



# **Potential of next-generation imaging spectrometers to detect and quantify methane point sources from space**

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## **Abstract**

We examine the potential for global detection of methane plumes from individual point sources with the new generation of spaceborne imaging spectrometers (EnMAP, PRISMA, EMIT, SBG) scheduled for launch in 2019-2025. These instruments are designed to map the Earth's surface with a sampling distance as fine as  $30 \times 30 \text{ m}^2$  but they have spectral resolution of 7-10 nm in the 2200-2400 nm band that should also allow useful detection of atmospheric methane. We simulate scenes viewed by EnMAP (10 nm spectral resolution, 180 signal-to-noise ratio) using the EnMAP End-to-End Simulation Tool with superimposed methane plumes generated by large-eddy simulations. We retrieve atmospheric methane and surface reflectivity for these scenes using the IMAP-DOAS optimal estimation algorithm. We find an EnMAP precision of 4-13% for atmospheric methane depending on surface type, allowing effective single-pass detection of 100+  $\text{kg h}^{-1}$  methane point sources depending on surface brightness, surface homogeneity, and wind speed. Successful retrievals over very heterogeneous surfaces such as an urban mosaic require finer spectral resolution. We simulated the EnMAP capability with actual plume observations over oil/gas fields in California from the airborne AVIRIS-NG sensor ( $3 \times 3 \text{ m}^2$  pixel resolution, 5 nm spectral resolution, SNR 200-400). We spectrally and spatially downsampled AVIRIS-NG images to match EnMAP instrument specifications and found that we could successfully detect point sources of  $\sim 100 \text{ kg h}^{-1}$  over bright surfaces. Estimated emission rates inferred with a generic Integrated Mass Enhancement (IME) method agreed within a factor of 2 between EnMAP and AVIRIS-NG. Better agreement may be achieved with a more customized IME method. Our results suggest that imaging



spectrometers in space could play a transformative role in the future for quantifying methane emissions from point sources on a global scale.

## 1 Introduction

Methane is a powerful greenhouse gas, yet its sources are highly uncertain. Quantifying methane emissions from different sources is critical for developing strategies to reduce atmospheric methane levels. Anthropogenic emissions originate from a large number of point sources (coal mine vents, oil/gas facilities, livestock operations, landfills, wastewater treatment plants) that are individually small, spatially clustered, often intermittent, and difficult to quantify (Allen et al., 2013; Frankenberg et al., 2016). Here we investigate the unique potential of new-generation satellite instruments designed to map the Earth's surface (imaging spectrometers) to also detect methane point sources in the shortwave infrared (SWIR).

There has been considerable interest in using SWIR satellite observations of atmospheric methane columns by solar backscatter to detect methane sources and test emission inventories (Jacob et al., 2016). These observations are traditionally made by atmospheric sensors with high spectral resolution ( $<1$  nm) to capture the fine structure of methane absorption lines (Table 1). The requirement of high spectral resolution has generally implied a coarse pixel resolution ( $>1$  km) to achieve satisfactory signal-to-noise ratios (SNR), but this limits the ability to localize and quantify individual point sources. Inverse analyses of observations from the SCIAMACHY instrument with 60 km pixel resolution, and from the GOSAT instrument with sparse sampling at 10 km pixel resolution, have quantified emissions over regional scales (Bergamaschi et al., 2009; Kort et al., 2014; Turner et al., 2015). The recently launched TROPOMI instrument with global daily coverage at 7 km pixel resolution (Hu et al., 2018) will refine the regional characterization but still cannot resolve point sources (Sheng et al., 2018). Planned instruments with  $\sim 1$  km pixel resolution (MethaneSat, Propp et al., 2017; Geo-FTS, Xi et al., 2016) should be able to detect large point sources after inversion of several days of observations (Cusworth et al., 2018; Turner et al., 2018) but would not resolve densely clustered or temporally variable sources.

Atmospheric sensors for methane generally focus on achieving high precision ( $<1\%$ ) and low relative bias ( $<0.3\%$ ), appropriate for regional characterization of sources (Buchwitz et al., 2015). However, these requirements can be relaxed if the focus is to observe individual plumes. Precision can be traded for pixel resolution because methane plumes are generally sub-kilometer in scale (Frankenberg et al., 2016), so that plume enhancements are larger when the pixel resolution is finer (Jacob et al., 2016). In a theoretical simulation study for the GHGSat microsatellite instrument



with 50×50 m<sup>2</sup> pixel resolution, Varon et al. (2018a) found that a 1-5% precision would be adequate for single-pass observation of plumes to quantify point sources of magnitude ~100 kg h<sup>-1</sup>. This would account for most of the total methane emitted by point sources in the United States reporting to the Greenhouse Gases Reporting Program (Jacob et al., 2016). A demonstration GHGSat instrument (GHGSat-D) launched in 2016 with an estimated precision of 13% limited by instrument imperfections, has proven able to detect large point sources in excess of 1000 kg h<sup>-1</sup> (Varon et al., 2018b).



Here we examine the potential of a different class of satellite instruments, imaging spectrometers, to provide global snapshots of individual methane point sources. These instruments are designed for global coverage of land surfaces, but they may be used for atmospheric sensing as well. They have fine spatial sampling, or pixel resolution (<100 m), with coarser spectral sampling to measure vibrational overtone absorptions in surface reflectance. Some current imagers such as Landsat (Roy et al., 2014) and WorldView-3 (<http://worldview3.digitalglobe.com>) have observing bands in the SWIR intended to infer soil moisture, mineral composition, and vegetation traits (Cleemput et al., 2018). However, the SWIR spectral resolutions for Landsat (100 nm) and WorldView-3 (40-50 nm) are too coarse to usefully observe methane. The Hyperion instrument onboard NASA Earth Observing-1 had 10 nm spectral resolution in the SWIR but a very low signal to noise ratio (SNR) of 20 (Folkman et al., 2001).

A new generation of imaging spectrometers set for launch over the next few years (EnMAP, PRISMA, EMIT, and an anticipated SBG investigation) will achieve ~10 nm or better spectral resolution in the SWIR with pixel resolution in the range 30-60 m and SNR of 180-400 or beyond (Table 1). Experience with airborne imaging spectrometers of comparable specifications suggests that these satellite instruments should be able to usefully observe methane plumes. The Airborne Visible/Infrared Imaging Spectrometer (AVIRIS-C), with a 10 nm spectral resolution and SNR of 70 (Green et al., 1998), was able together with Hyperion to detect the massive Aliso Canyon methane leak in California (Thompson et al., 2016). The next generation AVIRIS instrument (AVIRIS-NG), with a finer spectral resolution of 5 nm and SNR of 200 (Thorpe et al., 2014), was able to detect a range of methane plumes over the Four Corners region of New Mexico including from gas processing facilities, storage tanks, pipeline leaks, well pads, and coal mine venting shafts (Frankenberg et al., 2016). AVIRIS-NG was flown over 272000 potential methane emitting facilities in California between 2016 and 2018 (CARB, 2017).

