

1 Assessing the accuracy of low-cost optical particle  
2 sensors using a physics-based approach

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4 David H Hagan<sup>1,2</sup>, Jesse H Kroll<sup>1,3</sup>

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6 <sup>1</sup>Department of Civil and Environmental Engineering, Massachusetts Institute of Technology,  
7 Cambridge, MA 02139, USA

8 <sup>2</sup>QuantAQ, Inc., Somerville, MA 02143, USA

9 <sup>3</sup>Department of Chemical Engineering, Massachusetts Institute of Technology, Cambridge, MA  
10 02139, USA

11

12 Corresponding Author Emails: david.hagan@quant-aq.com or jhkroll@mit.edu

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1 **Abstract**

2 Low-cost sensors for measuring particulate matter (PM) offer the ability to understand  
3 human exposure to air pollution at spatiotemporal scales that have previously been  
4 impractical. However, such low-cost PM sensors tend to be poorly characterized, and  
5 their measurements of mass concentration can be subject to considerable error. Recent  
6 studies have investigated how individual factors can contribute to this error, but these  
7 studies are largely based on empirical comparisons, and generally do not examine the  
8 role of multiple factors simultaneously. Here, we present a new physics-based framework  
9 and open-source software package (*opcsim*) for evaluating the ability of low-cost optical  
10 particle sensors (optical particle counters and nephelometers) to accurately characterize  
11 the mass loading of aerosol particles. This framework, which uses Mie theory to calculate  
12 the response of a given sensor to a given particle population, is used to estimate the  
13 fractional error in mass loading for different sensor types, given variations in relative  
14 humidity, aerosol optical properties, and the underlying particle size distribution. Results  
15 indicate that such error, which can be substantial, is dependent on the sensor technology  
16 (nephelometer vs. optical particle counter), the specific parameters of the individual  
17 sensor, and differences between the aerosol used to calibrate the sensor and the aerosol  
18 being measured. We conclude with a summary of likely sources of error for different  
19 sensor types, environmental conditions, and particle classes, and offer general  
20 recommendations for choice of calibrant under different measurement scenarios.  
21

22 **1. Introduction**

23 Human exposure to aerosols is associated with adverse health impacts and increased  
24 mortality (Apte et al., 2018; Burnett et al., 2018; Cohen et al., 2017; Dockery et al., 1993).  
25 The source and composition of aerosols has been linked to a range of negative health  
26 impacts (Antonini et al., 2003; Hart et al., 2012; Henneberger and Attfield, 1997; Lipsett

1 and Campleman, 1999), with more than 4 million annual deaths worldwide attributed to  
2 ambient particulate matter pollution (Cohen et al., 2017). Accurate estimates of aerosol  
3 sources and health impacts rely critically on measurements of particulate matter mass  
4 concentrations across indoor and outdoor environments worldwide.

5

6 In many countries, particulate matter (PM) pollution is regulated by national or local  
7 government agencies (e.g., the US EPA in the United States) and is typically measured  
8 using federally approved reference methods that are high in accuracy and precision. The  
9 existing infrastructure is generally designed to measure regional-scale air pollution, in  
10 order to enforce (and assess the effectiveness of) air quality regulations. However, particle  
11 pollution can vary in space and time at much finer resolution than can be measured using  
12 standard monitoring technologies, due to their relatively high cost and size. Over the past  
13 several years, new technologies have emerged at price points (<\$2000) that allow PM  
14 measurements to be made with much higher spatiotemporal resolution, even down to  
15 the individual human level (Koehler et al., 2019; Tryner et al., 2019a, 2019b). These devices  
16 are physically small, use very little power, and can easily be deployed at scale. As a result,  
17 such sensors are ideally suited for use in dense distributed sensor networks, providing  
18 high-resolution air quality measurements, as well in as personal monitoring, providing  
19 individuals with the ability to measure and understand their exposure to harmful air  
20 pollutants. As with all low-cost sensors (LCS), accuracy is of paramount concern; as shown  
21 by a number of recent laboratory and field-based evaluation studies (Crilley et al., 2018;  
22 Dacunto et al., 2015; Di Antonio et al., 2018; Holstius et al., 2014; Levy Zamora et al., 2019;  
23 Malings et al., 2020; Northcross et al., 2013; Sousan et al., 2016b, 2016a; Wang et al.,  
24 2015), PM sensors can perform quite poorly without additional constraints or calibrations.

25

1 Most low-cost PM sensors measure particles via light scattering. Sampled particles  
2 intercept a beam of light (typically from a laser or LED with a wavelength between 405  
3 and 780 nm), and the scattered light is measured and correlated to a PM mass  
4 concentration. In this work, we refer to such instruments as optical particle sensors (OPSs).  
5 OPSs can be broken down into two main types, nephelometers and optical particle  
6 counters (OPCs). Nephelometers measure the particles as an ensemble, gathering light  
7 scattered by all particles across a wide range of angles, typically 7°-173° to avoid pure  
8 forward and backward scattering (Abu-Rahmah et al., 2006; Ahlquist and Charlson, 1967;  
9 Anderson et al., 1996). The total scattering amplitude is then correlated to a mass  
10 measurement made by a reference instrument. (Nephelometers that measure scattered  
11 light at a single angle are sometimes referred to as photometers; for the purposes of this  
12 work we consider photometers to be a subclass of nephelometer.) OPCs, by contrast,  
13 detect particles individually, providing information on their number and size. Light  
14 scattered by each individual particle is measured and each pulse is assigned to a size bin  
15 based on its total light intensity, resulting in a histogram which is converted to a mass  
16 loading once the entire distribution has been measured. While these technologies have  
17 been around for decades (Gucker et al., 1947; Patterson et al., 1926), they have recently  
18 become available at much lower cost due to the availability of small, inexpensive light  
19 sources and electronic components.

20

21 The use of light scattering introduces a number of fundamental limitations for making PM  
22 mass measurements. Many of these arise from environmental conditions and/or the  
23 properties of the aerosol being measured; these can be especially problematic when  
24 calibration is done using only a single aerosol type or condition. A number of recent  
25 empirical studies of OPSs have investigated some of these limitations. These issues  
26 include: (1) the inability to adapt to changes in the particle size distribution (Dacunto et

1 al., 2015; Wang et al., 2015); (2) the hygroscopic growth of particles due to changes in  
2 ambient relative humidity (Crilley et al., 2018; Di Antonio et al., 2018; Malings et al., 2020;  
3 Zheng et al., 2018); (3) changes in scattering efficiency due to differences in aerosol  
4 optical properties (Crilley et al., 2018; Di Antonio et al., 2018); and (4) the need for aerosol-  
5 specific correction factors to account for differences in density (Dacunto et al., 2015;  
6 Northcross et al., 2013). While these studies have examined how these individual effects  
7 in isolation may affect PM accuracy, to our knowledge there has not been a systematic,  
8 comprehensive investigation of all these factors together. Complicating matters is the  
9 fact that these individual properties are all intertwined – for example, when relative  
10 humidity increases, it can cause particles to take up water, which can change not only  
11 their size and mass but also their shape, refractive index, and density.

