

1 **Assimilation of lidar planetary boundary layer height**  
2 **observations.**

3 **Andrew Tangborn<sup>1</sup>, Belay Demoz<sup>1,2</sup>, Brian J. Carroll<sup>2</sup>, Joseph Santanello<sup>3</sup> and**  
4 **Jeffrey L. Anderson<sup>4</sup>**

5 <sup>1</sup>Joint Center for Earth Systems Technology, University of Maryland Baltimore County, Baltimore, MD,  
6 USA

7 <sup>2</sup>Dept. of Physics, University of Maryland Baltimore County, Baltimore, MD, USA

8 <sup>3</sup>Hydrological Sciences Laboratory, NASA Goddard Space Flight Center, Greenbelt, MD, USA

9 <sup>4</sup>National Center for Atmospheric Research, Boulder, CO, USA

## Abstract

Lidar backscatter and wind retrievals of the planetary boundary layer height (PBLH) are assimilated into 22 hourly forecasts from the NASA Unified - Weather and Research Forecast (NU-WRF) model during the Plains Elevated Convection at Night (PECAN) campaign on July 11, 2015 in Greensburg, Kansas, using error statistics collected from the model profiles to compute the necessary covariance matrices. Two separate forecast runs using different PBL physics schemes were employed, and comparisons with 6 independent radiosonde profiles were made for each run. Both of the forecast runs accurately predicted the PBLH and the state variable profiles within the planetary boundary layer during the early morning, and the assimilation had a small impact during this time. In the late afternoon, the forecast runs showed decreased accuracy as the convective boundary layer developed. However, assimilation of the Doppler lidar PBLH observations were found to improve the temperature and V velocity profiles relative to independent radiosonde profiles. Water vapor was overcorrected, leading to increased differences with independent data. Errors in the U velocity were made slightly larger. The computed forecast error covariances between the PBLH and state variables were found to rise in the late afternoon, leading to the larger improvements in the afternoon. This work represents the first effort to assimilate PBLH into forecast states using ensemble methods.

## 1 Introduction

The planetary boundary layer (PBL) plays an important role in weather, climate and pollution through its role in land-atmosphere interactions and mediation of Earth's water and energy cycles (Santanello et al. 2018). This layer is where the Earth's surface interacts with the atmosphere, exchanging momentum, heat, moisture and pollutants. The PBL height (PBLH) is central to these interactions and is controlled by the energy flux from the surface. Under certain conditions during daytime it defines the convective boundary layer (CBL) and during nighttime it is the stable (non-convective) boundary layer (SBL). Trace gases and aerosols emitted from the surface are rapidly transported within the CBL by turbulent atmospheric motion, and transfer of energy and mass into the free troposphere occurs across an interfacial layer at the top of the PBL. The PBL affects convection in the troposphere, which is generally initiated within the boundary layer and then penetrates its top (Hong and Pan, 1998; Browning, et al. 2007). Thus,

42 accurate knowledge of the PBLH is essential for both weather, pollution and climate fore-  
43 casting.

44 The PBLH is defined by thermodynamic properties such as a temperature inver-  
45 sion or hydrolapse which can be measured by radiosonde. Alternatively, the drop off in  
46 aerosol concentration that occurs across the top of the PBL is used, since aerosols are  
47 well mixed throughout the PBL when the CBL is present (Hicks, et al., 2019). Atmo-  
48 spheric models rely on parameterization schemes to define the structure of the PBL and  
49 compute PBLH. These are generally either local mixing schemes that use local turbu-  
50 lent kinetic energy (TKE, Janjic, 1994) or non-local flux schemes (Hong and Pan, 1996).  
51 Generally, these PBL parameterizations have systematically higher PBLH relative to ob-  
52 served values (Hegarty et al., 2018), and also have difficulties modeling the growth of the  
53 convective layer during the morning. The variety of definitions of PBLH make it diffi-  
54 cult to effectively evaluate existing models or develop new ones.

55 Observations of PBLH are traditionally made by radiosonde measurements, which  
56 have high vertical resolution but are expensive to launch frequently and are thus lim-  
57 ited to special experiments and/or ill-timed launches (*e.g.* 00/12 UTC National Weather  
58 Service launches) with respect to convective and stable PBL development. Likewise, space-  
59 borne measurements of the lower troposphere from passive and active instruments are  
60 severely limited in vertical, spatial, and/or temporal resolution (Wulfmeyer et al. 2015).  
61 Ground based measurement of PBLH has been proposed for an extensive network of ceilome-  
62 ters by adding to the functionality of instruments that were designed for measuring cloud  
63 heights (Hicks et al., 2016). The ceilometer measures the time required for a laser pulse  
64 to return to a receiver, from which the height of the scattering is determined. The in-  
65 tensity of the backscatter is correlated with the density of aerosols at a given height and  
66 the PBLH is inferred from the location of the maximum negative gradient of the backscat-  
67 ter intensity. Several algorithms employ wavelet transforms to identify the location of  
68 the negative gradient (*e.g.* Brooks, 2003; Knepp, *et al.*, 2017). This existing network of  
69 ceilometers could be used to create a relatively dense network of frequent PBLH obser-  
70 vations, as was recommended by the 2009 study from the National Research Council (NRC,  
71 2009) and the Thermodynamic Profiling Technologies Workshop (NCAR, 2012).

72 Since the ceilometer PBLH observations were not yet available for the time period  
73 we are studying, we employ Doppler lidar observations made at the Plains Elevated Con-

74 vection at Night (PECAN) site in Greensburg, Kansas, to demonstrate the methodol-  
75 ogy. PECAN was an intensive campaign to study organized Mesoscale convection sys-  
76 tems (MCSs) during the period June 1-July 15, 2015. It employed three aircraft and a  
77 large array of ground based lidar, radar and ground weather stations. The data we are  
78 using is from a Leosphere WINDCUBE-200S Doppler lidar owned and operated by the  
79 University of Maryland, Baltimore County (Delgado et al., 2016). This lidar operates  
80 at an infrared wavelength, and hence receives its strongest backscattered signal within  
81 the aerosol-laden PBL and is often below the measurement noise floor above the PBL.  
82 The Doppler shift of the backscattered signal is used to calculate wind speed as a func-  
83 tion of range, which can then be used to produce a multitude of wind and turbulence  
84 variables useful for PBL characterization (e.g. vertical velocity variance and signal-to-  
85 noise ratio variance). While Doppler lidars and ceilometers are similar in aerosol detec-  
86 tion, a Doppler lidar’s additional wind measurement capability makes it more broadly  
87 applicable and at times more accurate than a ceilometer for PBLH retrievals. The PBLH  
88 algorithm applied for this study combines several such aerosol and wind variables and  
89 each PBLH retrieval involves measurement of turbulence intensity, horizontal wind pro-  
90 files and backscatter intensity. The heights of steep gradients in these quantities are de-  
91 termined using empirical thresholds and wavelet transform techniques, and the three es-  
92 timates are combined using fuzzy logic. This is described at length in Bonin et al. (2018).  
93 Additional lidar parameters and the application of the algorithm to PECAN data were  
94 presented in Carroll et al. (2019). The PBLH retrievals were made from a repeating 25-  
95 minute lidar scan cycle. This Doppler lidar and PBLH algorithm combination are gen-  
96 erally well-suited for accurate and precise measurement of the PBLH during the daytime  
97 boundary layer, nocturnal boundary layer, and morning transition period (Bonin et al.  
98 2018, Carroll et al. 2019). The evening transition is the most challenging for this setup  
99 due to due to difficulties in defining a clear mixing layer during the decay of a turbulent  
100 daytime PBL (Lothon et al. 2014).

