Dear Dr. Alexander Myagkov,

problem:

Thank you for your new feedback. It was very useful to improve the article quality and clarity, improving the reliability, and thus the value, of our results. As before, we explain changes done to the manuscript here and attach a manuscript document with change-tracking.

Our previous responses are in blue, your new comments are in brown and our new responses are in black.

1.The main objective of the paper is to present a calibration methodology. The methodology itself is not affected by the use of a real or theoretical target RCS. Actually, once the real target RCS is retrieved, any possible bias in the results can be corrected without changing the calibration method. We now state this more clearly in lines 190-196. The company that manufactures the targets declares having a cutting accuracy better than 0.1 mm and an alignment precision better than 0.1°, therefore we can expect a bias but it should be on the order of 1-2 dBsm. We also include now how to account for the uncertainty of an eventual target characterization (eq. 6a and lines 231-234), and indicate that the uncertainty of the target calibration may increase the uncertainty in the results (lines 530-535). Finally, as future work we now include the need of a target characterization in an anechoic chamber to correct any bias introduced by the use of the theoretical model (lines 598-600).

Citation: "However, since at the writing time we do not have an experimental characterization for our targets, we rely on the: theoretical model. This is not a major issue because, once an experimental characterization of the target becomes available, it can be used to correct any calibration bias by rectifying the value of Gamma used in the calculations"

I do not agree with the authors. What is described in the manuscript is a method (i.e. description of steps to get knowledge), not a methodology (analysis of a set of methods). And in my opinion a calibration method is worth nothing without a proper characterization of a calibration target. I think this is the first thing one should do for the radar calibration – characterize the reference target. Currently it sounds to me, that after the proposed calibration procedure another calibration steps would be required (characterization of the target and application of another bias correction) when the target is measured in a chamber. The authors claim, "A detailed analysis enabled the design of a calibration methodology which can reach a cloud radar calibration uncertainty of 0.3 dB based on the equipment used in the experiment". This can be misleading for a reader. The authors do not reach the claimed value (0.3 dB) in the current work. As authors estimate, the real

- Authors characterize the target in a chamber and add these results (cross section and its uncertainties) in the manuscript.

uncertainty is not known at the moment and may be in the order of 2 dB (dBsm are not proper units here, since this value is unitless in linear scale). I suggest two ways to solve this

- Authors use sigma_rcs = 2 dB in Eq. 6a, reevaluate the results, and write explicitly in the abstract, main text, and conclusions that the uncertainty of the proposed method at the current stage is not better than ... dB due to uncharacterized reference target. Otherwise, it

is not honest to neglect a large uncertainty source just because it is not characterized. I would strongly recommend the authors to follow one of these ways.

We agree about the need of characterizing the target to provide final uncertainty results. However, we do not agree that the lack of the target characterization cancels the validity of the results, since we used a theoretical model of target RCS that has all the properties that a calibrated target would have, except the absolute values.

The results presented in the article enable the identification of several uncertainty sources, as well as their relative contribution to the experiment uncertainty. This information can be very valuable to design future calibration experiments based on reference reflectors, whether they are mounted on masts or held by other means, such as UAVs. The underlying principles remain the same.

We also add that our objective is not to claim we have a reference instrument, but to present all the information and advancements obtained from our experimental campaigns, specially in uncertainty characterization. For example, with the results we can quantitatively compare two different experimental setups, finding different factors limiting uncertainty for each (SCR for the 10 m mast, alignment for the 20 m mast). Additionally, we are not aware of any other published methodology of radar calibration that considers the bias introduced due to misalignment between target and radar. For us this work is a step towards more precise calibration methodologies, and we expect it to act as a reference to improve the preparation of future calibration experiments.

Because it is true that at this stage we can only do a rough estimation of RCS uncertainty, we agree to highlight this explicitly. This is now stated/included in:

Line 14 of the Abstract.

Lines 197-200

Line 229

Lines 507-509

Calibration result for all experiments (lines 512-531)

Table 2

Lines 585-587

This also implied modifications to some text in the article to remain consistent:

Lines 18-20 of the abstract.

Lines 54-55 of the Introduction

Lines 153-154

We also added an estimation of the maximum uncertainty in RCS characterization required to reach a calibration uncertainty of 0.5 dB in lines 592-597.

2. The problem with the power units arises because power output in the BASTA radar is in an

arbitrary power unit. We define this power unit as dB(AU) = 10log 10 (AU). The arbitrary unit

defined as AU is proportional to watts multiplied by a unitless digital gain k d , which depends on the digital signal processing configuration of the radar, such that dB(AU) = dBW +10log 10 (k d). Since the absolute calibration method will provide a calibration result that compensates this constant term, we did not work in transforming the power to standard physical units. We now explained this detail in lines 72-76 . For consistency, now every power unit is defined in dB(AU) units, and therefore the RCS calibration is now in dB(AU -1 m -2) and the reflectivity calibration is in dB(mm -6 m -5 AU -1). This way, when the term is multiplied by reflected power and distance to the corresponding power, the result will be in the correct units (dBsm or dBZ). All RCS values presented in the manuscript are now in dBsm units, both in text and figures. Line 84 also indicates that dBsm units are decibels referenced to a square meter. We also fixed a typo in Fig. 9 (prev fig 6). The maximum RCS indicated before in the label was of 28.28 dBsm, but it is actually 28.34 dBsm.

The introduced changes are even more confusing. The calibration terms characterize a ratio of a real measure over the calculated one. Therefore, calibration terms must be unitless in linear scale and in dB in the logarithmic scale. I do not understand what a calibration term in dB(AU^-1m^-2) means. The lines 72 – 76 are confusing. It is stated that kd is included to account for the units of the measured power which is in 10*log10(AU). One sentence later it is stated that kd is unitless. If it is unitless then the equation 1a has problems with units again. The nominator is unitless, the denominator has units of m^2*W as it was in the original version. I kindly ask the authors to carefully reconsider the units again. In fact, the previous version was better, the only problem was with units notation, i.e. dB was used instead

of dBm and dBsm (please see my previous comments). Please modify the units in such a way that the calibration factors are given in dB (unitless in linear scale). And please modify the units throughout the manuscript accordingly.

Following this suggestion, we now present the received power units to dBm and the calibration terms to dB. This led to the recalculation of the calibration terms absolute value.

Modifications:

- Because of this improvement, lines 79-80 and 82-84 explaining these arbitrary power units are no longer necessary and were removed.
- Eq. (2a) and line 542 were modified to remove the unnecessary digital gain term kd.
- Writing of the units corrected in lines 97, 112, 127, 142 and the Glossary.
- Absolute value of power/calibration terms corrected in lines 261, 264, 265, 267, 319-321, Table 1, and lines 512-531 of the calibration results.
- Additionally, Figures 3, 4(a), 5, 6 and 10 were modified to remain consistent with the power units.
- **4.** This change is included in every mention of receiver losses as L r (T, r), and is therefore propagated to the RCS and reflectivity calibration terms as well, which now depend on temperature and range (C Γ (T, r), C Z (T, r)).

I would recommend to use IF instead of r because for a different chirp configuration (slope) the relations between IF bins and range gates may change.

We originally used range because BASTA-Mini has only 4 standard operational modes, so it was not a very important distinction. Nevertheless, we agree that this formulation would improve the generality of the method, and consequently we modified all range dependent terms by terms depending on the IF frequency F b.

Besides, we added a short explanation of the F_b term in lines 75-81 and the new Equation (1), which indicates how F_b it is associated with the range r.

5. Section 5.5

In this newly added section the authors, as far as I understand, assume that during the 'passive' observations the power variability along IF depends only on gain changes. In general case this is not true:

 $Pr(IF) \sim G(IF)^*(Tsys(IF) + Tamb)$

Here Pr(IF) is the received power at IF in W, G(IF) is the linear gain of the receiver chain at IF (unitless), Tsys(IF) is system noise temperature at IF in K, Tamb is brightness temperature of the sky (or an object the radar was pointed to) in K, \sim is the proportionality sign. From this equation one can see that the received power depends on two parameters, namely the gain and the system noise temperature. If I understand right, the authors did so-called single point calibration. Using the single point calibration is it not possible to separate the gain and the system noise temperature. Therefore, typically two-point calibrations are used in radars and radiometers. Also the authors need to know Tamb (at least with respect to Tsys(IF)). I kindly ask the authors to clarify how they took these aspects into account to calibrate all the IF bins of the radar receiver.

In this section we do not intend to estimate the absolute value of gain at the IF, but rather to quantify relative gain changes with respect to the calibrated IF frequency (associated with the target position). To do this using passive observations, we had to make the assumption of a constant noise power, both from the system and from environment in the 12 MHz bandwidth of the receiver. This assumption is reasonable because components used in the receiver have a much larger bandwidth (for example for the Low Noise Amplifier (LNA) it is of 35 GHz).

We indicated briefly this assumption in the previous version, but we didn't mention the impact it could have on uncertainty. This is now estimated from the LNA specifications. Its variability of gain and noise figure in the 12 MHz bandwidth used is smaller than 0.1 dB. Since LNA are typically the main source of system noise in the receiver, we consider that 0.1 dB is a safe estimation of the uncertainty introduced by assuming a constant system noise in the IF bandwidth. This term dominates the RMSE between the fit and data, and the inter-period variability, thus we now define the IF correction function uncertainty to be of 0.1 dB.

Despite this, we agree that two point calibrations are highly desirable because they enable the retrieval of receiver absolute gain and system noise. This is now indicated in the text.

Thus, article changes are:

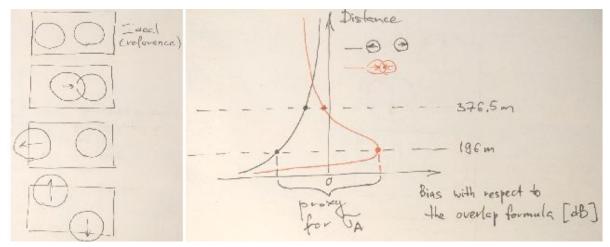
- More accurate explanation of passive observations (lines 365-368, equation 9a)
- Brief indication of the benefits of performing two point calibrations for receivers (lines 370-373)
- Explanation of the constant system noise assumption and introduced uncertainty (lines 374-383, equation 9b)
- Clarification in the explanation on how the IF correction function is retrieved (lines 384-396, line 406-408).
- Final uncertainty of the IF correction function (lines 413-414)

6. However, we did another revision of the scanning data and concluded that, at present, it is not possible to retrieve alignment information with an accuracy comparable to the antenna beam-width. This is now stated in lines 294-295. The reason is that the repeatability of the scanner positioning is not sufficient to allow a reliable retrieval under our current procedure. Additionally, we now include a discussion on how parallax errors can influence the measurements (286-290), and indicate that calibration results are compatible with parallax errors smaller than the radar beamwidth (296-298). Since we don't have information on the exact alignment, we now mention the parallel antennas only as an hypothesis (245, 299-300). Finally, we improved the calibration methodology by indicating how parallax errors can be taken into account, suggesting the addition of an additional range dependent correction function (300-301), and by introducing an uncertainty term representing the error in the antennas alignment estimation (eq. 6b, lines 243-245 and 301-302).

The assumption on parallel antennas can lead to large uncertainties. The problem with two antennas is that it is possible to measure the pattern of the receiving antenna with an external transmitter but it is often not possible to measure the transmitting antenna. Basically, with the proposed method only two points of the possible range dependent bias are characterized.

Instead of leaving this large uncertainty source untouched, I would encourage the authors to make a relatively simple estimation of possible impacts (just theoretical calculations, taking into account different divergences (magnitude and direction) of the two antennas and bias measurements at two distances). This would definitely improve the quality of the manuscript. The result of this theoretical estimation would give a proxy for sigma_a in Eq. 6a which is currently, if I understand it right, completely neglected.

Just to better understanding I give a couple of figures:



On the left figure you can see different divergence directions. On the right figure I illustrate the impact (qualitatively). The authors could perform such calculations and give an estimate for sigma a (maximum divergence from 0 dB line).