12

13 To examine the relative contribution of error by these various sources, we have  
14 developed a model that describes how a given sensor will respond to different aerosols  
15 under a wide range of conditions. This model is based entirely on the underlying physics  
16 of light scattering (Mie Theory) rather than empirical relationships obtained through  
17 laboratory or field measurements. While previous work has modeled nephelometers and  
18 OPCs in a similar way (Walser et al., 2017), we believe this is the first detailed treatment  
19 of light scattering as it relates specifically to LCS. We use this model to isolate the relevant  
20 sources of error and develop a better understanding of the limitations (as well as  
21 strengths) of different kinds of OPSs.

22

23 The modeling tool described here, which is open source and freely available, can be used  
24 for the systematic study of how different OPSs may detect various aerosol types under a  
25 range of environmental conditions. This enables new insights into the potential errors  
26 associated with a given PM measurement, optimal strategies for calibrating OPSs, the

1 development of algorithms for data analysis, and ultimately in the design of the sensors  
2 themselves. The objective of this work is to describe the model and software and to  
3 investigate broad influences on aerosol properties and sensor parameters on  
4 measurement performance. We do not investigate the performance of individual  
5 commercially available sensors under the full range of conditions expected in the  
6 atmosphere; but such studies are enabled by this modeling tool and are an important  
7 future extension of this work.

## 8 2. Methods

9 The modeling framework described in this section is available as an open-source (MIT  
10 license) python library (*opcsim*) and has been made available on GitHub. Detailed  
11 documentation, including installation instructions and examples, are available online  
12 (Hagan and Kroll, 2019). The framework, called “opcsim”, consists of two primary  
13 components: the code that models OPSs and implements the Mie Theory algorithms  
14 (Bohren and Huffman, 1983; Sumlin et al., 2018), and the code to build and evaluate  
15 aerosol distributions.

16  
17 We follow the same general modeling pattern regardless of sensor type. Steps include:  
18 (1) defining the device based on its key physical parameters; (2) calibrating the device to  
19 a specific aerosol type (for OPCs) or aerosol distribution (for nephelometers); and (3)  
20 evaluating each particle in an aerosol population by computing the scattered light signal  
21 using Mie theory and converting that signal to the sensor output based on its calibration.  
22 In the following sections we describe how the aerosol population is described by the  
23 model, followed by how the sensors themselves are treated.

24

## 1 2.1 Representing an aerosol distribution

2 We represent an aerosol distribution as the sum of  $n$  lognormal modes, where each mode  
3  $i$  is defined by its geometric mean particle diameter ( $\bar{D}_{pi}$ ), geometric standard deviation  
4 ( $\sigma_i$ ), and number concentration ( $N_i$ ). The aerosol distribution as a function of diameter  $D_p$   
5 ( $dN/d\log D_p$ ) is given by Equation 1 (Seinfeld and Pandis, 2006):

6

$$7 \quad \frac{dN}{d\log D_p} = \sum_{i=1}^n \frac{N_i}{\sqrt{2\pi} \log \sigma_i} \exp\left(-\frac{(\log D_p - \log \bar{D}_{pi})^2}{2 \log^2 \sigma_i}\right) \quad (\text{Equation 1})$$

8

9 Additionally, we define the composition of the aerosol distribution by defining the  
10 particle density ( $\rho_i$ ), hygroscopic growth factor ( $\kappa_i$ ), and complex refractive index ( $m_i$ ) for  
11 each mode. The role of these additional parameters is discussed in section 3, below.  
12 While more complex representations of the chemical makeup of the aerosol can be  
13 implemented using our modeling framework (i.e., core-shell representation of aerosols,  
14 complex aerosol mixtures, etc.), for the purposes of this manuscript we focus only on well-  
15 mixed homogeneous particle modes, as described by Eq. 1. The above number  
16 distribution can be converted to a mass distribution (or total mass concentration) by  
17 assuming all particles are spherical with a known density (Seinfeld and Pandis, 2006).

18

19

## 20 2.2 Representing Optical Particle Sensors

### 21 2.2.1 Optical Particle Counters (OPCs)

22 An OPC is defined by three instrument-specific parameters: (1) the wavelength of the  
23 light source ( $\lambda$ ), (2) the viewing angle for which the scattered light is collected, and (3) the  
24 number of discrete size bins and their widths. A bin, in this context, refers to a single  
25 "slice" of the aerosol size distribution, with a fixed width and units of particle diameter.

1 Typically, most low-cost OPCs have between 2-30 bins. These can be determined either  
 2 by looking up the parameters in the device’s datasheet provided by the manufacturer or  
 3 by making simple measurements. Bins are often chosen to reduce the uncertainty in  
 4 correct bin assignments within the bounds of what the sensor is capable of detecting.  
 5 Most low-cost OPCs have the smallest bin at  $D_{\min} \sim 500$  nm, with cost typically being the  
 6 driving factor – OPCs with lower  $D_{\min}$  employ more expensive, higher-quality optics and  
 7 photodetectors, allowing them to detect smaller particles. In this work, the bin  
 8 boundaries (and hence widths) used for a given OPC are taken from the manufacturer’s  
 9 spec sheets, if available; otherwise they are calculated by generating an array of  
 10 logarithmically-spaced bin boundaries for a set number of bins ( $n_{\text{bins}}$ ) between the  
 11 minimum and maximum defined diameters ( $D_{\min}$  and  $D_{\max}$ , respectively). Most often, a  
 12 light pulse generated by a single particle is assigned to exactly one bin; however, there  
 13 exist approaches where bin assignments are made using a probability distribution (Walser  
 14 et al., 2017); this is not implemented in this model but is an approach that could be added  
 15 in the future.. Table 1 lists bin widths and other parameters for a few commercially  
 16 available low-cost OPCs.

17

18 **Table 1.** Characteristics of a selection of commercially available low-cost optical particle  
 19 counters and nephelometers.

Manufacturer	OPS Type	Model	$\lambda$ (nm)	Viewing Angle ( $\varnothing_1, \varnothing_2$ )	# of Size Bins
Alphasense, Ltd.	OPC	OPC-N2	658	(32.0°, 88.0°)	16 (0.38 – 17.5 $\mu\text{m}$ )
Alphasense, Ltd.	OPC	OPC-N3	658	(32.0°, 88.0°)	24 (0.35 – 40.0 $\mu\text{m}$ )
Particle Plus	OPC		785	(58.0°, 118.0°)	6 (0.3 – 10.0 $\mu\text{m}$ )
NOAA/Handix	OPC	POPS	405	(38.0°, 142.0°)	16 (0.132 – 3.65 $\mu\text{m}$ )



Plantower	Nephelometer	PMS5003	~650	? <sup>1</sup>	6 (0.3 – 10+ $\mu\text{m}$ ) <sup>2</sup>
Sharp	Nephelometer	GP2Y1010AUOF	870-	? <sup>1</sup>	1 (?) <sup>1</sup>
	(Photometer)		980		
Shinyei	Nephelometer	PPD42NS	870-	? <sup>1</sup>	1 (>1 $\mu\text{m}$ )
	(Photometer)		980		
Samyoung	Nephelometer	DSM501A	870-	? <sup>1</sup>	1 (>1 $\mu\text{m}$ )
	(Photometer)		980		

1 <sup>1</sup> Unknown; not provided in the manufacturer’s technical data sheet or the technical literature

2 <sup>2</sup> The PMS5003 reports six bins; however these are not actual size bins, but rather software-  
3 computed results (He et al., 2020).

