1 **Towards Objective Identification and Tracking of Convective Outflow Boundaries in Next-**

2 **Generation Geostationary Satellite Imagery**

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29   **Abstract**

30   Sudden wind direction and speed shifts from outflow boundaries (OFBs) associated with deep

31   convection significantly affect weather in the lower troposphere. Specific OFB impacts include

32   rapid variation in wildfire spread rate and direction, the formation of convection, aviation hazards,

33   and degradation of visibility and air quality due to mineral dust aerosol lofting. Despite their

34   recognized importance to operational weather forecasters, OFB characterization (location, timing,

35   intensity, etc.) in numerical models remains challenging. Thus, there remains a need for objective

36   OFB identification algorithms to assist decision support services. With two operational next-

37   generation geostationary satellites now providing coverage over North America, high-temporal

38   and spatial resolution satellite imagery provides a unique resource for OFB identification. A

39   system is conceptualized here designed around the new capabilities to objectively derive dense

40   mesoscale motion flow fields in the Geostationary Operational Environmental Satellite (GOES)-

41   16 imagery via optical flow. OFBs are identified here by isolating linear features in satellite

42   imagery, and back-tracking them using optical flow to determine if they originated from a deep

43   convection source. This "objective OFB identification" is tested with a case study of an OFB

44   triggered dust storm over southern Arizona. Results highlight the importance of motion

45   discontinuity preservation, revealing that standard optical flow algorithms used with previous

46   studies underestimate wind speeds when background pixels are included in the computation with

47   cloud targets. The primary source of false alarms is incorrect identification of line-like features in

48   the initial satellite imagery. Future improvements to this process are described to ultimately

49   provide a fully automated OFB identification algorithm.

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1. Introduction

Downburst outflows from associated deep convection (Byers and Braham Jr., 1949; Mitchell and Hovermale, 1977) play a significant, dynamic role in modulation of the lower troposphere. Their direct impacts to society are readily apparent—capsizing boats on lakes and rivers with winds that seem to "*come out of nowhere*" (e.g. The Branson, MO duck boat accident; Associated Press 2018), causing shifts in wildfire motion and fire intensity that put firefighters in harm's way (e.g. the Waldo Canyon and Yarnell Hill Fires; Hardy and Comfort, 2015; Johnson et al., 2014), and threatening aviation safety at regional airports with sudden shifts from head to tail-winds and turbulent wakes (Klingle et al., 1987; Uyeda and Zrnić, 1986). In the desert southwest, convective outflows can loft immense amounts of dust, significantly reducing surface visibility and air quality for those within the impacted area (e.g. Idso et al. 1972; Raman et al. 2014). These outflows are commonly associated with rapid temperature, pressure, and moisture changes at the surface (Mahoney III, 1988). Furthermore, the collision of outflows from adjacent storms can serve as the focal point of incipient convection or the intensification of nascent storms (Rotunno et al., 1988; Mueller et al., 2003).

Despite the understood importance of deep convection and convectively driven outflows, high resolution models struggle to characterize and identify them (e.g. Yin et al. 2005). At present, outflow boundaries (OFBs) are instead most effectively monitored in real-time at operational centers around the world with surface, radar, and satellite data. Satellites often offer the only form of observation in remote locations. The most common method for detecting outflows via satellite data involves the identification of clouds formed by strong convergence at the OFB leading edge. When the lower troposphere is dry, OFBs may be demarcated by an airborne "dust front", after passing over certain surfaces prone to deflation by frictional winds (Miller et al., 2008). The task

73   of identifying OFBs can prove quite challenging and would benefit greatly from an objective

74   means of feature identification and tracking for better decision support services.

75   The Advanced Baseline Imager (ABI), an imaging radiometer carried on board the

76   Geostationary Operational Environmental Satellite (GOES)-R era systems, offers a leap forward

77   in capabilities for the real-time monitoring and characterization of OFBs. Its markedly improved

78   spatial (0.5 vs. 1.0 km visible, 2 km vs. 4 km infrared), spectral (16 vs. 5 spectral bands), and

79   temporal (5 min vs. 30 min continental U.S., and 15 min vs. 3 hr full disk) resolution provides new

80   opportunities for passive sampling of the atmosphere over the previous generation (Schmit et al.,

81   2016). The vast improvement of temporal resolution alone (which includes mesoscale sectors that

82   refresh as high as 30 s) allows for dramatically improved tracking of convection (Cintineo et al.,

83   2014; Mecikalski et al., 2016; Sieglaff et al., 2013), fires and pyroconvection (Peterson et al., 2015,

84   2017, 2018), ice flows, and synoptic scale patterns (Line et al., 2016). This higher temporal

85   resolution makes identification of features like OFBs easier as well because of greater frame-to-

86   frame consistency.

87   The goal of this work is to use ABI information towards objective identification of OFBs. One

88   of the notable challenges in satellite identification of OFBs over radar or models is the lack of

89   auxiliary information. When working with a radar or a numerical model framework, for example,

90   additional information is available on the flow, temperature, and pressure tendency of the

91   boundary. Without that information, however, forecasters must rely on their knowledge of gust

92   front dynamics to identify OFBs in satellite imagery. Here, we introduce the concept of objectively

93   derived motion using GOES-16 ABI imagery for feature identification via an advanced optical

94   flow method, customized to the problem at hand. A case study of a convectively triggered OFB

95   and accompanying haboob dust front is presented in 5-min GOES-16 contiguous United States

96   (CONUS) sector information, as a way of evaluating and illustrating the potential of the

97   framework.

98   This paper is outlined as follows. The background for objective motion extraction and OFB

99   identification is presented in Section 2. The optical flow methods developed for this purpose are

100   discussed in Section 3. Section 4 presents the case study test of the current algorithm, and Section

101   5 concludes the paper with a discussion on plans for future work in objective feature identification

102   from next-generation geostationary imagers of similar fidelity to the GOES-R ABI, which are

103   presently coming online around the globe.

104

105   **2.   Background**

106   *2.1 Previous Work in OFB Detection*

107   Objective identification of OFBs in meteorological data has been a topic of scientific inquiry

108   for more than 30 years. Uyeda and Zrnić (1986) and Hermes et al. (1993) use detections of wind

109   shifts in terminal Doppler radar velocity measurements to isolate regions of strong radial shear

110   associated with OFBs. Smalley et al. (2007) include the "fine line" reflectivity structure of

111   biological- and precipitation-sized particles to identify OFBs via image template matching.

112   Chipilski et al. (2018) considered the OFB objective identification in numerical models using

113   similar image processing techniques, but with additional dynamical constraints on vertical velocity

114   magnitudes and mean-sea level pressure tendency. Objective OFB identification has not been

115   demonstrated to date with the new ABI observations of the GOES-R satellite series. Identification

116   via satellite imagery would be valuable for local deep convection nowcasting algorithms which

117   use boundary presence as a predictor field (Mueller et al., 2003; Roberts et al., 2012), and for

118   operational centers around the world which may not have access to ground-based Doppler radar

119   data.

