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Radiative transfer acceleration based on the Principal Component Analysis and Look-Up Table of corrections: Optimization and application to UV ozone profile retrievals

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Abstract

18 In this work, we apply a principal component analysis (PCA)-based approach combined with look-up 19 tables (LUTs) of corrections to accelerate the VLIDORT radiative transfer (RT) model used in the 20 retrieval of ozone profiles from backscattered ultraviolet (UV) measurements by the Ozone Monitoring 21 Instrument (OMI). The spectral binning scheme, which determines the accuracy and efficiency of the 22 PCA-RT performance, is thoroughly optimized over the spectral range 265 to 360 nm with the 23 assumption of a Rayleigh-scattering atmosphere above a Lambertian surface. The high level of accuracy 24 $(\sim 0.03 \text{ \%})$ is achieved from fast-PCA calculations of full radiances. In this approach, computationally 25 expensive full multiple scattering (MS) calculations are limited to a small set of PCA-derived optical 26 states, while fast single scattering and 2-stream multiple scattering calculations are performed, for every 27 spectral point. The number of calls to the full MS model is only 51 in the application to OMI ozone 28 profile retrievals with the fitting window of 270-330 nm where the RT model should be called at fine 29 intervals (~0.03 nm with ~ 2000 wavelengths) to simulate OMI native measurements at 229 30 wavelengths (spectral resolution: 0.4-0.6 nm). We also developed a Look Up Table (LUT) to correct RT 31 approximations performed using a scalar RT model with 4 streams (discrete ordinates) and 24 layers, 32 thereby achieving the accuracy at the level attainable from simulations with a vector model with 12 33 streams and 72 layers; this speeds up the RT calculations by more than 2 orders of magnitude when 34 ignoring other overhead. Overall, we speed up our OMI retrieval by a factor of 3.3 over the previous 35 version, which has already been significantly sped up over line-by-line calculations due to various RT 36 approximations. Improved treatments for RT approximation errors using LUT corrections improve





spectral fitting (2-5 %) and hence retrieval errors, especially for tropospheric ozone by up to \sim 10%; the remaining errors due to the forward model errors are within 5 % in the troposphere and 3 % in the stratosphere.

40 1. Introduction

41 Optimal estimation-based inversions have become standard for the retrieval of atmospheric ozone 42 profiles from atmospheric chemistry UV-Vis backscatter instruments. This inversion model requires 43 iterative simulations of not only radiances, but also of Jacobians with respect to atmospheric and surface 44 variables, until the simulated radiances are sufficiently matched with the measured radiances. These 45 ozone profile algorithms face a computational challenge for use in global processing of high 46 spatial/temporal resolution satellite measurements, due to on-line radiative transfer (RT) computations 47 at many spectral points from 270 to 330 nm; it is computationally very expensive to perform full 48 multiple-scattering (MS) simulations with the polarized RT model. To reduce the computational cost, a 49 scalar RT model can be applied together with a polarization correction scheme based on a LUT (Kroon 50 et al., 2011; Miles et al., 2015). Another approach is to carry out on-line vector calculations at a few 51 wavelengths (Liu et al., 2010) together with other approximations (e.g., low-stream, coarse vertical 52 layering, Lambertian reflectance for surface and cloud, no aerosol treatment). However, the 53 computational speed is still insufficient to process one day of measurements from the Aura Ozone Monitoring Instrument (OMI) within 24 h (30 cross-track pixels × 1644 along-track pixels × 14 54 55 orbits) with reasonable computational resources. Consequently, only 20 % of the available OMI pixels are processed to generate the operational ozone profile (OMO3PR) product (Kroon et al., 2011), and 56 57 the spatial resolution is degraded by a factor of 4 to produce the research ozone profile (OMPROFOZ) 58 product (Liu et al., 2010). With the advent of sophisticated inversion techniques and superior 59 spaceborne remote sensing instruments, computational budgets have increased rapidly in recent years. 60 Joint retrievals combining UV and thermal infrared (~ 9.6 um) have been investigated to better 61 distinguish between upper- and lower-tropospheric ozone abundances from multiple instruments, e.g., 62 OMI + TES, OMI +AIRS, and GOME-2 + IASI (Fu et al., 2013; 2018; Cuesta et al., 2013). The 63 geostationary satellite instrument Tropospheric Emissions: Monitoring of Pollution (TEMPO), 64 scheduled for Launch in 2022, is specially designed for joint retrievals combining UV and visible (540-65 740 nm) radiances to enhance the performance of retrievals for ground-level ozone (Zoogman et al., 66 2017). Moreover, the temporal and spatial resolutions of upcoming geostationary satellite instruments 67 are being improved, leading to a tremendous increase in the data volume to be processed; for example, 68 daily measurements of TEMPO (with ~2000 N/S cross-track pixels \times ~1200 E/W mirror steps \times ~8





times a day) are ~30 times greater in volume than those of OMI. Therefore, accelerating RT simulations is one of the highest priority tasks to assure operational capability. For speed-up, LUTs have often been used in trace gas retrieval algorithms to serve as proxies for RT modeling or to perform corrections to on-line RT approximations. In recent years, applying neural network techniques and principal component analysis (PCA) to RT computational performance has received quite a lot of attention (e.g., Natraj et al., 2005; Spurr et al., 2013;2016; Liu et al., 2016; Yang et al., 2016; Loyola et al., 2018; Nanda et al., 2019; Liu et al., 2020).

76 The goal of this paper is to improve both computational efficiency and accuracy of RT 77 simulations in the OMI ozone profile algorithm (Liu et al., 2010) by combining a fast PCA-based RT 78 model with two kinds of correction techniques. The application of PCA to RT simulations was first 79 proposed by Natraj et al. (2005) by demonstrating a computational improvement of intensity simulation 80 in the $O_2 A$ band by a factor of 10 and with ~ 0.3 % accuracy compared to full line-by-line (LBL) 81 calculations. This scheme has been deployed to the UV-backscatter, thermal emission, and cross-over 82 régimes, and has been extended for the derivation of analytic Jacobians, for vector RT applications, and for bidirectional surface reflectances (Kopparla et al., 2016; 2017; Natraj et al., 2010; Somkuti et al., 83 84 2017; Spurr et al., 2013). The RT performance enhancement arises from a reduction in the number of 85 expensive full multiple scattering calculations; the PCA scheme uses spectral binning of the 86 wavelengths into several bins based on the similarity of their optical properties and the projection to 87 every spectral point of these full MS calculations which are executed for a small number of PCA-88 derived optical states. In addition to the adaption of a PCA-based RTM for our ozone profile retrieval, 89 we have adopted the undersampling correction from our previous implementation (Kim et al., 2009; 90 Bak et al., 2019); this enables us to use fewer wavelengths for further speed-up without much loss of 91 accuracy. Furthermore, we have developed a LUT-based correction to accelerate on-line RT simulations, 92 by starting with a lower-accuracy configuration (scalar RT with no polarization, 4 streams, 24 layers) 93 and then correcting the accuracy to the level attainable by means of a computationally more expensive 94 configuration (vector RT, 12 streams, 72 layers). The stream value refers to the number of discrete 95 ordinates in the full polar space; thus, for example, the term "12 streams" indicates the use of 6 96 upwelling and 6 downwelling polar cosine discrete ordinate directions. In previous work, PCA-based 97 RT calculations were assessed mostly against LBL calculations, independently from the inverse model. 98 Therefore, the PCA performance is likely to be overestimated in terms of operational capability, because 99 operational algorithms have their own speed-up strategies with many approximations; this is the case 100 for our ozone profile algorithm. As mentioned above, the PCA-based RT model is employed in this 101 work to make forward-model simulations of OMI measurements for the retrieval of ozone profiles.





- 102 Therefore, we evaluate the operational capability of our retrieval algorithm in terms of the retrieval
- 103 efficiency as well as the accuracy, and assess these relative to the current operational implementation.

This paper is structured as follows. Section 2 describes the current forward model scheme and evaluates the approximations made in RT calculations, with the determination of the configuration parameters for accurate simulations. The updated forward model scheme is introduced for the PCAbased RT model in Section 3.1, and the two kinds of correction schemes to use less spectral sampling and less accurate RT configuration are detailed in Section 3.2. The evaluation is performed in Section 4 and then we summarize and discuss the results in Section 5.