4

5 OPCs are calibrated by relating the scattered light intensity – a combination of the  
6 particle’s scattering cross section ( $C_{\text{scat}}$ ) and laser intensity – to the particle diameter.  
7 Practically, this is done by using calibration aerosols with known optical properties and  
8 size and generating a calibration curve between the test aerosol and the electronic pulse  
9 height generated by that aerosol. After repeating this process for many sizes, a calibration  
10 curve can be generated. Here, we compute the  $C_{\text{scat}}$  values using Mie theory using  
11 attributes of the calibration aerosol. To simplify the model, we make several assumptions,  
12 including: (1) all particles are spherical and homogeneous (well-mixed); (2) the laser  
13 intensity is constant, implying all particles are perfectly centered in the beam of the laser;  
14 and (3) the photodetector and electronics are 100% efficient, and so we do not consider  
15 the impact of signal-to-noise limitations.

16

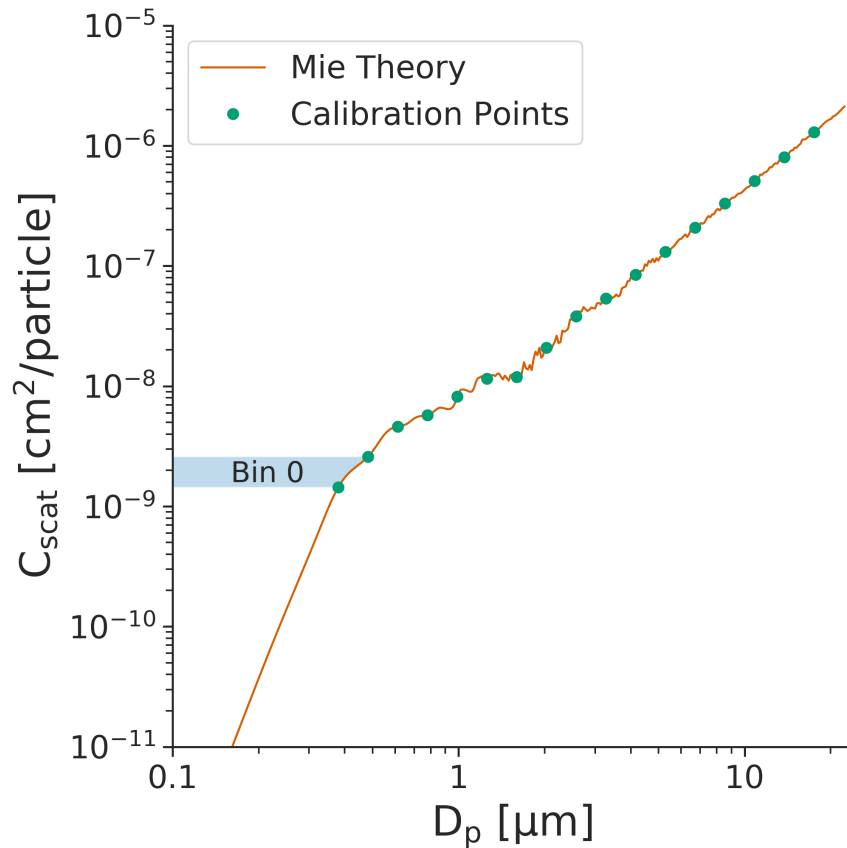
17 As most low-cost OPCs contain an elliptical re-focusing mirror to gather the scattered  
18 light across many angles, we compute the integrated light scattering intensity following  
19 a procedure first introduced by Jaenicke and Hanusch (Jaenicke and Hanusch, 1993). Mie  
20 theory calculations are implemented using equations by Bohren and Huffman (Bohren  
21 and Huffman, 1983). The scattering cross-section is calculated as:

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$$C_{\text{scat}} = \frac{\lambda}{4\pi} \int_{\theta_1}^{\theta_2} [i_1(\theta) + i_2(\theta)] \sin\theta \, d\theta \quad (\text{Equation 2})$$

where  $\lambda$  is the wavelength of incident light,  $\theta$  is the viewing angle (which ranges from  $\theta_1$  to  $\theta_2$ ), and  $i_1$  and  $i_2$  are the intensity distribution functions (Bohren and Huffman, 1983).

Figure 1 depicts the calibration curve generated for an OPC with the characteristics of the Alphasense OPC-N2 (Table 1), using polystyrene latex spheres (PSLs) of different diameters for calibration. Eq. 2 was used to compute the theoretical  $C_{\text{scat}}$  values (y axis), integrated across the entire viewing angle, for a range of particle diameters (x axis). The  $C_{\text{scat}}$  values at each bin boundary (green dots in Fig. 1) are then computed, and spline interpolation is used between each individual bin boundary to generate a mapping between the scattering amplitude and its corresponding bin assignment. In practice, this operates as a lookup table – a particle crossing the laser generates a scattering amplitude which is associated with a specific ‘bin’ via the calibration.



1  
 2 **Figure 1.** Calibration data for an OPC with 16 discrete size bins between 0.38 – 17.5  $\mu\text{m}$ . OPC  
 3 parameters were chosen to match the Alphasense OPC-N2 (wavelength of 658 nm, viewing angle  
 4 of 32-88°) using monodispersed polystyrene latex spheres ( $m = 1.592 + 0j$ ). The integrated  
 5 scattering amplitude calculated using Mie theory is shown as the solid line, with points depicting  
 6 the corresponding scattering amplitude at each of the bin boundaries. Shown as a shaded box is  
 7 the range of scattering amplitudes that is assigned to the smallest size bin.

8  
 9 For OPCs that measure scattered light across a wide angle,  $C_{\text{scat}}$  is generally a  
 10 monotonically increasing function of the particle size. However, there may be cases where  
 11 this is not true, typically due to the presence of Mie resonance (e.g., near  $D_p=1.5 \mu\text{m}$  in  
 12 Fig. 1). When the function is not monotonic, we apply a smoothing algorithm (Cerni, 1983;  
 13 Osborne et al., 2008) or merge together multiple bins (Pinnick et al., 1981; Walser et al.,  
 14 2017) and accept the tradeoff where we obtain a higher rate of correct bin assignment in  
 15 exchange for reduced bin resolution. This non-monotonicity is less of an issue as the

1 viewing angle becomes wider, as the larger range of angles will “smooth out” any Mie  
2 resonances (Figure S1). The wide viewing angle thus offers two key advantages: (1) the  
3 total signal (pulse height) is larger, making it easier to detect small particles using  
4 inexpensive electronics; (2) the calibration curve is less susceptible to small changes in  
5 particle scattering cross-section.

6

7 While an OPC sizes and counts individual particles, we generally are interested in  
8 evaluating the entire population of particles. To obtain the results for the entire  
9 population, we compute the scattering cross-section for each particle in the distribution,  
10 and assign it to a bin using the calibration curve generated previously – this results in a  
11 histogram with the total sum of particles in each discrete size bin over a period of time.  
12 Once we have the number distribution, we can compute the aerosol mass loading (PM)  
13 using Eq. (3):

14

$$15 \quad PM = \rho \sum_i N_i \frac{\pi}{6} D_{p,i}^3 \quad (\text{Equation 3})$$

16

17 where  $N_i$  is the number concentration for a given size bin,  $D_{p,i}$  is the geometric mean  
18 diameter for a given size bin, and  $\rho$  is the particle density, chosen to be constant. We can  
19 integrate mass loadings between different diameters by summing only across a sub-  
20 selection of bins (for example, if we intend to calculate the  $PM_{10}$  mass concentration, we  
21 would choose only the size bins corresponding to particles sized between 0-10  $\mu\text{m}$ ,  
22 whereas to calculate the  $PM_{2.5}$  mass concentration, we would use the bins corresponding  
23 to sizes between 0-2.5  $\mu\text{m}$ ). This approach for computing mass loadings is similar to that  
24 used by others (Di Antonio et al., 2018), though we use the geometric mean particle  
25 diameter as opposed to the mean particle diameter.

