

Error analyses of a multistatic meteor radar system to obtain a 3-dimensional spatial resolution distribution

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Abstract: In recent years, the concept of multistatic meteor radar systems has attracted the attention of the atmospheric radar community, focusing on the mesosphere and lower thermosphere (MLT) region. Recently, there have been some notable experiments using such multistatic meteor radar systems. Good spatial resolution is vital for meteor radars because nearly all parameter inversion processes rely on the accurate location of the meteor trail specular point. It is timely then for a careful discussion focussed on the error distribution of multistatic meteor radar systems. In this study, we discuss the measurement errors that affect the spatial resolution and obtain the spatial resolution distribution in 3-dimensional space for the first time. The spatial resolution distribution can both help design a multistatic meteor radar system and improve the performance of existing radar systems. Moreover, the spatial resolution distribution allows the accuracy of retrieved parameters such as the wind ~~fields~~ field to be determined.

1 Introduction

The mesosphere and lower thermosphere (MLT) is a transition region from the neutral to the partially ionized atmosphere. It is dominated by the effects of atmospheric waves, including planetary waves, tides and gravity waves. It is also a relatively poorly sampled part of the Earth's atmosphere by ground-based instruments. One widely used approach to sample this region is the meteor radar technique. The ablation of incoming meteors in the MLT region, i.e., $\sim 80 - 110$ km, creates layers of metal atoms, which can be observed from the ground by photometry or lidar (Jia et al., 2016; Xue et al., 2013). During meteor ablation, the trails caused by small meteor particles provide a strong atmospheric tracer within the MLT region that can be continuously detected by meteor ~~radars~~ radars, regardless of weather conditions. Consequently, the meteor radar technique has been a powerful tool for studying the MLT region for decades (Hocking et al., 2001; Holdsworth et al., 2004; Jacobi et al., 2008; Stober et al., 2013; Yi et al., 2018). Most modern meteor radars are monostatic, and this has two main limitations in

retrieving the complete wind fields. Firstly, limited meteor rates and relatively low measurement accuracies necessitate that all measurements in the same height range are processed to calculate a “mean” wind. Secondly, classic monostatic radars retrieve ~~windwinds~~ based on the assumption of a homogenous wind in ~~the~~ horizontal and ~~usually~~ a zero wind in ~~the~~ vertical direction.

The latter conditions can be partly relaxed if the count rates are high and the detections are distributed through a representative range of azimuths. If this is the case, a version of a Velocity Azimuth Display (VAD) analysis can be applied by expanding the zonal and meridional winds using a truncated Taylor expansion (Browning and Wexler, 1968). This is because each valid meteor detection yields a radial velocity in a particular ~~lookviewing~~ direction of the radar. The radar is effectively a multi-

beam Doppler radar where the “beams” are determined by the meteor detections. If there are enough suitably distributed detections in azimuth in a given observing period, the Taylor expansion approach using cartesian coordinates yields the mean zonal and meridional wind components (u_0, v_0) , the horizontal divergence $\left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y}\right)$, the stretching $\left(\frac{\partial u}{\partial x} - \frac{\partial v}{\partial y}\right)$ and the shearing $\left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x}\right)$ deformations of the wind fields from an analysis of the radial velocities. However, because the radar can only retrieve the wind projection in the radial direction as measured from the radar, the vorticity $\left(\frac{\partial v}{\partial x} - \frac{\partial u}{\partial y}\right)$ of the wind fields

is not available. This is common to all monostatic radar systems and a discussion of measurable parameters in the context of multiple fixed beam upper atmosphere Doppler radars is given by (Reid, 1987). Even by relaxing the assumption of a homogeneous wind fields and using the more advanced Volume Velocity Processing (VVP) (Philippe and Corbin, 1979) to retrieve the wind fields, the horizontal gradients of the wind fields cannot be recovered due to the lack of vorticity information.

To obtain a better understanding of the spatial variation of the MLT region wind fields, larger area observations (and hence higher meteor count rates) and ~~measurements of sampling~~ the ~~non homogenous wind fields~~ ~~observed area from different viewing angles~~ are needed. An extension of the classic monostatic meteor technique is required to satisfy these needs.

To resolve the limitations outlined above, the concept of multistatic meteor radar systems, such as MMARIA (multi-static and multi-frequency agile radar for investigations of the atmosphere) (Stober and Chau, 2015) and SIMO (single input multiple output) (Spargo et al., 2019), MIMO (multiple input multiple output radar) (Chau et al., 2019; DOREY et al., 1984) have been

designed and implemented (Stober et al., 2018). Multistatic systems can utilize the forward scatter of meteor trails, thus providing another perspective for observing the MLT. Multistatic meteor radar systems have ~~many~~several advantages over classic monostatic meteor radars, such as obtaining higher-order wind ~~fields~~field information and covering wider observation areas. There have been some particularly innovative studies using multistatic meteor radar systems in recent years. For example,

by combining MMARIA and the continuous wave multistatic radar technique (Vierinen et al., 2016), Stober ~~and Chau et al.~~
~~built a 5 station total 7 link multistatic radar network covering an approximately 600 km×600 km region in Germany to retrieve an arbitrary non homogenous wind fields with a 30 km×30 km horizontal resolution (Stober et al., 2018).~~ Chau et al.
~~(2018) built a 5-station 7-link multistatic radar network covering an approximately 600 km×600 km observing region over Germany to retrieve an arbitrary non-homogenous wind field with a 30 km×30 km horizontal resolution. Chau et al.~~

(2017) used two adjacent classic monostatic specular meteor radars in northern Norway to obtain horizontal divergence and vorticity (Chau et al., 2017). Other approaches, such as coded continuous wave meteor radar (Vierinen et al., 2019) and the compressed sense method in MIMO sparse signal recovery (Urco et al., 2019), are described in the corresponding references in these papers.

Analysing spatial resolution in interested areas limits is a fundamental but difficult topic for meteor radar systems. Meteor radar systems transmit radio waves and then receive radio waves reflected from meteor trails using a cluster of receiver/receiving antennas; commonly five antennas as in the Jones et al. (1998) configuration (Jones et al., 1998). By analysing the cross correlations of the signals received signals, one can determine on several pairs of antennas, the angle of arrivals (AoAs) which includes arrival (AoA) of each return can be determined. The AoA is described by the zenith angle θ and azimuth angle denoted as ϕ respectively. By measuring the wave propagation time, one can obtain from the meteor trail, range information can be determined. Most meteor radar systems rely on specular reflections from meteor trails. Thus, by combining the AoAs/AoA and the range information and then using geometric analysis, one can determine the location of a meteor trail trail can be determined. Accurately locating the meteor trail specular point (MTSP hereafter hereinafter) is important since atmospheric parameter retrieval (such as the wind fields field or the temperature) depends on the location information of meteor trails. The location accuracy, namely the spatial resolution, determines the reliability of the retrieved parameters. For multistatic meteor radar systems that which can relax the assumption of a homogenous horizontal wind fields field, the location accuracy becomes a more important issue because the horizontal spatial resolution affects the accuracy of the retrieved horizontal wind fields gradient field.

There are some discussions about measuring Although meteor radar systems have developed well experimentally in recent years, the reliability of the retrieved atmospheric parameters still requires further investigation for both the monostatic and multistatic meteor radar cases. In an attempt to investigate errors of the meteor radar. For example, a in two radar techniques, Wilhelm et al. (2017) compared 11 years of MLT region wind data from a partial reflection (PR) radar with collocated monostatic meteor radar winds and determined the 'correction factors' to bring the winds into agreement. Reid et al. (2018) reported a similar study for two locations for data obtained over several years. While the comparisons are interesting, partial reflection radars operating in the medium frequency (MF) and lower high frequency (HF) bands produce a height dependent bias in the measured winds (see e.g., Reid, 2015) which limits the ability to estimate errors in the meteor winds by comparing with them. However, the PR radar technique is one of very few that provides day and night coverage and data rates in the MLT comparable to that of meteor radars.

Meteor radars have largely replaced PR radars for MLT studies and are generally regarded as providing reference quality winds. It is essential then to know the reliability of atmospheric parameters determined by meteor radars and to do this, some quantitative error analyses are necessary.

A number of recent studies have discussed AoAs measuring AoA measurement errors for meteor radars (Kang, 2008; Vaudrin et al., 2018; Younger and Reid, 2017). However, those These studies focus on the phase errors in receiver antennas antenna pairs; Younger and seldom discuss Reid for the influence of monostatic case, and Vaudrin et al. for a more general case which

100 included multistatic configuration on the spatial resolutions meteor radars. Hocking (2018) used another approach and developed a vertical resolution analysis method ~~in a~~ for the 2-dimensional baseline vertical section (Hocking, 2018), which ~~bistatic case. Hocking's method (HM hereinafter)~~ simplifies the error propagation process in ~~receiver~~ the receiving antennas and ~~puts~~ emphasis on how a bistatic meteor radar configuration ~~would affect~~ affects the vertical resolution in a vertical section. ~~However, Hocking's method (HM hereinafter) can barely show bistatic configurations' influence on spatial resolution distribution due to ignore the discussion of It does not consider the~~ radial distance measuring error. ~~Moreover, HM is only a demo about vertical resolution in a specific vertical section, not in real three dimensional space. Hence, for practical purposes,~~ the ~~In this paper, we consider the more general~~ 3-dimensional case and determine the spatial distribution of both the horizontal resolution and vertical resolution ~~should be considered~~ ~~uncertainties~~.

105 ~~Although multistatic meteor radar systems have developed well experimentally in recent years, the reliability of retrieved atmospheric parameters lacks discussion both for monostatic and multistatic meteor radar. A large part of the reason is that no other measurement technology can provide contrast data for meteor radars in MLT region. On the one hand it proves that meteor radars are irreplaceable in MLT region as a measurement technology; on the other hand, to know the reliability of meteor radars obtained atmospheric parameters and to get better understanding of the dynamic process in MLT region, some quantitative error analyses are necessary and helpful. In this paper, we~~ We analyse the multistatic meteor radar resolution distribution in a three-dimensional space for both vertical and horizontal resolution for the first time. ~~And~~ This spatial resolution is a prerequisite for evaluating the reliability of retrieved atmospheric parameters, such as ~~the~~ wind ~~fields~~ field and ~~the~~ temperature.

2 Analytical Method

2.1 ~~brief~~ Brief introduction

110 The HM will be introduced ~~ahead in short~~ ~~briefly~~ here to help understand our ~~method, generalization~~. In the HM, ~~measuring~~ ~~measurement~~ errors that affect the vertical resolution can be classified into two types (Hocking, 2018): one ~~is those~~ ~~that~~ caused by the zenith angle measuring error $\delta\theta$ and ~~another is those that~~ ~~one~~ caused by the pulse-length effect on the vertical resolution. The ~~receiver~~ ~~receiving~~ array is a simple antenna pair that is collinear to ~~with~~ the baseline (figure 1). ~~The~~ HM ~~only calculate~~ ~~calculates~~ the vertical resolution in a two-dimensional vertical section which ~~pass through~~ ~~passes through~~ the baseline. The receiver antenna pair is equivalent to one receiver arm in a Jones configuration which is comprised of three collinear antennas ~~and is~~ usually in a 2λ \ 2.5λ configuration. Phasespacing. The phase difference of the received radio wave between the receiving antenna pair is denoted as $\Delta\Psi$. In meteor radar systems, there ~~are~~ is ~~generally~~ an 'acceptable' phase difference measuring error (PDME ~~hereafter~~ ~~hereinafter~~) $\delta(\Delta\Psi)$. A higher value of $\delta(\Delta\Psi)$ means that more detected signals will be judged as a meteor event ~~meanwhile~~ events, but with more misidentifications and bigger errors as well. $\delta(\Delta\Psi)$ is set to approximately 30° (Hocking, 2018; Younger and Reid, 2017) in most meteor radar systems. In the HM, the zenith angle

measuring error $\delta\theta$ is due to $\delta(\Delta\Psi)$ and $\delta(\Delta\Psi)$ is a constant. Therefore, the error propagation in the receiver is very simple,

130 and $\delta\theta$ is inversely proportional to the cosine of the zenith angle.

~~Now~~We now introduce our analytical method. Our method considers a multistatic system with multiple transmitters and one ~~receiver~~receiving array in 3-dimensional space ~~(as shown in~~ figure 2~~).~~ The ~~receiver~~receiving array is in the Jones configuration, ~~which can be that is~~ “cross-shaped”, ~~but it may also be~~ “T-shaped” or “L-shaped” ~~in plan view~~. The five receiver antennas are in the same horizontal plane and constitute two orthogonal antenna arms. To avoid a complex error propagation process in 135 ~~receiver~~the receiving array and to place emphasis on multistatic configurations, the PDMEs in the two orthogonal antenna arms ($\delta(\Delta\Psi_1)$ and $\delta(\Delta\Psi_2)$) are constants. Therefore, the ~~AoAsAoA~~ measuring errors (including ~~the~~ zenith and azimuth angle measuring errors $\delta\theta$, $\delta\phi$ respectively) can be expressed as ~~a simple~~ functionfunctions of zenith and azimuth angle. The radial distance is the distance between the MTSP and the receiver, which ~~is~~ denoted as R_s . R_s can be determined by combining the ~~AoAsAoA~~, baseline length d_i , and the radio wave ~~propagating~~propagation path length R (Stober and Chau, 2015). ~~See~~The 140 ~~geometry is shown in~~ figure 4(a). α ~~is~~ is the angle between the baseline (i.e., axis- X_i) and the line from the receiver to the MTSP, ~~), if (denoted as point A).~~ If α , d_i and R are known, R_s ~~will~~can be calculated easily using ~~the~~ Cosine Law as:

$$R_s = \frac{R^2 - d_i^2}{2(R - d_i \cos\alpha)} \quad (1)$$

~~α is the angle between the baseline (i.e. X_i axis) and the line from the receiver to the MTSP denoted as point A.~~

The multistatic configuration will influence the accuracy of R_s (denoted as δR_s). This is because ~~that~~ α , d and R are determined

145 by the multistatic configuration. We consider the error term δR_s in our method, which is ignored in the HM. δR_s is a function of the ~~AoAsAoA~~ measuring errors ($\delta\theta$ and $\delta\phi$) and the radio wave propagation path length measuring error (denoted as δR).

