

## Answers to Reviewer 2

**General comments:** This is a tough paper to review. On one hand the measurement setup and collected dataset is very interesting. On the other hand, there are several things that causes concerns.

I understand the problem the authors are facing and that is behind one of the weak points in the study. The authors are using the “soft spheroid” particle model in combination with T-matrix to compute single-scattering ice particle properties that are used for the retrievals. However, given a large number of articles published in the last decade that argue that such approach could lead to significant errors, this argument is not easy to make. It is not impossible, since one could argue that 35 GHz measurements are not suffering from the “non spheroidal” effect that much. However, the authors tried to avoid the existing literature and use not well supported arguments, see detailed comments below. I suggest that they improve this part of the paper. Please use more up-to-date literature and make your arguments using the current understanding of the problem.

Thank you a lot for your comment. We agree with you that our ice model is a weak point in this study. In this study, however, the authors major intension was to investigate the feasibility two combine two spatially separated radars to obtain microphysics information about the ice hydrometeors detected into clouds, above the melting layer. The combination of two radars located in different areas is a challenging task due to the different radar scanning geometries, time and spatial mismatch of the radar beams and different path integrated attenuation (PIA) along the radar beams. In our feasibility study we aim to examine the combination of radar measurements from a weather and a cloud radar within a common retrieval framework for particle size, mass and shape. For this purpose, we develop an ice microphysical retrieval for the estimation of the median size, the apparent shape and the ice water content of the detected ice populations making some a-priori assumptions about their particle size distribution and mass-size relation that they follow. Hence, to avoid the addition of more assumptions in the developed ice retrieval, we selected a simple ice particle model, i.e., the soft spheroid, which is known to require the least assumed parameters compared to more complex ice particle models.

We now dedicated a new subsection (1.2 *Representation of ice atmospheric hydrometeors using spheroids*) to include more up-to-date literature and to make a more convincing argument for the use of the simple soft spheroid approximation.

1) The revised manuscript now includes the following paragraph with a literature review:

*“Single scattering simulations are an indispensable tool to bridge the gap between microphysical properties of hydrometeors and polarimetric radar observations. In the case of ice particles, however, the calculation of scattering properties can be challenging due to their large complexity, variety in shape, structure, size and density. One of the most sophisticated methods, the Discrete-Dipole Approximation (DDA, Draine and Flatau, 1994), can be used to calculate the scattering properties of realistic ice crystals and aggregates. However, this approximation can be computational demanding. To reduce computation cost and complexity, ice particles are often assumed to be spheres and their scattering properties are calculated using the Mie theory or they are assumed to be spheroids using the T-Matrix method (Waterman, 1965) or the Self-Similar Rayleigh-Gans Approximation (SSRGA, e.g., Hogan and Westbrook, 2014; Hogan et al., 2017; Leinonen et al., 2018a) for scattering simulations. The calculations when SSRGA is used are known to be affected by the way that ice mass is distributed throughout the particle’s volume. As we aim for a simple ice particle model, we extensively used the T-Matrix method in this study, assuming the ice particles to be soft spheroids. It is a common approach in model studies that ice particles are represented by homogeneous spheroids with density equal or smaller of bulk ice. Due to its simplicity, the limitations of the spheroid approximation have been a heavily researched and*

*debated topic in the last decade. While Tyynelä et al. (2011) showed an underestimation of the backscattering for large snowflakes, Hogan et al. (2012) suggested that horizontally aligned oblate spheroids with a sphericity of 0.6 can reliably reproduce the scattering properties of realistic ice aggregates which are smaller than the radar wavelength. Nevertheless, the same study also concluded that for larger particles spheroids are an improvement to Mie spheres which can lead to a strong underestimation of  $Z_e$  and, in turn, strong overestimation of IWC. Leinonen et al. (2012) showed that the spheroidal model cannot always explain the radar measurements as more sophisticated particle models do, e.g., snowflake models. Later on, Hogan and Westbrook, (2014) indicated that the soft spheroid approximation underestimates the backscattered signal of large snowflakes (1 cm size) – measured with a 94 GHz radar – up to 40 and 100 times for vertical and horizontal incidence, respectively. In contrast, the simple spheroidal particle model could successfully explain measurements of slant-45° linear depolarization ratio, SLDR, as well as SLDR patterns on the elevation angles (Matrosov, 2015) during the Storm Peak Laboratory Cloud Property Validation Experiment (StormVEx). In Liao et al. (2016) it was found that randomly oriented oblate ice spheroids could reproduce scattering properties in Ku- and Ka-band similar to these from scattering databases when large particles were assumed to have a density of  $0.2 \text{ g cm}^{-3}$  and a maximum size up to 6 mm. Furthermore, although Schrom and Kumjian (2018) showed that some ice crystal shapes as branched planar particles could be better represented by plate crystals than spheroids, the simple spheroidal model has been used in recent studies to represent ice aggregates as in Jiang et al. (2019) or to retrieve shape from LDR as in Matrosov (2020). In all these studies, it is recognized that the spheroidal model requires less assumed parameters compared to more complex particle models.”*

2) We also wrote a new paragraph to support our decision to use the soft spheroid approximation for this study nonetheless. You can find our answer to one of your more specific comments further below.

The second problem is related to how reliable the retrieved values are. Because of the measurement setup the radar observations volumes are mismatched. Potentially because of this, the observed DWR and retrieved Dm values show artificial looking patterns. The authors use the retrievals to generate statistics of particle properties, see Fig 12, which is one of the main stated goals of the study. I would like to see a discussion what retrieved values can be trusted and why. Ideally, problematic data should be excluded.