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## 87 2 Imaging spectrometer spectra including methane plumes



88 The next generation of spaceborne imaging spectrometers in Table 1 includes PRISMA (launched March  
 89 2019), EnMAP (2020), EMIT (2022), SBG (2025-2027). The AMPS instrument (proposed) would bridge the gap  
 90 between surface imagers and methane sensors, by providing 1 nm SWIR spectral resolution while maintaining 30 m  
 91 spatial resolution (Thorpe et al., 2016). We will focus our baseline analysis on EnMAP, for which detailed  
 92 documentation is available (Guanter et al., 2015), and examine other instruments through sensitivity analyses. EnMAP  
 93 is a push-broom style instrument with 10 nm resolution in the SWIR and an expected 180 SNR at 2300 nm. PRISMA  
 94 (<http://www.prisma-i.it/>) has very similar instrument specifications as EnMAP. The EMIT instrument will fly on the  
 95 International Space Station. It is slated to have a 7-10 nm spectral resolution and 60 m pixel resolution (Green et al.,  
 96 2018). Other investigations, such as SBG, are called for in the NASA Earth Science and Applications Decadal Survey  
 97 (National Academies, 2018).

98 Figure 1 shows simulated transmission spectra in the weak (~1650 nm) and strong (~2300 nm) SWIR methane  
 99 absorption bands at the spectral resolutions of TROPOMI (0.25 nm FWHM), AVIRIS-NG (5 nm), and EnMAP (10  
 100 nm). EnMAP spectra are sampled following the precise wavelength positions given in Guanter et al. (2015). The 1650  
 101 nm methane band has the advantage of being near a CO<sub>2</sub> band, so that joint retrievals of methane and CO<sub>2</sub> can be  
 102 combined with independent knowledge of the CO<sub>2</sub> column mixing ratio to remove joint errors in surface reflectivity  
 103 and atmospheric scattering (the so-called “CO<sub>2</sub> proxy” method; Frankenberg et al. 2005a). However, the 1650 nm band  
 104 is much weaker than the 2300 nm band and only the 2v Q-branch could be detected at coarser spectral resolution.  
 105 Sampling the transmission spectra at the EnMAP spectral resolution yields only 8 data points in the 1650 nm band as  
 106 compared to 25 in the 2300 nm band. The 2300 nm band also exhibits more resolved structure. Our early attempts to  
 107 use the CO<sub>2</sub> proxy method in the 1650 nm band with EnMAP synthetic spectra were unsuccessful. In what follows we  
 108 focus on the 2300 nm band as sampled in the useful 2210 - 2410 nm range.

109 We examined the sensitivity of EnMAP to atmospheric methane variability by generating synthetic top of  
 110 atmosphere (TOA) EnMAP scenes with variable methane over a variety of surface types. We used for this purpose the  
 111 EnMAP End-to-End Simulation Tool (EeteS; Segl, 2012), developed to generate EnMAP TOA spectra with expected  
 112 instrument error included. EeteS takes surface information from another imaging instrument (e.g., SPOT-5), and passes  
 113 the image through spatial, atmospheric, spectral, and radiometric modules to generate EnMAP spectra. The  
 114 atmospheric module is based on the MODTRAN5 radiative transfer code. It assumes a horizontally invariant 1800 ppb  
 115 dry air column methane mixing ratio ( $X_{CH_4}$ ) and here we add methane plumes simulated with the Weather and Research  
 116 Forecasting Model Large Eddy Simulation (WRF-LES) at 30 × 30 m<sup>2</sup> resolution (Varon et al., 2018a).



Figure 2 shows a simulated red-blue-green (RGB) EeteS image over Berlin. We consider four scenes within this domain to add WRF-LES methane plumes and perform subsequent retrievals. The scenes - Grass, Dark (water), Bright, Urban - have mean SWIR surface reflectances of 0.09, 0.02, 0.30, 0.13, respectively. The urban scene is highly heterogeneous. The WRF-LES simulation is conducted with  $30 \times 30 \text{ m}^2$  resolution (the EnMAP pixel resolution),  $100 \text{ W m}^{-2}$  sensible heat flux (moderately unstable meteorological conditions), and a mean wind speed of  $3.5 \text{ m s}^{-1}$ . We generate an ensemble of 15 instantaneous plumes by sampling the WRF-LES simulation at five time slices and for three source rates of 100, 500, and  $900 \text{ kg h}^{-1}$ . This range is typical of large (but not unusually large) point sources (Jacob et al., 2016).



We compute the optical depth of the methane plume  $\tau(\lambda)$  at wavelength  $\lambda$  by multiplying HITRAN absorption cross sections ( $\sigma_H$ ; Kochanov et al., 2016) by the methane mixing ratio enhancement ( $\Delta VMR$ ) and density of dry air ( $VCD$ ) in the 72-layered atmosphere of the MERRA-2 meteorological reanalysis (Gelaro et al., 2017):

$$\tau(\lambda) = \sum_{i=1}^{72} \Delta VMR_i VCD_i \sigma_{H,i}(\lambda). \quad (1)$$

The plume transmission  $T(\lambda)$  is the negative exponential of  $\tau(\lambda)$  weighted by the geometric airmass factor  $A$  (AMF) for the backscattered solar radiation:

$$T(\lambda) = \exp\{-A\tau(\lambda)\}. \quad (2)$$

Each pixel's EeteS radiance spectrum is multiplied by this additional plume transmission. We do not add noise or aerosol effects to the plume transmission spectra because the EeteS scene already accounts for those in the computation of back-scattered radiances, and the plume transmission is just a multiplicative factor on these back-scattered radiances.



Figure 3 shows an example WRF-LES plume ( $500 \text{ kg h}^{-1}$  source rate) superimposed over the Grass and Urban scenes.

EnMAP has a specific spectral resolution and SNR. We examined the sensitivity of the retrieval to these parameters by generating synthetic spectra for different spectral resolutions and SNRs, thus extending our analysis to other new-generation imaging spectrometers (Table 1). For this purpose, we interpolated EeteS surface radiance spectra to the desired spectral resolution assuming no instrument noise. We then multiplied these radiance spectra by the standard atmosphere plus WRF-LES plume transmission spectra and added uncorrelated instrument noise as per the specified SNR.



To test our EnMAP retrievals on actual data, we also downsampled AVIRIS-NG images taken from aircraft over California (CARB, 2017) to match EnMAP spatial resolution, and further convolved these spectra with a 10 nm Gaussian filter to match EnMAP spectral resolution and wavelength positions. AVIRIS-NG flew at 3-4 km above the ground, so we simulated additional extinction at higher altitudes based on the U.S standard atmosphere (Kneizys et al., 1996). We compared the retrieved methane from AVIRIS-NG and the synthetic EnMAP to determine the ability of EnMAP to detect and quantify the methane point sources identified by AVIRIS-NG.

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### 152 3 Methane retrieval

We retrieved methane from the synthetic imaging spectrometer spectra by adapting the Iterative Maximum A Posteriori - Differential Optical Absorption Spectroscopy (IMAP-DOAS) algorithm developed for AVIRIS (Frankenberg et al., 2005b; Thorpe et al., 2017; Ayasse et al., 2018). DOAS retrievals isolate higher frequency features resulting from gas absorption from lower frequency features that include surface reflectance as well as Rayleigh and Mie scattering (Bovensmann et al., 2011). A polynomial term accounts for the low frequency features (Thorpe et al., 2017).

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#### 160 3.1 State vector

In addition to methane ( $\text{CH}_4$ ), the retrieval must account for variable absorption by water vapor ( $\text{H}_2\text{O}$ ) and nitrous oxide ( $\text{N}_2\text{O}$ ) over the 2210-2400 nm spectral region. We parameterize low frequency spectroscopic features as a sum of Legendre polynomials of order  $k = [0, K]$  with coefficients  $a_k$ . The state vector ( $\mathbf{x}$ ) optimized through the retrieval is therefore composed of the following elements:

$$\mathbf{x} = (s_{\text{CH}_4}, s_{\text{H}_2\text{O}}, s_{\text{N}_2\text{O}}, a_0, \dots, a_K)$$

where  $s$  is a scaling factor applied to the column mixing ratio of each gas from the U.S standard atmosphere (Kneizys et al., 1996). We do not include aerosols in the retrieval as they play little role at the relevant spatial and spectral resolution (Ayasse et al., 2018). Methane point sources generally do not co-emit aerosols.