120   Traditionally, forecasters have identified OFBs in satellite imagery by visually identifying

121   quasi-linear low-level cloud features and back-tracking them to an associated deep convection

122   source. Previous objective motion derivation algorithms are not designed to yield the dense wind

123   fields necessary for identifying and tracking features such as OFBs (Bedka et al., 2009; Velden et

124   al., 2005). In fact, the original image window-matching atmospheric motion vector (AMV)

125   algorithms produce winds only over targets deemed acceptable for tracking by pre-processing

126   checks on the number of cloud layers in a scene, brightness gradient strength, and patch coherency.

127   The targets are further filtered with post-processing checks on acceleration and curvature through

128   three-frame motion and deviation from numerical model flow (Bresky et al., 2012; Nieman et al.,

129   1997; Velden et al., 1997). These practices were followed for a very practical reason—AMV

130   algorithms were tailored for model data assimilation. In the formation of the model analysis,

131   observational data must be heavily quality-controlled with outliers removed, to minimize data

132   rejection. Here, information such as OFBs would be rejected due to the detailed space/time

133   structure of actual convection which is typically poorly represented by the numerical model.

134   Deriving two-dimensional flow information at every point in the imagery would require either

135   modification of previous AMV schemes or post-processing of the AMV data via objective analysis

136   (e.g. Apke et al. 2018). The latter typically will not capture motion field discontinuities, resulting

137   in incorrect flows near feature edges (Apke et al., 2016). To capture such discontinuities in a dense

138   flow algorithm, new computer vision techniques, such as the gradient-based methods of optical

139   flow, must be adopted.

140   *2.2 Optical Flow Techniques*

141   Optical flow gradient-based techniques derive motion within fixed windows, thus eliminating

142   the reliance on models for defining a search region. A core assumption of many optical flow

143   techniques is brightness constancy (Horn and Schunck, 1981). Considering two image frames,

144   brightness constancy states that the image intensity I at some point $\boldsymbol{x} = [x, y]^T$ is equal to the

145   image intensity in the subsequent frame at a new point, $\boldsymbol{x} + \boldsymbol{U}$, where $\boldsymbol{U} = [u, v]^T$ represents the

146   flow components of the image over the time interval (Δt) between the two images:

$$I(\boldsymbol{x}, t) = I(\boldsymbol{x} + \boldsymbol{U}, t + \Delta t) \tag{1}$$

147   Eq. (1) can be linearized to solve for the individual flow components, u and v:

$$\nabla I \cdot \boldsymbol{U} + I_t = 0 \tag{2}$$

148   Where $\nabla I = [I_x, \ I_y]$ represents the intensity gradients in the x and y direction, and I_t represents

149   the temporal gradient of intensity. For one image pixel, Eq. (2) contains two unknowns with a

150   simple translation model for $\boldsymbol{U}$; therefore, it cannot be solved point-wise. One well-known

151   approach to solving this so-called "aperture problem" is the Lucas-Kanade, hereafter LK method,

152   which considers a measurement neighborhood of the intensity space and time gradients (e.g.,

153   Baker and Matthews, 2004; Bresky and Daniels, 2006). Use of neighborhoods, or image windows,

154   to derive optical flow are called *local* approaches. Another seminal approach was introduced by

155   Horn and Schunck (1981; HS Method) which solves the aperture problem by adding an additional

156   smoothness constraint to the brightness constancy assumption, and minimizing an energy

157   magnitude between two images:

$$E(u, v) = \iint_{\Omega} (\nabla I \cdot \boldsymbol{U} + I_t)^2 + \alpha(|\nabla_2 u|^2 + |\nabla_2 v|^2) \ d\mathbf{x} \tag{3}$$

158   Where $E(u, v)$ represents an energy functional to be minimized over all image pixels Ω, α is a

159   constant weight used to control the smoothness of the flow components $u(\mathbf{x})$ and $v(\mathbf{x})$, and $\nabla_2 =$

160    $[\partial/\partial x, \partial/\partial y]^T$. This optical flow derivation is called a *global* approach. Eq. (3) is minimized in

161   Horn and Schunck (1981) by deriving the Euler-Lagrange equations, and numerically solving with

162   Gauss-Seidel iterations.

163   Linearizing the brightness constancy equation into Eq. (2) means that large and non-linear

164   displacements (typically > 1 pixel between images) will not be captured (Brox et al., 2004). Thus,

165   most optical flow computations initially spatially subsample images to where all displacements

166   are initially less than 1-pixel (Anandan, 1989; discussed more in Section 3.1), which can cause

167   fast moving small features to be lost. Note that reducing the temporal resolution of GOES imagery

168   (e.g. 10-min vs. 5-min scans) increases the displacement of typical meteorological features

169   between frames. Furthermore, constancy assumptions are more likely violated with reduced

170   temporal resolution since image intensity changes more through evaporation and condensation of

171   cloud matter over time. Thus, for the spatial resolution of ABI, it is impractical to consider optical

172   flow gradient-based methods at temporal resolutions coarser than 5-min for several mesoscale

173   meteorological phenomena, including OFBs. Very spatially coarse images do not need to be

174   initially used with faster scanning rates, such as super rapid scan 1-min information (Schmit et al.,

175   2013), or the 30-s temporal resolution mesoscale mode of ABI (Schmit et al., 2016).

176   While both the LK and HS methods are designed for deriving dense flow in satellite imagery,

177   neither account for motion discontinuities in fields. Hence, both suffer from incorrect flow

178   derivations near cloud edges, and would perform poorly for OFB detection and tracking. Black

179   and Anandan (1996) offer an intuitive solution to this problem, whereby the energy functional is

180   designed to minimize robust functions that are not sensitive to outliers:

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$$E(u,v) = \iint_\Omega \rho_d(\nabla I \cdot \boldsymbol{U} + I_t) + \rho_s(|\nabla_2 u|^2 + |\nabla_2 v|^2)d\mathbf{x} \tag{4}$$

181 The robust function data term for the standard HS approach is simply $\rho_d(r) = r^2$, and smoothness

182 $\rho_s(r) = r$ which implies that energy functionals increase quadratically for r outliers. Other robust

183 functions can also be minimized with similar gradient descent algorithms to Gauss-Seidel

184 iterations, while being less sensitive to outliers (Press et al., 1992; Black and Anandan, 1996).

185 Robust functions are popular in recent optical flow work (Brox et al., 2004; Sun et al., 2010), and

186 a similar approach was adopted here and is discussed further in the methodology section. The

187 reader is referred to works by Barron et al. (1994), Fleet and Weiss (2005), and Sun et al. (2010)

188 for a more comprehensive review on optical flow techniques.

189 The relevance of optical flow in satellite meteorological research continues to increase now

190 that scanning rates of sensors such as the ABI are routinely at sub 5-min time scales, making

191 motion easier to derive objectively (Bresky and Daniels, 2006; Héas et al., 2007; Wu et al., 2016).