101 The question remaining is how to assimilate these observations into a numerical  
102 weather prediction (NWP) model. A number of studies have explored assimilating bound-  
103 ary layer wind profile measurements from lidar (Hu et al. 2019, Coniglio et al. 2019, Degelia  
104 et al. 2019) and have shown that this increases the accuracy of forecasts due to improve-  
105 ments within the PBL. And further studies (Degelia et al. 2020; Chipilski et al. 2020)  
106 found that convective initiation (CI) was enhanced through the assimilation of thermo-

107 dynamic profiles within the PBL, though the former found that CI was degraded by the  
108 assimilation of kinematic (velocity) profiles. This work highlights the important role that  
109 the PBL plays in forecasting convective events, so that any observations that can im-  
110 prove estimation of the model state should be an important source of new information.  
111 We are interested assimilating the PBLH observations directly because the ceilometer  
112 network described above will focus on these retrievals, and satellite missions which mea-  
113 sure PBLH are also planned. PBLH is a diagnostic variable in NWP parameterized physics  
114 models. This means any correction to PBLH will be lost during the model forecast un-  
115 less the PBLH height observation is used to correct state variables such as temperature  
116 and moisture. This could be done either by adopting a variational data assimilation scheme,  
117 or through the use of an ensemble Kalman filter which would determine the error covari-  
118 ances between PBLH and state variables in the model. We choose the latter so as to avoid  
119 the task of linearizing the model physics. The structure of the covariance, and how the  
120 state variables are changed by assimilating PBLH, will depend on which PBL scheme  
121 is used. We will show how such a system could work by conducting a posteriori lidar PBLH  
122 observation impact experiments using forecast fields from a NASA Unified - Weather and  
123 Research Forecast (NU-WRF, Lidard-Peters, 2015) model runs for one day during the  
124 Plains Elevated Convection at Night (PECAN) campaign on July 11, 2015. The assim-  
125 ilation is done on 22 hourly WRF forecast fields throughout the day without cycling the  
126 analysis fields back into the model, using two different PBL parameterizations. In this  
127 paper, we demonstrate a new and promising method that uses the lidar-based aerosol  
128 backscatter and wind derived PBLH to correct model forecasted state variables. The pur-  
129 pose here is to show how ensemble computed error covariance can transfer observational  
130 information from PBLH to the state variable profiles.

## 131 **2 Methodology**

132 The assimilation methodology is based on the ensemble Kalman filter (EnKF)(Evenensen,  
133 1994; Burgers, et al. 1998; Evensen, 2009), where the analysis state is the estimate with  
134 a minimized error norm, relative to the given error statistics. It differs from the EnKF  
135 in that the analysis is not used as an initial state for the next model forecast. Rather,  
136 two existing one day NU-WRF forecasts, with different PBL physics schemes, are used  
137 when lidar measurements are available at a single location. These forecasts were produced  
138 as a part of the PECAN campaign in 2015, and we reuse them here to demonstrate the

139 assimilation algorithm that we have developed. These were not ensemble forecasts so we  
 140 cannot build a standard ensemble Kalman filter from them. Instead we use Ensemble  
 141 Optimal Interpolation (EnOI), in which profiles from neighboring model gridpoints are  
 142 used to obtain an estimate of error statistics (Oke, *et al.*, Evensen, 2003; 2010; Keppenne,  
 143 *et al.*, 2014). This approach will allow for the construction of the vertical component of  
 144 covariance, which is needed in order to understand how PBLH can be used to correct  
 145 atmospheric profiles through the use of profile and PBLH statistics. We use profiles from  
 146 nearby model grid points and have tested the system with varying numbers of grid points  
 147 in the ensemble. An ensemble Kalman filter would likely give different covariance infor-  
 148 mation, but the basic relationship between the state variable profiles and the PBLH are  
 149 determined by the model in the same manner here.

150 The NU-WRF simulations, taken from existing forecast runs used for the PECAN  
 151 campaign (Santanello *et al.*, 2019) are initialized using a National Center for Environ-  
 152 mental Prediction (NCEP) Global Forecast System (GFS) reanalysis. The two NU-WRF  
 153 simulations use the Mellor–Yamada–Janjic (MYJ)[Mellor and Yamada, 1974, 1982; Jan-  
 154 jic, 2002] and Mellor-Yamada-Nakanishi-Niino level 2.5 (MYNN) [Nakanishi and Niino,  
 155 2009] which are local 1.5 and 2.5 order turbulence closure schemes respectively. The PBLH  
 156 in each of these models is estimated using the turbulent kinetic energy (TKE) method.  
 157 The NU-WRF forecast state variables are temperature (T), specific humidity (Q) and  
 158 velocity (U,V), and we define the forecast vector  $\mathbf{x}^f = [T^f \ Q^f \ U^f \ V^f \ (PBLH)^f ]$ , where  
 159 we have combined PBLH with the state variables to enable the covariance calculation  
 160 between them. The vector  $\mathbf{x}$  is a column vector, so that the error covariance defined be-  
 161 low only includes vertical covariances. The forecast runs are initiated from the NOAA  
 162 global forecast system (GFS) reanalysis interpolated to the local domain of 30-48N and  
 163 84-110 W, with 220×220 lat/lon and 54 vertical levels, at 0 UTC. At this time, the ini-  
 164 tial state has assimilated all of the conventional and satellite observations globally. The two  
 165 WRF forecast experiments start at 0 UTC, and are run for 22 and 23 hours for the MYJ  
 166 and MYNN experiments, respectively. We use an ensemble of the 20×20 nearest grid-  
 167 points, so that all of the ensemble members are within about 30 km of the lidar obser-  
 168 vations (since the grid spacing is about 3 km). Generally, larger ensembles using grid-  
 169 points farther away will result in larger forecast error covariance because the geographic  
 170 variability. So this ensemble size was chosen as a balance between ensemble size and ge-  
 171 ographic localization. The forecast standard deviation for PBLH on the chosen ensem-

172 ble was around 27 m at 22 UTC. Lidar PBLH observations were made every 25 minutes  
 173 on that day in Greensburg, KS (37.6 N, 99.3 W), while balloon soundings were launched  
 174 from that location 6 times as part of the Plains Elevated Convection At Night (PECAN;  
 175 Gerts et al. 2017).