110 2. Current forward model scheme based on Vector LIDORT (VLIDORT) only

111 We first describe the current v1 SAO OMI ozone profile algorithm that was implemented in OMI 112 Science Investigator-led Processing Systems (SIPS) to generate the research OMPROFOZ ozone 113 profile product, publicly available at the Aura Validation Data Center (AVDC, 114 https://avdc.gsfc.nasa.gov/index.php?site=1620829979&id=74). It employs the OMI UV channel that 115 is divided into UV1 (270-310 nm) and UV2 (310-380 nm). The spatial resolution of UV1 is degraded 116 by a factor of 2 in order to increase the signal to noise ratio (SNR) in this spectral region. The full width 117 at half maximum (FWHM) of the instrument spectral response function (ISRF) is ~0.63 nm for UV1 118 and ~0.42 nm for UV2, with corresponding spectral intervals of 0.33 nm and 0.14 nm, respectively. The 119 total number of OMI wavelengths used in our spectral fitting for ozone profiles is 229, from 270-308 120 nm (UV1) and 312-330 nm (UV2). The RT model needs to simulate sun-normalized radiances as well as their derivatives with respect to the ozone profile elements and surface albedo. This simulation is 121 122 iteratively performed to ingest the atmospheric and surface variables adjusted through the physical 123 fitting between measured and simulated spectra and simultaneously the statistical fitting between the 124 state vector and the *a priori* vector. The retrieval is optimized within typically 2-3 iterations (up to 10 125 is permitted). The vertical grids of the retrieved ozone profiles in 24 layers are initially spaced in log (pressure) at $P_i = 2^{-\frac{i}{2} atm}$ for 0 (surface), 23 (~55 km) and with the top of atmosphere set for P_{24} 126 127 (~65 km). Each layer is thus approximately 2.5-km thick, except for the top layer (~ 10 km). A number 128 of RT approximations have already been applied in the current forward model to speed up the processing. 129 In the remainder of this section, the current forward model scheme is described, with its flow chart 130 depicted in the left panel of Fig. 1. An error analysis is performed for optimizing the RT model 131 configuration to maximize the simulation accuracy.





132 In the first step, we select 93 effective wavelengths with variable sampling intervals, 1.0 nm below 133 295 nm, 0.4 nm from 295-310 nm, and 0.6 nm above 310 nm. The number of the wavelengths is smaller 134 than the OMI native pixels (229 from 270-330 nm) by more than a factor of 2. The on-line radiative 135 transfer model is run to generate the full radiance spectrum (single + multiple scattering) at these 136 wavelengths in the scalar mode, with 8 streams and a Rayleigh atmosphere divided into 25 layers - a 137 grid that is similar to that for the retrieval, except for the top layer (~ 55 km to 65 km) which is further 138 divided into two layers. Note that the Vector Linearized Discrete Ordinate Radiative Transfer VLIDORT 139 model Version 2.8 (Spurr and Christi, 2019) is implemented in this study. In step 2, a polarization 140 correction is applied to the scalar calculations done in step 1 using the on-line vector calculation at 141 fourteen wavelengths (visually shown with the vertical lines in Fig. 3.a.2). In step 3 the simulation at 142 the effective wavelength grid is interpolated into 0.05 nm intervals with the undersampling correction, 143 and the result is finally interpolated/convolved into OMI native grids in step 4.

Figure 2.b shows approximation errors related to undersampling from 0.02 nm to 0.1 nm compared to the simulated radiance at the sampling rate of the ozone cross sections (0.01 nm) (Fig. 2.a). This illustrates that current forward model calculation has trivial errors (less than 0.01 %) except for 0.02 % around 310 nm if there is no error after undersampling correction to 0.05 nm. The correction applied in step 3 allows relaxation of the sampling rate without loss of the accuracy. This correction is based on the adjustment of the radiance due to the difference of the optical depth profiles between fine (λ_h) and coarse (λ_c) spectral grids assisted by application of the weighting functions ($\frac{dI}{dx}$) as follows:

152
$$I(\lambda_h)$$

153
$$= I(\lambda_c)$$

$$+\sum_{l=1}^{N} \frac{\partial I (\lambda_{c})}{\partial \Delta_{l}^{gas}} \left(\Delta_{l}^{gas}(\lambda_{h}) - \Delta_{l}^{gas}(\lambda_{c}) \right)$$

155
$$+ \frac{\partial I (\lambda_c)}{\partial \Delta_l^{ray}} \left(\Delta_l^{ray} (\lambda_h) - \Delta_l^{ray} (\lambda_c) \right), \quad (1)$$

151

where Δ_l^{gas} and Δ_l^{ray} are the optical depth profiles for trace gas absorption and Rayleigh scattering, $l = 1, \dots N_L$ (the number of atmospheric layers). Figure 2.c demonstrates that the undersampling correction works well for simulations at 0.2 nm intervals or less over the entire spectral range, but it can cause large errors when the simulations are performed at intervals of 1.0 nm, 0.4, and 0.6 nm for the spectral ranges, 270-295 nm, 295-310 nm, and 310-330 nm, respectively. Figure 3 shows the





161 approximations applied to on-line VLIDORT calculations, including (a.1) neglect of the polarization 162 effect; (b) use of 8 streams; and (c) use of a coarse 24-layer height grid. As we mentioned above, in the 163 v1 forward model the scalar model is used for all wavelengths, with the vector model at 14 wavelengths 164 for correcting the scalar simulations. However, Figure 3.a.2 illustrates that second order of polarization 165 correction errors (~ 0.2 %) could remain due to neglecting the dependence of polarization effects on the 166 fine structures of ozone absorption. Using 8 streams causes errors of ~ 0.05 % above 320 nm, whereas 167 using the 24 layers causes 1 % errors at shorter UV wavelengths. Based on the results shown in Fig 3, 168 we conclude that there is room for improving the simulation accuracy by increasing the number of 169 streams to 12, dividing the atmosphere into 72 layers and using more wavelengths in the polarization 170 correction.

171 **3.** The improved forward model scheme based on PCA-VLIDORT

172 The right panel of Fig. 1.2 illustrates the flow chart of the new forward model scheme (v2) which 173 employs the PCA-based RT model to perform on-line scalar simulations using 4 streams and a 24-layer 174 atmosphere for RT performance enhancement (step 1) and two kinds of correction schemes for 175 accounting for approximation errors (steps 2 and 3). Section 3.1.1 gives an overview on how the PCA 176 tool is combined with the VLIDORT Version 2.8 model; full theoretical details may be found in Spurr 177 et al. (2016) and Kopparla et al. (2017). Here, our paper gives details on how the PCA-based RT 178 configuration is optimized for the application to UV ozone profile retrievals for maximizing the speed-179 up in the section 3.1.2. Section 3.2 specifies the step 2 wherein the LUT-based correction is applied to 180 simulation errors due to the use of a scalar model, a smaller number of streams and coarser-resolution 181 vertical grid. In the step 3 the undersampling correction is adopted from the v1 implementation, but the 182 Rayleigh scattering term of the equation 1 is neglected for the speed up with trivial loss of accuracy.

183 **3.1.1 General PCA procedure**

The PCA-based RT process begins with a grouping of spectral points into several bins; atmospheric profile optical properties within each bin are similar. PCA is a mathematical transformation that converts a correlated mean-subtracted dataset into a series of principal components (PCs). To enhance RT performance, PCA is used to compress a binned set of correlated optical profile data into a small set of atmospheric profiles which capture the vast majority of the data variance within the bin. The layer extinction optical thickness Δ_{ni} and the single scattering albedos ω_{ni} are generally subjected to PCA, where *n* and *i* are indices for atmospheric layers $(n = 1, \dots N_L)$ and spectral points $(i = 1, \dots N_S)$,





191 respectively. For each bin, the optical profiles $\{\ln \Delta_{ni}, \ln \omega_{ni}\}$ is composed of $2N_L \times N_S$ matrix G in 192 log-space $(G_{n,i} = \ln \Delta_{ni}, G_{n+N_N,i} = \ln \omega_{ni})$. The mean-removed $2N_L \times 2N_L$ covariance matrix Y is 193 then:

194
$$Y = [G - \langle G \rangle]^{T} [G - \langle G \rangle], (2)$$

195 where \diamond denotes a mean-value over all grid points in a bin. This covariance matrix Y is decomposed 196 into eigenvalues ρ_k and unit eigenvectors X_k through solution of the eigenvalue problem $YX_k =$ 197 $\rho_k X_k$. The scaled eigenvectors of the covariance matrix are defined as the empirical orthogonal 198 function (EOFs), $W_k = \sqrt{\rho_k} X_k$, where the index k is ranked from 1 to $2N_L$ in descending order 199 staring with the largest eigenvalues. The principal components (PCs) are the projections of the original data onto the eigenvectors, $P_k = \frac{1}{\sqrt{\rho_k}} GW_k$. The original data set can then be expanded in terms of the 200 201 mean value and a sum over all EOFs. As inputs to the RT simulation, the PCA-defined optical states are defined as $F_0 = \exp[\langle G \rangle]$ and $F_k^{\pm} = F_0 \exp[\pm W_k]$, corresponding respectively to the mean value and 202 203 to positive and negative perturbations from the mean value by an amount equal to the magnitude of k^{th} 204 EOF. Therefore, $\Delta_{n,i}$ and $\omega_{n,i}$ (i=1...N_S) are expressed as followings:

205
$$\mathbf{F}_{o} = \left\{ \begin{array}{c} \Delta_{n,o} \\ \omega_{n,o} \end{array} \right\} \equiv \left\{ \begin{array}{c} \exp\left[\frac{1}{N_{s}}\sum_{i=1}^{N_{s}}\ln\Delta_{ni}\right] \\ \exp\left[\frac{1}{N_{s}}\sum_{i=1}^{N_{s}}\ln\omega_{ni}\right] \end{array} \right\}; \ \mathbf{F}_{k}^{\pm} = \left\{ \begin{array}{c} \Delta_{n,\pm k} \\ \omega_{n,\pm k} \end{array} \right\} \equiv \left\{ \begin{array}{c} \Delta_{n,o} \exp\left[\pm W_{n,k}\right] \\ \omega_{n,o} \exp\left[\pm W_{n+N_{L},k}\right] \end{array} \right\}.$$
(3)

206 For those optical quantities not included in the PCA reduction but still required in the RT simulations, 207 the spectral mean values for the bin are assumed, as long as they have smooth monotonic spectral 208 dependency or else are constant over the bin range. In our application, the phase functions and phase 209 matrices for Rayleigh scattering are derived from bin-average values of the depolarization factor. 210 Surface Lambertian albedos are constant in the RT simulation, but the calculated radiance is later 211 adjusted to account for 1th order wavelength dependency using surface albedo weighting functions. For 212 larger bins, it is possible to include the depolarization ratio or the Lambertian albedo as additional 213 elements in the optical data set subject to PCA; this has been investigated in another context by Somkuti 214 et al. (2017).