δR is caused by the measuring error of the wave propagation time δt , which is approximately $21\mu\text{s}$ (Kang, 2008). Thus, δR can be set as a constant and the default value in our program is $\delta R = c\delta t c\delta t = 6.3\text{km}$. It is worth noting that the maximum unambiguous range for pulse meteor radars is determined by the pulse repetition frequency (PRF) (Hocking et al., 2001;

150 Holdsworth et al., 2004). For multistatic meteor radars utilizing forward scatter, the maximum unambiguous range is c/PRF (where c is the speed of light). For the area where R ~~exceeds~~exceeds the maximum unambiguous range, δR is set to positive infinity.

2.2 ~~three~~Three kinds of coordinate systems and their transformations

To better depict the multistatic system configuration, three kinds of right-hand coordinate systems ~~(figure 3)~~ need to be

155 established, ~~which as shown in figure 3. These~~ are $X_0Y_0Z_0$, $X_iY_iZ_i$ and XYZ . $X_0Y_0Z_0$ is the ENU (east-north-up) coordinate system ~~and where~~ axis- X_0 , Y_0 , Z_0 represent the east, north, up directions respectively. Another two coordinate systems are established to facilitate different error propagations. All types of errors need to be transformed to the ENU coordinate system $X_0Y_0Z_0$ in the end. Coordinate system XYZ is established to depict the spatial configuration of the ~~receiver~~. ~~XYZ is fixed on~~

receiving array and has its the receiver. See figure 3, the coordinate origin of XYZ is on the receiver, there as shown in figure

160 3. Axis-Z is collinear with the antenna boresight and perpendicular to the receiver horizontal plane on which the receiving
array lies. Axis-X and axis-Y are collinear with the arms of the two orthogonal antenna arrays. AoAs will be represented
in XYZ for convenience. See Inspection of figure 4, indicates that it is convenient to analyse the range information in a plane
that goes through the baseline and the MTSP. Thus, a coordinate system $X_iY_iZ_i$ is established for a transmitter T_i . The
165 coordinate origins of $X_iY_iZ_i$ are all on the receiver, receiving array. We stipulate that axis- X_i points to transmitter i (T_i). Each
pair of T_i and the receiver R_X constitutes a radar link, which is referred to as L_i . The range related information for
each L_i will be calculated in $X_iY_iZ_i$. Different types of errors need to propagate to and be compared in $X_0Y_0Z_0$ which is
convenient for retrieving wind fields.

We stipulate that clockwise rotation satisfies the right-hand corkscrew rule. By rotating clockwise in order of $\psi_x^{X,i}$, $\psi_y^{Y,i}$ and
170 $\psi_z^{Z,i}$ about axis-X, Y and Z, respectively, one can transform XYZ to $X_iY_iZ_i$. It is worth mentioning that $X_iY_iZ_i$ is non-unique
because any rotation about axis- X_i can obtain another satisfactory $X_iY_iZ_i$. Hence, $\psi_x^{X,i}$ can be set to any value. Similarly,
by rotating clockwise in order of $\psi_x^{i,0}$, $\psi_y^{i,0}$ and $\psi_z^{i,0}$ about axis-X, Y and Z, respectively, one can transform $X_iY_iZ_i$ to
175 $X_0Y_0Z_0$. To realize the coordinate transformation between these three coordinate systems, a coordinate rotation matrix
 $A_R(\psi_x, \psi_y, \psi_z)$ is introduced. Using A_R , one can transform the coordinate point or vector presentation from one coordinate
system to another. The details of the coordinate rotation matrix $A_R(\psi_x, \psi_y, \psi_z)$ can be seen in Appendix (A.1).

175 2.3 two types of measuring errors

The analytical method of the spatial resolution of each radar link is the same. The difference between these radar links
180 are only the value of the six coordinates coordinate rotation angles ($\psi_x^{X,i}$, $\psi_y^{Y,i}$ and $\psi_z^{Z,i}$; $\psi_x^{i,0}$, $\psi_y^{i,0}$ and $\psi_z^{i,0}$) and the
baseline distance d_i . The spatial resolution related measurement errors which will cause location errors of the MTSP, can be
classified into two types: E_1 is caused by measurement errors at the receiver, and E_2 is due to the pulse length. These two
errors are mutually independent. Hence, the total error (E_{total}) can be expressed as:

$$E_{total}^2 = E_1^2 + E_2^2 \quad (2)$$

E_1 is related to three indirect measuring errors. They are zenith, azimuth and radial distance measuring errors, denoted as $\delta\theta$,
 $\delta\phi$ and δR_s respectively. In XYZ, E_1 can be decomposed into three orthogonal error vectors using $\delta\theta$, $\delta\phi$ and δR_s (see
185 figure 4(c)). Now we explain it in more detail. PDMEs, i.e., $\delta(\Delta\Psi_1)$ and $\delta(\Delta\Psi_2)$, are caused by some practical
factors, such as phase calibration mismatch and the fact that the specular point is not actually a point but has a few Fresnel
zones in length. A meteor radar system calculates phase difference of differences between different pairs of antennas though
cross-correlations and then fits them to get the most likely AoAs. Therefore, the system needs to set be assigned a
tolerant tolerance value of $\delta(\Delta\Psi_1)$ and $\delta(\Delta\Psi_2)$. Different meteor radar systems have different AoAs AoA-fit algorithms and
thus different AoAs AoA measuring error distributions. To analyse the spatial resolution for a SIMO

190 meteor radar system as commongenerally as possible and to avoid tedious error propagation in receiverat the receiving array, we start the error propagation from $\delta(\Delta\Psi_1)$ and $\delta(\Delta\Psi_2)$ and set them as constantAoAsconstants. AoA measuring errors, i.e. $\delta\theta$ and $\delta\phi$ can then be expressed as:

$$\delta\theta = \frac{\lambda}{2\pi D_1} \frac{\cos\phi}{\cos\theta} \delta(\Delta\Psi_1) + \frac{\lambda}{2\pi D_2} \frac{\sin\phi}{\cos\theta} \delta(\Delta\Psi_2) \quad (3)$$

$$\delta\phi = \frac{\lambda}{2\pi D_2} \frac{\cos\phi}{\sin\theta} \delta(\Delta\Psi_2) - \frac{\lambda}{2\pi D_1} \frac{\sin\phi}{\sin\theta} \delta(\Delta\Psi_1) \quad (4)$$

195 where λ is the radio wave lengthwavelength, D_1 and D_2 are the length of the two orthogonal antenna arms, and θ and ϕ are the zenith angle and the azimuth angle, respectively. The details can be seenfound in **Appendix (A.2)**. It is worth noting that $\delta\theta$ and $\delta\phi$ are not mutually independent. The Expectationexpectation value of their product is not identical to zero unless $\frac{E(\delta^2(\Delta\Psi_1))}{D_1^2}$ is equal to $\frac{E(\delta^2(\Delta\Psi_2))}{D_2^2}$.

δR_s can be expressed as a function of δR , $\delta\theta$ and $\delta\phi$ as:

$$200 \quad \delta R_s = F(\delta R, \delta\theta, \delta\phi) = f_R(\theta, \phi)\delta R + f_\theta(\theta, \phi)\delta\theta + f_\phi(\theta, \phi)\delta\phi \quad (5)$$

$f_R(\theta, \phi)$, $f_\theta(\theta, \phi)$ and $f_\phi(\theta, \phi)$ are the weightweighting functions of δR_s . The details about the weightweighting function and deduction can be found in **Appendix (A.3)**. SeeInspection of figure 4(c), indicates that E_1 can be decomposed into three orthogonal error vectors in coordinate XYZ, denoted as $\overrightarrow{\delta R_s}$, $\overrightarrow{R_s\delta\theta}$ and $\overrightarrow{R_s\sin\theta\delta\phi}$. These three vectors can be expressed in XYZ as:

$$205 \quad \overrightarrow{\delta R_s} = \delta R_s(\sin\theta\cos\phi, \sin\theta\sin\phi, \cos\theta)^T \quad (6)$$

$$\overrightarrow{R_s\delta\theta} = R_s\delta\theta(\cos\theta\cos\phi, \cos\theta\sin\phi, -\sin\theta)^T \quad (7)$$

$$\overrightarrow{R_s\sin\theta\delta\phi} = R_s\sin\theta\delta\phi(-\sin\phi, \cos\phi, 0)^T \quad (8)$$

E_2 is related to the radio wave propagatingpropagation path. A pulse might be reflected anywhere within a pulse length (see figure 4(b)). This causes a location error ofin the MTSP, represented as an error vector \overrightarrow{DA} . D is the median point of the 210 isosceles triangle ΔABC 's side BC. The representation of the error vector \overrightarrow{DA} can be solved in $X_iY_iZ_i$ by using geometryrelationshipgeometrical relationships as:

$$\overrightarrow{DA} = \left(\frac{(2-a_1-a_2)x_i+d_i(a_2-1)}{2}, \frac{(2-a_1-a_2)y_i}{2}, \frac{(2-a_1-a_2)z_i}{2} \right)^T \quad (9)$$

where S is the halfwave pulse length and $a_1 = \frac{R_s-S}{R_s}$. $a_2 = \frac{R_i-S}{R_i}$. d_i is the baseline length. (x_i, y_i, z_i) is the coordinate value of a MTSP (i.e. point A in figure 4) in $X_iY_iZ_i$. DetailsMore details can be seenfound in **Appendix (A4)**

Here thus far, two types of errors in different coordinate systems have been introduced. Now they need to be transformed to ENU coordinates $X_0Y_0Z_0$, which is convenient for comparing between in order to compare different radar links and analysing to analyse the wind fields. E_1 related error vectors, which are three orthogonal vectors $\overrightarrow{\delta R_s}$, $\overrightarrow{R_s\delta\theta}$ and $\overrightarrow{R_s\sin\theta\delta\phi}$ and represented in XYZ as eq.(6)-(8), and need to be transformed from XYZ to $X_0Y_0Z_0$. To project $\overrightarrow{\delta R_s}$, $\overrightarrow{R_s\delta\theta}$ and $\overrightarrow{R_s\sin\theta\delta\phi}$ towards axis- X_0, Y_0, Z_0 respectively, and reassemble them to form three new error vectors in axis- X_0, Y_0, Z_0 . Using the coordinate rotation matrix $A_R^{(XYZ, X_0Y_0Z_0)} = A_R(\Psi_x^{i,0}, \Psi_y^{i,0}, \Psi_z^{i,0}) \cdot A_R(\Psi_x^{X,i}, \Psi_y^{Y,i}, \Psi_z^{Z,i})$ and eq.(6)-(8), the unit vectors of those three vectors can be represented in $X_0Y_0Z_0$ as:

$$\begin{pmatrix} X'_0(\delta R_s) & X'_0(\delta\theta) & X'_0(\delta\phi) \\ Y'_0(\delta R_s) & Y'_0(\delta\theta) & Y'_0(\delta\phi) \\ Z'_0(\delta R_s) & Z'_0(\delta\theta) & Z'_0(\delta\phi) \end{pmatrix} = A_R^{(XYZ, X_0Y_0Z_0)} \cdot \begin{pmatrix} \sin\theta\cos\phi & \cos\theta\cos\phi & -\sin\phi \\ \sin\theta\sin\phi & \cos\theta\sin\phi & \cos\phi \\ \cos\theta & -\sin\theta & 0 \end{pmatrix} \quad (10)$$

$(X'_0(\delta R_s), Y'_0(\delta R_s), Z'_0(\delta R_s))^T$, $(X'_0(\delta\theta), Y'_0(\delta\theta), Z'_0(\delta\theta))^T$, $(X'_0(\delta\phi), Y'_0(\delta\phi), Z'_0(\delta\phi))^T$ are unit vectors of $\overrightarrow{\delta R_s}$, $\overrightarrow{R_s\delta\theta}$ and $\overrightarrow{R_s\sin\theta\delta\phi}$ in $X_0Y_0Z_0$ respectively. The 3×3 matrix on the left hand side of the eq.(10) is denoted as P_{ij} for $i, j = 1, 2, 3$.