192 The dense motion estimation within fine-temporal resolution data has yet to be used for feature

193 identification. Optimizing optical flow for this purpose, and its specific application to OFBs, is the

194 aim of this study. The next section outlines our approach to this end.

195

3. Methodology

197 *3.1 Optical Flow Approach*

198 A recent approach to handle piecewise and non-linear image changes in flow is introduced by

199 Brox et al. (2004) (Hereafter B04), where the brightness constancy assumption is no longer

200 linearized, i.e.

$$E(u, v) = \iint_\Omega \rho_d(|I(\boldsymbol{x} + \boldsymbol{U}, t + \Delta t) - I(\boldsymbol{x}, t)|^2$$

$$+ \gamma |\nabla_2 I(\boldsymbol{x} + \boldsymbol{U}, t + \Delta t) - \nabla_2 I(\boldsymbol{x}, t)|^2) \qquad (5)$$

$$+ \alpha \, \rho_s(|\nabla_2 u|^2 + |\nabla_2 v|^2) d\boldsymbol{x}$$

201 Following B04, within the data robust function, we now have also included a gradient constancy

202 assumption, which is weighted by a constant γ to make the derived flow more resilient to changes

203 in illumination. Avoiding linearization of constancy assumptions improves the identification of

204 large displacements between images. The Charbonnier penalty is used for the data and smoothness

205 robust functions following Sun et al. (2010),

$$\rho_d(r^2) = \rho_s(r^2) = \sqrt{r^2 + \epsilon^2} \qquad (6)$$

206 with ϵ representing a small constant present to prevent division by zero in minimization, set to

207 0.001. The values for $\boldsymbol{U}$ are found by solving the Euler-Lagrange equations of Eq. (5) with

208 numerical methods

$$E_u - \frac{dE_{u_x}}{dx} - \frac{dE_{u_y}}{dy} = 0 \qquad (7)$$

$$E_v - \frac{dE_{v_x}}{dx} - \frac{dE_{v_y}}{dy} = 0 \qquad (8)$$

209 with reflecting boundary conditions and subscripts that imply the derivatives. Eqs. (7) and (8) are

210 solved with a nested-fixed point successive over-relaxation iteration scheme described in B04 and

211 summarized in Fig. 1. The reader is referred to Chapter 4 of Brox (2005) for details on the full

212 discretization of the derivatives in the successive over-relaxation scheme. Here, only the spatial

213 dimensions are used for the smoothing term, though it is possible to include the time dimension

214 with this system as well.

215   A difficulty in solving Eqs. (7) and (8) is that the successive over-relaxation scheme may

216   converge on a local minimum of $E(u, v)$, rather than finding the global minimum. The typical

217   approach to find the global minimum is to compute optical flow with coarse- to fine-scale warping

218   iterations (e.g. Anandan, 1989). Coarse-to fine-scale warping iterations work by subsampling the

219   initial image at native resolution to a coarser spatial resolution and computing the flow initially at

220   the coarsest resolution in the image pyramid. The U results from the coarse image flow are then

221   used as the first guess field to for the next finest scale on the image pyramid (Fig. 2), and the

222   second image is warped accordingly. The warping step ensures that estimated displacements at

223   every step in the image pyramid remain small.

224   The B04 scheme includes coarse-to fine-scale warping iterations at every outer iteration k.

225   This means that the first iteration is run on a subsampled image, and the subsampling is reduced

226   by a scale factor at every k until the image reaches native resolution at the final $k = nK$. Images

227   at every k in this subsampling are found using a gaussian image pyramid technique with bicubic

228   interpolation. The flow values of the image at $k - 1$ are upscaled accordingly at k also with

229   bicubic interpolation (the initial flow guess is $u = v = 0$ at $k = 0$). For improved computation of

230   spatial derivatives, the initial image is also smoothed with a 9x9 pixel kernel gaussian filter with a

231   standard deviation set to 1.5 pixels. The specific settings used for the coarse- to fine- warped flow

232   scheme here are shown in Table 1.

233   *3.2 Objective OFB identification*

234   There are two steps to the objective OFB identification process. First, a linear feature or sharp

235   boundary is identified in visible or infrared imagery. In some cases, the first step alone is enough

236   to identify OFBs subjectively. The second step is tracking that feature back in time to see where

237   it originated from (typically, near an area with deep convection). In the case of near stationary

238   convection and low-level flow, a forecaster might also use radial like propagation in this decision-

239   making process, however, since convection geometry and low-level flow varies from storm to

240   storm, only the first two steps are considered here. This approach aims to mirror the subjective

241   process, leveraging the information content of optical flow to do so.

242   To handle the first step of line feature identification, a simple image line detection scheme was

243   performed using the sum of a set of two-dimensional convolution kernels:

244   $$a_1 = \begin{bmatrix} -1 & -1 & -1 \\ 2 & 2 & 2 \\ -1 & -1 & -1 \end{bmatrix} a_2 = \begin{bmatrix} -1 & 2 & -1 \\ -1 & 2 & -1 \\ -1 & 2 & -1 \end{bmatrix} a_3 = \begin{bmatrix} 2 & -1 & -1 \\ -1 & 2 & -1 \\ -1 & -1 & 2 \end{bmatrix} a_4 = \begin{bmatrix} -1 & -1 & 2 \\ -1 & 2 & -1 \\ 2 & -1 & -1 \end{bmatrix}$$

245

246   Applying these kernels to the gaussian smoothed ABI visible imagery (using a 21x21 kernel and

247   standard deviation of 5 pixels) results in high intensity values where line structures exist. A two-

248   dimensional convolution threshold of 0.02 was used with the visible imagery calibrated to

249   reflectance factor to isolate line features. This method was compared to a subjective interpretation

250   of boundary location for validation.

251   To address the second step of the process, the constrained optical flow approach described in

252   Section 3.1 was used to track the boundary pixels (both objectively and subjectively identified)

253   back in time for three hours. The values of motion at each step in the backwards trajectory were

254   determined with bilinear interpolation of the optical flow derived dense vector grid. If a back-

255   traced pixel of the linear feature arrived within 50 km of a convective area with a 10.35 μm

256   brightness temperature ($BT_{10.35}$) lower than 223 K (-50 °C; using previous satellite imagery

257   matched to the back-trajectory time), the original point was considered an OFB. While this

258   brightness temperature threshold is subjective and can vary from case to case, it was found to

259   produce a reasonable approximation of deep convection areas when compared to ground-based

260   radar information for the case study described in the subsequent sections.

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3.3 Data

261

262   The objective OFB identification methodology is tested using a case study from 5 July 2018

263   over the southwestern United States. This event featured a distinct OFB and associated dust storm

264   that was well-sampled by various ground- and space- based sensors. GOES-16 was in Mode-3,

265   generating one image over the study area every 5-min (continental U.S., or CONUS, ABI scan

266   domain). Optical flow computations employ the GOES-16 (GOES-East) ABI red band (0.64 μm;

267   ABI channel 2), provided at a nominal sub-satellite spatial resolution of 500 m, but closer to 1 km

268   at the case study location. This channel is used at native resolution, though it can be subsampled

269   with a low-pass filter such that future versions can implement color information from the blue and

270   near-infrared bands (e.g. Miller et al. 2012). This means that the optical flow approach here is

271   daytime only. A similar B04 approach can be used on infrared data as well for day/night

272   independent information, though for detecting OFBs in the low levels, proxy visible products

273   would perform best. As described above, the clean longwave infrared band (10.35 μm; ABI

274   channel 13) is used as first-order information on optically thick cloud-top heights and to assess the

275   convective nature of the observed scene ($BT_{10.35} < 223$ K).