176 For an EnKF the generalized analysis equations are:

$$\mathbf{x}^a = \mathbf{x}^f + \mathbf{K}(\mathbf{y}^o - H(\mathbf{x}^f)) \quad (1)$$

177 where  $\mathbf{x}^a$  is the analysis state,  $\mathbf{x}^f$  is the forecast state,  $\mathbf{y}^o$  is the observation vector and  
 178  $H$  is the non-linear observation operator. The gain matrix,  $\mathbf{K}$  is defined by:

$$\mathbf{K} = \mathbf{P}^f \mathbf{H}^T (\mathbf{H} \mathbf{P}^f \mathbf{H}^T + \mathbf{R})^{-1}, \quad (2)$$

179 and  $\mathbf{P}^f$  is the forecast error covariance,  $\mathbf{R}$  is the observation error covariance and  $\mathbf{H}$  is  
 180 the linearized observation operator. The matrices  $\mathbf{P}^f \mathbf{H}^T$  and  $\mathbf{H} \mathbf{P}^f \mathbf{H}^T$  are formed from  
 181 the ensemble of forecasts. In the present work, we use the EnOI method, and assimilate  
 182 observations one at a time using the the ensemble of profiles described above. In this case,  
 183  $\mathbf{x}^a$  and  $\mathbf{x}^f$  depend only on vertical level, and  $\mathbf{y}^o = y^o$ ,  $\mathbf{R} = (\sigma^o)^2$  and  $\mathbf{H} \mathbf{P}^f \mathbf{H}^T =$   
 184  $(\sigma^f)^2$  become scalar quantities. The analysis equations are then

$$\mathbf{x}^a = \mathbf{x}^f + \mathbf{K}(y^o - H(\mathbf{x}^f)) \quad (3)$$

185 and

$$\mathbf{K} = \mathbf{P}^f \mathbf{H}^T ((\sigma^f)^2 + (\sigma^o)^2)^{-1}, \quad (4)$$

186 The observation error standard deviation supplied by the lidar retrieval is  $\sigma^o$ , which is  
 187 determined from the combined uncertainty of the vertical velocity variance, velocity gra-  
 188 dient and backscatter gradient. Generally, when these quantities change rapidly at the  
 189 top of the PBL, then the estimated error is small. The error estimates are larger when  
 190 (during the evening), the gradients are much more gradual.  $\mathbf{H}$  is the linearized obser-  
 191 vation operator for PBLH. Because the PBLH is related to the state variables via the  
 192 two PBL physics schemes, determining  $\mathbf{H}$  would require linearizing the PBL physics at  
 193 every analysis time. Rather, here we use the EnOI described above to get:

$$\mathbf{P}^f \mathbf{H}^T \approx \langle (\mathbf{x}^f - \mu_{\mathbf{x}^f}) (H(\mathbf{x}^f - \mu_{\mathbf{x}^f}))^T \rangle \quad (5)$$

194 and

$$\mathbf{H} \mathbf{P}^f \mathbf{H}^T = (\sigma^f)^2 \approx \langle H(\mathbf{x}^f - \mu_{\mathbf{x}^f}) (H(\mathbf{x}^f - \mu_{\mathbf{x}^f}))^T \rangle \quad (6)$$

195 where  $\mu_{\mathbf{x}}^f$  is the mean forecast state of the ensemble of profiles. See Houtekamer and  
 196 Zhang (2016) for a review of ensemble Kalman filter techniques.

197 We expect the correlation between the air mass within the PBL and the free tro-  
 198 posphere to drop away rapidly, because of limited interactions between them. We found  
 199 that this can cause errors in the analysis profiles if error covariance between the state  
 200 variables and PBLH is allowed to continue into the troposphere. To reduce these errors  
 201 we have added an exponential decay starting at the model level closest to the PBLH ( $k_{PBLH}$ )  
 202 to define a vertical localization factor:

$$C_{loc} = exp \left[ -\alpha \left( \frac{k - k_{PBLH}}{k_{PBLH}} \right)^2 \right] \quad (7)$$

203 where  $k$  is the model level and  $\alpha = 8$  is an experimentally determined factor. The fac-  
 204 tor  $C_{loc}$  is multiplied by the vertical covariance in (5) to ensure that the covariance be-  
 205 tween the PBLH and the state variables becomes small within a couple of model levels  
 206 into the free troposphere.

207 Equations 3-4 are solved at each hour using the nearest lidar profile observation  
 208 in time, and the resulting analysis fields are compared to radiosonde profiles when the  
 209 latter are also available. There are 22 or 23 analyses (for each forecast run), and 6 times  
 210 where comparison with radiosonde profiles are made. We focus on the impact of the as-  
 211 simulation on the state variables T, Q, U and V rather than the PBLH because only the  
 212 state variables would be retained by a forecast.

### 213 **3 Results**

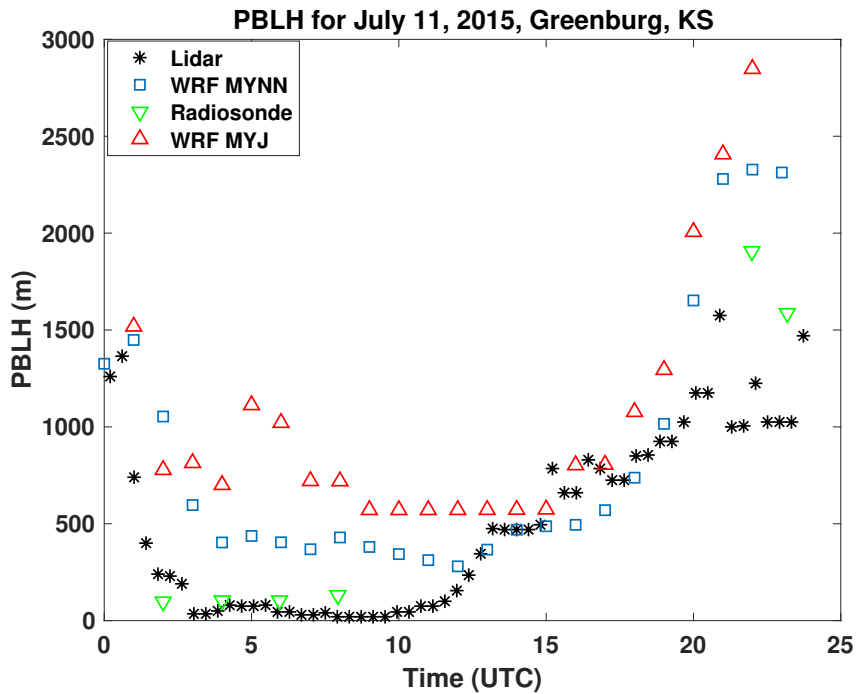
214 This section describes the NU-WRF simulation results, the assimilation of PBLH  
 215 into these forecasts, and the relationship between the assimilation impact and the time  
 216 varying forecast and observation error covariances.

#### 217 **3.1 NU-WRF simulations**

218 The one day NU-WRF simulations are presented in this section. Figure 1 shows  
 219 the PBLH during that day, derived from the two NU-WRF forecasts, lidar observations  
 220 and soundings. We have determined the sounding PBLH using the parcel method (Holz-  
 221 worth, 1964), which defines the top as the height where the potential temperature first  
 222 exceeds the ground temperature. The lidar PBLH (black \*, derived using the method



223 reported in Bonin, 2018) closely matches the radiosonde estimates (green triangles) in  
 224 the late evening to nighttime (2-7 UTC), while it is somewhat lower late afternoon to  
 225 early evening (18-24 UTC). The two NU-WRF forecasts differ from the observations de-  
 226 pending on the time of day. During nighttime and early morning the MYJ (red trian-  
 227 gles) and MYNN (blue squares) forecasts are higher than the observations, then rise less  
 228 than the lidar observations in the late morning and early afternoon (12-17 UTC, there  
 229 are no radiosonde measurements to compare to here) before rising much higher than the  
 230 observations in the late afternoon (18-24 UTC).



c

**Figure 1.** PBLH vs UTC time for July 11, 2015 for lidar backscatter (black \*), WRF model - MYJ (red triangles), WRF model - MYNN (blue squares), and radiosonde observations using parcel method (green triangles).