See From eq.(6)-(8) and figure 4(c), we see that the length of those three vectors, or (the) error values in other words, are δR_s , $R_s\delta\theta$, $R_s\sin\theta\delta\phi$ as the function of δR , $\delta\theta$, $\delta\phi$. In order to reassemble them to form new error vectors, transforming transformation of $\delta\theta$ and $\delta\phi$ into two independent errors $\delta(\Delta\Psi_1)$ and $\delta(\Delta\Psi_2)$ are is needed because $\delta\theta$ and $\delta\phi$ are not independent. Using eq. (3) and (4), one can transform vector $(\delta R, \delta\theta, \delta\phi)^T$ to three independent measuring errors δR , $\delta(\Delta\Psi_1)$ and $\delta(\Delta\Psi_2)$. And thus $(\delta R_s, R_s\delta\theta, R_s\sin\theta\delta\phi)^T$ can be expressed as:

$$\begin{pmatrix} \delta R_s \\ R_s\delta\theta \\ R_s\sin\theta\delta\phi \end{pmatrix} = \begin{pmatrix} f_R(\theta, \phi) & f_\theta(\theta, \phi) & f_\phi(\theta, \phi) \\ 0 & R_s & 0 \\ 0 & 0 & R_s\sin\theta \end{pmatrix} \cdot \begin{pmatrix} 1 & 0 & 0 \\ 0 & \frac{\frac{\lambda}{2\pi}\cos\phi}{\cos\theta D_1} & \frac{\frac{\lambda}{2\pi}\sin\phi}{\cos\theta D_2} \\ 0 & -\frac{\frac{\lambda}{2\pi}\sin\phi}{\sin\theta D_1} & \frac{\frac{\lambda}{2\pi}\cos\phi}{\sin\theta D_2} \end{pmatrix} \cdot \begin{pmatrix} \delta R \\ \delta(\Delta\Psi_1) \\ \delta(\Delta\Psi_2) \end{pmatrix} \quad (11)$$

The product of the first and the second term on the right hand side of eq.(11) is a 3×3 matrix, denoted as W_{ij} for $i, j = 1, 2, 3$. Seen From eq.(11), we see that the three error values δR_s , $R_s\delta\theta$, $R_s\sin\theta\delta\phi$ are the linear combinations of three basis δR , $\delta(\Delta\Psi_1)$, and $\delta(\Delta\Psi_2)$ with their corresponding linear coefficients W_{1j}, W_{2j}, W_{3j} and W_{3j} . Those three error values can be projected toward new directions (i.e.g. axis- X_0, Y_0, Z_0) by using P_{ij} . It worth noting that in a new direction, a same basis's projected linear coefficients from different error values should be used to calculate their sum of squares (SS). And then the square root of SS will be used as a new linear coefficient for that basis in the new direction. For example, in X_0 directions, basis $\delta(\Delta\Psi_1)$'s projected linear coefficients are $X'_0(\delta R_s)W_{12}$, $X'_0(\delta\theta)W_{22}$, $X'_0(\delta\phi)W_{32}$ from $\overrightarrow{\delta R_s}$, $\overrightarrow{R_s\delta\theta}$ and $\overrightarrow{R_s\sin\theta\delta\phi}$ respectively. Therefore, the new linear coefficient for $\delta(\Delta\Psi_1)$ in the X_0 direction is $W_{X_0}^{\delta(\Delta\Psi_1)} =$

$\pm\sqrt{(X'_0(\delta R_s)W_{12})^2 + (X'_0(\delta\theta)W_{22})^2 + (X'_0(\delta\phi)W_{32})^2}$. Similarly, one can get δR and $\delta(\Delta\Psi_2)$'s new linear coefficients in X'_0 , denoted as $W_{X'_0}^{\delta R}$ and $W_{X'_0}^{\delta(\Delta\Psi_2)}$. Thus, the true error values in the X_0 direction is $W_{X'_0}^{\delta R}\delta R + W_{X'_0}^{\delta(\Delta\Psi_1)}\delta(\Delta\Psi_1) + W_{X'_0}^{\delta(\Delta\Psi_2)}\delta(\Delta\Psi_2)$. Because δR , $\delta(\Delta\Psi_1)$, and $\delta(\Delta\Psi_2)$ are mutually independent, E_1 is related to the mean square error (MSE) values in the X_0 direction, denoted as $\delta_{(1)}X_0$, and can be expressed as $\delta_{(1)}X_0 =$

$$245 \quad \pm\sqrt{\left(W_{X'_0}^{\delta R}\delta R\right)^2 + \left(W_{X'_0}^{\delta(\Delta\Psi_1)}\delta(\Delta\Psi_1)\right)^2 + \left(W_{X'_0}^{\delta(\Delta\Psi_2)}\delta(\Delta\Psi_2)\right)^2}.$$

In short, E_1 related errors in ENU coordinate's three axis directions (denoted as $\delta_{(1)}X_0$, $\delta_{(1)}Y_0$ and $\delta_{(1)}Z_0$) can be expressed in the form of a matrix as:

$$\begin{pmatrix} \delta_{(1)}^2 X_0 \\ \delta_{(1)}^2 Y_0 \\ \delta_{(1)}^2 Z_0 \end{pmatrix} = P_{ij}^2 \cdot W_{ij}^2 \cdot \begin{pmatrix} \delta^2 R \\ \delta^2(\Delta\Psi_1) \\ \delta^2(\Delta\Psi_2) \end{pmatrix} \quad (12)$$

The E_2 related error vector \overrightarrow{DA} needs transformation from $X_iY_iZ_i$ to $X_0Y_0Z_0$. Therefore, E_2 related errors in the ENU coordinate's three axis directions (denoted as $\delta_{(2)}X_0$, $\delta_{(2)}Y_0$ and $\delta_{(2)}Z_0$) can be expressed in the form of a matrix as:

$$\begin{pmatrix} \delta_{(2)} X_0 \\ \delta_{(2)} Y_0 \\ \delta_{(2)} Z_0 \end{pmatrix} = \pm A_R(\Psi_x^{i,0}, \Psi_y^{i,0}, \Psi_z^{i,0}) \cdot \overrightarrow{DA} \quad (13)$$

E_1 and E_2 are mutually independent. By using eq.(1), the total MSE values in ENU coordinate's three axis directions (denoted as $\delta_{\text{total}}X_0$, $\delta_{\text{total}}Y_0$ and $\delta_{\text{total}}Z_0$) can be expressed in the form of matrix as:

$$\begin{pmatrix} \delta_{\text{total}}^2 X_0 \\ \delta_{\text{total}}^2 Y_0 \\ \delta_{\text{total}}^2 Z_0 \end{pmatrix} = \begin{pmatrix} \delta_{(1)}^2 X_0 \\ \delta_{(1)}^2 Y_0 \\ \delta_{(1)}^2 Z_0 \end{pmatrix} + \begin{pmatrix} \delta_{(2)}^2 X_0 \\ \delta_{(2)}^2 Y_0 \\ \delta_{(2)}^2 Z_0 \end{pmatrix} \quad (14)$$

255 In conclusion, for a radar link L_i and a MTSP represented as (x_0, y_0, z_0) in the ENU coordinate system $X_0Y_0Z_0$, as sketched in figure 4(a)-sketched, the location errors of this point in east, north and up directions ($\pm\delta_{\text{total}}X_0$, $\pm\delta_{\text{total}}Y_0$ and $\pm\delta_{\text{total}}Z_0$) can be calculated as follows: firstly, for a point (x_0, y_0, z_0) in $X'_0Y'_0Z'_0$, using A_R to transform it to $X_iY_iZ_i$ and denoted as (x_i, y_i, z_i) . Then in $X_iY_iZ_i$ calculate AoAs the AoA (θ and ϕ) and the range information (R_s and R_i). Details of AoAsAoA and range calculation can be seen in Appendix (A.5). It's worth noting that AoAs are the

260 AoA is given by the angles relative to the axes of XYZ. Secondly, in XYZ using AoAs the AoA and eq.(3)-(8) to calculate E_1 's three orthogonal error vectors as shown in figure 4(c)-sketched; in $X_iY_iZ_i$ use the range information and eq.(9) to calculate E_2 's error vector \overrightarrow{DA} as shown in figure 4(b)-sketched. Thirdly, project E_1 's three error vectors to $X_0Y_0Z_0$ by using eq.(10) and use eq.(11)-(12) to reassemble them to calculate E_1 related MSE values in the direction of

265 X_0, Y_0, Z_0 ; use eq.(13) to transform the E_2 error vector from $X_i Y_i Z_i$ to $X_0 Y_0 Z_0$. Finally, use eq. (14) to get the total location errors of a MTSP in (x_0, y_0, z_0) . Figure 5(a) describesshows the flow chart for the process above.—we have just described.

3 Results and Discussion

We wrote a The program to study the method we have described above. The program is written in the python language and is presented in the supplement. To calculate a special configuration of a multistatic radar system, we initially need to set six coordinate transformation angles ($\psi_x^{X,i}, \psi_y^{Y,i}$ and $\psi_z^{Z,i}$; $\psi_x^{i,0}, \psi_y^{i,0}$ and $\psi_z^{i,0}$) and the baseline length (i.e. d_i) for each radar link

270 L_i . For example, $\psi_x^{i,0} = \psi_y^{i,0} = 0$, $\psi_z^{i,0} = 30^\circ$ and $d_i = 250\text{km}$ means that transmitter T_i is 250km250 km, 30° east by south of the receiver R_X ; Further, $\psi_x^{X,i} = 5^\circ, \psi_y^{Y,i} = 0, \psi_z^{Z,i} = 0$ means one receiver arm (axis-Y) points to east by north 60° with 5° elevation. The interested-detection area of interest for a multistatic meteor radar is usually from 70km70 km to 110km110 km in height and lager than around $300\text{km} \times 300\text{km}$ in the horizontal. In our program, this area needs to be divided into a spatial grid for sampling. The default value of the sampling grid length is 1km1 km in height and 5km5 km in meridianthe meridional and zonal directions, respectively. After selecting the desired settings mentioned above, the program will traverse these steps through the sampling grid nodes and calculates the location errors ofat each node as described in figure 5(a). Figure 5(b) describesdescribes the parameter settings and traversesthe transversal calculation process above. For a given setting of radar link L_i , the program will output the squared values of E_1 related, E_2 related and total MSE (E_{total}^2 :

280 $\delta_{total}^2 X_0, \delta_{total}^2 Y_0, \delta_{total}^2 Z_0$; $E_1^2: \delta_{(1)}^2 X_0, \delta_{(1)}^2 Y_0, \delta_{(1)}^2 Z_0$; $E_2^2: \delta_{(2)}^2 X_0, \delta_{(2)}^2 Y_0, \delta_{(2)}^2 Z_0$). The location errors can be positive or negative and thus the spatial resolutions are twice the absolute value of the location errors. For an example, Seesee figure 5(c), for. For a detected MTSP represented as (x_0, y_0, z_0) in $X_0 Y_0 Z_0$, if with $\delta_{total}^2 X_0, \delta_{total}^2 Y_0, \delta_{total}^2 Z_0$ equalequal to 25, 16 and 9 km^2 , respectively, it means that the actual position of the MTSP could occur in an area which is $\pm 5 \text{ km}, \pm 4 \text{ km}, \pm 3 \text{ km}$ around (x_0, y_0, z_0) with equallyequal probability. TheConsequently, the zonal, meridianmeridional and vertical resolutionresolutions are 10 km, 8 km and 6 km respectively.

285 The HM analyses the vertical resolution (corresponding to δZ_0 in our paper) only in a 2-dimensional vertical section (corresponding to the $X_0 Z_0$ plane in our paper). To compare with Hocking's work, except $\psi_z^{i,0}$ is set to be -180° , and the other five coordinate transformation angles are all set to zero with d is equal to 300 km. The half wave-pulse length S is set to 2 km and $\delta(\Delta\Psi_1)$ to 35° . Calculating in only the $X_0 Y_0$ plane only should have degraded our method into Hocking's 2-dimensional analysis method, but the settings above doesn't work because Hocking'sthe HM method ignores δR_s . In fact,

290 Hocking'smethodthe HM considers only E_2 and $\overrightarrow{R_s \delta \theta}$ in the $X_0 Y_0$ plane. HenceConsequently, we need to further set $f_R(\theta, \phi)$, $f_\theta(\theta, \phi)$ and $f_\phi(\theta, \phi)$ to be zero. ThusWhen this is done, our method totally degrades into Hocking'smethodthe HM. Hocking's results are shown inas the absolute value of vertical location error normalized relative to the half wave-pulse width, i.e. $|\delta Z_0|/S$. HereafterHereinafter, $|E|/S$ is referred to as the normalized spatial resolutionresolutions such as $\delta_{(1)} X_0$

and $\delta_{total}Y_0$, where E represents the location errors in a direction. Thus, ~~Spatial~~ the spatial resolutions are $2S$ times the normalized spatial resolutions. The

Our normalized vertical resolution distributions are shown in figure 6(a). Our results and are the same as those presented in Hocking's work (Hocking, 2018). The distribution of $\overrightarrow{R_s\delta\theta}$ related, E_2 related, and total normalized vertical resolution distributions are shown in figure 6 from left to right, respectively. In most cases, E_2 is an order of magnitude smaller than $\overrightarrow{R_s\delta\theta}$. Only in the region directly above the receiver does E_2 have the same magnitude as $\overrightarrow{R_s\delta\theta}$. In other words, only in the region directly above the receiver can E_2 influence the total resolution. E_2 is related to the bistatic configuration, but $\overrightarrow{R_s\delta\theta}$ is not. Therefore, in the HM, the distribution of the total vertical resolution is changed varies slightly varying with \mathbf{d} . After adding the error term $\overrightarrow{\delta R_t}$, which is related to the bistatic configuration, the normalized total vertical spatial resolution distribution will change visibly varying changes more obviously with \mathbf{d} , as figure 7's first two rows show. The region between the two black lines represents a trustworthy the sampling volume for the receiver because where the elevation angle is beyond 30° with less influence of potential mutual antenna coupling or other obstacles in the surrounding. However, with. As the transmitter/receiver distance become longer, resolutions in this trustworthy sampling volume are not always acceptable. In figure 7's first row, the transmitter/receiver distance is 300 km and about half of the region between two black line have normalized vertical resolution values larger than 3 km. Because our analytical method can obtain spatial resolutions in 3-dimensional space, figure 7's third row shows a perspective to the horizontal section at 90 km altitude for figure 7's second row.