Thank you for this proposal, we believe it is a very good idea with good potential. Because of this, we performed several theoretical calculations to check if we could estimate a range of possible antenna misalignment angles and the associated uncertainty with our data. Summarizing, our results show that the experimental setup used is not appropriate for this measurement, but they also indicate us a path to perform such experiments in the future.

Since the targets used at 196 and 376.5 are different, the uncertainty in the **calibration coefficient difference** at both distances is very large (~3 dB). This uncertainty makes it impossible to bound the possible alignment within 1.5 degrees, which is our antenna characterization width. This large uncertainty comes mostly from the use of two different calibration targets. This decision was made because the experiment was designed to applicability of the absolute calibration method for different experimental setups, and because the proposed experiment was not considered at the time.

Given that the proposed experiment was not done during the campaigns presented in the article, we have no way to gather any additional information on antenna alignment for that period. Thus, we leave this section unchanged with respect to the previous version.

Yet, with the theoretical calculations we found that if we get an uncertainty in the order of 0.5 dB when comparing calibration constants at these two different distances, antenna misalignment could be constrained to values ranging in the order of tenths of degree. This could be achieved, for example, by using the 20 meter mast setup at both distances using the same reflector each time. It is worth noting that the tools developed for this analysis now enable the design of an experiment with optimized parameters for this retrieval.

Taking all this into consideration, in our opinion the potential of the proposed experiment indicate that it must be further studied for its implementation in future calibration campaigns.

7. To verify if data did follow a linear relationship, we did a new plot with the point density of all samples together. This figure has been added to the paper (Figure 7). In this figure it is easier to observe that deviated points are rather exceptional, with most points close to the regression. From this figure we think the 0.13 dB RMSE value is representative for most samples. We also modified Figure 6 (D). Now it is only used to introduce the data set, with the linear fit shown in new Figure 7. This produced text changes in lines 351-355, and 360-371.

In Fig. 7 the authors just masked the problem I am talking about. I agree that a majority of samples follow the linear model. But some complete iterations (like green points in Fig 6d) are off by more than 0.5 dB.

This happens because, to capture the widest range of possible temperatures, we had to use longer time series of data. The use of longer time series introduced some points measured under suboptimal conditions. For example, we have observed that high wind speeds lead to larger variabilities in the calibration value due to oscillations of both the mast and the radar. Meanwhile, drizzle adds a time dependent bias, most likely caused by changes in wet radome (and wet target) attenuation over time. An effort was done to clean the dataset, but inevitably some noisy data points remained. This is now stated more clearly in lines 345 to 348.

Therefore, to estimate the temperature correction function and its uncertainty we did an statistical analysis of data, and then used the RMSE between the model and data points as the estimation of model uncertainty.

To bound the uncertainty value we calculated the RMSE of the model for each degree of deviation from the reference temperature, obtaining a RMSE range between 0.07 to 0.23 for 0 and +3 degrees of deviation respectively. We also checked the bias per degree for each iteration and for all the dataset, and found out that its mean value is always within +- 0.2 dB with respect to the model. Therefore, we now state that the temperature correction function uncertainty is less or equal to 0.23. This change is reflected in lines 353-357.

This is also true for the mentioned case, were most points are covered by the other iterations data. The larger spread reaching deviations of 0.5 dB in this case are caused by short period of drizzle that happened in this iteration. However, since this data, and the rest of data that deviates from the model is also included in the calculation of RMSE, we think this is a reliable criteria for the estimation of the temperature correction function uncertainty.

Absolute Calibration method for FMCW Cloud Radars

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Abstract.

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This article presents a new Cloud Radar calibration methodology using solid reference reflectors mounted on masts, developed during two field experiments held in 2018 and 2019 at the SIRTA atmospheric observatory, located in Palaiseau, France, in the framework of the ACTRIS-2 research and innovation program.

The experimental setup includes 10 cm and 20 cm triangular trihedral targets installed at the top of 10 m and 20 m masts, respectively. The 10 cm target is mounted on a pan-tilt motor at the top of the 10 m mast to precisely align its boresight with the radar beam. Sources of calibration bias and uncertainty are identified and quantified. Specifically, this work assesses the impact of receiver compression, incomplete antenna overlap, temperature variations inside the radar, frequency dependent losses in the receiver IF, clutter and experimental setup misalignment. Setup misalignment is a source of bias previously undocumented in the literature, that can have an impact on the order of tenths of dB in calibration retrievals of W band Radars.

A detailed analysis enabled the design of a calibration methodology which can reach a quantification of the importance of each uncertainty source to the final cloud radar calibration uncertainty of 0.3 dB based on the equipment used in the experiment. Among different sources of uncertainty, the two largest terms are due to signal-to-clutter ratio and radar-to-target alignment. The . The dominant uncertainty source comes from the uncharacterized reference target, reaching 2 dB. Additionally, the analysis revealed that our 20 m mast setup with an approximate alignment approach is preferred to the 10 m mast setup with the motor-driven alignment system. The calibration uncertainty associated with signal-to-clutter ratio of the former is ten times smaller than for the latter.

Following the proposed methodology it is possible to reduce the added contribution from all uncertainty terms, excluding the target characterization, down to 0.4 dB. Therefore, this procedure should enable to achieve calibration uncertainties under 1 dB when characterized reflectors are available.

Cloud radar calibration results are found to be repeatable when comparing results from a total of 18 independent tests. Once calibrated the cloud radar provides valid reflectivity values when sampling mid-tropospheric clouds. Thus we conclude that the method is repeatable and robust, and that the uncertainties are precisely characterized. The method can be implemented under different configurations as long as the proposed principles are respected. It could be extended to reference reflectors held by other lifting devices such as tethered balloons or unmanned aerial vehicles.

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1 Introduction

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Clouds remain to this day one of the major sources of uncertainty in future climate predictions (Boucher et al., 2013; Myhre et al., 2013; Mülmenstädt and Feingold, 2018). This arises partly from the wide range of scales involved in cloud systems, where a knowledge of cloud micro-physics, particularly cloud-aerosol interaction, is critical to predict large scale phenomena such as cloud radiative forcing or precipitation.

To address this and other related issues, the ACTRIS Aerosols, Cloud and Trace Gases Research Infrastructure is establishing an state of the art ground based observation network (Pappalardo, 2018). Within this organization, the Centre for Cloud Remote Sensing CCRES is in charge of creating and defining calibration and quality assurance protocols for the observation of Cloud properties across the complete network.

One of the key instruments for cloud remote sensing stations is the Cloud Radar. Cloud radars enable retrievals of several relevant parameters for cloud research, including but not limited to liquid water and ice content profiles, cloud boundaries, cloud fraction, precipitation rate and turbulence (Fox and Illingworth, 1997; Hogan et al., 2001; Wærsted et al., 2017; Dupont et al., 2018; Haynes et al., 2009). Additionally, recent studies revealed the potential of cloud radars to support a better understanding of fog processes (Dupont et al., 2012; Boers et al., 2013; Wærsted et al., 2019).

However, calibration remains a crucial factor in the reliability of radar retrieved data (Ewald et al., 2019). Systematic differences of 2 dB have already been observed, for example, between the satellite based radar CloudSat and the Lindenerg MIRA (Protat et al., 2009). This is a very important issue, since calibration errors as small as 1 dB would already introduce uncertainties in liquid water and ice content retrievals in the order of 15-20% (Fox and Illingworth, 1997; Ewald et al., 2019).

Since the objective of the CCRES is to guarantee a network of high quality observations, it is essential to develop standardized and repeatable calibration methods for its instrumental network.

This paper presents an absolute calibration method for W band radars. It has been developed based on results from two experimental calibration campaigns performed at the SIRTA Atmospheric Observatory, located in Palaiseau, France (Haeffelin et al., 2005). The SIRTA observatory hosts part of the ACTRIS CCRES infrastructure. For the experiments we used a BASTA-Mini W band Frequency Modulated Continuous Wave (FMCW) Radar, with scanning capabilities (Delanoë et al., 2016). Nevertheless, the principles, procedures and limitations presented here should be applicable for any radar with similar characteristics, even when operating in another frequency band.

The method consists on an end-to-end calibration approach, consisting in retrieving the radar calibration coefficient by sampling the power reflected from a reference reflector mounted on top of a mast (Chandrasekar et al., 2015). A detailed analysis of uncertainty and bias sources is performed, with the objective of reducing uncertainty under determining how to improve the experiment to reach a calibration uncertainty of $0.5 \, dB$. This low uncertainty in the calibration would not only be useful for high quality retrievals, but also enables the use of the radar as a reliable reference for calibration transfer to other ground or space based cloud radars (Bergada et al., 2001; Protat et al., 2011; Ewald et al., 2019).

The article is structured as follows: Section 2 present the ecuations and theoretical considerations involved in the calibration exercise. Section 3 shows the experimental setup, complemented by section 4 where the experimental procedure and data treat-

60 ment is presented. Section 5 presents an analysis of the sources of uncertainty and bias involved in our calibration experiment. Section 6 presents the final calibration results, the uncertainty budget and an analysis of the variability in the calibration bias correction, followed by the conclusions.

2 Equations used in Radar Calibration

The absolute calibration of a radar consists in determining the ealibration terms RCS Calibration Term C_{Γ} and the Radar Equivalent Reflectivity Calibration Term C_{Z} . They enable the calculation of Radar Cross Section $\Gamma(r)$ (RCS) or Radar Equivalent Reflectivity Z_{e} respectively, from the power backscattered by a punctual or distributed target towards the radar (Bringi and Chandrasekar, 2001).

Equation (2a) presents an expression for the RCS calibration term $C_{\Gamma}(T, F_b)$ of a FMCW radar as a function of its internal parameters. The deduction of this expression is shown in the supplementary material. G_t and G_r are the maximum gain of the transmitting and receiving antennas respectively, dimensionless. λ is the wavelength of the carrier wave in meters and p_t is the power emitted by the radar in watts.

The gain of solid state components changes with variations in their temperature. Thus we make this dependence explicit in the receiver loss budget $L_r(T,r)$ $L_r(T,F_b)$ and in the transmitter loss budget $L_t(T)$. The loss budget is the product of all losses divided by the gain terms at the end of the receiver or emitter chain, and has no dimensions.

Additionally, a range dependence is included in $L_r(T,r)L_r(T,F_b)$ to account for variations in the receiver IF loss for different beat frequency F_b values. The beat frequency is used in FMCW radars to calculate is proportional to the distance between the radar and reflective targets (Delanoë et al., 2016). Changes instrument and the backscatterer element (Delanoë et al., 2016). Thus, changes in the IF loss for different beat frequencies introduce a range dependent loss.

The term k_d is included to account for the units of radar measured power. Radar measured power is in arbitrary power units defined as $dB(AU) = 10\log_{10}(AU)$ bias. For the 12.5 meter resolution mode used in this calibration excercise, F_b ranges between 168 and 180 MHz, and can be related to r (in meters) using Eq. (1).

These units are related to physical power units by the equation $dB(AU) = dBW + 10\log_{10}(k_d)$, where dBW is the physical power expressed in decibels relative to 1 watt and k_d is a unitless constant determined by the digital signal processing configuration.

$$5 \quad r = 500 \cdot (F_b - 168[\underbrace{MHz}]) \tag{1}$$

In theory, $C_{\Gamma}(T,r) C_{\Gamma}(T,F_b)$ can be calculated by characterizing the gains and losses of every component inside the radar system and adding them. This can be very challenging, depending on the complexity of the radar hardware and the available radio frequency analysis equipment. In addition, with this procedure it is not possible to quantify losses due to interactions between different components, specially changes in antenna alignment or radome degradation (Anagnostou et al., 2001). This

motivates the implementation of an end-to-end calibration, which consists on the characterization of the complete radar system at once by using a reference reflector and Eq. (2b).

$$C_{\Gamma}(T,\underline{r}F_b) = 10\log_{10}\left(\frac{L_t(T)L_r(T,r)(4\pi)^3}{k_dG_tG_r\lambda^2p_t}\frac{L_t(T)L_r(T,F_b)(4\pi)^3}{G_tG_r\lambda^2p_t}\right)$$
(2a)

$$\Gamma(r) = C_{\Gamma}(T, \underline{r}F_b) + 2L_{at}(r) + 40\log_{10}(r) + P_r(r)$$
(2b)

Equation (2b) links the calibration term $C_{\Gamma}(T,r)$ $C_{\Gamma}(T,F_b)$ to the RCS $\Gamma(r)$ of a target at a distance r. $\Gamma(r)$ is expressed in dBsm units (decibels referenced to a square meter), $L_{at}(r)$ is the atmospheric attenuation between the object and the radar in dB, which can be calculated using a millimeter-wave attenuation model (for ex. (Liebe, 1989)), $P_r(r)$ is the power received from the target in dB(AU) and $C_{\Gamma}(T,r)$ dBm and $C_{\Gamma}(T,F_b)$ is the RCS calibration term in $dB(m^{-2}AU^{-1}) = 10\log_{10}(m^{-2}AU^{-1})$ units dB. The units in the RCS calibration term compensate the radar power units, guaranteeing the retrieval of physical RCS values. The explicit temperature and range dependency of the calibration term has the function of compensating gain changes in $P_r(r)$ introduced by temperature effects and variations in the IF loss with distance.