276   High frequency Automated Surface Observing Stations (ASOS; NOAA 1998), recording

277   temperature, pressure, wind speed and direction once every minute, complement the satellite

278   imagery. The Weather Surveillance Radar-1988 Doppler (Crum and Alberty, 1993) dual-

279   polarimetric data also sampled the OFB event from the KIWA radar near Phoenix, AZ. To

280   highlight the OFBs and the presence of dust, horizontal reflectivity and correlation coefficient are

281   used (Van Den Broeke and Alsarraf, 2016). Finally, for information on the full 3D dynamics of

282   the case study, a numerical model representation of the environment was collected from the High

283   Resolution Rapid Refresh system (HRRR, Benjamin et al. 2016). The combination of these model

284   and observation datasets is employed to confirm the presence of a distinct convective OFB, rather

285   than some other quasi-linear feature, such as a bore or elevated cloud layer, etc.

286

287   **4.   Case Study Description**

288   Convection was observed in south central Arizona on 5 July 2018 after 1800 UTC. A large

289   and well-defined linear structure emerged from below the convective cloud cover at 2200 UTC to

290   6 July 2018 0100 UTC propagating westward in GOES-16 imagery (Fig. 3). This linear structure,

291   demarcated by roll (arcus) clouds on the northern side and lofted dust on the southern side, was

292   apparent with strong visible reflectance contrast against the relatively dark surface and $BT_{10.35}$ ~

293   10 K cooler than the underlying surface. The dust lofted by this outflow produced low visibility

294   and hazardous driving conditions near Phoenix, AZ. Dust storm warnings were issued by the local

295   National Weather Service (NWS) forecast office by 2300 UTC. The structure's observed radial

296   propagation away from nearby deep convection and associated cloud and dust features lends to its

297   interpretation as a convective OFB.

298   The OFB was also captured in radar scans from KIWA at 2200 UTC (Fig. 4). The coincidence

299   of low correlation coefficient (< ~0.5) and moderate to high reflectivity (near 20 dBZ) imply that

300   the OFB was associated with airborne coarse dust particles. Surface observations taken at the

301   ASOS station reveal temperatures exceeding 317 K (44 °C) ahead of the OFB, with calm winds

302   (Fig. 5). Temperatures dropped by 4 K, wind speeds changed direction and increased sharply, and

303   dew points increased rapidly as the OFB crossed the station at ~2316 UTC. The rapid change in

304   low-level meteorology is consistent with convective OFBs sampled in previous studies (e.g.

305   Mahoney III, 1988; Miller et al., 2008).

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306   The HRRR model captured the broad characteristics of this event (Fig. 6), showing moderate

307   low-level winds in excess of 10 m s^{-1} (Fig. 6a), cooler temperatures (Fig. 6b), and simulated

308   cumulus clouds from forced ascent (Fig. 6c). Model cross sections (Fig. 6d) indicated a moderate

309   increase in vertical motion ahead of the numerically derived boundary, and a sharp decrease in

310   virtual potential temperature behind the boundary. The shape of the virtual potential temperature

311   profile is consistent with other model observations of OFBs (e.g. Chipilski et al., 2018). The

312   observation and model data all show that the linear structure observed in Fig. 3 was modifying the

313   dynamics of the surface in a manner consistent with OFBs, and not some other linear cloud feature

314   type that is decoupled from the surface and may be misidentified by the satellite. Since such low-

315   level linear features are often obscured by cloud layers at higher altitudes, this case study in some

316   respects represents a best-case-scenario for evaluating optical flow capabilities towards identifying

317   OFBs.

318

319   **5.   Results**

320   The first step in OFB identification requires identification of a feature that appears linear in

321   the imagery. Compared to the subjective boundary identification (considered as truth here; Fig 7a,

322   blue dots), the convolution method gives a reasonable approximation to where the OFB is located

323   within the higher intensity points in the convolution (Fig. 7b). Unfortunately, the simply-applied

324   convolution is also sensitive to linear features associated with the deep convection itself (the blue

325   shading in Fig. 7b). Hence, false alarms appear east of the boundary. These issues can be filtered

326   out using either cloud-top height or brightness temperature thresholding from separate infrared

327   channels. Alternatively, the storm-relative motion (here > 15 m s^{-1}) from optical flow was used

328   here to filter the false alarms (the red shading in Fig. 7b).

329   The second step requires these linear fast-moving features to be traced backward to a deep

330   convection source using the optical flow computation (Fig. 8). To the west of the boundary, near

331   stationary optical flow vectors highlight the background (or ground) pixels. The boundary itself

332   exhibits a westward movement near 15-20 m s^{-1} (~30-40 kts). The feature also appears to bow

333   outwards after faster motions are observed, near 33° N, -112° E during 2338-2358 UTC (Figs. 8b,

334   c). Similar westward motion is derived in the wake of the OFB, within the convective cold pool.

335   This results from the presence of airborne dust particles, which facilitate the computation of optical

336   flow vectors in this region.

337   The backwards trajectories of the subjectively and objectively identified OFB pixels in Figs.

338   7c and d (B04 method) show that many of the linear cloud features, particularly those associated

339   with the central arcus cloud, indeed originated near deep convection. However, when the

340   backwards trajectories of the B04 method were compared to other optical flow methods, such as

341   the approach by Wu et al. (2016), most were unsuccessful at obtaining coincidence between linear

342   cloud features along the OFB and a deep convection source. Wu et al. (2016) used an approach

343   introduced to the community by Farnebäck (2001), which is a *local* window method for optical

344   flow.

345   Example points 1–7 examined within the subjectively identified OFB backward trajectories

346   highlight an issue with *local* window approaches for this application (Fig. 9). The B04 approach

347   (Fig. 9, blue/yellow) produced motions that were relatively consistent with the true boundary

348   motion. Thus, many points that are lost in the *local* approaches are successfully backtracked to

349   the initial deep convection (e.g. points 3–5). With the Wu et al. approach (Fig. 9, orange/red),

350   OFB targets move slower than the actual boundary, and, over a three-hour tracking period,

351   eventually become stuck within the stationary background pixels. This tracking issue stems from

352 an assumption made in many *local* approaches that pixels within an image window all move in the

353 same direction with the same speed. When background pixels are included within an image

354 window containing clouds or dust, the resulting optical flow speed would then be underestimated.

