

Testing the efficacy of atmospheric boundary layer height detection algorithms using uncrewed aircraft system data from MOSAiC

5 Gina Jozef^{1,2,3}, John Cassano^{1,2,3}, Sandro Dahlke⁴, Gijs de Boer^{2,5,6}

¹Dept. of Atmospheric and Oceanic Sciences, University of Colorado Boulder, Boulder, CO, USA

²Cooperative Institute for Research in Environmental Sciences, University of Colorado Boulder, Boulder, CO, USA

³National Snow and Ice Data Center, University of Colorado Boulder, Boulder, CO, USA

⁴Alfred Wegener Institute Helmholtz Centre for Polar and Marine Research, Potsdam, Germany

10 ⁵NOAA Physical Sciences Laboratory, Boulder, CO, USA

⁶Integrated Remote and In Situ Sensing, University of Colorado Boulder, Boulder, CO, USA

Correspondence to: Gina Jozef (gina.jozef@colorado.edu)

Abstract. During the Multidisciplinary drifting Observatory for the Study of Arctic Climate (MOSAiC) expedition, meteorological conditions over the lowest 1 km of the atmosphere were sampled with the DataHawk2 (DH2) fixed-wing uncrewed aircraft system (UAS). These in situ observations of the central Arctic atmosphere are some of the most extensive to date and provide unique insight into the atmospheric boundary layer (ABL) structure. The ABL is an important component of the Arctic climate, as it can be closely coupled to cloud properties, surface fluxes, and the atmospheric radiation budget. The high temporal resolution of the UAS observations allows us to manually identify the ABL height (Z_{ABL}) for 65 out of the total 89 flights conducted over the central Arctic Ocean between 23 March and 26 July 2020 by visually analyzing profiles of virtual potential temperature, humidity, and bulk Richardson number. Comparing this subjective Z_{ABL} with Z_{ABL} identified by various previously published automated objective methods allows us to determine which objective methods are most successful at accurately identifying Z_{ABL} in the central Arctic environment. The objective methods we use are the Liu-Liang, Heffter, virtual potential temperature gradient maximum, and bulk Richardson number methods. In the process of testing these objective methods on the DH2 data, numerical thresholds were adapted to work best for the UAS-based sampling. To determine if conclusions are robust across different measurement platforms, the subjective and objective Z_{ABL} determination processes were repeated using the radiosonde profile closest in time to each DH2 flight. For both the DH2 and radiosonde data, it is determined that the bulk Richardson number method is the most successful at identifying Z_{ABL} , while the Liu-Liang method is least successful. The results of this study are expected to be beneficial for upcoming observational and modeling efforts regarding the central Arctic ABL.

1 Introduction

The transfer of energy between the Earth's surface and the overlying atmosphere, particularly at high latitudes, remains an area of substantial uncertainty in our understanding of the global climate system (de Boer et al., 2012; Tjernström et al., 2012; Karlsson and Svensson, 2013). The consequences of this uncertainty are significant, with global climate model projections of present-day sea ice demonstrated to fall short of simulating the observed rate of change (Stroeve et al., 2007; Stroeve et al., 2012). The thermodynamic structure of the lower atmosphere plays a central role in

regulating cloud lifecycle and radiative transfer, and their influence on atmospheric energy transport (Tjernström et al., 2004; Karlsson and Svensson, 2013; Brooks et al., 2017). Significant insight can be gained by measurements collected over the central Arctic Ocean pack ice, focused on the structure of the lower atmosphere, its spatial and temporal variability, the intensity of turbulent energy fluxes, and its connection to surface features. To provide such measurements, uncrewed aircraft were deployed in the lower atmosphere during legs 3 (March through May 2020) and 4 (June through August 2020) of MOSAiC (Multidisciplinary drifting Observatory for the Study of Arctic Climate; Shupe et al. 2020), a year-long expedition that took place from October 2019 to September 2020 in which the icebreaker RV *Polarstern* (Alfred-Wegener-Institut Helmholtz-Zentrum für Polar- und Meeresforschung, 2017) was frozen into the central Arctic ocean sea ice pack and allowed to passively drift across the central Arctic for an entire year (Fig. 1). Additional information on measurements taken of the atmosphere and sea ice during MOSAiC can be found at Shupe et al. (2022) and Nicolaus et al. (2022) respectively.

