

# 1 **A multi-wavelength numerical model in support to quantitative** 2 **retrievals of aerosol properties from automated-lidar-ceilometers** 3 **and test applications for AOT and PM10 estimation**

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9 **Abstract.** Knowledge of the vertical distribution of aerosol particles is a key factor in the study of climate, air pollution,  
10 and meteorological processes. The use of automated lidar-ceilometers (ALC) systems for the aerosol vertically-resolved  
11 characterization has increased in the recent years thanks to their low construction and operation costs, and to their  
12 capability at providing continuous, unattended measurements. At the same time there is a need to convert the ALC  
13 signals into usable geophysical quantities. In fact, the quantitative assessment of the aerosol properties from ALC  
14 measurements and the relevant assimilation in meteorological forecast models is amongst the main objectives of the EU  
15 COST Action TOPROF (Towards Operational ground-based PROFiling with ALCs, doppler lidars and microwave  
16 radiometers). Concurrently, the E-PROFILE program of the European Meteorological Services Network (EUMETNET)  
17 focuses on the harmonization of ALC measurements and data provision across Europe. Within these frameworks, we  
18 implemented a model-assisted methodology to retrieve key aerosol properties (extinction coefficient, surface area and  
19 volume) from elastic lidar and/or ALC measurements. The method is based on results from a large set of aerosol  
20 scattering simulations (Mie-theory) performed at UV, visible and near IR wavelengths using a "Monte-Carlo" approach  
21 to select the input aerosol microphysical properties. An average 'continental aerosol type' (i.e. clean-to-moderately  
22 polluted continental aerosol conditions) is addressed in this study. Based on the simulation results, we derive mean  
23 functional relationships linking the aerosol backscatter coefficients to the above-mentioned variables. Applied in the  
24 data inversion of single wavelength lidars and/or ALCs, these relationships allow quantitative determination of the  
25 vertically-resolved aerosols backscatter, extinction, volume and surface area and, in turn, of the extinction-to-  
26 backscatter ratio (i.e., the lidar-ratio, LR) and of extinction-to-volume conversion factor ( $c_v$ ) at 355, 532, 1064 nm.  
27 These variables provide valuable information for visibility, radiative transfer and air quality applications. This study  
28 also includes 1) validation of the model simulations with real measurements and 2) test applications of the proposed  
29 model-based ALC inversion methodology. In particular, our model simulations were compared to backscatter and  
30 extinction coefficients independently retrieved by Raman lidar systems operating at different continental sites within the  
31 European Aerosol Research Lidar NETwork (EARLINET). This comparison shows good model-measurements  
32 agreement, with LR discrepancies below 20%. The model-assisted quantitative retrieval of both aerosol extinction and  
33 volume was then tested using raw data from three different ALCs systems (CHM15k-Nimbus), operating within the  
34 Italian Automated Lidar-ceilometer Network (ALICENET). To this purpose, a one-year-record of the ALCs-derived  
35 aerosol optical thickness (AOT) at each site was compared to direct AOT measurements performed by co-located sun-  
36 sky photometers. This comparison shows an overall AOT agreement within 30% at all sites. At one site, the model-  
37 assisted ALC estimation of the aerosol volume and mass (i.e., PM10) in the lowermost levels was compared to values

38 measured at the surface-level by co-located in situ instrumentation. Within this exercise, the ALC-derived daily-mean  
39 mass concentration was found to well reproduce the corresponding (EU regulated) PM<sub>10</sub> values measured by the local  
40 Air Quality agency in terms of both temporal variability and absolute values. Although limited in space and time, the  
41 good performances of the proposed approach in these preliminary tests suggest it could possibly represent a valid option  
42 to extend the capabilities of ALCs at providing quantitative information for operational air quality and meteorological  
43 monitoring.

## 44 **1 Introduction**

45 Due to the impact of atmospheric aerosols on both air quality and climate, substantial efforts have been made to expand  
46 our knowledge of their sources, properties and fate. Aerosol particles affect the Earth's radiation budget mainly by two  
47 different processes: 1) by scattering and absorbing both solar and terrestrial radiation (aerosol direct effect, Haywood  
48 and Boucher, 2000 and aerosol semi-direct effect, Johnson et al., 2004) and 2) by serving as cloud and ice condensation  
49 nuclei (aerosol indirect effect, Lohmann and Feichter, 2005, Stevens and Feingold, 2009 and Feingold et al., 2016). The  
50 complexity of these processes and the extreme spatial and temporal variability of the aerosol sources, physical and  
51 chemical properties and atmospheric processing make the quantification of their impacts very difficult. Aerosols have  
52 also proven detrimental effects on human health (e.g., D'Amato et al., 2013, World Health Organization, 2013,  
53 Lelieveld et al., 2015). In fact, their concentration (often evaluated in terms of particulate matter mass, or PM) is  
54 regulated by specific air quality legislation worldwide. In Europe, the Air Quality Directive 2008/50 defines the  
55 'objectives for ambient air quality designed to avoid, prevent or reduce harmful effects on human health and the  
56 environment as a whole' (EC, 2008).

57 Among the aerosol observational systems, the LIDAR technique has been proved to be the optimal tool to provide  
58 range-resolved, accurate aerosol data necessary in radiative transfer computations (e.g. Koetz et al., 2006, Tosca et al.,  
59 2017) and is often usefully employed in supporting air quality studies (e.g. Menut et al., 1997, He et al., 2012). With a  
60 spectrum of different system types (elastic backscatter, Raman, High Spectral Resolution, and multi-wavelength lidars),  
61 each with specific pro and cons (Lolli et al., 2018), this technique allows retrievals of aerosol and cloud optical  
62 properties and relevant distribution within the atmospheric column at several ground-based observational sites (Fernald  
63 et al., 1972; Klett, 1981; Shipley et al., 1983, Kovalev and Eichinger, 2004, Heese and Wiegner, 2008; Ansmann et al.,  
64 2012). Since 2006, the Cloud Aerosol Lidar and Infrared Pathfinder Satellite Observation (CALIPSO) platform (Winker  
65 et al. 2003) also provides a unique, global view of aerosol and cloud vertical distributions through space-based  
66 observations (at the operating wavelengths of 532 and 1064 nm). Recently, within the Cloud-Aerosol Transport System  
67 (CATS) mission, a lidar was also installed at the International Space Station (ISS, McGill et al., 2015 and York et al.,  
68 2016). Space-borne lidar observations are however affected by some drawbacks, as: 1) limited temporal resolution and  
69 spatial coverage (the CALIPSO spatial distance between two consecutive ground tracks is about 1000 kilometers and  
70 each track has a footprint of 70 m), 2) the contamination of unscreened clouds, and 3) difficulties in quantitatively  
71 characterizing the aerosol properties in the lowermost troposphere (Pappalardo et al., 2010). Ground-based lidar  
72 networks thus still represent key tools in integrating space-borne observations to study aerosol properties and their 4D  
73 distribution. An example of these networks is the European Aerosol Research Lidar NETwork (EARLINET,  
74 <http://www.earlinet.org/>), which, since 2000, provides an extensive collection of ground-based data for the aerosol  
75 vertical distribution over Europe (Bösenberg et al., 2003, Pappalardo et al., 2014). The advanced multi-wavelength  
76 elastic and Raman lidars employed in this network allows independent retrieval of aerosol extinction ( $\alpha_a$ ) and  
77 backscattering coefficient ( $\beta_a$ ) profiles. Yet, despite their unsurpassed potential in data accuracy, advanced lidar

78 networks such as EARLINET have the unsolved problems of the sparse spatial and temporal sampling and of the  
79 complexity of operations. In fact, the typical distance between the EARLINET stations is of the order of several  
80 hundreds of kilometers and regular measurements of EARLINET are only performed on selected days of the week  
81 (Mondays and Thursdays) and for a few hours (mainly at nighttime, due to low signal-to-noise ratio of the Raman  
82 signal in daylight). Furthermore, these systems are complicated to be operated, require specific expertise and are  
83 therefore unsuitable for operational applications.

84 Nowadays, hundreds of single channel Automated Lidar Ceilometers (ALCs) are in operation over Europe and  
85 worldwide. Although such simple lidar-type instruments were originally designed for cloud base detection only, the  
86 recent technological advancements make now these systems reliable and affordable for aerosol measurements,  
87 increasing the interest in using this technology in different aerosol-related sectors (e.g. air quality, aviation security,  
88 meteorology, etc.). Recent studies showed that the ALC technology is now mature enough to be used for a quantitative  
89 evaluation of the aerosol physical properties in the lower atmosphere (Wiegner, M. and A. Geiß, 2012, Wiegner et al.,  
90 2014) and the exploitation of the full potential of ALCs in the aerosol remote sensing is a current matter of discussion in  
91 the lidar community (e.g. Madonna et al., 2015, 2018). The evaluation of ALC capabilities at providing quantitative  
92 aerosol information is among the main objectives of the EU COST Action ES1303, TOPROF (Towards Operational  
93 ground-based PROFiling with ALCs, doppler lidars and microwave radiometers). An effort in this direction is also  
94 underway in the framework of E-PROFILE, one of the observation programs of the EUROpean METeorological services  
95 NETwork (EUMETNET). In fact, several ALC stations are progressively joining E-PROFILE to develop an operational  
96 network to produce and exchange ALC-derived profiles of attenuated backscatter. A recent project funded by the EU  
97 LIFE+ program (DIAPASON, Desert-dust Impact on Air quality through model-Predictions and Advanced Sensors  
98 ObservatioNs, LIFE+2010 ENV/IT/391) also prototyped and tested an ALC system with an additional depolarization  
99 channel, capable of discriminating non spherical aerosol types, such as desert dust (Gobbi et al., 2018). Such upgraded  
100 ALC systems could further improve the capabilities of the operational aerosol profiling in a near future.

101 Given the necessity to couple advancement in instrumental technology with tools capable of translating raw data into a  
102 robust, quantitative and usable information, we propose and characterize here a methodology to be applied to elastic  
103 backscatter lidars and/or ALC measurements to retrieve, in a quasi-automatic way, vertically-resolved profiles of some  
104 key aerosol optical and microphysical properties. This effort is intended to contribute better exploit these systems  
105 potential in integrating data collected by more advanced lidar systems/networks. In particular, the ALC-derived aerosol  
106 properties addressed in this study are: backscatter ( $\beta_a$ ,  $\text{km}^{-1} \text{sr}^{-1}$ ), extinction ( $\alpha_a$ ,  $\text{km}^{-1}$ ), surface area ( $S_a$ ,  $\text{cm}^2 \text{cm}^{-3}$ ) and  
107 volume ( $V_a$ ,  $\text{cm}^3 \text{cm}^{-3}$ ), the latter being convertible into aerosol mass concentration ( $\mu\text{g m}^{-3}$ ) via assumption on particle  
108 density. To this purpose, we developed a numerical aerosol model to perform a large set of aerosol scattering  
109 simulations. Based on results from this numerical model, we derive mean functional relationships linking  $\beta_a$  to  $\alpha_a$ ,  $S_a$   
110 and  $V_a$ , respectively. These relationships are then applied in the ALC data inversion and analysis. A similar approach  
111 was applied in past studies for lidar-based investigations of stratospheric (Gobbi, 1995) and tropospheric aerosols  
112 (maritime, desert dust and continental type) at visible and UV lidar wavelengths, (Barnaba and Gobbi, 2001, Barnaba  
113 and Gobbi, 2004a, hereafter BG01, BG04a, respectively, Barnaba et al., 2004). Here we extend this approach to all the  
114 Nd:YAG laser harmonics commonly used by both advanced lidars and ALC systems (i.e. 355, 532, 1064 nm  
115 wavelengths) and specifically address an ‘average-continental’ aerosol type, intended to represent clean-to-moderately  
116 polluted continental aerosol conditions (see Section 2.1). In fact, despite the known differences that can be encountered

117 across the continent both in the short and the long-term (e.g., Putaud et al. 2010), this aerosol type is expected to  
118 climatologically dominate over most of Europe.

119 Overall, this investigation is organized as follows: in Section 2 we describe the aerosol model set up to reproduce clean  
120 to moderately polluted continental conditions, and the Monte Carlo methodology followed to compute the  
121 corresponding bulk optical and physical properties. Section 3 shows and discusses the results of the numerical model,  
122 and presents the model-based, mean functional relationships linking the different variables at 355, 532 and 1064 nm. In  
123 Section 4 we evaluate both the model simulations capability to reproduce real measurements in continental aerosol  
124 conditions, and the capability of the model-based ALC inversion approach to derive quantitative geophysical  
125 information. The EARLINET database was used for the first task while tests on the accuracy of the model-based ALC  
126 inversion were performed evaluating both the ALC-derived aerosol volume and optical thickness (AOT, i.e. the  
127 vertically integrated aerosol extinction). To this purpose we applied the proposed methodology to three ALC systems  
128 operating within the Italian Automated Lidar-Ceilometer NETwork (ALICE-NET, [www.alice-net.eu](http://www.alice-net.eu)). In particular, the  
129 ALC-derived AOT and aerosol volume (plus mass) were compared, respectively, to reference measurements performed  
130 by ground-based sun photometers and in situ aerosol instruments (optical counters and PM10 samplers).

131 Section 5 summarizes the developed approach and main results, critically examining strengths and weaknesses. It also  
132 includes discussion on the perspectives of the application of this (or similar) methodology in operational ALC  
133 networks.

134

## 135 **2 The aerosol model**

136 A numerical aerosol model was set up to calculate mean functional relationships between the aerosol backscatter ( $\beta_a$ )  
137 and some relevant aerosol properties, as  $\alpha_a$ ,  $S_a$  and  $V_a$ . This is done in a two-step procedure (Figure 1), following an  
138 approach similar to that developed by BG01 and BG04a:

139 1) Generation of a large set (here 20000) of aerosol optical and physical properties by randomly varying, within  
140 appropriate ranges, the microphysical parameters describing the aerosol size distribution and composition (blue box in  
141 Fig. 1);

142 2) Based on results at point 1), determination of mean functional relationships linking such key variables (grey box in  
143 Fig. 1).

144 The following Section describes rationale and set-up of the first step, the second one being thoroughly discussed in  
145 Section 3.

### 146 **2.1 Selection of the aerosol microphysical parameters**

147 As anticipated, an average ‘continental’ aerosol type (i.e. describing clean to moderately polluted continental  
148 conditions, e.g. Hess et al., 1998) was targeted in this study, this being the aerosol type expected to dominate over  
149 Europe. Based on a scheme originally proposed by d’Almeida et al. (1991), and on a large set of following  
150 observational evidences (e.g. Van Dingenen et al., 2004), in this work its size distribution is described as an external  
151 mixture of three size modes. These are (in order of increasing size range): 1) a first ultrafine mode; 2) a second fine  
152 mode, mainly composed of water-soluble particles; 3) a third mode of coarse particles.

153 A three-mode lognormal size distribution described by Eq. (1) is employed to this purpose:

$$n(r) = \frac{dN}{d \log r} = \sum_{i=1}^3 \frac{N_i}{\sqrt{2\pi} \log \sigma_i} \exp \left[ -\frac{(\log r - \log r_i)^2}{2(\log \sigma_i)^2} \right] \quad (1).$$

In Eq. (1),  $r_i$ ,  $\sigma_i$  and  $N_i$  are respectively the modal radius, the width and the particle number density of the  $i^{\text{th}}$  aerosol mode ( $i = 1, 2, 3$ ). At each computation,  $r_i$  and  $\sigma_i$  are randomly chosen within a relevant variability range. Values of  $N_i$  are conversely obtained by firstly randomly choosing the total number of particles,  $N_{\text{tot}}$ , to be included in the whole size distribution ( $N_{\text{tot}} = N_1 + N_2 + N_3$ ), and then by applying specific rules for the number mixing ratio,  $x_i = N_i/N_{\text{tot}}$ , of each component to this total. To reproduce clean to moderately polluted continental conditions, the value of  $N_{\text{tot}}$  is made variable between  $10^3$  and  $3 \times 10^4 \text{ cm}^{-3}$  (e.g. Hess et al., 1998; Van Dingenen et al., 2004). Being the result of different sources/processes, the three modes are also assumed to have a different composition, this impacting the optical computations through the relevant particle refractive index ( $m_i$ ), with both its real and imaginary component ( $m_i = m_{r_i} - i \times m_{\text{im}_i}$ ). The Mie theory for spherical particles of radius  $r_i$  and refractive index  $m_i$  is then used to compute the extinction and backscatter coefficients (see below).

A description of the assumptions made for each mode and relevant parameter, mostly based on literature data (Table 1), is given hereafter, the summary of the relevant variability chosen for each parameter being provided in Table 2.

#### 1) First Mode

This ultrafine mode is the one more directly simulating fresh, anthropogenic emissions. The number mixing ratio  $x_{i=1}$  ( $N_{i=1}/N_{\text{tot}}$ ) of this mode is let variable between 10% (rural conditions, Van Dingenen et al., 2004) and 60% (more polluted conditions, Hess et al., 1998). The variability of its modal radius ( $r_1 = 0.005 - 0.03 \mu\text{m}$ ) is chosen to include from nucleation mode particles to Aitken mode particles. To take into account the wide variability of species within this ultrafine mode, from non-absorbing (e.g., inorganic particles) to highly absorbing materials (e.g. black carbon), wide ranges of variability has been set for its refractive indexes (at  $\lambda=355 \text{ nm}$ :  $m_{r_1}$  in the range 1.40 - 1.8, and  $m_{\text{im}_1}$  in the range 0.01 - 0.47, see Table 2 for the corresponding values at  $\lambda=532$  and 1064 nm).

#### 2) Second mode

The second aerosol mode accounts for 40-90% of  $N_{\text{tot}}$ , with (dry)  $r_2$  between 0.03 and 0.1  $\mu\text{m}$ . Its composition ( $m_{r_2}$ , and  $m_{\text{im}_2}$ ) is also made highly variable so to include water soluble inorganic and organic particles (Hess et al., 1998; BG04a; Dinar et al. 2008). In this case, at  $\lambda=355 \text{ nm}$ ,  $m_{r_2}$  is in the range 1.40 - 1.7 and  $m_{\text{im}_2}$  is in the range 0.0001 - 0.01 (Table 2).

#### 3) Third mode

This coarser aerosol mode (modal radius  $r_3$  in the range 0.3 - 0.5  $\mu\text{m}$ ) is mainly intended to account for soil derived (dust-like) particles that are a primary continental emission. A quite narrow variability is thus fixed for its  $m_{r_3}$  and  $m_{\text{im}_3}$  values (1.5 - 1.6 and 0.0001 - 0.01, respectively at 355 nm). The relevant number mixing ratio  $x_3$  ( $N_3/N_{\text{tot}}$ ) is set variable between 0.001% and 0.5 %, this mode contributing mostly to the total aerosol volume (thus mass) rather than to the total number of particles.

As mentioned, refractive indexes were also made wavelength dependent, as this feature is also typically observed as linked to the different particle composition. In particular, for the second mode (water-soluble particles) we include an increase with the wavelength of the upper boundary values of  $m_{\text{im}_2}$  and a decrease of  $m_{r_2}$  at  $\lambda=1064 \text{ nm}$  (d'Almeida et al., 1991). For the (dust-like) third-mode particles, the upper boundary values of  $m_{\text{im}_3}$  are set to decrease with increasing wavelengths (Gasteiger et al., 2011, Wagner et al., 2012).