This principle can be used in an end-to-end calibration by installing a target with a known RCS Γ_0 at a known distance r_0 and sampling the power $P_r(r_0)$ reflected back to calculate $C_{\Gamma}(T,r)C_{\Gamma}(T,F_b)$. However, some additional considerations must be made to perform this retrieval.

In Eq. (2a) we state that the calibration value has a temperature and a range dependency. Experimental results indicate that the temperature dependency of $C_{\Gamma}(T,r)$ $C_{\Gamma}(T,F_b)$ can be approximated by a linear relationship, as shown in Eq. (3). Here n is the temperature dependency term in $dB \, {}^{\circ}C^{-1}$, T the internal radar temperature in ${}^{\circ}C$ and T_0 a reference temperature value in ${}^{\circ}C$. More details about the temperature correction can be found in Sect. (5.4).

The range dependency of $C_{\Gamma}(T,r)$ dependence of $C_{\Gamma}(T,F_b)$ is treated independently, by defining a IF loss correction function $f_{IF}(F_b)$, in dB units. This function is introduced to compensate for relative loss variations at different distances of frequencies. The IF loss correction function is studied in Sect. 5.5.

From the afforementioned observations, we divide $C_{\Gamma}(T,F_b)$ in three components, shown in Eq. (3). This separation consists of a constant calibration coefficient C_{Γ}^0 , in $dB(m^{-2}AU^{-1})dB$, and the two correction functions $n(T-T_0)$ and $f_{IF}(F_b)$.

$$C_{\Gamma}(T, rF_b) = C_{\Gamma}^0 + n(T - T_0) + f_{IF}(rF_b)$$
(3)

As $f_{IF}(r) f_{IF}(F_b)$ corrects for relative variations in receiver loss with distance, we define $f_{IF}(r_0) = 0$ at the reference $f_{IF}(F_0) = 0$ at the IF frequency value F_0 , associated to the reflector position r_0 (linked by Eq. (1)). Using this and Eqs. (2b) and (3), we get Eqs. (4a) and (4b).

$$C_{\Gamma}(T, \mathbf{r}F_0) = C_{\Gamma}^0 + n(T - T_0)$$
 (4a)

$$C_{\Gamma}(T, \mathbf{r}_{\Gamma_0}) = \Gamma_0 - 40\log(r_0) - 2L_{at}(r_0) - P_r(r_0) \tag{4b}$$

Equation (4a) shows how the calibration term $C_{\Gamma}(T,r_0)$ $C_{\Gamma}(T,F_0)$ at position r_0 is related to the calibration coefficient C_{Γ}^0 and the temperature correction $n(T-T_0)$. Meanwhile, Eq. (4b) indicates how experimental $P_r(r_0)$ measurements can be associated with a $C_{\Gamma}(T,r_0)$ $C_{\Gamma}(T,F_0)$ value, using in-situ information to calculate $2L_{at}(r_0)$. Then, using Eq. (4a), we can compute C_{Γ}^0 by substracting the temperature correction function $n(T-T_0)$. This temperature correction is derived independently in Sect. (5.4). Knowing C_{Γ}^0 and the temperature correction, $C_{\Gamma}(T,r)$ $C_{\Gamma}(T,F_0)$ is calculated by adding the IF correction function, independently retrieved in Sect. 5.5.

Once $C_{\Gamma}(T,r)$ $C_{\Gamma}(T,F_b)$ is known, we can calculate the radar Equivalent Reflectivity calibration term $C_{Z}(T,r)$, in $dB(mm^6m^{-5}AU^{-1}\sin dB)$, with Eq. (5a) (Yau and Rogers, 1996). This relationship assumes the radar has two identical parallel antennas with a Gaussianly shaped main lobe. θ is the antenna beamwidth in radians, $m\delta r$ is the radar distance resolution in meters and $|K| = |(\epsilon_r - 1)/(\epsilon_r + 2)|$ is the dielectric factor. This factor is related to the relative complex permittivity ϵ_r of the scattering particles, and can be calculated, for example, using the results of Meissner and Wentz (2004).

 $C_Z(T,r)$ $C_Z(T,F_b)$ enables the calculation of the Radar Equivalent Reflectivity Z_e in dBZ units of a distributed target located at a distance r, by using the relationship of Eq. (5b).

$$C_Z(T, \underline{\underline{r}F_b}) = 10\log_{10}\left(\frac{8\ln(2)\lambda^4 10^{18}}{\theta^2 \pi^6 K^2 \delta r}\right) + C_{\Gamma}(T, r)$$
(5a)

$$Z_e(r) = C_Z(T, rF_b) + 2L_{at}(r) + 20\log_{10}(r) + P_r(r)$$
(5b)

135 3 Experimental setup

Two calibration campaigns, that lasted one month each, were performed in May-June of 2018 and March-April of 2019 at the SIRTA observatory, located in Palaiseau, France (Haeffelin et al., 2005). The observatory has a 500 meter long grass field in an area free of buildings, trees or other sources of clutter, well suited to install our calibration setup, shown in Fig. 1.

The instrument used for the calibration experiments is a BASTA-Mini. BASTA-Mini is a 95 GHz FMCW radar with scanning capabilities and two parallel Cassegrain antennas (Delanoë et al., 2016). The antennas are separated by 35 cm, and have a Fraunhofer far field distance of $\approx 50~m$ with a Gaussianly shaped main lobe (verified experimentally in Sect. 5.2). Transmitted power is fixed to 0.5~W 500 mW, and is under constant monitoring using a diode with an uncertainty of $\approx 0.4~dB$. The diode enable the monitoring of $L_t(T)$ variations, yet our experiments have shown that T is a better indicator to capture

the variability of $C_{\Gamma}(T,r)C_{\Gamma}(T,F_b)$. This is likely because internal temperature changes affect both $L_r(T,r)L_r(T,F_b)$ and $L_t(T)$ simultaneously, and therefore the information provided by the diode is not sufficient to capture the behavior of the whole system. The results of the temperature dependency study for our radar is shown in Sect. 5.4.

This radar also includes hardware to enable the tuning of the carrier wave frequency within a range of $\approx 1~GHz$, centered at 95~GHz. During the experiments we fixed the BASTA-Mini base frequency at 95.64~GHz to avoid any interference with the other two W band radars operating in parallel at the same site.

Our reference targets are two Triangular Trihedral Reflectors (also known as Corner Reflectors) composed by three orthogonal triangular conducting plates. Trihedral targets have a large RCS for their size and a low angular variability of RCS around their boresight (Atlas, 2002; Doerry and Brock, 2009; Chandrasekar et al., 2015). One reflector has a size parameter of 10 cm, with a maximum theoretical RCS at our radar operation frequency of 16.30 dBsm. The other is 20 cm with a maximum theoretical RCS of 28.34 dBsm (Brooker, 2006). These targets were mounted on top of masts B and C in Fig. 1 respectively.

Only mast C was used in the 2018 campaign, while both were used in 2019.

To align the system first we aim the radar towards the approximate position of the target. Second, we aim the target by slowly changing pan-tilt angles in the motor on mast B, or axially rotating the tube of mast C to maximize the power $P_r(r_0)$ measured at the radar. Third, radar aiming is tuned around target position until the maximum reflected power is found. Finally, we repeat the second step, after which we have the system ready to sample $P_r(r_0)$.

It must be mentioned that this procedure does not guarantee a perfect alignment. In fact, it is impossible to have every element perfectly adjusted because of limits in the positioner resolution or uncertainties introduced when installing each element. Sections 4 and 5.6 explain how we deal with these limitations.

4 Methodology

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This section describes the procedure followed when performing calibration experiments using the setup described in Sect. 3.

The methodology has the objective of quantifying and correcting when possible all sources of uncertainty to enable a reliable estimation of the calibration terms $C_{\Gamma}(T,r)$ and $C_{Z}(T,r)C_{\Gamma}(T,F_{b})$ and $C_{Z}(T,F_{b})$.

A challenge we found when using targets mounted on masts to estimate $C_{\Gamma}(T,r)$ $C_{\Gamma}(T,F_b)$ is that the value of the target RCS Γ_0 may vary depending on how components are aligned. Our studies have shown that for the feasible alignment accuracy we can get when installing our setup, this effect is in the order of tenths of dB, and therefore not negligible. Additionally, we concluded that if we leave this uncertainty source uncorrected, we would introduce a bias in the calibration result (see Sect. 5.6).

The flow chart of Fig. 2 illustrates the calibration procedure. To quantify the bias introduced by alignment uncertainty we decided to divide each calibration expreriment in N iterations. Each iteration consists on a system realignment, followed by sampling of the target signal $P_r(r_0)$ for at least one hour. Then, we select the data from the contiguous hour with the lowest variability as the iteration result.

The period chosen to perform the sampling is important, because it will have an incidence on how stable is the calibration value. To minimize uncertainty it is recommended to perform calibration iterations when the atmosphere is clear, there is no rain and wind speed is under $1\ ms^{-1}$. However, these requirements may change depending on how robust is each setup to atmospheric conditions.

FMCW radars have a discrete distance resolution. Consequently, power measurements vs distance are resolved in finite discrete points, usually named gates. Because of this resolution limitation, power received from a point target is spread between the gates closer to its position (Doviak and Zrnić, 2006). This phenomena is known as spectral leakage. To reduce leakage BASTA-Mini uses a Hann time window (Richardson, 1978; Delanoë et al., 2016).

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To correctly asses the total reflected power we set the radar resolution to 12.5 meters (chirp bandwidth of 12 MHz), and its integration time to 0.5 seconds. This resolution is high enough to accurately identify the reference reflector signal while avoiding the introduction of additional clutter from the trees located behind the mast (see Fig. (3)).

To calculate $P_r(r_0)$ we add five gates: the target's gate plus two before and two after the target's position. Adding more contiguous gates increase the power value by less than $0.01 \ dB$, thus we conclude that these five gates concentrate almost all the power reflected back from the target.

Then $P_r(r_0)$ is corrected considering compression effects and antenna overlap losses (Sects. 5.1 and 5.2). For each corrected $P_r(r_0)$ sample we proceed to calculate a single C_{Γ}^0 value with Eq. (4a) and the temperature correction function. This single sample is defined as $C_{\Gamma s}^0$ to differentiate it from the final calibration coefficient C_{Γ}^0 of Eq. (3). Atmospheric attenuation $L_{at}(r_0)$ is calculated using in-situ atmospheric observations and the model published by Liebe (1989).

The target effective RCS Γ_0 is calculated using a theoretical RCS model, considering the beam incidence angle on the targetbut neglecting possible alignment errors. It is worth noting that echo. Echo chamber measurements have shown that real targets RCS can be deviated from the theoretical value depending on the manufacturing precision. Our corner reflectors have an angular manufacturing precision better than 0.1° , therefore RCS deviations could be on the order of 1-real RCS uncertainty with respect to the model can be roughly estimated to be approximately 2 dBsm (Garthwaite et al., 2015). However, since at the writing time we do not have an experimental characterization for our targets, we rely on the theoretical model. This is not a major issue because, once dB (Garthwaite et al., 2015). Once an experimental characterization of the target becomes available, it can be used to correct any calibration bias and to reduce uncertainty by rectifying the value of Γ_0 used in the calculations.

