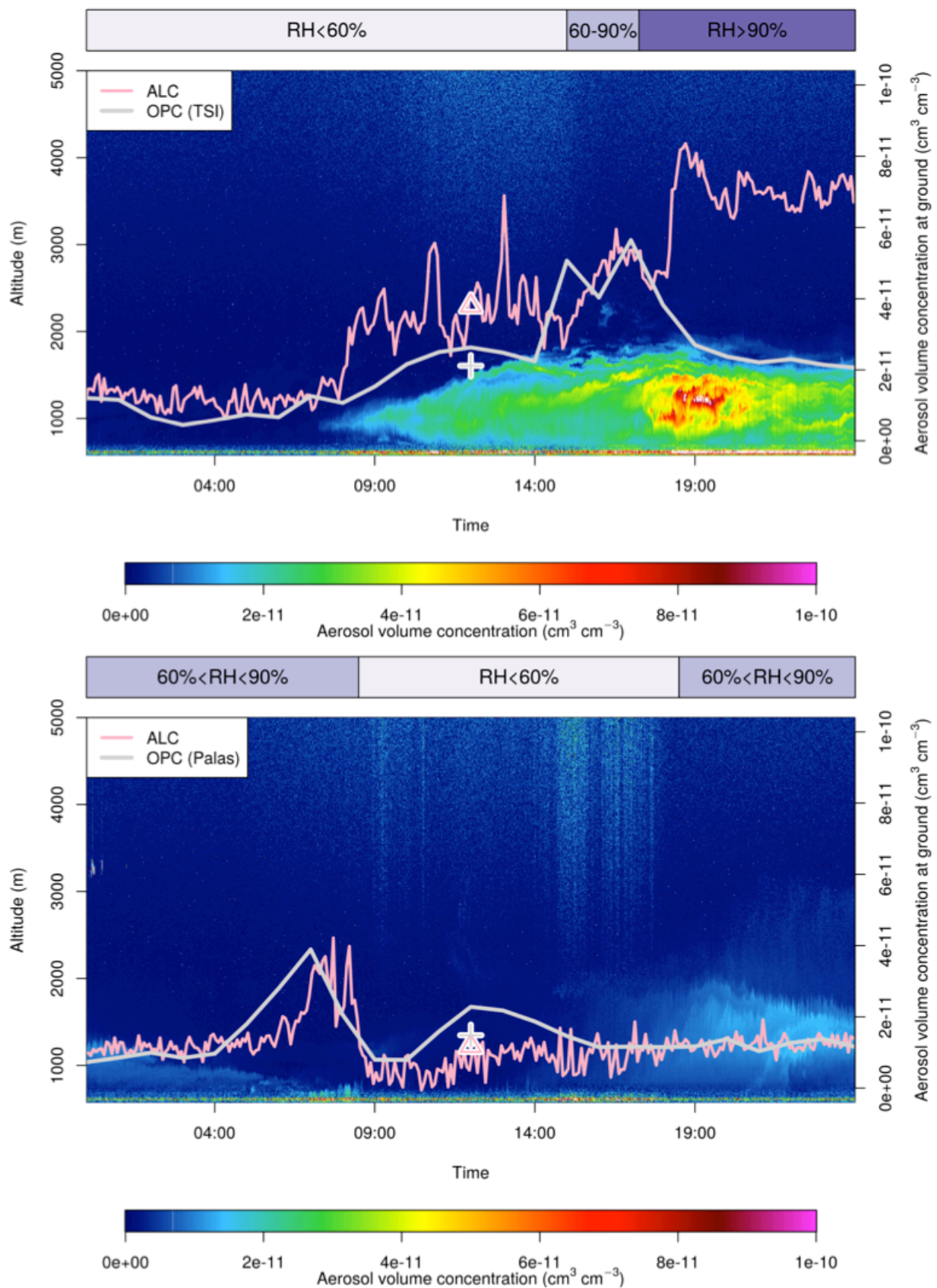
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Figure 8. Time-height cross-section of the aerosol volume concentration at Aosta San Christophe for 29 December 2016 (upper panel) and 05 September 2017 (lower panel). The right y-axis reports the volume concentration measured at surface through TSI and Fidas®200s OPCs (upper and lower panels, grey curves) and the ALC-derived volume concentration at 75

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1035 **m (pink curves). The grey crosses and the pink triangles refer to the daily mean aerosol volume value derived by OPCs and**

1036 **ALC measurements, respectively. The horizontal bars in the upper part of the panelse indicate the ranges (RH<60%,**