26

## 1 2.2.2 Integrating nephelometers

2 Nephelometers gather the light scattered by an aerosol population across a wide range  
3 of angles to gather as much of the scattered light as possible, while avoiding the near-  
4 forward and near-backward scattered light. Here, we define a nephelometer by the  
5 wavelength of its light source ( $\lambda$ ) and its viewing angle.

6

7 In practice, nephelometers are calibrated empirically by correlating the total scattered  
8 light signal to a reference mass measurement (Dacunto et al., 2015; Sousan et al., 2016b;  
9 Wang et al., 2015). Within our model, we do the same by computing the total scattered  
10 light signal using Mie theory and then take the ratio of the scattered light to a calculated  
11 mass loading. The total scattered light signal is calculated by integrating Eq. 2 across the  
12 entire particle size distribution, resulting in a single scattered light intensity for a given  
13 aerosol distribution. The calibration factor is then calculated by taking the ratio of this  
14 value and the mass loading of the aerosol distribution, which is calculated by integrating  
15 the volume distribution and multiplying by the particle density (Equation 3). Once we  
16 have computed the calibration factor, we can calculate the mass loading for any aerosol  
17 distribution by multiplying the calibration factor by the calculated total scattered light  
18 signal.

19

## 20 3. Results and discussion

21 We use the model described above to isolate the relative source of error associated with  
22 various differences in physical and optical properties of aerosols, as well as with the  
23 devices themselves. We include both simple, targeted experiments probing the effects  
24 of changes in isolated properties, as well as more complex, realistic experiments that  
25 attempt to mimic real-world scenarios. In the latter case, we include a variety of aerosol  
26 types in our model runs to resemble real-world use-cases; aerosol types include urban

1 aerosol, wildfire emissions, marine aerosol, dust, and continental background. The  
 2 physical and optical properties for these aerosols are summarized in Table 2. We discuss  
 3 these results in the context of three particle sensors chosen to be representative of low-  
 4 cost OPSs: a nephelometer, which uses a 658 nm light source and has a viewing range of  
 5 7°-173°, and two OPCs, both with 16 equally-spaced bins, a 658 nm light source, and a  
 6 viewing angle of 32-88°. The two OPCs differ only in the minimum particle size measured:  
 7 the 'low-cost OPC' is representative of commercial OPCs currently on the market and  
 8 measures particles in the 0.38-17.5  $\mu\text{m}$  size range; and the 'high-end OPC', representing  
 9 an idealized OPC that can measure much smaller particles, with a detection range of 0.1-  
 10 17.5  $\mu\text{m}$ . We note that many expensive OPCs cannot measure particles down to 100 nm;  
 11 this lower size cutoff was chosen as an approximate smallest particle size that an optical  
 12 sensor can detect.

13

14 **Table 2.** Aerosol optical and chemical properties used in this work.

Aerosol Type	Refractive index	Hygroscopicity parameter $\kappa^6$	Density ( $\text{g cm}^{-3}$ )
Urban <sup>1</sup>	1.525+0.020j	0.40	1.35
Background <sup>2</sup>	1.520+0.008j	0.25	1.45
Marine <sup>3</sup>	1.384+0.001j	1.10	2.16
Dust <sup>4</sup>	1.555+0.003j	0.03	2.60
Wildfire <sup>5</sup>	1.570+0.002j	0.10	1.58

15

16 <sup>1</sup> (Chen et al., 2019; Cheung et al., 2019; Hussein et al., 2004; Jurányi et al., 2013; Raut  
 17 and Chazette, 2007; Rissler et al., 2014; Shepherd et al., 2018; Wehner and  
 18 Wiedensohler, 2003)

19 <sup>2</sup> (Levoni et al., 1997; Wang et al., 2014; Yin et al., 2015)

20 <sup>3</sup> (Levoni et al., 1997; Ueda et al., 2016; Zieger et al., 2017)

21 <sup>4</sup> (Koehler et al., 2009; Petzold et al., 2009; Rocha-Lima et al., 2018)

22 <sup>5</sup> (Bougiatioti et al., 2016; Laing et al., 2016; McMeeking, 2004; Shepherd et al., 2018)

23 <sup>6</sup> (Petters and Kreidenweis, 2007)

24

1 We begin by investigating the impact that water uptake, driven by changes in the ambient  
2 relative humidity, has on the ability of all three OPSs to infer  $PM_{2.5}$  mass. Next, we explore  
3 the impact of aerosol optical properties (namely, the complex RI), followed by the impact  
4 that perturbations in the underlying particle size distribution can have in the OPS's ability  
5 to infer mass loadings. Finally, we summarize our results into general recommendations  
6 about each OPS type. Throughout, to provide a simple metric for the accuracy of OPS  
7 measurements, we present our results in terms of the ratio of the inferred or measured  
8  $PM_{2.5}$  mass concentration ( $M_m$ ) to the actual  $PM_{2.5}$  mass concentration ( $M_a$ ) at 0% relative  
9 humidity. An  $M_m/M_a$  ratio of greater than one implies we are overestimating the  $PM_{2.5}$   
10 loading, whereas a value less than one implies we are underestimating it.

11

### 12 3.1 Relative humidity and hygroscopic growth

13 One of the most widely discussed sources of error for OPS measurements is that caused  
14 by water uptake (Crilley et al., 2018; Di Antonio et al., 2018; Malings et al., 2020; Wang et  
15 al., 2015; Zheng et al., 2018). As relative humidity increases, hygroscopic particles (those  
16 with non-zero hygroscopic growth parameters,  $\kappa$ ) become larger as they take up water  
17 (Petters and Kreidenweis, 2007), leading to an increase in scattering caused by their  
18 increase in size. Additionally, water uptake changes the optical and chemical properties  
19 of the aerosol (e.g., RI, density, etc.), which can complicate any corrections. The EPA  
20 requires  $PM_{2.5}$  measurements to be made at relative humidities between 30-40% (Chow  
21 and Watson, 1998) to minimize the effects of hygroscopic growth on samples; however,  
22 since very few low-cost OPSs control for relative humidity (for example, with an in-line  
23 dryer), this can often lead to discrepancies when performing a calibration by co-location  
24 or when comparing results between instrument types.