One important indicator of the extent to which energy may be transferred between the Earth's surface and overlying atmosphere is the atmospheric boundary layer (ABL) height. The ABL is the turbulent lowest part of the atmosphere that is directly influenced by the Earth's surface (Stull, 1988; Marsik et al., 1995). In the central Arctic, the ABL is impacted by interactions between the atmosphere and sea ice pack features, which can cause either buoyantly or mechanically produced turbulence. The generation of buoyant turbulence can occur through surface energy fluxes emitted from open water regions such as leads (Lüpkes et al., 2008), cold air advection (Vihma et al., 2005), or turbulent mixing below cloud base due to cloud top radiative cooling (Tjernström et al., 2004). Mechanical generation, which is the dominant driver of turbulence in the central Arctic (Brooks et al., 2017), can occur due to the presence of a low-level jet (Brooks et al., 2017; Banta, 2003), sea ice features such as ridges and ice edges (Andreas et al., 2010) or oceanic waves (Jenkins et al., 2012). Solar heating of the Earth's surface and the subsequent formation of buoyant thermals, which is a dominant forcing of the ABL in most parts of the planet (Marsik et al., 1995), does not often play a role in the central Arctic due to the relatively reflective surfaces found there.

The Arctic ABL is usually either stable or near-neutral (Brooks et al., 2017; Esau and Sorokina, 2009). A stable boundary layer forms when there is a deficit of radiation at the surface or when warmer air is advected over a cooler surface, and can range from being nearly well-mixed with moderate turbulence to nearly laminar (Stull, 1988). A neutral boundary layer occurs when air at the surface is neutrally buoyant (Sivaraman et al., 2013) due primarily to mechanically generated turbulence which mixes air between the surface and above atmosphere (Brooks et al., 2017). A convective boundary layer forms when convective thermals create positive buoyancy (Liu and Liang, 2010) and an air parcel at the surface rises adiabatically until becoming neutrally buoyant; this phenomenon is rare in the central Arctic. While the various forms that the Arctic ABL may take are complex, most of the time, the Arctic ABL is capped by a potential temperature inversion, marking the entrainment zone, which is a stable layer between the ABL and free atmosphere (Stull, 1988). One important difference between the Arctic ABL and that in the mid-latitudes is that there is usually no residual layer above a stable Arctic ABL, due to the lack of a diurnal cycle. Additionally, the Arctic ABL is typically much shallower than that at mid-latitudes (Esau and Sorokina, 2009). These discrepancies cause certain ABL height detection methods to fail when applied to Arctic data.

75 Knowing the height of the Arctic ABL is important for many applications. First, it is a metric which represents the
altitude up to which the atmosphere is directly impacted by surface processes. This can then inform the extent to which
the surface interacts with atmospheric features such as clouds (and their influence on radiative transfer in the lower
atmosphere), low level jets (LLJs), and temperature inversion layers, which all have important implications for Arctic
warming (Serreze and Barry, 2011). For example, a shallow, stable ABL is more likely to be observed with clear skies
above (Brooks et al., 2017), which promotes longwave cooling of the surface and decoupling from the above
80 atmosphere. In this instance, a surface-based temperature inversion is likely to constrain warming to the surface, which
contributes to Arctic amplification (Lesins et al., 2012). ABL height (hereafter Z_{ABL}) plays an important role in many
other applications including transfer of air pollutants and weather forecasting (Garratt, 1994). Since any determination
of Z_{ABL} is simply an approximation, the most value can be gained if this approximation is as accurate as possible. The
goal of the current work is to determine which methods can best accomplish this.

85 The depth of the ABL has been previously defined using a variety of approaches that involve visualizing the profiles
of different variables. Table 1 lists some of the thermodynamic and kinematic variables that have previously been
used to identify Z_{ABL} , as well as some examples of associated literature that references use of that variable. Each of
these profiles typically exhibits a distinct change in structure at the top of the ABL. Additional methods may exist,
such as analyzing the vertical gradient of aerosol content, but are not listed since the current study focuses on Z_{ABL}
90 determination using thermodynamic and kinematic processes.

Table 1: List of quantities previously used to identify Z_{ABL} , as well as some associated literature in which each
variable is referenced.