215 In the PCA-based RT package, three independent RT models are combined in order to generate the 216 full scattering intensity field (I_{Full}) at each spectral point λ_i in a single bin





217
$$I_{\text{full}}(\lambda_i) \cong [I_{2s}(\lambda_i) + I_{\text{FO}}(\lambda_i)]C(\lambda_i). (4)$$

Two fast RT models, the "First-Order" (FO) and 2STREAM (2S), are used to generate an accurate single scatter (SS) field (I_{FO}) and an approximate multiple scatter (MS) field (I_{2S}), respectively, for every spectral point. The scalar 2S model computes the radiation field with 2 discrete ordinates only. To derive the correction factors $C(\lambda_i)$, we first compute (logarithmic) ratios of the full-scatter and 2Sbased intensity fields calculated with PCA-derived optical states F_0 and F_k^{\pm} :

223
$$J_o = \ln \left[\frac{I_{VLD}(F_o) + I_{FO}(F_o)}{I_{2S}(F_o) + I_{FO}(F_o)} \right] ; J_k^{\pm} = \left[\frac{I_{VLD}(F_k^{\pm}) + I_{FO}(F_k^{\pm})}{I_{2S}(F_k^{\pm}) + I_{FO}(F_k^{\pm})} \right].$$
(5)

Intensity ratios at the original spectral points $J(\lambda_i)$ are then obtained using a second-order central difference expansion based on the PCA principal components P_{ki} :

226
$$J(\lambda_i) = J_o + \sum_{k=1}^{N_{EOF}} \frac{(J_k^+ - J_k^-)}{2} P_{ki} + \frac{1}{2} \sum_{k=1}^{N_{EOF}} (J_k^+ - 2J_o + J_k^-)^2 P_{ki}^2 .$$
(6)

227 The correction factors $C(\lambda_i) = \exp[J(\lambda_i)]$ are then applied to the approximate simulation $[I_{2s}(\lambda_i) + I_{FO}(\lambda_i)]$ according to Equation 4 above. More details can be found in the literature (Natraj et al., 2005, 2010; Spurr et al., 2013; 2016; Kopparla et al., 2017).

230 So far, we have discussed generation of total *intensity* field, using values $I_{FO}(\lambda_i)$ and $I_{2s}(\lambda_i)$ from full-spectrum FO and 2S model calculation, and PCA-derived values $I_{VLD}(F)$, $I_{2S}(F)$ and 231 232 $I_{FO}(F)$ based on PCA-derived optical states $F = \{F_o, F_k^{\pm}\}$. The above procedure works with 233 VLIDORT operating in scalar or vector mode; however, the 2S model is purely scalar, and cannot be 234 used if we want to establish PCA-RT approximations to the Q and U components of the Stokes vector 235 with polarization present. Instead, we rely on just the VLIDORT and FO models, and develop a PCA-236 RT scheme based on the differences between the VLIDORT and FO Q/U values for monochromatic and 237 PCA-derived calculations, with an additive correction factor instead of the logarithmic ratios in 238 Equation (6) above. This was first introduced in Natraj et al. (2010), and is discussed in detail in Spurr 239 et al. (2016).

240 Of greater importance for us is the need to derive PCA-RT approximations to profile Jacobians 241 (weighting functions of the total intensity with respect to ozone profile optical depths). A PCA-RT





Jacobians scheme was developed by Spurr et al. (2013) for total column Jacobians in connection with the retrieval of total ozone; this scheme involved formal differentiation of the entire PCA-RT system as outlined above for the intensity field. This is satisfactory for bulk property Jacobians, but for profile Jacobians it is easier to write (Efremenko et al., 2014; Spurr et al., 2016):

246
$$\mathbf{K}_{Full}^{(\xi)}(\lambda_i) \cong \left[\mathbf{K}_{2S}^{(\xi)}(\lambda_i) + \mathbf{K}_{FO}^{(\xi)}(\lambda_i)\right] D^{(\xi)}(\lambda_i), \tag{7}$$

Here, $K_{2S}^{(\xi)}(\lambda_i) \equiv \frac{\partial I_{2S}(\lambda_i)}{\partial \xi}$, with similar definitions for the FO and VLIDORT partial derivatives. The Jacobian correction factor $D^{(\xi)}(\lambda_i) = \exp[L^{(\xi)}(\lambda_i)]$ is determined using the same central-difference expansion as that in Equation (6), but with quantities

250
$$L_{0}^{(\xi)} = \ln \left[\frac{K_{VLD}^{(\xi)}(F_{o}) + K_{FO}^{(\xi)}(F_{o})}{K_{2S}^{(\xi)}(F_{o}) + K_{FO}^{(\xi)}(F_{o})} \right] ; L_{\pm k}^{(\xi)} = \left[\frac{K_{VLD}^{(\xi)}(F_{k}^{\pm}) + K_{FO}^{(\xi)}(F_{k}^{\pm})}{K_{2S}^{(\xi)}(F_{k}^{\pm}) + K_{FO}^{(\xi)}(F_{k}^{\pm})} \right]$$
(8)

251 in place of J_o and J_k^{\pm} in Equation (5).

252 **3.1.2** The binning scheme

253 The major performance saving is achieved by limiting full-MS VLIDORT calculations to those 254 based on the reduced set of PCA-derived optical states F_0 and F_k^{\pm} . A general binning scheme has been 255 developed over the shortwave region from 0.29 to 3.0 µm (Kopparla et al. 2016), whereby the entire 256 region is divided into 33 specially-chosen sub-windows encompassing the major trace-gas absorption 257 signatures; in each such sub-window there are 11 bins for grouping optical properties, and up to four 258 EOFs for each PCA bin treatment; with this scheme, radiance accuracies of 0.1% can be achieved 259 throughout the region. However, the binning scheme should be tuned to the specific application to get 260 additional computational saving, and here, we investigate the optimal set for spectral binning and the 261 number of EOFs in the Hartley and Huggins ozone bands (265-360 nm).

262 Optical properties within each bin must be strongly correlated to reduce the number of EOFs 263 required to attain a given accuracy. According to Kopparla et al. (2016), the UV region is divided at 340 264 nm, beyond which O_2 - O_2 absorption must be considered. In our application, the spectral region 340-265 360 nm is further divided at 350 nm: in the first sub-window, ozone absorption is much stronger than 266 O_2 - O_2 , while for the second (350-360 nm), O_2 - O_2 absorption becomes dominant. The binning criteria





267 are generally determined by similarities in total optical depth of gas absorption profiles τ_{ij} as defined 268 below:

269
$$\Gamma_{\rm g} = -\ln \sum_{n=1}^{N_L} \sum_{j=1}^{N_g} \tau_{nj}, \qquad (9)$$

270 where N_L and N_g denote the number of atmospheric layers and atmospheric trace gases.