To get an intuitionistic a perspective to on the spatial resolution distribution in 3-dimensional space, figure 8 shows the normalized zonal, meridional and vertical spatial resolution distribution of distributions for a multistatic radar link whose transmitter/receiver separation is 180 km away and the transmitter is south by east 30° of the receiver. ~~Classic~~ The

classic monostatic meteor radar is a special case of a multistatic meteor radar system whose baseline length is zero. By setting

the transmitter/receiver distance to be zero in our program, a monostatic meteor radar's spatial resolution can also be obtained.

The In this case, the spatial resolution distributions are highly symmetrical and correspond to the real characteristics of monostatic meteor radar (this is not shown in the text, here, but can be seen found in the supplement SF1). In the discussion

above, the receiver and transmitter antennas are all coplanar. By setting varying $\psi_x^{x,i}$, $\psi_y^{y,i}$ and $\psi_z^{z,i}$ in our program, the non-

coplanar receiver/transmitter-antennas situations can also be studied. Slightly tilting of the receiver horizontal plane (for

example, set $\psi_x^{x,i} = \psi_y^{y,i} = 5^\circ$) will cause causes the horizontal spatial distributions to change (seen see SF2 and SF3 in the

supplement). In practical applications, like practice, the Earth's curvature and local topography will lead to tilts in the receiver

horizontal plane tilting. Thus, this. This kind of tilting tilt should also be taken into account for multistatic meteor radar systems.

The and details of relating to the parameter setting selections for this can be seen found in the supplement.

The AoAsAoA error propagation process in the receiver has been simplified to yield eq.(3)-(4) by using the constant PDMEs

as the start of error propagation. This is for the sake of the adaptable providing the most general example of our method. If

analysing AoAs the analysis of AoA errors starts were to start from the original received voltage signals, (e.g., Vaudrin et al.,

330 2018), the error propagation process will change with a would depend on the specific receiver interferometer configuration and at the specific signal processing method. In practical situations for an unusual The approach used here can be applied to different receiver antenna configuration configurations or new original signal processing algorithm, an error propagation process based on the specific circumstances needs to be established. Substitute algorithms. This would involve substitution of $\delta(\Delta\Psi_1)$ and $\delta(\Delta\Psi_2)$ into other mutually independent measuring errors in a practical situation, to suit the experimental arrangement and then establishing a new AoAsAoA error propagation to obtain $\delta\theta$ and $\delta\phi$. Or in other words, rewrite This means rewriting the second and third term in eq. (11) to the determine a new established AoAsAoA error propagation matrix and new mutually independent measuring errors, respectively. Our analytical method can still work.

335 It worth noting that except for using the PDMEs as the start of the error propagation, all the analytical processes are built on the mathematical error propagations. PDMEs include the uncontrolled errors, such as the those resulting from the returned wave being scattered wave from a few Fresnel zones along the meteor trail, phase calibration inaccuracy, and noises noise. However, there are other error sources in practical situation practice. For example, planes aircraft or lightning may make troubles for meteor radar's discrimination system. And interference of fading clutter from obstacles in surroundings 340 will can cause further measurement errors of AoAs in the AoA. These issues are related to actual physical situations and beyond the scope of this text.

345 The trustworthy sampling volume Knowing the valid observational volume for meteor detections and the errors associated with each detection is vital for a meteor radar system and as it determines the detection area and which meteors could can be used into calculate wind retrievals velocities and also the uncertainties associated with the winds themselves. To avoid reduce the influence of the mutual antenna coupling or the ground clutter, the elevation angle of a detection should beyond be above a threshold, for example and 30° in general. The spatial resolution is another thing that affects the trustworthy sampling typically used, and this sets the basic valid observational volume. See Within this, the normalised vertical resolution varies, and in Figure 7 and SF4 in the supplement, only the area areas of normalized vertical resolution with values below 3 km are shown, which we argue represents an acceptable sampling volume. With In addition, as the transmitter/receiver distance 350 increasing, this increases, the sampling volume becomes smaller along with and the vertical resolution in this volume is reduced. This fact effect limits the practically usable transmitter/receiver distance distances for multistatic meteor radar. Measurement response is important for measuring meteor trails' Doppler shift caused by the background wind. Theradars.

355 The geometry of the multistatic meteor radar case also impacts on the ability of the radar to measure the Doppler shifts associated with drifting meteor trails within the observational volume. This is because the measured Doppler shift is caused produced by the component of the wind fields field in the direction of the Bragg Vector, which in the multistatic configuration is divergent from the receiver's line of sight (see e.g., Spargo et al., 2019). The smaller the angle between the Bragg vector and the wind fields is, the lager this larger is the Doppler shift is and meanwhile the higher SNR. The Bragg vector of the multistatic configuration is divergent from the receiver's line of sight. Monostatic meteor radars can only detect winds in radial direction, thus only the mean wind can be solved. By synthesizing monostatic and multistatic the high order component of the wind fields can be solved. The bigger the angle between the Bragg vector and radial direction is, or more

diversified Bragg vectors in other words, the more complete and accurate the wind fields will be observed. In short, the trustworthy sampling volume, measurement response and the SNR). This means that within the observational volume, the angular diversity of the Bragg vector should both be taken into account in the wind retrievals. The retrieval process. A discussion of wind retrievals is beyond the scope of this text and will be considered in a future work.

365

4 Conclusion

In this study, we have presented the preliminary results of our from an analytical error analytic method. Our analysis of multistatic meteor radar system measurements of angles of arrival. The method can calculate the spatial resolution (the spatial uncertainty) in the zonal, meridional and vertical directions for an arbitrary receiving antenna array configuration in three-dimensional space. A given detected MTSP can locate is located within the spatial resolution volume with an equal probability. Higher values of spatial resolution mean that this region needs more meteor counts or longer averaging to obtain a reliable accuracy. Our method shows that the spatial configuration of a multistatic system will greatly influence the spatial resolution distribution in ENU coordinates and thus will in turn influence the retrieval accuracy of atmospheric parameters such as the wind fields. The multistatic meteor radar system's spatial resolution analysis is a key point in analysing the accuracy of retrieved wind and other parameters. The influence of the spatial resolution on wind retrieval will be discussed in the future work.

Multistatic radar systems come in many types, and our the work in this paper considers only single-input (single-antenna transmitter) and multi-output (5-antenna interferometric receiver) pulse radar systems. Although the single-input multi-output (SIMO) pulse meteor radar is a classic meteor radar system, other meteor radar systems, such as continuous wave radar systems and MISO (multiple-antenna transmitter and single-antenna receiver), also show good experimental results and have some advantages over SIMO systems. Using different types of meteor radar systems to constitute the meteor radar network is the future trend and so we will add the spatial resolution analyses of other system to the frame of types using our method in the future. We will also validate and apply the error analyses of spatial resolution analysis in the horizontal wind determination into a multistatic meteor radar system, which that will be built soon be installed in China.

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Code availability. The program to calculate the 3D spatial resolution distributions are is available in the supplement.

Author contributions: W.Z, X.X, W.Y designed the study. W.Z deduced the formulas and wrote the program. W.Z wrote the paper for the first version. X.X supervised the work and provided valuable comments. I.M.R revised the paper. All of the authors discussed the results and commented on the paper.

Competing interest. The authors declare no conflicts of interests

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400 Reference

Browning, K. A., and Wexler, R.: The Determination of Kinematic Properties of a Wind Field Using Doppler Radar, Journal of Applied Meteorology, 7, 105-113, 10.1175/1520-0450(1968)007<0105:Tdokpo>2.0.Co;2, 1968.

Ceplecha, Z., Borovička, J., Elford, W. G., ReVelle, D. O., Hawkes, R. L., Porubčan, V., and Šimek, M.: Meteor Phenomena and Bodies, Space Science Reviews, 84, 327-471, 10.1023/A:1005069928850, 1998.

405 Chau, J. L., Stober, G., Hall, C. M., Tsutsumi, M., Laskar, F. I., and Hoffmann, P.: Polar mesospheric horizontal divergence and relative vorticity measurements using multiple specular meteor radars, Radio Science, 52, 811-828, 10.1002/2016rs006225, 2017.

Chau, J. L., Urco, J. M., Vierinen, J. P., Volz, R. A., Clahsen, M., Pfeffer, N., and Trautner, J.: Novel specular meteor radar systems using coherent MIMO techniques to study the mesosphere and lower thermosphere, Atmos. Meas. 410 Tech., 12, 2113-2127, 10.5194/amt-12-2113-2019, 2019.

DOREY, J., BLANCHARD, Y., and CHRISTOPHE, F.: Le projet RIAS, une approche nouvelle du radar de surveillance aérienne, Onde électrique, 64, 15-20, 1984.

Hocking, W. K., Fuller, B., and Vandepoor, B.: Real-time determination of meteor-related parameters utilizing modern digital technology, Journal of Atmospheric and Solar-Terrestrial Physics, 63, 155-169, 10.1016/s1364-415 6826(00)00138-3, 2001.

Hocking, W. K.: Spatial distribution of errors associated with multistatic meteor radar, Earth, Planets and Space, 70, 93, 10.1186/s40623-018-0860-2, 2018.

Holdsworth, D. A., Reid, I. M., and Cervera, M. A.: Buckland Park all-sky interferometric meteor radar, Radio Science, 39, 10.1029/2003rs003014, 2004.

420 Jacobi, C., Hoffmann, P., and Kärschner, D.: Trends in MLT region winds and planetary waves, Collm (52°N, 15°E), Annales Geophysicae (ANGE), 2008.

Jia, M. J., Xue, X. H., Dou, X. K., Tang, Y. H., Yu, C., Wu, J. F., Xu, J. Y., Yang, G. T., Ning, B. Q., and Hoffmann, L.: A case study of A mesoscale gravity wave in the MLT region using simultaneous multi-instruments in Beijing, Journal of Atmospheric and Solar-Terrestrial Physics, 140, 1-9, 10.1016/j.jastp.2016.01.007, 2016.

425 Jones, J., Webster, A. R., and Hocking, W. K.: An improved interferometer design for use with meteor radars, Radio Science, 33, 55-65, 10.1029/97rs03050, 1998.

Kang, C.: Meteor radar signal processing and error analysis, 2008.

Philippe, W., and Corbin, H.: On the Analysis of Single-Doppler Radar Data, Journal of Applied Meteorology - J APPL METEOROL, 18, 532-542, 10.1175/1520-0450(1979)018<0532:OTAOSD>2.0.CO;2, 1979.

430 Reid, I. M.: SOME ASPECTS OF DOPPLER RADAR MEASUREMENTS OF THE MEAN AND FLUCTUATING COMPONENTS
OF THE WIND-FIELD IN THE UPPER MIDDLE ATMOSPHERE, Journal of Atmospheric and Terrestrial Physics,
49, 467-484, 10.1016/0021-9169(87)90041-9, 1987.

Spargo, A. J., Reid, I. M., and MacKinnon, A. D.: Multistatic meteor radar observations of gravity-wave-tidal interaction
over southern Australia, Atmos. Meas. Tech., 12, 4791-4812, 10.5194/amt-12-4791-2019, 2019.

435 Stober, G., Sommer, S., Rapp, M., and Latteck, R.: Investigation of gravity waves using horizontally resolved radial
velocity measurements, Atmos. Meas. Tech., 6, 2893-2905, 10.5194/amt-6-2893-2013, 2013.

Stober, G., and Chau, J. L.: A multistatic and multifrequency novel approach for specular meteor radars to improve
wind measurements in the MLT region, Radio Science, 50, 431-442, 10.1002/2014rs005591, 2015.

440 Stober, G., Chau, J. L., Vierinen, J., Jacobi, C., and Wilhelm, S.: Retrieving horizontally resolved wind fields using multi-
static meteor radar observations, Atmos. Meas. Tech., 11, 4891-4907, 10.5194/amt-11-4891-2018, 2018.

Urco, J. M., Chau, J. L., Weber, T., Vierinen, J. P., and Volz, R.: Sparse Signal Recovery in MIMO Specular Meteor Radars
With Waveform Diversity, Ieee Transactions on Geoscience and Remote Sensing, 57, 10088-10098,
10.1109/tgrs.2019.2931375, 2019.

Vaudrin, C. V., Palo, S. E., and Chau, J. L.: Complex Plane Specular Meteor Radar Interferometry, Radio Science, 53, 112-
445 128, 10.1002/2017rs006317, 2018.

Vierinen, J., Chau, J. L., Pfeffer, N., Clahsen, M., and Stober, G.: Coded continuous wave meteor radar, Atmospheric
Measurement Techniques, 9, 829-839, 10.5194/amt-9-829-2016, 2016.