Quantity Used	Previous Literature
virtual potential temperature	Heffter, 1980; Stull, 1988; Seibert et al., 2000; Pesenson, 2003; Dai et al., 2011; Sivaraman et al., 2013; Dai et al., 2014; Zhang et al., 2014
vertical gradient of virtual potential temperature	Heffter, 1980; Stull, 1988; Liu and Liang, 2010; Dai et al., 2011; Sivaraman et al., 2013; Dai et al., 2014
vertical gradient of temperature	Stull, 1988; Collaud Coen et al., 2014
bulk Richardson number	Stull, 1988; Seibert et al., 2000; Zilitinkevich and Baklanov, 2002; Steeneveld et al., 2007; Dai et al., 2011; Sivaraman et al., 2013; Dai et al., 2014; Zhang et al., 2014; Collaud Coen et al., 2014
total wind speed	Stull, 1988; Seibert et al., 2000; Liu and Liang, 2010; Steeneveld et al., 2007; Sivaraman et al., 2013; Zhang et al., 2014
wind shear	Dai et al., 2011; Dai et al., 2014; Zhang et al., 2014
liquid water content and absolute humidity	Seibert et al., 2000; Pesenson, 2003; Dai et al., 2014
turbulent kinetic energy	Stull, 1988; Dai et al., 2014; Zhang et al., 2014

95 Due to the different atmospheric dynamics involved in each of the above approaches, the definition of Z_{ABL} is often
debatable amongst experts. Depending on one's purpose for knowing Z_{ABL} , different approaches may be most
applicable. Of these methods, some of the most widely used ones, and the ones applied in the current analysis of a
central Arctic dataset to determine Z_{ABL} , are the ones that involve analysis of virtual potential temperature (θ_v), vertical

100 gradient of virtual potential temperature ($d\theta_v/dz$), humidity (relative and absolute), bulk Richardson number (Ri_b), and
wind speed profiles. The current focus is on these variables because the physical basis for each one as an indication
of Z_{ABL} is relevant for the Arctic atmosphere. Specifically, θ_v helps identify the entrainment zone above the ABL, the
vertical gradient of humidity either decreases or increases noticeably above the ABL (Dai et al., 2014), Ri_b helps
105 identify where turbulence (usually caused by strong wind shear or surface roughness in the Arctic ABL (Grachev et
al., 2005)) ceases above the ABL, and wind speed helps identify the top of the ABL when it is capped by an LLJ as
the ABL top is often at or just below the LLJ core (Stull, 1988). Other methods, such as that using temperature
inversion top to identify Z_{ABL} (Collaud Coen et al., 2014), do not perform well in the Arctic where a weak temperature
110 inversion can extend well above the ABL. Though turbulent kinetic energy is recognized as perhaps the most valuable
profile for Z_{ABL} identification (Stull, 1988; Siebert et al., 2000; Dai et al., 2014; Zhang et al., 2014), these data are not
available to aid in the current study.

High resolution data collected by the DataHawk2 uncrewed aircraft system (UAS) allows for determination of Z_{ABL}
115 with high accuracy through manual visual analysis. However, visually determining Z_{ABL} case-by-case is time
consuming for processing a large dataset. Therefore, the UAS-derived dataset is leveraged to compare manually (or
'subjectively') determined Z_{ABL} with that identified through previously published automated (or 'objective') methods.
While this subjective Z_{ABL} may not necessarily be the 'true' ABL top, as the definition of this quantity can be debatable
among experts and is not constant over time, it is the best estimate of Z_{ABL} given the available data. This evaluation is
120 completed to identify objective methods that can accurately diagnose Z_{ABL} across a larger dataset of central Arctic
atmospheric conditions.