191 For convenience, the aerosol parameters boundaries summarized in Table 2 refer to dry particles and to ground level.  
 192 However, the effect of a variable RH, its variability with altitude as well as the generally observed decrease of particle  
 193 number with altitude is also considered in the model. More specifically, the number of particles in each mode,  $N_i$ , and  
 194 RH are both made altitude-dependent through the following equations (Patterson et al., 1980, BG01):

$$195 \quad N_i(z) = N_i(0) \times \exp\left(\frac{-z}{H_i}\right), \quad (2)$$

$$196 \quad RH(z) = 70 \times \exp\left(\frac{-z}{5.5 \text{ km}}\right) \times (1 + dRH), \quad (3)$$

197 the altitude  $z$  being variable here between 0 and 5 km.  $N_i(0)$  and  $H_i$  in eq.2 are the number of particles at the ground and  
 198 the scale height for each mode, respectively.

199 To describe the altitude effect, in eq. (2) an exponential decrease with height of the particle number density is assumed.  
 200 To rescale the particle number density of the different modes,  $H_{i=1=2}$  is set equal to 5.5 km (Barnaba et al., 2007) while  
 201  $H_{i=3}$  (coarse particles) is set to 0.8 km (Barnaba et al., 2007). In eq. (3), the additional term  $(1+dRH)$  allows a further  
 202 variability with respect to the mean  $RH(z)$  profile assumed, here  $dRH$  is randomly chosen between -60 and +60). Values  
 203 of RH greater than 95% are discarded to avoid divergence.

204 Additionally, while first and third modes are assumed to be water insoluble, the second mode ( $i=2$ ) is fully hygroscopic.  
 205 Aerosol humidification is thus considered to act on both particle size and refractive indices of the second aerosol mode  
 206 (e.g., BG01), as:

$$207 \quad r_{2\_RH} = r_{2\_0} \sqrt{\frac{2-0.01RH}{2(1-0.01RH)}}, \quad (4)$$

$$208 \quad m_{2\_RH} = m_w + (m_{2\_0} - m_w) \left(\frac{r_{2\_0}}{r_{2\_RH}}\right)^3, \quad (5).$$

209 In eq. 4 and 5,  $r_{2\_RH}$  and  $m_{2\_RH}$  are the RH-corrected modal radius and refractive index for the second mode,  
 210 respectively;  $r_{2\_0}$  and  $m_{2\_0}$  are the particle dry modal radius and refractive index, respectively;  $m_w$  is the water refractive  
 211 index (assumed as equal to  $1.34 - i7e^{-9}$ ,  $1.33 - i1.3e^{-9}$   $1.33 - i2.9e^{-6}$  at 355, 532 and 1064 nm, respectively).

212 Finally, following Barnaba et al. (2007), an increase of the width of the size distribution with altitude (eq. 6) has been  
 213 introduced for the first and second aerosol mode:

$$214 \quad \sigma_{1,2}(z) = \sigma_{1,2,z0} \times \exp\left(\frac{z}{30}\right). \quad (6)$$

215 In fact, Barnaba et al., (2007) showed that this was necessary to better reproduce the observed decrease of the Lidar  
 216 Ratio (LR) with altitude, and likely related to a broadening of the particle size distribution with aging.

217 Once the value of each microphysical parameter is randomly selected within its relevant variability range, and once  
 218 corrections are applied following eqs. (2) – (6), each resulting aerosol size and composition-resolved distribution is used  
 219 to compute the aerosol  $S_a$  and  $V_a$ , as well as to feed a Mie code (assumption of spherical particles, Bohren and  
 220 Huffman, 1983) to compute  $\beta_a$ , and  $\alpha_a$ , (BG01, see also Fig. 1). Overall, the equations used are as follows:

$$221 \quad \beta_a = \int Q_{bsc}(r, \lambda, m) \pi r^2 \frac{dN}{d \log r} \frac{1}{r \ln 10} dr \quad (7)$$

$$222 \quad \alpha_a = \int Q_{ext}(r, \lambda, m) \pi r^2 \frac{dN}{d \log r} \frac{1}{r \ln 10} dr \quad (8)$$

223 
$$S_a = 4\pi \int r^2 \frac{dN}{d \log r} \frac{1}{r \ln 10} dr \quad (9)$$

224 
$$V_a = \frac{4}{3}\pi \int r^3 \frac{dN}{d \log r} \frac{1}{r \ln 10} dr, \quad (10)$$

225 where  $Q_{\text{bsc}}(r_i, \lambda, m_i)$  and  $Q_{\text{ext}}(r_i, \lambda, m_i)$  are, respectively, the backscatter and the extinction efficiencies. As mentioned,  
 226 the optical computations are made at the three different wavelengths: 355, 532, 1064 nm (i.e., those of Nd:YAG laser  
 227 harmonics, the most common wavelengths used by ground-based and space-borne aerosol lidars).

228 Since in our simulations the third aerosol mode is intended to represent dust-like particles, an empirical correction for  
 229 non-sphericity is also applied to the Mie-derived optical properties of this mode. This procedure is based on BG01,  
 230 which uses the results of Mishchenko et al. (1997) obtained for surface-equivalent mixtures of prolate and oblate  
 231 spheroids.

## 232 2.2 Model simulation results

233 In Figure 2 we show the results of 20000 simulations of continental aerosol optical and physical properties derived  
 234 randomly varying the relevant aerosol size distributions and compositions as described in the previous section. In  
 235 particular, the results for  $\alpha_a$ ,  $S_a$  and  $V_a$  are shown as a function of  $\beta_a$  in Figure 2a, b, c (blue crosses) referring to  $\lambda =$   
 236 1064 nm. For each variable (A), average value per bin of  $\beta_a$  and relevant standard deviations ( $\langle A \rangle \pm dA$ ) are shown as  
 237 red dots and vertical bars, respectively. Note that 10 equally spaced bins per decade of  $\beta$  have been considered, and that  
 238  $\langle A \rangle \pm dA$  are only shown for bins containing at least 1% of the total number of pairs. Corresponding relative errors  
 239 ( $dA/\langle A \rangle$ ) are depicted in Figure 2d, e, f. Some sensitivity tests of these model outputs to the variability of the input  
 240 microphysical parameters employed are provided in Appendix A.

241 Based on these results, at step-two of the procedure (see scheme in Figure 1), we derive aerosol-specific mean  
 242 relationships linking aerosol extinction, surface area and volume ( $\alpha_a$ ,  $S_a$  and  $V_a$ ) to its backscatter ( $\beta_a$ ). To this purpose,  
 243 we used a seventh-order polynomial fit in log-log coordinates. The choice of a seventh-order polynomial fit was made  
 244 for homogeneity with BG01 and BG04a. These relationships are shown as green lines in Figure 2a, b, c while the  
 245 relevant fit parameters are reported in Table 3 referring to  $\lambda = 1064$  nm (fit parameters related to computations at  $\lambda =$   
 246 355 and 532 nm, are given in Table A1 and Table A2, Appendix B).

247 The red vertical bars of Figure 2 also highlight the ranges of  $\alpha_a$ ,  $S_a$  and  $V_a$  which are statistically significant, i.e. those in  
 248 which, at  $\lambda = 1064$  nm, the model provides at least 1% of the total points per corresponding bin of  $\beta_a$ . These are:  $10^{-4} -$   
 249  $10^{-1} \text{ km}^{-1}$ ,  $10^{-7} - 10^{-5} \text{ cm}^2/\text{cm}^3$  and  $10^{-13} - 10^{-10} \text{ cm}^3/\text{cm}^3$ , for  $\alpha_a$ ,  $S_a$  and  $V_a$  respectively, corresponding to the backscatter  
 250 range  $9 \times 10^{-5} \leq \beta_a \leq 4 \times 10^{-3} \text{ km}^{-1} \text{ sr}^{-1}$ . In terms of aerosol properties variability, the relative errors associated to  $\alpha_a$  and  $V_a$   
 251 show almost no dependence on  $\beta_a$ , with values between 30% and 40%. Conversely, the modeled aerosol surface area  
 252 exhibits a larger dispersion, with relative error values spanning the range 40% - 70%, and decreasing as  $\beta_a$  increases.

253 A key parameter for the inversion of lidar signals is the so-called Lidar Ratio ( $LR$ ), i.e. the ratio between  $\alpha_a$  and  $\beta_a$   
 254 (Ansmann et al., 1992). In Figure 3 we thus show the results of our simulations in terms of  $LR$  vs  $\beta_a$  at the three  $\lambda$  (355,  
 255 532 and 1064 nm, Figure 3a, b, c, respectively) and relevant  $dLR/LR$  values (Figure 3d, e, f, respectively). The color  
 256 code is the same of Fig. 2. Additional horizontal black lines have been inserted representing values (solid central lines)  
 257 of the ‘weighted- $LR$ ’  $\pm 1$  standard deviation (dotted side lines), i.e. the  $LR$  weighted by the number of simulated points  
 258 in each considered backscatter bin. The ‘weighted- $LR$ ’ values derived at 355, 532 and 1064 nm, are  $50.1 \pm 17.9$  sr,  $49.6$   
 259  $\pm 16.0$  sr and  $37.7 \pm 12.6$  sr, respectively. Figure 3 also allows showing that the statistically significant regions of

260 simulated backscatter values shifts towards smaller values with increasing  $\lambda$  (e.g. at  $\lambda = 355$ , the  $\beta_a$  extending regions is  
261  $4 \times 10^{-5} - 2 \times 10^{-2} \text{ km}^{-1} \text{ sr}^{-1}$ , whereas, at 532 nm, it ranges between  $2 \times 10^{-5} - 1 \times 10^{-2} \text{ km}^{-1} \text{ sr}^{-1}$ ). Furthermore, Figure 3  
262 reveals a quite different shape of the LR vs  $\beta_a$  functional relationships (green curves) at different wavelengths. At 355  
263 and 532 nm the curve is concave, with quite similar LR maxima of the fitting curve (54.3 and 53.8 sr at approximately  
264  $\beta_a = 4 \times 10^{-4} \text{ km}^{-1} \text{ sr}^{-1}$  and  $2 \times 10^{-3} \text{ km}^{-1} \text{ sr}^{-1}$ , respectively). At 1064 nm the curve is conversely monotonic, with a flex  
265 point at  $\beta_a = 3\text{-}4 \times 10^{-4} \text{ km}^{-1} \text{ sr}^{-1}$ . A larger data dispersion also characterizes the results at  $\lambda = 355$  and 532 nm (LR values  
266 from 10 to 90 sr) in comparison to  $\lambda = 1064$  nm (LR in the range 18 – 80 sr, except for a minor number of outliers).  
267 This translates into different LR relative errors at UV, VIS and infrared (IR) wavelengths. At 1064, dLR/LR slightly  
268 decreases for increasing backscatter, with values around 35%. At the shorter wavelengths, it increases as a function of  
269  $\beta_a$ , with a large (>40%) relative error for values of  $\beta_a > 2 \times 10^{-3} \text{ km}^{-1} \text{ sr}^{-1}$ .

270 To insert our results into a more general context, we compared the derived, model-based weighted-LR values to some  
271 LR data reported in the literature (Table 4). In particular, we selected some of the works using the aerosol model  
272 developed to invert the Calipso lidar data (Omar et al., 2009). This latter considers six different aerosol sub-types: clean  
273 continental (CC), clean marine (CM), dust (D), polluted continental (PC), polluted dust (PD), and smoke (S). Our  
274 model-derived LR at 532 nm falls in the middle of the range (35-70 sr) fixed by the Calipso CC and PC aerosol classes.  
275 The work by Papagianopoulos et al. (2016), in which the LR values are adjusted accordingly to EARLINET  
276 observations, reports a LR range at 532 nm of 47-62 sr. At the same wavelength, the aerosol range defined by the  
277 LIVAS climatology (Lidar climatology of Vertical Aerosol Structure for space-based lidar simulation studies, Amiridis  
278 et al., 2015), is 54-64 sr. In both cases, our model seems to be closer to the LR values of CC aerosol type, which is  
279 compatible to our intention to simulate clean-to-moderately polluted continental aerosol type. At 532 nm, our LR value  
280 is also reasonably in between the CC and PC LR values derived by Omar et al. (2009), but again closer to the CC LR  
281 value. The very small decrease of LR values between 532 and 355 nm estimated by LIVAS for the CC aerosol is also  
282 consistent with our results. Similarly, our model predicts a lower mean LR in the near IR with respect to the green, in  
283 agreement with results of Amiridis et al. (2015) in CC conditions and not to those in polluted conditions. Table 4 also  
284 includes the continental aerosol LR values estimated in the work of Düsing et al. (2018) through comparison between  
285 airborne in situ and ground-based lidar measurements. Our model is in good agreement with their LR values at 355 and  
286 532 nm. At 1064 nm, the algorithm developed by Düsing et al. (2018) provided a value of LR around 15 sr. On the  
287 other hand, in the same study the authors found that, rather, a value of LR = 30 sr gives the better accord between their  
288 Mie and lidar-based  $\alpha_a$ , this value being closer to our model-derived one at 1064 nm (LR = 37.7). The difference  
289 between these two values is explained by the authors to be probably due to the estimation of the aerosol particle number  
290 size distribution, a critical parameter for a reliable modeling of aerosol particle backscattering.

291 As a last added value of the outcome from our model-based results, we derive here and provide in Table 5 extinction-to-  
292 volume conversion factors,  $c_v = V_a/\alpha_a$  (e.g., Ansmann et al., 2010) at three different wavelengths (355, 532 1064 nm),  
293 and compare these to similar outcomes from other studies. To our knowledge, values of continental particles  $c_v$  at three  
294 wavelengths are only available in Mamouri and Ansmann (2017). Note that  $c_v$ , is also proportional, through the particle  
295 density  $\rho_a$ , to the inverse of the so-called ‘mass-to-extinctions efficiency’ (MEE, i.e.  $\alpha_a/(V_a*\rho_a)$ ) a parameter important in  
296 several aerosol-related applications (e.g. the estimation of particulate matter mass from satellite AOT or in modules of  
297 global circulation and chemical transport models to compute aerosol radiative forcing effects, Hand and Malm, 2007).  
298 For convenience, model-derived MEE values are also included in Table 5.

299



### 300 **3 Evaluation of the model performances and potential of its application**

301 In this section, we evaluate the capability of the model results to reproduce ‘real’ aerosol conditions and explore the  
302 potential of the proposed model-based ALC inversion in producing quantitative geophysical information. In particular:

303 - In Section 3.1 we compare our simulations to real observations of independent backscatter and extinction coefficients  
304 made by different EARLINET Raman lidars (Bösenberg et al., 2001, Pappalardo et al., 2014).

305 - In Section 3.2, our model results are used to invert measurements acquired by some ALCs systems operating within  
306 ALICE-NET, which networks several ALCs systems (Nimbus CHM15k by Lufft) located across Italy and run by  
307 Italian research institutions and environmental agencies. Here we use data from some of these systems to derive the  
308 aerosol optical and physical properties (e.g. the aerosol optical thickness, AOT, and the aerosol volume and mass).

#### 309 **3.1 Comparison of the modelled aerosol optical properties to EARLINET measurements**

310 As mentioned EARLINET Raman stations perform coordinated measurements two days per week following a schedule  
311 established in 2000 (Bösenberg et al., 2003). Overall, the EARLINET database includes the following categories:  
312 ‘climatology’, ‘CALIPSO’, ‘Saharan dust’, ‘volcanic eruptions’, ‘diurnal cycles’, ‘cirrus’, and ‘others’ (forest fires,  
313 photo smog, rural or urban, and stratosphere). To be comparable to our results, we used EARLINET  $\beta_a$  and  $\alpha_a$   
314 coefficients at 355 nm and at 532 nm within the quality assured (QA) ‘climatology’ category (Pappalardo et al., 2014).  
315 However, note that additional data filtering was necessary to screen out residual, likely unreliable values within this  
316 QA-‘climatology’ category. In particular, we only selected those EARLINET QA data further satisfying the following  
317 criteria:

- 318 -  $\beta_a$  and  $\alpha_a$  coefficients evaluated independently, i.e. only obtained using the Raman method (Ansmann et al.,  
319 1992);
- 320 -  $\beta_a$  and  $\alpha_a > 0$ ;
- 321 -  $LR < 100$ ;
- 322 - Relative errors on  $\beta_a$  and  $\alpha_a < 30\%$ .

323

324 Then, we selected those sites in Europe expected to be mostly impacted by ‘continental’ aerosols and having the largest  
325 datasets (e.g., at least 100 points) at 355 and 532 nm. Overall, 5 sites satisfied these conditions (Table 6), and namely  
326 Madrid (Spain), Potenza and Lecce (Italy), Leipzig and Hamburg (Germany). Finally, being interested in continental  
327 conditions here, we filtered out those measurements dates affected by desert dust at the measuring sites, i.e. we removed  
328 from our ‘model-measurement comparison data set’ all the dates within the EARLINET ‘climatology’ category also  
329 belonging to the EARLINET ‘Saharan dust’ category.

330 Figure 4 depicts the results of the model-measurements comparison at the sites fulfilling our requirements in terms of  
331 LR vs  $\beta_a$  at  $\lambda=355$  nm (the corresponding results at  $\lambda=532$  nm, including Madrid in place of Hamburg, are given in  
332 Appendix C, Figure C1). The colored area represents the model-simulated data range, while the color code indicates the  
333 absolute number of simulated values (i.e. counts) in each  $\beta_a$  - LR pair. The EARLINET-measured values are reported as  
334 black open circles. Note that, being the model simulations performed over an altitude range 0-5 km (see Section 2.1)  
335 only those simulations corresponding to the altitude range ( $\Delta z$ ) covered by the measurements at each EARLINET  
336 station was taken into account here. Figure 4 shows the model results to well encompass the measured LR vs  $\beta_a$  data,  
337 with few measurements outside the modeled range (most of the exceptions are found for Potenza). Statistically, the

338 highest number density of simulated data well fits the observations, with the exception of Hamburg (Figure 4a), which  
339 however has the lowest number of measured data (it is not an EARLINET station any longer, see Table 6).

340 In Figure 5 the previous results at  $\lambda=355$  nm are converted in terms of ‘mean’ LR per bin of  $\beta_a$  for both model (blue)  
341 and observations (red, again, only  $\beta_a$  bins containing at least 1% of the total modeled data were considered). This view  
342 shows that there is a general good agreement between the modeled and the measured LR values, and in their variation  
343 with  $\beta_a$ . Some major deviations are found for Potenza and are further discussed in the following. The model-  
344 measurements accordance shown in Figure 5 was evaluated in quantitative terms by computing mean LR relative  
345 differences at both  $\lambda = 355$  and  $532$  nm, i.e., we derived  $([(LR_{mod} - LR_{meas}) / LR_{meas}] * 100)$  values, where  $LR_{mod}$  and  
346  $LR_{meas}$  are the lidar ratio values computed by model and derived by lidar measurements, respectively. These values are  
347 reported in Table 7 for each considered EARLINET station, together with the measurements-based mean LR in each  
348 observational site (computed weighting the number of observations per  $\beta_a$ -spaced bins).

349 Results in Figure 5 and Table 7 also give some hints on the capability of the aerosol type assumed (and its admitted  
350 ranges of variability) to reproduce ‘real’ continental aerosol conditions in different sites across Europe. In fact, the four  
351 continental sites selected with our criteria are still expected to be partially impacted by different aerosol types.