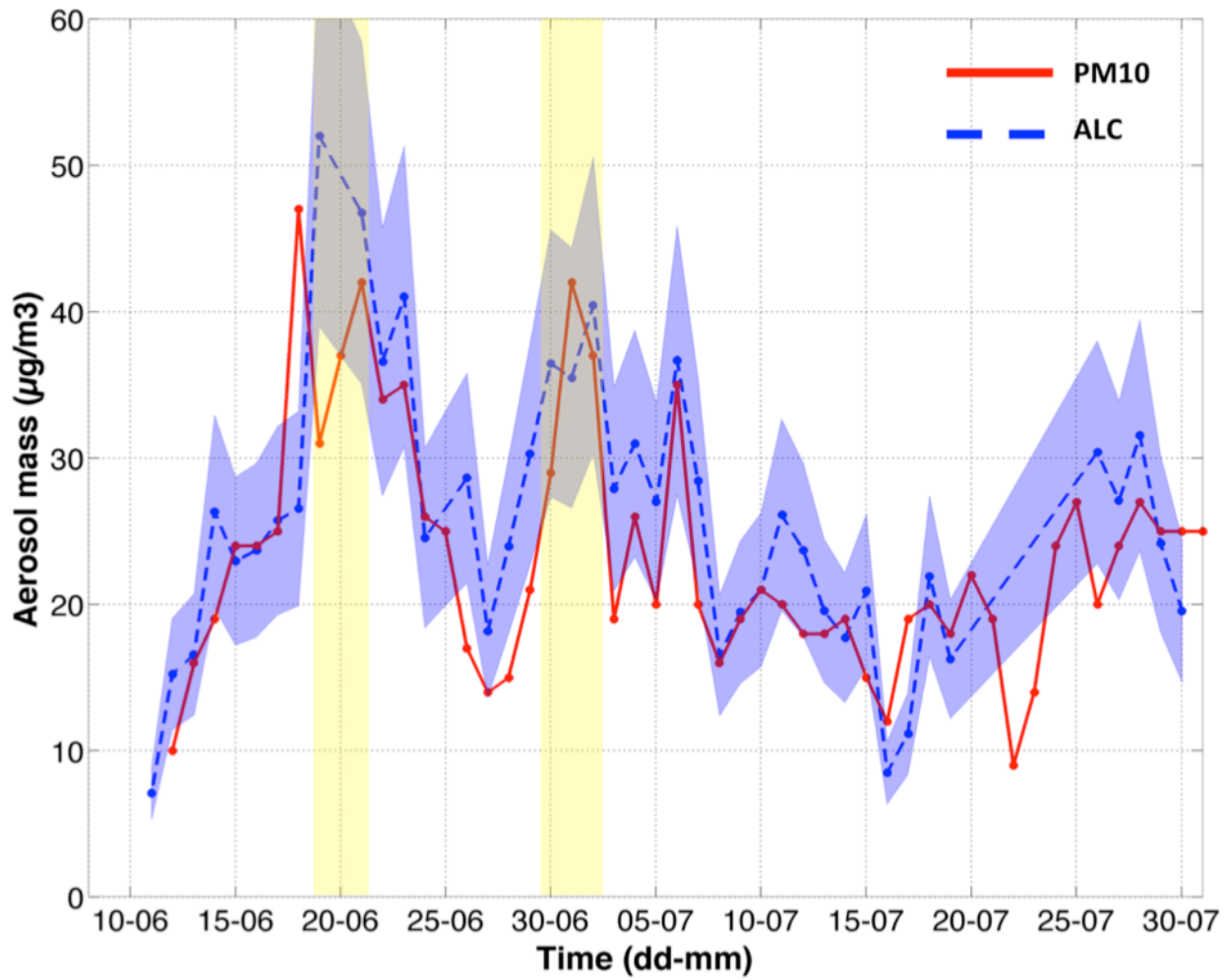1037 **60%<RH<90% and RH>90%, respectively) of the measured in-situ RH during the analyzed days.**

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1040 **Figure 9. Daily-resolved aerosol mass concentration at SPC, for the period June – July 2012, estimated from ALC-derived**
1041 **aerosol volume data at 225 m a.s.l. converted into mass using a fixed particle density ρ_a = 2 µg/m^3 (blue dotted line) and a**
1042 **variable ρ_a between 1.5 -2.5 µg/m^3 (shaded blue area). The red solid line is the daily PM10 concentration as measured by the**
1043 **local Air Quality agency (ARPA). Vertical yellow shaded stripes indicate the presence of dust events.**

Appendix A: Model sensitivity tests

1045 To evaluate the proposed continental model configuration (hereafter CM0) and discuss its sensitivity to the variability
1046 of the employed parameters, an overview of the impact on the model results produced by changing the limit of the
1047 variability ranges of these parameters (i.e. using different model configuration, CMX) is given in this section.

1048 The varied model (CMX-CM0) mean difference on the considered optical property (OP) has been quantified through
1049 the following equation:

1050 $$< \frac{dOP}{OP} > = \left(\frac{1}{Nbin}\right) \cdot \sum_{i=1}^{Nbin}[(< OP_{CMX} >_i - < OP_{CM0} >_i)/< OP_{CM0} >_i], \qquad (A.1)$$

1051 where N_{bin} is the total number of defined bins of β_a.

1052 The results of the mean differences of α_a and LR for different ranges of β_a and for the whole β_a interval are reported on
1053 table A1, where relevant sensitivity cases (i.e. relative mean difference greater than 1%) at λ=355 nm have been taken
1054 into account.