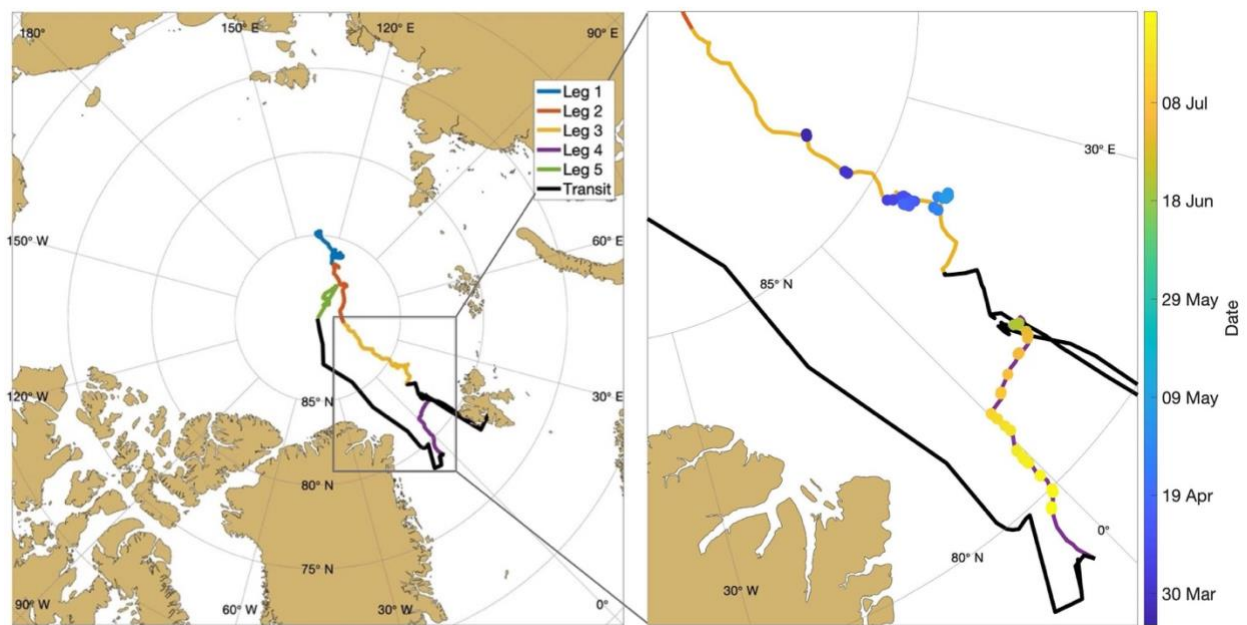
To subjectively identify Z_{ABL} in each atmospheric profile from DH2 data, the stability regime of the ABL (stable,
neutral, or convective) is categorized and Z_{ABL} is visually identified through combined evaluation of θ_v , humidity
(both relative humidity (RH) and mixing ratio), and Ri_b profiles. Objective identification of Z_{ABL} is derived through
125 the application of four previously published methods: the Liu-Liang method (Liu and Liang, 2010), the Heffter method
(Heffter, 1980), the virtual potential temperature gradient maximum (TGRDM) method (Dai et al., 2014), and the Ri_b
method (Sivaraman et al., 2013), adapted to best suit the DH2 profiles examined. Then, statistical comparisons
between the objective and subjective Z_{ABL} are conducted. Next, the objective methods are applied in their adapted
form to radiosonde profiles nearest in time to each DH2 flight to determine if these methods are robust across different
130 measurement platforms for central Arctic conditions. Finally, discussion is included on the features that do or do not
lend themselves to accurate identification of Z_{ABL} by the objective methods, and findings are summarized to support
future studies seeking to identify Z_{ABL} quickly, objectively, and accurately across large atmospheric datasets collected
in the central Arctic.

2 Data and methods

130 2.1 The DataHawk2

Data presented in this study were obtained between 23 March and 26 July 2020 using the University of Colorado DataHawk2 (DH2) UAS (de Boer et al. submitted). Flights were conducted from the sea ice alongside the *Polarstern*, known as the MOSAiC floe, ranging in location from 86.2° N, 15.8° E on 23 March, to 79.8° N, 1.9° W on 26 July 2020 (Fig. 1). Throughout this period, the MOSAiC floe evolved from snow-covered rigid ice situated in the high Arctic to being covered with melt ponds and leads close to the sea ice edge. The surface atmospheric temperatures also transitioned from nearly -35 °C at the beginning of leg 3 to hovering near 0 °C throughout the entirety of leg 4.