To evaluate the performance of the PCA approximation, the "exact-RT" model is executed in order to calculate a fully accurate multiple-scattering spectrum using the FO model for an accurate single scattering field and VLIDORT model for an accurate multiple scattering field:

275
$$I_{\text{exact}}(\lambda_i) = I_{\text{VLD}}(\lambda_i) + I_{\text{FO}}(\lambda_i). \quad (10)$$

274

276 We first evaluate the impact of applying different binning steps and numbers of EOFs in Fig. 4. 277 where the residuals $(I_{PCA} - I_{EXACT})$ are plotted as a function of Γ_g for the spectral window 265-340 278 nm at small and large SZAs, respectively. In this evaluation, the bins are equally spaced in Γ_g for the 279 five steps from 0.20 to 1.0. For $\Gamma_g < 1$, where the extinction is strong enough that radiances are very 280 small, the residuals are effectively reduced by having more bins rather than increasing the number of 281 EOFs. In this optical range, using the first EOF is enough to capture the vast majority of the spectral 282 variance, with the optimization of the binning step. However, the bins should be narrowly spaced with 283 Γ_{g} intervals of at least 0.3-0.4 for those spectral grids for which Γ_{g} is less than -2. These spectral grids 284 are correlated with the Hartley band above ~300 nm, where radiance values rapidly increase due to 285 decreasing ozone absorption, but the spectral variations are almost unstructured. The rest of our 286 spectral region corresponds to the Huggins band above 310 nm, where spectral variations are distinctly 287 influenced by local maxima and minima of ozone absorption. In this spectral region, PCA 288 approximation errors can be greatly reduced by increasing the number of EOFs. However, it is 289 interesting to note that the PCA approximation is not further improved by using 4 EOFs instead of 3 290 (not shown here). Figure 4 also illustrates the dependence of the PCA performance on SZA in the 291 spectral range below 340 nm: For example, when 2 EOFs are applied with the binning step 0.4, errors 292 are within ± 0.02 % at smaller SZA, but increase up to ± 0.03 % at larger SZA. Therefore, as listed 293 in Table 1, two sets of binning criteria are determined to keep the accuracy within 0.05 % for any





viewing geometry. Based on the experiments shown in Fig. 5, the binning criteria are determined for the other sub-windows listed in Table 1, namely 340-350 nm and 350-360 nm: The former is set with bins at intervals of 1 and using the first two EOFs, while the latter is divided into a single bin with the first four EOFs. Figure 6 illustrates the binning criteria thus determined, demonstrating that the PCA performance keeps accuracies within 0.03 % when various sets of SZAs, ozone profiles, and vertical layers are implemented.

300 **3.2 LUT-based correction**

301 Two sets of LUTs are created, for high accuracy (LUT_H: vector/12 streams/72 layers), and low 302 accuracy (LUT_L: scalar/4 streams/24 layers) configurations. The on-line PCA-VLIDORT model is 303 configured to run in the "LUT_L" mode. The correction spectrum is straightforwardly calculated as 304 the ratio of the LUT-based spectrum (LUT_H/LUT_L), but the radiance correction term is additionally 305 adjusted to account for the different gas optical depth profiles used in on-line and LUT simulations. The 306 RT results are corrected as follows.

$$308 \qquad I_{on} = I_{on,L} \times \exp\left(\ln(I_{LUT_H}/I_{LUT_L}) + \sum_{n=1}^{N_L} \left[\left(\frac{\partial \ln I}{\partial \tau_{LUT_H}} - \frac{\partial \ln I}{\partial \tau_{LUT_L}} \right) \times (\tau_{on} - \tau_{LUT}) \right](n) \right]; \quad (11a)$$

307

309
$$\frac{\partial I}{\partial A_{son}} = \frac{\partial I}{\partial A_{son,L}} \times \frac{\frac{\partial I}{\partial A_{s}}}{\frac{\partial I}{\partial A_{s}}}; \quad (11b)$$

310
$$\frac{\partial I}{\partial \tau_{on}} = \frac{\partial I}{\partial \tau_{on,L}} \times \frac{\partial I}{\partial \tau_{LUT_H}} / \frac{\partial I}{\partial \tau_{LUT_L}}, \quad (11c)$$

311 where A_s and τ_n represent the surface albedos and gas absorption optical depths (n is the layer index). 312 To construct LUTs, RT calculations are performed using the VLIDORT version 2.8 model for sets of 313 geometrical configurations (θ , θ_o ; solar zenith angle, viewing zenith angle), surface pressures for 314 22 climatological ozone profiles and 92 wavelengths (265-345 nm) as listed in Table 2. The azimuth 315 dependence is treated exactly using the 0-2 Fourier intensity components in a Rayleigh scattering 316 atmosphere in conjunction with the associated cosine-azimuth expansion of the full intensity; see the





317 discussion below. The 22 ozone profiles are constructed from the GOME ozone profile product (Liu 318 et al. 2005), where the ozone profile shapes vary according to three latitude regimes and with the total 319 column ozone amounts at 50 DU intervals. The 92 wavelengths are regularly sampled at 5 nm intervals 320 below 295 nm and at 1.0 nm intervals up to 310 nm in the Hartley band, while irregularly sampled 321 based on the local minima and maxima of the ozone absorption structures in the Huggins band. The results of these RT calculations are separated into two components: the path radiance I_{atm} and the 322 323 surface reflectance term Isfc according to Chandrasekhar (1960), so that the following relationship 324 is employed to recover the full radiance:

325
$$I(\theta, \theta_o, \varphi - \varphi_o, A_s) = I_{atm}(\theta, \theta_o, \varphi - \varphi_o) + I_{sfc}(\theta, \theta_o, A_s).$$
(12)

326

327 I_{atm} represents the purely atmospheric contribution to the radiance in the presence of a dark surface 328 (zero albedo), and in a Rayleigh scattering atmosphere, this is given as a Fourier expansion in the cosine 329 of the relative azimuth angle.

331
$$I_{atm}(\theta, \theta_o, \varphi - \varphi_o) = I_o(\theta, \theta_o) + \cos(\varphi - \varphi_o)I_1(\theta, \theta_o) + \cos(\varphi - \varphi_o)I_2(\theta, \theta_o).$$
(13)

330

332 However, it is more convenient to write this in the form:

334
$$I_{atm} = I_o(\theta, \theta_o) \left(1 + \mathrm{aq}_1 \cos(\varphi - \varphi_o) Z_1(\theta, \theta_o) + \mathrm{aq}_2 \cos^2(\varphi - \varphi_o) Z_2(\theta, \theta_o) \right); \quad (14a)$$

333

335
$$Z_1(\theta, \theta_o) = \frac{1}{\mathrm{aq}_1} \frac{I_1(\theta, \theta_o)}{I_0(\theta, \theta_o)}; \qquad Z_2(\theta, \theta_o) = \frac{1}{\mathrm{aq}_2} \frac{I_2(\theta, \theta_o)}{I_0(\theta, \theta_o)}; \quad (14b)$$

337
$$aq_1 = \frac{3}{8}\cos\theta\sin\theta\sin\theta_o; \ aq_2 = \frac{3}{32}\frac{(\sin\theta\sin\theta_o)^2}{\cos\theta_o}.$$
 (14c)

336

In the LUTs, the three coefficients (I_o , Z_1 , and Z_2) are stored instead of I_{atm} . Note that the use of terms aq₁ and aq₂ is taken from Dave (1964); most of the angular variability in components I₁ and I₂ are captured analytically with these functions. In other words, Z_1 and Z_2 are angularly smooth and well-behaved (non-singular) functions, which helps improve





342 angular interpolation accuracy with fewer points in the angular grids. The surface term is

343
$$I_{sfc}(\theta, \theta_o, A_s) = \frac{A_s T(\theta, \theta_o)}{1 - A_s s^*}.$$
 (15)

In the LUTs, we store the transmission term $T(\theta, \theta_o)$, which is the product of the atmosphere 344 345 downwelling flux transmittance for a solar source with the upwelling transmittance from a surface illuminated isotropically from below, and the geometry-independent term s^* which is 346 347 the spherical albedo from such a surface. This is the so-called "planetary problem" calculation (Chandrasekhar, 1960), and the code to obtain T and s^* is now implemented in VLIDORT 348 Version 2.8 (Spurr, 2019). One of the key features of the VLIDORT code is its ability to 349 350 generate simultaneously (along with the Stokes vector radiation field) any set of Jacobians with 351 respect to atmospheric and surface optical properties. VLIDORT also contains an analytical 352 linearization of the planetary problem. Indeed, in our Rayleigh-based application, we require 353 Jacobians with respect to the albedo A_s and the ozone profile elements τ . First for the albedo 354 weighting function we have straightforward differentiation from Equation (15) as following

355
$$\frac{\partial I}{\partial A_s} = T(\theta, \theta_o) \left(\frac{qr}{A_s}\right)^2; \qquad qr = A_s/(1 - A_s s^*).$$
(16)

356

357 For the optical depth derivative, $\partial I/\partial \tau$ is calculated from

359
$$\frac{\partial I}{\partial \tau} = \frac{\partial I_o}{\partial \tau} + \operatorname{aq_1} \cos(\varphi - \varphi_o) \frac{\partial Z_1}{\partial \tau} + \operatorname{aq_2} \cos 2(\varphi - \varphi_o) \frac{\partial Z_2}{\partial \tau} + qr. \frac{\partial T}{\partial \tau} + T. (qr)^2. \frac{\partial s^*}{\partial \tau}. (17)$$

All partial derivatives in this expression are returned automatically by VLIDORT. For a given
ozone profile, wavelength, and surface pressure, the number of the LUT values specified in
Table 3 is 770 (nVar × n
$$\theta$$
 × n θ_0 + S_b + $\frac{dS_b}{d\tau}$, nVar = 8: I_o , $Z_1, Z_2, T, \frac{dI_o}{d\tau}, \frac{dZ_1}{d\tau}, \frac{dZ_2}{d\tau}, \frac{dT}{d\tau}$), which is
much smaller than that of a LUT with dependence on 8 relative azimuth angles and 5 surface
albedo values (11,520 =nVar × n θ × n θ_0 × n($\varphi - \varphi_0$) × nA_s, nVar = 3: I, $\partial I/\partial \tau$, $\partial I/\partial A_s$). LUT-based
simulated radiances are evaluated against on-line simulations: The LUT interpolation errors





366 are mostly less than 0.2-0.3 % (not shown here), except for extreme path length scenarios (e.g., 367 ~1% at $\theta_o = 87.0^\circ$) as shown in Fig. 7 a.b, However the interpolation errors are quite similar 368 to each other for LUT_H and LUT_L. Therefore, those errors are canceled out when performing 369 corrections using the two LUTs and thereby the overall error after LUT correction is much smaller than ~ 0.05 % (Fig. c). Note that the accuracy is completely maintained with respect to 370 both $\varphi - \varphi_o$ and A_s, while the size of a LUT is reduced by a factor of 15. However, LUT 371 372 corrections still contain ozone profile shape errors due to the use of 22 representative total 373 ozone-dependent ozone profiles in the LUT. Figure 8 shows an example of the correction 374 spectrum as a function of SZA, showing that polarization errors are mostly dominant, except 375 at the high SZAs above 310 nm, where errors due to use of a low number of streams become 376 significant, and for wavelengths below 300 nm where the use of the coarse vertical layering 377 scheme becomes the main source of uncertainty.