25

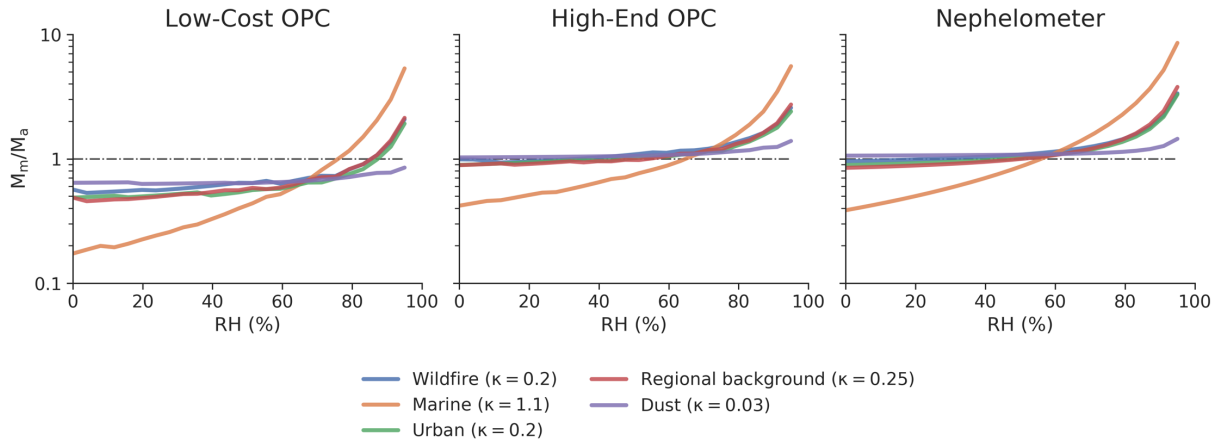
1 Figure 2 shows the impact that RH can have on the accuracy of an OPS. There is little  
2 effect until relative humidity reaches the deliquescence point of the aerosol, which  
3 depends on aerosol composition. At higher relative humidities, OPSs will tend to  
4 overestimate  $PM_{2.5}$  mass, especially for aerosols comprised of hygroscopic materials.  
5 When relative humidity approaches 95%, such overestimates in  $PM_{2.5}$  mass become  
6 exceedingly large: the OPCs observe a similar effect, with errors ranging from 100%-500%  
7 depending on the hygroscopicity of the aerosol. Nephelometers see a more pronounced  
8 effect with errors as high as 750% for extremely hygroscopic aerosols and 200%-300%  
9 errors for less hygroscopic aerosols.

10

11 The larger error of the nephelometer is caused in part by the fact that the inferred  $PM_{2.5}$   
12 mass is directly proportional to the total scattered light, which has no upper limit. For the  
13 OPCs, particles that take up significant water can be assigned to larger size bins and thus  
14 will not be integrated in the  $PM_{2.5}$  mass calculation. At moderate humidities (50%-80%),  
15 errors for both the nephelometers and OPCs can vary by as much as 20%-50%, which is  
16 in agreement other published experimental studies on the subject (Crilley et al., 2018; Di  
17 Antonio et al., 2018; Malings et al., 2020; Zheng et al., 2018). In addition to overestimating  
18 mass loadings at high relative humidity due to hygroscopic growth, the OPCs  
19 underestimate the mass loadings across all relative humidities. This is not caused by  
20 relative humidity or a lack of hygroscopic growth, but instead is a result of the “missing  
21 mass” below the detectable threshold of the OPC. The low-cost OPC, which cannot  
22 detect particles smaller than 380 nm, misses between 30%-90% of the mass, whereas the  
23 high-end OPC, which can detect particles larger than 100 nm, misses very little mass for  
24 most aerosol types. The only exception is the marine aerosol, which has a refractive index  
25 that is substantially different than the aerosol with which the instrument was calibrated.

26





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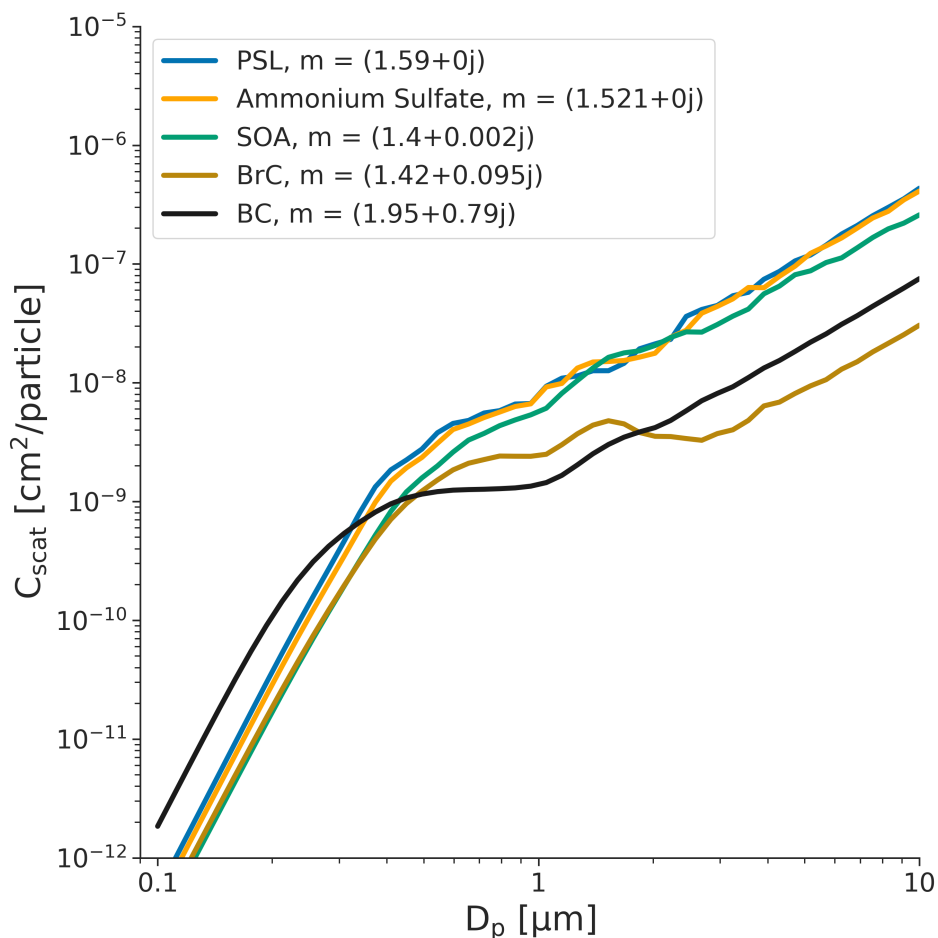
2 **Figure 2.** The accuracy in  $PM_{2.5}$  mass loading for a given particle sensor ( $M_m/M_a$ ) as a function of  
 3 relative humidity, for common aerosol types. All three particle sensors were calibrated with  
 4 ammonium sulfate (number-weighted Geometric Mean (GM) = 200 nm, Geometric Standard  
 5 Deviation (GSD) = 1.65). Details on the physical and optical properties of the various aerosols can  
 6 be found in Table 2.

7

### 8 3.2 Choice of calibration material and aerosol optical properties

9 OPCs are calibrated by correlating the scattering amplitude of known particle sizes for a  
 10 particles of a given composition (Gao et al., 2013). The relationship between scattering  
 11 amplitude and bin assignment (i.e., particle size) is heavily dependent on the aerosol's  
 12 complex refractive index (RI). Figure 3 shows the Mie scattering curve for a range of  
 13 common calibration materials, including both absorbing and non-absorbing materials.  
 14 For a given particle size, the RI of the particle can result in a range of scattered light  
 15 intensities ( $C_{scat}$ ) that vary by as much as an order of magnitude. This can have pronounced  
 16 effects on the calculated size (and hence mass) of a particle. In particular, the Mie curve  
 17 for black carbon (BC) is substantially different than those of non-absorbing materials, with  
 18 higher real (scattering) and imaginary (absorbing) components. As a result, for an OPC  
 19 calibrated with a non-absorbing material (such as PSLs), smaller BC particles (diameters  
 20 < 300 nm) will be overestimated in size (due to higher scattering), whereas larger BC

1 particles ( $> 300$  nm) will be underestimated (due to absorption) (Bohren and Huffman,  
2 1983).  
3  
4 Even small changes in the scattering (real) component of the RI of the calibration material  
5 can lead to particles being assigned to the incorrect bin: an RI higher than that of the  
6 calibration material will generally cause particles to be assigned to bins that are too large  
7 (overestimating in size and mass), and an RI lower than that of the calibration material will  
8 generally cause particles will be assigned to bins that are too small (underestimating in  
9 size and mass). This effect is examined in greater detail in the Supporting Information  
10 (Figs S2 and S3). Considering that bins are often hundreds of nm in width, the impact of  
11 such bin mis-assignment on reported mass can be large. For both OPCs and  
12 nephelometers, this will lead to large errors in inferred mass, though it can be more  
13 pronounced for OPCs, since the error for nephelometers is proportional to the increase  
14 in scattering and is not affected by the mis-assignment of individual particles to a  
15 particular size bin.