1055 CM1 refers to a model configuration without the first aerosol mode (N_1%=0). The overall decrease on the values of α_a
1056 and LR (around 3-4%) is due to the sum of significant and opposite effects for low and high values of β_a where
1057 $<d\alpha_a/\alpha_a>$ and $<dLR/LR>$ are of the order of -6% and 8%, respectively. Removing the coarser aerosol mode (N_3%=0),
1058 causes positive mean values for $<d\alpha_a/\alpha_a>$ and $<dLR/LR>$ of the order of 5% (sensitivity case CM2). In this case, the
1059 largest impact is observed for the β_a range between 2×10^{-4} and 2×10^{-3} km^{-1}sr^{-1}.

1060 An opposite result is obtained by decreasing the upper bound of the r_2 variability range (r_2=0.03 – 0.05 µm, CM3). In
1061 fact also this model configuration leads to lower α_a and LR ($<d\alpha_a/\alpha_a>$ and $<dLR/LR>$ are equal to -6%, approximately).
1062 In this case, the variation on the r_2 parameter affects the higher ranges of β_a (β_a=2×10^{-4}-2×10^{-2} km^{-1}sr^{-1}). Higher modal
1063 radii for the coarse-mode particle (r_3=1 – 1.2 µm) in CM4 configuration leads to the increase of the contribution of
1064 model-generated points with higher β_a and causes lower values of α_a and LR ($<d\alpha_a/\alpha_a>$ and $<dLR/LR>$ are equal to -5%,
1065 approximately) only for high values of β_a (β_a=2×10^{-3}-2×10^{-2} km^{-1}sr^{-1}), whereas the effect over the whole β_a range is
1066 around -1%.

1067 The CM5 configuration accounts for the presence of more absorbing particles in the first aerosol mode, where the lower
1068 bound of m_{1im} has been increased by a factor of 10 (m_{1im} =0.1-0.47). This produces a significant effect only for the
1069 lower values of β_a (β_a=2×10^{-5}-2×10^{-4} km^{-1}sr^{-1}), with an increase of α_a and LR of approximately 4%. On the contrary,
1070 increasing the lower bound of the real part of the second aerosol mode refractive index (m_{2r} =1.55-1.70) has a large
1071 impact on the considered parameters. In fact, the CM6 configuration largely underestimates both α_a and LR (around -
1072 15% for both parameters) for all β_a ranges.

1073 The CM7 configuration refers to the impact of the total number of particles at the ground (N_{tot}). In this case, decreasing
1074 the upper bound of the variability range of N_{tot} by a factor of 2 (N_{tot} =500-5000 cm^{-3}) lowers the mean values of α_a and
1075 LR of around 5%. Nevertheless, this effect is totally due to the contribution of the β_a values between 2×10^{-3} and 2×10^{-2}
1076 km^{-1}sr^{-1}, where $<d\alpha_a/\alpha_a>$ and $<dLR/LR>$ are around -10%. Assuming no increase with altitude for $\sigma_{1,2}$ (sensitivity case
1077 CM8) produces relevant differences on the mean values of α_a and LR. In CM8, the overall overestimation of these two
1078 parameters is quite limited ($<d\alpha_a/\alpha_a>$ = 6.3 and $<dLR/LR>$ =6.4), whereas a large and opposite impact is observed for

1079   low and high values of β_a. In fact, $<d\alpha_a/\alpha_a>$ ($<dLR/LR>$) is equal to -14.1 (-13.9) and 18.5 (19.0) for $\beta_a= 2\times10^{-5}$ - 2×10^{-4}

1080   and $\beta_a= 2\times10^{-5}$ - 2×10^{-4} km^{-1}sr^{-1}, respectively. As explained by Barnaba et al. (2007), the dependence of $\sigma_{1,2}$ to the

1081   altitude can be associated to the fact that, when increasing the distance from the main aerosol sources, the particle

1082   processing is more efficient.

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1084   **Table B1. Mean differences of α_a and LR between different model sensitivity cases and the proposed continental model**
1085   **configuration.**