We performed one calibration experiment with 6 iterations during the 2018 campaign using the 20 m mast. In the 2019 campaign we did two experiments: one with 10 iterations using the 10 m mast and another with 2 iterations on the 20 m mast (Fig. 1).

The retrieval of the temperature dependency coefficient n and the reference temperature T_0 is done simultaneously with the calibration coefficient experiment, by extending the sampling period beyond one hour when using the 20 m mast. This is done to capture the temperature effect in the variability of $C_{\Gamma s}^0$, by capturing a larger part of the temperature daily cycle. The results of this experiment can be seen in Sect. 5.4. Likewise, the retrieval of the IF correction function $f_{IF}(F_b)$ is an independent experiment based on sampling noise with the radar to get the IF amplification curve of the receiver. The details of this experiment are in Sect. 5.5.

From each iteration we get a distribution of resulting $C_{\Gamma s}^0$ values with a small spread introduced by second order effects. The average value of each iteration i is named $C_{\Gamma i}^0$, and its corresponding standard deviation is named σ_i . With this information we proceed to calculate the bias corrected calibration coefficient C_{Γ}^0 , by using Eq. (6). $\tilde{\Lambda}$ is the bias correction term. The method used to calculate $\tilde{\Lambda}$ relies on simulating the probability distribution of Γ_0 for a given set of uncertainties in the setup parameters. More detail can be found in Section 5.6 and Section S3 of the supplementary material.

$$C_{\Gamma}^{0} = \frac{1}{N} \sum_{i=1}^{N} C_{\Gamma i}^{0} - \tilde{\Lambda} \tag{6}$$

Equations (7a) and (7b) show the uncertainties δC_{Γ} and δC_{Z} associated with the estimation of $C_{\Gamma}(T, F_{b})$ and $C_{Z}(T, F_{b})$ respectively.

 σ_T is the uncertainty term associated with the temperature correction function $n(T-T_0)$.

 σ_{IF} is the uncertainty term associated with the IF loss correction function $f_{IF}(F_b)$.

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The term $\sum \sigma_i^2$ comes from the averaging operation in the estimation of $C_{\Gamma i}^0$ (Eq. 6). Since the $C_{\Gamma i}^0$ terms are corrected using the temperature correction function, the uncertainty of the later must be propagated as well, hence the term σ_T^2/N appears.

 σ_{Λ} the uncertainty of the bias correction calculation. It is calculated from the standard deviation σ_i . This procedure is explained in Section S3 of the supplementary material.

 σ_{SCR} is the uncertainty introduced by clutter. Clutter is the presence of unwanted echoes which affect our reading of $P_r(r_0)$, coming from reflections on other objects in the environment. The method to quantify the uncertainty σ_{SCR} uses a parameter named Signal to Clutter Ratio (SCR), explained in detail in Sect. 5.3.

 σ_{Γ_0} is the uncertainty of the reference target RCS. In this work we use a theoretical model to calculate the target effective RCS, which has an uncertainty of approximately 2 dB based on the manufacturing characteristics. The inclusion of an experimental characterization of the target RCS can improve the estimation of C_{Γ}^0 and δC_{Γ} . Because we do not have this information available at the writing date, we neglect this by reducing this uncertainty term.

 σ_K is the uncertainty in the estimation of the backscattering particles dielectric factor. Because our objective is to calculate the calibration term of the radar, we reference this value to |K| = 0.86, corresponding to pure water at 5 °C and neglect the δ_K uncertainty term. However, the value of K and its uncertainty σ_K must be considered when performing radar retrievals (e.g. Sassen (1987); Liebe et al. (1989); Gaussiat et al. (2003)).

 σ_A is the uncertainty introduced in the estimation of θ and from parallax errors and deviations from a Gaussian beam shape (Sekelsky and Clothiaux, 2002). For this work we make the assumption of parallel antennas with a Gaussian beam shape, thus we neglect this term. This problem is dissused discussed more in depth in Section 5.2.

Since both σ_K and σ_A are neglected, we get $\delta C_{\Gamma} \approx \delta C_Z$.

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$$\delta C_{\Gamma} = \sqrt{\frac{1}{N^2} \sum_{i=1}^{N} \sigma_i^2 + \frac{\sigma_T^2}{N} + \sigma_{IF}^2 + \sigma_T^2 + \sigma_{SCR}^2 + \sigma_{\Lambda}^2 + \sigma_{\Gamma_0}^2}$$
 (7a)

$$\delta C_Z = \sqrt{\delta C_\Gamma^2 + \sigma_K^2 + \sigma_A^2} \tag{7b}$$

5 Sources of uncertainty and bias in Absolute Calibration with corner reflectors

In this section we identify and quantify the uncertainty and bias introduced by several terms in Eq. (2b). Following the recommendations in the work of Chandrasekar et al. (2015), we study the impact of receiver saturation, signal to clutter ratio, antenna lobe shape and antenna overlap. Additionally, we consider the impact of temperature fluctuations inside the radar box, loss changes with distance due to uneven amplification at the receiver IF and the effects of imperfect alignment of the reference target.

5.1 Receiver compression

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It is advisable to design calibration experiments which avoid the appearence of compression effects. If this is not possible, compression must be considered in the data treatment so that the retrieved calibration remains valid in the receiver's linear regime, where it usually operates during cloud sampling (Scolnik, 2000).

For studying how these effects could affect our calibration, we retreived the radar's receiver power transfer curve. Receiver characterization was done by removing the radar's antennas and connecting the emitter's end to the receiver's input, with two attenuators in between. The first was a $40 \ dB$ fixed attenuator, while the second was a tunable attenuator covering the range between $50 \ \text{and} \ 1 \ dB$ of losses. The adjustable attenuator enabled the retrieval of the power transfer curve by varying the attenuation and sampling the power at the receiver's end (digital processing included). Our retrieved power transfer curve is shown in Fig. 4 (a).

Compression effects must be considered in calibration, or a bias will be introduced. In consequence, we include compression correction in every sample of reflected power, which consists on projecting their value to the ideal linear response using the power transfer curve.

For example, the power received from the 20 cm target on the 20 m mast returned was $\frac{198.7 \, dB}{4.1 \, dBm}$ in average, before corrections. The power transfer curve shows that at this power values we have a loss caused by compression of $\approx 0.3 \, dB$. After correcting each power sample by compression with the power transfer curve, we obtain a corrected power average value of $\frac{199.1 \, dB}{4.5 \, dBm}$. Meanwhile, for the 10 cm target on the 10 m mast the average power value before corrections is $\frac{197.83.2 \, dBm}{4.5 \, dBm}$. As this value is lower than what is obtained the 20 m mast, the associated compression effect is also smaller, of $\approx 0.2 \, dB$. After aplying this correction to each power sample we end with a new corrected power average of $\frac{198.0 \, dB}{4.5 \, dBm}$.

5.2 Antenna Properties

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Manufacturer specifications indicate that antenna beamwidth should be of 0.8° . However, data from an experimental characterization done by the same manufacturer in an anechoic chamber indicate that antenna beam shape is better approximated by a Gaussian with a Half Power Beam Width (HPBW) of $\theta \approx 0.88^{\circ}$. The total gain difference between the experimental curve and the Gaussian approximation of $\approx 0.0003~dB$ in the HPBW region. Therefore, we conclude that the contribution to uncertainty introduced by assuming a Gaussian beam shape is negligible. The Antenna beam shape and Gaussian curve are shown in Fig. 4 (b).

Another source of bias introduced by the antennas is the parallax error. Antenna parallax errors introduce a range dependent bias, determined by the antenna beamwidth and the relative angles of deviation between the antennas boresight. This bias is usually larger in the first few hundred meters closest to the radar. For example, for a deviation of half the antenna beamwidth, losses would be on the order of 10 dB and would vary significantly over the first hundreds of meters, decreasing with distance to about 1 dB at a approximately 4 kilometers (Sekelsky and Clothiaux, 2002).

To study this effect we took advantage of our experimental setup and the scanning capabilities of the radar, to check if the radar antennas were properly aligned. This was done by using the target on the 20 m mast. Results are shown in Fig. 4 (b). After analyzing the results we observed that the aiming uncertainty is in the same order of magnitude of the antennas beamwidth. Since the correction of the parallax error requires a very precise measurement of antenna alignment, we conclude that it is not possible to correct for antenna deviations directly with this information.

However, the relatively small difference of $0.5 \ dB$ in the estimation of C_{Γ}^0 during the calibration experiments of 2019, obtained using two masts in the most sensitive distance range (placed at 196 and 376.5 meters of distance respectively), indicate that antennas are unlikely to have a deviation comparable to their beamwidth (calibration results in Sect. 6).

Therefore, for the present version of this calibration methodology we assume that both antennas are parallel and that they have a Gaussian beam lobe. Once a reliable method for antenna pattern retrieval is developed for W band radars, it can be directly incorporated into the calibration term by adding an additional correction function $f_A(r)$ to Eq. (3). The uncertainty in this alignment estimation can also be included in the uncertainty budget, with the term σ_A of Eq. (7b).

Even if antennas are parallel, it is necessary to include a correction for the loss $L_o(r)$ caused by incomplete antenna overlap. The correction, shown in Eq. (8), accounts for the loss of power that would be received from a point target compared to a monostatic system (Sekelsky and Clothiaux, 2002). This loss occurs because a point target cannot be in the center of two non-concentric parallel antenna beams.

$$L_o(r) = \exp\left(\frac{2\arctan(\frac{d}{2r})^2}{0.3606\ \theta^2}\right) \tag{8}$$

Equation (8) assumes that the radar has two identical, parallel antennas with Gaussian beam lobes. Their main axis is separated by a distance d, and the point target is located at a distance r, facing the geometrical center of the radar, where the

gain is maximum. For the BASTA-Mini d = 35 cm. This introduces a loss of 0.08 dB for the target at $r_0 = 196$ meters of distance, and of 0.02 dB for the target at $r_0 = 376.5$ meters.

5.3 Signal to Clutter Ratio

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The power sampled from our reference reflector is an addition of the power from the target (signal) and unwanted reflections on other elements in the environment, such as the ground or the mast (clutter). We observed that this clutter dominates above the radar noise, and thus becomes the main source of interference in our calibration signal.

To quantify the impact of clutter we use the Signal to Clutter Ratio (SCR) parameter. It is calculated as the ratio of total power received from the target to power received from clutter under the same configuration, but with the reference reflector removed. SCR enables the uncertainty σ_{SCR} introduced by clutter in the sampled $P_r(r_0)$ values to be computed (Chandrasekar et al., 2015).

Clutter power is sampled and corrected following the same methodology used for reflector $P_r(r_0)$ retrievals, but in an scanning pattern mode to capture clutter around the mast area. Figure 5 shows our results from scanning around the 10 and 20 m masts with targets removed.

We observe that the 10 m mast is more reflective than the 20 m one. This may be caused by its smaller height (more ground clutter) and its larger geometrical cross-section. We can also see that the signal in the 10 m is stronger where absorbing material is not present (below $\approx 1.5^{\circ}$ of elevation). In both cases we did not detect any signal from the nearby trees close to the target's position.

To calculate SCR we compare the average power received from each target during the calibration experiments with the maximum clutter power observed in a region of 0.125° around the target's coordinates, vertically and horizontally. The value is taken from the radar's positioner resolution.

The average reflected power from power received from the 10 cm target on the 10 m mast is 198.0 dB 3.4 dBm. This provides an SCR value of $19.4 \cdot 19.4 \cdot 19.4 \cdot 19.4 \cdot 19.4 \cdot 19.4 \cdot 19.1 \cdot 19$

5.4 Temperature correction

BASTA-Mini has a regulation system to control temperature fluctuations inside the radar box. However, since the radar is based on solid state components, even small temperature fluctuations may impact the performance of the transmitter and receiver, and therefore affect the calibration stability. To account for this effect we introduced a temperature dependency in the calibration term, shown in Eq. (3).

During the experiments we verified the need of this correction by observing that the retrieved calibration term $C_{\Gamma}(T, r_0)$ and $C_{\Gamma}(T, F_0)$ has a consistent change depending on the time of the day, and that this change is strongly correlated to the temperature inside the radar.