### 231 3.2 Impact of assimilation on state variables

232 Since we are primarily interested in the impact of the assimilation on state vari-  
 233 ables within the boundary layer, in Figures 2 and 3 we plot the RMS difference between  
 234 the model and the independent (unassimilated) radiosonde profiles from the surface to  
 235 roughly the top of the boundary layer in the late afternoon. This corresponds to the first

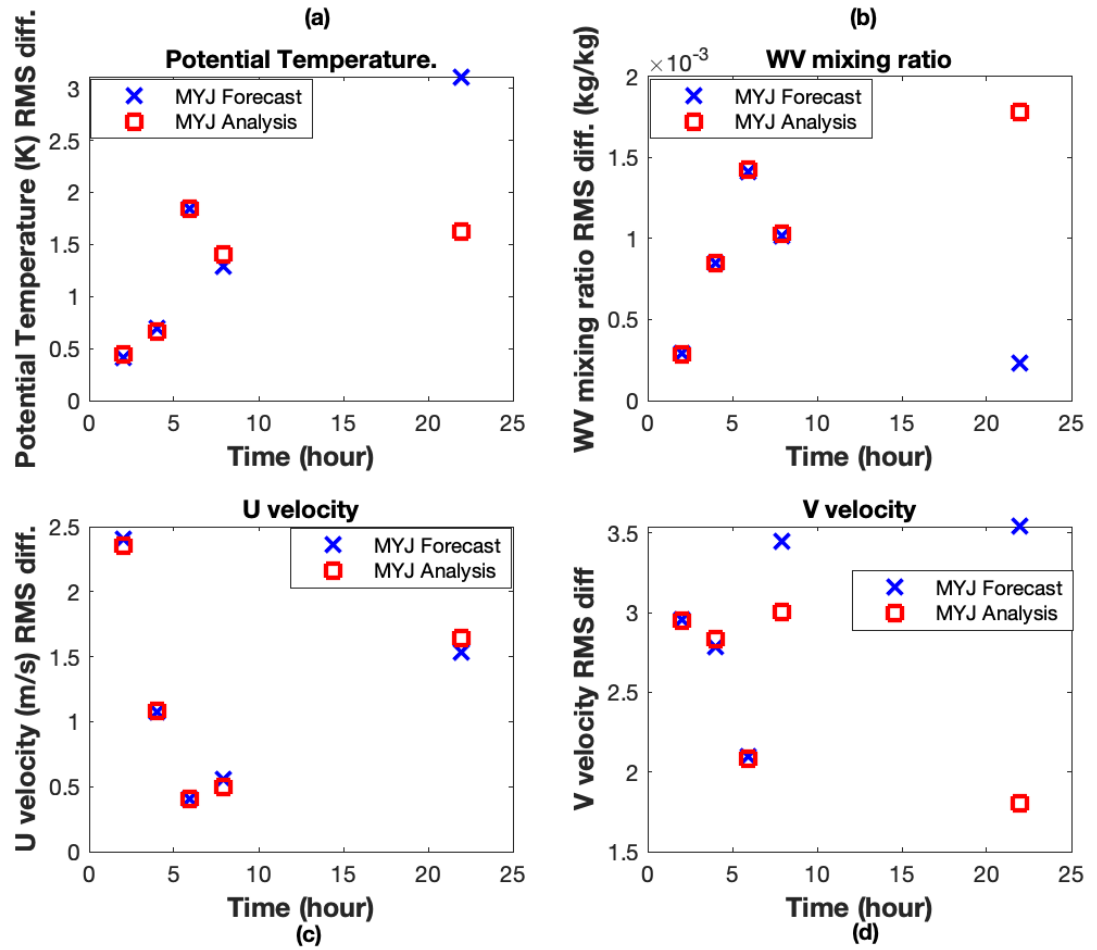
236 8 layers, or about 800 mb. We use a fixed number of layers so as to make the compar-  
 237 isons of the RMS differences consistent during the day, rather than computing the RMS  
 238 over a different number of layers as the PBL grows during the day. For the temperature  
 239 forecast, the RMS difference would is

$$RMS(t_a) = \left[ \frac{1}{8} \sum_{i=1}^8 (T_i^f - T_i^{sonde})^2 \right]^{1/2} \quad (8)$$

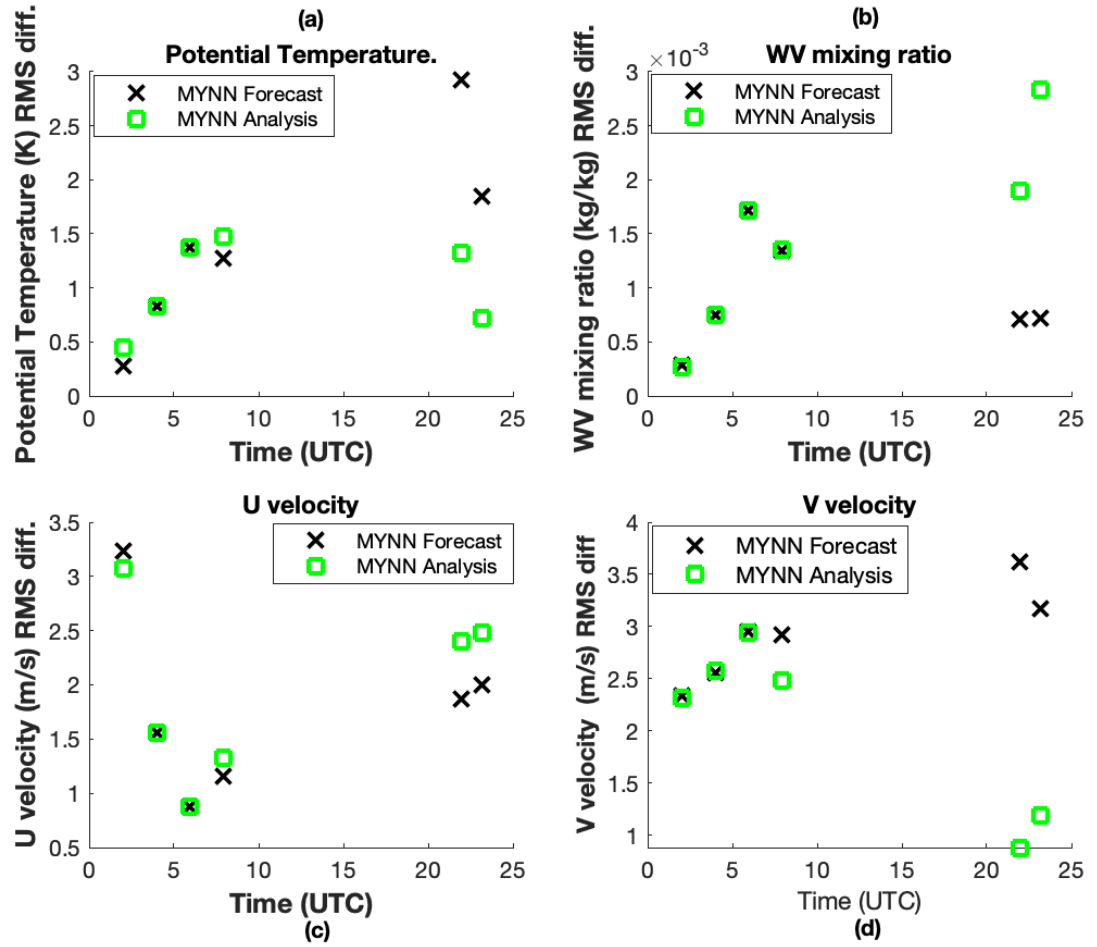
240 where  $t_a$  is the analysis time and "i" represents the model level. Figures 2 and 3 show  
 241 the RMS differences with the radiosonde profiles throughout the day for the forecasts  
 242 (blue x) and analyses (red squares) for potential temperature (a), water vapor mixing  
 243 ratio WV (b) and the U (c) and V (d) components of velocity.