Vierinen, J., Chau, J. L., Charuvil, H., Urco, J. M., Clahsen, M., Avsarkisov, V., Marino, R., and Volz, R.: Observing
450 Mesospheric Turbulence With Specular Meteor Radars: A Novel Method for Estimating Second-Order
Statistics of Wind Velocity, Earth and Space Science, 6, 1171-1195, 10.1029/2019ea000570, 2019.

Xue, X. H., Dou, X. K., Lei, J., Chen, J. S., Ding, Z. H., Li, T., Gao, Q., Tang, W. W., Cheng, X. W., and Wei, K.: Lower
thermospheric-enhanced sodium layers observed at low latitude and possible formation: Case studies, Journal
of Geophysical Research-Space Physics, 118, 2409-2418, 10.1002/jgra.50200, 2013.

455 Yi, W., Xue, X. H., Reid, I. M., Younger, J. P., Chen, J. S., Chen, T. D., and Li, N.: Estimation of Mesospheric Densities at
Low Latitudes Using the Kunming Meteor Radar Together With SABER Temperatures, Journal of Geophysical
Research-Space Physics, 123, 3183-3195, 10.1002/2017ja025059, 2018.

Younger, J. P., and Reid, I. M.: Interferometer angle-of-arrival determination using precalculated phases, Radio Science,
52, 1058-1066, 10.1002/2017rs006284, 2017.

460 Appendix

A.1 Coordinates rotation matrix

For a right-handed rectangular coordinate system XYZ , we rotate clockwise Ψ_x about the axis-x to obtain a new coordinate. We specify that clockwise rotation satisfies in the right-hand screw rule. A vector in XYZ , denoted as $(x, y, z)^T$, is represented as $(x', y', z')^T$ in the new coordinate. The relationship between $(x, y, z)^T$ and $(x', y', z')^T$ is:

465
$$\begin{pmatrix} x' \\ y' \\ z' \end{pmatrix} = A_x(\psi_x) \begin{pmatrix} x \\ y \\ z \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos\psi_x & \sin\psi_x \\ 0 & -\sin\psi_x & \cos\psi_x \end{pmatrix} \begin{pmatrix} x \\ y \\ z \end{pmatrix} \quad (A1.1)$$

Similarly, we rotate clockwise Ψ_y about the axis-y to obtain a new coordinate. The presentation for a vector in new coordinates and the original can be linked by a matrix, $A_y(\psi_y)$:

$$A_y(\psi_y) = \begin{pmatrix} \cos\psi_y & 0 & -\sin\psi_y \\ 0 & 1 & 0 \\ \sin\psi_y & 0 & \cos\psi_y \end{pmatrix} \quad (\text{A1.2})$$

we rotate clockwise Ψ_z about axis-z to obtain a new coordinate. The presentation for a vector in new coordinates and original

470 can be linked by a matrix $A_z(\psi_z)$:

$$A_z(\psi_z) = \begin{pmatrix} \cos\psi_z & \sin\psi_z & 0 \\ -\sin\psi_z & \cos\psi_z & 0 \\ 0 & 0 & 1 \end{pmatrix} \quad (\text{A1.3})$$

For any two coordinate systems XYZ and $X'Y'Z'$ with co-origin, one can always rotate clockwise Ψ_x , Ψ_y and Ψ_z in order of axis-X, Y, Z respectively, transforming XYZ to $X'Y'Z'$ (figure A.1). The presentation for a vector in $X'Y'Z'$ and XYZ can be linked by a matrix, $A_R(\psi_x, \psi_y, \psi_z)$:

475 $A_R(\psi_x, \psi_y, \psi_z) = A_z(\psi_z)A_y(\psi_y)A_x(\psi_x) =$

$$\begin{pmatrix} \cos\psi_y\cos\psi_z & \sin\psi_x\sin\psi_y\cos\psi_z + \cos\psi_x\sin\psi_z & -\cos\psi_x\sin\psi_y\cos\psi_z + \sin\psi_x\sin\psi_z \\ -\cos\psi_y\sin\psi_z & -\sin\psi_x\sin\psi_y\sin\psi_z + \cos\psi_x\cos\psi_z & \cos\psi_x\sin\psi_y\sin\psi_z + \sin\psi_x\cos\psi_z \\ \sin\psi_y & -\sin\psi_x\cos\psi_y & \cos\psi_x\cos\psi_y \end{pmatrix} \quad (\text{A1.4})$$

We call $A_R(\psi_x, \psi_y, \psi_z)$ as the coordinates rotation matrix.

A.2 AoAsAoA measuring errors

In coordinate XYZ , AoAs includes zenith angle θ and azimuth angle ϕ . In the plane wave approximation, the radio wave is

480 at angle γ_1 and γ_2 with an antenna array (figure A.2). There is a phase difference $\Delta\Psi_1$ and $\Delta\Psi_2$ between two antennas (figure 1). See figure 1, $\Delta\Psi_1$ and $\Delta\Psi_2$ can be expressed as:

$$\Delta\Psi_1 = \frac{2\pi D_1 \cos\gamma_1}{\lambda} \quad (\text{A2.1})$$

$$\Delta\Psi_2 = \frac{2\pi D_2 \cos\gamma_2}{\lambda} \quad (\text{A2.2})$$

Using γ_1 , γ_2 the AoAsAoA can be expressed as:

485 $\cos^2\gamma_1 + \cos^2\gamma_2 + \cos^2\theta = 1 \quad (\text{A2.3})$

$$\tan\phi = \frac{\cos\gamma_2}{\cos\gamma_1} \quad (\text{A2.4})$$

Or in another expression:

$$\cos\gamma_1 = \sin\theta\cos\phi \quad (\text{A2.5})$$

$$\cos\gamma_2 = \sin\theta\sin\phi \quad (A2.6)$$

490 substitute $\cos\gamma_1$ and $\cos\gamma_2$ in (A2.3) and (A2.4) by using (A2.1) and (A2.2):

$$\cos^2\theta = 1 - \left(\frac{\lambda}{2\pi}\right)^2 \left(\frac{\Delta^2\Psi_1}{D_1^2} + \frac{\Delta^2\Psi_2}{D_2^2}\right) \quad (A2.7)$$

$$\ln(\tan\phi) = \ln(D_1\Delta\Psi_2) - \ln(D_2\Delta\Psi_1) \quad (A2.8)$$

| (A2.7) and (A2.8) link the phase difference with the AoAsAoA and expanding θ and ϕ , $\Delta\Psi_1$ and $\Delta\Psi_2$ to first order:

$$2\cos\theta\sin\theta\delta\theta = \left(\frac{\lambda}{2\pi}\right)^2 \left[\frac{2\Delta\Psi_1\delta(\Delta\Psi_1)}{D_1^2} + \frac{2\Delta\Psi_2\delta(\Delta\Psi_2)}{D_2^2}\right] \quad (A2.9)$$

$$495 \quad \delta\phi = \frac{\sin\phi\cos\phi}{\Delta\Psi_2} \delta(\Delta\Psi_2) - \frac{\sin\phi\cos\phi}{\Delta\Psi_1} \delta(\Delta\Psi_1) \quad (A2.10)$$

For (A2.9) and (A2.10), substitute $\Delta\Psi_1$ and $\Delta\Psi_2$ using (A2.1), (A2.2) and (A2.5), (A2.6) to the functions of θ , ϕ . Now, eq. (3) and eq. (4) have been proven. If the zenith angle $\theta = 0^\circ$, we stipulate that $\frac{\cos\phi}{\sin\theta}$ and $\frac{\sin\phi}{\sin\theta}$ are 1.

A.3 Radial distance measuring error

Expand R_s , R and $\cos\alpha$ in eq.(1) to first order, δR_s can be expressed as a function of δR and $\delta(\cos\alpha)$:

$$500 \quad \delta R_s = \frac{R^2 - 2R\cos\alpha + d^2}{2(R-d\cos\alpha)^2} \delta R + \frac{d(R^2 - d^2)}{2(R-d\cos\alpha)^2} \delta(\cos\alpha) \quad (A3.1)$$

α is the angle between R_s and axis- X_i . We denote the zenith and azimuth angles in coordinate- $X_iY_iZ_i$ as θ' and ϕ' , respectively. And the relationship between α and θ' , ϕ' is

$$\cos\alpha = \sin\theta'\cos\phi' \quad (A3.2)$$

| Using coordinates rotation matrix $A_R(\psi_x^{X,i}, \psi_y^{Y,i}, \psi_z^{Z,i})$, $\sin\theta'\cos\phi'$ can be expressed as the function of AoAsAoA:

$$505 \quad \sin\theta'\cos\phi' = A_{11}\sin\theta\cos\phi + A_{12}\sin\theta\sin\phi + A_{13}\cos\theta$$

$$(A3.3)$$

A_{ij} are represent the elements in matrix $A_R(\psi_x^{X,i}, \psi_y^{Y,i}, \psi_z^{Z,i})$ for $i, j = 1, 2, 3$.

Using (A3.2) and (A3.3), $\delta(\cos\alpha)$ can be expressed as a function of $\delta\theta$ and $\delta\phi$ as:

$$\delta(\cos\alpha) = (A_{11}\cos\theta\cos\phi + A_{12}\cos\theta\sin\phi - A_{13}\sin\theta)\delta\theta + (-A_{11}\sin\theta\sin\phi + A_{12}\sin\theta\cos\phi)\delta\phi \quad (A3.4)$$

510 Finally, δR_s can be expressed as the function of $\delta R, \delta\theta, \delta\phi$ as:

$$\delta R_s = F(\delta R, \delta \theta, \delta \phi) = f_R(\theta, \phi) \delta R + f_\theta(\theta, \phi) \delta \theta + f_\phi(\theta, \phi) \delta \phi$$

(A3.5)

For:

$$f_R(\theta, \phi) = \frac{d^2 + R^2 - 2Rd(A_{11}\sin\theta\cos\phi + A_{12}\sin\theta\sin\phi + A_{13}\cos\theta)}{2[R - d(A_{11}\sin\theta\cos\phi + A_{12}\sin\theta\sin\phi + A_{13}\cos\theta)]^2}$$

515 (A3.6)

$$f_\theta(\theta, \phi) = \frac{d(R^2 - d^2)(A_{11}\cos\theta\cos\phi + A_{12}\cos\theta\sin\phi - A_{13}\sin\theta)}{2[R - d(A_{11}\sin\theta\cos\phi + A_{12}\sin\theta\sin\phi + A_{13}\cos\theta)]^2}$$

(A3.7)

$$f_\phi(\theta, \phi) = \frac{d(R^2 - d^2)(-A_{11}\sin\theta\sin\phi + A_{12}\sin\theta\cos\phi)}{2[R - d(A_{11}\sin\theta\cos\phi + A_{12}\sin\theta\sin\phi + A_{13}\cos\theta)]^2}$$

(A3.8)

520 A.4 True error of E_2

See figure 4 (b); the total length of side AC and side AB represents the pulse width. Side AC equals side CB and they are both equal to half of the pulse width S. In $X_iY_iZ_i$, the presentation of point A is (x_i, y_i, z_i) , the receiver is $(0, 0, 0)$ and T_i is $(d, 0, 0)$. The distance between T_i and A is $R_i = R - R_s$. We denote that the presentation of point B and C in $X_iY_iZ_i$ is (x_B, y_B, z_B) and (x_C, y_C, z_C) , respectively. We use vector collinear to establish equations for B and C. Therefore, one can 525 obtain the coordinates of point B and C by the following equations:

$$(x_B, y_B, z_B)^T = \frac{R_s - S}{R_s} (x_i, y_i, z_i)^T \quad (A4.1)$$

$$(x_C - d, y_C, z_C)^T = \frac{R_i - S}{R_i} (x_i - d, y_i, z_i)^T \quad (A4.2)$$

$$(x_i - d, y_i, z_i)^T \quad (A4.2)$$

For isosceles triangle ABC, the perpendicular line AD intersects side CB ~~in middle point at the midpoint~~ D. Then, we obtain

530 the coordinate value of D in $X_iY_iZ_i$ as:

$$(x_D, y_D, z_D) = \frac{1}{2}(x_B + x_C, y_B + y_C, z_B + z_C) = \frac{1}{2}((a_1 + a_2)x_i - a_2 d + d, (a_1 + a_2)y_i, (a_1 + a_2)z_i) \quad (A4.3)$$

We denote $a_1 = \frac{R_s - S}{R_s}$, $a_2 = \frac{R_i - S}{R_i}$. Finally, one can obtain the error vector of E_2 as vector \overrightarrow{DA} in $X_iY_iZ_i$:

$$\overrightarrow{DA} = \left(\frac{(2-a_1-a_2)x_i + d(a_2-1)}{2}, \frac{2-a_1-a_2}{2}y_i, \frac{2-a_1-a_2}{2}z_i \right)^T \quad (A4.4)$$

535 A.5 ~~alculate AoAs~~Calculate AoA and range information in $X_iY_iZ_i$

For a space point (x_i, y_i, z_i) in $X_iY_iZ_i$ which ~~represent~~represents a MTSP, R_s can be solved easily as:

$$\vec{R}_s = (x_i, y_i, z_i)$$

$$R_s = \sqrt{x_i^2 + y_i^2 + z_i^2} \quad (A6.1)$$

The distance between transmitter T_i and receiver R_X is d_i as ~~sketched~~shown in figure 4(a). Thus, the coordinate value of