355 The slow bias is observed in plots of optical flow speeds along the OFB (Fig. 10), where the Wu

356 et al. approach was ~5-10 m s^{-1} slower than the B04 approach. While not shown, we found similar

357 backward trajectory issues using the LK approach. Full loops of the optical flow in Fig. 8 and

358 trajectories in Fig. 9 are included as supplementary material to this manuscript.

359 For all approaches tested, however, methods struggled to backtrack the newly formed cumulus

360 to the north and the dust front to the south. With the cumulus to the north, the issues with each

361 algorithm appear to result from rapid cumulus development between frames (e.g. points 1 and 2 in

362 Figs. 9a, b). Condensation like what is observed here is unfortunately not considered in the

363 brightness constancy assumption. Thus, condensing cloud features would only be tracked back to

364 when they initially form (after Fig. 9b) without additional dynamic constraints to Eq. (5). An

365 example can be seen when points 1 and 2 become stuck in Fig. 9c. This has important implications

366 on limitations of backtracking OFB features to deep convection with optical flow from imagery.

367 If no cloud or dust feature exists to visualize an OFB in satellite imagery, some of the feature

368 propagation may be lost.

369 The dust to the south appears in the satellite imagery as early as 2200 UTC, though it was quite

370 transparent relative to the ground. It is therefore possible the stationary background pixels may be

371 dominant in the optical flow computation at points 6 and 7, resulting in slower wind speeds than

372 the true OFB propagation. Points 6 and 7 are also located near cumulus moving across the OFB

373 motion to the south. This dust front tracking could be improved using multispectral techniques

374 designed to highlight dust features over ground pixels, or by using additional color spectrum

375   information to discourage flow smoothness in Eq. (5) across the dust front from the cumulus to

376   the south (e.g. Sun et al., 2010).

377   Many line-like targets east of the OFB in Fig. 7d also originated from the deep convection,

378   which constitute false alarms. These false alarms can be reduced by further improving the OFB

379   targeting step in the objective process in future studies. For this case study, it may have been

380   possible to use convergence methods, analogous to radar-based objective OFB identification, to

381   isolate the boundary. However, convergence as derived from the optical flow information here

382   would only work because of local, stationary background pixels ahead of the OFB. This means

383   that convergence would almost always be inferred from OFB motion when the background is

384   involved. This approach would also be sensitive to nearby cloud structures ahead of the OFB

385   which would exhibit different (non-stationary) motion from the background. It is for this reason

386   that a backwards trajectory approach was elected instead of basing the detection on local horizontal

387   convergence. The optical flow approach used here does help highlight the OFB when storm

388   motion alone was considered in addition to convolution, showing how additional tools can be used

389   in synergy to arrive at a more comprehensive objective feature identification approach in future

390   studies.

391

6.   Conclusions and Future Outlook

393   A new method for the objective identification of outflow boundaries (OFBs) in GOES-16

394   Advanced Baseline Imager (ABI) data was developed using optical flow motion derivation

395   algorithms and demonstrated with provisional success on a dust storm case study. An optical flow

396   system constructed for this purpose shows promise in identifying and backtracking object events

397   to their source over traditional flow derivation methods, which can potentially be used to isolate

398 convective OFB features. To the best of the authors' knowledge, this study represents a first

399 attempt to objectively identify OFBs in geostationary satellite imagery.

400   The primary conclusions of this study are that optical flow approaches are now a viable option

401 to acquire meso-scale flows relevant to OFB tracking and detection in 5-min geostationary satellite

402 imagery, though the successful backtracking of OFB features requires use of flow algorithms that

403 can handle the presence of motion discontinuities and stationary background flow. The optical

404 flow algorithm tested in this study produced a dense motion field that was closer than other

405 methods to the true OFB motion and provided valuable information towards full objective OFB

406 identification in new products.

407   While several OFB related image pixels were successfully identified, the algorithm here is

408 relatively immature and remains fraught with false alarms where linear features are incorrectly

409 identified, and where correct features were not successfully backtracked to deep convection. The

410 algorithm is still limited by the assumptions made within optical flow, which only account for

411 changes in image brightness intensity resulting from pure feature advection. Therefore, if no

412 features (e.g. clouds) exist to highlight an OFB boundary within the imagery, the method proposed

413 here would not function properly. The method also struggles to resolve true OFB motions with

414 transparent dust movement, where a textured background beneath the dust may dominate the

415 motion estimate within a scene. Also, while infrared brightness temperature was enough to

416 identify deep convection in this case study, convection may be missed by brightness temperature

417 imagery if it is obscured by a higher cloud layer, or if the minimum cloud-top brightness

418 temperature exceeds an arbitrarily set threshold.

419   Given these limitations, future studies will explore more advanced systems for linear structure

420 identification to identify candidate features for tracking towards full objective OFB identification.

421   A machine learning system will be used to determine which linear characteristics of the image

422   should be backtracked instead of using two-dimensional convolution. Optical flow can be used to

423   precondition training information for a machine learning approach, if motion or semi-Lagrangian

424   fields are needed. Furthermore, it will be prudent to use deep convection correspondence through

425   optical flow backtracking as one of many fields in future products, such as radial propagation away

426   from storms and near surface meteorological properties, to probabilistically decide if an image

427   pixel is associated with an OFB. To better identify deep convection areas, the GOES Lightning

428   Mapper (GLM) can be used, which provides information on lightning location and energy at 8 km

429   resolution with a 2 ms frame rate.

430   Feature identification with optical flow is not restricted to OFBs alone. For example, the

431   above-anvil cirrus plume (Bedka et al., 2018) over deep convection has been identified as an

432   important indicator of severe weather at the ground, yet no objective means of identification exists

433   today. The properties from optical flow could be used as an additional source of information in

434   such algorithm designs, allowing researchers to backtrack features to their apparent source (the

435   overshooting top in the case of the above-anvil cirrus plume) and monitor cloud temperature and

436   visible texture trends, or to simply use the dense motion itself to achieve better results. This

437   method will also be applicable to other cold pool outflow phenomena, such as bores, for which

438   new algorithms could utilize numerical model or surface observations for further clarification of

439   linear feature type.

440   Motion discontinuity preserving optical flow will also benefit several current algorithms for

441   monitoring deep convection in satellite imagery. Objective deep convection cloud-top flow field

442   algorithms (Apke et al., 2016, 2018) will benefit particularly when sharp cloud-edges and ground

443   pixels are present in an image scene. Systems that use infrared cloud-top cooling or emissivity

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444   differences for deep convection nowcasting will also improve with better estimates of pre-

445   convective cumulus motion (Cintineo et al., 2014; Mecikalski and Bedka, 2006).

446   While the utility of a backwards trajectory approach was considered here, many other possible

447   methods exist for exploiting the semi-Lagrangian properties of time-resolved observations in

448   satellite imagery (e.g. Nisi et al., 2014). Use of fine-temporal resolution information will improve

449   optical flow estimates, and in turn the estimates of brightness temperature, reflectance, or cloud-

450   property changes in a moving frame of reference. We will explore these and other refinements in

451   ongoing and future work on this exciting frontier of next-generation ABI-enabled science.