Thank you for pointing this out. The different locations of the two radars are an opportunity and challenge at the same time as we aim to obtain an oblique perspective on the cross-section area between both instrument. For this reason, we conducted a sub-study investigating all possible sources of errors and uncertainties in our measurement dataset, i.e., radar reflectivity, differential reflectivity and dual-wavelength ratio. An extended sub-section can be now found in our Sect. 3.1.2:

#### *“3.1.2 Assessment of radar observations errors*

*Radar measurements are often affected by systematic or random errors. To assess their impact on the ice microphysics retrieval developed in this study we need to investigate possible errors in POLDIRAD and MIRA-35 observations as well as all their sources.*

*The absolute radiometric calibration of both instruments is an important error source in DWR measurements. While the error of the absolute radiometric calibration of POLDIRAD is estimated to be  $\pm 0.5 \text{ dB}$  following the validation with an external device (Reimann, 2013), the budget laboratory calibration of MIRA-35 following Ewald et al. (2019) is estimated to be  $\pm 1.0 \text{ dB}$ .*

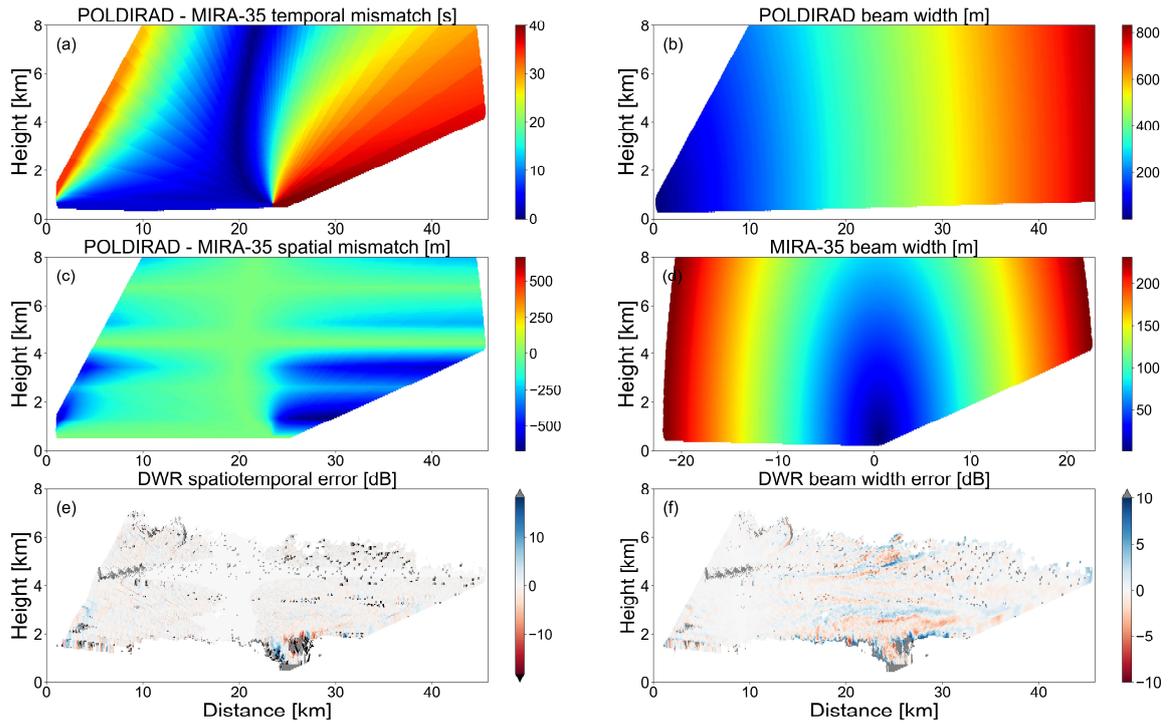
*In order to test for a systematic ZDR bias, we exploited POLDIRAD measurements during vertically pointing scans (also known as birdbath scans, e.g., Gorgucci et al., 1999) in a liquid cloud layer performed on the 4<sup>th</sup> April 2019. The measurements indicated that ZDR has an offset of about  $+0.15 \text{ dB}$  as ZDR values are expected to be near  $0 \text{ dB}$  for this case due to the spherical apparent*

shape of liquid droplets. In Fig. A1 (Appendix A), examples of radar reflectivity  $Z_e$ , differential reflectivity ZDR as well as a scatter plot showing the ZDR offset are presented.

Another error that should be considered is the random error, especially for ZDR measurements at low signal levels. To detect and filter out regions with high ZDR noise we compare the local (3 range gates) standard deviation  $ZDR_{stdv}$  with the local mean  $ZDR_{mean}$ . Subsequently, we only include regions where the signal  $ZDR_{mean}$  exceeds the noise  $ZDR_{stdv}$  by one order of magnitude. An example of this approach can be found in Fig. A2 (Appendix A). While we apply the retrieval to all cloud regions, the described ice mask and noise filters are used during the statistical aggregation of retrieval results.

In our case of spatially separated radar instruments, an azimuthal misalignment between both instruments had to be excluded to obtain meaningful DWR measurements. To this end, we performed several solar scans with both instruments in spring 2019 to confirm their azimuthal pointing accuracy (e.g., Reimann and Hagen, 2016). Here, we found an azimuth offset of  $-0.2^\circ$  for POLDIRAD and an azimuth offset of  $+0.1^\circ$  for MIRA-35. Consecutive solar scans confirmed the azimuthal pointing accuracy within  $\pm 0.1^\circ$ . Despite the small azimuthal misalignment, the radar beam centroids of both instruments were clearly within the respective other beam width during our measurement period in 2019.