352 - A good agreement between the model and the observations in terms of LR mean values is found for Hamburg (Figure  
353 5a), with mean LR differences of the order of 5% (Table 7). Still, the measured LR values have a high variability and  
354 their distribution is positioned towards high values of  $\beta_a$  ( $1 \times 10^{-3}$  to  $4 \times 10^{-3}$   $\text{km}^{-1} \text{sr}^{-1}$ ). This could be due to the presence  
355 of different aerosols types as slightly polluted marine and polluted aerosol (Matthias and Bösenberg, 2002).

356 - A good accord for Leipzig (Fig. 5c) also indicates that this site is mostly dominated by ‘pure’ continental particles. In  
357 fact, the distribution of observed LR points in Fig. 4, which covers  $\beta_a$  values ranging from  $2 \times 10^{-4}$  to  $3 \times 10^{-3}$   $\text{km}^{-1} \text{sr}^{-1}$ , is  
358 well centered to the modeled simulations highest density (counts > 40). Table 7 shows that at both wavelengths mean  
359 discrepancies with LR measurements keep well below 10%.

360 The highest differences in Fig. 5 are found in some southern Europe EARLINET sites:

361 - In Lecce (Fig. 5b), the best agreement between model and observations is found for the lowest values of  $\beta_a$  (between  
362  $9 \times 10^{-4}$  to  $1 \times 10^{-3}$   $\text{km}^{-1} \text{sr}^{-1}$ , see Table 7). Also, the increase from 10% to 18% in the discrepancies at 355 and 532 nm  
363 indicates some model problems in correctly reproducing the spectral variability of the optical properties, suggesting  
364 some mismatch between modeled and real aerosol sizes in this site (see discussion below).

365 - In Potenza (Fig. 5d), a significant difference between the mean LR curves emerges for  $\beta_a$  values  $> 6 \times 10^{-4}$   $\text{km}^{-1} \text{sr}^{-1}$ ,  
366 with observed LR values lower than those simulated here.

367 These discrepancies could be due to the influence of marine aerosols at both stations (De Tomasi et al., 2006, Mona et  
368 al., 2006, Madonna et al., 2011), which is expected to produce lower LR values for high values of  $\beta_a$  (e.g. BG01). In  
369 fact, Madrid shows better performances, with dLR/LR comparable to those in Leipzig.

370 To provide some insight into the reasons of the model-measurements differences at LC and PO sites, some specific  
371 model sensitivity tests have been performed and are reported in Appendix D. In particular, for Lecce, we found that  
372 better agreement between the observed and simulated LR vs  $\beta_a$  behavior at 355 nm is obtained by reducing the  
373 variability range of  $N_{tot}$  (from  $500 - 10000 \text{ cm}^{-3}$  to  $500 - 5000 \text{ cm}^{-3}$  at ground). This indicates that LC is likely affected  
374 by cleaner continental aerosol type conditions. The sensitivity simulations done for understanding the mismatches with  
375 Potenza measurements show that an extension of the variability range of the coarse mode radius is needed to reproduce

376 the observed decrease of LR for increasing backscatter (Figure 5d). This suggests the presence of coarse particles larger  
 377 than those assumed in such clean continental environment (Appendix D). This is compatible with the suspect of marine  
 378 air contamination, although at this stage we are not able to exclude additional contamination of coarser particle of soil  
 379 origin.

380 Overall, mean LR differences between our ‘average-continental’ model and data at selected continental sites in Europe  
 381 keep lower than 20% (Table 7), this indicating it reasonably well reproduces the clean-to-moderately polluted  
 382 continental aerosol conditions we intended to simulate.

### 383 3.2 Model results application to Nimbus CHM15-k ALC measurements

384 To test and validate the model-based inversion methodology, we used the derived functional relationships (Section 2.2)  
 385 to invert and analyze the measurements of some ALICENET ALCs (Lufft CHM15k systems). These instruments are  
 386 biaxial ceilometers that emit laser pulses at 1064 nm (Nd:YAG-laser, class M1) with a typical pulse energy of 8  $\mu$ J and  
 387 a pulse repetition rate of about 6500 Hz. The instruments have a specified range of 15 km and full overlap at around  
 388 1500 m (Heese et al., 2010). The manufacturer provides the overlap correction functions ( $O(z)$ ) for each system. As  
 389 shown recently by Wiegner and Geiß (2012) and Wiegner et al. (2014), a promising strategy to retrieve the aerosol  
 390 backscatter coefficient from ALC measurement is adopting the forward solution of the Klett inversion algorithm (Klett  
 391 1985). This solution requires a known calibration constant of the system (i.e. absolute calibration,  $c_L$ ) and an  
 392 assumption on the LR. The advantage with respect to the backward solution is that calibration is not affected by the low  
 393 SNR in the upper troposphere and it is needed occasionally. Furthermore, starting close to the surface, the data retrieval  
 394 allows resolving aerosol layers in the boundary layer even if their optical depth is high. The forward solution of the  
 395 Klett inversion algorithm is thus adopted here. For convenience, we report here the equations used within our procedure  
 396 to obtain  $\beta_a$  from ALC measurements, which are also described in Wiegner and Geiß (2012, equations 1 – 3):

$$397 \quad \beta_a(z) = \frac{Z(z)}{LR N(z)} - \beta_m(z) \quad (11)$$

398 with

$$399 \quad Z(z) = LR z^2 P(z) \exp\left[-2 \int_0^z (LR \beta_m - \alpha_m) dz'\right] \quad (12)$$

400 and

$$401 \quad N(z) = c_L - 2 \int_0^z Z(z') dz'. \quad (13)$$

402 Here,  $\beta_m$  and  $\alpha_m$  are the molecular backscatter and extinction coefficients calculated from climatological, monthly air  
 403 density profiles and  $z^2 P(z)$  is the ALC range ( $z$ ) corrected signal ( $P$ ) (also referred to as RCS), that is the raw data  
 404 obtained by the considered ALCs. As anticipated, knowledge of the calibration constant  $c_L$  is needed to solve eq. 13  
 405 (and thus 11, forward solution). In our analysis of ALC daily records, the constant  $c_L$  has been obtained by the  
 406 “backward approach” (Rayleigh calibration) applied to night-time, cloud-free ALC signal averaged over 1 or 2 hours at  
 407 75 m height resolution. This allows for using the best  $c_L$  retrieval (that is the night-time, lowest noise one), in the  
 408 forward solution of the lidar equation, which guarantees operating over the best signal to noise range of the ALC signal.

### 409 3.2.1 Model-based retrieval of aerosol optical properties

410 Operatively, inversion of the aerosol properties,  $\alpha_a(z)$  and  $\beta_a(z)$ , is performed using an iterative technique, since we need  
411 to correct the backscatter signal at each altitude  $z$  for extinction losses. The iterative procedure is stopped when  
412 convergence in the integrated aerosol backscatter ( $IAB = \sum_0^{z_{cal}} \beta_a(z)$ ) is reached (e.g. BG01). At each step, aerosol  
413 extinction is derived using the functional relationship  $\alpha_a = \alpha_a(\beta_a)$  of Table 3.

414 An example of the outcome of this retrieval methodology is depicted in Figure 6. It shows the time-height (24h, 0 - 6  
415 km) contour plot of  $\alpha_a$  retrieved at 1064 nm during a whole day of measurements (June 26, 2016) performed by the  
416 ALICENET system of Aosta San Christophe (ASC, 45.8°N, 7.4°E 570 m a.s.l., Northern Italy, Figure 7a). Time and  
417 altitude resolutions are 1 min and 15 m, respectively. Note that ALC data are cloud-screened using the cloud mask of  
418 the Lufft firmware.

419 The aerosol optical thickness (AOT) is obtained vertically integrating the ALC-derived  $\alpha_a(z)$  from the surface up to a  
420 fixed height  $z_{AOT}$ , above which the aerosol contribution is assumed to be negligible. In Figure 6, the ALC-derived AOT  
421 values at 1064 nm (pink curve, with a temporal resolution of 5 min) is superimposed to the extinction contour.  
422 Reference AOT values from a co-located sun-sky radiometer (a Prede POM-02 system) are shown by orange circles.  
423 These were extrapolated at 1064 nm from the instrument 1020 nm-channel using the Angström exponent derived fitting  
424 AOT values at all the radiometer wavelengths. This example illustrates the very good performances of our model-  
425 assisted inversion scheme, and the capability of this approach to extend to nighttime the (daylight-only) radiometer  
426 observations.

427 To evaluate the performances of our model-assisted retrieval of  $\alpha_a(z)$  over a more statistically significant dataset, the  
428 same approach illustrated in Figure 6 was applied to a longer record in the ASC site, plus Nimbus CHM-15k ALC  
429 datasets from two additional ALICENET sites: San Pietro Capofiume (SPC, 44°39N, 11°37E, 10 m a.s.l.) and Rome  
430 Tor Vergata (RTV, 41.88°N, 12.68°E, 100 m a.s.l.). The location of the instruments is shown in Figure 7a (red circles),  
431 while some information on system types and site characteristics is given in Table 8. The data analyzed here were  
432 collected during the following periods: April 2015 – June 2017, June 2012 – June 2013 and February 2014 – September  
433 2015, for ASC, SPC and RTV, respectively.

434 In those sites, reference AOTs were collected by three co-located sun-sky radiometer, and namely using two SKYNET  
435 Prede sun-sky radiometers at ASC and SPC (POM-02L and POM-02, respectively, [www.euroskyrad.net](http://www.euroskyrad.net)) and an  
436 AERONET Cimel CE 318-2 instrument operational at RTV (<https://aeronet.gsfc.nasa.gov>, Rome Tor Vergata station,  
437 data level 2.0). Only AOT values between 0.01 and 0.2 at 1064 nm were considered. This range allows for excluding  
438 the data points with 1064 nm-AOT lower than the sunphotometer expected accuracy ( $dAOT=0.01$ ) and those where we  
439 found aerosol extinction to cause significant deterioration of our ALC signal. Overall a total of 1237, 268, 850 AOT  
440 pairs were analyzed at ASC, SPC and RTV, respectively.

441 Also note that, although CHM-15k data are already corrected for the  $O(z)$  function provided by the manufacturer, the  
442 variation of the ALC internal temperature was shown to lead to  $O(z)$  differences up to 45% in the first 300 m above  
443 ground (Hervo et al., 2016). For this reason, in our analyses the lowest valid altitude of the CHM-15k for both the SPC  
444 and RTV systems was fixed to be about 400 m. A linear fit of the first two valid ALC points is then used to extrapolate  
445  $\alpha_a(z)$  down to the ground ( $z_0$ ). Conversely, due to the optimal characterization down to the ground of  $O(z)$  provided by  
446 Lufft for the CHM-15k system installed at ASC, values at  $z_0$  at this site are not those extrapolated but actually those

447 measured. The maximum altitude of aerosol extinction vertical integration to derive the AOT,  $z_{\text{AOT}}$ , was selected as the  
448 first height above 4000 m where the range corrected signal (RCS) has a SNR < 1.

449 Results of the long-term AOT comparison are summarized in Figure 7 and Table 9. For each site under investigation,  
450 Figure 7 shows the histograms of the AOT differences between the hourly-mean coincident AOTs as derived by the  
451 ALCs and measured by the sun-photometers (red curve, corresponding AOT vs AOT scatter plots at the three  
452 considered sites are given in Appendix E). To evaluate the advantage of our approach with respect to more standard  
453 lidar inversions, we also computed AOT differences using two fixed-LR values. In particular, we used LR = 52 sr (i.e.  
454 the value suggested by the E-Profile network, black lines) and LR = 38 sr (i.e. the weighted mean LR value derived  
455 from our model, see Section 3, blue lines). Figure 7 shows that the best agreement is found at ASC. The distribution of  
456 AOT difference has a maximum around 0 for each of the three inversions schemes, with very low dispersion. The full  
457 width at half maximum, FWHM, is in fact around 0.015, and approximately 55% of the data are included in the interval  
458 -0.01 – 0.01, which is even within the expected error of photometric measurement. For SPC and RTV, the red and blue  
459 histograms are peaked around 0, whereas the black ones are shifted, with maxima around 0.01-0.02 and 0.02-0.03 for  
460 SPC and RTV, respectively. These two sites have higher dispersion (FWHM = 0.03), and approximately 30% of the  
461 data are included in the interval -0.01 – 0.01 for the red and blue histograms at both sites, which is probably due to the  
462 different aerosol load affecting the different ALICENET stations. As pointed out by the low value of the average AOT  
463 computed at ASC for the analyzed dataset ( $\langle \text{AOT} \rangle = 0.027$ ), low pollution levels generally characterize this site, with  
464 some exceptions due to wind-driven aerosol transport from the nearby Po valley (Diémoz et al, 2018a, 2018b this issue).  
465 On the contrary, RTV ( $\langle \text{AOT} \rangle = 0.044$ ) and, especially, SPC in the Po Valley ( $\langle \text{AOT} \rangle = 0.076$ ) are characterized by  
466 higher aerosol content and pollution levels, which explain the larger histogram dispersions. Note that the high frequency  
467 of fog events in winter markedly reduces the number of analyzed AOT pairs at SPC site, while some desert dust  
468 affected days at both SPC (e.g., Bucci et al., 2018) and RTV (e.g., Barnaba et al., 2017) were removed from our  
469 datasets (no desert-dust affected dates in ASC).

470 Table 9 summarizes the long-term performances of the model-based procedure in deriving quantitative AOT from the  
471 ALC systems at the three investigated sites. It includes values of the average differences between the ALC-derived and  
472 sunphotometers-measured AOT (both bias,  $\langle \text{dAOT} \rangle$ , and absolute difference  $\langle |\text{dAOT}| \rangle$ , with associated standard  
473 deviations) obtained using both the proposed model-based approach and the fixed-LR inversions. For SPC and RTV  
474 sites, these numbers show that the best ALC–photometer accordancy is reached when employing either the model-based  
475 or the fixed LR=38 sr inversion scheme. In fact, these two approaches have similar performances in terms of mean  
476  $\text{dAOT}$  values ( $\langle |\text{dAOT}| \rangle = 0.011, 0.013$  and  $0.013, 0.014$  for SPC and RTV, respectively), mean percent error  
477 ( $\langle |\text{dAOT}| \rangle / \langle \text{AOT} \rangle = 0.16, 0.19$  and  $0.31, 0.33$ ) and a very low mean relative bias ( $\langle \text{dAOT} \rangle / \langle \text{AOT} \rangle = -0.043, 0.005$   
478 and  $0.088, 0.11$ ). On the other hand, the fixed LR=52 sr retrieval produces an overestimation of AOT in both SPC and  
479 RTV ( $\langle \text{dAOT} \rangle / \langle \text{AOT} \rangle = 0.33$  and  $0.44$ ) with larger discrepancies between retrieved and observed AOTs ( $\langle |\text{dAOT}| \rangle$   
480  $= 0.021$  and  $0.026$ ,  $\langle |\text{dAOT}| \rangle / \langle \text{AOT} \rangle = 0.38$  and  $0.49$ ). For the ASC site, due to the low aerosol content, the  
481 differences among the inversion schemes are almost negligible.

482 Overall, for the three sites, the statistics over the long-term datasets employed showed good results of the model-based  
483 approach with similar behavior of the retrievals with a fixed LR of 38 sr, while a fixed LR value of 52 sr produces an  
484 overestimation of the AOT at SPC and RTV. As different sites have different (and not known a-priori) characteristic LR  
485 values, these results highlight the potential of the model-based approach to derive quite accurate  $\beta_a$  and  $\alpha_a$  coefficients  
486 without the need to choose and fix an arbitrary LR value.

### 487 3.2.2 Model-based retrieval of aerosol volume (and mass)

488 In this section we provide examples of the applicability of the proposed approach to derive air-quality relevant  
489 parameters. In particular, we use the ALC,  $\beta_a$ -retrieved data and the 7<sup>th</sup>-order polynomial fit linking  $\beta_a$  (at  $\lambda = 1064$   
490 nm) to  $V_a$  (see also Table 3 and Figure 2c) to derive the aerosol volume (and mass).

491 The ALC-estimates were firstly compared to aerosol volume derived in situ at the ASC site by two different optical  
492 particle counters (OPCs) on 29<sup>th</sup> December 2016 and 5<sup>th</sup> September 2017. For the case of the 29<sup>th</sup> December 2016, a  
493 TSI Optical Particle Sizer (OPS) 3330 was employed. This instrument has 16 channels that can be programmed to  
494 provide the number concentration at different (and logarithmically spaced) diameter size ranges within the interval 0.3 -  
495 10  $\mu\text{m}$ . Further details can be found in the TSI manual (2011). For the case of the 5<sup>th</sup> September 2017, the Fidas®200s  
496 OPC was used. This spectrometer is able to retrieve high-resolution particle spectra (size measurements between 0.15  
497 and 27  $\mu\text{m}$ , with 32 channels/decade, Pletscher et al., 2016). For both dates, Figure 8 shows the time (x-axis, 24h) vs.  
498 height (left y-axis) contour plots of the ALC-based retrieval of the aerosol volume concentration ( $\text{cm}^3/\text{cm}^3$ ). The OPC-  
499 derived aerosol volume concentration measured at ground-level is reported as a function of time (x-axis) on the right y-  
500 axis (grey curve). The corresponding ALC-derived volume concentration (integrating the ALC data between 0 and 75  
501 m) is shown by a pink curve (same right y-axis). Daily mean volume concentration values derived by OPCs and by  
502 ALC are also plotted (grey cross and pink triangle, respectively). The horizontal bar in the upper part of the figure  
503 indicates the ranges of RH measured in-situ during the analyzed cases.

504 The OPC-to-ALC comparison is certainly affected by intrinsic factors, as differences on the atmospheric layer sampled  
505 (at ground and integrated between 0 and 75 m, for OPC and ALC, respectively) and on the probing methods (in-situ and  
506 remote sensing, dried air sampled by OPC and ambient conditions sampled by the ALC). Furthermore, as mentioned in  
507 Section 4.2.1, a major critical issue of ALC retrievals at low levels is the correction for the overlap function, which  
508 needs to be experimentally characterized and verified for each instrument.

509 These issues are visible in the given example of Figure 8. In fact, in the upper panel, the agreement between the ALC-  
510 derived and the TSI-OPC aerosol  $V_a$  values is good between 0 and 7 UTC. In the following hours both instruments  
511 register an increase of the aerosol volume, although with some discrepancies in absolute values. Starting from 18 UTC,  
512 the ALC derives an aerosol volume concentration higher than the OPC one by a factor of 3-3.5. This disagreement  
513 could be related to both the presence/arrival of fine particles ( $<0.3 \mu\text{m}$ ) not measured by the optical counter (see for  
514 example Diémoz et al., 2018a), or to aerosol hygroscopic effects (increase of volume associated to hygroscopic growth  
515 seen by ALC but not by the OPC which dries the air samples). This latter effect is confirmed by the large RH values  
516 ( $\text{RH} > 90\%$ ) measured after 18 UTC. The lower panel shows a good agreement between the ALC-derived and the Fidas  
517 OPC  $V_a$  values, in particular until 4 UTC and after 16 UTC. Some differences emerge around 7 UTC and between 11  
518 and 15 UTC, where the ALC volume is lower by a factor of 2 compared to the in situ Fidas  $V_a$  values. The smaller  
519 minimum detectable size of the Fidas OPC instrument with respect to the OPS is likely the reason for the better accord  
520 between ALC and OPC  $V_a$  values in this test date. In this case, the effect of RH seems to be less important, and indeed  
521 RH values keep lower than 90%.