378 4. Evaluation

379 The PCA-RT model developed as described in this paper is implemented as the forward model 380 component of an iterative OE based inversion for retrieving ozone profile from OMI measurements. In 381 previous studies, the PCA-RT performance was evaluated against a suite of exact monochromatic 382 baseline of fully accurate VLIDORT simulations. However, such exact RT calculations cannot be 383 applied in the operational data processing system, especially when thousands of spectral points are 384 involved; in other words, the operational capability of the PCA-RT approach has been overestimated in 385 previous studies. Therefore, we evaluate the RT model developed against the existing forward model 386 where many RT approximations are applied to meet the computational budget in the operational system.

387 Table 4 contains sets of configurations for 7 forward models. OMI spectra are simulated at the under-388 sampled ("US") intervals specified in the first column of this table and then interpolated at high-389 resolution ("HR") intervals (second column) with the undersampling correction before convolution with 390 OMI slit functions. In the v1 forward model, the US spectral intervals were set at 1.0 nm/0.4 nm 391 intervals below/above 295 nm and 0.6 nm above 310 nm, while the HR spectral interval was set at 0.05 392 nm. In the updated RT model, the spectral points are selected at 0.3 nm (0.1 nm) intervals below (above) 393 305 nm and the HR interval is set as 0.03 nm, which enables us to achieve very high-accuracy, better 394 than 0.01 %, as shown in Fig. 2.c. In the reference configuration (abbreviated to "Ref"), VLIDORT is 395 run in vector mode with 12 streams and 72 atmospheric layers, so that the RT approximation errors are





396 significantly eliminated. The VLIDORT-based forward model is run with five sets of configurations 397 (abbreviated to "VLD" in Table 4) to quantify the impact of RT approximations on ozone retrievals. 398 Figure 10 compares the mean biases of the retrieved ozone profiles between VLD/PCA and Ref for 399 three SZA regimes. VLD⁰ represents the v1 forward model configuration, demonstrating that the ozone 400 retrieval errors due to the entire forward model errors range from ~ 2 % for the large SZA regime to \sim 401 5 % for the small SZA regime at the lower atmospheric layers, but ~ 1 % at the upper layers. The 402 configuration VLD¹ assesses the impact of undersampling errors on the retrievals, causing negative 403 biases of up to 2.0 % below \sim 20 km. Compared to the use of 12 streams, using 8 streams causes 404 negligible impacts on ozone retrievals (VLD²) as the corresponding RTM approximation errors are 405 negligible, except for extreme viewing geometries where the ozone retrieval errors are overwhelmed 406 by instrumental measurement errors (a few %), rather than the forward model errors of ~ 0.05 % as 407 shown in Fig. 3. The VLD³ based RT calculation is applied to ozone retrievals for evaluating on-line 408 polarization correction, showing that the corresponding errors in tropospheric ozone retrievals are 409 estimated as \pm 2 % at small SZAs. The evaluation for VLD⁴ demonstrates that the use of coarse 410 atmospheric layering causes the largest errors (~4.5 % in the troposphere, ~ 1.5 % in the stratosphere). 411 PCA^o represents the v2 forward model configuration while PCA¹ is done in the highest accurate 412 configuration except for PCA approximation. Retrieval errors due to PCA approximation are negligible 413 except for the bottom few layers at smaller solar zenith angles (up to ~ 1.5 %). Differences between 414 PCA^o and PCA¹ represents the ozone retrieval errors due to LUT errors, mostly related to the profile 415 shape errors between LUT and on-line calculations. In Fig. 10, the comparison between VLD0 (v1 416 PROFOZ) and PCA0 (v2 PROFOZ) is performed for individual ozone profile retrievals. The large systematic errors of ~ 5-15 % due to v1 forward model errors are greatly eliminated below 30 km. In 417 418 addition, the random-noise errors are significantly eliminated over the entire layers at high solar zenith 419 angles. However, there are still some remaining retrieval errors up to - 5% in troposphere and 3 % in 420 stratosphere due to v2 forward model simulation errors. Figure 11 further evaluates the v2 421 implementation. First of all, the comparison of the runtime (Fig 11.a) demonstrates a 3.3-fold-increase 422 in speed on average, thanks to switching the forward model from v1 to v2. Some spectral fit residuals 423 are eliminated in the UV 1 band over the middle area of the swath (low latitudes), where the SZAs are 424 relatively small, by up to ~2 %; the corresponding improvements are found in the stratospheric column 425 ozone. The amount of the stratospheric column ozone deviated from the reference is reduced by ~ 0.2 % 426 with v2 implementation. On the other hand, the tropospheric column ozone retrievals show 427 improvements for most cases, whereas the fit residuals of the UV2 band are slightly worse in the low 428 latitudes, but significantly better (2-5 %) in the Northern high latitudes (OMI along-track number $> \sim$





429 1300).

430 5. Summary and Conclusions

431 We have extended the PCA-based fast RT method to overcome computational challenges for OE-432 based SAO OMI ozone profile retrievals from ultraviolet measurements requiring iterative calculations 433 of the radiance and its Jacobian derivatives, to match the simulated spectrum to the measured spectrum 434 The PCA-RT model is designed to perform MS calculations for a few EOF-derived optical states which 435 are developed from spectrally binned sets of inherent optical properties that possess some redundancy. 436 To maximize the performance enhancement, we carefully tuned the binning scheme for the UV ozone 437 fitting window from 265 nm to 360 nm in such a way as to choose the number of EOFs to be as small 438 as possible for each bin, rather than always using the first four EOFs for all bins selected in previous 439 studies. The spectral windows are divided into three sub-windows: 1) 265-340 nm, 2) 340-350, and 3) 440 350-360 nm. Then, optical profiles are grouped into bins according to criteria based on the total gas 441 optical depth, as specified in Table 1. Spectral bins correlated to the Hartley ozone band use only the 442 first EOF, but 2 or 3 EOFs are required for those bins related to the Huggins ozone band. The MS model 443 is executed 85 times for the entire wavelength range (265-360 nm), and only 51 times for the OMI 444 ozone fitting window (270-330 nm). We demonstrated that the PCA approximation errors are within 445 0.03 % for any viewing geometry, optical depth profile, and vertical layering. The existing (v1) forward 446 model calculations are evaluated to determine the optimal configuration for the v2 forward model. RT 447 approximation errors exist due to the use of 24 quite coarse vertical layers (2.5 km thick), which can 448 cause radiance simulation errors of up to ~ 1 % below 320 nm and this leads to ozone retrieval errors of 2-449 4 % in the troposphere and 1 % in the stratosphere. Eight-stream calculations can result in radiance residuals 450 of ~ 0.05 % or less except at extreme viewing geometries, which causes trivial errors on ozone retrievals 451 compared to other error factors. In spite of accounting for polarization errors using vector and scalar 452 differences at 14 wavelengths, the retrieval accuracies are systematically worse by ~ 2 % due to neglecting 453 second-order polarization errors which are strongly correlated with ozone absorption features. We found that 454 72 atmospheric layers (~ 0.7 km thick) and 12 streams should be used at least for fully accurate RT 455 calculations comparable to those with 99 atmospheric layers and 32 streams. The OMI spectral fit uses 229 456 wavelengths at OMI native grids, but the existing RT model simulates 93 wavelengths and is then 457 interpolated onto 0.05 nm grids with undersampling correction. However, we found room to improve our 458 retrievals (~ 1.5 % on avearge) by simulating 244 wavelength grids selected at intervals of 0.3 nm/0.1 nm 459 below/above 305 nm and then performing the undersampling correction to 0.03 nm. Applying the PCA-RT 460 approach allows us to reduce the number of MS calculations from the high-resolution optical dataset to 51





461 sets of EOF-derived optical states, but the performance savings are not enough to improve over previous RT 462 approximations. To improve both efficiency and accuracy, we have developed a LUT-based correction for 463 eliminating the RT approximation errors arising from the vector vs scalar, 12 vs. 4 streams, and 72 vs. 24 464 layers. In conclusion, the updated PCA-based RT model combined with LUT corrections makes ozone 465 profile retrievals faster than the v1 forward model by a factor of 3.3 on average. Improvements in fitting accuracies are also achieved in the UV1 band by 2 % and in the UV2 band by 2-5 %. Correspondingly, 466 467 the ozone profile retrievals are significantly improved, especially in the troposphere by \sim up to 10 %. 468 However, there are still some remaining retrieval errors of up to - 5% in troposphere and 3 % in 469 stratosphere due to the LUT correction errors and PCA approximation errors in the v2 implementation. 470 The updated forward model is in preparation for reprocessing all OMI measurements (2004 -current) 471 for the next version of the PROFOZ product.