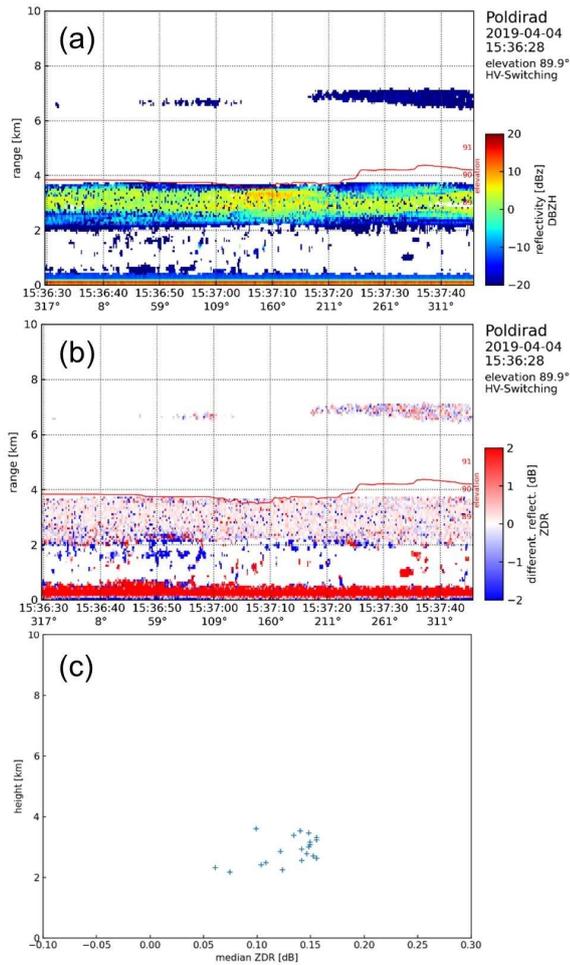
Besides an azimuthal misalignment, we also analyzed the temporal mismatch between both RHIs as well as the volumetric mismatch in the context of non-uniform beam filling. Although the RHIs from the radars were scheduled to be executed simultaneously, regions within the RHIs are measured at slightly different times by both instruments. This temporal mismatch can lead to slightly different  $Z_e$  radar observations from both radars in the context of horizontal advection of an inhomogeneous cloud scene. In the following we used this temporal mismatch to estimate the resulting DWR error for the example case shown in Fig. 3. Using wind data (Fig. 4) from the Oberschleißheim sounding station (source: Deutscher Wetterdienst, data provided by University of Wyoming; <http://weather.uwyo.edu/upperair/sounding.html>, last access: 10 June 2021), we converted the temporal mismatch (Fig. 5a) between the radar measurements for each pixel in the common radar grid to a spatial difference (Fig. 5c). To estimate the impact of this spatiotemporal mismatch (hereafter spatiotemporal error) we subsequently used these spatial differences to calculate DWR values between pixels in the spatially higher resolved MIRA-35  $Z_e$  measurements (Fig. 5e). Concluding the DWR error assessment, we also analyzed the volumetric mismatch caused by the different beam widths of the two radars. For spatially heterogeneous scenes, this volumetric mismatch can lead to artificial DWR signatures caused by a non-uniform beam filling. Here, the spatially higher resolved MIRA-35  $Z_e$  measurements (30 m range gate length) along the RHI cross section were used as a proxy to obtain the spatial heterogeneity of  $Z_e$  perpendicular to the RHI cross section. In a first step, the local beam diameters for each pixel in the common grid are calculated for POLDIRAD (Fig. 5b) and MIRA-35 (Fig. 5d). Then, moving averages along the  $Z_e$  cross sections from MIRA-35 are performed using the corresponding local beam diameters. Hence, at each pixel of the common radar grid two averaged MIRA-35  $Z_e$  values are obtained; one corresponding to the local beam diameter of MIRA-35 and one corresponding to the local beam diameter of POLDIRAD. Subtracting the averaged  $Z_e$  for each pixel, we were able to estimate the error caused by the volumetric mismatch between both radar beams (Fig. 5f).



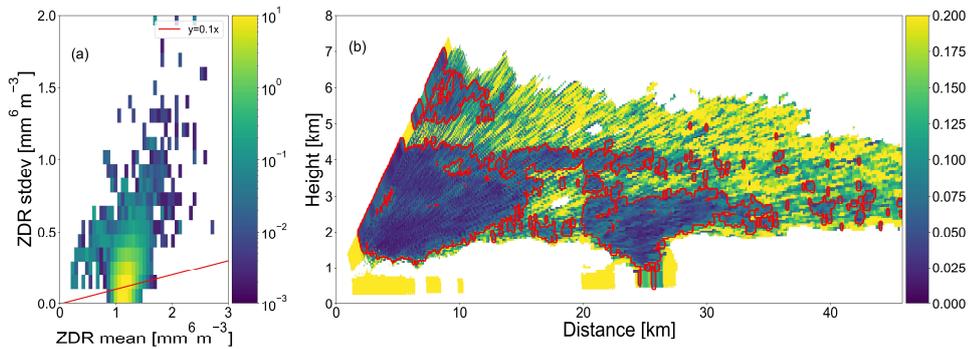
**Figure 1:** DWR error assessment due to temporal mismatch (left panels) and volumetric mismatch (right panels). In (a), (c) and (e) panels, the POLDIRAD and MIRA-35 temporal mismatch, the POLDIRAD and MIRA-35 spatial mismatch and the spatiotemporal error in dB are plotted. In (b) and (d) panels, the POLDIRAD and MIRA-35 beam widths are presented, while in panel (f) the estimated DWR error due to the volumetric mismatch is shown. For this plot the data from 30<sup>th</sup> January 2019 at 10:08 UTC are used. The ice masked and noise/error filtered values (except for random ZDR error) in (e) and (f) are plotted with grey color. Black color in panel (e) denotes the additional missing values due to the spatial shift of the radar grid. For better visualization purposes the isotherms of  $-5\text{ }^{\circ}\text{C}$ ,  $-15\text{ }^{\circ}\text{C}$  and  $-25\text{ }^{\circ}\text{C}$  are not plotted here. “

This sub-section is accompanied with a respective Appendix A:

“Appendix A: Radar measurements error assessment

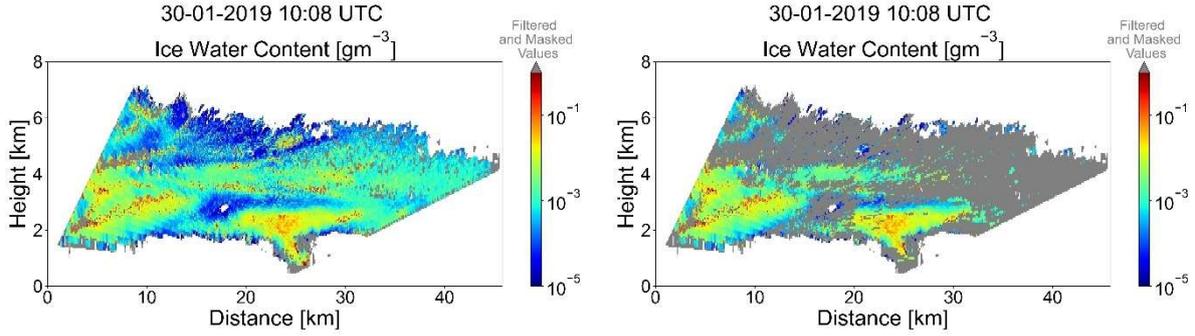


**Figure A1:** POLDIRAD (a)  $Z_e$  and (b) ZDR measurements for different times and azimuth angles of a liquid cloud layer with a vertical pointing antenna on 04 April 2019. Panel (c) shows the offset of the averaged ZDR for the range where the liquid layer was detected.



**Figure A2:** (a) the local standard deviation of ZDR is plotted as a function the local mean of ZDR. (b) the ratio  $a = ZDR_{stdev}/ZDR_{mean}$  can be used to filter out noisy ZDR measurements. In the red encircled areas ( $a < 0.1$ ), the retrieval results are considered to be reliable enough to be aggregated into statistical results.”