522 In general, high RH values ( $\text{RH} \geq 90\%$ ) are known to markedly affect the aerosol mass estimation from remote  
523 sensing techniques and its relationship with ‘reference’ PM<sub>2.5</sub> or PM<sub>10</sub> measurements methods, usually performed in  
524 dried conditions (e. g. Barnaba et al., 2010; Adam et al., 2012, Li et al., 2016, Li et al., 2017). This theme is also  
525 discussed in Diémoz et al. 2018a for the ALC measurement site of Figure 8. Nevertheless, even with the mentioned

526 limitations, results in Fig. 8 well show the potential of the developed method in providing sound values of aerosol  
527 volume, and hence, mass, in average-RH regimes, giving support to more standard PM10 air quality monitoring.

528 To give a further example in this direction, the model-assisted retrievals of aerosol mass over a longer time period were  
529 used to derive daily-mean aerosol mass concentrations (PM10), a measurement typical of air quality stations. To this  
530 purpose, for the two-months period June-July 2012, we derived daily mean values of aerosol volume at the SPC site  
531 using the functional relationships  $V_a = V_a(\beta_a)$ , and then converted these into mass (PM10) using typical values of  
532 aerosol densities ( $\rho_a$ ). Results are shown in Figure 9. It compares the daily average PM10 concentration measured in  
533 situ at SPC by the Italian Regional Environmental Protection Agency (ARPA, red solid curve) and the model-assisted,  
534 ALC-derived daily mass concentration obtained assuming both a fixed particle density  $\rho_a = 2 \text{ g/cm}^3$  (blue dotted curve),  
535 and a range of it between 1.5-2.5  $\text{g/cm}^3$  (shaded area), this range covering approximately the typical  $\rho_a$  values at the  
536 SPC site. Yellow shaded areas indicate the presence of dust events (e.g. Bucci et al., 2018) that are excluded from the  
537 results reported in the next paragraph.

538 More in detail, the daily-mean, ALC-derived mass concentrations were estimated in two steps: 1) estimation of hourly  
539 mass values for the selected height; 2) computation of the daily values through the median of the hourly values. To  
540 guarantee a good daily representativeness, the second step is applied only to those days in which at least 50% of the  
541 hourly values is available in all the following temporal ranges: 00 - 05 UTC, 06 - 11 UTC, 12 - 17 UTC, 18 - 23 UTC.  
542 Note that, due to the uncertainties associated to the  $O(z)$  in the first hundreds of meters (as previously mentioned, the  
543 ALC system at SPC has an old firmware, and its overlap function is not optimally characterized), we used the 225 m  
544 level as more trustworthy to estimate ALC mass concentration. On the other hand, during the considered period of the  
545 year (i.e. June and July), the comparison to ground-level PM10 at SPC is expected to be only slightly affected by this  
546 height difference, particularly in daytime, due to the strong convection within the mixing layer. Possible exception  
547 could be in nocturnal conditions when vertical gradients in the lowermost hundreds of meters can occur. However, our  
548 statistical (3-year) ALC records show the mixing layer height at SPC to descend below 250 m only 4-5 hours per day in  
549 July (usually between 22 and 3 UTC, i.e., when emissions are at a minimum). Overall, Figure 9 confirms a good  
550 agreement between the ALC-derived and the ARPA reference PM10 values, with a correlation coefficient (R) of 0.64.  
551 In fact, mean, absolute mean and relative differences, between the two series are:  $\langle \text{dPM10} \rangle = 2.3 \pm 6.0 \text{ g/cm}^3$ ,  
552  $\langle |\text{dPM10}| \rangle = 4.8 \pm 4.3 \text{ g/cm}^3$  and  $\langle (\text{dPM10}/\text{PM10}) \rangle = 0.14 \pm 0.27$ . This agreement attests that SPC site can indeed be  
553 considered an ‘average’ continental site and suggests the potential of this approach to derive information on aerosol  
554 volume and mass. Still, due to the specificity of each site and to the limited period considered here, these results cannot  
555 be taken as representative of all continental sites at all times. Further studies at different places and over longer time  
556 periods would be necessary to better assess the uncertainty of the proposed retrieval, including uncertainties due to the  
557 variability of ‘continental’ conditions (in terms of particle size distribution, compositions, hygroscopic effects, etc...),  
558 but also of the instrument-dependent performances (e.g. overlap corrections, etc...).

#### 559 4 Summary and Discussion

560 Thanks to their low construction/operation costs and to their capability at providing continuous, unattended  
561 measurements, the use of automated-lidar-ceilometers (ALCs) for aerosol characterization has increased in the recent  
562 years. Several numerical approaches were recently proposed to estimate the aerosol vertical profile either using  
563 ceilometer measurement only, or coupling these with ancillary measurements (e.g., Stachlewska et al., 2010; Flentje et  
564 al., 2010; Wiegner et al., 2012; Wiegner et al., 2014; Cazorla et al., 2017, Román et al., 2018).

565 This work proposes a methodology to retrieve key aerosol properties (as extinction coefficient, surface area and  
566 volume, thus mass) from lidar/ALC measurements using in support the results from a specifically developed aerosol  
567 numerical model to drive the retrievals. In particular, the numerical model uses a "Monte-Carlo" approach to simulate a  
568 large set (20000) of aerosol microphysical properties intended to reproduce the variability of 'average' (clean-to-  
569 moderately polluted) continental conditions, i.e., those expected to dominate over Europe. Based on the assumption of  
570 particle sphericity, relevant computations of aerosol physical (surface area and volume,  $S_a$  and  $V_a$ ) and optical  
571 (backscattering and extinction coefficients,  $\beta_a$  and  $\alpha_a$  through Mie scattering theory) properties were performed at three  
572 commonly used lidar wavelengths (i.e., at the Nd:YAG laser harmonics 355, 532, 1064 nm). Fitting procedures of this  
573 large set (20,000) of  $\beta_a$  vs.  $\alpha_a$ ,  $S_a$  and  $V_a$  data-pairs were then used to derive mean functional relationships linking  $\beta_a$  to  
574  $\alpha_a$ ,  $S_a$  and  $V_a$ , respectively. The model's statistical uncertainties (i. e., those related to the variability of the  
575 microphysical parameters used in input to the computations of the bulk physical/optical properties) associated to these  
576 so-derived mean relationships were found to be within 30% and 40% for  $\beta_a$  vs  $\alpha_a$  and  $\beta_a$  vs  $V_a$ , respectively, while  $\beta_a$  vs  
577  $S_a$  exhibits a larger dispersion (relative standard uncertainty of 40%-70%, depending on  $\beta_a$ ). It is worth mentioning that  
578 these are higher than those associated to the retrievals of aerosol bulk parameters using the complete set of Raman  
579 lidars observations (three aerosol backscattering and two extinction coefficients, i.e., the so called 3+2 approach),  
580 assuming, as in our case, no random uncertainty in the lidar input data. For example, Veseloski et al. (2012) found a  
581 maximum uncertainty of 15% for particle volume and surface area estimation, in the case of 0% random uncertainty in  
582 the lidar input data. Note however, that such multi-wavelength lidar systems are 10 to 20 times more expensive than  
583 ALC systems, need to be operated by highly trained operators, and are rarely run all day round.

584 The model results also allowed exploring the expected dependence of the (continental aerosol) lidar ratio (LR) on  $\beta_a$  at  
585 355, 532 and 1064 nm, and in turn, the mean, 'weighted'-LR value at each wavelength (found to be  $50.1 \pm 17.9$  sr,  $49.6$   
586  $\pm 16.0$  sr and  $37.7 \pm 12.6$  sr, at 355, 532 and 1064 nm respectively). Availability in literature of LR values at 1064 nm  
587 are scarce and its monotonic increase with  $\beta_a$  found in this work (Figure 3) suggests that the use of a fixed LR value for  
588 the inversion of ALC signals should be done with caution and carefully evaluated case by case. A similar, non-  
589 monotonic behavior characterizes the shapes of LR vs  $\beta_a$  curve at 355 and 532 nm.

590 We tested the reliability of our model results in two ways: 1) the model numerical computations were compared to  
591 'real' lidar measurements (specifically selected within the EARLINET database), and 2) the model-assisted retrievals of  
592 aerosol optical (AOT) and physical ( $V_a$ , PM10) properties by real, operational ALC systems were compared to  
593 corresponding 'reference' measurements performed by co-located, independent instrumentation.

594 In particular, in task 1) our simulations were compared to backscatter and extinction coefficients at 532 and 355 nm  
595 independently retrieved by advanced Raman lidar systems operating at different EARLINET sites in Europe (namely  
596 Hamburg and Leipzig in Germany, Madrid in Spain, Lecce and Potenza in Italy). The model simulations were found to  
597 statistically well match the observations (Figures 4, 5 and C1). Mean discrepancies between model and measurement-  
598 based LR were found to be lower than 20%, suggesting a good capability of the assumed aerosol model (and admitted  
599 range of variability) to represent 'real', 'average continental' aerosol conditions in different sites across Europe. Some  
600 differences emerged for so southern Italy EARLINET sites, possibly affected by the influence of marine aerosols,  
601 leading to lower LR values for high values of  $\beta_a$ .

602 For task 2) we applied the proposed model-based inversion to different ALC systems (Lufft CHM-15k), part of the  
603 Italian ALICENET network. We firstly tested the ability of the proposed approach to derive aerosol extinction by  
604 comparing hourly-mean, vertically-integrated  $\alpha_a$  (i.e., hourly mean AOT) derived by three ALC systems to  
605 corresponding AOT measurements from co-located sun-photometers (ALICENET sites of Aosta San Cristophe (ASC),



606 San Pietro Capofiume (SPC) and Rome Tor Vergata (RTV), Figure 7). ALC-sun photometer agreement was found to be  
607 within 30%. Tests on the use of fixed LR were also performed to investigate the advantage of the proposed approach  
608 with respect to more standard ones. To this purpose, we used the (1064 nm) fixed-LR value suggested by the E-Profile  
609 EUMETSAT Program and the ‘weighted mean’ derived from our model (52 sr and 38 sr, respectively). While for the  
610 ASC site negligible differences were found among the three retrieval schemes, for both SPC and RTV sites the best  
611 ALC – sun photometer accordance in AOT is reached when employing the model-based or the fixed LR=38 sr  
612 inversion schemes, with a mean error around 16-19 % and 31-33 % for SPC and RTV, respectively. Applying the fixed  
613 LR value of 52 sr produces an overestimation of the AOTs, with a mean relative bias equal to 33 % and 44 % at SPC  
614 and RTV, respectively. This suggests that, at 1064 nm, the LR value for continental aerosol is lower than the one  
615 assumed by the E-Profile procedure and, more in general, this highlights the advantage of a procedure not requiring an  
616 a-priori, and to some extent arbitrary, choice of the LR value.

617 As a second test in task 2, values of aerosol volume (and mass) derived using the model-assisted ALC retrieval were  
618 compared to in situ aerosol measurements performed by OPCs and PM10 analyzers. A continuous, two-months  
619 comparison (June – July 2012) between daily average aerosol mass concentration as measured in situ and derived by  
620 ALC (in the lowest altitudes) at SPC, showed a mean relative difference of around 15% (Figure 9).

621 Overall, the good results obtained in our validation efforts are encouraging but necessarily related to the specific  
622 conditions at the measuring sites considered and to the characteristics of the instruments employed. They are therefore  
623 not necessarily representative of results obtainable in all European continental sites, and at all times. Further tests using  
624 wider datasets covering a variety of sites and ALC instrumentation would be desirable to better understand potential  
625 and limits of the applicability of the proposed method over the larger scale. An obvious intrinsic limitation is that the  
626 method is dependent on the considered aerosol type which in this study was tuned to reproduce average continental  
627 aerosol conditions. Errors associated to the application of the derived functional relationship might be larger if more  
628 ‘specific’ aerosol conditions (e.g. contamination by sea salt or desert dust particles) affect a given site. In the future, the  
629 information coming from ALC systems with an additional depolarization channel (as tested in the DIAPASON Project,  
630 Gobbi et al., 2018) could be used to force the retrieval to different model schemes (e.g. switching from ‘no dust’ to  
631 ‘dust’ schemes conditions) in the same vertical profile. This will enhance the capabilities of ALCs to operatively  
632 estimate and characterize the aerosol optical properties (e.g. Gasteiger and Freudenthaler, 2014).

633 Additionally, although our validation exercises returned results well within the uncertainties related to the model  
634 statistical variability alone (i.e., the relative errors associated to the mean functional relationships), the expected total  
635 uncertainty to be associated to the method should include terms that have not been specifically addressed in this work,  
636 as for example the instrumental error itself.

637 On the other hand, the proposed approach has the main advantage of allowing the operational (i.e. 24/7) retrieval of  
638 fairly reliable, remote sensing profiles of aerosol optical ( $\beta_a$ ,  $\alpha_a$ ) and physical ( $S_a$ ,  $V_a$ ) properties (with associated  
639 uncertainties and limitations) by means of relatively simple and robust instruments. This could temporally and spatially  
640 complement the information coming from more advanced lidar networks (for example, the Raman channel of multi-  
641 wavelength system cannot be used in daylight conditions) and, more in general, could represent a valid option to  
642 deliver, in quasi real time, the 3D aerosol fields useful for operational air quality (e.g. integration of the in situ surface  
643 measurements) and for meteorological and climate monitoring (e.g. aerosol-cloud interaction and aerosol transport and  
644 dispersion processes).

645

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662

663 **Data availability**

664 AERONET Rome-Tor Vergata sun photometer AOT data were downloaded from the AERONET web page  
665 (AERONET, 2018). SKYNET sun photometer AOT data were downloaded from the SKYNET webpage (SKYNET,  
666 2018). EARLINET backscattering and extinction coefficients were downloaded from the EARLINET webpage  
667 (EARLINET, 2018). ALICENET ALC raw data are available upon request at [alicenet@isac.cnr.it](mailto:alicenet@isac.cnr.it).

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951 **Table 1. Aerosol parameter values as reported in literature for continental-type aerosols.**

<i>Reference</i>	$r_1$ ( $\mu\text{m}$ )	$r_2$ ( $\mu\text{m}$ )	$r_3$ ( $\mu\text{m}$ )	$N_1/N_{\text{tot}}$	$N_2/N_{\text{tot}}$	$N_3/N_{\text{tot}}$	$m_{r_1}$	$m_{r_2}$	$m_{r_3}$	$N_{\text{tot}}$	Aerosol type
	$\sigma_1$	$\sigma_2$	$\sigma_3$	(%)	(%)	(%)	$m_{\text{im}_1}$	$m_{\text{im}_2}$	$m_{\text{im}_3}$	( $\text{cm}^{-3}$ )	
<i>Whitby (1978)</i> <sup>1</sup>	0.008 1.6	0.034 2.1	0.46 2.2	0.56	0.44	$4 \times 10^{-4}$	-	-	-	1800	Clean continental
<i>D'Almeida et al. (1991)</i> <sup>2</sup>	0.012 2.0	0.029 2.24	0.471 2.51	0.06	0.94	$2 \times 10^{-6}$	1.75 0.44	1.53 0.012	1.53 0.008	20000	Average continental
<i>Hess et al. (1998)</i> <sup>2</sup>	0.012 2.0	0.021 2.24	0.471 2.51	0.56	0.44	$0.3 \times 10^{-4}$	1.75 0.44	1.53 0.012	1.53 0.008	15300	Average continental
<i>Barnaba and Gobbi (2004a)</i> <sup>1</sup>	0.007-0.012 1.7-2.0	0.021-0.077 2.03-2.24	0.403-0.5 2.11-2.24	6.1-54.2	45.8-93.9	(2-26.1) $\times 10^{-4}$	1.25-2.00 0.07-1.00	1.53 $6 \times 10^{-3}$	1.53 $8 \times 10^{-3}$	$10^3 - 10^4$	
<i>Omar et al. (2009)</i> <sup>1</sup>	-	0.093-0.10 1.53-1.61	0.68-0.76 1.9-2.1	-	0.999-1	(0.02-3) $\times 10^{-4}$	-	1.38-1.40 (0.1-6.3) $\times 10^{-3}$	1.40-1.46 (3.4-6.3) $\times 10^{-3}$	-	Clean and polluted continental
<i>Levy et al. (2007)</i> <sup>2</sup>	0.018 2.0	0.005 2.97	0.5 2.97	1	$1 \times 10^{-7}$	$1 \times 10^{-13}$	1.75 0.44	1.53 $6 \times 10^{-3}$	1.53 $8 \times 10^{-3}$	-	
<i>Barnaba et al. (2007)</i> <sup>1</sup>	-	0.05-0.1 1.35-1.70	0.4-0.5 1.5-2.0	-	0.98-0.99	0.01-0.02	-	1.35-1.55 (2.5-20) $\times 10^{-3}$	1.53-1.6 (1.0-80) $\times 10^{-4}$	$1-3 \times 10^3$	Continental - coastal
<i>Amiridis et al. (2015)</i> <sup>1</sup>	-	0.03-0.9 1.6-2.2	0.47-0.69 1.9-2.5	-	1	(4-8) $\times 10^{-7}$	-	1.42-1.45 (2.3-6) $\times 10^{-3}$	1.45-1.53 (2.3-6) $\times 10^{-3}$	-	Clean and polluted continental

952 <sup>1</sup>The refractive index is at  $\lambda=532$  nm.

953 <sup>2</sup>The refractive index is at  $\lambda=550$  nm.

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956 **Table 2. Variability ranges used in this study. Values refer to ground and dry conditions (see text for details).**

Parameter	Mode I	Mode II	Mode III
$r_i$ ( $\mu m$ )	0.005 - 0.03	0.03 - 0.1	0.3 - 0.5
$\sigma$	1.35 - 1.7	1.35 - 1.7	1.5 - 2.4
$N/N_{i\ tot}$ (%)	10 - 60	40 - 90	0.01 - 0.5
$m_{r_i}$ (355 nm)	1.40 - 1.80	1.40 - 1.70	1.50 - 1.60
(532 nm)	1.40 - 1.80	1.40 - 1.70	1.50 - 1.60
(1064 nm)	1.42 - 1.82	1.37 - 1.66	1.50 - 1.60
$m_{im_i}$ (355 nm)	$1 \times 10^{-2} - 0.47$	$1 \times 10^{-4} - 0.010$	$1 \times 10^{-4} - 0.02$
(532 nm)	$9 \times 10^{-3} - 0.44$	$1.2 \times 10^{-4} - 0.012$	$1 \times 10^{-4} - 0.01$
(1064 nm)	$9 \times 10^{-3} - 0.44$	$1.5 \times 10^{-4} - 0.015$	$1 \times 10^{-4} - 0.005$
$N_{tot}$ ( $cm^{-3}$ )	500 - 10000		

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**Table 3. Parameters of the Seventh-Order Polynomial Fits ( $y=a_0+a_1x+a_2x^2+a_3x^3+a_4x^4+a_5x^5+a_6x^6+a_7x^7$ ) for  $\lambda = 1064$  nm, with  $x=\log(B_a)$  (in  $\text{km}^{-1} \text{sr}^{-1}$ ) and  $y=\log(\alpha_a, S_a, \text{ or } V_a)$  in ( $\text{km}^{-1}, \text{cm}^2/\text{cm}^3$  and  $\text{cm}^3/\text{cm}^3$ , respectively).**

<b>Functional relationship at 1064 nm</b>	<i>Extinction coefficient</i>	<i>Surface area</i>	<i>Volume</i>
$a_0$	3.797837507651898	12.019452592845141	-5.314834128998254
$a_1$	3.294032541389781	30.825966279368547	2.500484347793244
$a_2$	0.962603336867675	24.518531616019207	-1.196109537503000
$a_3$	0.241796629870675	10.625241994796593	-1.583236058579546
$a_4$	0.064609145804688	2.634051072085453	-0.681801883947768
$a_5$	0.017721752150233	0.373150843707711	-0.145232662646142
$a_6$	0.002722551625862	0.027971628176431	-0.015471229968392
$a_7$	0.000157245409783	0.000854381337164	-0.000658925756875

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963 Table 4. Mean weighted LR at 355, 532 and 532 nm derived in this work and comparison to the corresponding aerosol  
 964 subtypes (clean continental, CC, and polluted continental, PC) from relevant literature.