472

- *Author contributions.* JB and XL designed the research; RS provided oversight and guidance for using
 both VLIDORT and PCA-based VLIDORT; KY developed the LUT creation and interpolation scheme;
 XL contributed to analyzing ozone profile retrievals with different forward model approaches; JB
 conducted the research and wrote the paper; CN, CM, GA, and KC contributed to the analysis and
 writing; CM and GA contributed to managing the computational resources.
- 478 *Competing interests.* The authors declare that they have no conflicts of interest.

479 Data availability. OMI Level1b radiance datasets are available at
480 <u>https://aura.gesdisc.eosdis.nasa.gov/data/Aura_OMI_Level1/</u> (last access: 31 AUG 2020). The LUT
481 database are attainable upon request.

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References

491 492 493 494 495 496 497 498	 Chandrasekhar, S.: Radiative Transfer, Dover Publications, Mineola, New York, 1960. Cuesta, J., Eremenko, M., Liu, X., Dufour, G., Cai, Z., Höpfner, M., von Clarmann, T., Sellitto, P., Foret, G., Gaubert, B., Beekmann, M., Orphal, J., Chance, K., Spurr, R., and Flaud, JM.: Satellite observation of lowermost tropospheric ozone by multispectral synergism of IASI thermal infrared and GOME-2 ultraviolet measurements over Europe, Atmos. Chem. Phys., 13, 9675–9693, https://doi.org/10.5194/acp-13-9675-2013, 2013. Dave, J.V.: Meaning of Successive Iteration of the Auxiliary Equation in the Theory of Radiative Transfer, The Astrophysical Journal, 140, 1292-1303, 1964.
499	Efremenko D. S., Loyola, D. G., Spurr, R., and Doicu, A.: Acceleration of Radiative Transfer Model
500	Calculations for the Retrieval of Trace Gases under Cloudy Conditions. J. Quant. Spectrosc. Radiat.
501	Transfer., 135, 58-65, 2014.
502	Fu, D., Kulawik, S. S., Miyazaki, K., Bowman, K. W., Worden, J. R., Eldering, A., Livesey, N. J.,
503	Teixeira, J., Irion, F. W., Herman, R. L., Osterman, G. B., Liu, X., Levelt, P. F., Thompson, A. M.,
504	and Luo, M.: Retrievals of tropospheric ozone profiles from the synergism of AIRS and OMI:
505	methodology and validation, Atmos. Meas. Tech., 11, 5587-5605, https://doi.org/10.5194/amt-11-
506	5587-2018, 2018.
507	Fu, D., Worden, J. R., Liu, X., Kulawik, S. S., Bowman, K. W., and Natraj, V.: Characterization of ozone
508	profiles derived from Aura TES and OMI radiances, Atmos. Chem. Phys., 13, 3445-3462,
509	https://doi.org/10.5194/acp-13-3445-2013, 2013.
510	Kim, P. S., Jacob, D. J., Liu, X., Warner, J. X., Yang, K., Chance, K., Thouret, V., and Nedelec, P.: Global
511	ozone–CO correlations from OMI and AIRS: constraints on tropospheric ozone sources, Atmos.
512	Chem. Phys., 13, 9321–9335, https://doi.org/10.5194/acp-13-9321-2013, 2013.
513	Kopparla, P., Natraj, V., Limpasuvan, D., Spurr, R., Crisp, D., Shia, R.L., Somkuti, P., Yung, Y.L.; PCA-
514	based radiative transfer: Improvements to aerosol scheme, vertical layering and spectral binning.
515	J. Quant. Spectrosc. Radiat. Transf., 198, 104–111, 2017.
516	Kopparla, P., Natraj, V., Spurr, R., Shia, R.L., Crisp, D., Yung, Y.L.; A fast and accurate PCA based
517	radiative transfer model: Extension to the broadband shortwave region. J. Quant. Spectrosc. Radiat.
518	Transf., 173, 65–71, 2016.
519 520 521 522 523 524	 Kroon, M., Petropavlovskikh, I., Shetter, R., Hall, S., Ullmann, K., Veefkind, J. P., McPeters, R. D., Browell, E. V., and Levelt, P. F.: OMI total ozone column validation with Aura-AVE CAFS observations, J. Geophys. Res., 113, D15S13, doi:10.1029/2007JD008795, 2008. Liu, C., Yao, B., Natraj. V., Kopparla, P., Weng., F., Le., T., Shia., R., and Yung, Y. L.: A spectral Data Compression (SDCOMP) Radiative Transfer Model for High-Spectral-Resolution Radiation Simulations, J. Atmos., Sci., 77, 2055-2066, https://doi.org/10.1175/JAS-D-19-0238.1, 2020.
525 526	Liu, X., Bhartia, P.K, Chance, K, Spurr, R.J.D., and Kurosu, T.P.: Ozone profile retrievals from the ozone monitoring instrument. Atmos. Chem. Phys., 10, 2521–2537, 2010.
527	Liu, X., Yang, Q., Li, H., Jin, Z., Wu, W., Kizer, S., Zhou., D. K., and Yang., P.; Development of a fast
528	and accurate pcrtm radiative transfer model in the solar spectral region. Applied optics, 55(29),
529	8236–8247., 2016.
530 531 532 533 534 535	 Loyola, D. G., Gimeno García, S., Lutz, R., Argyrouli, A., Romahn, F., Spurr, R. J. D., Pedergnana, M., Doicu, A., Molina García, V., and Schüssler, O.: The operational cloud retrieval algorithms from TROPOMI on board Sentinel-5 Precursor, Atmos. Meas. Tech., 11, 409–427, https://doi.org/10.5194/amt-11-409-2018, 2018. Miles, G. M., Siddans, R., Kerridge, B. J., Latter, B. G., and Richards, N. A. D.: Tropospheric ozone and ozone profiles retrieved from GOME-2 and their validation, Atmos. Meas. Tech., 8, 385–398,





36	https://doi.org/10.5194/amt-8-385-2015, 20	015

537	Nanda, S., de Gra	af, M., Ve	efkind, J.	P., ter Linder	ı, M., Snee	ep, M., de	e Haan, J.,	and Lev	velt, P. F.: A
538	neural netwo	rk radiati	ve transf	er model ap	oroach app	olied to	the Tropo	spheric	Monitoring
539	Instrument	aerosol	height	algorithm,	Atmos.	Meas.	Tech.,	12,	6619-6634,
540	https://doi.org	g/10.5194/	amt-12-60	519-2019, 201	9.				