Therefore, only measurements fulfilling the aforementioned criteria are considered for statistical results of the ice microphysics retrieval. For visualization purposes in the case study, we only applied the ice mask and noise/error filter using grey colors, and left later filtered regions due to ZDR error. In the following figures we demonstrate the effect of the ZDR filter, with IWC results shown throughout the manuscript (left panel) and after the filter application of the ZDR noise filter (right panel). Only the valid values of the right plot are these considered in the statistics of this paper.



Finally, there are statements in the manuscript that are not correct. A good example is the definition of the reflectivity factor. These should be corrected.

Thank you for your correction. These formulas are now better described as:

*“The radar reflectivity factor  $z$  is defined as the sixth moment of the particle size distribution  $N(D)$  and is thus designed to be proportional to the to the Rayleigh scattering cross section of small – size much smaller comparing to the radar wavelength – liquid spheres:*

$$z [\text{mm}^6 \text{m}^{-3}] = \int_0^{\infty} N(D)D^6 dD \quad (1a)$$

*where,*

*$z$ : the radar reflectivity in linear scale,*

*$N$ : the number concentration,*

*$D$ : the geometric diameter of the particles.*

*This formula can be also expressed in logarithmic terms:*

$$Z [\text{dBZ}] = 10 \cdot \log_{10}(z). \quad (1b)$$

*This definition, however, cannot be directly applied to snow due to the varying density, the irregular shape and larger size of ice particles which cause deviations from the Rayleigh into the Mie scattering regime. Moreover,  $N(D)$  for ice particles is referring to the size distribution of their melted diameters. Nevertheless, an equivalent radar reflectivity factor  $Z_e$  can be derived from the measured radar reflectivity  $\eta$  ( $\eta = \sum_{Vol} \sigma_n$ ; normalized to a specific volume summation of backscattering cross-section of all detected hydrometeors) when the dielectric factor of water  $|K|^2 = 0.93$  is assumed:*

$$z_e [\text{mm}^6 \text{m}^{-3}] = \frac{\lambda^4}{\pi^5 |K|^2} \cdot \eta \text{ and } Z_e [\text{dBZ}] = 10 \cdot \log_{10}(z_e) \quad (1c)$$

*where:*

*$\lambda$ : the radar wavelength. In the Rayleigh regime, the radar reflectivity factor  $Z$  or the equivalent radar reflectivity factor  $Z_e$  (for simplicity reasons referred also as radar reflectivity in this paper) is proportional to the sixth power of the particle size, while in the Mie regime  $Z_e$  scales with the second power of the particle size. In both regimes  $Z_e$  scales linearly with the particle number concentration.”*

Overall, in my opinion the manuscript needs significant improvements before it can be considered for the publication. I encourage the authors to revise it significantly and resubmit.

**Specific comments:**

Line 30 "Ice clouds can cause a cooling effect at the surface by reflecting the shortwave, incoming solar radiation but they can also contribute to warming of the atmosphere by trapping the longwave, terrestrial radiation (Liou, 1986)."

Are you sure about this statement? To my knowledge, ice clouds have a net warming effect. It is possible that in some particular cases they would lead to cooling, but the warming effect is more common.

Thank you a lot for this comment. The sentence is now rephrased providing more general information about the role of ice clouds in the Earth's energy budget:

*"Ice clouds are known to reflect the shortwave, incoming solar radiation, but they can also trap the longwave, terrestrial radiation interfering to the Earth's energy budget (Liou, 1986). Their influence on the radiation budget of the climate system strongly depends on their top height as well as on ice crystals habits and effective ice crystal size (Zhang et al., 2002)."*

Line 49-50: "Another way to gain microphysics information is to use multi-frequency radar observations as they exploit the scattering properties of ice particles in both Rayleigh and Mie regime."

Strictly speaking, Mie regime is not a correct term, a better one is the resonance scattering regime or non-Rayleigh scattering. The Mie solution is only applicable to spherical droplets.

Thank you for pointing this out. The sentence is now rephrased as "Another way to gain microphysics information is to use multi-frequency radar observations as they exploit the scattering properties of ice particles in both Rayleigh and non-Rayleigh regime."

Equation 1a – This equation is derived assuming small (much smaller than the wavelength) spherical water droplets.

The phrase "small – size much smaller comparing to the radar wavelength –" is now added to the sentence describing Eq. 1a.

Line 64-65 "Similarly, the equivalent radar reflectivity factor  $Z_e$  can be calculated when the radar reflectivity  $\eta$  is measured as well as the dielectric factor of water  $|\delta - \frac{3}{4}|^2 = 0.93$  and Rayleigh scattering is assumed: "

This is incorrect. Please check the definition of the radar reflectivity factor and equivalent reflectivity. Also, what is  $\eta$ , you have not defined it.

Thank you a lot for your comment. The sentence is now rephrased and the definition of measured radar reflectivity  $\eta$  is now added.

Line 96, Straka's paper is not presenting a method to distinguish hydrometeor types but summarizes characteristics of hydrometeors and corresponding dual-polarization radar signatures. There are many papers that actually present the method.

Thank you, that sentence is now rephrased and states that Straka et al. (2000) had summarized the characteristics of atmospheric hydrometeors at a wavelength of 10 cm.

Line 107 "Snow hydrometeors with size in the millimeters order of magnitude are found to have low densities and therefore, the Self-Similar Rayleigh-Gans Approximation (SSRGA, e.g. Hogan and

Westbrook, 2014; Hogan et al., 2017; Leinonen et al., 2018a), that is applicable for "soft spheres", can be used. ”

Actually, the exact opposite is argued in (Hogan et al., 2017; Leinonen et al., 2018a) and many other studies i.e. by Kneifel et al. (2015). SSRGA is introduced exactly because we cannot use the “soft sphere or spheroid” approximation. SSRGA takes into account internal distribution of ice particle mass that in its turn affect RGA’s form-factor and therefore the scattering properties. “Soft-sphere or spheroid” approximation assumes that mass is uniformly distributed, and inclusions are much smaller than the wavelength and hence an effective media approximation, such as Maxwell Garnett, can be used. These approximations seem to fail at mm-wavelengths.