452

453   **7. Data Availability**

454   Data used in this study are available upon request by contacting the lead author.

455

456   **8. Author Contributions**

457   JMA developed the primary code for the optical flow approach used here. He also co-

458   developed the objective outflow boundary identification techniques, and related Figs. (1, 2, 3, 7-

459   10) in the manuscript. He co-wrote much of the text and led the efforts of interpretation, analysis,

460   and presentation of the results.

461   KAH was responsible for case study identification, and collection of surface, radar, and HRRR

462   data relevant to this case study. He developed Figs. 4, 5 and 6. He also co-developed the objective

463   OFB identification process and co-wrote the text.

464   SDM was the PI of the Multidisciplinary University Research Initiative (MURI) research

465   project and was responsible for managing the development of the optical flow code and outflow

466   boundary case study information involved in Figs. (1, 2, 3, 7-10). He co-developed the objective

467   identification process and maintained and managed the satellite data necessary to complete the

468   study. He also co-wrote much of the text within the manuscript.

469   DAP co-developed the objective OFB identification scheme. He also co-developed the

470   conceptual Figs. 1 and 2 to add clarity to the optical flow process used here and co-wrote the text

471   within the manuscript.

472

473   **9.   Competing Interests**

474   The authors declare that they have no conflicts of interest.

475

476   **10.   Acknowledgements**

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638 **12. List of Tables**

639   **Table 1.** Settings used in the Brox et al. (2004) successive over-relaxation scheme.

640

13. List of Figures

Figure 1. Flow chart of the B04 optical flow approach used here. Note that SF, nK, nL and nM are defined in Table 1.

Figure 2. Schematic of coarse- to fine-scale warping optical flow in GOES-imagery. The largest displacements are found in the initial coarse grid (yellow arrow at the top of the pyramid), which are used as initial displacements for the next levels (red and blue arrows). The final displacement is the sum of each displacement estimate (white arrow). In this schematic, an example scale factor of 0.5 was used over 3 pyramid levels, in this work, a scale factor of 0.95 for 77 levels was used.

Figure 3. The 6 July 2018 0023 UTC GOES-16 0.64-µm visible reflectance (top) and $BT_{10.35}$ (bottom) over south-central AZ, centered on an OFB of interest.

Figure 4. The KIWA Radar 2244 UTC 0.5° horizontal reflectivity (top) in dBZ and correlation coefficient (bottom). Range rings in grey indicate every 30° azimuth and 50 km in range.

Figure 5. Surface High Frequency METAR observations of temperature (K; top left), dewpoint (K; top right), mean sea level pressure (middle left), wind direction (° from N; middle right), wind speed (m s^{-1}; bottom left), and wind gusts (m s^{-1}; bottom right). The surface station was located at (32.95 °N -111.77 °E). The red line indicates the approximate time of boundary passage over the station.

Figure 6. Four panel of HRRR output of OFB event, including a) wind speed, b) temperature, c) simulated infrared brightness temperature, and d) a cross section along the black line in c with virtual potential temperature θv in black contours (K), omega in color shaded pixels, and regions of relative humidity > 90% highlighted with dark shading (bottom right).

663 **Figure 7.** The 0023 UTC GOES-16 0.64-µm visible channel shown with a) subjectively

664 identified OFB (blue dots) and b) objectively identified linear features (blue shading). Also

665 shown are linear features that contained fast storm-relative motion (red shading). The results of

666 backtracking the c) subjectively and d) objectively identified OFB features are also shown,

667 where blue dots represent targets tracked back within 50 km of a deep convection event, and

668 orange dots are targets that were not.

669 **Figure 8.** GOES-16 0.64-µm visible channel imagery on 5 July 2018 at a) 2258 UTC, b) 2338

670 UTC, c) 2358 UTC, and d) 0023 UTC over central Arizona shown with every 20th optical flow

671 vector in the x and y directions (subsampled for image clarity) illustrated with yellow wind barbs

672 (knots). Circles represent motion < 5 kts, which commonly occur over ground pixels.

673 **Figure 9.** The GOES-16 0.64-µm visible imagery shown with image targets backtracked from

674 subjective identification in Fig. 7a at 0023 UTC 6 July 2018 using the B04 method (blue/yellow)

675 and the Wu et al. (2016) approach (orange/red) at a) 0023 UTC, b) 2358 UTC, c) 2338 UTC and

676 d) 2213 UTC. Individual points are highlighted from each approach (yellow and red dots; see

677 text).

678 **Figure 10.** Color shaded wind speed for 0023 UTC 6 July 2018 over central Arizona shown

679 from a) the B04 optical flow method and b) the Wu et al. (2016) flow, shown with respective

680 flow vectors and the subjective position of the front edge of the OFB (blue line).
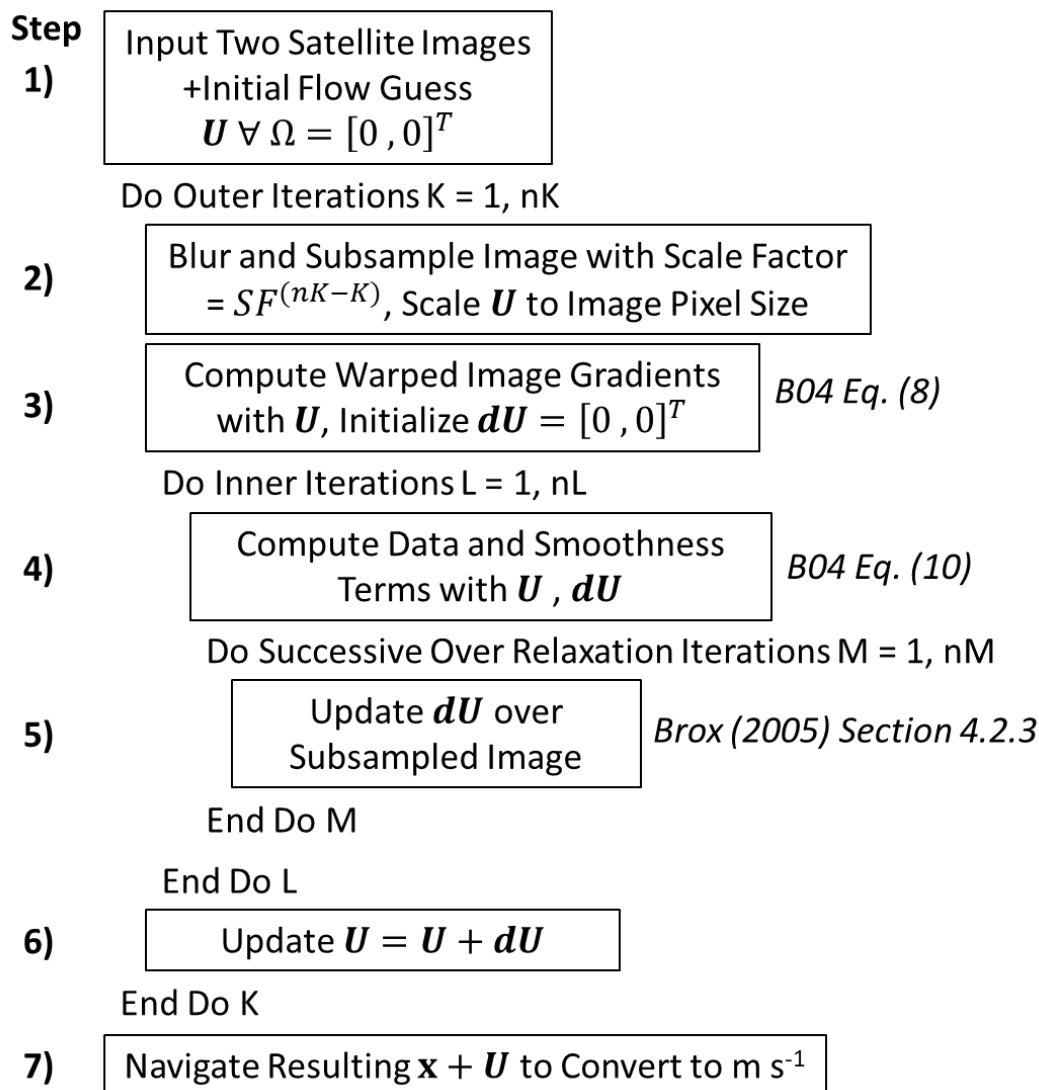
681