540 T_i in $X_iY_iZ_i$ is $(d_i, 0, 0)$ and R_i can be solved as:

$$R_i = \sqrt{(x_i - d_i)^2 + y_i^2 + z_i^2} \quad (A6.2)$$

Before we calculate the AoAs in $X_iY_iZ_i$, the representation of unit vectors of axis-X, Y, Z in $X_iY_iZ_i$ ~~need~~needs to ~~know~~be known. In XYZ those unit vectors are easily represented as $(1,0,0)^T$, $(0,1,0)^T$, $(0,0,1)^T$. Though ~~coordinat~~the coordinate rotation matrix $A_R(\psi_x^{X,i}, \psi_y^{Y,i}, \psi_z^{Z,i})$, one can get those unit vector's representation in $X_iY_iZ_i$ as:

545 $\vec{n}_x = (A_{11}, A_{21}, A_{31})^T$

$$\vec{n}_y = (A_{12}, A_{22}, A_{32})^T$$

$$\vec{n}_z = (A_{13}, A_{23}, A_{33})^T \quad (A6.3)$$

For \vec{n}_x , \vec{n}_y and \vec{n}_z are unit vectors of Axis-X, Y, Z respectively. ~~And, and~~ A_{ij} are the elements ~~in~~in 3×3 matrix $A_R(\psi_x^{X,i}, \psi_y^{Y,i}, \psi_z^{Z,i})$ for $i, j = 1, 2, 3$. Now AoA as the AoA can get be obtained as:

550 $\cos \theta = \frac{\vec{R}_s \cdot \vec{n}_z}{R_s} \quad (A6.4)$

$$\sin \theta = \sqrt{1 - \cos^2 \theta}$$

(A6.5)

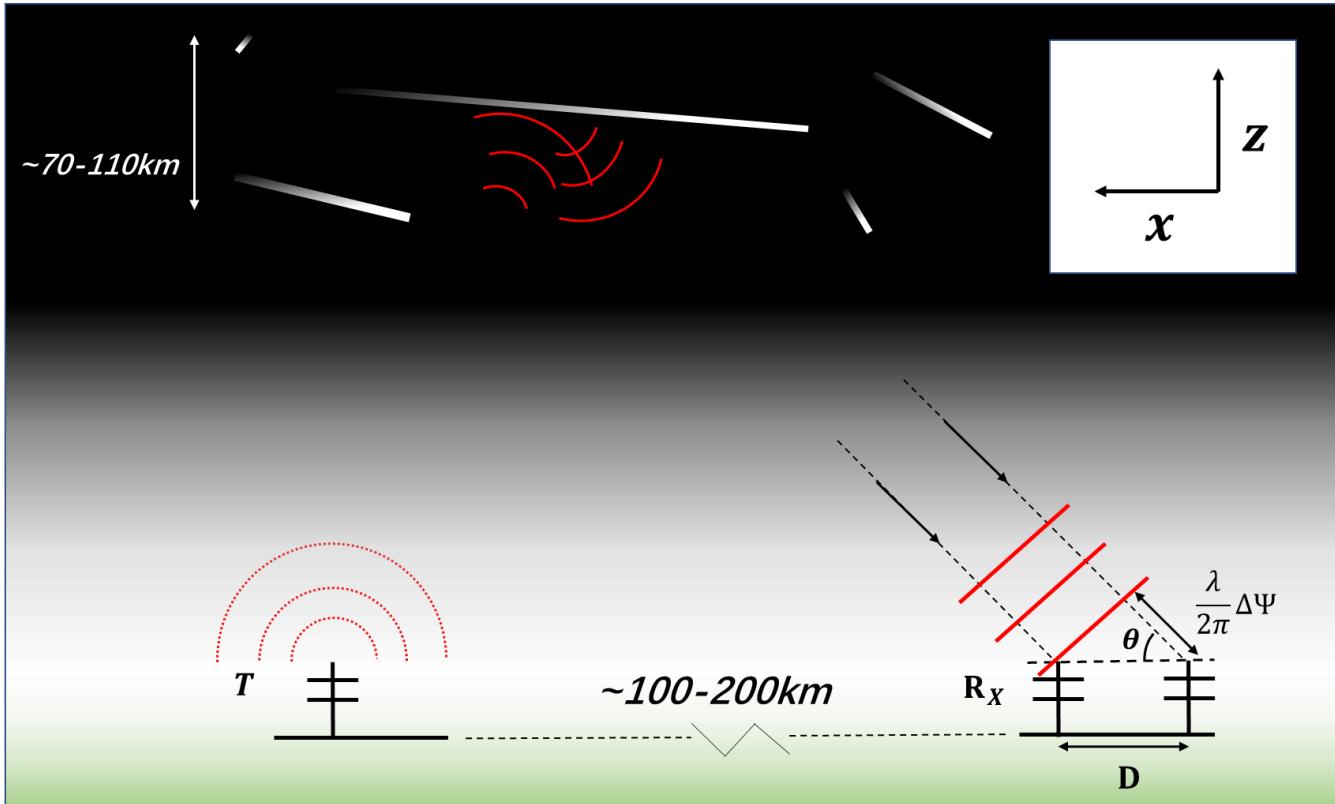
$$\cos \phi = \frac{\vec{R}_s \cdot \vec{n}_x}{R_s} \frac{\vec{n}_x}{\sin \theta}$$

(A6.6)

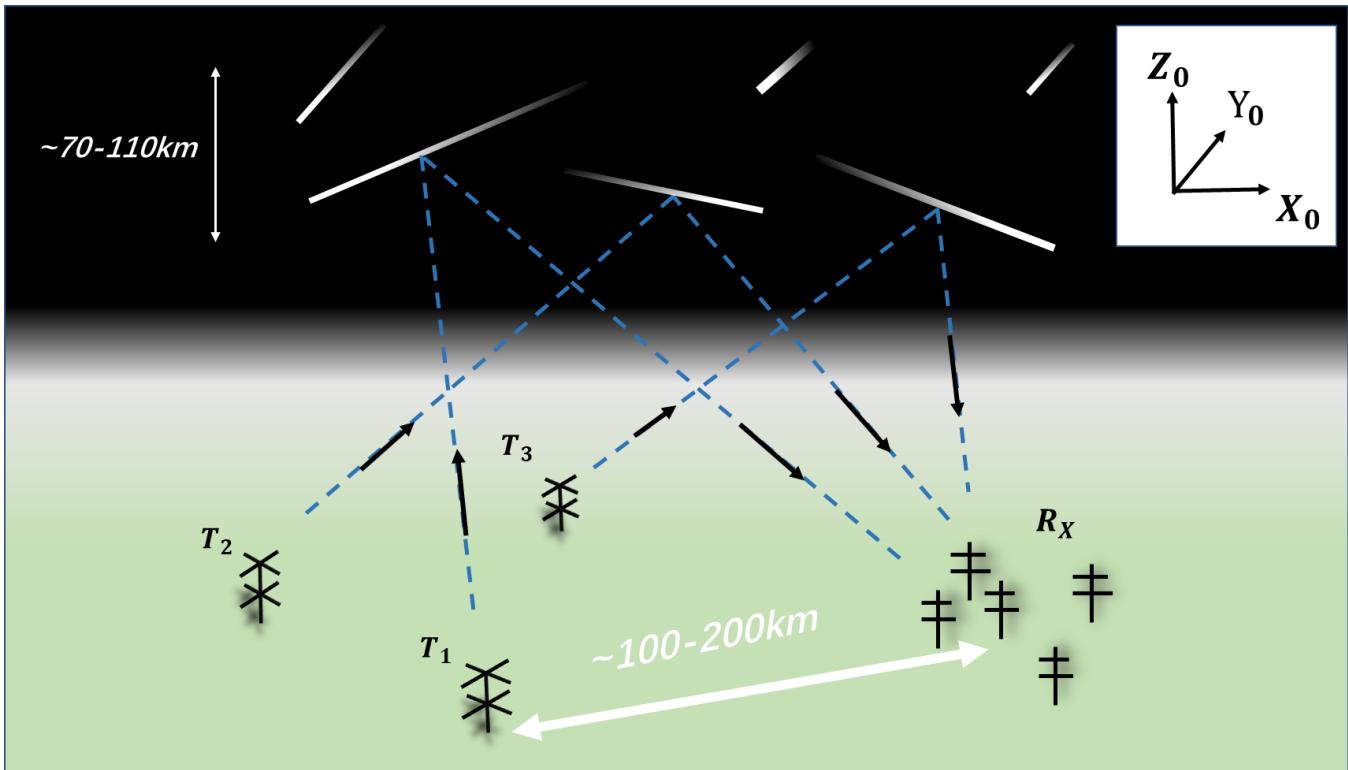
555 $\sin \phi = \frac{\vec{R}_s \cdot \vec{n}_y}{R_s} \frac{\vec{n}_y}{\sin \theta} \quad (A6.7)$

For $0^\circ < \theta < 180^\circ$ and $0^\circ \leq \phi < 360^\circ$. When $\theta = 0^\circ$, we handle it as same as in **Appendix (A.2)**.

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565 Figure 1: Schematic diagram of the simplified bistatic configuration used in Hocking's vertical resolution analysis (Hocking, 2018).
 The two receiverreceiving antennas and a transmitterthe transmitting antenna are collinear. The analysis is in a 2-dimensional vertical section through the baseline—joining the antennas. The radio wave is scattered byfrom a few Fresnel zones of several
 570 kilometres longlength around thespecular point inon the meteor trail and received by receiverthe receiving antennas. The cross-correlation analysis between receiverthe receiving antennas can be used to solve for the AoAs. The fact thatBecause the radio wave
bounced backis reflected from a region a few Fresnel zones willcauseinlength the measured phase difference between the receiver
 antenna pairdeviatingpairs todeviates from the ideal phase difference. The ideal phase difference will solve an AoAs pointing to
MTSP. This deviation from the ideal phase difference is one of the error sources ofin the PDME. In this work, we solve for the ideal
phase difference associated with the AoA directed to the MTSP.

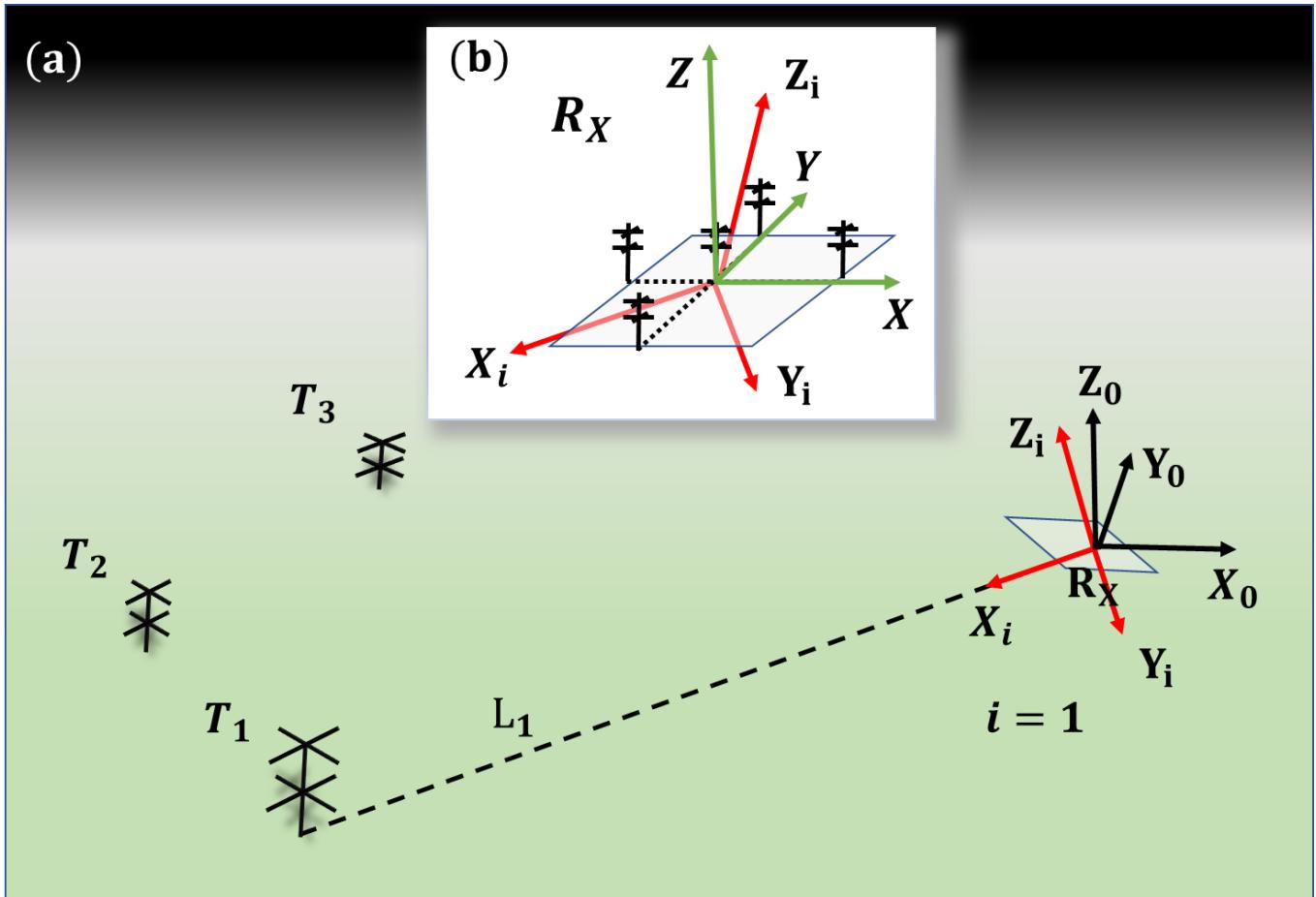


580 **Figure 2: Schematic diagram of a multistatic meteor radar system using SIMO (single-input and multi-output). There are three transmitters (T_1, T_2 and T_3) and one receiver (R_X) in the picture. The transmitter/receiver distance is usually approximatelytypically 100-200 km. X_0, Y_0, Z_0 represents the east, north and up directions of the receiver-receiving antenna. Over 90% of the received energy comes from about one kilometre around the specular point of the meteor trail, which is slightly less than the length of the central Fresnel zone (Ceplecha et al., 1998).**