135



140

Figure 1: (Left) The drift track of the *Polarstern*, separated by color into the 5 different legs. The black “transit” line indicates when the ship was travelling under its own power between legs 3 and 4 and between legs 4 and 5. **(Right)** The zoomed in portion of the *Polarstern* drift during which DH2 flights were conducted (legs 3 and 4). The locations of all of the DH2 flights are overlaid on the drift track and color coded by date, with blue-tinted dots indicating flights conducted during leg 3 and yellow-tinted dots indicating flights conducted during leg 4.

145

The DH2 (Hamilton et al., submitted) is a fixed-wing, battery powered UAS (1.1 m wingspan, 1.8 kg weight, 40 min endurance) carrying various meteorological sensors, which measure the state of the atmosphere in Earth-relative coordinates. Instrumentation includes a fine wire array providing high frequency (800 Hz) information on temperature and air speed, multiple sensors for temperature and relative humidity (Vaisala RSS421 measuring at 5 Hz and SHT-85 measuring at 100 Hz), and up- and downward looking thermopile sensors to provide infrared brightness temperatures of the sky and surface. Air pressure is measured at 5 Hz by the Vaisala RSS421 sensor. Measurements of attitude, airspeed, ground speed, and altitude support the derivation of high-frequency (10 Hz) 3D wind estimates, which are time-lag-corrected to account for differences in time constants of the GPS system and onboard inertial measurement unit (IMU), true airspeed (TAS)-corrected to eliminate errors in the measurement of TAS, and corrected

150

for potential angular offsets between the aircraft, IMU, and pitot probe. Additional detail on the derivation of these winds, as well as processing methods of the additional variables are described in detail in the metadata for the DH2 dataset used for the current study (Jozef et al., 2021), the manuscript describing UAS data collection during MOSAiC (de Boer et al., submitted) and the technical manuscript describing the DH2 aircraft (Hamilton et al., submitted). Despite these corrections, near-surface wind measurements (below ~30 m) are usually unreliable due to manual, rather than autopilot, takeoffs and landings, which conflict with the measurements and calculations of the winds. Therefore, we do not use DH2 wind speeds below 30 m in the current study.

Combined, these sensors provide a comprehensive picture of atmospheric thermodynamic and kinematic state along with some context on the surface and sky condition under which these measurements were obtained. Table 2 lists the resolution, repeatability (standard deviation of difference between two successive repeated calibrations), and response time for the Vaisala RSS421 sensor. Uncertainty in the wind speed estimation is not provided, as determining this is still in progress.

Table 2: Accuracy and reliability of the variables recorded by the Vaisala RSS421 sensors used in this study.

Variable	Resolution	Repeatability	Response Time
Pressure	0.01 hPa	0.4 hPa	-
Temperature	0.01 °C	0.1 °C	0.5 s
Humidity	0.1 %RH	2 %RH	<0.3 s (at 20 °C) to <10 s (at -40 °C)

Measurements collected by the DH2 are logged at different frequencies, requiring the implementation of a time alignment process to assure that the time index for each datapoint of each variable is consistent with all other measurements. Additionally, the wind measurements have been filtered to remove the impact of the angle and ground speed of the aircraft to provide estimates of the true wind speed (de Boer et al., submitted). Data collected by the DH2 during MOSAiC are available for public download through the National Science Foundation Arctic Data Center at <https://doi.org/10.18739/A2Z60C34R> (Jozef et al., 2021).

During MOSAiC, DH2 flights were conducted whenever flight weather criteria were met and when the team was able to access the ice alongside the *Polarstern*. The weather criteria include wind speeds with a sustained average below 10 m s⁻¹, and gusts below 14 m s⁻¹, as well as sufficient visibility to maintain visual contact with the aircraft at all times during flight. In addition, DH2 flights required coordination with other MOSAiC activities, especially those impacting air space over the MOSAiC floe, including manned helicopter flights and other UAS and tethered operations.

The most common flight pattern conducted with the DH2, and the flight pattern from which data for this analysis were acquired, was a profiling flight in which the plane flew a spiral ascent and descent pattern, with a radius of 75-100 m between the surface and 1 km altitude (or cloud base, if lower than 1 km), with the aircraft ascending and descending at a rate of 2 m s⁻¹ and flying at an airspeed of 14-18 m s⁻¹. Each profiling flight lasted an average of 30 min, with some shorter flights when the air temperature was at its coldest (~-35 °C) near the beginning of leg 3, and some longer

185 flights when the air temperature was much warmer (~ 0 °C) during leg 4. Throughout the measurement period, 89 flights were conducted with the DH2. In the present study, 65 of these flights are found to have a clearly identifiable Z_{ABL} within the altitude range sampled. The remaining flights sampled only the lowest portion of the atmosphere due to cloud cover or other environmental conditions and therefore did not observe the full depth of the ABL.