<i>LR (sr)</i>	$\lambda = 355 \text{ nm}$	$\lambda = 532 \text{ nm}$	$\lambda = 1064 \text{ nm}$
<i>Omar et al., (2009) (Calipso aerosol model)</i>	-	70 ± 25 (PC)	30 (PC)
	-	35 ± 16 (CC)	30 (CC)
<i>Amiridis et al. (2015) (LIVAS database)</i>	59.5* (PC)	64 (PC)	-
	56.5* (CC)	54 (CC)	-
<i>Papagiannopoulos et al., (2016) (EARLINET measurements)</i>	-	62 ± 10 (PC)	-
	-	47 ± 4 (CC)	-
<i>Düsing et al., (2018) (in-situ and lidar measurements)</i>	55	55	30; 15**
<i>This work</i>	50.1 ± 17.9	49.6 ± 16.0	37.7 ± 12.6

965 \* derived using the extinction-related and backscatter-related Ångström exponents given by Amiridis et al. (2013)

966 \*\* see the explanation in the text for the two different values

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**Table 5. Extinction-to-volume conversion factors,  $c_v = V_a/\alpha_a$  (and corresponding ‘mass-to-extinctions efficiency’ values,  $MEE = \alpha_a/(V_a * \rho_a)$ , given assuming  $\rho_a = 2 \text{ g/cm}^3$ ) of continental particles as derived from our model at different wavelengths**

<i>Reference</i>	<i>c<sub>v</sub> [10<sup>6</sup>m] (corresponding MEE, [m<sup>2</sup>g<sup>-1</sup>])</i>			<i>Notes</i>	
	<i>Wavelength [nm]</i>	355	532		1064
<i>Hess et al. (1998)</i>	-	-	0.35 (1.43)	-	OPAC, clean continental model
<i>Hess et al. (1998)</i>	-	-	0.28 (1.79)	-	Opac, polluted continental model
<i>Barnaba and Gobbi (2004b)</i>	-	-	0.18 (2.78)	-	Continental model
<i>Ansmann et al. (2011b)</i>	-	-	0.18 (2.78)	-	Germany, fine aerosol fraction
<i>Lewandosky et al. (2010)</i>	-	-	-	0.77 – 2 (0.25-0.65)	Mexico city basin
<i>Sicard et al. (2012)</i>	-	-	0.26 (1.92)	-	AERONET, Spain
<i>Mamouri and Ansmann (2017)</i>	0.17 (2.94)	0.30 (1.67)	0.96 (0.52)	-	Germany, continental anthropogenic pollution
<i>Mamouri and Ansmann (2017)</i>	0.23 (2.17)	0.41 (1.22)	1.41 (0.35)	-	Cyprus, continental anthropogenic pollution
<i>Mamali et al. (2018)</i>	-	0.14, 0.24 (3.57, 2.03)	-	-	Cyprus, fine non-dust aerosol fraction
<i>This work</i>	0.12 (4.17)	0.19 (2.63)	0.60 (0.83)	-	Continental (clean to moderately polluted)

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**Table 6. Main characteristics of the dataset of the EARLINET continental sites considered in this study. The listed dataset refers to the data downloaded from the EARLINET site (last access on the 11<sup>th</sup> of January 2018).**

<i>Station</i>	<i>Number of points at 355 and at 532 nm</i>	<i>Altitude range (<math>\Delta z</math>, in km)</i>	<i>Period</i>
<i>Lecce (LC)</i> 40.33 N, 18.10 E, 30 m a.s.l.	1012 – 109	1 – 4	Aug2007 – Oct2013
<i>Leipzig (LE)</i> 51.35 N, 12.43 E, 90 m a.s.l.,	5186 – 4549	1.5 – 4	Aug2008 – Sept2016
<i>Potenza (PO)</i> 40.6 N, 15.72 E, 760 m a.s.l.	1244 – 219	1.5 – 4	May2000 – Aug2009
<i>Hamburg (HH)</i> 53.57 N, 9.97 E, 25 m a.s.l.	243 – n.a.	0.5 – 4	Apr2001 – Oct2002
<i>Madrid (MA)</i> 40.45 N, 3.73 E, 669 m a.s.l.	n.a. – 492	0.5 – 4	Jun2006 – Jun2008

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976 **Table 7 Mean LR discrepancies between our model results and EARLINET measurements and weighted LR at 355 and 532**  
 977 **nm for the considered EARLINET stations.**

<i>Station</i>	$[(LR_{mod} - LR_{meas})/LR_{meas}] * 100$		<i>EARLINET</i> weighted LR (sr)	
	$\lambda = 355 \text{ nm}$	$\lambda = 532 \text{ nm}$	$\lambda = 355 \text{ nm}$	$\lambda = 532 \text{ nm}$
<i>LC</i>	10	18	51.8	44.5
<i>LE</i>	6	9	52.6	51.0
<i>PO</i>	17	7	44.9	57.2
<i>HH</i>	5	-	53.3	-
<i>MA</i>	-	6	-	54.2

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980 **Table 8. Main characteristics of the ALC and co-located sun-sky radiometer equipment located at the considered**  
 981 **ALICENET sites.**

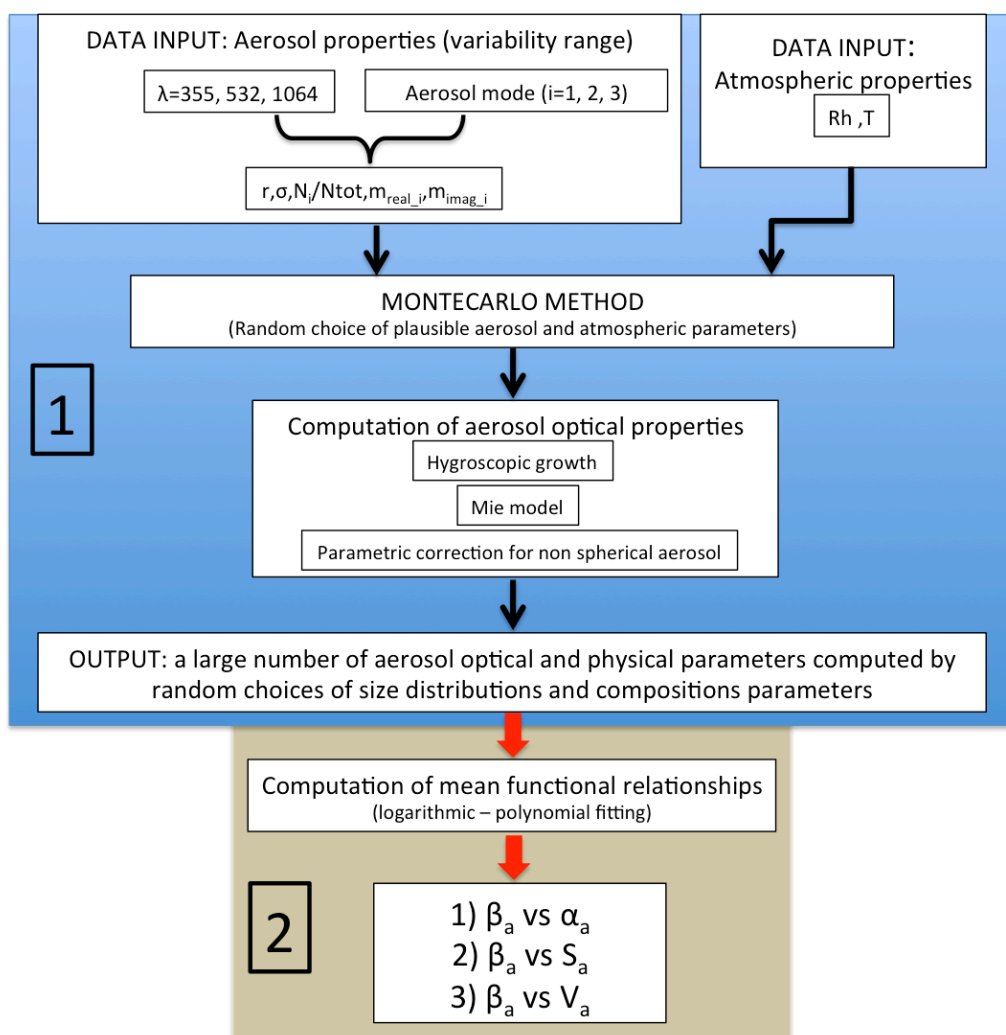
	<i>Site type</i>	<i>ALC model</i>	<i>ALC firmware</i>	<i>Sun photometer model</i>
<i>ASC</i>	alpine	Nimbus CHM150104	0.743	POM-02
<i>SPC</i>	rural	Nimbus CHM110115	0.556	POM-02L
<i>RTV</i>	semi-rural	Nimbus CHM070052	0.720	CIMEL CE-318

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984 **Table 9. Results of the comparison between the AOT measured by sun-photometers and the one derived by ALCs (model-**  
 985 **based and fixed LR inversion schemes) at three ALICENET stations. Mean differences (expressed in terms of  $\langle dAOT \rangle =$**   
 986  **$\langle (AOT_{\text{ceil}} - AOT_{\text{phot}}) \rangle$ ,  $\langle |dAOT| \rangle$  (module),  $\langle dAOT/AOT \rangle$  and  $\langle |dAOTI/AOT| \rangle$ ) are reported with their standard deviations.**

<i>ALICENET sites</i>	$\langle dAOT \rangle$	$\langle  dAOT  \rangle$	$\langle dAOT/AOT \rangle$	$\langle  dAOTI/AOT  \rangle$
<b>ASC</b>				
<i>Variable LR from our model</i>	$-0.004 \pm 0.015$	$0.010 \pm 0.013$	$-0.25 \pm 0.57$	$0.31 \pm 0.35$
<i>LR = 52 sr</i>	$0.002 \pm 0.021$	$0.009 \pm 0.015$	$0.31 \pm 0.58$	$0.33 \pm 0.35$
<i>LR = 38 sr</i>	$-0.004 \pm 0.014$	$0.009 \pm 0.012$	$-0.23 \pm 0.43$	$0.30 \pm 0.32$
<b>SPC</b>				
<i>Variable LR from our model</i>	$-0.001 \pm 0.020$	$0.013 \pm 0.016$	$-0.005 \pm 0.28$	$0.19 \pm 0.20$
<i>LR = 52 sr</i>	$0.021 \pm 0.026$	$0.026 \pm 0.02$	$0.33 \pm 0.35$	$0.38 \pm 0.26$
<i>LR = 38 sr</i>	$-0.003 \pm 0.019$	$0.011 \pm 0.014$	$-0.043 \pm 0.24$	$0.16 \pm 0.18$
<b>RTV</b>				
<i>Variable LR from our model</i>	$0.004 \pm 0.020$	$0.014 \pm 0.014$	$0.11 \pm 0.49$	$0.33 \pm 0.30$
<i>LR = 52 sr</i>	$0.016 \pm 0.023$	$0.021 \pm 0.018$	$0.44 \pm 0.59$	$0.49 \pm 0.45$
<i>LR = 38 sr</i>	$0.003 \pm 0.019$	$0.013 \pm 0.013$	$0.088 \pm 0.460$	$0.31 \pm 0.27$

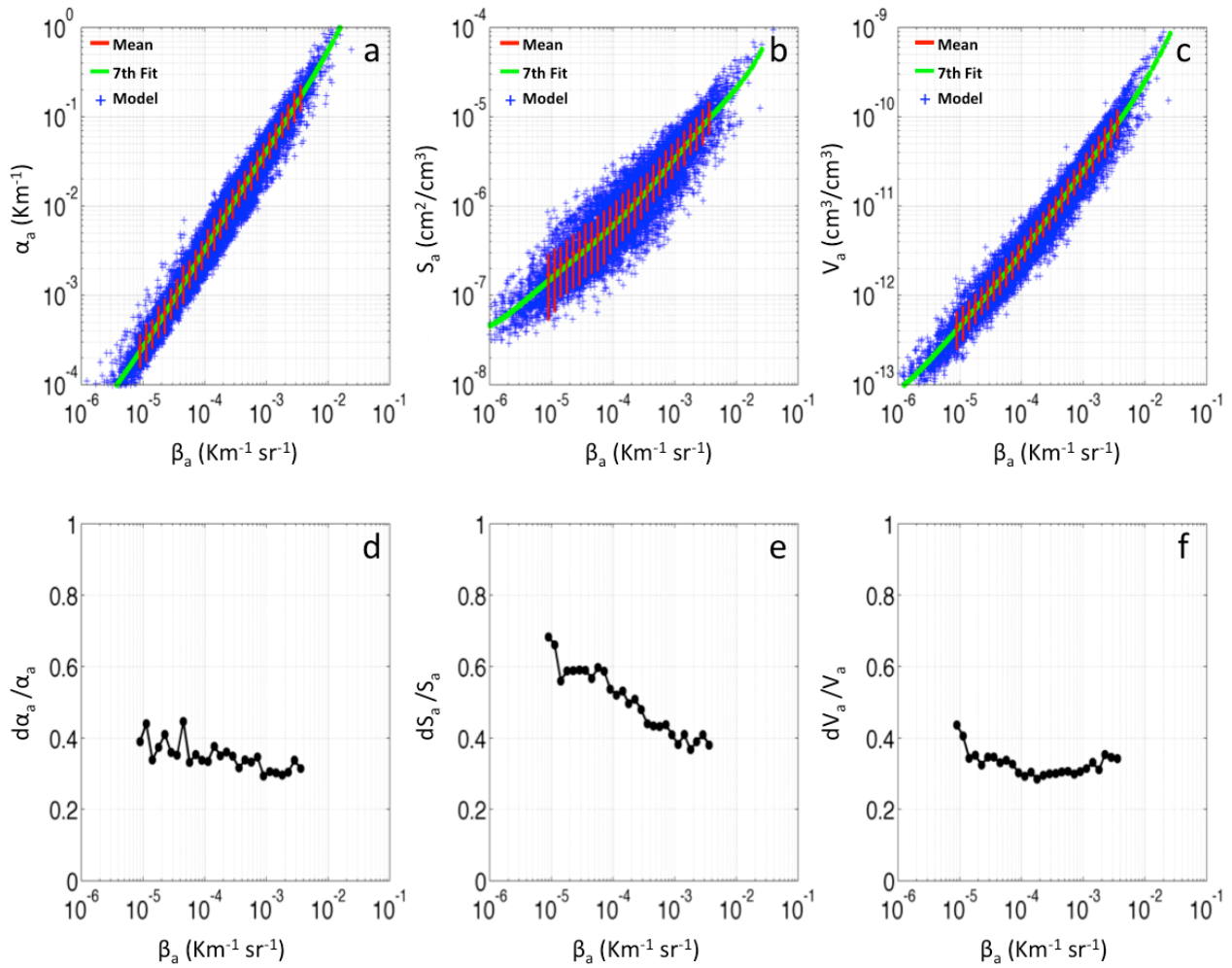
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991 **Figure 1. Schematic of the two-step model structure developed to obtain, as a result, functional relationships between the**  
 992 **aerosol backscatter ( $\beta_a$ ) and the aerosol extinction, surface area and volume ( $\alpha_a$ ,  $S_a$  and  $V_a$ , respectively).**

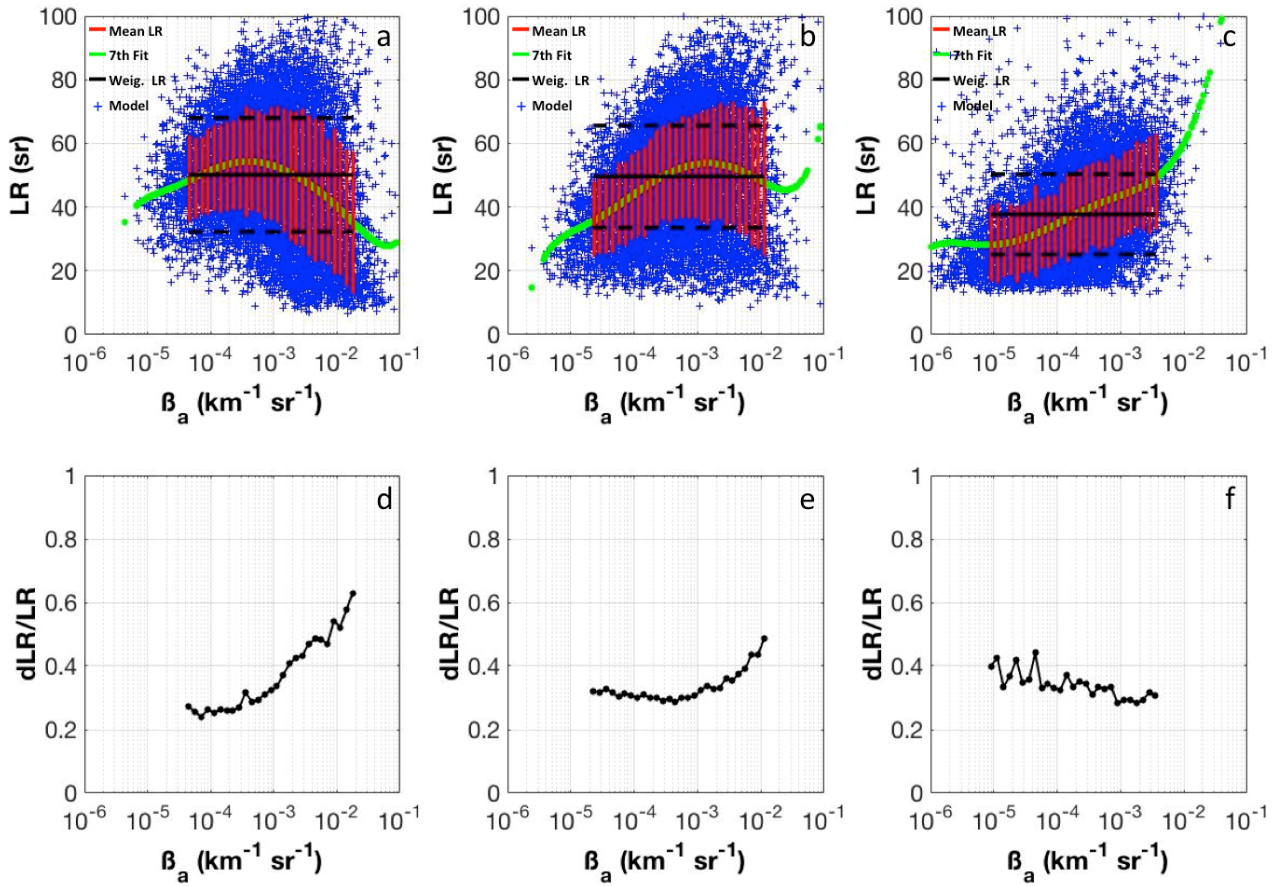
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995 **Figure 2.** Scatterplots of a)  $\alpha_a$  ( $\text{km}^{-1}$ ), b)  $S_a$  ( $\text{cm}^2/\text{cm}^3$ ) and c)  $V_a$  ( $\text{cm}^3/\text{cm}^3$ ) vs backscatter  $\beta_a$  ( $\text{km}^{-1} \text{sr}^{-1}$ ) and relevant relative  
 996 errors (panels d, e, f, respectively) as derived from 20000 model computations (blue points) at  $\lambda = 1064 \text{ nm}$ . Red dots and  
 997 error bars are the average values per decade of  $\beta$  and their standard deviations, green lines are the 7th-order polynomial fit  
 998 curve of the 20000 points.