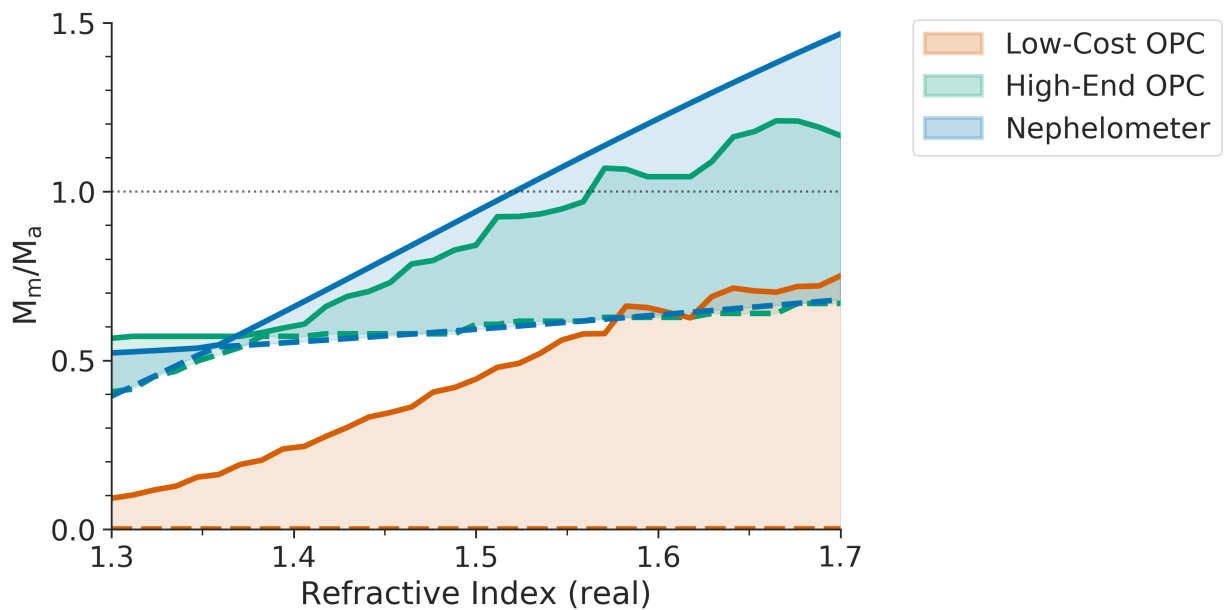


1  
 2 **Figure 3.** Mie curves (integrated over a viewing angle of 32°-88°) for a select group of common  
 3 calibration materials. Materials shown include polystyrene latex spheres (PSLs), ammonium  
 4 sulfate, secondary organic aerosol (SOA), and black carbon (BC). Small differences in the  
 5 refractive index of a measured material can lead to drastic bin mis-assignment (see SI), depending  
 6 on where bin boundaries are set at the time of calibration.

7  
 8 The effect of differences in refractive index on inferred PM<sub>2.5</sub> mass measurements is shown  
 9 in Fig. 4. Results are shown for a single aerosol distribution, in which the only parameter  
 10 allowed to vary is the RI. The real component of the refractive index is shown on the x  
 11 axis, with the upper and lower bounds being determined by the imaginary part of the  
 12 refractive index; the imaginary component ranges from 0 (non-absorbing) to 0.79 (black  
 13 carbon). The nephelometer (blue swatch in Fig. 4) is calibrated using ammonium sulfate  
 14 ( $m = 1.521 + 0j$ ). When the nephelometer is evaluated at this exact RI (and a constant size

1 distribution), it measures mass accurately ( $M_m/M_a = 1$ ). However, if the real component of  
 2 the aerosol being evaluated is higher than that of the calibration standard, the total  
 3 scattering is greater, resulting in the inferred  $PM_{2.5}$  mass being larger than the actual  $PM_{2.5}$   
 4 mass ( $M_m/M_a > 1$ ). Similarly, as the absorbing component becomes larger, less of the  
 5 incoming light is scattered, resulting in a substantial underestimation of the mass loading.

6  
7



8  
9 **Figure 4.** The accuracy of OPSs as a function of the refractive index of the aerosol being  
 10 measured. The real component of the RI is on the x axis, and the width of each swatch bounded  
 11 by the absorption/imaginary component, which spans from 0 (non-absorbing, solid line) to 0.79  
 12 (black carbon, dashed line). Results are shown for a nephelometer (blue), and the two OPCs  
 13 (orange and green). All results are for a generic particle size distribution with number-weighted  
 14  $GM=200$  nm and  $GSD=1.65$  and the OPCs were calibrated with PSLs whereas the nephelometer  
 15 was calibrated with ammonium sulfate.

16  
17 Also shown are the results for two OPCs. The high-end OPC (green) is sensitive to  
 18 particles as small as 100 nm, whereas the low-cost OPC (red) is sensitive to particles as  
 19 small as 380 nm. As the absorbing component of the refractive index becomes larger, the  
 20 scattering amplitude across the entire distribution is too small for the OPC to detect,

1 resulting in a mass reading of zero. Both OPCs exhibit this effect, but for the high-end  
2 OPC, fewer particles will fall below the size cutoff of the OPC than for the low-cost OPC,  
3 resulting in a less dramatic underestimation of the mass. Most commercially available  
4 OPCs are more similar to the low-cost OPC, with lower limits of detection of around 500  
5 nm. If operating in an environment in which the aerosol is strongly absorbing, large  
6 underestimates in  $PM_{2.5}$  should be expected. Even under conditions where the aerosol is  
7 not absorbing, the low-cost OPC largely underestimates the mass due to its high  
8 minimum size cutoff. For nephelometers, the errors are not as drastic, but still do depend  
9 strongly on the RI of the calibration aerosol used.