244 During the night (2-9 UTC), the assimilation has a relatively smaller impact on  
 245 the potential temperature RMS differences (upper left) in the early morning (6 and 8  
 246 UTC), and the two forecasts have similar accuracy. By late afternoon (22 and 23 UTC,  
 247 note that the MYJ forecast stops at 22 UTC) the radiosonde comparisons show that the  
 248 assimilation reduces RMS differences in the potential temperatures by around 1.5K for  
 249 MYJ and 2K for MYNN. The water vapor mixing ratio (upper right) also has little im-  
 250 pact from the assimilation between 2 and 8 UTC, but at 22 UTC (the next radiosonde  
 251 profile) the RMS differences for both MYJ and MYNN analyses increase by at least  $1.5 \times$   
 252  $10^{-3} kg/kg$  in the late afternoon. The U-velocity profiles (lower right) show small dif-  
 253 ferences between the MYJ and MYNN through 8 UTC (3 a.m. local time) and the as-  
 254 similation increases the RMS differences with radiosonde profiles by nearly  $1m/s$  start-  
 255 ing at 22 UTC for both models. The V-velocity profiles (d) begin to differ between MYJ  
 256 and MYNN for the forecasts at 8 UTC ( $0.5m/s$  decrease), and assimilation also decreases  
 257 the RMS differences with radiosondes in late afternoon by  $1.5 - 2m/s$ .

258 We would like to understand why there is a smaller impact during night time and  
 259 early morning, whereas there are decreases in the RMS differences in temperature and  
 260 V velocity and increases in moisture and U velocity in the late afternoon. To this end,  
 261 we plot the forecast, analysis and radiosonde profiles (T, Q, U and V) at 4 UTC (11 p.m.  
 262 local time) and 22 UTC (5 p.m. local time) in Figures 4-7. At 4 UTC, (Figures 4,5) these  
 263 clearly indicate that there are small corrections made by the assimilation, as the red and



**Figure 2.** RMS difference for lowest 8 layers, vs. time of forecast (blue x) and analysis (red square) with radiosonde profiles for potential temperature (a), water vapor (b), U velocity (c) and V velocity (d).

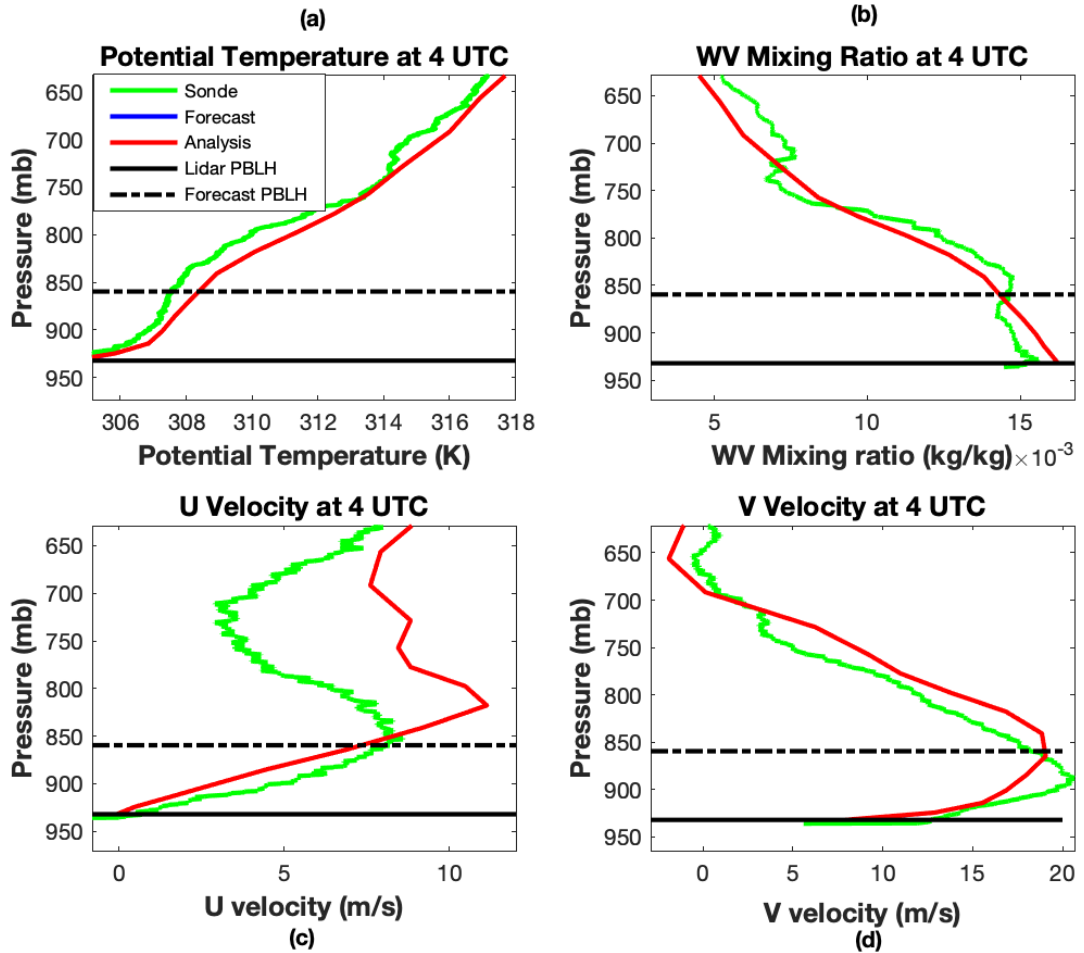


**Figure 3.** Same as Figure 2, but for MYNN PBL model, with forecast (black x) and analysis (blue square).

264 blue profiles closely overlap. But it also shows that the profiles (particularly tempera-  
265 ture and moisture) more accurately follow the radiosonde profiles (except for the U ve-  
266 locity above the PBL), meaning that that any substantial corrections would have made  
267 the profiles worse relative the the radiosonde profiles and ultimately degrade the next  
268 PBLH forecast. In contrast, Figure (1) shows that the forecast PBLH (particularly MYJ)  
269 is quite a bit higher than the lidar observation at 4 UTC. In the late afternoon Figures  
270 6 and 7 indicate that there are large differences between the forecasts and radiosonde  
271 profiles for all of the state variables. The forecast PBLH values differ substantially from  
272 the lidar measurements as well. The correction to the forecast profiles generally pushes  
273 the analyses towards the independent radiosonde profiles, particularly for temperature  
274 and V velocity. So the forecasts that predicted both PBLH and state variables with rel-  
275 atively greater accuracy in the early morning were not corrected, while the less accurate  
276 afternoon forecast was drawn towards the independent radiosonde measurements. The  
277 assimilation also made changes to the vertical velocity (W) in the afternoon, but there  
278 is no independent data to compare with so we have not included it.