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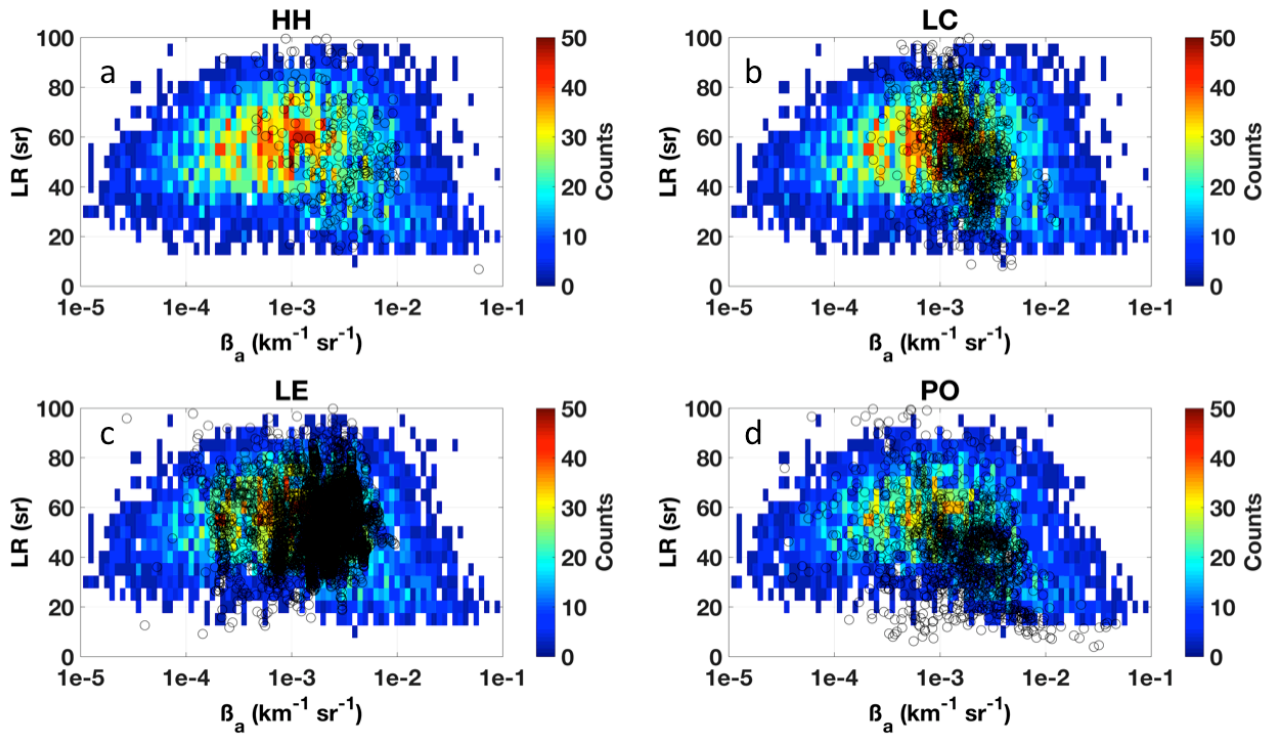
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Figure 3. Upper plots: scatterplots of LR (sr) versus  $\beta_a$  ( $\text{km}^{-1} \text{sr}^{-1}$ ) at: a) 355 nm; b) 532 nm; c) 1064 nm (blue points). The 7th-order polynomial fit curve (green lines) and the average values per decade of  $\beta$  together with their standard deviations (red points and red vertical bars, respectively) are also reported. Horizontal black lines are mean values of the ‘weighted-LR’ and  $\pm 1$  s. d. (solid and dotted lines, respectively). Lower plots: relative errors associated with the model-derived LR at d) 355 nm; e) 532 nm; f) 1064 nm.



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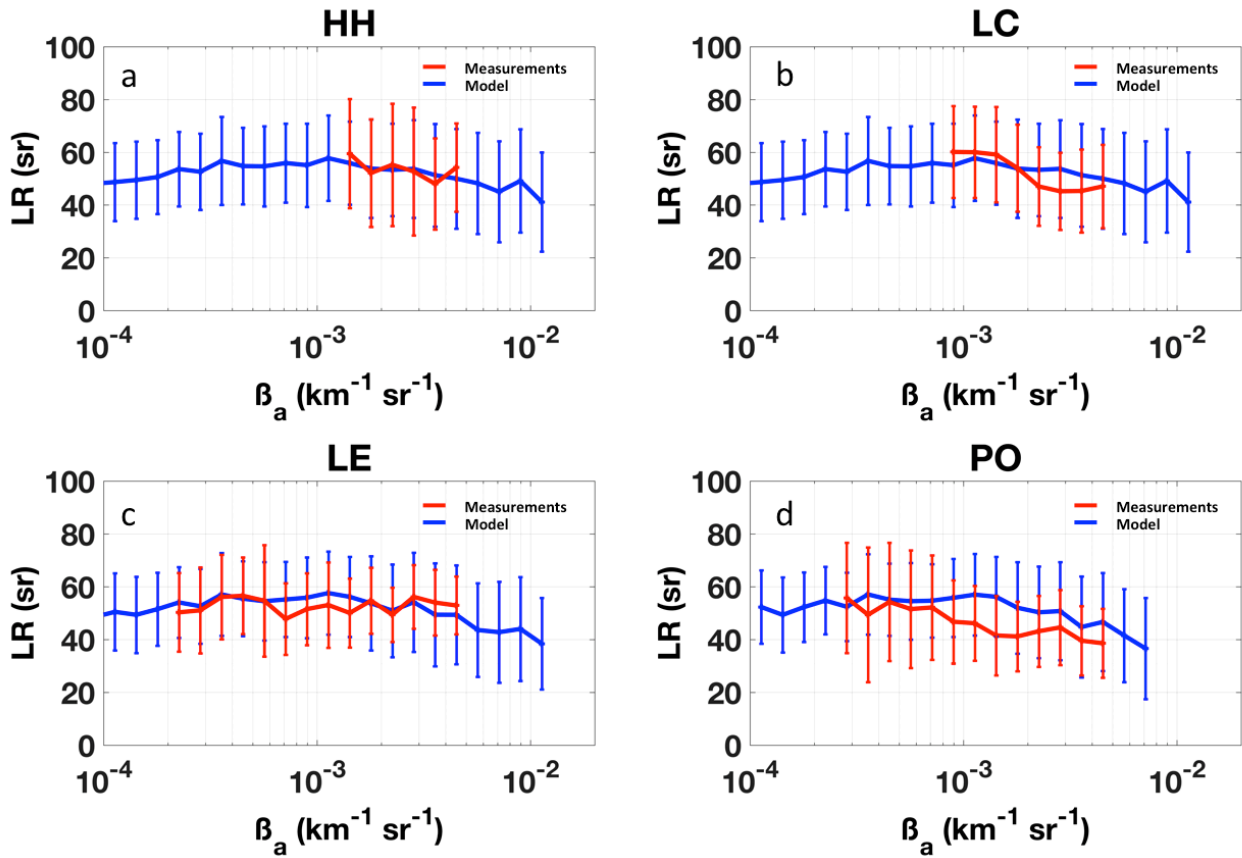
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Figure 4. Scatterplots of LR (sr) versus  $\beta_a$  ( $\text{km}^{-1} \text{sr}^{-1}$ ) at 355 nm as simulated by our model (colored region) and measured by EARLINET lidars (black open circles) in Hamburg (Germany) (a), Lecce (Italy) (b), Leipzig (Germany) (c) and Potenza (Italy) (d). The color area is the region of simulated values, the color code indicating the number of simulated values in each  $\beta_a$  -LR pair (see legend). In particular, the color-2D histogram is computed using a semi-logarithmic box consisting of 10 equally spaced bins per decade of  $\beta_a$  in the x-axis and 5 spaced LR values in the y-axis.



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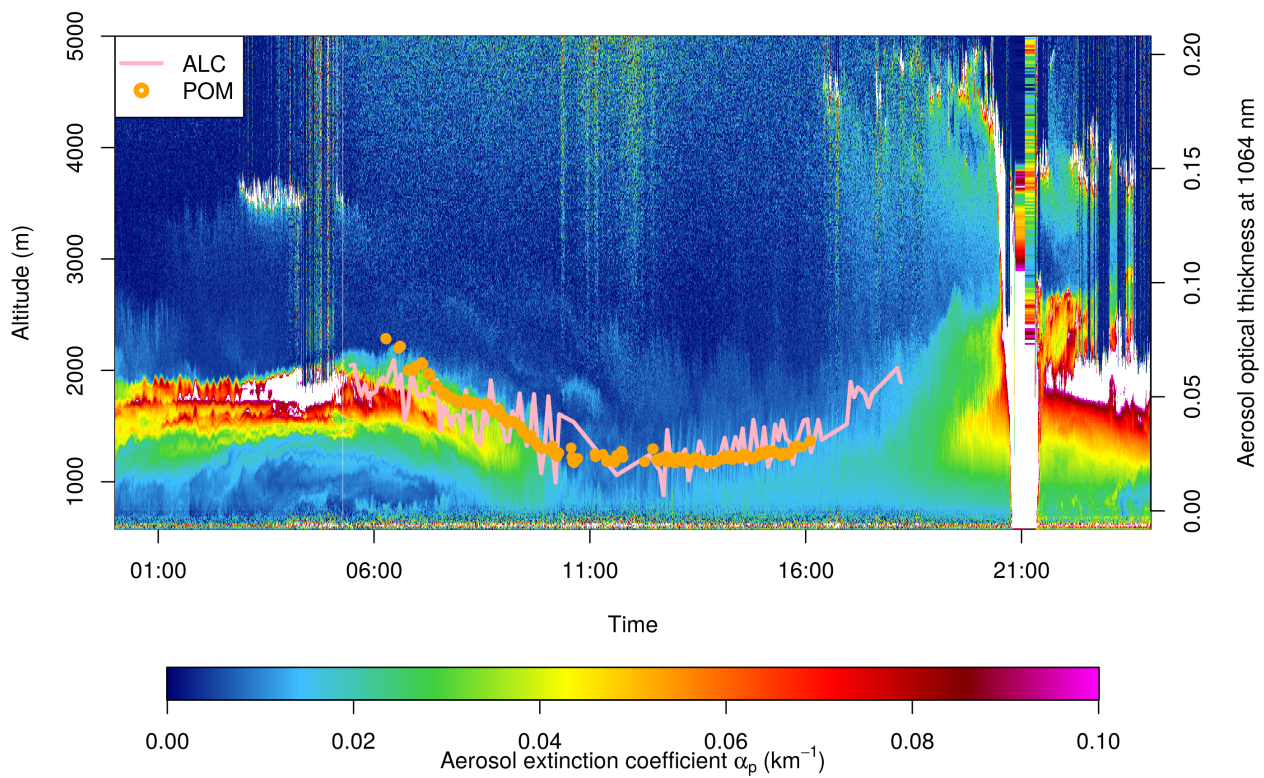
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Figure 5. Model-simulated (blue) and lidar measured (red) LR vs  $\beta_a$  mean curves at 355 nm calculated per 10 equally spaced bins per decade of  $\beta_a$  in a) Hamburg, b) Lecce, c) Leipzig, and d) Potenza EARLINET lidar station. Vertical bars are the associated standard deviations.





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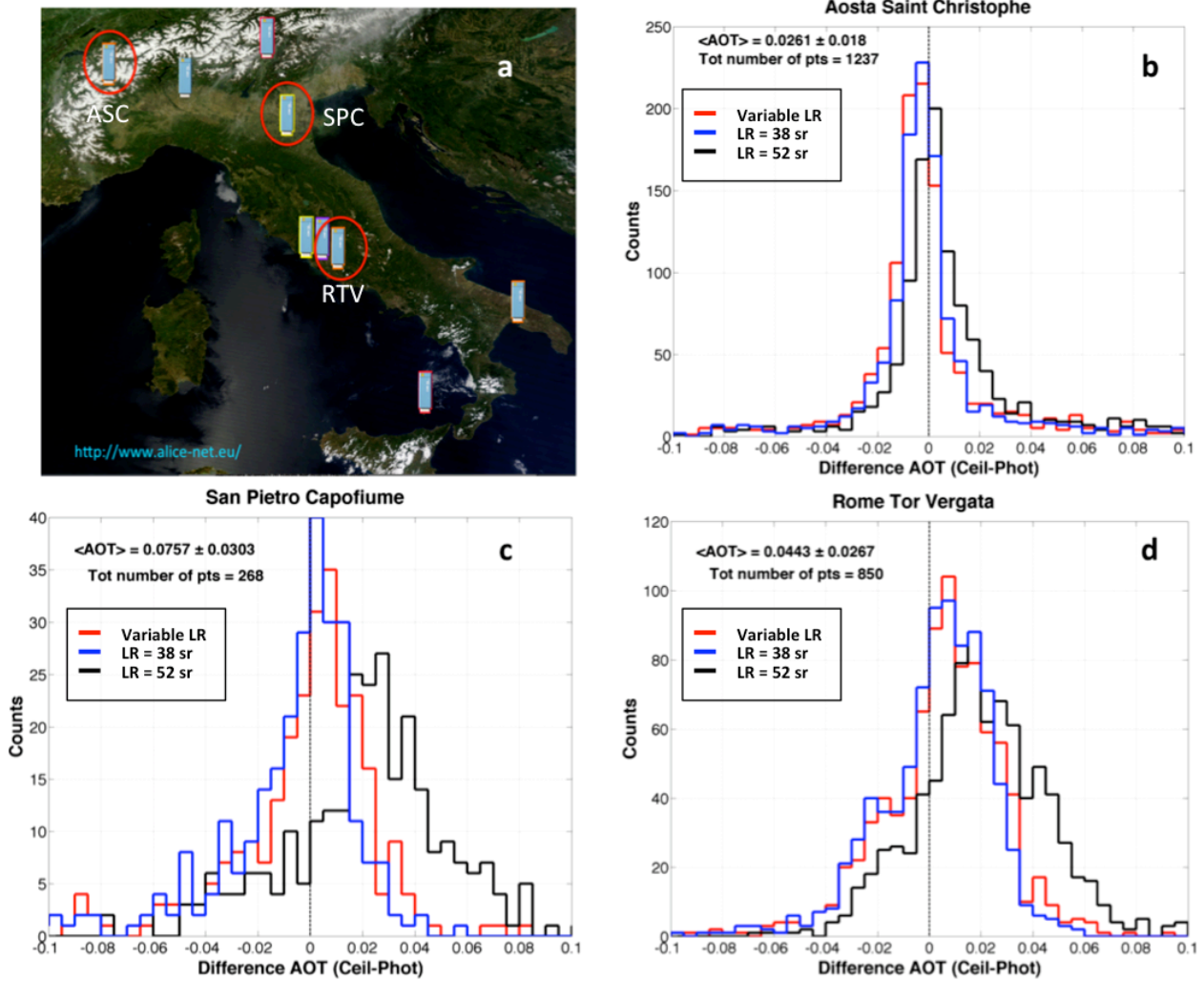
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**Figure 6. Time-height cross-section of the aerosol extinction coefficients  $\alpha_a$  [ $\text{km}^{-1}$ ], as derived at 1064 nm on 26 June 2016 by the ALICENET ALC of Aosta San Christophe (Northern Italy). The orange circle points and the pink line are the AOT values (right y-axis, panel b) measured by a co-located POM-02L radiometer and estimated from the ALC following our approach.**



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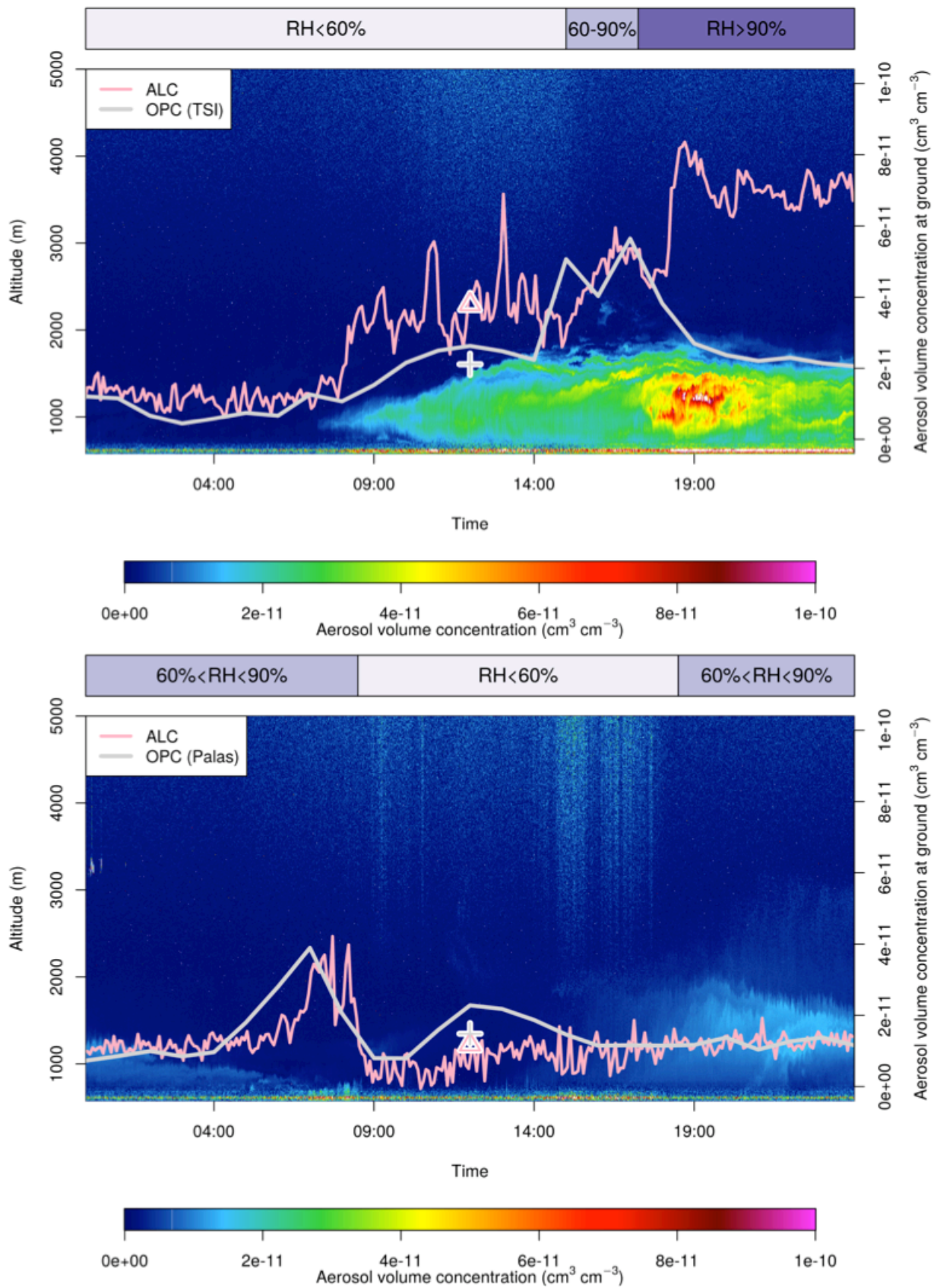
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Figure 7. a) geographical map of the ALC network ALICENET. The red circles highlight the selected sites for this study: Aosta San Christophe (ASC), San Pietro Capofiume (SPC) and Rome Tor Vergata (ASC). b-d) histograms of the differences between the hourly-mean coincident AOTs at 1064 nm as derived by ALCs and measured by photometers, at ASC, SPC and RTV, respectively. The different colors (red, blue and black) depict the different inversion schemes: model-based inversion scheme, LR = 38 sr and LR = 52 sr, respectively. In each panel the values of the average measured AOT (and its associated standard deviation) and of the number of considered pairs are also reported.