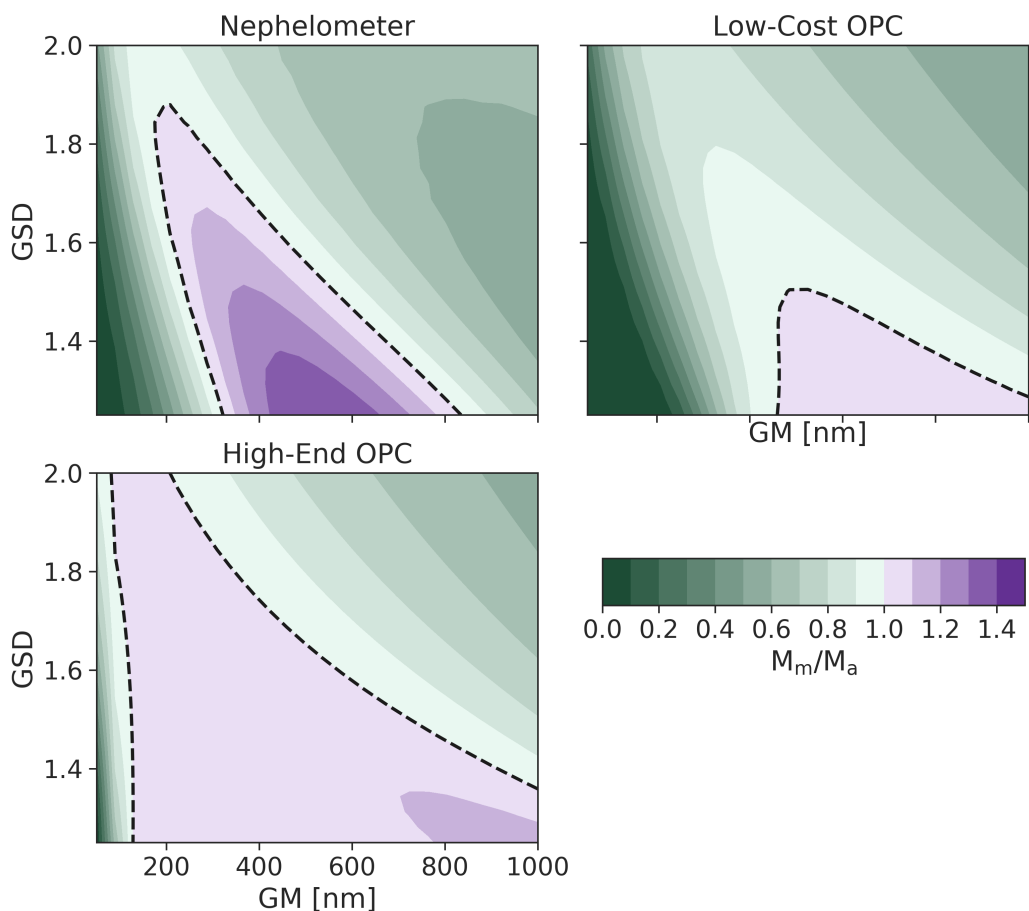
10

### 11 3.3 Changes in the Particle Size Distribution (PSD)

12 The ability of optical particle sensors to adapt to perturbations in the underlying particle  
13 size distribution (PSD) is important because PSDs can be highly variable over short  
14 periods of time, especially in urban areas with varying contributions from various local  
15 sources. Fig. 5 shows the accuracy of all three OPSs as the function of the PSD of the  
16 particles being measured. These calculations assume a single lognormal mode with all  
17 other properties of the aerosols (density, refractive index, and hygroscopicity) held  
18 constant. For the purpose of the model, the OPCs were calibrated using PSLs at each bin  
19 boundary, and the nephelometer was calibrated using ammonium sulfate ( $N=1e4 \text{ cm}^{-3}$ ,  
20  $GM=400 \text{ nm}$ , and  $GSD=1.65$ ). The entire population of ammonium sulfate particles is  
21 then evaluated while varying the number-weighted mean particle diameter ( $GM$ ) and the  
22 width of the distribution ( $GSD$ ). For each PSD, we compute the relative accuracy of each  
23 device and plot the results in Fig. 5, in which the color and contours correspond to the  
24  $M_m/M_a$  metric.

25

1 The nephelometer substantially underestimates the mass concentration (by 50%-70%) for  
2 most PSDs, since it is calibrated to a single PSD. As the PSD changes, the ratio of total  
3 scattered light to integrated mass changes, causing the accuracy to change as well. (The  
4 accuracy depends somewhat on viewing angle, as shown in Figure S4, but errors are still  
5 substantial across a wide range of PSDs.) OPCs are potentially better since they measure  
6 the size of the particles and can theoretically account for changes in the PSD; however,  
7 they are still subject to errors given their limitations in detected size range. In particular,  
8 the low-cost OPC considerably underestimates the mass (by 60%-90%) for most PSDs as  
9 the bulk of the mass is below the detectable size limit of the OPC. As the geometric mean  
10 diameter increases in size, or the width of the distribution becomes larger, a larger  
11 fraction of the particles enters the detectable range, slightly improving the results for the  
12 low-cost OPC. The high-end OPC is most able to adapt to the changes in the PSD due  
13 to its significantly smaller  $D_{\min}$  (100 nm); there is roughly a 20% difference across the entire  
14 range of PSDs shown. Unlike the low-cost OPC, a majority of the mass falls within the  
15 detectable range of the high-end OPC, resulting in little to no effect of changes to the  
16 PSD on accuracy of the mass concentration measurement.



1  
 2 **Figure 5.** Mass concentration accuracy ( $M_m/M_a$ ) of OPSs for a range of particle size distributions  
 3 (PSDs). Accuracy is shown for all combinations of PSDs with number-weighted geometric mean  
 4 diameters (GMs) between 100 - 1000 nm and geometric standard deviations (GSDs) between 1.2  
 5 - 2.0. Perturbations in the PSD can lead to large errors for nephelometers and optical particle  
 6 counters with high minimum particle size cutoffs. All results are shown for ammonium sulfate  
 7 particles; the OPCs were calibrated with PSLs and the nephelometer was calibrated with  
 8 ammonium sulfate ( $N=1e4\text{ cm}^{-3}$ ,  $GM=400\text{ nm}$ , and  $GSD=1.5$ ). A black dashed line indicates the  
 9 "1:1" line, where  $M_m/M_a = 1$ .

10  
 11 While previous work has highlighted the importance of the varying PSD and its effect on  
 12 making accurate mass measurements with OPSs (Di Antonio et al., 2018; Gao et al., 2013;  
 13 Malings et al., 2020), the effect of 'missing mass' – the mass below the lowest size bin of  
 14 an OPC – has received relatively little attention. The standard way to treat this missing  
 15 mass is to empirically correct via regression analysis (Dacunto et al., 2015; Malings et al.,

1 2020). While this can mitigate absolute errors, it requires the assumption that the PSD is  
2 constant in shape, varying only in magnitude. With particle loadings that are mostly below  
3 10's of  $\mu\text{g m}^{-3}$  throughout the United States, this assumption is unlikely to be a large  
4 source of absolute error. However, if the same approach were used in highly polluted  
5 environments where sub-300 nm aerosol loadings can easily reach hundreds of  $\mu\text{g m}^{-3}$   
6 (Bhandari et al., 2020; Gani et al., 2019, 2020), changes in the PSD are likely to lead to  
7 large errors (in both an absolute and relative sense) in mass loading measurements.  
8 Overall, nephelometers and OPCs with high minimum size cutoffs are prone to  
9 substantial uncertainties as the underlying PSD changes, whereas for OPCs with low  
10 minimum size cutoffs this effect is relatively minor.