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#### 170 3.2 Optimal estimation

To retrieve the state vector from the Eetes TOA radiances, we use a forward model similar to previous IMAP-DOAS algorithms (Thorpe et al., 2017, Ayasse et al., 2018), with a modification to the polynomial term for surface reflectance:

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$$F^h(\mathbf{x}, \lambda) = I_0(\lambda) \exp \left( -A \sum_{n=1}^3 s_n \sum_{l=1}^{72} \tau_{n,l} \right) \sum_{k=0}^K a_k P_k(\lambda) \quad (3)$$

Here  $F^h$  is the high-resolution backscattered TOA radiance at wavelength  $\lambda$ ,  $I_0(\lambda)$  is the incident TOA solar intensity,  $\tau_{n,l}$  is the default optical depth from the US standard atmosphere for trace gas element  $n = [1,3]$  of the state vector at vertical level  $l = [1,72]$ ,  $s_n$  is the scaling factor to that default optical depth optimized in the retrieval,  $P_k(\lambda)$  is the  $k^{\text{th}}$  Legendre polynomial, and the  $a_k$  are coefficients optimized in the retrieval. The optical depth  $\tau_{n,l}$  is computed in the same fashion as Equation 1, using information from the MERRA-2 reanalysis and HITRAN absorption cross sections. For satellite retrievals, the AMF is a scalar describing the optical path through the atmosphere. In Section 4.3, we apply the IMAP-DOAS algorithm to airborne AVIRIS-NG scenes and use a vector-valued AMF that depends on the height of the aircraft.

Previous IMAP-DOAS algorithms used a simple polynomial approximation for the surface reflectance, but here we use Legendre polynomials to exploit their orthogonality. We find that  $K = 4$  provides sufficient spectral resolution whereas previous applications using simple polynomials required  $K = 6$  (Ayasse et al., 2018).

We compute the TOA backscattered radiances  $F^h(\mathbf{x}, \lambda)$  over the 2210-2410 nm spectral range at 0.02 nm resolution, and assemble these in a vector  $\mathbf{F}^h(\mathbf{x})$  representing the high-resolution spectrum as simulated by the forward model for a given  $\mathbf{x}$ . We convolve this spectrum with the instrument FWHM and then sample at the known wavelength positions. For example, for EnMAP, we convolve  $\mathbf{F}^h(\mathbf{x})$  with a 10 nm FWHM and sample the resulting spectra at EnMAP's 10 nm intervals to get the low-resolution  $\mathbf{F}(\mathbf{x})$ . We also explored performing separate convolutions on the high resolution transmission and polynomial terms in Equation 3, and then multiplying them together to get  $\mathbf{F}(\mathbf{x})$ . We found little difference in the results between methods.

Observed backscattered TOA radiances ( $\mathbf{y}$ ) can be represented as

$$\mathbf{y} = \mathbf{F}(\mathbf{x}) + \boldsymbol{\epsilon} \quad (4)$$

where the observational error  $\boldsymbol{\epsilon}$  is the sum of instrument and forward model errors. As is commonly done for satellite retrievals, we assume that the forward model error is small compared to the instrument error characterized by the SNR. The forward model is non-linear so that the solution must be obtained iteratively. A Jacobian matrix is calculated for each iteration of the state vector

$$\mathbf{K}_i = \left. \frac{\partial \mathbf{F}}{\partial \mathbf{x}} \right|_{\mathbf{x}=\mathbf{x}_i} \quad (5)$$

and we employ a Gauss-Newton iteration to solve iteratively for the optimal state vector (Rodgers, 2000):



$$\mathbf{x}_{i+1} = \mathbf{x}_A + (\mathbf{K}_i^T \mathbf{S}_0^{-1} \mathbf{K}_i + \mathbf{S}_A^{-1})^{-1} \mathbf{K}_i^T \mathbf{S}_0^{-1} [y - \mathbf{F}(\mathbf{x}_i) + \mathbf{K}_i(\mathbf{x}_i - \mathbf{x}_A)] \quad (6)$$

Here  $\mathbf{S}_0 = [\epsilon\epsilon^T]$  is the observation error covariance matrix defined by the instrument SNR,  $\mathbf{x}_A$  is the prior estimate of the state vector, and  $\mathbf{S}_A$  is the prior error covariance matrix. We set a weak prior error variance for methane,  $\sigma_{\text{CH}_4}^2 = 5$ , to accommodate large plume enhancements. The prior  $X_{\text{CH}_4}$  estimate is 1800 ppb. The iterative analytical solution to the inverse problem as described by equation (6) also provides the posterior error covariance matrix ( $\hat{\mathbf{S}}$ ) as part of the solution:

$$\hat{\mathbf{S}} = (\mathbf{K}_i^T \mathbf{S}_0^{-1} \mathbf{K}_i + \mathbf{S}_A^{-1})^{-1} \quad (7)$$

$\hat{\mathbf{S}}$  gives information on the error correlation between retrieved methane and surface reflectivity, which is a major concern for methane retrievals (Butz et al., 2012).

## 4. Results and Discussion

### 4.1 EnMAP plume retrievals over different surfaces

Figure 3 shows examples of the IMAP-DOAS retrievals of 500 kg h<sup>-1</sup> and 900 kg h<sup>-1</sup> WRF-LES plumes over the Grass and Urban scenes. Near the emission source, the 500 kg h<sup>-1</sup> plume is clearly defined in the Grass scene. It is also detectable in the Urban scene but obscured by surface retrieval artifacts. The 900 kg h<sup>-1</sup> plume is better captured over both surfaces, though major retrieval artifacts remain in the Urban scene.

Varon et al. (2018a) previously estimated the theoretical ability of a satellite instrument to quantify source rates from point sources as a function of instrument precision, assuming a uniform surface reflectance. They concluded that an instrument with 1-5% precision on  $X_{\text{CH}_4}$  would be able to quantify point sources with an error of 70-170 kg h<sup>-1</sup>. Here we characterize the EnMAP instrument precision as the relative root-mean squared-error (RRMSE) between the true and retrieved column methane concentrations for individual 30 × 30 m<sup>2</sup> pixels in the scenes of Figure 2 including the WRF-LES plumes. Figure 4 summarizes the results for the four scenes of Figure 2. We find precisions of 8.2 ± 0.7% for Grass, 13 ± 0.7% for Urban, and 3.7 ± 0.5% for Bright scenes. The standard deviations refer to the RRMSEs computed for the 15 different realizations of the WRF-LES plumes and for the 3 source rates of 100, 500, and 900 kg h<sup>-1</sup>. The Dark scene was consistently unsuccessful, with error of at least 100% for each realization, and we do not discuss it further. The Bright scene performs the best because of the large backscattered photon flux. The Urban scene performs worse than the Grass scene, even though its average SWIR surface reflectance is larger, due to the larger variability in





229 reflectance over the scene including dark pixels. As illustrated in Figure 3, the 8% precision over the relatively uniform  
 230 grass surface should enable EnMAP to successfully quantify 500 kg h<sup>-1</sup> point sources in a single pass.