682

683 **14. Tables**

684   **Table 1.** Settings used in the Brox et al. (2004) successive over-relaxation scheme.

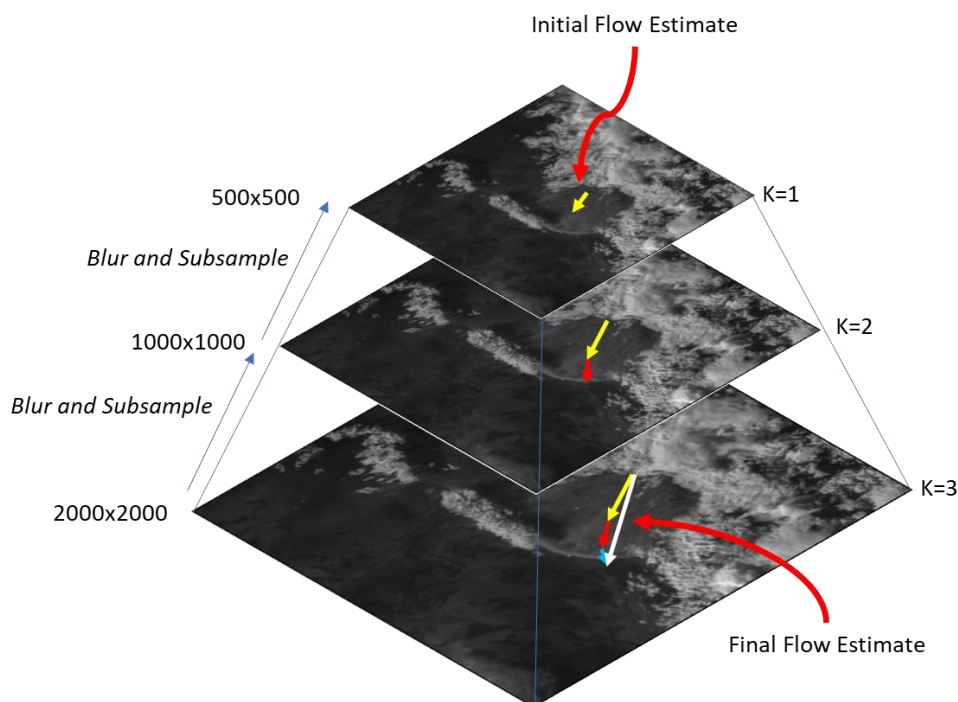| Parameter	Value
Outer Iterations (Pyramid Levels, nK)	77
Inner Iterations (nL)	10
Successive Over-Relaxation Iterations (nM)	5
Successive Over-Relaxation Parameter	1.99
Pyramid Scale Factor (SF)	0.95
γ	10
α	50

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686 **15. Figures**

Step 1) Input Two Satellite Images +Initial Flow Guess $U \forall \Omega = [0, 0]^T$

Do Outer Iterations K = 1, nK

2) Blur and Subsample Image with Scale Factor $= SF^{(nK-K)}$, Scale U to Image Pixel Size

3) Compute Warped Image Gradients with U, Initialize $dU = [0, 0]^T$ — *B04 Eq. (8)*

Do Inner Iterations L = 1, nL

4) Compute Data and Smoothness Terms with U, dU — *B04 Eq. (10)*

Do Successive Over Relaxation Iterations M = 1, nM

5) Update dU over Subsampled Image — *Brox (2005) Section 4.2.3*

End Do M

End Do L

6) Update $U = U + dU$

End Do K

7) Navigate Resulting $x + U$ to Convert to m s^{-1}

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688 **Figure 1.** Flow chart of the B04 optical flow approach used here. Note that SF, nK, nL and nM
689 are defined in Table 1.
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Figure 2. Schematic of coarse- to fine-scale warping optical flow in GOES-imagery. The largest displacements are found in the initial coarse grid (yellow arrow at the top of the pyramid), which are used as initial displacements for the next levels (red and blue arrows). The final displacement is the sum of each displacement estimate (white arrow). In this schematic, an example scale factor of 0.5 was used over 3 pyramid levels, in this work, a scale factor of 0.95 for 77 levels was used.

Figure 3. The 6 July 2018 0023 UTC GOES-16 0.64-µm visible reflectance (top) and $BT_{10.35}$ (bottom) over south-central AZ, centered on an OFB of interest.

Figure 4. The KIWA Radar 2244 UTC 0.5° horizontal reflectivity (top) in dBZ and correlation coefficient (bottom). Range rings in grey indicate every 30° azimuth and 50 km in range.

Figure 5. Surface High Frequency METAR observations of temperature (K; top left), dewpoint (K; top right), mean sea level pressure (middle left), wind direction (° from N; middle right), wind speed (m s^{-1}; bottom left), and wind gusts (m s^{-1}; bottom right). The surface station was located at (32.95 °N -111.77 °E). The red line indicates the approximate time of boundary passage over the station.

Figure 6. Four panel of HRRR output of OFB event, including a) wind speed, b) temperature, c) simulated infrared brightness temperature, and d) a cross section along the black line in c with virtual potential temperature θ_v in black contours (K), omega in color shaded pixels, and regions of relative humidity > 90% highlighted with dark shading (bottom right).