- Natraj, V., Jiang, X., Shia, R.-L., Huang, X., Margolis, J. S., and Yung, Y. L.; Application of principal component analysis to high spectral resolution radiative transfer: A case study of the O₂-A band, *J. Quant. Spectrosc. Radiat. Transf.*, 95(4), 539–556. <u>https://doi.org/10.1016/j.jqsrt.2004.12.024</u>, 2005.
- Natraj, V., Shia, R. L., and Yung, Y. L.;. On the use of principal component analysis to speed up radiative
 transfer calculations, *J. Quant. Spectrosc. Radiat. Transf.*, 111(5), 810–816.
 https://doi.org/10.1016/j.jqsrt.2009.11.004, 2010.
- Somkuti, P., Boesch, H., Vijay, N., and Kopparla, P.; Application of a PCA-based fast radiative transfer
 model to XCO2 retrievals in the shortwave infrared, *J. Geophys. Res*: Atmospheres, 122, 10,477–
 10,496. https://doi.org/10.1002/2017JD027013., 2017.
- Rodgers, C. D.; Inverse methods for atmospheric sounding—Theory and practice. In C. D. Rodgers
 (Ed.), Series on atmospheric oceanic and planetary physics (Vol. 2). Oxford: World Scientific
 Publishing Co. Pte. Ltd. <u>https://doi.org/10.1142/9789812813718</u>, 2000.
- Spurr, R., Natraj, V., Lerot, C., Van Roozendael, M., and Loyola, D.; Linearization of the principal component analysis method for radiative transfer acceleration: Application to retrieval algorithms and sensitivity studies, *J. Quant. Spectrosc. Radiat. Transf.*, 125, 1–17. https://doi.org/10.1016/j.jqsrt.2013.04.002, 2013.
- Spurr, R., Natraj, V., Kopparla, P., and Christi, M.; Application of Principal Component Analysis (PCA)
 to Performance Enhancement of Hyperspectral Radiative Transfer Computations, in "Principal
 Component Analysis: Methods, Applications and Technology", NOVA publishers, 2016
- Spurr, R.; User's Guide VLIDORT version 2.8.1, RT Solutions, Inc., 9 Channing Street, Cambridge,
 MA 02138, USA., 2019.
- Spurr, R. and Christi, M. : The LIDORT and VLIDORT Linearized Scalar and Vector Discrete Ordinate
 Radiative Transfer Models: An Update for the last 10 Years. Light Scattering Reviews, Volume 12,
 ed. A. Kokhanovsky, Springer, 2019.
- Yang, Q., Liu, X., Wu, W., Kizer, S., and Baize, R. R.: Fast and accurate hybrid stream PCRT
 M-SOLAR radiative transfer model for reflected solar spectrum simulation in the cloudy atm
 osphere, Opt. Express 24, A1514-A1527, 2016.
- Zoogman, P., Liu, X., Suleiman, R. M., Pennington, W. F., Flittner, D. E., Al-Saadi, J. A., Hilton, B. B., 569 570 Nicks, D. K., Newchurch, M. J., Carr, J. L., Janz, S. J., Andraschko, M. R., Arola, A., Baker, B. D., 571 Canova, B. P., Chan Miller, C., Cohen, R. C., Davis, J. E., Dussault, M. E., Edwards, D. P., Fishman, 572 J., Ghulam, A., González Abad, G., Grutter, M., Herman, J. R., Houck, J., Jacob, D. J., Joiner, J., 573 Kerridge, B. J., Kim, J., Krotkov, N. A., Lamsal, L., Li, C., Lindfors, A., Martin, R. V., McElroy, 574 C. T., McLinden, C., Natraj, V., Neil, D. O., Nowlan, C. R., O'Sullivan, E. J., Palmer, P. I., Pierce, 575 R. B., Pippin, M. R., Saiz-Lopez, A., Spurr, R. J. D., Szykman, J. J., Torres, O., Veefkindz, J. P., Veihelmann, B., Wang, H., Wang, J., and Chance, K.: Tropospheric Emissions: Monitoring of 576 Spectrosc. 577 (TEMPO), J. Quant. Pollution Radiat. Transf., 186. 17-39. 578 https://doi.org/10.1016/j.jqsrt.2016.05.008, 2017.
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- 585 Table 1. The PCA-RT configuration optimized over the UV spectral range 265-360 nm. The
- 586 optical depth of the total gas column (Γ_g defined in eq. 9) is used to set the criteria for the spectral
- 587 binning; for example, one or more bins are created at intervals $(\Delta \Gamma_g)$ in the range Γ_g^{min} to Γ_g^{max} .
- 588 For each bin, the optical states are expanded in terms of the first few number of EOFs (nEOF).

	265-340 nm								
SZA or VZA < 70 °					SZA or VZA $\geq 70^{\circ}$				
List	$\Gamma_{\rm g}^{lower}, \ \Gamma_{\rm g}^{upper}$	$\Delta\Gamma_{g}$	nEOF	List	$\Gamma_{\rm g}^{lower}, \ \Gamma_{\rm g}^{upper}$	$\Delta\Gamma_{\rm g}$	nEOF		
1	∞ to -1.7	2	1	1	∞ to - 1.5	2.0	1		
2	-1.7 to -1.2	0.5	1	2	-1.5 to -0.7	1.2	1		
3	-1.2 to 0.0	0.4	1	3	-0.7 to 0.4	0.35	1		
4	0.0 to 0.5	0.5	1	4	0.4 to 0.7	0.3	1		
5	0.5 to 3.5	0.6	2	5	0.7 to 2.5	0.6	3		
6	3.5 to 4.5	1.0	2	6	2.5 to 3.5	1.0	3		
7	4.5 to ∞	2.0	2	7	3.5 to 4.5	1.0	2		
8				8	4.5 to ∞	2.0	2		
340-350 nm 350-360 nm									
List	$\Gamma_{\rm g}^{lower}, \ \Gamma_{\rm g}^{upper}$	$\Delta\Gamma_{\rm g}$	nEOF	List	$\Gamma_{\rm g}^{lower}, \Gamma_{\rm g}^{upper}$	$\Delta\Gamma_{g}$	nEOF		
1	∞ to ∞	1.0	2	1	∞ to ∞	8	4		

589

590 Table 2. LUT parameter specification. Note that the relative azimuth dependence is taken

591 into account explicitly through the Fourier coefficients of path radiance (Table 3) and the

592	surface albedo	dependence is ta	iken into account l	bv the planeta	rv problem.

Parameter	Symbo l	Ν	Grid Values
Ozone Profile ⁺	0 ₃ P	22	 Low-latitude (30°S-30°N) L200,L250,L300,L350 Mid-latitude (30°-60°N/S) M200,M250,M300,M350, M400,M450, M500,M550 High-latitude (60°-90°N/S) H100,H150,H200,H250,H300,H350, H400,H450, H500,H550
Wavelength	ength λ 92		265-345 nm
Solar Zenith Angle	θο,	12	0, 16, 31, 44, 55, 64, 71, 76.5, 80.5, 83.5, 86, 88°
Viewing Zenith Angle	θ	8	0, 15, 30, 43, 53, 61, 67, 72°
Surface albedo	A _s	1	0.0
Surface pressure	P_s	12	100, 150, 200, 300, 400, 500, 600, 700, 800, 900, 1013.25, 1050 hPa

⁺Total ozone-based ozone profiles for three latitude regimes. The grid values represent the





594 amount of total ozone (DU).

595

596 Table 3. LUT variable specification

Table 5. LUT variable specification								
Variable	Dimensions	Variable	Dimensions					
I _o ^a	$n\lambda, n\theta, n\theta_o, nP_s$	$dI_o/d\tau$	$n\lambda, n\theta, n\theta_o, nz, nP_s$					
Z_1^a	$n\lambda, n\theta, n\theta_o, nP_s$	$dZ_1 d\tau$	$n\lambda, n\theta, n\theta_o, nz, nP_s$					
$\boldsymbol{Z_2}^{\mathrm{a}}$	$n\lambda, n\theta, n\theta_o, nP_s$	$dZ_2/d\tau$	$n\lambda, n\theta, n\theta_o, nz, nP_s$					
T^{b}	$n\lambda, n\theta, n\theta_o, nP_s$	dT/dτ	$n\lambda, n\theta, n\theta_o, nz, nP_s$					
S _b ^c	$n\lambda$, nP_s	$dS_b/d\tau$	$n\lambda$, nz , nP_s					
$oldsymbol{ au}^{ m d}$	$n\lambda$, nz^+							

^a Fourier coefficients of path radiance with respect to relative azimuth angle

^b total transmission of the atmosphere

^c spherical albedo of the atmosphere

600 ^d total gas absorption optical depth profile

601 ⁺ Number of atmospheric layers

602

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604	Table 4. List of configurations used in evaluating the different forward model
605	calculations for OMI ozone profile retrievals. The reference, VLIDORT, and PCA-RT
606	models are abbreviated as Ref, VLD, and PCA, respectively.

RT models	US SI (nm) ^a	HR SI (nm) ^b	Nstream ^c	Nlayer ^d	Polarization ^e	RT corr ^f
Ref	0.3(<305nm) 0.1 (≥305 nm)	0.03	12	72	True	False
VLDº	1.0(<295nm) 0.4(between) 0.6 (≥310 nm)	0.05	8	24	False	On-line polcorr
VLD ¹	1.0(<295nm) 0.4(between) 0.6 (≥310 nm)	0.05	12	72	True	False
VLD ²	0.3(<305nm) 0.1 (≥305 nm)	0.03	8	72	True	False
VLD ³	0.3(<305nm) 0.1 (≥305 nm)	0.03	12	72	False	On-line polcorr
VLD ⁴	0.3(<305nm) 0.1 (≥305 nm)	0.03	12	24	True	False
PCA ⁰	0.3(<305nm) 0.1 (≥305 nm)	0.03	4	24	False	LUT
PCA ¹	0.3(<305nm) 0.1 (≥305 nm)	0.03	12	72	True	False

^a Under-sampled (US) spectral intervals (nm) used to define wavelengths at which RT is actually executed.

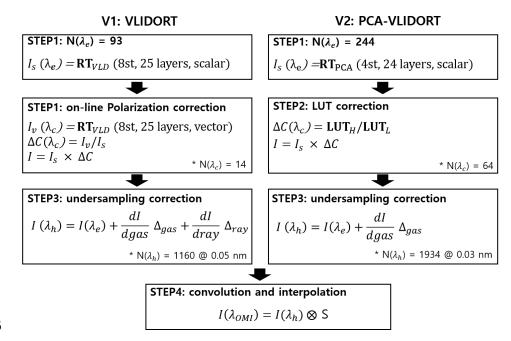
 ⁶⁰⁷^bHigh-resolution (HR) spectral intervals (nm) used to define wavelengths where under-sampled simulations are interpolated before spectral convolution.