231 Beyond the precision for the methane retrieval, an additional limitation for retrieving point sources is the error  
 232 correlation with variable surface reflectance. This is illustrated in Figure 3 with the retrieved  $X_{CH_4}$  enhancements over  
 233 Grass and Urban scenes relative to the background. In the case of the Grass scene with a 500 kg h<sup>-1</sup> source, the 8%  
 234 precision limits the ability to observe the downwind plume but there is a clear enhancement over background at the  
 235 source location. With a 900 kg h<sup>-1</sup> source the downwind plume becomes well-defined against the background. In the  
 236 case of the Urban scene, the detection of the 500 kg h<sup>-1</sup> plume is far more problematic because of large positive artifacts  
 237 over dark (water) pixels. The 900 kg h<sup>-1</sup> plume is still difficult to distinguish from the artifacts and would require prior  
 238 knowledge of source location to be identified and quantified. The error correlation between methane and surface  
 239 reflectance in the retrieval can be reduced by increasing the spectral resolution of the instrument as discussed in Section  
 240 4.2.

#### 242 4.2 Sensitivity to instrument spectral resolution and SNR

243 Here we examine the potential of future instruments with improved spectral resolution and SNR relative to  
 244 EnMAP (Table 1) to achieve improved retrievals of point sources. Figure 5 shows the change in retrieval precision as  
 245 we vary the spectral resolution from 10 to 1 nm and the SNR from 100 to 500. The precision estimates are calculated  
 246 using two methods. First, we estimate the precision by evaluating the RRMSEs averaged over the Grass, Urban, and  
 247 Bright scenes of Figure 2, for 3 source rates and 15 instantaneous plume realizations, following the procedure of  
 248 Section 4.1. Since SNR varies on a per-pixel basis, the plotted SNRs for this method represent the mean scene SNR.  
 249 Specifications of the instruments in Table 1 are identified on the plot. Precision improves as spectral resolution and  
 250 SNR increase, as expected. The dependencies are not linear, and the contours are concave, meaning that precision is  
 251 more effectively improved by increasing spectral resolution by a certain factor than by increasing SNR by the same  
 252 factor. Increasing the spectral resolution improves precision through multiple independent factors: by increasing the  
 253 number of independent measurements across the methane interval; by increasing the effective squared depth of the  
 254 sharpest methane absorptions, for improved spectral contrast relative to the continuum; and by better resolution of the  
 255 unique methane absorption shape, which improves discrimination against potential surface confusers.

256 Second, we estimate theoretical precision in Figure 5 by extracting the associated  $X_{CH_4}$  posterior error  
 257 covariance term of  $\hat{S}$  from Equation 7. Here we find that instrument precision improves more as a function of SNR than



spectral resolution, which is a different result than the first precision method. Issues with the surface retrieval drive the contrasting results between the two methods. This underscores the difficulty in assigning a single retrieved  $X_{CH_4}$  uncertainty value for different instrument configurations. For a spaceborne AVIRIS-NG instrument, multiple along-track samples would increase the SNR as a function of  $\sqrt{N}$ , where  $N$  = number of along-track frames. For the second precision method, doing multiple along-track samples improves the theoretical precision from 5% to 1%. Varon et al. (2018a) found that an instrument with 5% precision could constrain most anthropogenic point sources above 170 kg h<sup>-1</sup>. Using both the RRMSE and theoretical precision methods of Figure 5, we find that a spaceborne AVIRIS-NG instrument (spectral resolution 5 nm, SNR 200-400) would have a precision of 5.5 – 1.0%, meaning that such an instrument could constrain a majority of anthropogenic methane point sources.

A benefit of increasing spectral resolution is to improve decoupling of surface and methane spectroscopic features. We saw in Figure 3 that this was a major source of error over inhomogeneous surfaces such as the Urban scene. It is manifested in the retrieval by an error correlation between state vector elements  $s_{CH_4}$  (scaling factor for methane column mixing ratios) and  $a_k$  (coefficients for the surface reflectivity described by Legendre polynomials). This error correlation is described by the posterior error covariance matrix  $\hat{\mathbf{S}}$  obtained as part of the retrieval (Equation 6). For example, the error correlation decreases significantly between EnMAP ( $r = -0.33$ ) and AMPS ( $r = -0.19$ ). This driven by the increase in spectral resolution from 10 nm to 1 nm. A separate test shows that simply increasing the SNR to 300 (as for SBG) does not improve the error correlation.

An important implication of decoupling  $X_{CH_4}$  from the surface reflectance in the retrieval is to improve the capability for plume pattern recognition, which is necessary to convert observed plume methane enhancements into source rates (Varon et al., 2018a). Figure 6 illustrates this for the Grass and Urban scenes of Figure 3 including the plume from the 500 kg h<sup>-1</sup> point source. Following Varon et al. (2018a), we define the plume for the retrieved scenes with a plume mask that applies median and Gaussian filters to pixels above the 80<sup>th</sup> percentile of  $X_{CH_4}$  within the scene. Retrievals are performed with the specifications of the EnMAP instrument (10 nm spectral resolution, SNR 180), SBG (10 nm, 300), and AMPS (1 nm, 400).

For the Grass scene we find that all three instruments can discern the plume pattern near the emission source and separate it from surface features. SBG and AMPS capture larger plume domains because of their higher precisions (Figure 5), but a source rate can still be estimated successfully with EnMAP by taking into account the dependence of the retrieved plume extent on instrument precision (Varon et al., 2018a). For the Urban scene, EnMAP plume detection is swamped by surface artifacts. Simply increasing the SNR as in the SBG instrument does not improve the situation.



287 Increasing the spectral resolution to 1 nm as in the AMPS instrument enables detection of the plume though  
 288 quantification would still be prone to surface artifacts.

289

#### 290 4.3 Evaluation with AVIRIS-NG observations

291 To test the EnMAP retrieval capability with actual observations, we downsampled AVIRIS-NG spectra taken  
 292 over California methane emitting facilities (CARB, 2017). We chose three scenes observed by AVIRIS-NG on  
 293 different days over oil and gas facilities. Figure 7 shows the RGB images, the AVIRIS-NG plume retrievals performed  
 294 by applying the method of Section 3 with a variable AMF, and the downsampled EnMAP retrievals. Plume masks  
 295 have been applied in the same way as for Figure 6. At the altitudes used for the California survey, AVIRIS-NG has 3×3  
 296 m<sup>2</sup> pixel resolution and hence features much sharper methane enhancements than EnMAP (note the different scales for  
 297 the middle and right panels).

298 We see from Figure 7 that EnMAP is able to detect the same plumes as AVIRIS-NG (two plumes in the  
 299 bottom panels). This is facilitated by the brightness of the surfaces. The surface reflectivities retrieved simultaneously  
 300 with the methane enhancements in our IMAP-DOAS algorithm are 0.39-0.49, brighter than the Bright EeteS scene in  
 301 Section 4.1.