Figure 6 (a), (b) and (c) show the results of a representative experiment done in the 2018 campaign. Here we left the radar sampling the target signal for several hours, to observe the variability of $C_{\Gamma}(T,r_0)$ $C_{\Gamma}(T,F_0)$ during the day. (a) shows the raw result in the RCS calibration term $C_{\Gamma}(T,r_0)C_{\Gamma}(T,F_0)$. There is a spread of almost 1 dB between the maximum and minimum values during the whole timeseries. (b) is a fourier transform of this raw timeseries. Here we can see that most of the variability happens in the timescale of hours. (c) presents the timeseries of (a), but in a daily cycle perspective. Here we plot hourly means of the deviation of $C_{\Gamma}(T,r_0)$ $C_{\Gamma}(T,F_0)$ with respect to the total average, with its hourly standard deviation as errorbars. We also superimposed atmospheric attenuation and the radar amplifier temperature to show that the first has a much smaller impact in calibration variability compared to the second.

Figure (d) shows the raw results of plotting variations in $C_{\Gamma}(T, r_0)$ $C_{\Gamma}(T, F_0)$ to temperature changes around $T_0 = 26.5$ °C. These variations are calculated independently for each iteration, by substracting the constant term of the linear fit of $C_{\Gamma}(T, r_0)$ $C_{\Gamma}(T, F_0)$ with respect to temperature. This operation removes the effect introduced by differences in alignment between different iterations. The reference T_0 value is chosen because it is approximately the average internal temperature when considering all the experiments.

To maximize the range of temperatures covered we choose to not limit the sampling period to one hour. This decision has the drawback of increasing the noise of the dataset due to the inclusion of some data taken under suboptimal conditions, for example with wind speed velocities above 1 meter per second or with the presence of drizzle. Yet, this step is necessary to enable the retrieval of the temperature correction function for the widest range of temperatures possible.

To retrieve the temperature dependency we use a linear regression perform a linear regression over the results from all the experiments done in 2018 and 2019, as shown in Fig. 7. In this case the data used was not limited to one hour, to maximize the range of temperatures covered. 7. The regression shows that the variability in the calibration term has an almost linear relationship with internal radar temperature, in the dB scale, and it is the same for both campaigns. This analysis allows us to estimate the value n = 0.093 dB ° C^{-1} for the temperature correction function of Eq. (3). To estimate the uncertainty of the temperature correction function we calculate the RMSE between the linear regression model and the whole dataset, for each degree of deviation in temperature. The RMSE value for the complete dataset is of 0.13 dB, while its value per degree ranges between 0.07 and all data, obtaining the value $\sigma_T = 0.13$ 0.23 dB for a deviation of 0 and +3 °C respectively. These results enable us to conclude that the temperature correction function uncertainty σ_T is ≤ 0.23 dB.

5.5 IF loss correction function $f_{IF}(F_b)$

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FMCW radars rely on estimating the beat frequency of the received signal to estimate the distance of an object. This signal may suffer uneven amplification depending on its frequency, because of a frequency dependent gain function in the amplifiers of the IF chain of the radar. Since there is a direct relationship between the beat frequency IF frequency F_b and the target distance r, this dependency on the beat frequency introduces a gain variability with respect to the target distance r. As introduced in Sect. 2, this distance dependency is compensated in the calibration term with a IF correction function $f_{IF}(F_b)$

To retrieve the $f_{IF}(F_b)$ we turn off the radar emitter and sample the environmental noise with the radar operating in its ealibration mode (12.5 meters distance resolution and 0.5 seconds integration time). The power $P_r(r)$ measured by the radar

under these conditions receiver when no active signal is input corresponds to the system noise temperature $N_s(F_b)$ plus the environmental noise power density N_0 , reduced by the radar total receiver loss $\log_{10}(L_r(T,r))\log_{10}(L_r(T,F_b))$, as indicated in Eq. (??9a) (Pozar, 2009). Noise power density is assumed to be

$$P_r(r) \equiv P_r(F_b) = N_0 + N_s(F_b) - 10\log(L_r(T, F_b)) \tag{9a}$$

$$\approx N_c - 10 \tag{9b}$$

The standard way to retrieve each of this terms is to perform a two point calibration. This requires the use of two noise sources at significantly different and well known temperatures. Usually, the temperatures of the loads are environmental temperature (298 K) and liquid nitrogen (77 K) (Rodríguez Olivos, 2015).

However, this approach was not practical with the equipment available on site, thus a different approach is used. BASTA-mini radar has a IF bandwith of 12 *MHz*, while the receiver Low Noise Amplifier (LNA) has a bandwidth of 35 *GHz*. Since the operational bandwidth is much narrower than the receiver full bandwidth, it is reasonable to assume that system noise is constant in the 12 IF bandwidth (168 to 180 *MHz* bandwith of this operation mode.).

$$P_r(r) = N_0 - 10\log(L_r(T, r))$$

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This enables the approximation of Eq. (9b). Here, we assume that system and environmental noise are constant in the IF frequency range, and that their addition can be expressed as a constant noise power N_c . To estimate the uncertainty introduced by this approximation, we observe that the gain and noise figure of the radar LNA have variations smaller than 0.1 dB in the operational frequency range. This indicates that uncertainty coming from assuming a constant system and environmental noise power for the IF bandwidth should be in the order of 0.1 dB or lower.

Then, to retrieve the $f_{IF}(F_b)$ we turn off the radar emitter and sample the environmental noise with the radar operating in its calibration configuration (12.5 meters distance resolution and 0.5 seconds integration time). After retrieving a significant amount of noise samples we calculate the average value of the difference $P_r(r_0) - P_r(r)$ for each distance $P_r(F_0) - P_r(F_b)$ for each IF frequency F_b , to remove the effect of the unknown noise power density. This operation is done to quantify relative gain variations around the calibrated frequency F_0 .

By using Eqs. (2a) and (3), we get that the difference $P_r(r_0) - P_r(r) - P_r(F_0) - P_r(F_b)$ is equivalent to the difference between $C_{\Gamma}(T, F_b)$ and $C_{\Gamma}(T, F_0)$, and therefore it is equivalent to the IF correction function $f_{IF}(F_b)$ (Eq. (10)). The temperature effect in gain is removed because both $P_r(r_0)$ and $P_r(r) - P_r(F_0)$ and $P_r(F_b)$ are sampled simultaneously, and therefore under the same temperature conditions.

$$\overline{P_r(F_0) - P_r(F_b)} = -10\log(L_r(T, F_0)) + 10\log(L_r(T, F_b)) = -C_\Gamma(T, F_0) + C_\Gamma(T, F_b) = f_{IF}(F_b)$$
(10)

For this experiment only, $P_r(r_0)P_r(F_0)$ corresponds to the power measured at the gate closer to the reference target position, without integrating other gates. This is done because there is no significant leakage and, as results of Fig. (8) show, $L_r(T,r)$ $L_r(T,F_b)$ changes are negligible in the five gates used for integration.

$$\overline{P_r(r_0) - P_r(r)} = -10\log(L_r(T, r_0)) + 10\log(L_r(T, r)) = -C_{\Gamma}(T, r_0) + C_{\Gamma}(T, r) = f_{IF}(r)$$

It must be noted that the IF correction function retrieved in this way is only valid for the operation mode used during its retrieval. If the radar operates in another mode, this function should be retrieved again, or a broader characterization of the IF loss must be used.

Figure 8 shows the results of the IF correction function retrieval referenced to $P_r(r_0)$ at $P_r(F_0)$, using F_0 associated to the target distance $r_0 = 376.5$ meters (corresponding to the 20 m mast experiment setup). We can observe that all functions retrieved in 2019 are in close agreement, without significant variations between different dates or time of the day chosen for the plots. The 2018 function is different because hardware was modified between both calibration campaigns. Additionally, in 2018 the emitter was not turned off to perform the noise sampling. Rather, we resorted to use a sampling period with clear sky conditions to respect the assumption assumptions of Eq. (??). 9b). To avoid the effect of crosstalk, we only consider gates farther than 200 meters from the radar.

A sixth degree polynomial is used to fit $f_{IF}(F_b)$ between 150 and 6000 meters, which is the maximum valid range for the 12.5 meter resolution mode (Delanoë et al., 2016). For both 2018 and all 2019 curves, the fit has a RMSE < 0.03 dB. Furthermore, the standard deviation between results from the four periods of 2019 has a maximum value of 0.04 dB for any gate. Both results indicate that the uncertainty introduced by the IF correction function is $\leq 0.04 dB$. Finally, the IF correction function retrieved for the 10 m mast setup in 2019 (with $r_0 = 196$ meters) is almost identical to the 20 m mast results. These functions are presented in Sect. 6. Considering these low RMSE values, we decide to select as the IF correction function uncertainty the uncertainty introduced by assuming a constant system noise, thus $\sigma_{IF} = 0.1 dB$.

415 5.6 Misalignment Bias

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The retrieval of $C_{\Gamma}(T, r_0)$ $C_{\Gamma}(T, F_0)$ using Eq. (4b) requires a precise knowledge of the reference target effective RCS Γ_0 . Each dBsm of difference between the theoretical value used in calculations and the effective target RCS will introduce a bias of the same magnitude in the estimation of the calibration coefficient C_{Γ}^0 , and thus in $C_{\Gamma}(T, r_0)C_{\Gamma}(T, F_0)$.

The effective reflector RCS is the actual physical value that would be measured by a perfectly calibrated radar. It is different from the target intrinsic RCS which only depends on its physical properties. Effective RCS changes when the experimental setup is modified. For example, if the point target is not exactly in the beam center, antenna gain will not be maximum and therefore the effective RCS will decrease compared to the intrinsic value. Effective RCS also changes when the incidence angle of the radar beam is modified. This latter effect may increase or decrease effective RCS depending on the original situation.

A common approach in these type of experiments is to set Γ_0 to be the maximum theoretical RCS of the target, assuming misalignment will cause a negligible deviation from this value. This procedure can be refined for cases where the system default

configuration does not have the target boresight aligned with the radar position. In these cases, effective RCS can be calculated using equations derived from geometrical optics (more complex optical calculations may be necessary for other wavelengths or target sizes). For example, we use the equations published by Doerry and Brock (2009) when calculating the effective RCS of our Triangular Trihedral target on the 20 m mast.

Unfortunately, this approach does not correct the impact of alignment uncertainties. We observed that random errors in the element positioning will statistically impact the effective Γ_0 in a single direction. Thus, simply taking the average of many target sampling iterations would result in a biased estimation of the calibration.

With the objective of quantifying the impact of alignment uncertainties we developed a geometrical simulator of effective RCS. This simulator receives as input the position of each element in the setup and calculates the effective RCS considering the beam incidence angle and antenna gain variations when the target is not in the center of the beam. The degrees of freedom included in the simulator are shown in Fig. 9 (a). It enables the modification of the radar aiming angles, the mast dimensions and the positioning and orientation of the target. The equations used in the simulator can be found in the article support material.

We now use the simulator to study how uncertainty in alignment can affect the value of Γ_0 . For this, we model an example experiment based on the 20 m mast setup. In this model we separate input variables between known and uncertain. Known terms can be fixed or measured very precisely in the field experiment, hence they are set as fixed values. Meanwhile, uncertain terms represent the parameters that cannot be fixed or measured very precisely, and for that reason are better expressed as probability distributions (terms defined in Fig. 9 (a)).

- Known terms:

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$$-x_r = 376.5 \text{ m}$$

$$-h_r = 5.3 \text{ m}$$

$$-\rho = 20 \text{ m}$$

$$-\alpha = 48^\circ$$

$$-\text{Target Size} = 20 \text{ cm}$$

- Variables with uncertainty:

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$$-\theta_r = \mathcal{N}(\theta_r^*, \sigma_{\theta_r}^2)$$

$$-\phi_r = \mathcal{N}(\phi_r^*, \sigma_{\phi_r}^2)$$

$$-\theta = \mathcal{N}(0, \sigma_{\theta}^2)$$

$$-\phi = \mathcal{U}([0^\circ, 360^\circ))$$

$$-\tau = \mathcal{N}(\tau^*, \sigma_{\tau}^2)$$

In the uncertain variables, $\theta_r^* = 87.82^\circ$, $\phi_r^* = 0^\circ$ and $\tau^* = 0^\circ$ represent the nominal alignment angles, which are the values expected under an ideal field experiment where the radar aims directly to the target and the mast is perfectly vertical. To these

nominal values we associate a distribution shape and the uncertainty set $\sigma_{\theta_r} = 0.075^{\circ}$, $\sigma_{\phi_r} = 0.075^{\circ}$, $\sigma_{\theta} = 1.5^{\circ}$, $\sigma_{\tau} = 5^{\circ}$. Each term, known and uncertain, is estimated from observations done during the experimental field work.