279 The WV is shown to be increased by the assimilation (since WV and PBLH are  
280 negatively correlated and higher PBLH corresponds to lower WV levels in the PBL mod-  
281 els), but the analysis overshoots the radiosonde WV profile for MYNN, hence causing  
282 the increase in the water vapor RMS difference in Figures 2 and 3. The MYJ forecast  
283 for WV is mostly too high, so the analysis also increases the RMS difference. Compared  
284 to temperature, WV is highly variable in time and space and it has been shown in the  
285 past that slanted balloon trajectories underestimate the WV present (Demoz et al 2006;  
286 Crook, 1996). The U velocity difference with the radiosonde is larger for the analysis,  
287 but this correction is more difficult because the differences (at least for MYJ) are both  
288 positive and negative and the PBLH observation only contains a single piece of infor-  
289 mation. The V velocity is, on the other hand, greatly improved by the assimilation. These  
290 analysis profiles show that, for this one analysis time, the assimilation is pushing the state  
291 variables in the proper direction for temperature, V velocity and moisture, though the  
292 moisture correction overshoots the radiosonde profile. PBLH is not a prognostic vari-  
293 able, so that the analysis PBLH values are not retained and therefore cannot directly  
294 affect the next forecast. But it is important to note that the temperature and moisture  
295 profiles are changed by the assimilation in a way that indicates that the next forecast  
296 is likely to have a more accurate PBLH estimate. Figures 6 and 7 both show that the

297 level at which the potential temperature begins to rise and the WV mixing ratio begins  
 298 to drop has been moved to a level much closer to that observed by the lidar. We do not  
 299 make forecasts from the analysis fields, but these profiles show promise for improved PBLH  
 300 forecasts when cycling experiments are done in a future implementation.



**Figure 4.** Profiles from radiosonde (green), forecast (blue) and analysis (red) for potential temperature (a), water vapor mixing ratio (b), u-velocity (c) and v-velocity (d) at 4 UTC, July 11, 2015 in Greensburg, KS. The model uses the MYJ physics parameterization.

301 **3.3 Ensemble error covariances**

302 The increasing differences between PBLH and profile forecasts from early morn-  
 303 ing to late afternoon only partly explain the much larger impact of the assimilation at  
 304 22 UTC. We can also analyze the assimilation by investigating the error covariance be-

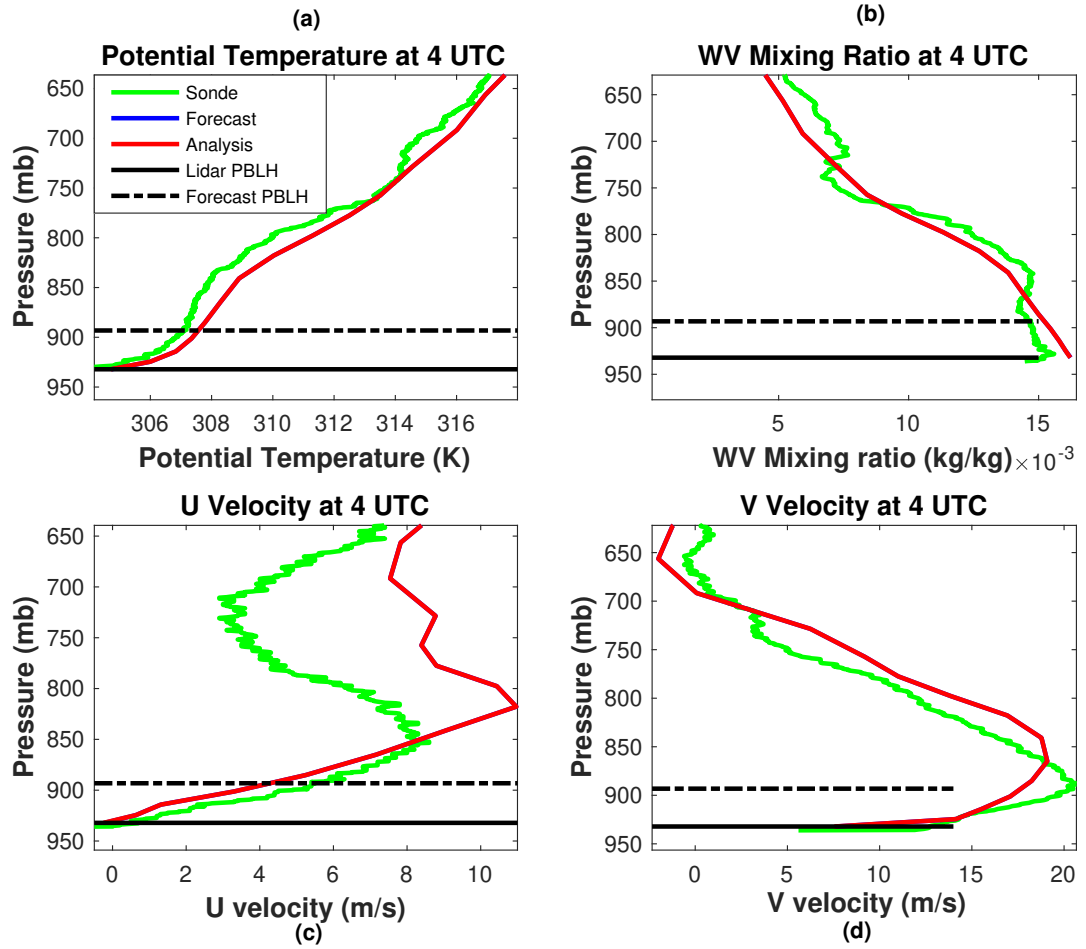
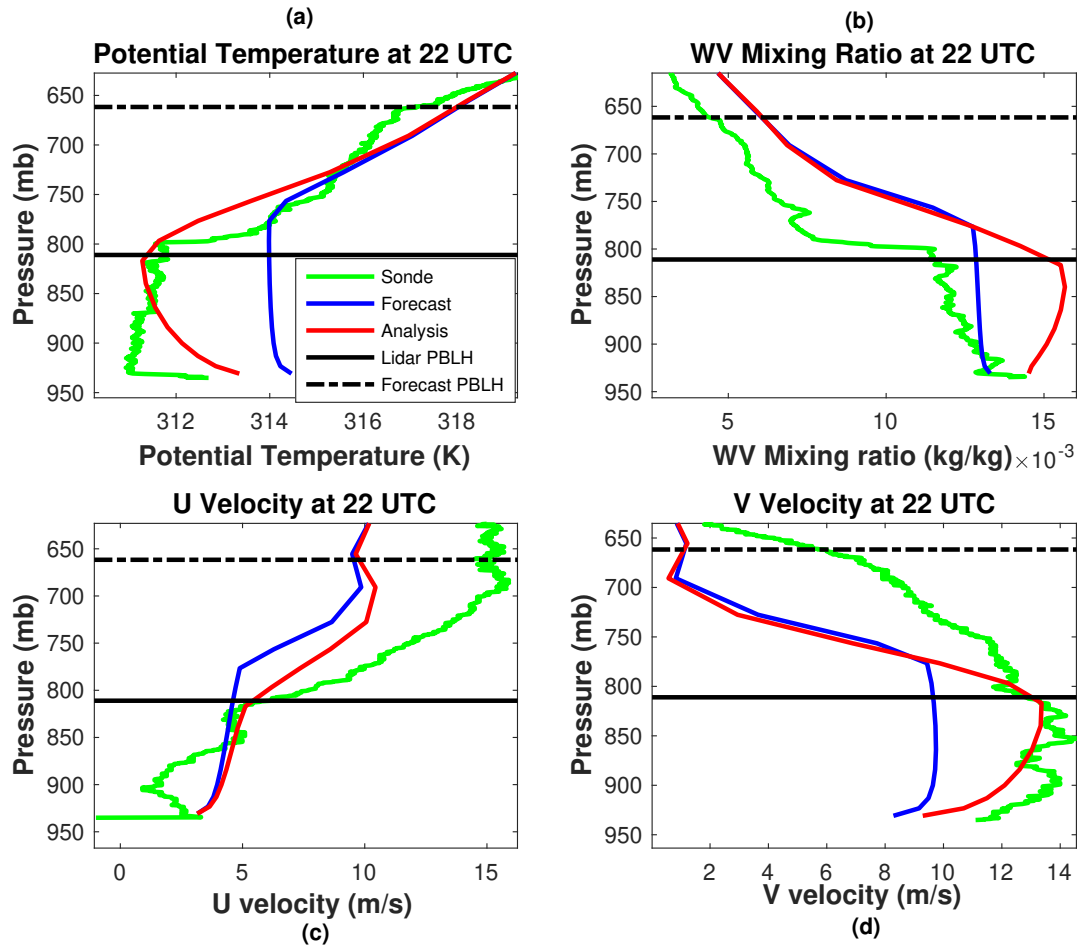
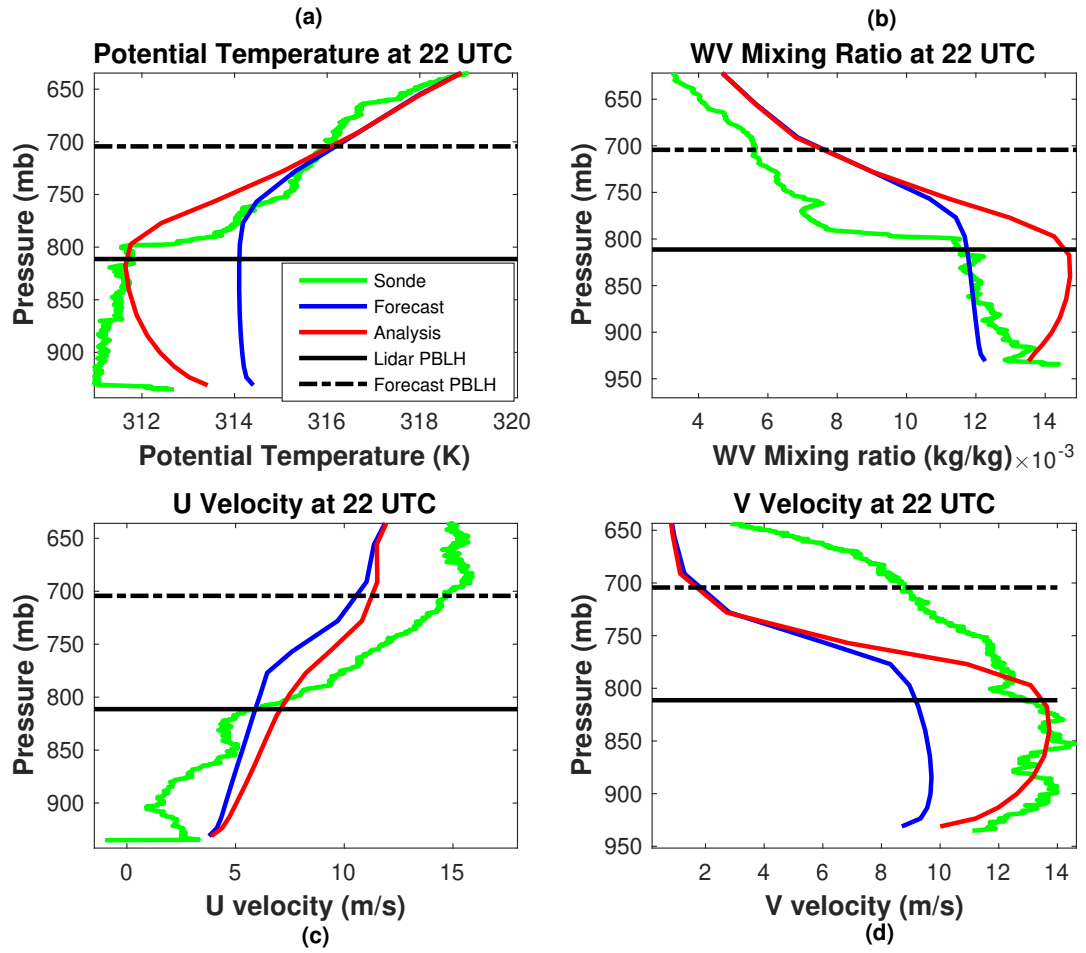