302 The plume observations can be related to the corresponding source rates by computing the integrated mass  
 303 enhancements (IME) within the plume mask (Frankenberg et al., 2016; Varon et al., 2018a). The IME is calculated as:

$$304 \quad \text{IME} = \sum_{i=1}^N \Delta\Omega_i \Lambda_i \quad (7)$$

305 where  $\Delta\Omega_i$  is the plume mass enhancement in pixel  $i$  relative to background (kg m<sup>-2</sup>),  $\Lambda_i$  is the corresponding area of  
 306 the pixel, and the summation is over the  $N$  pixels within the plume mask. The point source rate  $Q$  is then inferred from  
 307 the IME as (Varon et al., 2018a)

$$308 \quad Q = \frac{U_{eff}}{L} \text{IME} \quad (8)$$

309 where  $L = \sqrt{\sum_{i=1}^N \Lambda_i}$  is a characteristic plume size and  $U_{eff}$  is an effective wind speed describing the rate of turbulent  
 310 dissipation of the plume ( $L/U_{eff}$  is the lifetime of the plume against turbulent dissipation to below the detection limit).  
 311 Varon et al. (2018a) relate  $U_{eff}$  to the 10-m wind speed ( $U_{10}$ ) by fitting to WRF-LES simulations. Here we use their  
 312 relationship derived for the a 50 m pixel resolution, 5% precision instrument (Varon et al., 2018), and apply it as a  
 313 rough approximation to the AVIRIS-NG and downsampled EnMAP plumes:

$$314 \quad U_{eff} = 1.1 \log U_{10} + 0.6 \quad (9)$$



where  $U_{eff}$  and  $U_{10}$  are in units of  $[m\ s^{-1}]$ . We obtain  $U_{10}$  from the HRRR-Reanalysis at 3-km hourly resolution (<https://rapidrefresh.noaa.gov/>).

Figure 7 shows the source rates inferred from the AVIRIS-NG and EnMAP retrievals for each point source. The AVIRIS-NG source rates are a factor of 1.2-3.0 greater (average 1.9) than the EnMAP source rates. There could be several factors behind this discrepancy including error correlation with surface reflectivity in the EnMAP retrieval that would cause some loss of the plume, and use of a generic plume mask and IME algorithm for both instruments. As pointed out by Varon et al. (2018a), the  $U_{10}$ - $U_{eff}$  relationship needs to be tailored to the pixel resolution and precision of the particular instrument, and to the choice of plume mask. Nevertheless, the results do confirm that EnMAP should be able to detect plumes and quantify source rates down to  $\sim 100\ kg\ h^{-1}$  when the scene is sufficiently bright.



## 5 Conclusions

We examined the potential of next-generation spaceborne imaging spectrometers (EnMAP, PRISMA, EMIT, SBG,) for observing atmospheric methane plumes from point sources and inferring the corresponding source rates. These instruments have launch dates of 2019-2025 and focus on observing the Earth surface with fine pixel resolution ( $30 \times 30\ m^2$ ), but they also have observing channels at 2200-2400 nm with 7-10 nm spectral resolution that could be used to retrieve methane plumes. This would achieve much finer spatial resolution than the standard satellite instruments designed to measure atmospheric methane, and would provide a unique resource for global mapping of individual methane point sources.

We focused our baseline analysis on EnMAP (spectral resolution 10 nm, SNR 180, 2020 launch date) as its specifications are well documented (Guanter et al, 2015). We created synthetic spectra using the EnMAP End-to-End Simulation Tool (EeteS) to simulate various surface scenes (Grass, Urban, Bright) with instrument errors and with superimposed methane plumes generated by a WRF Large Eddy Simulation (LES). We then retrieved these scenes for atmospheric methane together with surface reflectivities using the Iterative Maximum A Posteriori - Differential Optical Absorption Spectroscopy (IMAP-DOAS) approach. The resulting precisions for methane are 8% for the Grass scene, 13% for Urban, and 4% for Bright. A  $500\ kg\ h^{-1}$  methane plume (typical of very large point sources) is readily detected over the relatively homogeneous Grass surface. The highly heterogeneous Urban surface is much more challenging because of retrieval artifacts.

The limitation of EnMAP in detecting methane plumes over heterogeneous surfaces is caused by error correlation between methane and surface reflectivity in the retrieval. We examined how precision and error correlation



could be improved by increasing spectral resolution and SNR. We find that spectral resolution reduces error correlation more important than SNR. The proposed Atmospheric Methane Plume Spectrometer (AMPS), which bridges the gap between imaging spectrometers and atmospheric sensors (1 nm spectral resolution, SNR 400), can greatly decrease surface artifacts and detect a 500 kg h<sup>-1</sup> plume even over the heterogeneous Urban surface. Alternative surface parameterizations might also improve  $X_{CH_4}$  and surface separation. For example, a channelwise representation with reflectances tied through an empirical covariance structure (Thompson et al., 2018) has been used previously to improve consistency in water vapor estimations. Alternative algorithms, such as matched filter approaches (Ong et al., 2019) may show different  $X_{CH_4}$  sensitivities, and in particular may be better able to represent structured reflectances of more complex surfaces.

We tested the EnMAP capability with actual observations by downsampling AVIRIS-NG images taken from aircraft (3 × 3 m<sup>2</sup> pixels, 5 nm spectral resolution, SNR 200) over California methane emitting facilities (CARB, 2017). We showed that these EnMAP-like images are able to detect and quantify actual plumes of magnitude ~100 kg h<sup>-1</sup> over relatively bright surfaces. Source rates inferred from the plumes with a generic Integrated Mass Enhancement (IME) method are a factor of 1.2 to 3 lower for EnMAP than for AVIRIS-NG, which could be due in part to unaccounted dependence of the IME method on instrument pixel size and precision.

In summary, our analysis shows that future spaceborne imaging spectrometers designed to map the Earth surface in the SWIR also have considerable potential for detecting methane plumes from point sources and quantifying source rates. The detection capability of 100-500 kg h<sup>-1</sup> over relatively bright or homogeneous land surfaces would allow accounting for a wide range of point sources. The fine spatial resolution of these instruments should make them a unique resource to contribute to tiered observing systems for greenhouse gases (Duren and Miller, 2012).

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**Table 1.** Shortwave infrared (SWIR) remote sensors for observing methane point sources

Instrument	Pixel size (km <sup>2</sup> )	SWIR spectral range (nm) <sup>a</sup>	Spectral resolution (nm) <sup>b</sup>	Signal-to-noise ratio (SNR) <sup>c</sup>	Observing epoch
<i>Aircraft</i>					
AVIRIS-NG <sup>d</sup>	0.003 × 0.003	1600–1700; 2200–2510	5.0	200–400 <sup>e</sup>	Campaigns
<i>Satellite</i>					
<i>Atmospheric sensors</i>					
SCIAMACHY <sup>f</sup>	30 × 60	1630–1670	1.4	1500	2002–2012
GOSAT <sup>g</sup>	10 × 10	1630–1700	0.06	300	2009–
GHGSat <sup>h</sup>	0.05 × 0.05	1600–1700	0.1	TBD	2016–
TROPOMI <sup>i</sup>	7 × 7	2305–2385	0.25	100	2017–
AMPS <sup>j</sup>	0.03 × 0.03	1990–2420	1.0	200–400	Proposed
<i>Imaging spectrometers</i>					
PRISMA <sup>k</sup>	0.03 × 0.03	1600–1700; 2200–2500	10	180	2019–
EnMAP <sup>l</sup>	0.03 × 0.03	1600–1700; 2200–2450	10	180	2020–
EMIT <sup>m</sup>	0.06 × 0.06	1600–1700; 2200–2510	7–10	200–300	2022–
SBG <sup>n</sup>	0.03 × 0.03	1600–1700; 2200–2510	7–10	200–300	2025–

<sup>a</sup>Methane has absorption bands near 1650 and 2300 nm.

5 <sup>b</sup>Spectral resolution is represented by the full-width at half-maximum (FWHM).

<sup>c</sup>For SCIAMACHY and GOSAT, SNR is for CO<sub>2</sub> band used in the CO<sub>2</sub>-proxy method retrieval. For other instruments, SNR is at 2300 nm.

10 <sup>d</sup>Airborne Visible/Infrared Imaging Spectrometer – Next Generation (Thorpe et al., 2017). AVIRIS-NG provides roughly a ground sampling distance (GSD) of 1 m per km altitude. The Frankenberg et al. (2016) and Duren et al. (2019) campaigns operated at 3–4 km altitude.