With these input parameters we sample the Γ_0 distribution that would arise after a large amount of experimental iterations. 460 Figure 9 (b) shows the results from this sampling. The black dashed line shows the effective RCS under our experimental configuration, when each element is in its nominal position. We can see that this effect cannot be neglected in our case, since its value is $0.8 \ dB$ lower than the maximum theoretical RCS.

However, this single correction does not suffice. The results of the model show that the addition of uncertainty into the process induces another bias of $\approx 0.3 \ dB$, in average. Since this is withing the order of magnitude of our desired uncertainty in the calibration, the example clearly illustrates the need of including a bias correction step in our calibration methodology.

The standard deviation σ_{ϵ} between N experimental retrievals of $C^0_{\Gamma i}$ cannot be used directly as an estimation of uncertainty because the RCS distribution shape is not Gaussian. The uncertainty introduced by this variability is studied by sampling a large set of possible RCS distributions based on our experimental configuration, and selecting the candidates matching our observed spread σ_{ϵ} . This set provides an estimation of the expected bias correction $\tilde{\Lambda}$ and of the effective RCS uncertainty σ_{Λ} . The uncertainty of the C^0_{Γ} estimator of Eq. (6) will correspond to the uncertainty of each $C^0_{\Gamma i}$ estimation propagated through the calculation of their average (terms $\sum \sigma_i^2/N^2$ and σ_T^2/N of Eq. (7a)) plus the effective RCS uncertainty σ_{Λ} . The details on how this estimator works and how the RCS distribution sampling is done are fully explained in Sect. S3 of the supplementary material.

6 Results

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In 2018 we used the 20 m mast only, performing six iterations. For 2019 we did 10 iterations using the 10 m mast and 2 iterations with the 20 m mast. The distributions of C_{Γ}^{0} obtained in each iteration and experiment is shown in Fig. 10.

The radar hardware changed between 2018 and 2019 campaigns due to experiments required to retrieve the power transfer curve and perform maintenance operations. This implies that we cannot compare absolute calibration values between both campaigns. What remains valid is to compare properties such as the variability, and the results from both experiments of 2019.

In the results we can notice a difference in $C_{\Gamma i}^0$ spread when comparing the 10 and 20 m masts. The 6 iterations of 2018 (Fig. 10 (A)) have an spread of $\sigma_{\epsilon}=0.33~dB$, while the spread of the 10 iterations of 2019 is 0.11 dB (Fig. 10 (B)). This happens because the 10 m mast has a motor on top which enables a much finer adjustment of the target position, improving the repeateability of the experiments.

There is also a small difference in the spread of the curves. The $C_{\Gamma i}^0$ values retrieved in experiment (B) have a smaller spread σ_i . This is because we took all the samples during one single night, with very clear conditions and an average wind speed below 1 m/s. A great advantage was the presence of the motor that enables target alignment in ≈ 5 minutes. Meanwhile, for experiment (A) curves were sampled during different days, because the 20 m mast setup requires more time to align (≈ 2 hours). The different conditions in each day led to a more varied shape in the retrieved curves. This effect is specially noticeable in experiment (C), where the iterations were performed during daytime, when atmospheric conditions are more

dynamic, specially wind speed variability. The introduced variability was not fully compensated by our corrections and thus bimodal distributions remained. However, individual spread is still small, within ≈ 0.1 dB, so we decided to accept these samples for calibration purposes.

To study the dependency of the bias correction on the amount of iterations we calculate the bias correction term $\tilde{\Lambda}$ and its uncertainty σ_{Λ} of experiments (A) and (B) with different amounts of repetitions. The order of the iterations used in each row match the sequential order indicated in Fig. 10. The results are shown in Table 1. For both cases we have the best estimate when we use all the samples available for each experiment, and thus we use this bias correction and uncertainty when computing the calibration coefficient.

For experiment (C) we followed a different approach. Because we only have two samples, the calculated $\sigma_{\epsilon} = 0.2 \ dB$ is very likely to be underestimated. Consequently, and because the experimental procedure was identical to what was done in 2018, we assume our parameters σ_{ϵ} , $\tilde{\Lambda}$ and σ_{Λ} to be equal to the best estimation of experiment (A). This is possible because in our methodology we assume that the bias probability distribution of a given system is unique, even if it is unknown, and what is done by performing many iterations is to successively restrict the possible sets of uncertainties that can generate results consistent with the observations. This latter hypothesis is consistent with the decrease in uncertainty for the bias correction when increasing the amount of iterations.

Table 2 contains a summary of all known bias corrections and uncertainty contributions, introduced in Sect. 4. With the afformentioned results, we use Eqs. (6), (3), (7a) and (7b) to estimate the RCS and Reflectivity calibration terms $C_{\Gamma}(T,r)$ and $C_{Z}(T,r)C_{\Gamma}(T,F_{b})$ and $C_{Z}(T,F_{b})$, alongside their uncertainty. $C_{Z}(T,r)$ Since the term $\sigma_{\Gamma_{0}}$ is much larger than all other uncertainty sources, we calculate a partial calibration uncertainty including all but this term, to simplify the comparison of uncertainty contributions between different experimental setups. This term is then added for the calculation of the final result.

510 $C_{Z}(T,F_{b})$ is calculated for the range resolution $\delta r=12.5~m$, which is the same mode used for target sampling. T is the radar amplifier temperature in ${}^{o}C$ and $f_{TF}(r)$ $f_{TF}(F_{b})$ is the IF loss correction function.

- (A) 20 m mast - 2018:

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$$\diamond \ C_{\Gamma}(T,r) = -275.6 + 0.093(T - 26.5) + f_{IF}(r) \left[dB(m^{-2}AU^{-1}) \right] \pm 0.3 \left[dB \right] C_{\Gamma}(T,F_b) = -80.98 + 0.093(T - 26.5) + f_{IF}(F_b) = -80.98 + 0.003(T - 26.5) + f_{IF}(F_b) = -80.98 + 0.003(T - 26.5) + f_{IF}(F_b) = -80.98 +$$

$$515 \qquad \diamond \ C_Z(T,r) = -191.5 + 0.093(T-26.5) + f_{IF}(r) \ [dB(mm^6m^{-5}AU^{-1})] \\ \pm 0.3 \ [dB] \ C_Z(T,F_b) = 3.05 + 0.093(T-26.5) + f_{IF}(F_b) \\ + 0.093(T-26.5) + 0.093(T-26.5) + 0.093(T-26.5) \\ + 0.093(T-26.5) + 0.093(T-26.5) + 0.093(T-26.5) \\ + 0.093(T-26.5) + 0.093(T-26.5) + 0.093(T-26.5) \\ + 0.093(T-26.5) + 0.003(T-26.5) \\ + 0.003(T-26.5) + 0.003(T-26.5) \\ + 0.003(T-$$

- (B) 10 m mast - 2019:

$$520 \qquad \diamond \ C_{\Gamma}(T,r) = -274.4 + 0.093(T - 26.5) + f_{IF}(r) \left[dB(m^{-2}AU^{-1}) \right] \pm 0.9 \left[dB \right] C_{\Gamma}(T,F_b) = -79.76 + 0.093(T - 26.5) + f_{IF}(F_b) = -274.4 + 0.093(T - 26.5) + f_{IF}(F_b) = -2$$

Table 1. Bias correction $\tilde{\Lambda}$ and its uncertainty σ_{Λ} calculated using a different amount of iterations, for the experiments of 2018 and 2019 calibration campaigns (for ex. 3 iterations means we used iterations 1, 2 and 3 of the experiment). We include the average and spread σ_{ϵ} between the retrieved $C_{\Gamma i}^0$ for each case. This variability σ_{ϵ} is introduced in the bias estimation procedure to determine the bias correction $\tilde{\Lambda}$ and its uncertainty σ_{Λ} .

(A) 20 m mast 2018	Exp. Results		Bias Correction	
N° of iterations	$\frac{1}{N} \sum C_{\Gamma i}^0$	σ_{ϵ} [dB]	$\tilde{\Lambda}$ [dB]	σ_{Λ} [dB]
2	-275.11 - <u>-80.51</u>	0.38	0.98	1.78
3	-275.19 <u>-80.59</u>	0.33	0.65	0.86
4	-275.25 <u>-80.65</u>	0.31	0.51	0.50
5	-275.24 <u>-80.64</u>	0.28	0.40	0.33
6	-275.14 <u>-80.54</u>	0.33	0.44	0.28
(B) 10 m mast 2019				
N° of iterations				
2	-274.15 -79.55	0.15	0.78	1.65
3	-274.16 <u>-79.56</u>	0.12	0.42	0.70
4	-274.17 -79.57	0.11	0.27	0.34
5	-274.20 - <u>79.60</u>	0.12	0.24	0.20
6	-274.23 -79.62	0.12	0.22	0.13
7	-274.23 -79.63	0.11	0.19	0.10
8	-274.22 -79.62	0.11	0.18	0.07
9	-274.21 -79.61	0.11	0.17	0.06
10	-274.20 -79.60	0.11	0.16	0.05
(C) 20 m mast 2019				
N° of iterations				
2	-273.41 - <u>78.81</u>	-	0.44	0.28

$$\diamond \ C_Z(T,r) = -190.3 + 0.093(T - 26.5) + f_{IF}(r) \ [dB(mm^6m^{-5}AU^{-1})] \pm 0.9 \ [dB] \ C_Z(T,F_b) = 4.28 + 0.093(T - 26.5) + f_{IF}(F_b) + f_{IF}(F_b) = 4.28 + 0.093(T - 26.5) + f_{IF}(F_b) + f_{IF}(F_b) = 4.28 + 0.093(T - 26.5) + f_{IF}(F_b) =$$

$$\diamond \ f_{IF}(r) = 4.87 \cdot 10^{-22} r^6 - 9.79 \cdot 10^{-18} r^5 + 6.35 \cdot 10^{-14} r^4 - 1.18 \cdot 10^{-10} r^3 - 1.16 \cdot 10^{-7} r^2 + 1.09 \cdot 10^{-4} r - 2.08 \cdot 10^{-2} \ [dB]$$

$$f_{IF}(F_b) = 7.60 \cdot 10^{-6} F_b^6 - 7.97 \cdot 10^{-3} F_b^5 + 3.48 F_b^4 - 8.10 \cdot 10^2 F_b^3 + 1.06 \cdot 10^5 F_b^2 - 7.40 \cdot 10^6 F_b + 2.15 \cdot 10^8 \ [dB]$$

- (C) 20 m mast - 2019:

$$\diamond \ C_{\Gamma}(T,r) = -273.9 + 0.093(T - 26.5) + f_{IF}(r) \left[dB(m^{-2}AU^{-1}) \right] \pm 0.4 \left[dB \right] C_{\Gamma}(T,F_b) = -79.25 + 0.093(T - 26.5) + f_{IF}(F_b) = -273.9 + 0.003(T - 26.5) + f_{IF}(F_b) = -273.9 +$$

$$\diamond \ C_Z(T,r) = -189.8 + 0.093(T - 26.5) + f_{IF}(r) \left[dB(mm^6m^{-5}AU^{-1})\right] \pm 0.4 \left[dB\right] C_Z(T,r) = 4.79 + 0.093(T - 26.5) + f_{IF}(F_b) +$$

These results enable the analysis of relative uncertainty contributions from different sources, however the total calibration uncertainty may be underestimated. As indicated in Sects. 4 and 5, some bias terms remain unknown. Specifically, target physical RCS must be measured in an echo chamber to improve the misalignment bias estimation. In addition, the method to characterize antenna alignment must be improved to determine if there is a need for an additional distance correction function (Sect. 5.2). The uncertainty of these retrievals will impact the total uncertainty value, however, it is possible to quantify this effect through the terms σ_{Γ_0} and σ_A of Eq. (7b).