Thank you for this valuable comment. After an extended research on that topic, we corrected our statement. We have now made clear the difference between the Self-Similar Rayleigh-Gans Approximation and PyTMatrix. This part of the text is now modified as follows:

*“One of the most sophisticated methods, the Discrete-Dipole Approximation (DDA; Draine and Flatau, 1994), can be used to calculate the scattering properties of realistic ice crystals and aggregates. However, this approximation can be computational demanding. To reduce computation cost and complexity, ice particles are often assumed to be spheres and their scattering properties are calculated using the Mie theory or they are assumed to be spheroids using the T-Matrix method (Waterman, 1965) or the Self-Similar Rayleigh-Gans Approximation (SSRGA; e.g., Hogan and Westbrook, 2014; Hogan et al., 2017; Leinonen et al., 2018a) for scattering simulations. The calculations when SSRGA is used are known to be affected by the way that ice mass is distributed throughout the particle’s volume. As we aim for a simple ice particle model, we extensively used the T-Matrix method in this study, assuming the ice particles to be soft spheroids.”*

Line 110: ” A well-proven approach is the soft spheroid particle model which uses the effective medium approximation (EMA) to model the refractive index of ice crystals and aggregates, e.g. the Bruggeman or Maxwell-Garnett models as in Garnett and Larmor (1904). ”

What are the assumptions and for which conditions EMA are proven to work? This is a blank statement implying that EMA always work regardless of conditions.

Thank you for this comment. More information about EMA along with references is now added in Sect. 3.2.1 *Soft spheroid model*. In particular, the following paragraph is now added:

*“Our soft spheroid model uses the effective medium approximation (EMA) to model the refractive index of the composite material as an ice matrix with air inclusions following the Maxwell-Garnett (MG) mixing formula given in Garnett and Larmor (1904):*

$$\frac{e_{eff}-e_i}{e_{eff}+2e_i} = f_i \frac{e_i-e_m}{e_i+2e_m} \quad (4)$$

*with,*

*$e_m, e_i$ : the permittivities of the medium and the inclusion, respectively,*

*$e_{eff}$ : the effective permittivity,*

*$f_i$ : the volume fraction of the inclusions.*

*The complex refractive index,  $m$ , is then calculated from  $m = \sqrt{e_{eff}}$ . In the framework of the EMA, the electromagnetic interaction of an inhomogeneous dielectric particle (components with different refractive indices) can be approximated with one effective refractive index of a*

*homogeneous particle (e.g., Liu et al., 2014; Mishchenko et al., 2016). In Liu et al. (2014), internal mixing was proven to best represent the scattering properties of hydrometeors. Here, the refractive index is modelled as an internal mixing of ice with air inclusions which are arranged throughout the ice particle. The same work also pointed out that the size parameter  $D_{crit} = \frac{\pi d}{\lambda}$  for each of these air inclusions should not be larger than 0.4 (with  $d$  as the diameter of the inclusion). “*

Line 112 : “Many studies, e.g. Hogan et al. (2012), have ...”

There are many studies after that which argue the opposite. As I have mentioned above the development of SSRGA was motivated by the inability of a “soft spheroid” particle model to reproduce the observations. You should use more recent literature to argue the point and make more convincing argument.

Thank you for pointing out this deficiency in the old version of the manuscript. As we already mentioned to one of your general remarks, we now included more recent literature in a new, dedicated subsection (Sect. 1.2). We would like to answer your second remark to make a more convincing argument why we use the soft spheroid approximation together with your next remark:

Line 122 and up to line 116 “Moreover, using spheroids we can better understand the ambiguities between the aforementioned degrees of freedom. Here, more sophisticated models of specific ice crystals could introduce additional challenges to sort a collection of ice shapes along these degrees of freedom or to define variables like the aspect ratio.”

The aspect ratio and orientation angle are part of the “soft-spheroid” particle model. These parameters are irrelevant for a more complex particle representation. One should keep this in mind, while interpreting observations and using different particle models for such interpretations. If you have selected “soft-spheroid” as your model of a more complex ice particle, then you should expect that your model is a (possibly over) simplified representation. Whether this is an advantage or not, it is a matter of discussion. So please make a stronger argument of your point?

Thanks for stressing this point! The new Sect. 1.2 *Representation of ice atmospheric hydrometeors using spheroids* also includes a paragraph in which we make a stronger argument why we use the soft spheroid approximation instead of a more complex particle representation:

*“Although more complex ice particle and scattering models are available, this work will use the soft spheroid approximation out of the following reasons: (1) In this work we aim to provide a feasibility study to combine two spatially separated radars to better constrain the ice crystal shape in microphysical retrievals using simultaneous DWR and ZDR observations from an oblique angle. Besides instrument coordination, the actual measurements and the assessment of measurement errors, the ice crystal and scattering model are just one component. Due to its simple and versatile setup, this work will utilize the soft spheroid approximation to study the benefit of additional ZDR measurements and the role of the observation geometry. (2) More importantly, to our knowledge, the more accurate SSRGA described by Hogan and Westbrook (2014) does not (yet) provide polarimetric variables used in this study, namely the ZDR. (3) In anticipation of a prognostic aspect ratio of ice crystals in bulk microphysical models (e.g., the adaptive habit prediction; Harrington et al., 2013), we aim to keep a minimal set of degrees of freedom to remain comparable with these modelling efforts. (4) Using ice spheroids we are able to vary parameters such as median size, aspect ratio and ice water content independently, which serve as degrees of freedom of the ice spheroids, and calculate their optical properties without*

*much computational cost as in other scattering algorithms (e.g., DDA) that are used in more realistic ice crystal shapes simulations. Moreover, using spheroids we can better understand the ambiguities between these simple, aforementioned degrees of freedom.”*

**Aspect ratio**, page 8: Please use commonly used definitions of AR. At the moment, it is very confusing.

Thank you for this comment. The text is now modified and contains aspect ratio (AR) and sphericity (S) definitions. Particularly, it has been updated as:

*“For the scattering simulations we assumed that ice hydrometeors can be represented by ice spheroids. The shape of the particles is defined using the aspect ratio parameter. In PyTMatrix this is the ratio of the horizontal to rotational axis of the particle. From the description of the simulated ice spheroids in Fig. 6, it is obvious that oblate (shaped like lentil) and prolate particles (shaped like rice) have AR larger and lower than 1.0, respectively, as z axis is selected to be the rotational axis. Using this principle, the representative value of sphericity = 0.6 for oblate ice spheroids from Hogan et al. (2012) is calculated as  $AR = 1.67$  for oblate ice particles in this study and therefore, this number was used as a reference value for the simulation plots (Fig. 7, Fig. 8 and Fig. 9a). In this study, we used S additionally to AR to compare retrieval results when the oblate and prolate shape assumption is used. S for oblates and prolates is found to be smaller than 1, while for spheres is equal to 1.”*

Line 276 “Furthermore, small oscillations out of this plane with a standard deviation of 20° are included to consider the flutter of ice crystals. ”

How? Please explain what you mean.