<sup>e</sup>Along-track oversampling increases SNR by  $\sqrt{N}$  where  $N$  = number of along-track frames. AVIRIS-NG routinely achieves  $N > 4$  so AVIRIS-NG effective SNR at 2300 nm can be as much as 400.

<sup>f</sup>Scanning Imaging Absorption SpectroMeter for Atmospheric CHartography (Frankenberg et al., 2006)



<sup>e</sup>TANSO-FTS instrument aboard the Greenhouse gases Observing SATellite (Kuze et al., 2016). Pixels are circles of 10 km diameter.

<sup>b</sup>GreenHouse Gases Satellite (McKeever et al., 2017). Revisit times are for selected  $12 \times 12 \text{ km}^2$  scenes. The demonstration GHGSat-D instrument presently in space has additional instrument imperfections that limit its precision to 13% (McKeever

5 et al. 2017).

<sup>i</sup>TROPOspheric Monitoring Instrument (Hu et al., 2018)

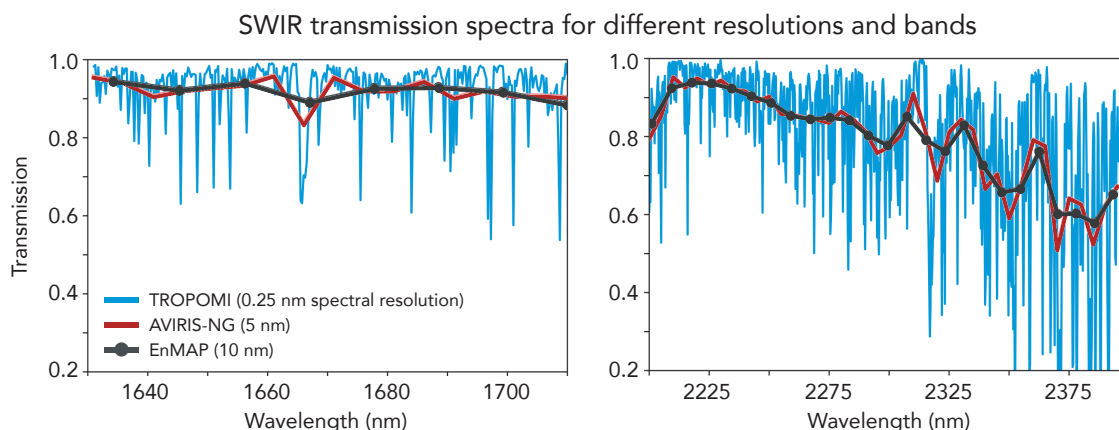
<sup>j</sup>Airborne Methane Plume Spectrometer (Thorpe et al., 2016)

<sup>k</sup>PRecursore IperSpettrale della Missione Applicativa (<http://prisma-i.it>)

<sup>l</sup>Environmental Mapping and Analysis Program (Guanter et al., 2015)

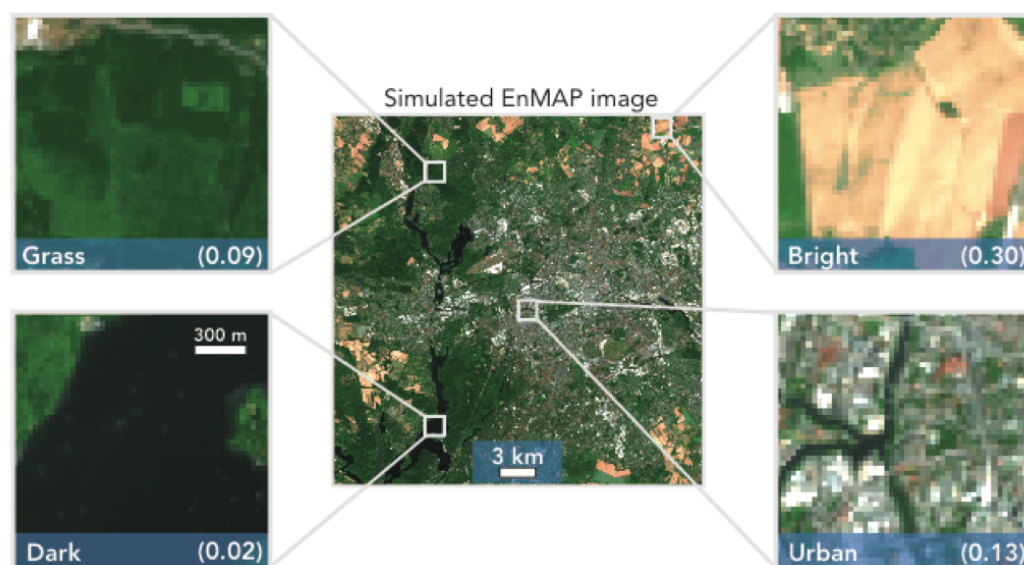
10 <sup>m</sup>Earth Surface Mineral Dust Source Investigation (Green et al., 2018)

<sup>n</sup>Surface Biology and Geology, previously called HypsIRI (Hochberg et al., 2015)

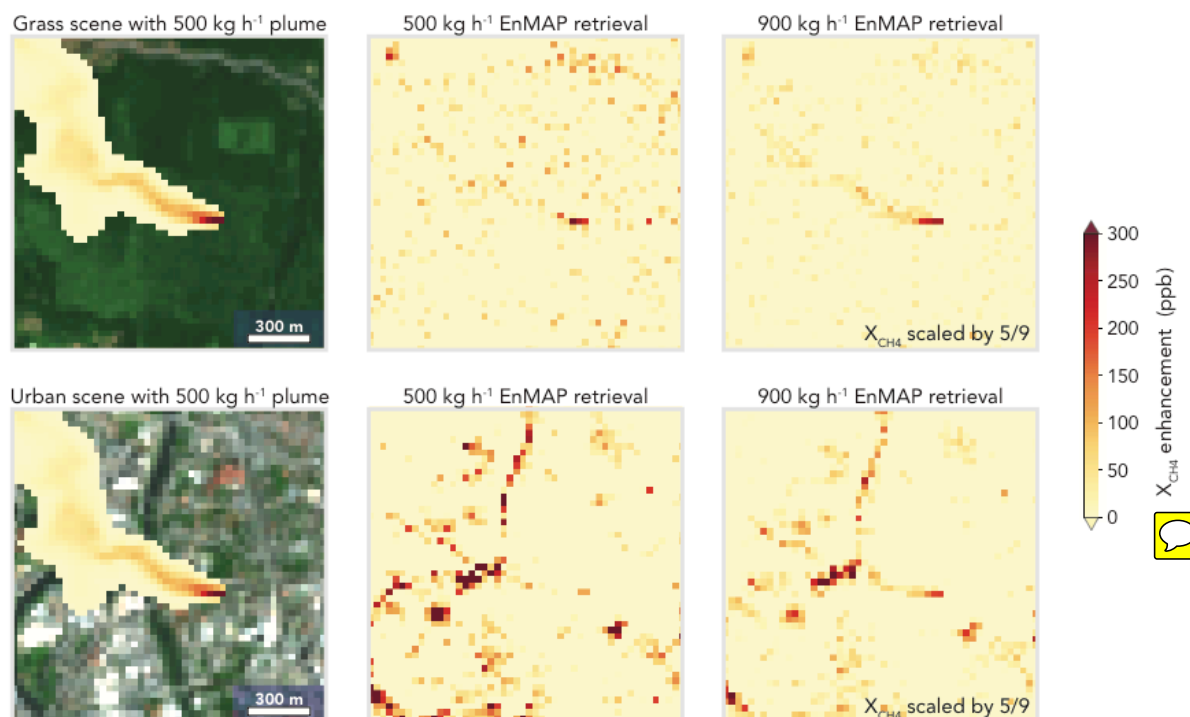


**Figure 1.** Simulated top of the atmosphere (TOA) transmission spectra for different spectral resolutions (FWHM = full-width at half-maximum) in the 1650 nm (left panel) and 2300 nm (right panel) shortwave infrared (SWIR) bands. High-resolution spectra were simulated for the U.S. Standard Atmosphere with 1800 ppb total column methane using the HITRAN spectroscopic database and the HITRAN Application Programming Interface (HAPI) tool (Kochanov et al., 2016), and were then sampled with spectral resolutions of 0.25 nm (TROPOMI), 5 nm (AVIRIS-NG), and 10 nm (EnMAP) at the appropriate wavelength positions.