To finalize, we perform a test of the calibration results by measuring a altostratus cloud in both campaigns (Fig. 11). The sampling was done with the 25~m resolution, and thus 6~dB had to be substracted from the $C_Z(T,r)$ $C_Z(T,F_b)$ calibration calculated for the 12.5~m resolution. In this correction, 3~dB come from the change in the distance resolution term δr (Eq. (5a)), and the other 3~dB are substracted to compensate the additional digital gain k_d coming from doubling the amount of points in the chirp fourier transform (Delanoë et al., 2016). A Signal to Noise Ratio threshold of 8~dB is used to remove noise samples. We observe that for both campaigns the reflectivity measured in altostratus cloud is within -30 - 0~dBZ, which are typical values reported in literature (Uttal and Kropfli, 2001).

7 Conclusions

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This study presents a cloud radar calibration method that is based on cloud radar power signal backscattered from a reference reflector. We study the validity of the method and variability of the results by performing measurements in two experimental setups and analyzing the associated results. In the first experimental setup we use a scanning BASTA-Mini W-band cloud radar, that aims towards a 20-cm triangular trihedral target installed at the top of a 20-m mast, located 376.5 m from the radar. For the second experimental setup, we use the same radar, aimed towards a 10-cm triangular trihedral target mounted on a pan-tilt motor at the top of a 10-m mast. The mast is located 196 m from the radar.

The first consideration in the design of the experimental setup is the need to avoid excessive compression or saturation in the radar receiver. This must be checked before any calibration attempt by comparing measurements of radar backscattered power with the radar receiver power transfer curve. In both our setups we find losses due to compression on the order of $0.2 \sim 0.3$ dB. There is a compensating effect between target RCS and radar-to-target distance (Eq. 2b). Since the compression effect is small, we correct it using our receiver power transfer curve. However, in cases where the radar is operating close to saturation, or when compression effects are larger than the calibration uncertainty goal, it is advisable to compensate by reducing target size or by positioning the target farther away from the radar.

Secondly, the reflector must be positioned far enough from the radar to be outside the antennas near-field distance and to ensure that the received power has low antenna-overlap losses. The BASTA-Mini cloud radar has a Fraunhofer near-field distance of 50 m. The estimated maximum overlap loss is less than $0.1 \ dB$ for the closest (10-m) mast setup. Thus we conclude that the target positioning is far enough for both setups.

Thirdly, the experimental setup should strive to reduce clutter in the radar measurements. This can be achieved by operating in an open field that is several hundred meters in length and free of trees or other signal-inducing obstacles. It is also advisable to perform radar measurements under clear conditions, without fog or rain, with wind speed below $1 ms^{-1}$ and low turbulence.

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Next, the proposed calibration method requires performing several iterations in the same setup configuration. In each iteration the setup is first realigned, followed by approximately one hour of sampling of the reference reflector backscattered power. The sampled power is then corrected for compression effects, incomplete antenna overlap, variations in radar gain due to temperature and atmospheric attenuation, before being used to estimate a RCS calibration term value. Once all iterations are completed, the final RCS and Equivalent Reflectivity calibration terms can be computed with their respective uncertainties.

Iterations are necessary because they enable the quantification of bias introduced by inevitable system misalignment. Our experiments indicate that, for our setup, at least 5 iterations are necessary to reach convergence in the calculation of bias and uncertainty associated with misalignment. We find a bias correction of $\approx 0.4 \pm 0.3~dB$ for the 20-m mast, and of $\approx 0.2 \pm 0.1~dB$ for the 10-m mast. This difference can be explained by the more precise alignment attainable with the pan-tilt motor installed on the 10 m mast.

Calibration is also impacted by changes in the gain of radar components associated with internal temperature variations. For the radar used in our experiment, these changes reach up to $\pm 0.6~dB$. Our experiments enabled us to retrieve a correction function for the temperature dependence and to reduce the temperature uncertainty contribution to $\sigma_T = 0.13~\sigma_T = 0.23~dB$. This result indicates that lower calibration uncertainties can be achieved by studying temperature effects, especially for solid state radars.

Another necessary consideration is the inclusion of gain variations with distance, introduced by frequency dependent losses in the IF of the radar receiver. We found calibration variations with distance up to $0.9 \ dB$ for the 2019 campaign. Therefore, characterizing the IF loss is a necessary step to validate the calibration results for all ranges.

Our analyses reveal that the predominant sources of uncertainty in our experimental setups are due to source of uncertainty for all experiments is the reference target RCS, reaching approximately $2\ dB$ due to the use of a theoretical model instead of an experimental characterization. The next most important contributions to uncertainty come from the levels of clutter and alignment precision. These two effects have different magnitudes in our two experimental setups (10-m and 20-m masts). The 20-m mast setup uncertainty is limited by the uncertainty contribution of the alignment bias estimation $\sigma_{\Lambda} = 0.28\ dB$. This result reveals that there is a tradeoff between better target alignment and additional clutter introduced by the alignment motor.

The complete uncertainty budget enables us to conclude that the proposed calibration method can yield uncertainties as low as to reach a calibration uncertainty of $0.5 \, dB$, it is necessary to have a target RCS characterization with an uncertainty of $0.3 \, dB$ with our current equipment. This result, based on the accumulated uncertainty of all terms, except target RCS, of $0.4 \, dB$. This uncertainty was obtained using the 20-cm target on the 20 m mast during the 2018 experiment, where six target sampling iterations were performed. Additionally, in 2019 two completely different calibration setups were used with the same radar hardware, and in both cases, we obtain the same calibration result within uncertainty bounds.

Finally, because of cloud radar hardware evolutions in the fall of 2018, the calibration coefficients found in May 2018 and March 2019 differ by $1.2 \ dB$. We compare cloud radar measurements of altrostratus clouds performed in May 2018 and March 2019. The reflectivity distributions of the two events are consistent and compatible with values previously registered in literature. The two distributions yield median values that differ by $0.3 \ dB$.

For future work we envisage to develop a technological solution to allow target orientation without introducing additional clutter. Another interesting prospect is to improve the accuracy of the radar positioner, to enable direct retrieval of antenna pattern directly with the radar, following the method proposed by Garthwaite et al. (2015). This retrieval would improve bias correction arising from parallax errors. We also plan to perform an echo chamber characterization of our reference targets, to remove any possible bias caused by manufacturing imprecision, to reduce its contribution to total uncertainty and to improve the estimation of our misalignment bias correction.

Further, there is ongoing research on calibration and antenna pattern characterization methods based on reference targets held by Unmanned Aerial Vehicles (UAVs) (Duthoit et al., 2017; Yin et al., 2019). Since the underlying principle is the same, most considerations written here should be directly applicable in these new experiments. Here the UAV takes the role of the mast, holding the reflector (usually a sphere), and therefore it is important to characterize the UAV RCS and verify that it does not interfere with the experiment. The main difference would be in the procedure necessary to estimate bias, because the reference target (usually a sphere) will be always moving due to wind. Here an adaptation of the effective RCS simulator would be necessary to account for the target type and different alignment protocol.

615 Author contributions.

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All authors contributed to the planning of the campaigns and the design of the calibration experiments.

Author Julien Delanoë was responsible of radar installation and operation.

Authors Jean-Charles Dupont and Felipe Toledo worked in the preparation, development and operation of the necessary infrastructure for the experiments.

Authors Julien Delanoë and Felipe Toledo retrieved the Power Transfer Curve of the Radar Receiver.

Data analysis and the establishment of the calibration methodology presented in the paper was done by Felipe Toledo.

Authors Martial Haeffelin and Felipe Toledo worked in defining the paper structure and content.

Authors Felipe Toledoand Susana Jorquera, Susana Jorquera and Cristophe Le Gac worked in developing the method to retrieve the IF correction function, and in its calculation.

Author Christophe Le Gac contributed with technical information about the radar.

All authors reviewed the paper.

Competing interests.

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Symbols

Symbol	Description	Units
$C_{\Gamma}(T,F_b)$	RCS Calibration Term	dB
$C_{\Gamma}(T,F_0)$	RCS Calibration Term at the IF frequency F_0 , associated to the reference	dB
	target position r_0	
C_{Γ}^0	RCS Calibration Coefficient	dB
$C^0_{\Gamma s}$	Single sample of the Calibration Coefficient C_{Γ}^0	dB
$C_Z(T, F_b)$	Radar Equivalent Reflectivity Calibration term	dB
δC_{Γ}	RCS calibration uncertainty	dB
δC_Z	Reflectivity calibration uncertainty	dB
F_b	Signal frequency at the radar receiver IF	MHz
$f_{IF}(F_b)$	IF loss correction function	dB
$\Gamma(r)$	Radar Cross Section of reflections at a distance r	dBsm
Γ_0	Radar Cross Section of the reference target	dBsm
k_d	Digital gain of the radar receiver	
$ ilde{\Lambda}$	Misalignment bias correction	dB
λ	Radar carrier wavelength	m
N	Number of iterations performed in a calibration experiment	
$P_r(r_0)$	Power received from the target position r_0	dBm
$P_r(r)$	Power received from distance r	dBm
p_t	Radar transmitted power	W
r	Distance from the radar	m

Symbol	Description	Units
r_0	Distance between the reference target and the radar	m
σ_A	Uncertainty of the antenna properties (beam shape and alignment)	dB
σ_ϵ	Standard deviation between the N mean RCS Calibration Coefficients $C^0_{\Gamma i}$,	dB
	used to calculate the bias correction	
σ_{Γ_0}	Uncertainty of the reference target RCS	dB
σ_i	Standard deviation of the RCS Calibration Coefficient samples $C^0_{\Gamma s}$ for iter-	dB
	ation i	
σ_{IF}	Uncertainty of the IF loss correction function	dB
σ_{Λ}	Uncertainty of the misalignment bias correction	dB
σ_{SCR}	Uncertainty introduced by clutter at the target position	dB
σ_T	Uncertainty of the temperature correction function	dB
θ	Antenna beamwidth	rad
Z_e	Radar Equivalent Reflectivity	dBZ

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Table 2. Summary of all known corrections and uncertainty contributions in the calculation of $C_{\Gamma}(T,r)C_{\Gamma}(T,F_b)$. The absolute correction terms have a sign associated with the direction in which they impact the final calibration calculation. For the receiver compression correction we present the average magnitude and for the temperature correction we present the range of possible values. The Partial Calibration Uncertainty is the addition of all uncertainty terms except σ_{Γ_0} . This term is later added to calculate the Total Calibration Uncertainty.