2.1.1 Preparing the DataHawk2 data for analysis

190 The primary profiles of interest for subjective and objective Z_{ABL} identification are θ_v , humidity (RH and mixing ratio), wind speed, Ri_b , and $d\theta_v/dz$. θ_v was calculated using RSS421 temperature, pressure, and RH. Differences in response times of the RSS421 temperature and RH sensors has a negligible impact on the calculation of θ_v because the moisture content in the Arctic atmosphere is so low that θ and θ_v values typically differ on the order of less than 1 K. Regardless, the addition of humidity does not change the structure and location of features in the θ_v profile, which is what is important for Z_{ABL} identification. For proper visualization, we average the θ_v , humidity, and wind speed variables over 1 m altitude bins throughout the entire flight to further eliminate the effects of differences in sensor response times during ascent and descent (e.g., values at 10.5 m are averaged from 10 to 11 m). This also mitigates the effect of changes in atmospheric conditions near the surface throughout the span of a flight, though the near-surface observations largely remained constant during a given flight. 1 m is chosen as an averaging bin because using a greater bin value would eliminate much of the fine scale detail in the θ_v and humidity profiles which the DH2 provides, and which makes its data a valuable resource in honing Z_{ABL} detection methods. However, since fine scale fluctuations in wind speeds evident at the 1 m scale are usually artifacts of the wind estimation routines applied to a circular flight pattern, we additionally apply a 60 m running mean, which eliminates small-scale wiggles while retaining the important large-scale features. Next, we exclude periods of manual flight during takeoff and landing (this is usually at altitudes below 5 m) since measurements during manual flight are prone to inaccuracies due to the irregular flight pattern. Lastly, we exclude the first 5 seconds of flight, as the initial measurements after takeoff may be faulty due to hysteresis associated with the sensor sitting still at the surface before launch.

200 Using the 1 m averaged θ_v and wind speed component profiles, we calculate the Ri_b profile. Ri_b is calculated at altitude, z , using the following equation from Stull (1988):

$$205 Ri_b(z) = \frac{\left(\frac{g}{\overline{\theta_v}}\right)\Delta\theta_v \Delta z}{\Delta u^2 + \Delta v^2} \quad (1)$$

210 where g is acceleration due to gravity, $\overline{\theta_v}$ is mean virtual potential temperature over the altitude range being considered, z is altitude, u is zonal wind, v is meridional wind, and Δ represents the difference over the altitude range used to calculate Ri_b throughout the profile. The only way that Ri_b can be negative is if the value for $\Delta\theta_v$ is negative, indicating a convective atmosphere with buoyancy-driven generation of the turbulence. Ri_b profiles are created by calculating Ri_b over a 30 m altitude range (Δ), at 5 m resolution (i.e., between 30 and 60 m, then between 35 and 65 m, and so on). Since we do not use DH2 winds below 30 m, and intermediate Ri_b value between the surface and 30 m is calculated using an assumed zero wind at the surface. It is not crucial to consider the drift speed of the ice for the

calculation of this initial Ri_b value since the ice drift speed during MOSAiC was on average less than 0.1 m s^{-1} (Krumpen et al., 2021), a value much smaller than that of the observed wind speeds. Lastly, the $d\theta_v/dz$ profile is similarly created by calculating $d\theta_v/dz$ over an altitude range of 30 m, at 5 m resolution.

220 The above profiles are used to determine stability regime, visually identify Z_{ABL} using criteria founded in this manuscript, and objectively identify Z_{ABL} using the four published methods. For the remainder of this manuscript, Z_{ABL} determined from manual visual identification is referred to as the ‘subjective’ Z_{ABL} and that determined by the published methods (which are automated algorithms performed by computers) are referred to as ‘objective’ Z_{ABL} . These terms are used as a simplification to differentiate between manual and automated methods, though they both rely on much of the same underlying physical processes for Z_{ABL} identification.