Thank you for your point. After your comment we have now replaced this sentence with the following part:

*“Here, all prolate ice particles were assumed to fall with their maximum diameter aligned to the horizontal plane. Hence, all ice prolates (hereafter referred as horizontally aligned prolates or horizontally aligned prolate ice spheroids) are rotated 90° (mean canting angle) in the yz plane (Fig. 6), while ice oblates have a 0° mean canting angle. The canting angle, i.e. angle between the particle’s major dimension and the horizontal plane, of the falling hydrometeors has been the topic of several studies. This value in nature is not so easy to estimate and thus, a standard deviation (e.g., 2°–23° as in Melnikov, 2017) is often additionally used. Here, we used a fixed standard deviation of 20° to describe the oscillations of the particles maximum dimension around the selected canting angle. Then the calculation of the scattering properties is performed using an adaptive integration technique for all possible particle’s geometries, ignoring the Euler angles alpha and beta of the scattering orientation.”*

Line 304, equation 6: What was the motivation of using the melted snowflake diameter?

Thank you for this comment. In Eq. 6,  $D_{eq}$  was only used as an interim parameter to obtain an  $m(D_{max})$  relationship since Yang et al (2000) only provided a fit for an equivalent volume sphere. The formula is now modified as:

*“... The  $D_{eq}$  is used to describe the diameter of a spherical water particle with the same mass as an ice particle with maximum dimension  $D_{max}$ .*

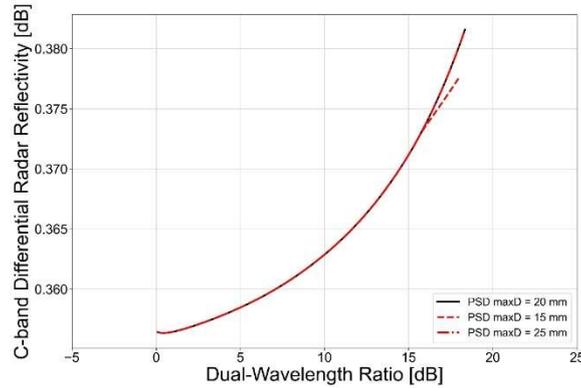
$$m(D_{max}) = \frac{\pi \cdot \rho_w \cdot D_{eq}^3}{6} = \frac{\pi \cdot \rho_w}{6} e^{\sum_{n=0}^4 b_n (\ln(D_{max}))^n} \quad (7)$$

where  $b_n$  is taken from Table 2 in Yang et al. (2000), the water density  $\rho_w = 1 \text{ g cm}^{-3}$  and  $D_{eq}$  as well as  $D_{max}$  are in microns.”

Page 11. Look-up table structure

Given that you are using the “soft spheroid” particle model and T-matrix to compute single scattering ice particle properties, I was expecting a discussion on how the refractive index is defined for different AR values. Are you preserving particle mass or density (and therefore the refractive index)? This should be explicitly discussed. What are PSD integration limits used in calculations? How did you select the maximum D value? The selection of maximum D has a direct influence on Zdr. Is that the reason why you are having an issue with reproducing Zdr observations?

Thank you for your comment on this topic. In our study a soft spheroid is a homogeneous mixture of ice and air, it’s ice mass is evenly distributed all over the spheroid’s volume and EMA is used for the refractive index calculation. As already quoted to your question regarding EMA above, we now included a more detailed description how we calculate the refractive index. To answer your question here, we are preserving particle mass and not density / refractive index when increasing the AR until reaching the bulk density of ice. Then particle mass is clipped. Regarding integration limits: for our work we a-priori chose the maximum diameter of PSD ( $D_{max}$ ) to be 20 mm, as we are interested to retrieve ice microphysics for ice particles that are detected within the cloud and above the melting layer and not for large aggregated snowflakes that reach the ground. In the following plot we calculated ZDR and DWR for AR = 1.67, shape parameter of PSD  $\mu = 0$  and median mass diameter,  $D_m$ , to vary between 0.1–3.02 mm.



We show that for even larger  $D_{max}$ , i.e.,  $D_{max} = 25 \text{ mm}$ , the simulated ZDR-DWR don’t differ from the results using  $D_{max} = 20 \text{ mm}$ . Minor differences in ZDR and DWR are observed when  $D_{max}$  is selected to be lower, i.e., 15 mm.

The ice spheroids construction is now added in the text (Sect. 3.2.1 *Soft spheroid model: Mass-size relation*). In particular:

“The maximum dimension,  $D_{max}$ , and the sphericity values for the spheroids were a-priori defined and their mass was calculated according to the formula that describes the relation between mass and  $D_{max}$  (mass-size relation), providing information about the mass of the ice crystals and therefore, their effective density with respect to their size. Mass  $m$  of an ice particle is usually connected to its maximum diameter  $D_{max}$  with a power-law formula,

$$m(D_{max}) = aD_{max}^b \quad (5)$$

where,

$a$ : the prefactor of the  $m(D_{max})$ , refers to the density scaling at all particles sizes,

*b: the exponent of the  $m(D_{max})$ , relates to the particles shape and growth mechanisms. With the mass and the spheroid dimensions known, the density of the ice spheroid was calculated. In the special case when the density was found to exceed the density of solid ice ( $0.917 \text{ g cm}^{-3}$ ), the mass of the spheroid was clipped and its density was set equal to the ice density. With the ice spheroid mass known, we also calculated the mass of the bulk ice spheroid (same-dimensions spheroid with density of solid ice) and from the ratio  $mass_{sph}/mass_{bulk\_sph}$  we calculated the ice fraction of the simulated spheroid. Using the ice fraction, we then calculated the particle's refractive index from the Maxwell-Garnet mixing formula."*

Fig 11, page 15 How physical are high  $D_m$  values closer to the cloud top? Please explain what data can be trusted. Ideally you should mask questionable data. This affects the results presented in Fig. 12.

Thank you for pointing this out. In the latest version of our manuscript in the already mentioned Sect. 3.1.2: *Assessment of radar observations errors* as well as in our Appendix A: *Radar measurements error assessment* you may find a detailed description of our approach to consider only reliable measurements in our ice microphysics retrieval and thus, for the development of Fig. 12 (current Fig. 14). As we already mentioned to one of your previous answers, the ZDR noise filter is not yet applied to the case study plots shown in Fig. 12 which would remove most of these large  $D_m$  values closer to cloud tops. For the statistical results shown in Fig. 14, we applied all described filters.