Absolute Corrections	Term [dB]	(A) 20 m mast 2018	(B) 10 m mast 2019	(C) 20 m mast 2019
Compression	Fig. 4 (a)	-0.3 in avg.	-0.2 in avg.	-0.3 in avg.
Partial Antenna Overlap	$L_o(r_0)$	-0.02	-0.08	-0.02
Temp. Corr. ($T_0=26.5~^{\circ}\mathrm{C}$)	$n(T-T_0)$	within ± 0.6	within ± 0.6	within ± 0.6
Misalignment Bias	$ ilde{\Lambda}$	-0.44	-0.16	-0.44
IF loss correction	$f_{IF}(r)$ - $f_{IF}(F_b)_{\sim}$	≤ 10.61	≤ 10.91	≤ 10.91
Uncertainty Sources	Term [dB]			
$C^0_{\Gamma i}$ estimation	$\sqrt{rac{1}{N^2}\sum\sigma_i^2}$	0.03	0.01	0.09_0.07_
Temp. Corr. in $C^0_{\Gamma i}$ retrievals	$rac{\sigma_T}{\sqrt{N}}$	0.05_0.09	0.04 _0.07	0.090.16
Temp. Corr. in $C_{\Gamma}(T,r)$, $C_{Z}(T,r)$ $C_{\Gamma}(T,F_{b})$, $C_{Z}(T,F_{b})$	σ_T	0.13_0.23	0.13_0.23	0.130.23
Signal to Clutter Ratio	σ_{SCR}	0.09	0.93	0.09
Bias Correction	σ_{Λ}	0.28	0.05	0.28
IF loss correction	σ_{IF}	< 0.04 0.1	< 0.04 0.1	< 0.040.1
Partial Calibration Uncertainty		0.40	0.97	0.43
Reflector RCS Uncertainty	$\sigma_{\Gamma_{\mathbf{Q}}}$	2	2	2~
Total Calibration Uncertainty	$\delta C_{\Gamma}; \delta C_{Z}$	0.33-2.04	0.94-2.22	0.35 -2.04

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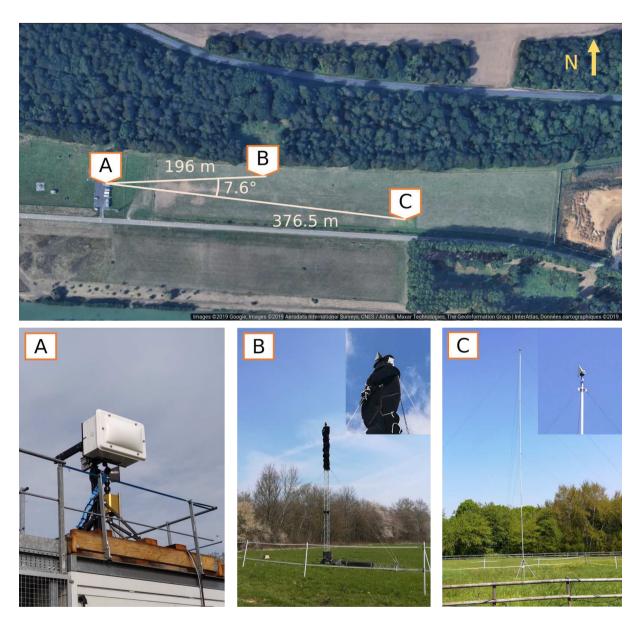


Figure 1. Experimental setup for 2018 and 2019 calibration experiments: (A) Scanning BASTA-Mini radar located in a reinforced platform 5 m above the ground. (B) 10 mast with a 10 cm triangular trihedral target mounted on a pan-tilt motor with an angular resolution and repeateability better than 0.1° . This mast has microwave absorbing material wrapped to it to reduce its RCS (clutter). The 10 m mast was only installed in the 2019 calibration campaign. (C) 20 m mast with a 20 cm triangular trihedral target. The target aiming is fixed relative to the mast. This mast is used in both 2018 and 2019 calibration campaigns. Angular separation between the masts is enough to sample both targets without mutual interference.

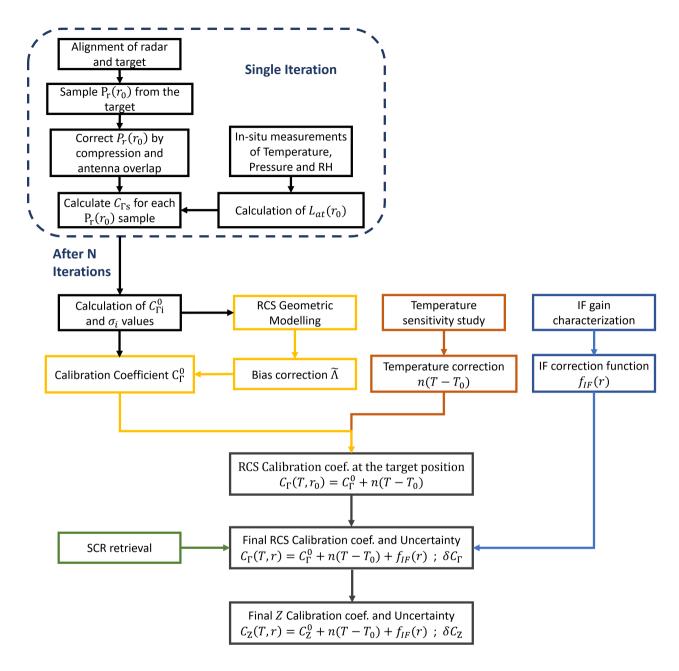


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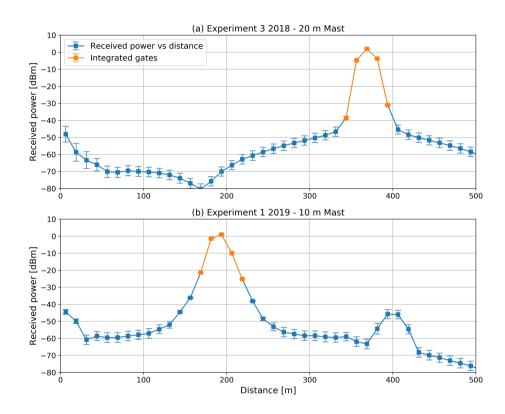


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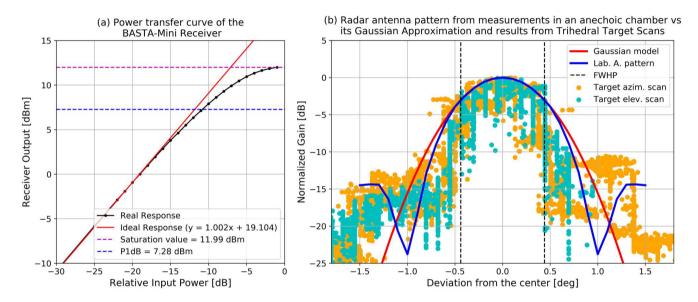


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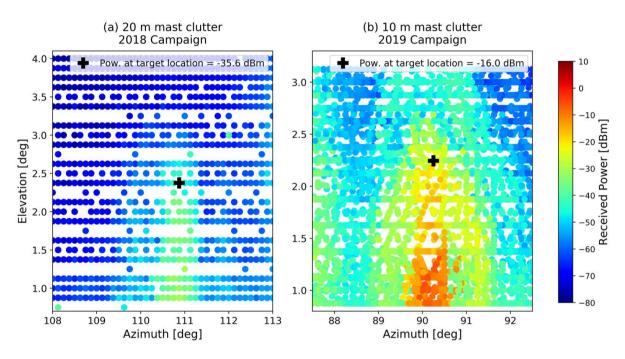


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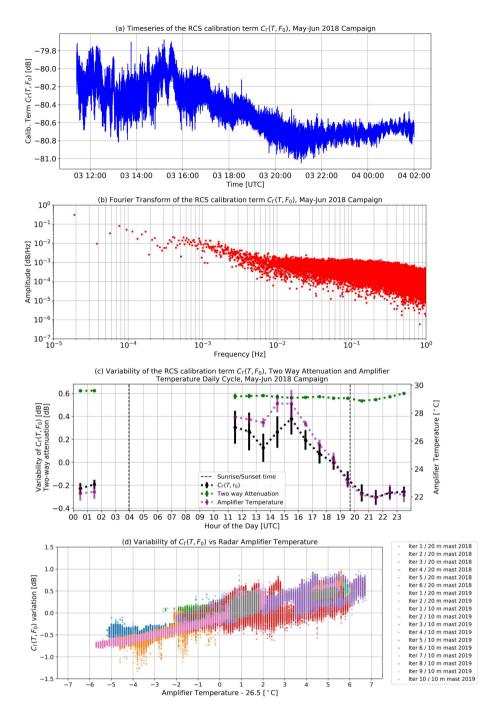


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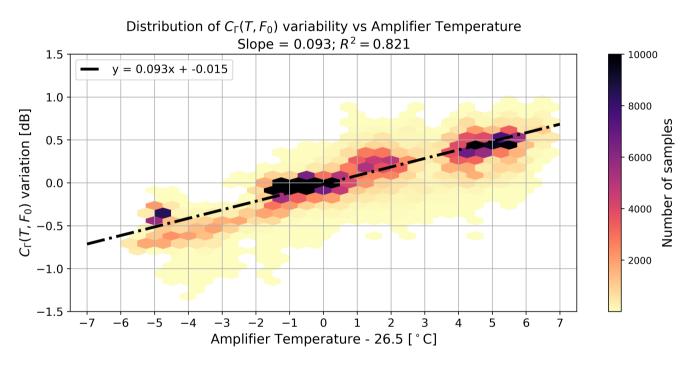


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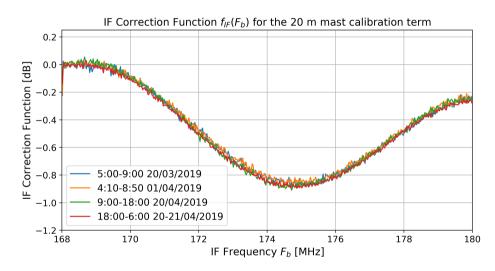
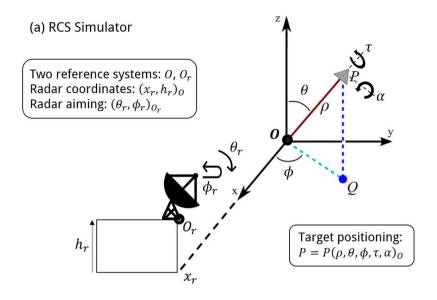


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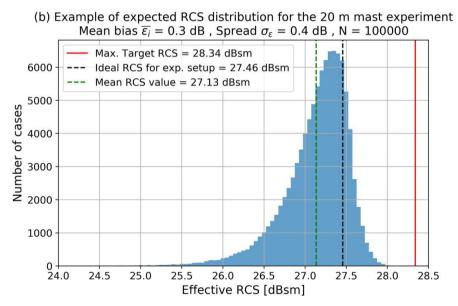


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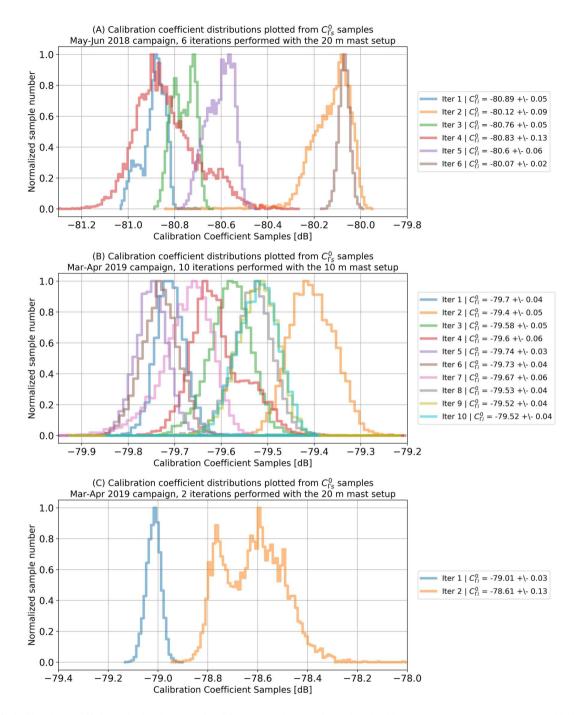


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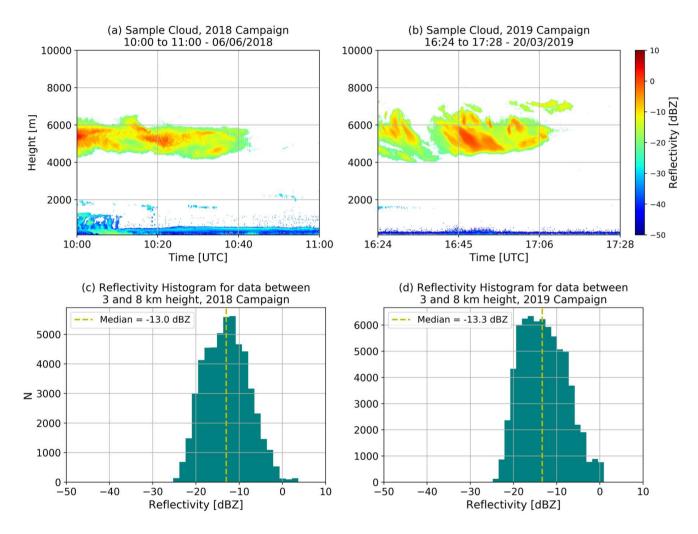


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