225 2.2 Determining stability regime

Some of the methods for both subjectively and objectively identifying Z_{ABL} differ depending on the stability regime, so the sampled regime is first identified for each DH2 flight. The three possible stability regimes considered include a convective boundary layer (CBL), stable boundary layer (SBL), and neutral boundary layer (NBL; Liu and Liang, 2010). In a CBL, θ_v near the surface is greater than that of the overlying ABL (Stull, 1988). In an SBL, the vertical
230 gradient of θ_v is positive (Stull, 1988). In an NBL, θ_v at the surface is approximately the same value as that of the overlying remainder of the ABL (Stull, 1988).

Therefore, stability regimes are identified by comparing θ_v between the lowest altitude sampled by the DH2 (‘ i ’ in the below equations; typically $\sim 5\text{m}$ since altitudes below this are usually sampled with manual flight) and 40 m above, using Eq. (2)-(4) below adapted from Liu and Liang (2010).

$$235 \quad \theta_{v_{i+40m}} - \theta_{v_i} < -\delta_s = \text{CBL} \quad (2)$$

$$\theta_{v_{i+40m}} - \theta_{v_i} > +\delta_s = \text{SBL} \quad (3)$$

$$-\delta_s \leq \theta_{v_{i+40m}} - \theta_{v_i} \leq +\delta_s = \text{NBL} \quad (4)$$

In these equations, δ_s is a stability threshold that represents the minimum positive or negative vertical gradient of θ_v near the surface necessary for the ABL to qualify as an SBL or CBL respectively. If this minimum is not either
240 negatively (in the case of a CBL) or positively (in the case of an SBL) reached, the ABL is identified as an NBL (Liu and Liang, 2010). In an idealized case, δ_s would be zero. However, in practice it must be specified as a small positive number, and this number depends on the surface characteristics as well as inherent uncertainties or noise in the measurements. For profiles over ocean/ice, this number has been defined to be 0.2 K (Liu and Liang, 2010).

245 While Liu and Liang (2010) compare θ_v between pressure levels that equate to approximately 40 and 160 m in the conditions we sampled, this range was found to be inadequate for differentiating between an SBL, NBL or CBL in the Arctic, where the top of the ABL is often below 160 m, and sometimes even below 40 m. Therefore, considering the

θ_v change below ~45 m more accurately reflects the stability regime of the Arctic ABL. Once the stability regime is identified, criteria based on the θ_v , humidity, and Ri_b profiles are applied to subjectively determine Z_{ABL} . For the current dataset, 31 SBL cases, 32 NBL cases, and 2 CBL cases were identified.

250 2.3 Subjective identification of atmospheric boundary layer height

There is no one best method for subjectively identifying Z_{ABL} that is agreed upon throughout the scientific community, evident by the many methods outlined in Table 1, and therefore a subjectively determined Z_{ABL} is prone to error. The best we can do to increase the confidence in a subjectively determined Z_{ABL} is to take into account several of the most commonly used methods and establish criteria which are applied consistently across all profiles. We describe these
255 criteria below.

To subjectively identify Z_{ABL} , the θ_v profile is first analyzed, as the θ_v profile changes structure above the ABL (Stull, 1988). For a CBL and NBL, above the ABL, θ_v changes from decreasing or constant with height, to increasing with height, marking the entrainment zone (Stull, 1988). The structure of an SBL, however, can vary a lot more (Mayer et al., 2012; Steeneveld et al., 2007; Zilitinkevich and Baklanov, 2002). In an ideal SBL case, the θ_v inversion is at its
260 strongest (greatest vertical gradient of θ_v) near the surface and transitions to the free atmosphere (nearly constant or gradually increasing θ_v with altitude) above the SBL, with no entrainment zone (Stull, 1988). Z_{ABL} is then identified as the altitude of the shift from the surface-based θ_v inversion to the free atmosphere (Stull, 1988). In reality, the structure of an SBL is often not that simple, and the height of an SBL can be difficult to identify based on θ_v alone (Stull, 1988; Zhang et al., 2014). SBLs in the DH2 dataset often include a weaker surface-based θ_v inversion capped
265 by a layer of enhanced stability (stronger θ_v inversion), reminiscent of an entrainment zone, likely because of surface-drag induced turbulence close to the surface. ABLs with this structure form as the near-surface atmosphere fluctuates between weakly stable and near-neutral (Brooks et al., 2017). In more difficult cases such as these, the top of the SBL can be better determined by supplementing the θ_v profile with the RH and mixing ratio profiles, which usually have an obvious transition at the top of the ABL (Dai et al., 2014). This transition can manifest as either a shift from zero
270 or positive to negative vertical gradient of humidity, or as a humidity inversion. Use of the humidity profiles can also increase the confidence in identification of CBL and NBL height.