Page 17. Unknown mass-size relationship

Please explain, what was the logic behind the selection of the  $m(D)$  relations. Are they representative of the events you have observed, i.e. represent particles observed in this temperature regime? Are they sufficiently different to cover a possible range of values? On lines 290–300 you just state that you use them without much discussion why.

You are right, the motivation for BF and Yang was not really described until now. As we laid this out as a feasibility study of the measurement combination, we aimed for two sufficiently different  $m(D)$  relationships to get a better understanding of the limitations when soft spheroids are used for that task. Since BF95 was heavily used in previous studies we started off with their  $m(D)$  to represent more 2D ice crystals ( $b=1.9$ ) resulting in soft spheroids of lower density. To contrast them, we choose the  $m(D)$  analog to the aggregates from Yang et al (2000) which result in soft spheroids with a higher and constant density. The corresponding introduction now reads:

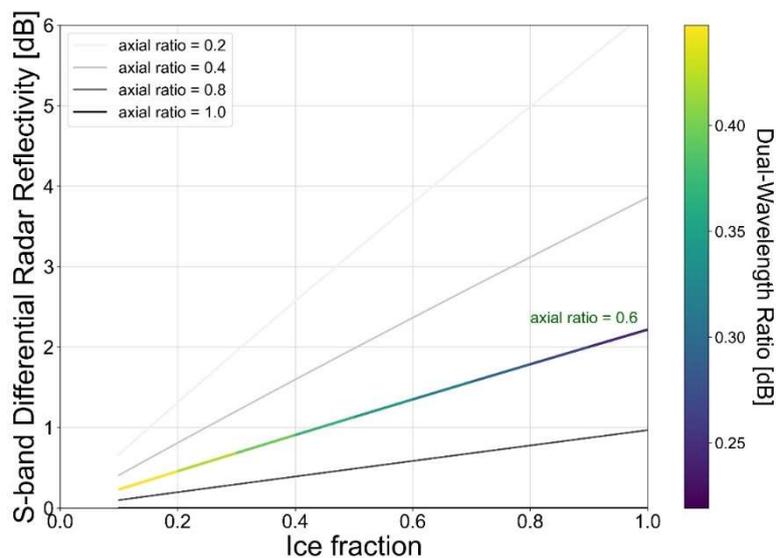
*"While the effective density of a spheroid decreases strongly with its size due to the exponent  $b=1.9$  in BF95, we contrast this with a second  $m(D_{max})$  with a higher and constant density. To that end we borrowed the  $m(D_{max})$  from the irregular aggregate model from Yang et al. (2000) to create soft spheroids with an analog mass-size ratio."*

Our choice to continue the study with this  $m(D)$  can also be understood when answering your next question:

"The BF95 mass-size relation was found to model to low ice particle densities ... BF95 prescribes near zero density values for large particles. This leads to very low simulated ZDR values with increasing particles size (Fig. 6)"

Hogan et al (2012) was able to reproduce observed Zdr values. Could you please explain what is the difference between their and your study?

Simulations using the BF95 mass-size relation and the typical  $AR = 1.67$  in current Fig. 8 (previously Fig. 6) show that BF95 cannot reproduce our radar dataset. Especially for large particles the ZDR signal is very low due to the low density. Although we can achieve higher ZDR values like in Hogan et al (2012) when spheroids with higher density are used, these spheroids are then so compact that they cannot simultaneously achieve large DWR values. With our soft spheroids we can reproduce ZDR simulations similar to Fig. 5 from Hogan et al., 2012 for S-band and different axial ratio and ice fraction values. For the typical  $AR = 1.67$  (axial ratio 0.6), the maximum DWR for  $D_{max} = 1$  mm was, however, not larger than 0.45 dB when ice fraction was 0.2.



To conclude, we further used the aggregates mass-size relation after a radiative closure study from Ewald et al. (2021) where it was shown that aggregate habit reconciled better the lidar, radar, and solar radiance measurements against columns or plates. We also noticed that the aggregates could better explain our ZDR-DWR measurement space than the well-known BF95. Hence, we continued our ice studies using the aggregates mass-size relation from Yang et al. (2000). In this way we managed to obtain simultaneously larger ZDR and DWR that better match our radar measurements.

Fig. 16 page 20. I think this is the most interesting plot of the paper. I am not sure how practical this is, but the difference between a and b panels indicates that there is extra information that can be retrieved by using different measurement geometries. Very interesting.

Thank you for your interest in this plot (current Fig. 18). Our intention was indeed to show the different radar geometries effect on the retrieved parameters. Moreover, the valuable contribution of ZDR in different areas of the cloud cross-section is highlighted in this plot. For very low radar elevation angles from both radars, e.g., at the lowest part of the cloud cross-section, and also above MIRA-35, i.e., when POLDIRAD is scanning and MIRA-35 is pointing to zenith, ZDR can not only provide apparent shape information but also helps to narrow down the solution space for the size retrieval.

#### References

Draine, B. T. and Flatau, P. J.: Discrete-Dipole Approximation For Scattering Calculations, *J. Opt. Soc. Am. A*, *JOSAA*, 11, 1491–1499, <https://doi.org/10.1364/JOSAA.11.001491>, 1994.

Ewald, F., Groß, S., Hagen, M., Hirsch, L., Delanoë, J., and Bauer-Pfundstein, M.: Calibration of a 35 GHz airborne cloud radar: lessons learned and intercomparisons with 94 GHz cloud radars, *Atmos. Meas. Tech.*, 12, 1815–1839, <https://doi.org/10.5194/amt-12-1815-2019>, 2019.

Ewald, F., Groß, S., Wirth, M., Delanoë, J., Fox, S., and Mayer, B.: Why we need radar, lidar, and solar radiance observations to constrain ice cloud microphysics, *Atmos. Meas. Tech.*, 14, 5029–5047, <https://doi.org/10.5194/amt-14-5029-2021>, 2021.

Garnett, J. C. M. and Larmor, J.: XII. Colours in metal glasses and in metallic films, *Philosophical Transactions of the Royal Society of London. Series A, Containing Papers of a Mathematical or Physical Character*, 203, 385–420, <https://doi.org/10.1098/rsta.1904.0024>, 1904.

Gorgucci, E., Sarchilli, G., and Chandrasekar, V.: A procedure to calibrate multiparameter weather radar using properties of the rain medium, *IEEE Trans. Geosci. Remote Sens.*, 37, 269–276, <https://doi.org/10.1109/36.739161>, 1999.