Figure 5. Same as figure 4 except using MYNN model.



**Figure 6.** Same as figure 4 except using except at time 22 UTC.



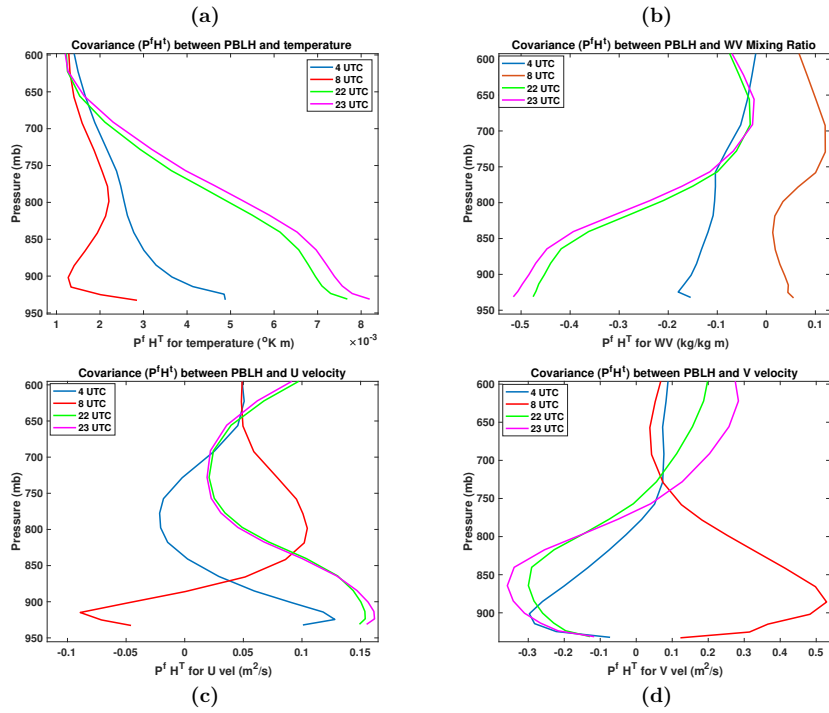


**Figure 7.** Same as figure 6 except using MYNN model.

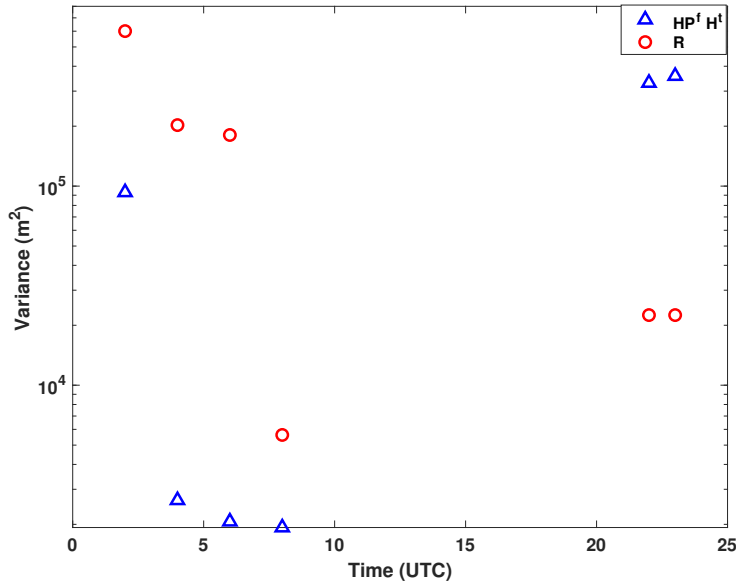
305 tween PBLH and each of the state variables ( $\mathbf{P}^f \mathbf{H}^T$ ) and the relative error variances in  
 306 observation space ( $\mathbf{H} \mathbf{P}^f \mathbf{H}^T$  and  $\mathbf{R}$ ). We show  $\mathbf{P}^f \mathbf{H}^T$  in Figure 8 for the MYNN PBL  
 307 physics model at the 6 radiosonde times. The covariance with temperature is always pos-  
 308 itive, and grows by a factor of 4 by late afternoon near the surface. The covariance with  
 309 WV is mostly negative and grows by roughly a factor of 5, while the covariance with the  
 310 two components of velocity oscillate between positive and negative and shows less con-  
 311 sistent growth. Thus, the largest impact of assimilation on temperature and moisture  
 312 occurs in late afternoon while more limited velocity corrections are largely constrained  
 313 by the correlations determined by the ensemble of model forecast states. In addition, the  
 314 covariance between PBLH and the U velocity are substantially smaller than those with  
 315 the V velocity. This means that spurious correlations between PBLH and U might be  
 316 present, given the relatively small ensemble and the geographic variation of the ensem-  
 317 ble members. The error variances are also plotted at the radiosonde times in Figure 9,  
 318 which shows that the observation errors are much larger than the forecast errors dur-  
 319 ing evening and early morning times (2,4,6,8 UTC) and then become relatively smaller  
 320 in the late afternoon (22,23 UTC). This is an additional contributing factor to the min-  
 321 imal impact of PBLH observations early in the day and the much larger impact in the  
 322 afternoon.

#### 323 4 Discussion and Conclusions

324 These offline data assimilation experiments indicate that assimilating ground based  
 325 lidar backscatter and wind measurements of PBLH into a regional NWP model will likely  
 326 lead to corrections to profiles within the PBL, particularly when, in the future, this ap-  
 327 proach is applied to an EnKF assimilation system with cycling. Using two NU-WRF fore-  
 328 casts over a period of one day with different PBL physics models, we show how the state  
 329 variables, T, WV, U and V can be corrected using an assimilation system with ensem-  
 330 ble based error covariances. During the night and early morning the assimilation has rel-  
 331 atively little impact on the state variables, but by late afternoon the temperature field  
 332 is drawn closer to independent radiosonde measurements. We have shown that the lack  
 333 of data impact early in the day is due to the relatively higher accuracy of the model  
 334 and lack of correlation between the forecast PBLH and temperature profiles at that time.  
 335 Later in the day, when the model is less accurate in predicting the growth of the bound-  
 336 ary layer, the data begins to draw the analyses mostly toward the independent radiosonde



**Figure 8.** Covariance  $\mathbf{P}^f \mathbf{H}^T$  between PBLH and temperature (a), water vapor (b), U velocity (c) and V velocity (d), at times 4, 8, 22 and 23 UTC, for PBL physics model MYHH.



**Figure 9.** Forecast ( $\mathbf{HP}^f \mathbf{H}^t$ ) and observation ( $\mathbf{R}$ ) error covariance for the PBL physics model MYHH at the 6 radiosonde times.

337 profiles. The assimilation overcorrected the water vapor mixing ratio in the direction of  
 338 radiosonde data, and this could likely be tuned in an assimilation system. And it cor-  
 339 rected the the V velocity component by a smaller amount, and reduced differences with  
 340 the radiosonde profiles for the V velocity. These corrections are the result of ensemble  
 341 computed error covariances between the PBLH and the state variable profiles within the  
 342 PBL. The results here indicate that this approach has some potential to be used in a fore-  
 343 cast system in a way that that the PBLH observational information could be carried for-  
 344 ward in time so as to impact the forecast accuracy within the PBL. An additional value  
 345 of assimilating PBLH is its close connection with the PBL scheme used in the model.  
 346 The ensemble covariances between PBLH and the different state variables are controlled  
 347 through the PBL physics scheme. This has an impact on the corrections made to the  
 348 profiles within the PBL, which can be used as another way to evaluate the physics pa-  
 349 rameterizations. For example, the MYJ and MYNN result in forecast profiles that dif-  
 350 fer, particularly in WV in the late afternoon. And the differences in reponse to assim-  
 351 ilation are an indication of how the two different PBL schemes affect the covariance be-  
 352 tween PBLH and the state variables. However, a full evaluation would require that the  
 353 assimilation be implemented into a cycling data assimilation system.

354 This work is intended only to demonstrate a necessary first step in terms of how  
355 ensemble statistics can help to constrain profiles within the PBL by assimilating PBLH  
356 observations. A more complete demonstration of this approach will require the construc-  
357 tion of an EnKF, which should be run over many days with a variety of weather patterns,  
358 including significantly warmer(cooler) and wetter(drier) days. This is needed to show  
359 how the assimilated PBLH observations will impact future forecasts within the PBL. More  
360 of the PBL physics schemes need to be investigated as well, since the correlation between  
361 PBLH and state variables will vary widely depending on which scheme is used. An es-  
362 timate of the forward operator error should be included in the algorithm as well. There  
363 are also differences in the way PBLH is computed in the PBL physics schemes, and the  
364 methods used for radiosonde observations (see Hegarty, et al., 2018). This will impact  
365 the manner in which the assimilation and resulting forecasts are validated. The larger  
366 uncertainty in the lidar PBLH retrievals during nighttime (Figure 9) mean that the as-  
367 similation will not significantly constrain the model state within the PBL during this pe-  
368 riod. So it would be very helpful to complement PBLH observations with thermodynamic  
369 and kinematic profile observations, particularly overnight. The fact that PBLH is a non-  
370 negative variable means that the O-F statistics will likely be non-Gaussian so that the  
371 assimilation algorithm would need to include an extension to handle this possibility (e.g.  
372 Cohn, 1997).

373 In addition, a cycling EnKF will involve spatial covariances in both horizontal and  
374 vertical directions, and will allow for both inflation and horizontal localization. This will  
375 enable further tuning of the system to optimize the analysis state relative to the inde-  
376 pendent radiosonde observations. The PBLH assimilation within the EnKF framework  
377 could be done in any of numerous existing EnKF assimilation systems that connect with  
378 WRF, including NU-WRF (Lidard-Peters *et al.*, 2015) and WRF-DART (Anderson *et*  
379 *al.*, 2009). Future development of PBLH assimilation algorithms will also need to ad-  
380 dress the effect of the different definitions of PBLH, such as the TKE method used the  
381 physics schemes and the backscatter and wind profile method used in the retrievals.

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## 391 **6 Data Sets**

392 PECAN ([https://data.eol.ucar.edu/master\\_list/?project=PECAN\verb](https://data.eol.ucar.edu/master_list/?project=PECAN\verb)) data are  
393 archived by NCAR/EOL, which is funded by NSF. The forecast and analysis fields pro-  
394 duced for this work are stored at <https://alg.umbc.edu/pecan/>.

## 395 **7 Competing Interests**

396 The authors declare that they have no conflict of interest.

## 397 **8 Author Contributions**

398 Andrew Tangborn built the assimilation system, with input from Jeffrey Anderson on  
399 the algorithm. Belay Demoz and Brian Carroll provided the lidar observations. Joseph  
400 Santanello provided background information on PBL physics. All of the authors contributed  
401 to writing and revising the paper.

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