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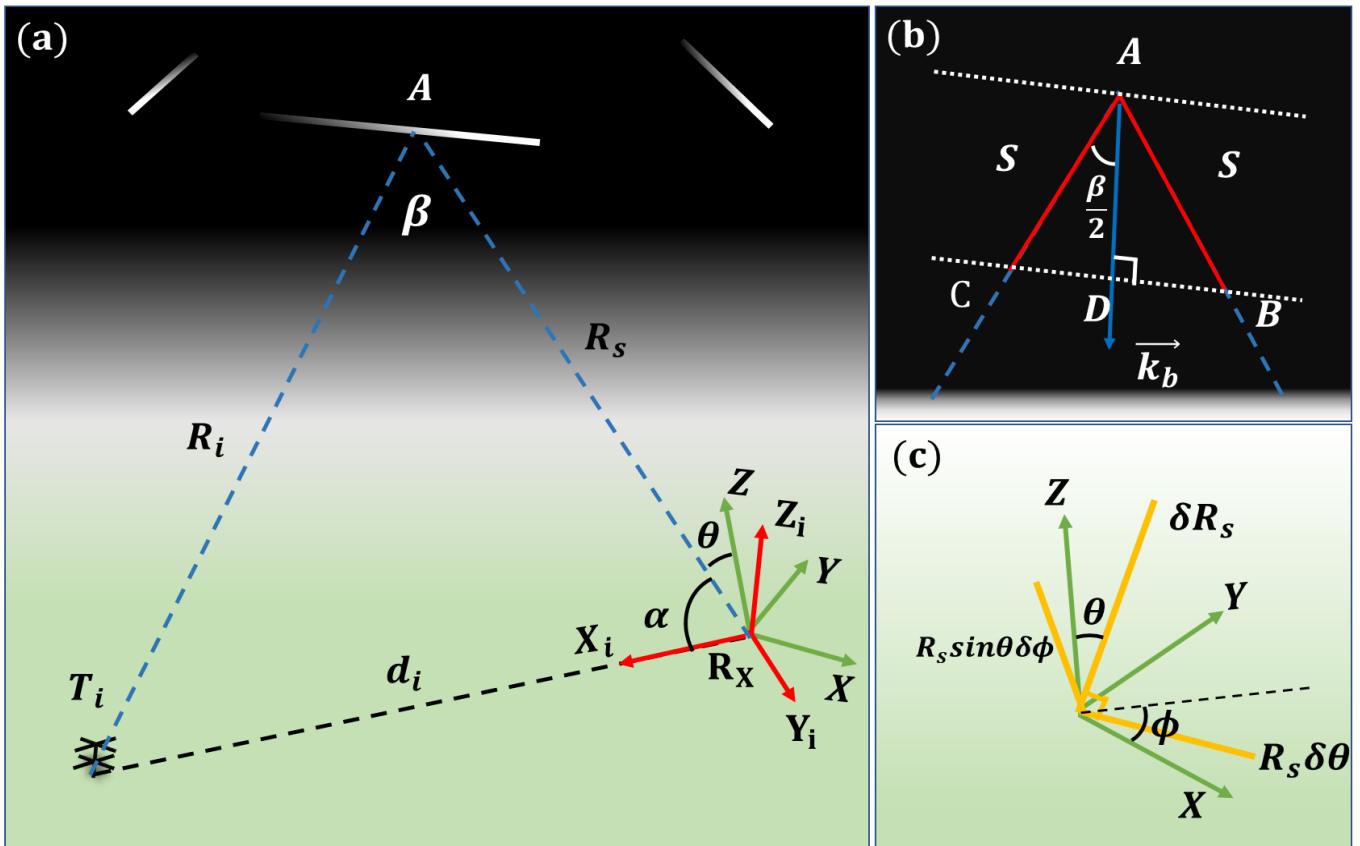
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600 Figure 3: (a) Schematic diagram of the three introduced-coordinate systems used in this work. $X_i Y_i Z_i$ are a class of coordinate systems whose axis- X_i pointpoints to transmitter T_i , with, i . And in this picture, i are = 1,2,3. $X_0 Y_0 Z_0$ is the ENU coordinate system and to which all errors will beare compared in this coordinate. (b) Magnified plot of the receiver-receiving array. XYZ is fixed on the receiver horizontal plane. Axis-X and Y are collinear to with the two arms of the antenna arraysarray.

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615 Figure 4: (a) Schematic diagram of ~~at the~~ forward scatter geometry for the radar link between T_i and R_X . Point-A is the MTSP. (b) Magnified plot of specular point A. The red line represents a radio wave pulse, and S is the half ~~wave~~-pulse length. \vec{k}_b is the Bragg vector which halves ~~the~~ forward scatter angle β . (c) Schematic diagram of E_1 in XYZ, which can be decomposed into three orthogonal vectors.

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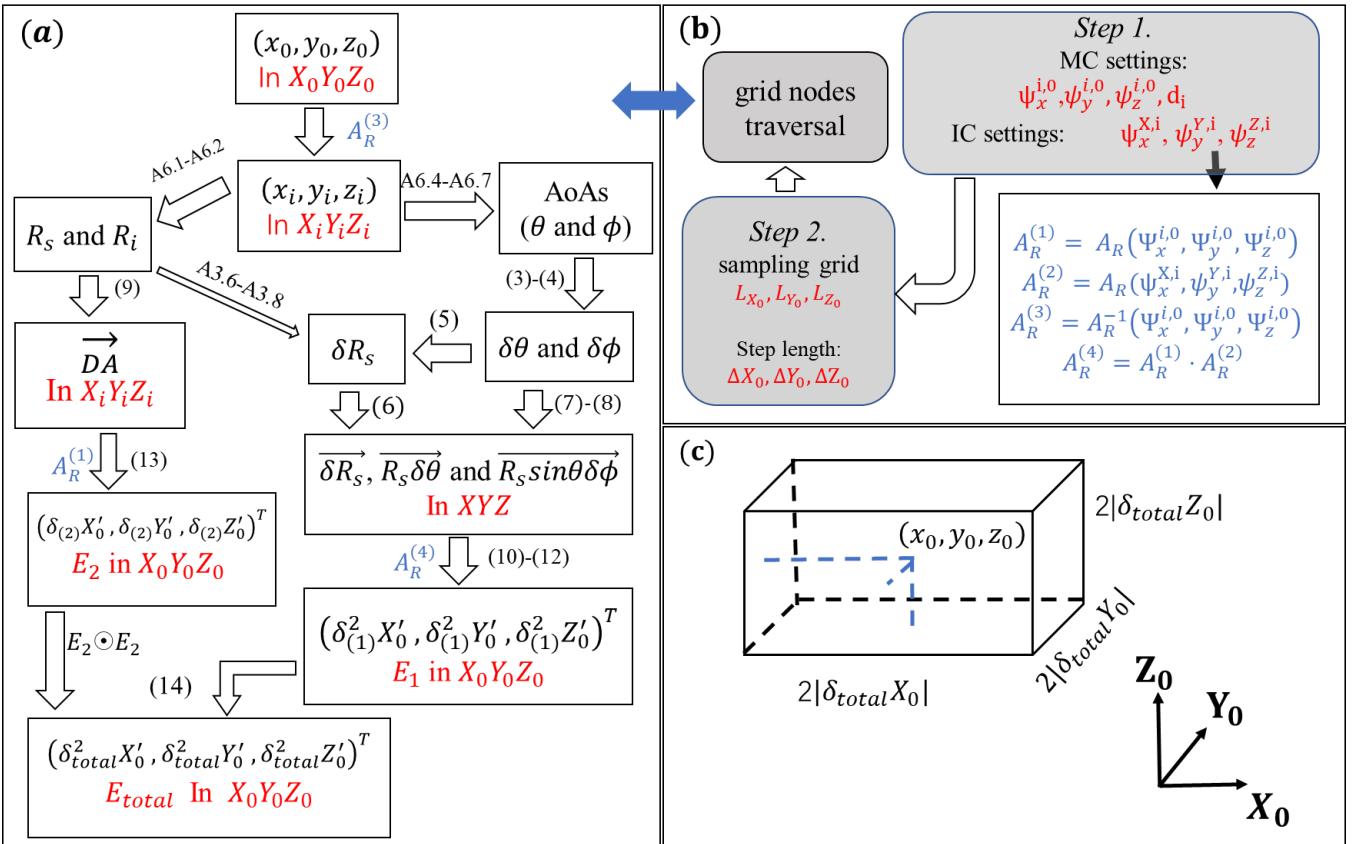


Figure 5: (a) the flow chart of the location error calculation process for a point in $X_0Y_0Z_0$. The marksnotation beside arrows represent the corresponding equations (black) or coordinate rotation matrix (blue) in the paper. “ \odot ” is the Hadamard product.

630 Thus $E_2 \odot E_2$ will getyield $(\delta_{(2)}^2 X_0, \delta_{(2)}^2 Y_0, \delta_{(2)}^2 Z_0)^T$. (b) the flow chart of the program to calculate the location errors distributions for a radar link L_i . This process includes parameters settings for a radar link, generating; the generation of the sampling grid nodes and the traversing of all the nodes. For each node, the program uses the calculation method described in (a). MC: is the multistatic configuration, IC: is the interferometer (receiver antennasreceiving antenna) configuration. (c) Schematic diagram of the relationship between the spatial resolution and the total location errors of the MTSP. For a detected point in space, the MSE of MTSP’s location errors is $\pm|\delta_{total}X_0|$, $\pm|\delta_{total}Y_0|$, $\pm|\delta_{total}Z_0|$ in the zonal, meridianmeridional and vertical directions, respectively. This means that the actual specular point might occur in a region which formforms a $2|\delta_{total}X_0| \times 2|\delta_{total}Y_0| \times 2|\delta_{total}Z_0|$ cube and the detected point is on the centralcentroid of this cube.

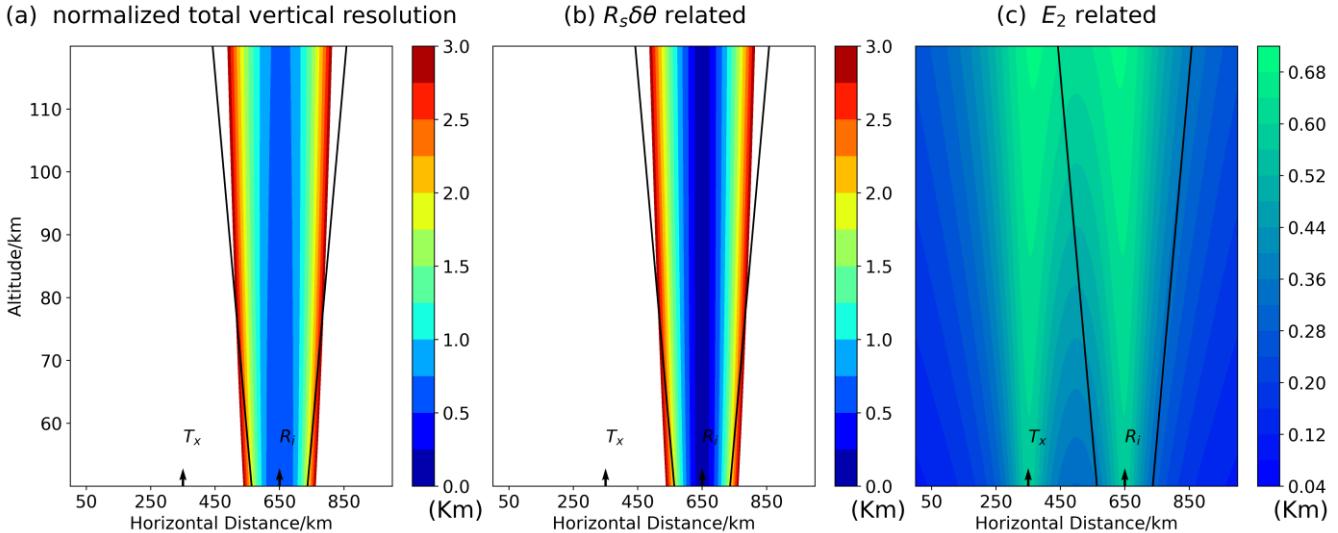


Figure 6: the normalized vertical resolution distribution in a vertical section from 50 km to 120 km height when ignore the error term “ δR_s ” is ignored. Panels (a), (b), and (c) are the total, $R_s\delta\theta$ related, and E_2 related normalized resolution distribution distributions, respectively. The These results is the same as those produced in Hocking’s work (Hocking, 2018). Two The two black arrows represent the positions right above the transmitter (Tx) and the receiver (Rx) and the transmitter/receiver separation is 300 km away. The region between the two black oblique lines is a the trustworthy sampling volume for the receiver receiving array because the elevation angle is beyond 30° with little to reduce influence of from potential mutual antenna coupling or from other obstacles in the surrounding area. Except the region in at large elevation angle angles (i.e., 90°), the E_2 related resolution values are much lower than the $R_s\delta\theta$ related errors. The $R_s\delta\theta$ related resolution distribution is depends only depend on the receiver receiving antennas. Thus, the total vertical resolution distribution is nearly unchanged with the variation of the transmitter/receiver distance varying. The normalized resolution values that exceed 3 km (which correspond 12 km vertical resolution aren’t) are not shown.