Model configuration	β_a (km^{-1}sr^{-1}) 2×10^{-5}-2×10^{-4}		β_a (km^{-1}sr^{-1}) 2×10^{-4}-2×10^{-3}		β_a (km^{-1}sr^{-1}) 2×10^{-3}-2×10^{-2}		β_a (km^{-1}sr^{-1}) 2×10^{-5}-2×10^{-2}	
	$<d\alpha_a/\alpha_a>$ (%)	$<dLR/LR>$ (%)	$<d\alpha_a/\alpha_a>$ (%)	$<dLR/LR>$ (%)	$<d\alpha_a/\alpha_a>$ (%)	$<dLR/LR>$ (%)	$<d\alpha_a/\alpha_a>$ (%)	$<dLR/LR>$ (%)
CM1 ($N_1\%=0$)	-6.2	-6.4	3.1	3.2	7.8	7.9	-3.7	-3.5
CM2 ($N_3\%=0$)	4.7	4.9	8.6	8.9	2.8	2.7	5.3	5.4
CM3 ($r_2=0.03 - 0.05$ μm)	-2.0	-1.7	-10.3	-10.2	-8.9	-8.2	-6.7	-6.4
CM4 ($r_3=1.0 - 1.2$ μm)	<1	<1	-2.1	-2.0	-5.24	-5.3	-1.2	-1.0
CM5 ($m_{1im}=0.1-0.47$)	4.3	4.2	<1	<1	<1	<1	1.8	1.8
CM6 ($m_{2r}=1.55 - 1.70$)	-10.9	-10.9	-16.2	-16.3	-18.9	-19.1	-15.3	-15.3
CM7 ($N_{TOT}=500-5000$)	<1	<1	<1	<1	-11.2	-10.7	-3.7	-3.5
CM8 (σ_1, σ_2 constant)	-14.1	-13.9	6.4	6.1	18.5	19.0	6.3	6.4

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1088 **Appendix B: Model-based functional relationships at 355 and 532 nm**

1089 The parameters of the seventh-order polynomial fit used to derive the functional relationships between log(x) and log(y)

1090 (where x = β_a and y =α_a, S_a or V_a) at λ = 355 and 532 nm are reported in Tab. A1 and Tab. A2, respectively.

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1092 **Table A1. Parameters of the Seventh-Order Polynomial Fits (y = a_0+ a_1x+a_2x^2+a_3x^3+a_4x^4+a_5x^5+a_6x^6+a_7x^7) for λ = 355 nm, with**

1093 **x=log(β_a) (in km^{-1} sr^{-1} unit) and y=log(α_a, S_a, or V_a) in (km^{-1}, cm^2/cm^3 and cm^3/cm^3, respectively).**

| Functional relantionship at 355 nm	Extinction coefficient	Surface area	Volume
a_0	3.797837507651898	12.019452592845141	-5.314834128998254
a_1	3.294032541389781	30.825966279368547	2.500484347793244
a_2	0.962603336867675	24.518531616019207	-1.196109537503000
a_3	0.241796629870675	10.625241994796593	-1.583236058579546
a_4	0.064609145804688	2.634051072085453	-0.681801883947768
a_5	0.017721752150233	0.373150843707711	-0.145232662646142
a_6	0.002722551625862	0.027971628176431	-0.015471229968392
a_7	0.000157245409783	0.000854381337164	-0.000658925756875

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1097 **Table A2. Parameters of the Seventh-Order Polynomial Fits (y= a_0+ a_1x+a_2x^2+a_3x^3+a_4x^4+a_5x^5+a_6x^6+a_7x^7) for λ = 532 nm, with**

1098 **x=log(β_a) (in km^{-1} sr^{-1} unit) and y=log(α_a, S_a, or V_a) in (km^{-1}, cm^2/cm^3 and cm^3/cm^3, respectively).**

Functional relantionship at 532 nm	Extinction coefficient	Surface area	Volume
a_0	3.797837507651898	12.019452592845141	-5.314834128998254
a_1	3.294032541389781	30.825966279368547	2.500484347793244
a_2	0.962603336867675	24.518531616019207	-1.196109537503000
a_3	0.241796629870675	10.625241994796593	-1.583236058579546
a_4	0.064609145804688	2.634051072085453	-0.681801883947768
a_5	0.017721752150233	0.373150843707711	-0.145232662646142
a_6	0.002722551625862	0.027971628176431	-0.015471229968392
a_7	0.000157245409783	0.000854381337164	-0.000658925756875

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Appendix C: Model – EARLINET comparison at 532 nm

1103   Figure C1 depicts the result of the comparison between EARLINET stations and our developed model (red and blue

1104   curves, respectively) in terms of 'mean' LR per bin of β_a at λ=532 nm. Note that only β_a bins containing at least 1% of

1105   the total modeled data were considered. Similarly to the results at 355 nm shown in section 4.1, a general good

1106   agreement between the modeled and the measured LR values is found. As attested by the low value of the mean

1107   discrepancy of Table 6, the modeled curve well fits with Madrid observations. Some major deviations are found for