10

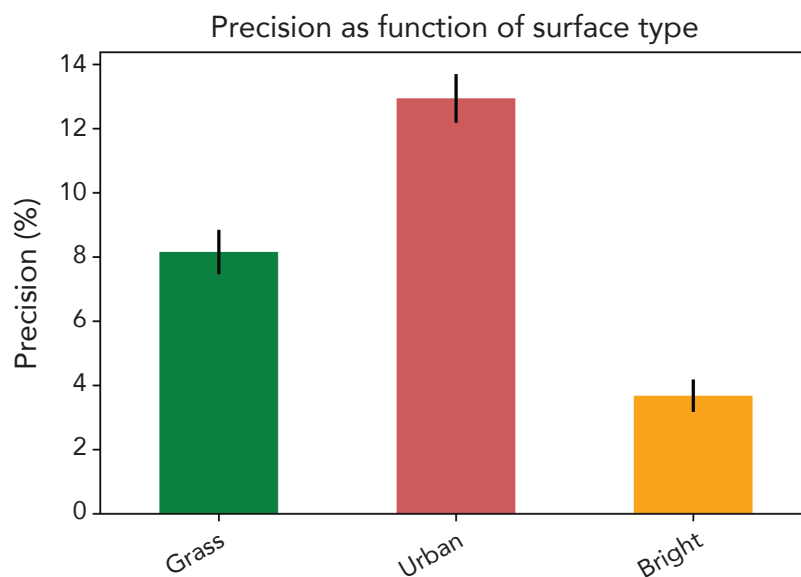


**Figure 2.** RGB image of a synthetic EnMAP scene simulated using the EnMAP End-to-End Simulation Tool (EetsS) over Berlin. Four scenes with  $30 \times 30 \text{ m}^2$  pixel resolution are shown (Grass, Dark, Bright, Urban) with average surface reflectances in the SWIR (2210–2410 nm) given in parentheses. These different scenes are used in Section 3 to evaluate the sensitivity of EnMAP to atmospheric methane.

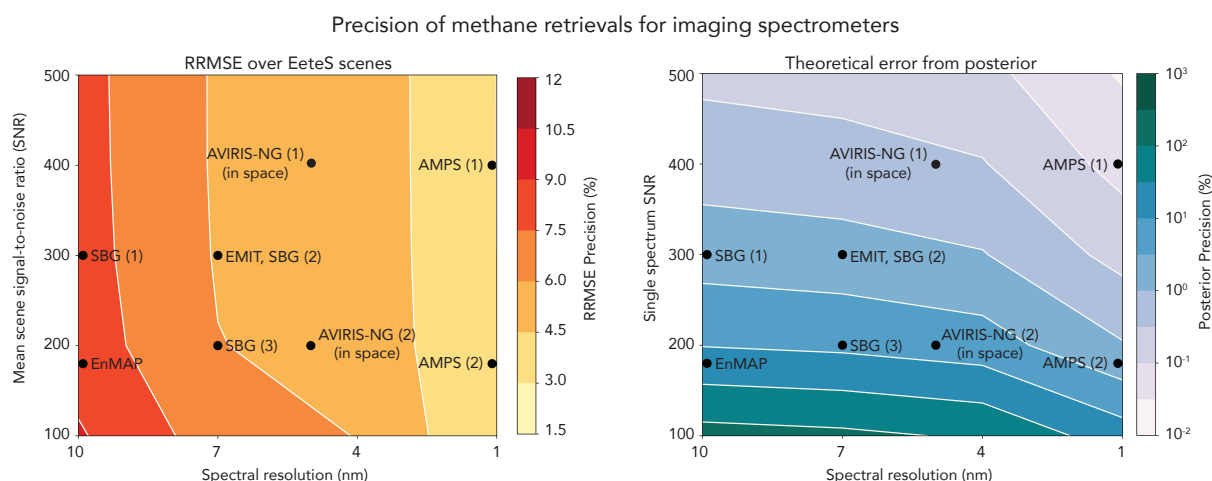


**Figure 3.** Retrieval of a methane plume over grass and urban EnMAP scenes. The plume was generated by WRF-LES with a source rate of either 500 kg h<sup>-1</sup> or 900 kg h<sup>-1</sup>. The left panels show the dry air column mixing ratio enhancements relative to the 1800 ppb background for a 500 kg h<sup>-1</sup> methane plume superimposed on the RGB images of Figure 2. The middle panels show the retrieval of those enhancements using the IMAP-DOAS retrieval algorithm applied to the EnMAP instrument specifications. The right panels show the retrieval of the 900 kg h<sup>-1</sup> plume. The  $X_{\text{CH}_4}$  enhancements in the right panels are scaled by 5/9 to be comparable with the other panels. Negative enhancements are reset to equal the background.