- 610 ^c the number of discrete ordinates in the full polar space; ^dNumber of atmospheric layers
- 611 °RT model is run in the vector (scalar) mode if polarization is true (false).
- 612 ^fOn-line polarization correction as described in Section 2, which is originally developed from Liu et al. (2010).
- 613 LUT-based correction introduced in Section 3.3, which is developed in this work to account for RT approximation
- 614 errors due to neglecting polarization as well as using 4 streams and 24 layers.

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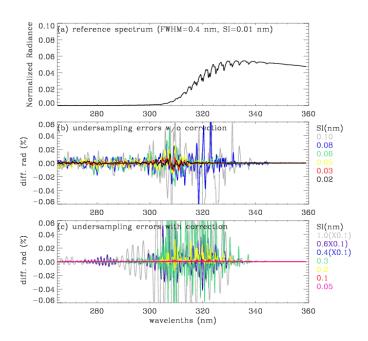


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617 **Fig.1**. Schematic flowcharts of VLIDORT (v1) and PCA-VLIDORT (v2) based forward 618 models, respectively. Note that VLIDORT was used in the generation of the OMPROFOZ v1 619 dataset, while PCA-VLIDORT is in preparation for OMPROFOZ v2 production. The number 620 of wavelengths used in each process is denoted as N(λ) when the spectral window 270-330 621 nm is applied. λ_e represents the wavelength grids used for RT calcluation, while λ_c and λ_h 622 are grids used in RT approximation correction and undersampling correction, respectively. 623 See text for definition of other variables.







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626 Fig.2. (a) Reference (truth) normalized radiance spectrum simulated at the spectral intervals 627 (SIs) of 0.01 nm in 265-360 nm (solar zenith angle = 65° , viewing zenith angle = 30° , relative 628 azimuth angle = 120°), which is used for evaluating the simulations in Figs. (b) and (c). (b) 629 Impact of under-sampling on the simulation. (c) is similar to (b), but now the under-sampling correction has been applied. In Fig 2(c), the under-sampling errors are divided by 10 at SIs \geq 630 0.4 nm. Note that individual radiances simulated at different SIs are interpolated to 0.01 nm 631 632 and then convolved with the Gaussian function (FWHM: 0.4 nm) which represents the OMI 633 instrument spectral response function.

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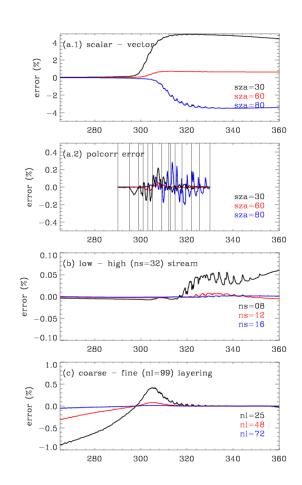
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Fig. 3. Errors of the radiance simulation due to the RT approximation used in v1, arising from (a.1) neglecting the
polarization effect for different solar zenith angles (sza), (a.2) polarization correction errors, (b) using a low
number of streams (ns), and (c) using a coarse number of vertical grids (nl). Note that vertical lines in Fig. 3.a.2
indicate wavelengths used in deriving the on-line polarization correction spectrum.

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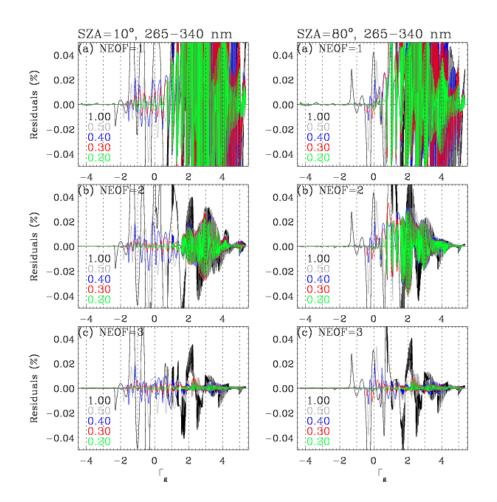


Fig. 4. Residuals (%) of the PCA-RT radiance in the wavelength range 265-340 nm compared to the exact-RT calculations, for different binning steps (different colors) and number of EOFs (a, b, c). Results are plotted as a function of Γ_g (logarithm of the total gas optical depth), for solar zenith angles (SZAs) of (left)10° and (right) 80°.





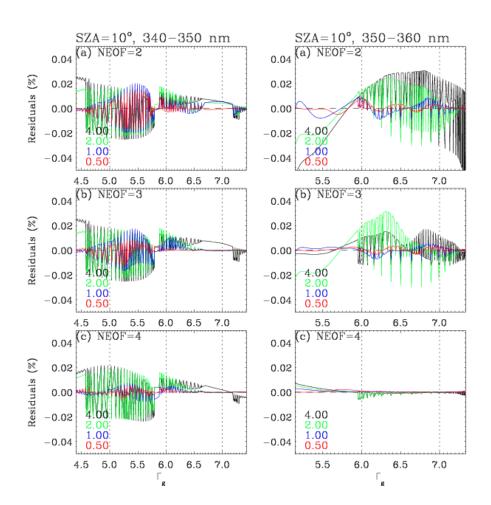


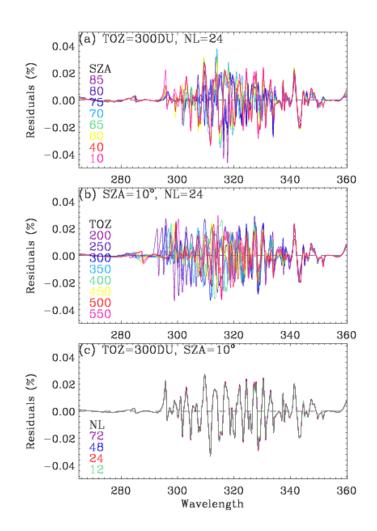


Fig. 5. Same as Fig. 4, but for different windows, (left) 340-350 nm and (right) 350-360 nm, respectively.

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Fig. 6. Residuals (%) of the PCA-RT radiances with the binning scheme given in Table 1, for
various sets of (a) solar zenith angles, (b) ozone profiles with different total ozone columns
(TOZs), and (c) number of atmospheric layers.





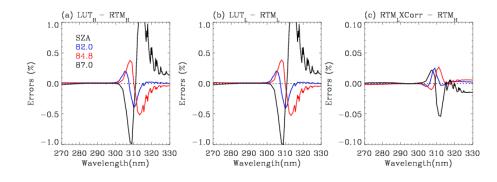
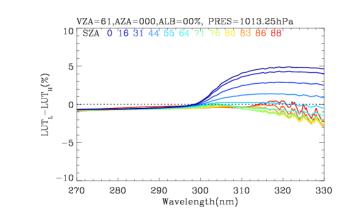


Fig. 7. Evaluation of simulations with respect to extreme SZAs at VZA = 61° , AZA = 0° , ALB=0 %, and surface pressure = 1013.25 hPa. LUT and RTM represent LUT and on-line radiative transfer model (RTM) based calculations, respectively, with the subscripts H and L indicating high and low accuracy configurations. RTM_L X Corr radiances are simulated using on-line RTM with low accuracy configuration, but corrected using LUT_H/LUT_L.

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Fig. 8. Example of LUT-based correction spectrum.





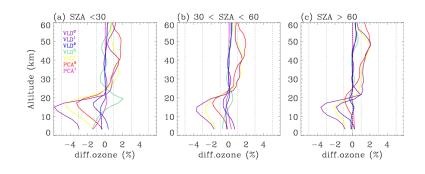
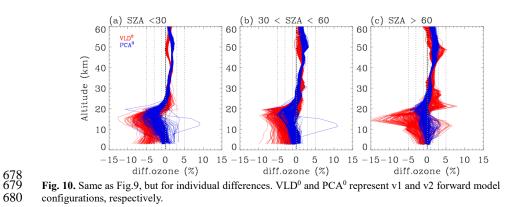


Fig. 9. Mean biases of ozone profile retrievals with different configurations compared to thosewith the reference configuration. Each configuration is given in Table 4.

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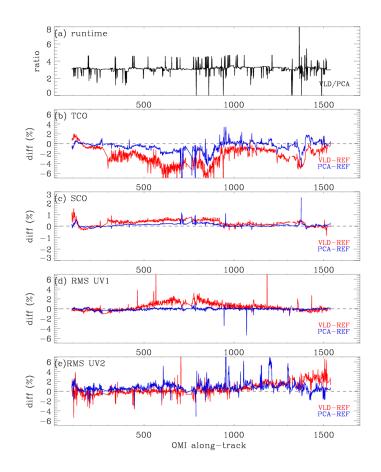
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Fig. 11. Same as Fig. 10, but for (a) runtime, (b) tropospheric column ozone (TCO), (c)
stratospheric column ozone (SCO), and (d) UV1 (270-310 nm)/(e) UV2 (310-330 nm) fitting
residuals, along with the OMI along-track position (1-1644) at nadir cross-track. Note that the
fitting residuals are estimated as root mean square (RMS) errors for differences between
measured and simulated spectra relative to the measurement error. VLD and PCA represent
v1 and v2 forward model configurations, respectively.

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