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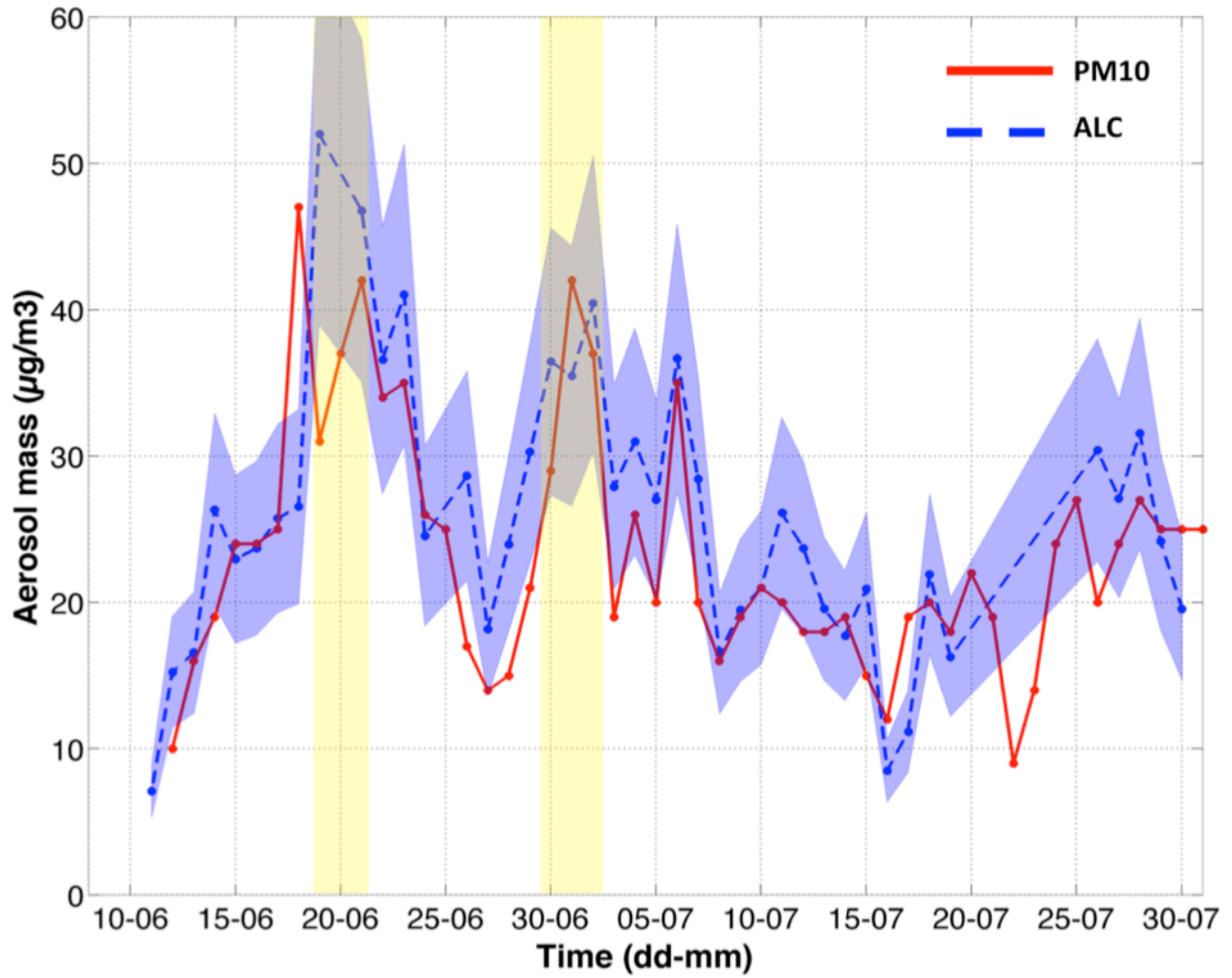
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Figure 8. Time-height cross-section of the aerosol volume concentration at Aosta San Christophe for 29 December 2016 (upper panel) and 05 September 2017 (lower panel). The right y-axis reports the volume concentration measured at surface through TSI and Fidas@200s OPCs (upper and lower panels, grey curves) and the ALC-derived volume concentration at 75

1038 m (pink curves). The grey crosses and the pink triangles refer to the daily mean aerosol volume value derived by OPCs and  
1039 ALC measurements, respectively. The horizontal bars in the upper part of the panel indicate the ranges (RH<60%,  
1040 60%<RH<90% and RH>90%, respectively) of the measured in-situ RH during the analyzed days.  
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1043 Figure 9. Daily-resolved aerosol mass concentration at SPC, for the period June – July 2012, estimated from ALC-derived  
 1044 aerosol volume data at 225 m a.s.l. converted into mass using a fixed particle density  $\rho_a = 2 \mu\text{g}/\text{m}^3$  (blue dotted line) and a  
 1045 variable  $\rho_a$  between  $1.5 - 2.5 \mu\text{g}/\text{m}^3$  (shaded blue area). The red solid line is the daily PM10 concentration as measured by the  
 1046 local Air Quality agency (ARPA). Vertical yellow shaded stripes indicate the presence of dust events.

## 1047 **Appendix A: Model sensitivity tests**

1048 To evaluate the proposed continental model configuration (hereafter CM0) and discuss its sensitivity to the variability  
1049 of the employed parameters, an overview of the impact on the model results produced by changing the limit of the  
1050 variability ranges of these parameters (i.e. using different model configuration, CMX) is given in this section.

1051 The varied model (CMX-CM0) mean difference on the considered optical property (OP) has been quantified through  
1052 the following equation:

$$1053 \langle \frac{dOP}{OP} \rangle = \left( \frac{1}{N_{bin}} \right) \cdot \sum_{i=1}^{N_{bin}} [(\langle OP_{CMX} \rangle_i - \langle OP_{CM0} \rangle_i) / \langle OP_{CM0} \rangle_i], \quad (A.1)$$

1054 where  $N_{bin}$  is the total number of defined bins of  $\beta_a$ .

1055 The results of the mean differences of  $\alpha_a$  and LR for different ranges of  $\beta_a$  and for the whole  $\beta_a$  interval are reported on  
1056 table A1, where relevant sensitivity cases (i.e. relative mean difference greater than 1%) at  $\lambda=355$  nm have been taken  
1057 into account.

1058 CM1 refers to a model configuration without the first aerosol mode ( $N_1\%=0$ ). The overall decrease on the values of  $\alpha_a$   
1059 and LR (around 3-4%) is due to the sum of significant and opposite effects for low and high values of  $\beta_a$  where  
1060  $\langle d\alpha_a/\alpha_a \rangle$  and  $\langle dLR/LR \rangle$  are of the order of -6% and 8%, respectively. Removing the coarser aerosol mode ( $N_3\%=0$ ),  
1061 causes positive mean values for  $\langle d\alpha_a/\alpha_a \rangle$  and  $\langle dLR/LR \rangle$  of the order of 5% (sensitivity case CM2). In this case, the  
1062 largest impact is observed for the  $\beta_a$  range between  $2 \times 10^{-4}$  and  $2 \times 10^{-3}$   $\text{km}^{-1}\text{sr}^{-1}$ .

1063 An opposite result is obtained by decreasing the upper bound of the  $r_2$  variability range ( $r_2=0.03 - 0.05$   $\mu\text{m}$ , CM3). In  
1064 fact also this model configuration leads to lower  $\alpha_a$  and LR ( $\langle d\alpha_a/\alpha_a \rangle$  and  $\langle dLR/LR \rangle$  are equal to -6%, approximately).  
1065 In this case, the variation on the  $r_2$  parameter affects the higher ranges of  $\beta_a$  ( $\beta_a=2 \times 10^{-4}$ - $2 \times 10^{-2}$   $\text{km}^{-1}\text{sr}^{-1}$ ). Higher modal  
1066 radii for the coarse-mode particle ( $r_3=1 - 1.2$   $\mu\text{m}$ ) in CM4 configuration leads to the increase of the contribution of  
1067 model-generated points with higher  $\beta_a$  and causes lower values of  $\alpha_a$  and LR ( $\langle d\alpha_a/\alpha_a \rangle$  and  $\langle dLR/LR \rangle$  are equal to -5%,  
1068 approximately) only for high values of  $\beta_a$  ( $\beta_a=2 \times 10^{-3}$ - $2 \times 10^{-2}$   $\text{km}^{-1}\text{sr}^{-1}$ ), whereas the effect over the whole  $\beta_a$  range is  
1069 around -1%.

1070 The CM5 configuration accounts for the presence of more absorbing particles in the first aerosol mode, where the lower  
1071 bound of  $m_{1im}$  has been increased by a factor of 10 ( $m_{1im}=0.1$ - $0.47$ ). This produces a significant effect only for the  
1072 lower values of  $\beta_a$  ( $\beta_a=2 \times 10^{-5}$ - $2 \times 10^{-4}$   $\text{km}^{-1}\text{sr}^{-1}$ ), with an increase of  $\alpha_a$  and LR of approximately 4%. On the contrary,  
1073 increasing the lower bound of the real part of the second aerosol mode refractive index ( $m_{2r}=1.55$ - $1.70$ ) has a large  
1074 impact on the considered parameters. In fact, the CM6 configuration largely underestimates both  $\alpha_a$  and LR (around -  
1075 15% for both parameters) for all  $\beta_a$  ranges.

1076 The CM7 configuration refers to the impact of the total number of particles at the ground ( $N_{tot}$ ). In this case, decreasing  
1077 the upper bound of the variability range of  $N_{tot}$  by a factor of 2 ( $N_{tot}=500$ - $5000$   $\text{cm}^{-3}$ ) lowers the mean values of  $\alpha_a$  and  
1078 LR of around 5%. Nevertheless, this effect is totally due to the contribution of the  $\beta_a$  values between  $2 \times 10^{-3}$  and  $2 \times 10^{-2}$   
1079  $\text{km}^{-1}\text{sr}^{-1}$ , where  $\langle d\alpha_a/\alpha_a \rangle$  and  $\langle dLR/LR \rangle$  are around -10%. Assuming no increase with altitude for  $\sigma_{1,2}$  (sensitivity case  
1080 CM8) produces relevant differences on the mean values of  $\alpha_a$  and LR. In CM8, the overall overestimation of these two  
1081 parameters is quite limited ( $\langle d\alpha_a/\alpha_a \rangle = 6.3$  and  $\langle dLR/LR \rangle = 6.4$ ), whereas a large and opposite impact is observed for

1082 low and high values of  $\beta_a$ . In fact,  $\langle d\alpha_a/\alpha_a \rangle$  ( $\langle dLR/LR \rangle$ ) is equal to -14.1 (-13.9) and 18.5 (19.0) for  $\beta_a = 2 \times 10^{-5} - 2 \times 10^{-4}$   
 1083 and  $\beta_a = 2 \times 10^{-5} - 2 \times 10^{-4} \text{ km}^{-1} \text{sr}^{-1}$ , respectively. As explained by Barnaba et al. (2007), the dependence of  $\sigma_{1,2}$  to the  
 1084 altitude can be associated to the fact that, when increasing the distance from the main aerosol sources, the particle  
 1085 processing is more efficient.

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1087 **Table B1. Mean differences of  $\alpha_a$  and LR between different model sensitivity cases and the proposed continental model**  
 1088 **configuration.**

Model configura- tion	$\beta_a$ ( $\text{km}^{-1} \text{sr}^{-1}$ ) $2 \times 10^{-5} - 2 \times 10^{-4}$		$\beta_a$ ( $\text{km}^{-1} \text{sr}^{-1}$ ) $2 \times 10^{-4} - 2 \times 10^{-3}$		$\beta_a$ ( $\text{km}^{-1} \text{sr}^{-1}$ ) $2 \times 10^{-3} - 2 \times 10^{-2}$		$\beta_a$ ( $\text{km}^{-1} \text{sr}^{-1}$ ) $2 \times 10^{-5} - 2 \times 10^{-2}$	
	$\langle d\alpha_a/\alpha_a \rangle$ (%)	$\langle dLR/LR \rangle$ (%)	$\langle d\alpha_a/\alpha_a \rangle$ (%)	$\langle dLR/LR \rangle$ (%)	$\langle d\alpha_a/\alpha_a \rangle$ (%)	$\langle dLR/LR \rangle$ (%)	$\langle d\alpha_a/\alpha_a \rangle$ (%)	$\langle dLR/LR \rangle$ (%)
CM1 ( $N_1\%=0$ )	-6.2	-6.4	3.1	3.2	7.8	7.9	-3.7	-3.5
CM2 ( $N_3\%=0$ )	4.7	4.9	8.6	8.9	2.8	2.7	5.3	5.4
CM3 ( $r_2=0.03 - 0.05 \mu\text{m}$ )	-2.0	-1.7	-10.3	-10.2	-8.9	-8.2	-6.7	-6.4
CM4 ( $r_3=1.0 - 1.2 \mu\text{m}$ )	<1	<1	-2.1	-2.0	-5.24	-5.3	-1.2	-1.0
CM5 ( $m_{lim}=0.1-0.47$ )	4.3	4.2	<1	<1	<1	<1	1.8	1.8
CM6 ( $m_2=1.55 - 1.70$ )	-10.9	-10.9	-16.2	-16.3	-18.9	-19.1	-15.3	-15.3
CM7 ( $N_{TOT}=500-5000$ )	<1	<1	<1	<1	-11.2	-10.7	-3.7	-3.5
CM8 ( $\sigma_1, \sigma_2$ constant)	-14.1	-13.9	6.4	6.1	18.5	19.0	6.3	6.4

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1091 **Appendix B: Model-based functional relationships at 355 and 532 nm**

1092 The parameters of the seventh-order polynomial fit used to derive the functional relationships between  $\log(x)$  and  $\log(y)$   
 1093 (where  $x = \beta_a$  and  $y = \alpha_a, S_a$  or  $V_a$ ) at  $\lambda = 355$  and  $532$  nm are reported in Tab. A1 and Tab. A2, respectively.

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1095 **Table A1. Parameters of the Seventh-Order Polynomial Fits ( $y = a_0 + a_1x + a_2x^2 + a_3x^3 + a_4x^4 + a_5x^5 + a_6x^6 + a_7x^7$ ) for  $\lambda = 355$  nm, with**  
 1096  **$x = \log(\beta_a)$  (in  $\text{km}^{-1} \text{sr}^{-1}$  unit) and  $y = \log(\alpha_a, S_a, \text{ or } V_a)$  in ( $\text{km}^{-1}, \text{cm}^2/\text{cm}^3$  and  $\text{cm}^3/\text{cm}^3$ , respectively).**

Functional relationship at 355 nm	Extinction coefficient	Surface area	Volume
$a_0$	3.797837507651898	12.019452592845141	-5.314834128998254
$a_1$	3.294032541389781	30.825966279368547	2.500484347793244
$a_2$	0.962603336867675	24.518531616019207	-1.196109537503000
$a_3$	0.241796629870675	10.625241994796593	-1.583236058579546
$a_4$	0.064609145804688	2.634051072085453	-0.681801883947768
$a_5$	0.017721752150233	0.373150843707711	-0.145232662646142
$a_6$	0.002722551625862	0.027971628176431	-0.015471229968392
$a_7$	0.000157245409783	0.000854381337164	-0.000658925756875

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**Table A2. Parameters of the Seventh-Order Polynomial Fits ( $y = a_0 + a_1x + a_2x^2 + a_3x^3 + a_4x^4 + a_5x^5 + a_6x^6 + a_7x^7$ ) for  $\lambda = 532$  nm, with  $x = \log(B_a)$  (in  $\text{km}^{-1} \text{sr}^{-1}$  unit) and  $y = \log(\alpha_a, S_a, \text{ or } V_a)$  in ( $\text{km}^{-1}, \text{cm}^2/\text{cm}^3$  and  $\text{cm}^3/\text{cm}^3$ , respectively).**

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$a_7$	0.000157245409783	0.000854381337164	-0.000658925756875

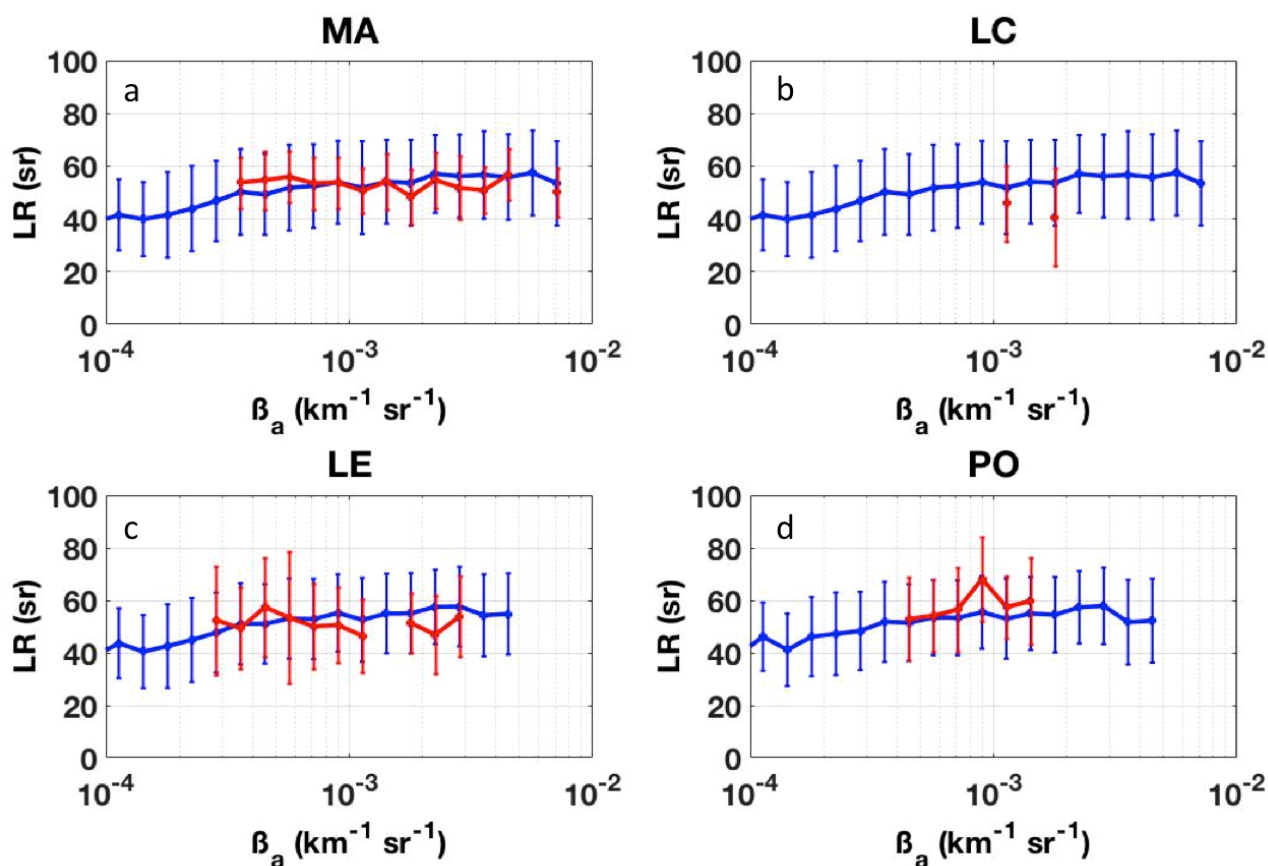
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1105 **Appendix C: Model – EARLINET comparison at 532 nm**

1106 Figure C1 depicts the result of the comparison between EARLINET stations and our developed model (red and blue  
1107 curves, respectively) in terms of ‘mean’ LR per bin of  $\beta_a$  at  $\lambda=532$  nm. Note that only  $\beta_a$  bins containing at least 1% of  
1108 the total modeled data were considered. Similarly to the results at 355 nm shown in section 4.1, a general good  
1109 agreement between the modeled and the measured LR values is found. As attested by the low value of the mean  
1110 discrepancy of Table 6, the modeled curve well fits with Madrid observations. Some major deviations are found for  
1111 Lecce, which, however, at 532 nm, has a very low number of considered points (i.e. 109).