11

#### 12 4. Implications and future work

13 In this work, we have laid out a framework for understanding the sensitivity of low-cost  
14 optical particle sensors to the various physical and optical properties of aerosols. We  
15 described a new Mie theory-based software package (*opcsim*) for modeling the response  
16 of OPSs to various aerosols and demonstrated its use for better understanding the  
17 strengths and limitations of various low-cost particle sensors. We also used the model to  
18 investigate how various potential pitfalls (e.g., changes to environmental conditions,  
19 mismatches between calibration particles and particles being measured) may contribute  
20 to errors in mass concentration measurements. A summary of these results is given in  
21 Table 3.

22

23

24



1 **Table 3.** Effects of changing environmental/aerosol parameters on the relative error in  
 2 measured mass loading by different OPS types.

Parameter changed	OPS type		
	Low-Cost OPC	High-End OPC	Nephelometer
RH & Hygroscopicity (Figure 2)	Very high for (20-200%) for hygroscopic materials when RH > ~75%		
Optical Properties (Figure 4) <sup>1</sup>	High (30 – 100%)	Medium (20 – 60%)	Medium (20 – 75%)
Particle Size Distribution (Figure 5)	High 60 - 90%	Low < 20%	High 50 – 70%

3  
 4 <sup>1</sup> Primarily a source of error when an OPS calibrated with non-absorbing particles measures absorbing  
 5 particles (or vice versa).

6  
 7 Consistent with previous studies, our results suggest that relative humidity is a large  
 8 source of uncertainty for all OPSs when the aerosol is hygroscopic and relative humidities  
 9 are above the deliquescence point, typically around 75%; additionally, the error  
 10 introduced by relative humidity is highly sensitive to the aerosols’ affinity for water. This  
 11 is correctable, at least to first order, limiting the impact of RH error on final results (Crilley  
 12 et al., 2018; Di Antonio et al., 2018; Malings et al., 2020). Differences in aerosol optical  
 13 properties can also lead to substantial errors in mass loadings for low-cost OPCs, with  
 14 smaller but still substantial errors for high-end OPCs and nephelometers. This is  
 15 especially important when the aerosol is strongly absorbing, as the amount of scattered  
 16 light can make small particles undetectable with inexpensive optical detectors. If it were  
 17 possible to measure some proxy for aerosol composition in real time, it would be possible  
 18 to vastly reduce this error and improve the accuracy of mass measurements using OPSs  
 19 for real-time data collection. Additionally, this work highlights the importance of the  
 20 underlying particle size distribution: differences in size distributions can lead to large

1 errors in mass measurements from low-cost OPCs and nephelometers, while being only  
2 a small source of error for high-end OPCs that can properly count and size particles at  
3 low sizes. The ability of a given OPC to measure small particles is particularly important,  
4 with marginal improvements in minimum detected size leading to large gains in ability to  
5 accurately infer mass. Finally, the choice of calibrant is shown to be extremely important  
6 for both nephelometers and OPCs. Ensuring that OPSs are calibrated intelligently (i.e.,  
7 using particles similar to the aerosol to be detected) can lead to significant improvements  
8 in expected performance.

9  
10 Table 4 summarizes these results within the context of measurements of representative  
11 real-world aerosol types. It provides an overview of the potential errors associated with  
12 different types of optical particle sensors under various scenarios, with recommendations  
13 for the type of calibration particles that would minimize errors in PM<sub>2.5</sub> mass  
14 measurements. Generally, in environments where small particles (< 300 nm) comprise a  
15 large percentage of the total mass, low-cost OPCs will be subject to considerable error.  
16 This will also be the case in environments with substantial levels of light-absorbing  
17 aerosol, such as wildfires or soot-heavy environments. (Sensor calibration using  
18 absorbing particles could help mitigate this effect, though this would introduce new  
19 errors when measuring non-absorbing aerosol.) In environments in which the underlying  
20 aerosol size distribution is variable (especially on short (sub-hour) timescales), such as  
21 urban environments or evolving wildfire plumes, nephelometers and low-cost OPCs will  
22 struggle to keep up with the changes in the relationship between the total scattered light  
23 and mass loading, leading to large variance in mass estimates.

24  
25 The estimates and recommendations given in Table 4 are not intended to be  
26 comprehensive, but rather serve as a starting point for characterizing the strengths and

1 limitations of low-cost OPSs using Mie theory (and specifically the *opcsim* software  
2 package). Additional *opcsim* simulations carried out across a range of sensor designs,  
3 calibrant particles, and measured particle types could provide more comprehensive and  
4 quantitative estimates of errors in measured particle sizes and mass loadings, including  
5 for individual sensors and individual use-cases. Future improvements to *opcsim* could be  
6 made to allow for the simulation of more complex aerosols (e.g., externally-mixed  
7 populations, other particle morphologies) or the inclusion of more complex bin-  
8 assignment algorithms; comparison with laboratory studies (in which  $M_m/M_a$  is measured  
9 rather than just estimated) would also be useful. Additionally, co-located data with size-  
10 resolved measurements would allow for improved validation of the OPC component of  
11 this model. It is hoped that the Mie-theory-based approach described here will lead to  
12 an improved understanding of the errors associated with low-cost optical PM  
13 measurements, insight into calibration techniques that minimize such errors, and  
14 ultimately guidance into the design of new PM sensors for improved low-cost  
15 measurements of air quality and human exposure.

16

17

**Table 4.** Summary of expected performance and recommendations for calibration materials for use of low-cost optical particle sensors to measure the particulate mass loadings of different aerosol types.

AEROSOL TYPE	AEROSOL PROPERTIES	SUGGESTED CALIBRANT	SENSOR PERFORMANCE BY OPS TYPE <sup>1</sup>		
			Low-Cost OPC	High-End OPC	Nephelometer
FOSSIL-FUEL COMBUSTION	Very small GMD, mostly non-hygroscopic, moderate absorbing RI <sup>2</sup>	Calibrate with aerosols closer in RI, such as from combustion sources	Will perform poorly due to the small GMD and absorption component of the aerosol	Will perform moderately well though will miss ultrafine particles	Can perform moderately well if calibrated using appropriate materials
WILDFIRE	Varying PSD, moderate absorbing component of RI <sup>3</sup>	Calibrate with aerosol of similar optical properties and GMD (ideally biomass smoke)	Will likely undersize and underestimate mass due to the absorbing component of the aerosol; the GMD will change with proximity to the source leading to changes in accuracy	Will perform moderately well, though may mis-size the particles as the properties of the aerosol change as the plume evolves	Can perform well under certain circumstances; moderate error should be expected as the GMD of the wildfire plume evolves
URBAN	Varying GMD, moderate hygroscopicity <sup>4</sup>	Calibrate with NIST urban aerosol or collected urban aerosol	Performance depends on uniformity of sources; large errors will occur as aerosol source (and GMD) changes	Will perform moderately well to well, though will miss ultrafine particles	Will perform moderately well if averaged over a long period of time to normalize the GMD
DUST	Large GMD, non-hygroscopic <sup>5</sup>	Calibrate with Arizona Road dust or collected dust	Likely to perform well, given the large particle sizes	Likely to perform well, given the large particle sizes	Likely to perform well

<sup>1</sup> Based on properties of the aerosol only, and not external environmental parameters (i.e., RH).

<sup>2</sup> (Bond et al., 2002; Bond and Bergstrom, 2006; Chang et al., 2004)

<sup>3</sup> (Laing et al., 2016; McMeeking, 2004)

<sup>4</sup> (Hussein et al., 2004; Wehner and Wiedensohler, 2003)

<sup>5</sup> (Petzold et al., 2009; Rocha-Lima et al., 2018)

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10 package.

11

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