Harrington, J. Y., Sulia, K., and Morrison, H.: A Method for Adaptive Habit Prediction in Bulk Microphysical Models. Part I: Theoretical Development, *J. Atmos. Sci.*, 70, 349–364, <https://doi.org/10.1175/JAS-D-12-040.1>, 2013.

Hogan, R. J. and Westbrook, C. D.: Equation for the Microwave Backscatter Cross Section of Aggregate Snowflakes Using the Self-Similar Rayleigh–Gans Approximation, *J. Atmos. Sci.*, 71, 3292–3301, <https://doi.org/10.1175/JAS-D-13-0347.1>, 2014.

Hogan, R. J., Tian, L., Brown, P. R. A., Westbrook, C. D., Heymsfield, A. J., and Eastment, J. D.: Radar Scattering from Ice Aggregates Using the Horizontally Aligned Oblate Spheroid Approximation, *J. Appl. Meteorol. Clim.*, 51, 655–671, <https://doi.org/10.1175/JAMC-D-11-074.1>, 2012.

Hogan, R. J., Honeyager, R., Tyynelä, J., and Kneifel, S.: Calculating the millimetre-wave scattering phase function of snowflakes using the self-similar Rayleigh–Gans Approximation, *Q. J. R. Meteorol. Soc.*, 143, 834–844, <https://doi.org/10.1002/qj.2968>, 2017.

Jiang, Z., Verlinde, J., Clothiaux, E. E., Aydin, K., and Schmitt, C.: Shapes and Fall Orientations of Ice Particle Aggregates, *J. Atmos. Sci.*, 76, 1903–1916, <https://doi.org/10.1175/JAS-D-18-0251.1>, 2019.

Leinonen, J., Kneifel, S., Moisseev, D., Tyynelä, J., Tanelli, S., and Nousiainen, T.: Evidence of nonspheroidal behavior in millimeter-wavelength radar observations of snowfall, *J. Geophys. Res.: Atmos.*, 117, <https://doi.org/10.1029/2012JD017680>, 2012.

Leinonen, J., Kneifel, S., and Hogan, R. J.: Evaluation of the Rayleigh–Gans approximation for microwave scattering by rimed snowflakes, *Q. J. R. Meteorol. Soc.*, 144, 77–88, <https://doi.org/10.1002/qj.3093>, 2018.

Liao, L., Meneghini, R., Tokay, A., and Bliven, L. F.: Retrieval of Snow Properties for Ku- and Ka-Band Dual-Frequency Radar, *J. Appl. Meteorol. Clim.*, 55, 1845–1858, <https://doi.org/10.1175/JAMC-D-15-0355.1>, 2016.

Liou, K.-N.: Influence of Cirrus Clouds on Weather and Climate Processes: A Global Perspective, *Mon. Weather Rev.*, 114, 1167–1199, [https://doi.org/10.1175/1520-0493\(1986\)114<1167:IOCCOW>2.0.CO;2](https://doi.org/10.1175/1520-0493(1986)114<1167:IOCCOW>2.0.CO;2), 1986.

Liu, C., Lee Panetta, R., and Yang, P.: Inhomogeneity structure and the applicability of effective medium approximations in calculating light scattering by inhomogeneous particles, *J. Quant. Spectrosc. Radiat. Transfer*, 146, 331–348, <https://doi.org/10.1016/j.jqsrt.2014.03.018>, 2014.

Matrosov, S. Y.: Evaluations of the Spheroidal Particle Model for Describing Cloud Radar Depolarization Ratios of Ice Hydrometeors, *J. Atmos. Oceanic Technol.*, 32, 865–879, <https://doi.org/10.1175/JTECH-D-14-00115.1>, 2015.

Matrosov, S. Y.: Ice Hydrometeor Shape Estimations Using Polarimetric Operational and Research Radar Measurements, 11, 97, <https://doi.org/10.3390/atmos11010097>, 2020.

Melnikov, V.: Parameters of Cloud Ice Particles Retrieved from Radar Data, *J. Atmos. Oceanic Technol.*, 34, 717–728, <https://doi.org/10.1175/JTECH-D-16-0123.1>, 2017.

Mishchenko, M. I., Dlugach, J. M., and Liu, L.: Applicability of the effective-medium approximation to heterogeneous aerosol particles, *J. Quant. Spectrosc. Radiat. Transfer*, 178, 284–294, <https://doi.org/10.1016/j.jqsrt.2015.12.028>, 2016.

Reimann, J.: On Fast, Polarimetric Non-Reciprocal Calibration and Multipolarization Measurements on Weather Radars, 2014.

Reimann, J. and Hagen, M.: Antenna Pattern Measurements of Weather Radars Using the Sun and a Point Source, *J. Atmos. Oceanic Technol.*, 33, 891–898, <https://doi.org/10.1175/JTECH-D-15-0185.1>, 2016.

Schrom, R. S. and Kumjian, M. R.: Bulk-Density Representations of Branched Planar Ice Crystals: Errors in the Polarimetric Radar Variables, *J. Appl. Meteorol. Clim.*, 57, 333–346, <https://doi.org/10.1175/JAMC-D-17-0114.1>, 2018.

Tyynelä, J., Leinonen, J., Moisseev, D., and Nousiainen, T.: Radar Backscattering from Snowflakes: Comparison of Fractal, Aggregate, and Soft Spheroid Models, *J. Atmos. Oceanic Technol.*, 28, 1365–1372, <https://doi.org/10.1175/JTECH-D-11-00004.1>, 2011.

Waterman, P. C.: Matrix formulation of electromagnetic scattering, *Proc. IEEE*, 53, 805–812, <https://doi.org/10.1109/PROC.1965.4058>, 1965.

Yang, P., Liou, K. N., Wyser, K., and Mitchell, D.: Parameterization of the scattering and absorption properties of individual ice crystals, *J. Geophys. Res.: Atmos.*, 105, 4699–4718, <https://doi.org/10.1029/1999JD900755>, 2000.

Zhang, Y., Li, Z., and Macke, A.: Retrieval of Surface Solar Radiation Budget under Ice Cloud Sky: Uncertainty Analysis and Parameterization, *J. Atmos. Sci.*, 59, 2951–2965, [https://doi.org/10.1175/1520-0469\(2002\)059<2951:ROSSRB>2.0.CO;2](https://doi.org/10.1175/1520-0469(2002)059<2951:ROSSRB>2.0.CO;2), 2002.