1108   Lecce, which, however, at 532 nm, has a very low number of considered points (i.e. 109).
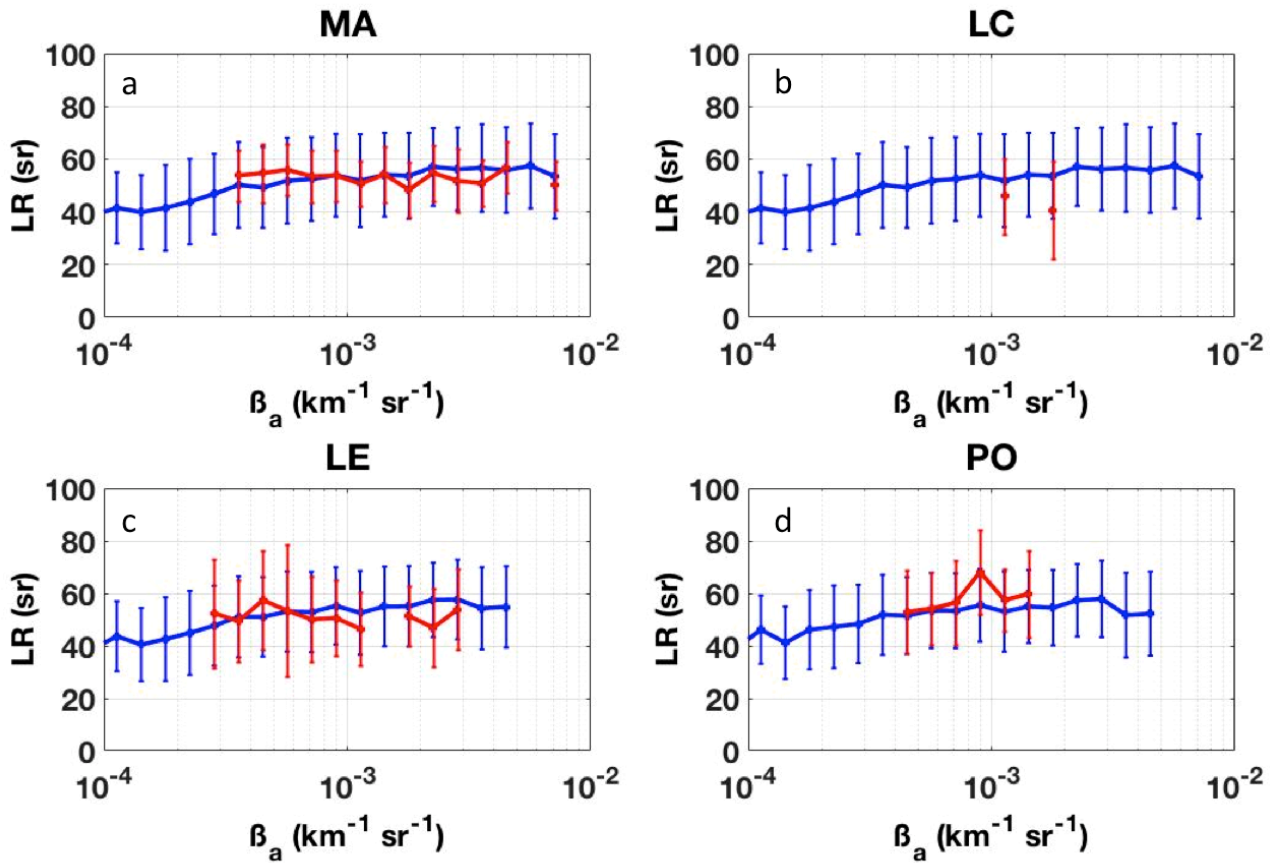

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1110   **Figure C1. Model-simulated (blue) and lidar measured (red) LR vs β_a mean curves at 532 nm calculated per 10 equally**

1111   **spaced bins per decade of β_a in a) Madrid, b) Lecce, c) Leipzig, and d) Potenza EARLINET lidar station. Vertical bars are**

1112   **the associated standard deviations.**

Appendix D: Model sensitivity tests for optimal configurations at LC and PO sites

1114   According to the results reported in Tab. B1, two model configurations (CM0a and CM0b) have been set up to better
1115   reproduce the EARLINET observations of LR vs β_a at LC and PO sites, respectively. The comparison between these
1116   two configurations, the EARLINET measurements and the CM0 set-up are illustrated in Fig. B1 (panel a and b for LC
1117   and PO, respectively) in terms of LR mean value curves per 10 equally spaced bins per decade of β_a. Blue and red
1118   colors have the same meaning of Fig. 5 (i.e. CM0 model and observation curves, respectively), black curves refer to the
1119   LR vs β_a estimated through the CM0a and CM0b model versions for LC and PO stations, respectively. Vertical bars are
1120   the associated standard deviations.

1121   The only difference between CM0a and CM0 configuration consists in the upper bound of the variability range of N_{tot}
1122   (5000 vs 10000 cm^{-3} at ground, respectively). This modification seems to fit the observed LR vs β_a behavior at 355 nm.
1123   The upper bound N_{tot} value is similar to the one (i.e. N_{tot} upper bound =3000 cm^{-3} at ground) used in the work of
1124   Barnaba et al. (2007) to characterize the optical properties of the continental aerosol present over southeastern Italy.
1125   The computed mean model-measurement LR relative difference between CM0a configuration and LC Earlinet
1126   measurements is around 5%.