Figure 7. The 0023 UTC GOES-16 0.64-µm visible channel shown with a) subjectively identified OFB (blue dots) and b) objectively identified linear features (blue shading). Also shown are linear features that contained fast storm-relative motion (red shading). The results of backtracking the c) subjectively and d) objectively identified OFB features are also shown, where blue dots represent targets tracked back within 50 km of a deep convection event, and orange dots are targets that were not.
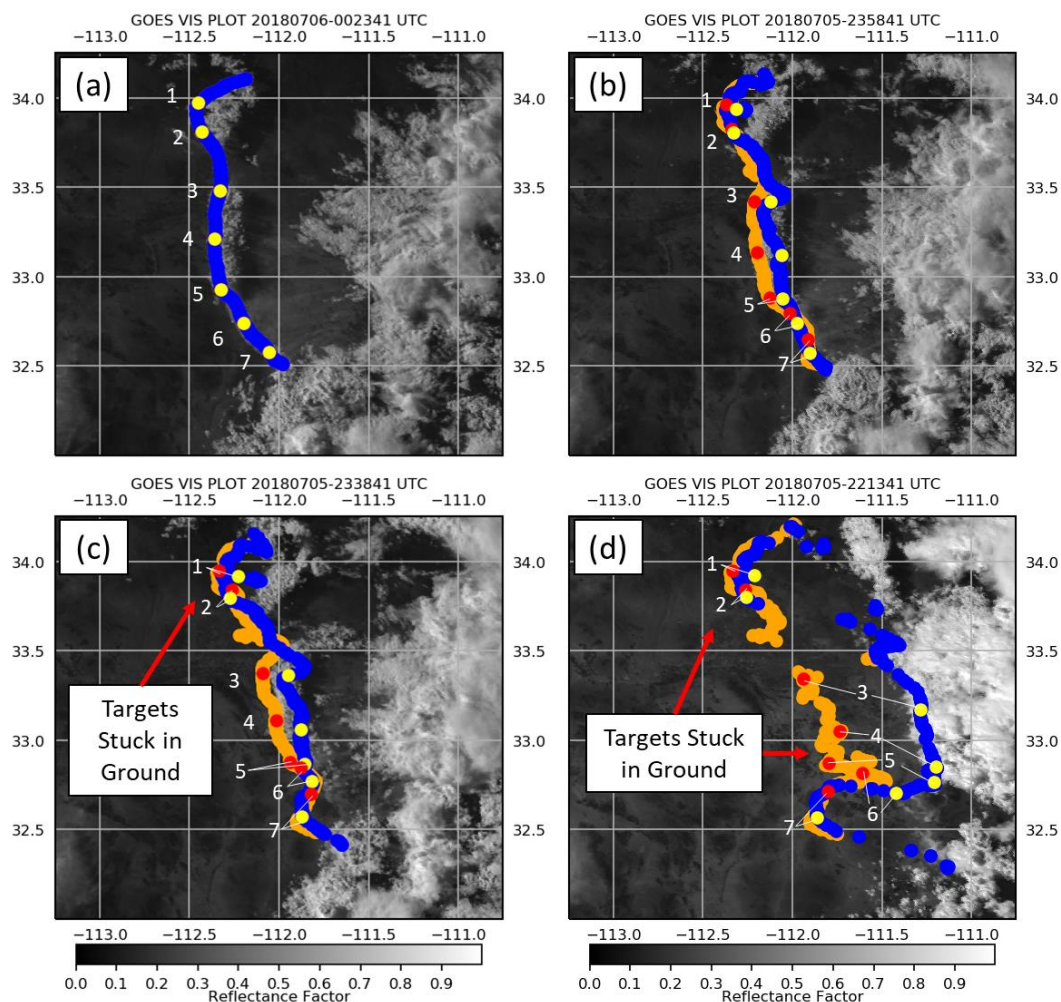
Figure 8. GOES-16 0.64-µm visible channel imagery on 5 July 2018 at a) 2258 UTC, b) 2338 UTC, c) 2358 UTC, and d) 0023 UTC over central Arizona shown with every 20th optical flow vector in the x and y directions (subsampled for image clarity) illustrated with yellow wind barbs (knots). Circles represent motion < 5 kts, which commonly occur over ground pixels.

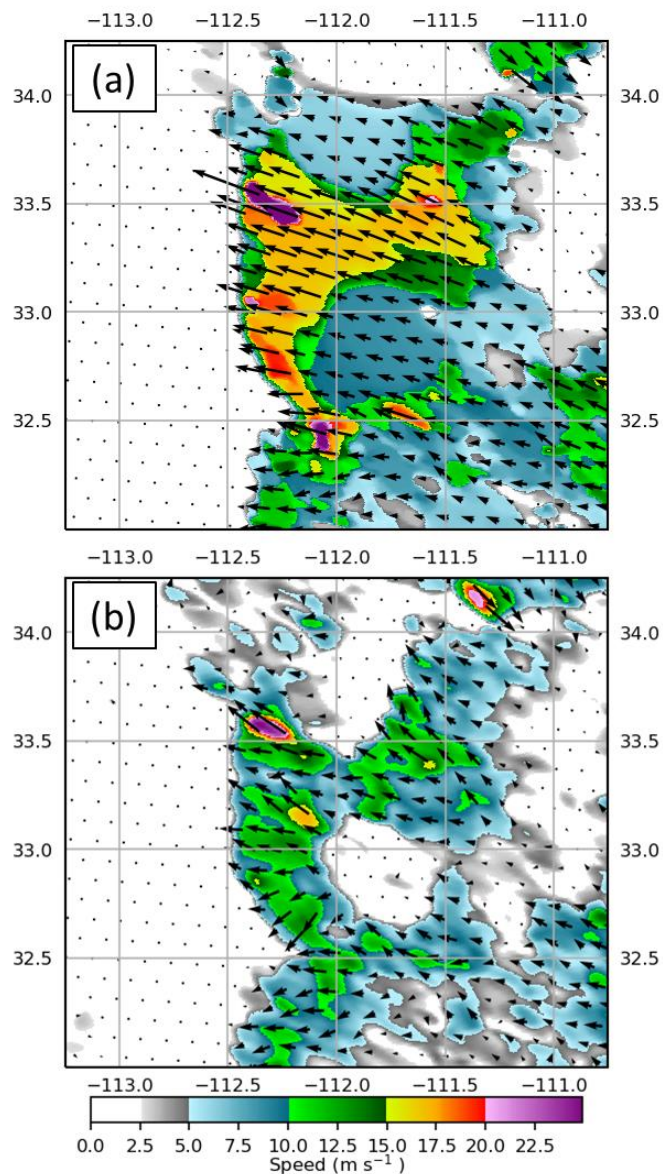**Figure 9.** The GOES-16 0.64-µm visible imagery shown with image targets backtracked from subjective identification in Fig. 7a at 0023 UTC 6 July 2018 using the B04 method (blue/yellow) and the Wu et al. (2016) approach (orange/red) at a) 0023 UTC, b) 2358 UTC, c) 2338 UTC and d) 2213 UTC. Individual points are highlighted from each approach (yellow and red dots; see text).

Figure 10. Color shaded wind speed for 0023 UTC 6 July 2018 over central Arizona shown from a) the B04 optical flow method and b) the Wu et al. (2016) flow, shown with respective flow vectors and the subjective position of the front edge of the OFB (blue line).

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