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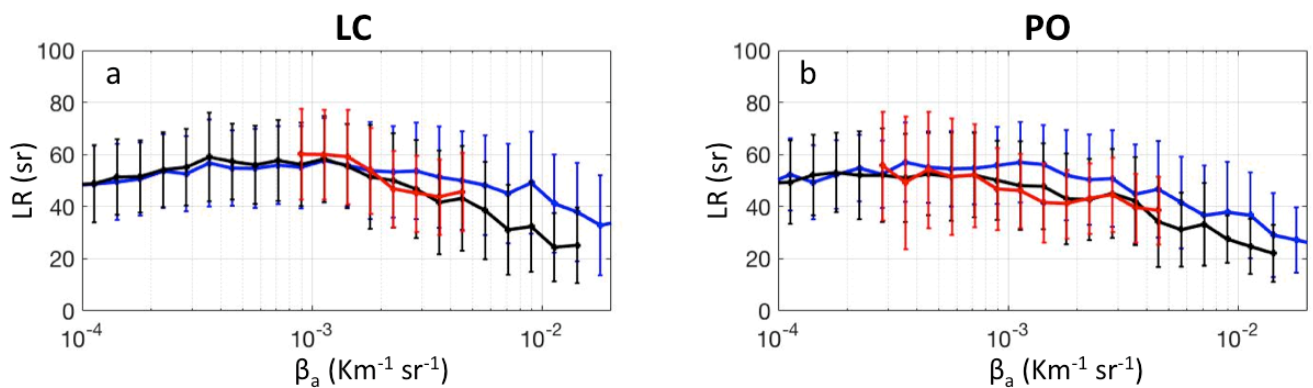
1113 **Figure C1. Model-simulated (blue) and lidar measured (red) LR vs  $\beta_a$  mean curves at 532 nm calculated per 10 equally**  
1114 **spaced bins per decade of  $\beta_a$  in a) Madrid, b) Lecce, c) Leipzig, and d) Potenza EARLINET lidar station. Vertical bars are**  
1115 **the associated standard deviations.**

1116 **Appendix D: Model sensitivity tests for optimal configurations at LC and PO sites**

1117 According to the results reported in Tab. B1, two model configurations (CM0a and CM0b) have been set up to better  
 1118 reproduce the EARLINET observations of LR vs  $\beta_a$  at LC and PO sites, respectively. The comparison between these  
 1119 two configurations, the EARLINET measurements and the CM0 set-up are illustrated in Fig. B1 (panel a and b for LC  
 1120 and PO, respectively) in terms of LR mean value curves per 10 equally spaced bins per decade of  $\beta_a$ . Blue and red  
 1121 colors have the same meaning of Fig. 5 (i.e. CM0 model and observation curves, respectively), black curves refer to the  
 1122 LR vs  $\beta_a$  estimated through the CM0a and CM0b model versions for LC and PO stations, respectively. Vertical bars are  
 1123 the associated standard deviations.

1124 The only difference between CM0a and CM0 configuration consists in the upper bound of the variability range of  $N_{tot}$   
 1125 ( $5000$  vs  $10000 \text{ cm}^{-3}$  at ground, respectively). This modification seems to fit the observed LR vs  $\beta_a$  behavior at  $355 \text{ nm}$ .  
 1126 The upper bound  $N_{tot}$  value is similar to the one (i.e.  $N_{tot}$  upper bound  $=3000 \text{ cm}^{-3}$  at ground) used in the work of  
 1127 Barnaba et al. (2007) to characterize the optical properties of the continental aerosol present over southeastern Italy.  
 1128 The computed mean model-measurement LR relative difference between CM0a configuration and LC Earlinet  
 1129 measurements is around 5%.

1130 Similarly, the CM0b configuration uses the same value for the upper bound of  $N_{tot}$  variability range and, in addition,  
 1131 higher values of the  $r_3$  variability range of ( $1.0 - 1.2 \mu\text{m}$  vs  $0.3 - 0.5 \mu\text{m}$ , respectively). As highlighted by the panel b of  
 1132 Fig. B1, this model configuration allows well reproducing the LR vs  $\beta_a$  behavior derived by EARLINET lidar Raman  
 1133 measurements at  $355 \text{ nm}$ . This result seems to indicate the presence of coarser aerosols in a clean continental  
 1134 environment. In comparison to the CM0 model, the mean model-measurement LR relative difference decreases from  
 1135 17% to 6%.

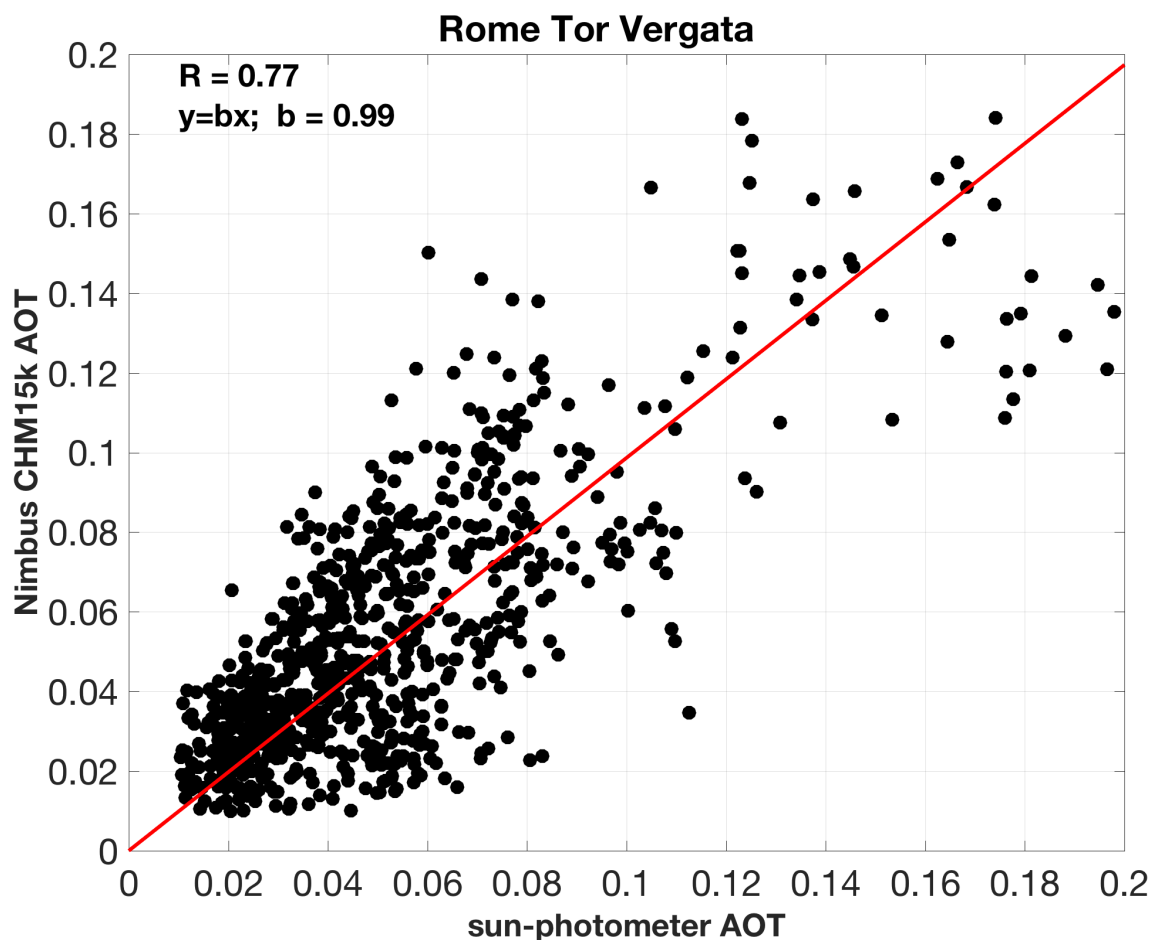


1136  
 1137 **Figure D1. Model-simulated (blue and black lines) and lidar measured (red lines) LR vs  $\beta_a$  mean curves at  $355 \text{ nm}$  calculated**  
 1138 **per 10 equally spaced bins per decade of  $\beta_a$  for the LC and PO EARLINET lidar stations (panel a and b, respectively). Blue**  
 1139 **color refers to CM0 model configuration, black color to CM0a and CM0b model configurations adapted to LC and PO sites,**  
 1140 **respectively.**

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1142 **Appendix E: ALC vs sun-photometer AOTs**

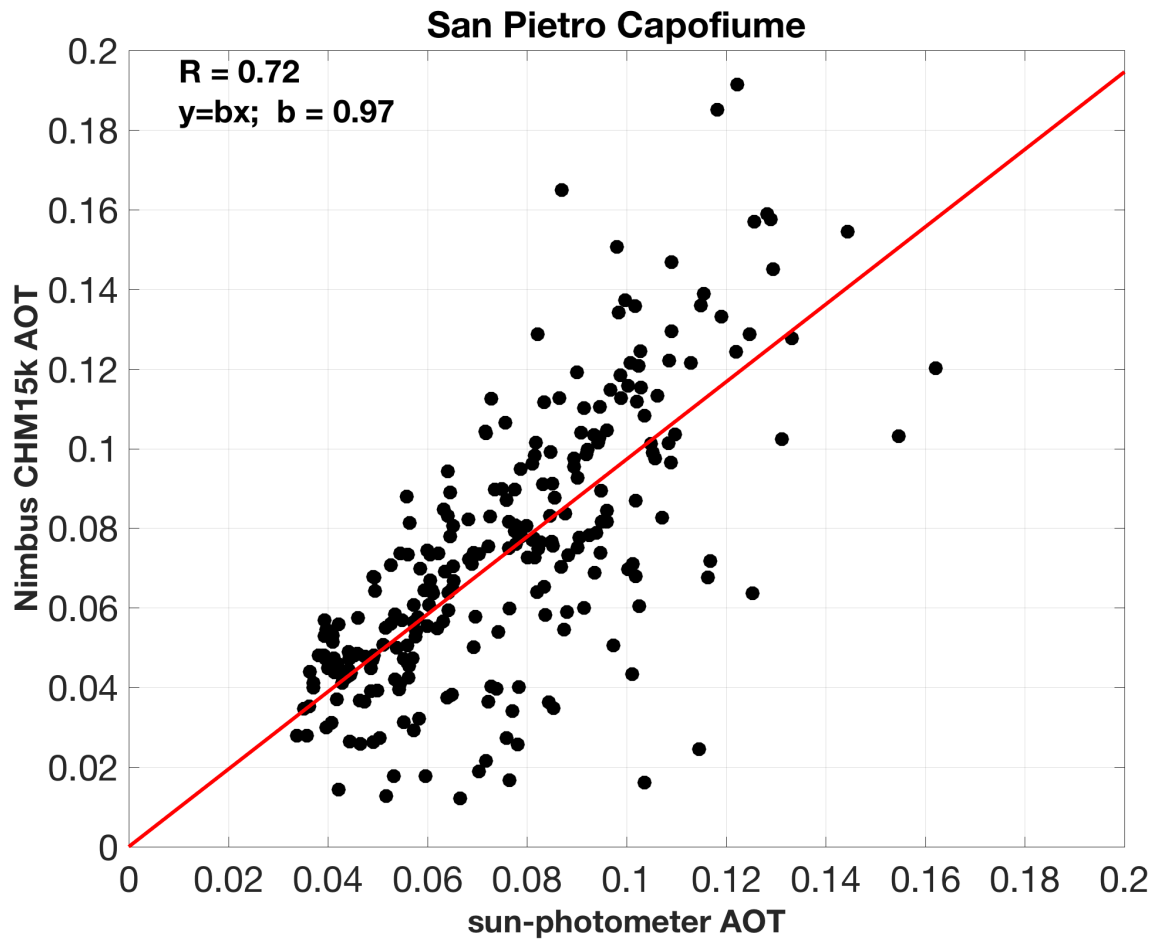
1143 To have sense of both absolute and relative errors of AOT, we reported in this section the scatter plots between the  
1144 hourly-mean coincident AOTs at 1064 nm as derived by ALC model-based approach and those measured at 1020 nm  
1145 by the sun-photometers installed at RTV, SPC and ASC, respectively (Figure E1, E2 and E3). The corresponding linear  
1146 fit  $y = bx$  (red line), where  $x = \text{sun-photometer AOT}$ ,  $y = \text{Nimbus CHM15k AOT}$  are also shown in the plots. The  
1147 values of the correlation coefficients for the three sites ( $R = 0.77$ ,  $R=0.72$  and  $R=0.73$  for RTV, SPC and ASC,  
1148 respectively) attest a relatively good agreement between the two AOT measurements.



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1150 **Figure E1. Scatter plot between the hourly-mean coincident AOTs at 1064 nm as derived by the ALC model-based approach**  
1151 **and measured at 1020 nm by the AERONET photometer at RTV. The red line represents the linear fit  $y = bx$  between the**  
1152 **two datasets, where  $x = \text{sun-photometer AOT}$ ;  $y = \text{Nimbus CHM15k AOT}$ .**

1153



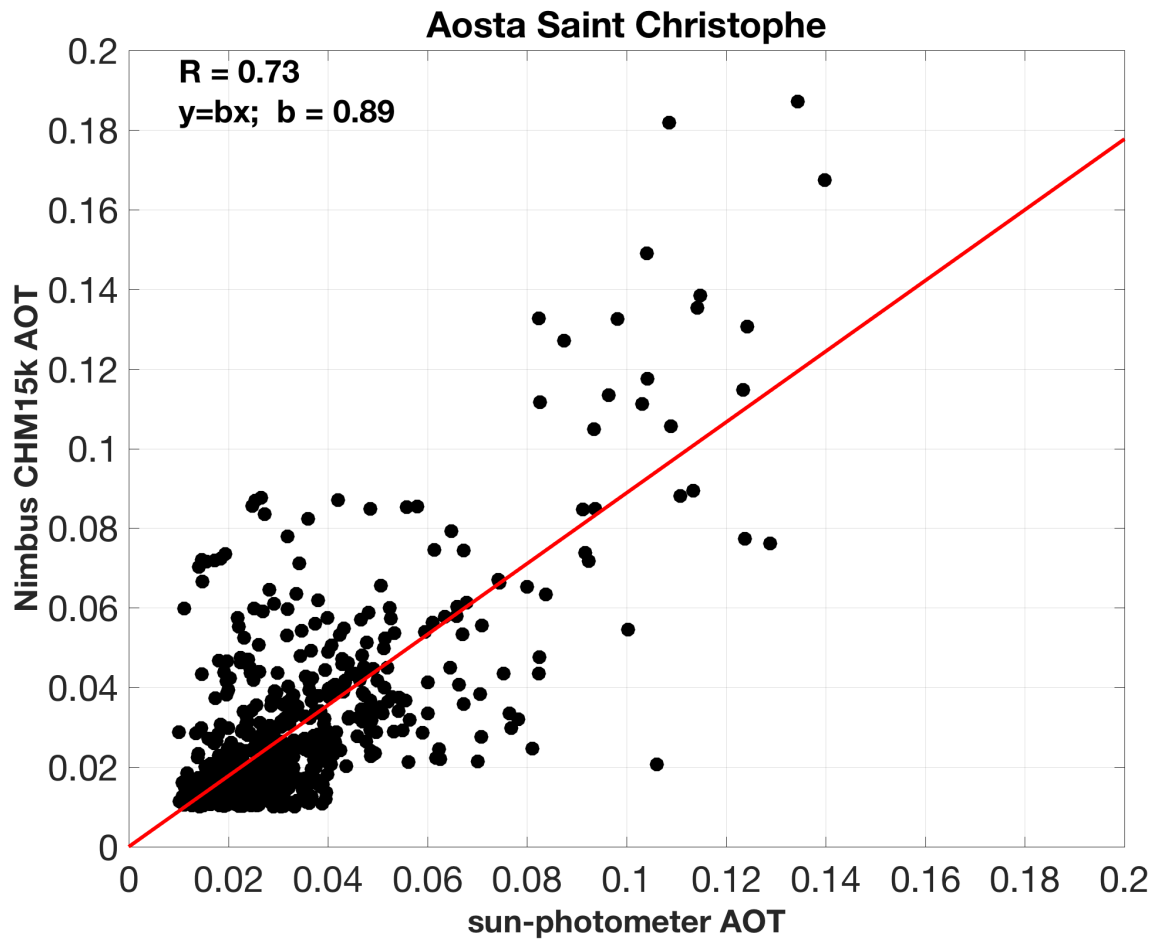
1154

1155 **Figure E2. Scatter plot between the hourly-mean coincident AOTs at 1064 nm as derived by the ALC model-based approach**  
 1156 **and measured at 1020 nm by the SKYRAD photometer at SPC. The red line represents the linear fit  $y = bx$  between the two**  
 1157 **datasets, where  $x = \text{sun-photometer AOT}$ ;  $y = \text{Nimbus CHM15k AOT}$ .**

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1162 Figure E3. Scatter plot between the hourly-mean coincident AOTs at 1064 nm as derived by the ALC model-based approach  
1163 and measured at 1020 nm by the SKYRAD photometer at ASC. The red line represents the linear fit  $y = bx$  between the two  
1164 datasets, where  $x = \text{sun-photometer AOT}$ ;  $y = \text{Nimbus CHM15k AOT}$ .

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