1127   Similarly, the CM0b configuration uses the same value for the upper bound of N_{tot} variability range and, in addition,
1128   higher values of the r_3 variability range of (1.0 – 1.2 μm vs 0.3 - 0.5 μm, respectively). As highlighted by the panel b of
1129   Fig. B1, this model configuration allows well reproducing the LR vs β_a behavior derived by EARLINET lidar Raman
1130   measurements at 355 nm. This result seems to indicate the presence of coarser aerosols in a clean continental
1131   environment. In comparison to the CM0 model, the mean model-measurement LR relative difference decreases from
1132   17% to 6%.
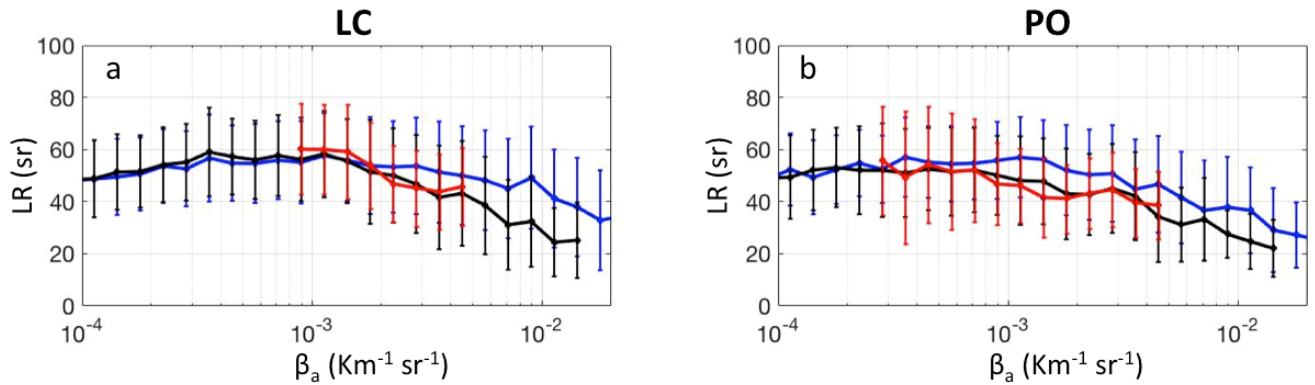
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Figure D1. Model-simulated (blue and black lines) and lidar measured (red lines) LR vs β_a mean curves at 355 nm calculated per 10 equally spaced bins per decade of β_a for the LC and PO EARLINET lidar stations (panel a and b, respectively). Blue color refers to CM0 model configuration, black color to CM0a and CM0b model configurations adapted to LC and PO sites, respectively.

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1139 **Appendix E: ALC vs sun-photometer AOTs**