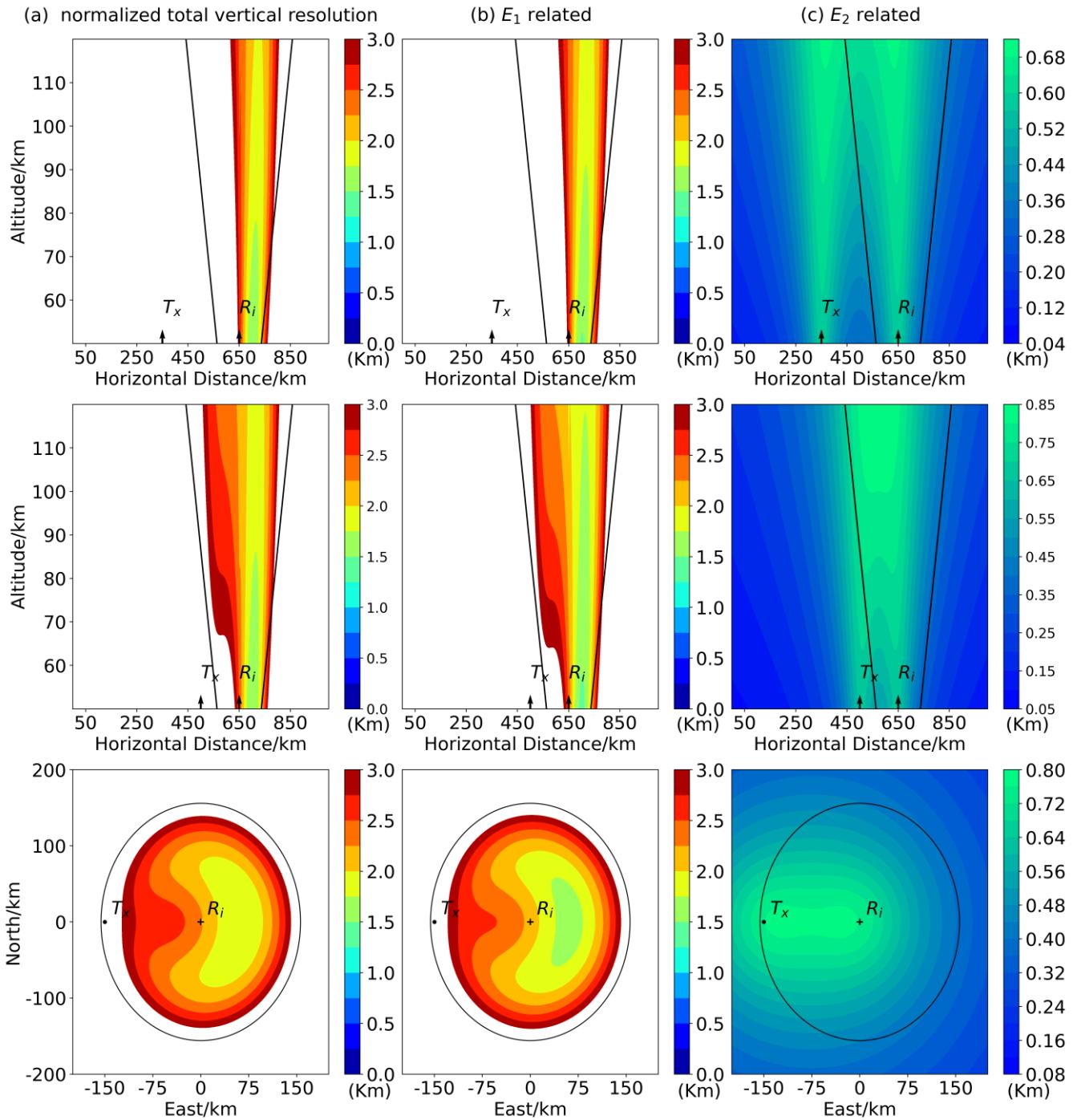
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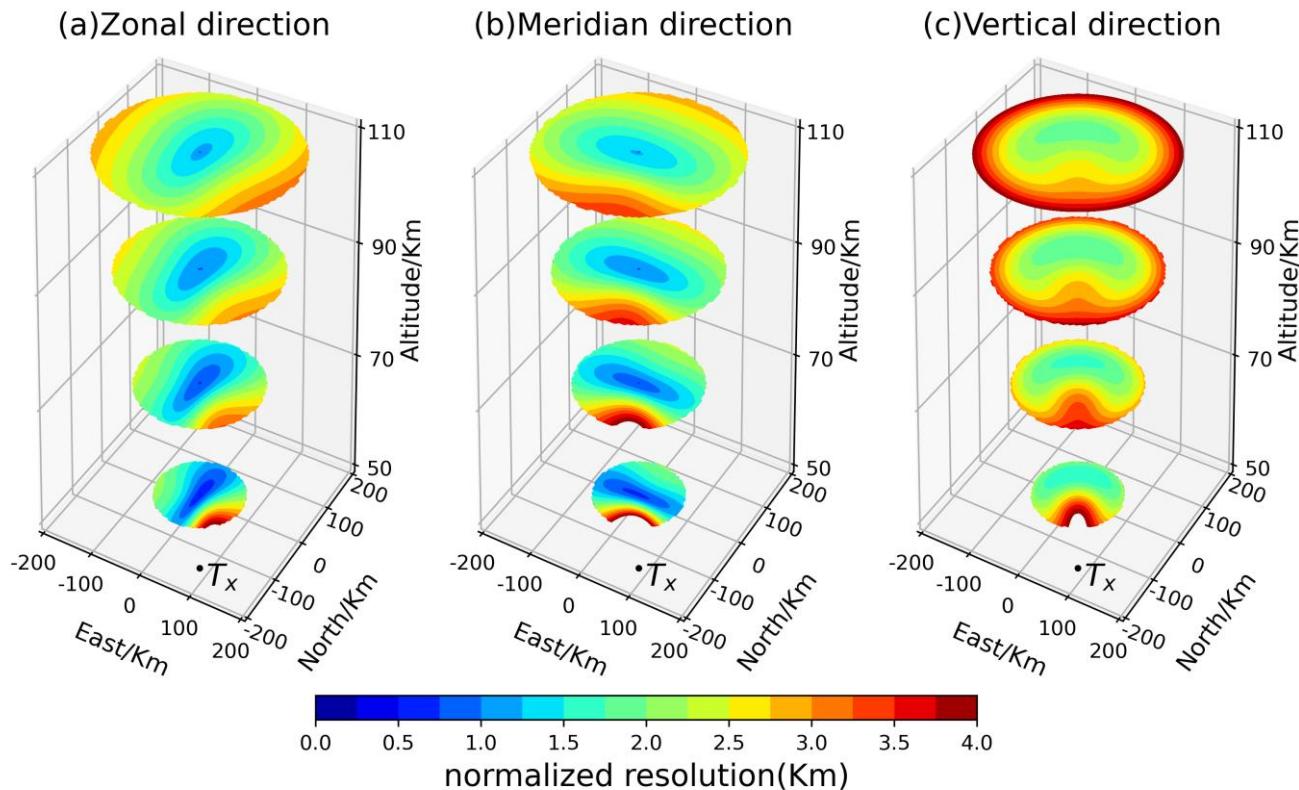


670 **Figure 7:** the normalized vertical resolution distribution using the analytical method [described](#) in this paper. The first and second [rows](#) represent a vertical section [from](#) height [from](#) 50 km to 120 km. The third row [represents](#) the horizontal section [at](#) 90 km and the [receiver](#) [array](#) is [on](#) the origin with positive coordinate [value](#) [represent](#) [east](#) [values](#) [representing](#) the [eastward](#) or [north](#) [direction](#), [northward](#) [directions](#), [respectively](#). The first row has the same parameters settings as Figure 6 and is used to compare with Figure 6. [The](#) E_1 [related](#) resolution will change with [the](#) transmitter/receiver configuration because it

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considers the error term “ δR_s ”. Thus, the total vertical resolution will change with the transmitter/receiver configuration. With the transmitter/receiver distance varying from 300 km (the first row) to 150 km (the second row), the total vertical resolution distribution is clearly changed. The third row is the perspective to the horizontal section at 90 km altitude for the second row. The normalized resolution values that exceed 3 km aren't are not shown.

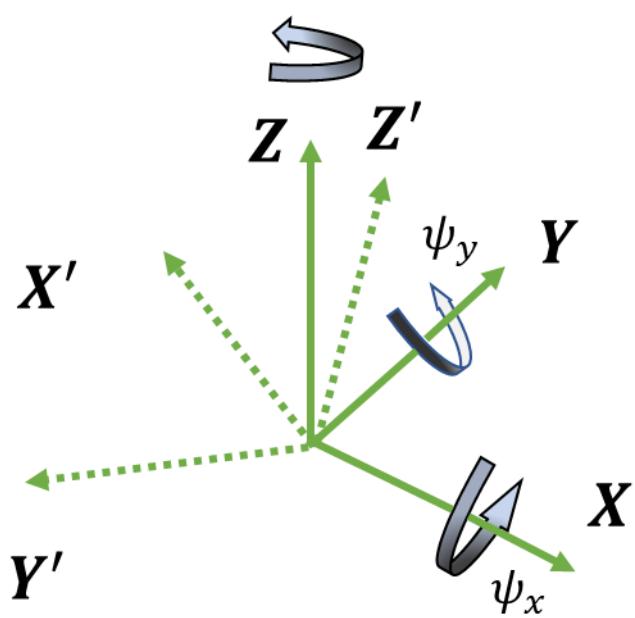
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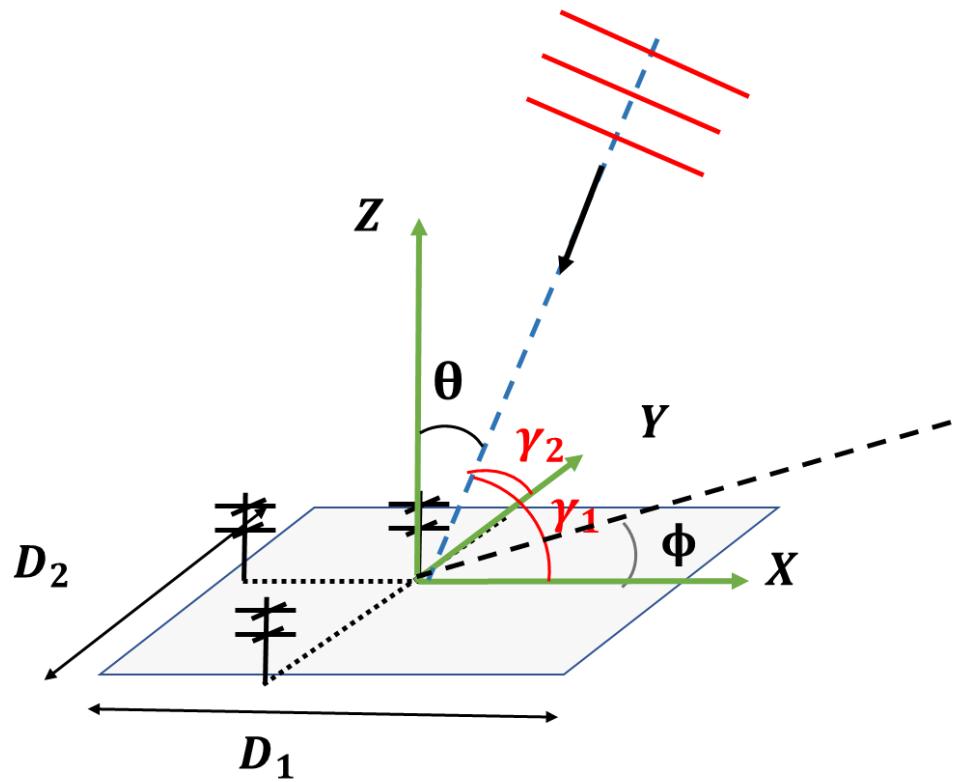
Figure 8: the 3D contour plot of the normalized resolution distribution for a multistatic radar link whose baseline length is 180 km and whose transmitter is south by east 30° of the receiver. The black dots represent the position right above the transmitter and the receiver receiving array is on at the origin of the axes. (a), (b) and (c) are the normalized resolution distribution distributions in the zonal, meridian and vertical directions, respectively. The subplot's four slice circles from bottom to top are the horizontal section in 50 km, 70 km, 90 km and 110 km height, respectively. The region whose elevation angle of the receiver is less than 30° isn't is not shown and therefore the slice circles become larger from the bottom to the top. The normalized resolution values that exceed 4 km (which correspond corresponds to 16 km resolution aren't are not) are not shown.

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Figure A.1



Receiver

700 Figure A.2 (two The receiving array geometry (only three antennas are not shown for concise clarity))