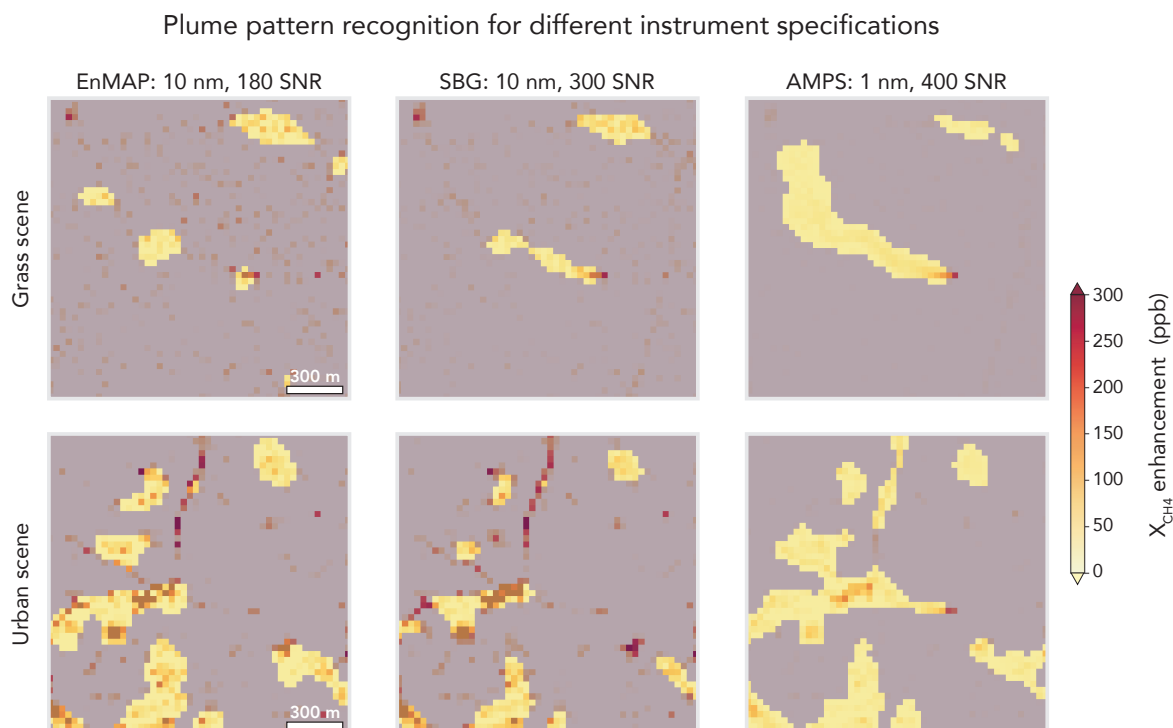




**Figure 4.** Precision of atmospheric methane retrievals from the EnMAP instrument (Table 1) over different surfaces. The precisions are the relative root-mean squared errors (RRMSE) between the “true” methane columns in synthetic scenes and values obtained from the IMAP-DOAS retrieval applied to the EnMAP top-of-atmosphere (TOA) backscattered radiances. The error bars represent the standard deviation over 15 WRF-LES plume realizations and 3 source magnitudes for the plume (100, 500, 900 kg h<sup>-1</sup>).



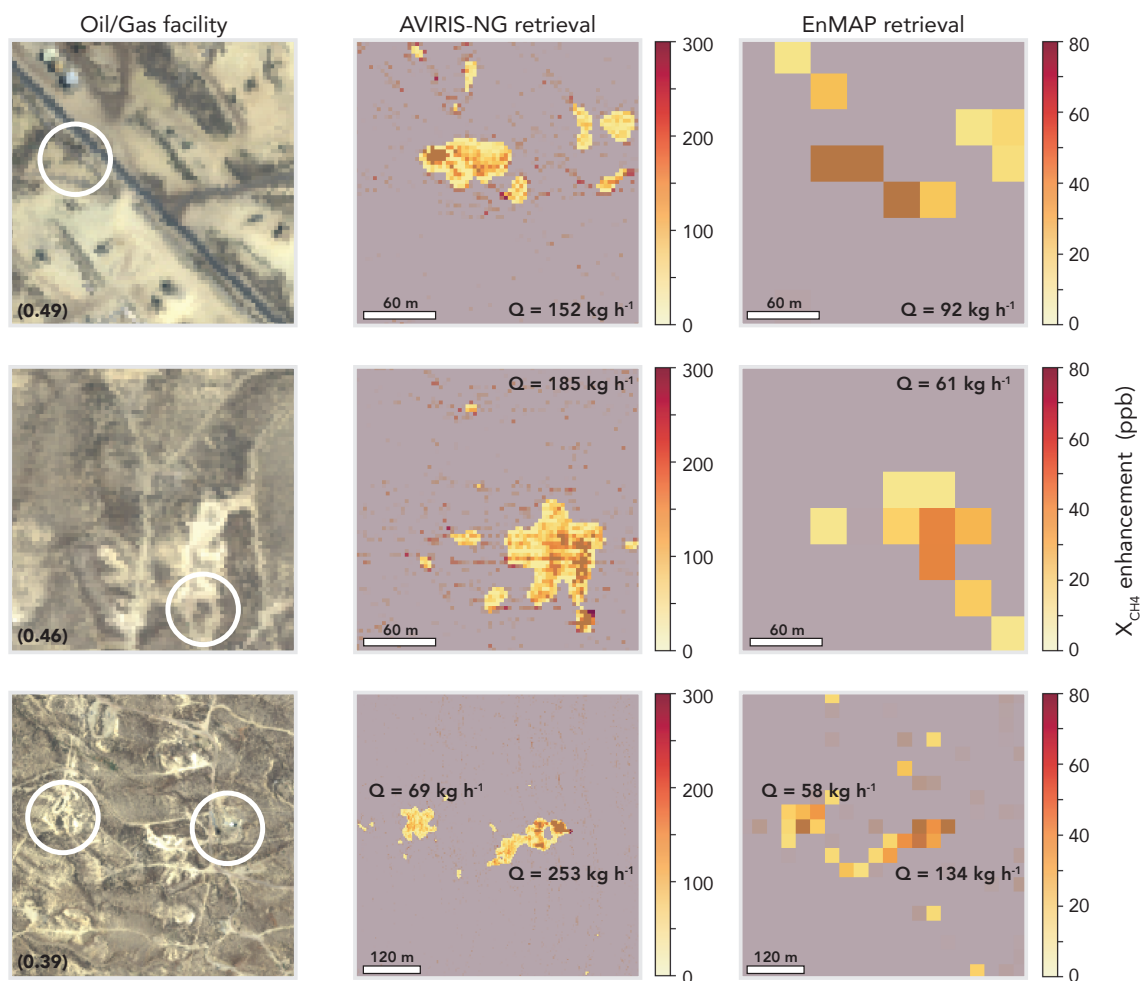
**Figure 5.** Precision of methane retrievals for spaceborne imaging spectrometers observing in the SWIR (2210-2400 nm), as a function of instrument signal-to-noise (SNR) and full-width half-maximum (FWHM) spectral resolution. The left panel shows precision expressed as the relative root-mean-square error (RRMSE) for synthetic retrievals over three scenes of Figure 2 (Grass, Urban, Bright) including a point source of  $100\text{--}900\text{ kg h}^{-1}$  and 15 different WRF-LES plume realizations. The SNR in the left panel represents the mean SNR over all three EteS scenes. The right panel shows theoretical precision expressed from the posterior error covariance matrix in Equation 7. Black dots show different instrument specifications from Table 1. Specifications for the SBG and AMPS instruments are still at the design stage and values shown here are for the ranges under consideration. Results given for AVIRIS-NG are for a satellite instrument with  $30\times 30\text{ m}^2$  pixel resolution, with (1) or without (2) along-track oversampling, and with other specifications (spectral resolution, SNR) the same as the airborne instrument.



**Figure 6.** Plume pattern recognition applied to a point source of  $500 \text{ kg h}^{-1}$  over Grass and Urban scenes as shown in Figure 3. The plume pattern is defined by applying median and Gaussian filters to pixels above the 80<sup>th</sup> percentile of  $X_{\text{CH}_4}$  in the scene. Areas excluded by the mask are shown in gray. The panels show retrievals from the EnMAP, SBG, and AMPS instruments.



### Methane retrievals over oil/gas facilities in California



**Figure 7.** Retrieval of atmospheric methane plumes from oil/gas facilities imaged by the AVIRIS-NG instrument at 3–4 km altitude over California (CARB, 2017). The left panels show the RGB images mapped by AVIRIS-NG with the oil/gas facilities of interest circled. Inset in the bottom left corner is the mean retrieved SWIR surface reflectivity for the scene. The middle panels show the IMAP-DOAS retrieval applied to the AVIRIS-NG images with  $3 \times 3 \text{ m}^2$  pixel resolution and 5 nm spectral resolution. The right panels show the IMAP-DOAS retrieval applied to spectra that were spatially and spectrally



downsampled to match EnMAP instrument specifications ( $30 \times 30$  m<sup>2</sup> pixels, 10 nm spectral resolution). Note the difference in color scale for the methane enhancements in the AVIRIS-NG and EnMAP retrievals, reflecting the coarser pixel resolution of EnMAP. The plume mask is overlaid on each. The source rates for each plume obtained from the IME method are inset.