1140 To have sense of both absolute and relative errors of AOT, we reported in this section the scatter plots between the
1141 hourly-mean coincident AOTs at 1064 nm as derived by ALC model-based approach and those measured at 1020 nm
1142 by the sun-photometers installed at RTV, SPC and ASC, respectively (Figure E1, E2 and E3). The corresponding linear
1143 fit y = bx (red line), where x = sun-photometer AOT, y = Nimbus CHM15k AOT are also shown in the plots. The
1144 values of the correlation coefficients for the three sites (R = 0.77, R=0.72 and R=0.73 for RTV, SPC and ASC,
1145 respectively) attest a relatively good agreement between the two AOT measurements.
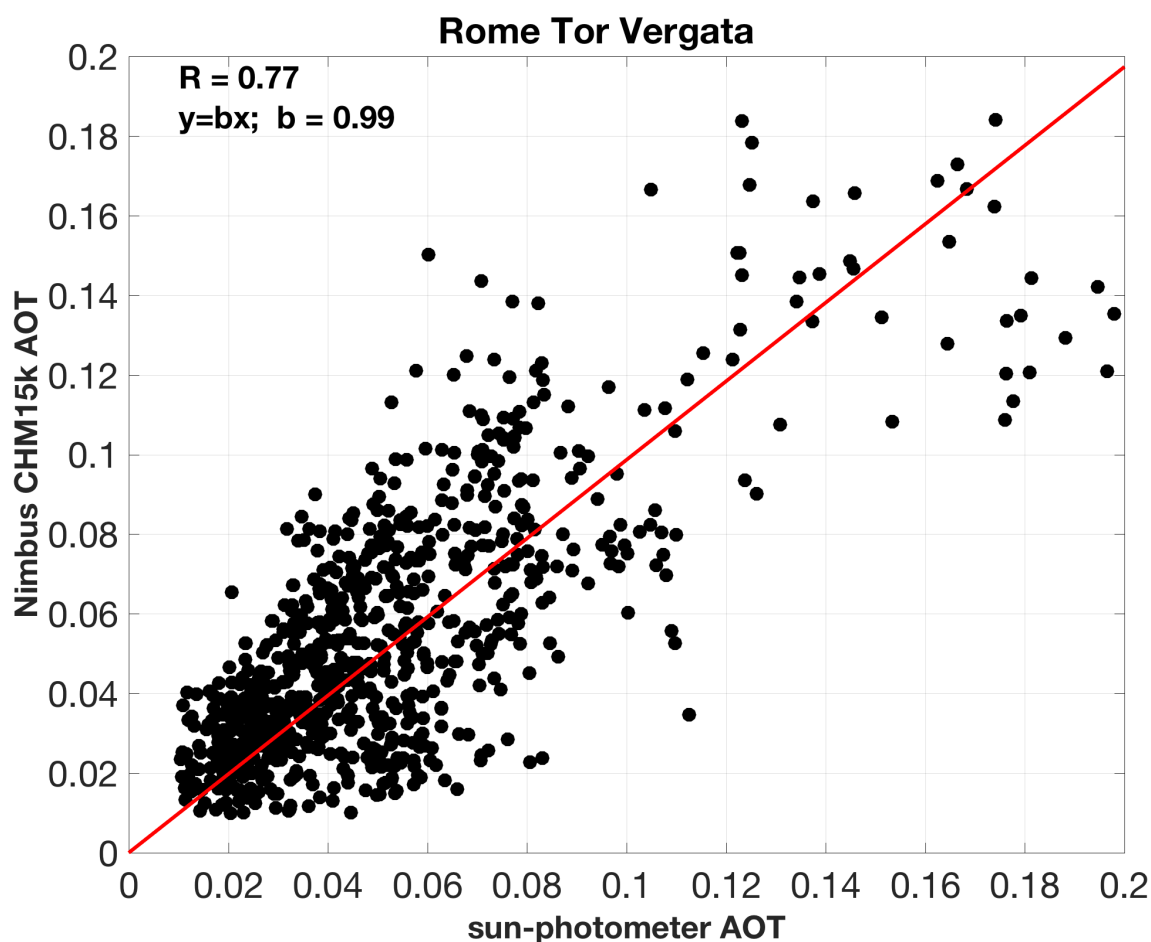
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1147 **Figure E1. Scatter plot between the hourly-mean coincident AOTs at 1064 nm as derived by the ALC model-based approach**
1148 **and measured at 1020 nm by the AERONET photometer at RTV. The red line represents the linear fit y = bx between the**
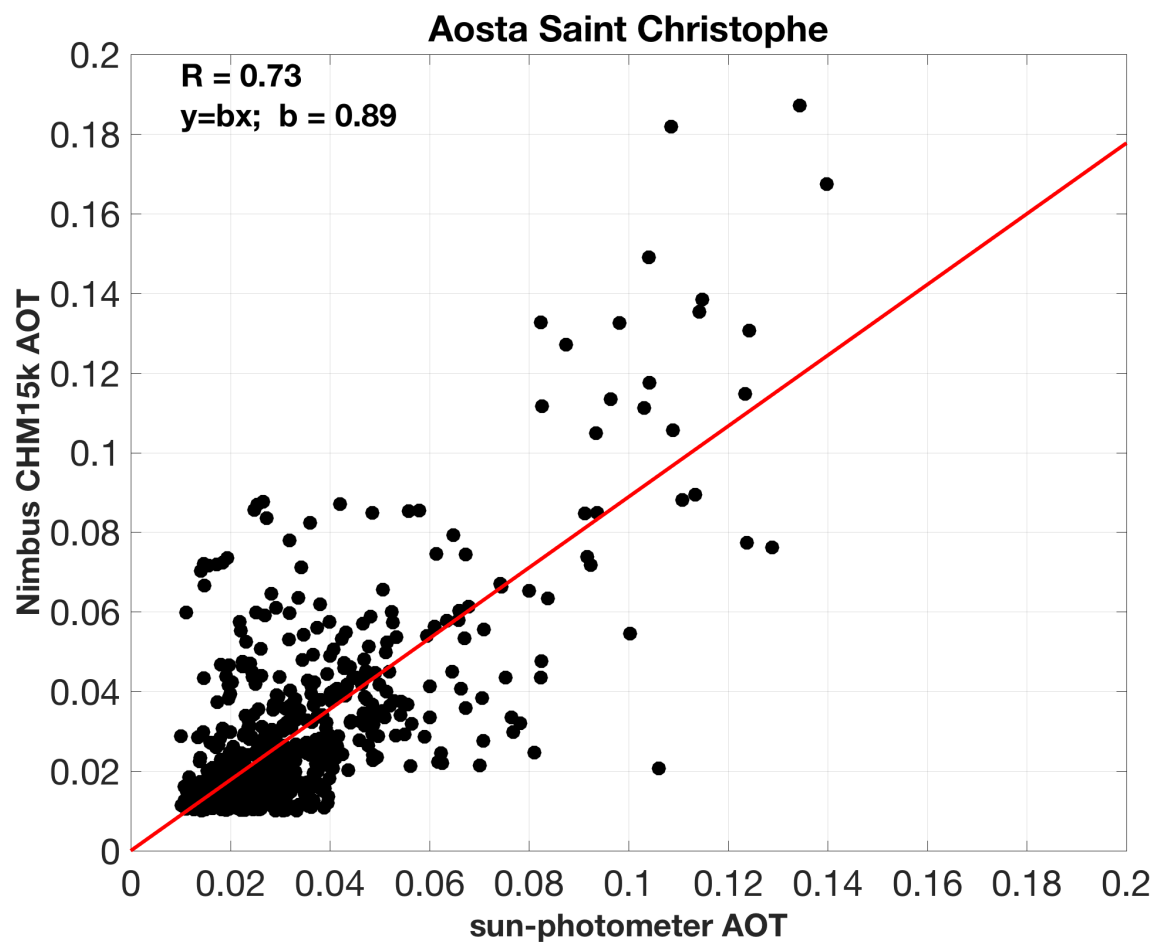1149 **two datasets, where x = sun-photometer AOT; y = Nimbus CHM15k AOT.**

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San Pietro Capofiume

R = 0.72
y=bx; b = 0.97

(x-axis: sun-photometer AOT; y-axis: Nimbus CHM15k AOT)
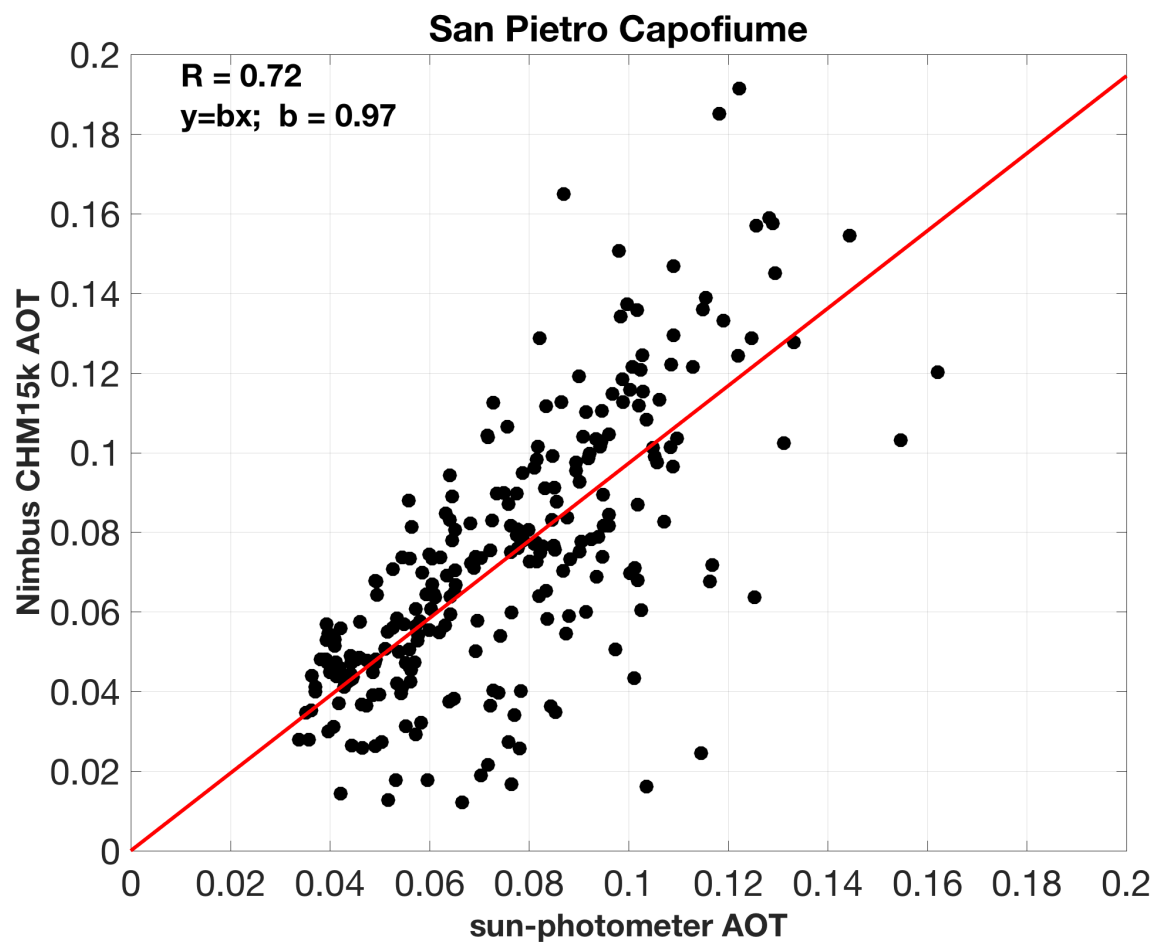
Figure E2. Scatter plot between the hourly-mean coincident AOTs at 1064 nm as derived by the ALC model-based approach and measured at 1020 nm by the SKYRAD photometer at SPC. The red line represents the linear fit y = bx between the two datasets, where x = sun-photometer AOT; y = Nimbus CHM15k AOT.

Aosta Saint Christophe

R = 0.73
y=bx; b = 0.89

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Figure E3. Scatter plot between the hourly-mean coincident AOTs at 1064 nm as derived by the ALC model-based approach and measured at 1020 nm by the SKYRAD photometer at ASC. The red line represents the linear fit y = bx between the two datasets, where x = sun-photometer AOT; y = Nimbus CHM15k AOT.

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