In addition, the Ri_b profile can aid in Z_{ABL} identification (Zhang et al., 2014). Ri_b is the ratio between buoyantly produced (from thermals) or suppressed (from static stability) turbulence, and mechanically produced turbulence (from wind shear; Sivaraman et al., 2013). Therefore, Ri_b can help to identify the top of the ABL under the assumption
275 that turbulence ceases above the ABL (Stull, 1988). In the limit of layer thickness becoming small, Ri_b can be compared to a critical value of ~0.25 (Stull, 1988), with Ri_b below the critical value indicating an atmosphere that is likely to become or remain turbulent, and Ri_b above the critical value indicating that an already laminar layer will not become turbulent, as static stability is strong enough to suppress mechanically generated turbulence. However, Ri_b does not assume a small layer thickness, so a critical value is not well defined for Ri_b . Thus, for Ri_b near the critical
280 value, there is uncertainty in the likelihood of turbulence (AMS Glossary of Meteorology). Different studies have found the appropriate critical Richardson number to range from as low as 0.15 to as high as 7.2 in coarse resolution

models (Dai et al., 2014), but across the board, lower Ri_b is expected in the ABL, and higher Ri_b is expected above the ABL (Seibert et al., 2000). This increase in Ri_b above the ABL is in large part due to the decrease in wind shear. By examining Ri_b profiles for the DH2 flights, this transition from low values (near zero) to high values (with an increase of a few digits above the lower altitude values) can aid in identifying the top of the ABL.

Table 3 below outlines the criteria applied to determine Z_{ABL} depending on stability regime, which are separated depending on how many kinks there are in the θ_v profile that might indicate the entrainment zone. The term ‘kink’ refers to a dramatic shift in slope (i.e., drastic change in vertical gradient). The primary methods applied to determine Z_{ABL} are those in which there are either one or two θ_v kinks, where we rely most heavily on the θ_v profile, and secondarily on the humidity profile. Only in few especially difficult cases were the Ri_b profiles used heavily.

Table 3: Subjective criteria for identifying Z_{ABL} , depending on stability regime.

	One θ_v kink	Multiple θ_v kinks	No clear θ_v kinks
Convective boundary layer (CBL)	<p>Z_{ABL} is the altitude at which the vertical gradient of θ_v is positive, and is the bottom of a layer of enhanced stability (greater vertical gradient of θ_v), corresponding to a kink in the humidity profiles and an increase in Ri_b.</p> <p>Example: Fig. 2a</p>		
Neutral boundary layer (NBL)	<p>Z_{ABL} is the altitude of the singular θ_v kink marking the bottom of the lowest θ_v inversion.</p> <p>Example: Fig. 2b</p>	<p>Z_{ABL} is the altitude of the θ_v kink near the bottom of the lowest θ_v inversion which corresponds to a kink in the humidity profiles and an increase in Ri_b.</p> <p>Example: Fig. 2c</p>	<p>Z_{ABL} is the altitude of a faint θ_v slope shift which is identified via a corresponding kink in the humidity profiles and increase in Ri_b.</p> <p>Example: Fig. 2d</p>
Stable boundary layer (SBL)	<p>Z_{ABL} is the altitude of the θ_v kink marking the bottom of a layer of enhanced stability (greater vertical gradient of θ_v), corresponding to a kink in the humidity profiles and sometimes an increase in Ri_b.</p> <p>Example: Fig. 2e</p>		<p>Z_{ABL} is the altitude of a faint θ_v slope shift which is identified via a corresponding kink in the humidity profiles and sometimes an increase in Ri_b.</p> <p>Example: Fig. 2f</p>

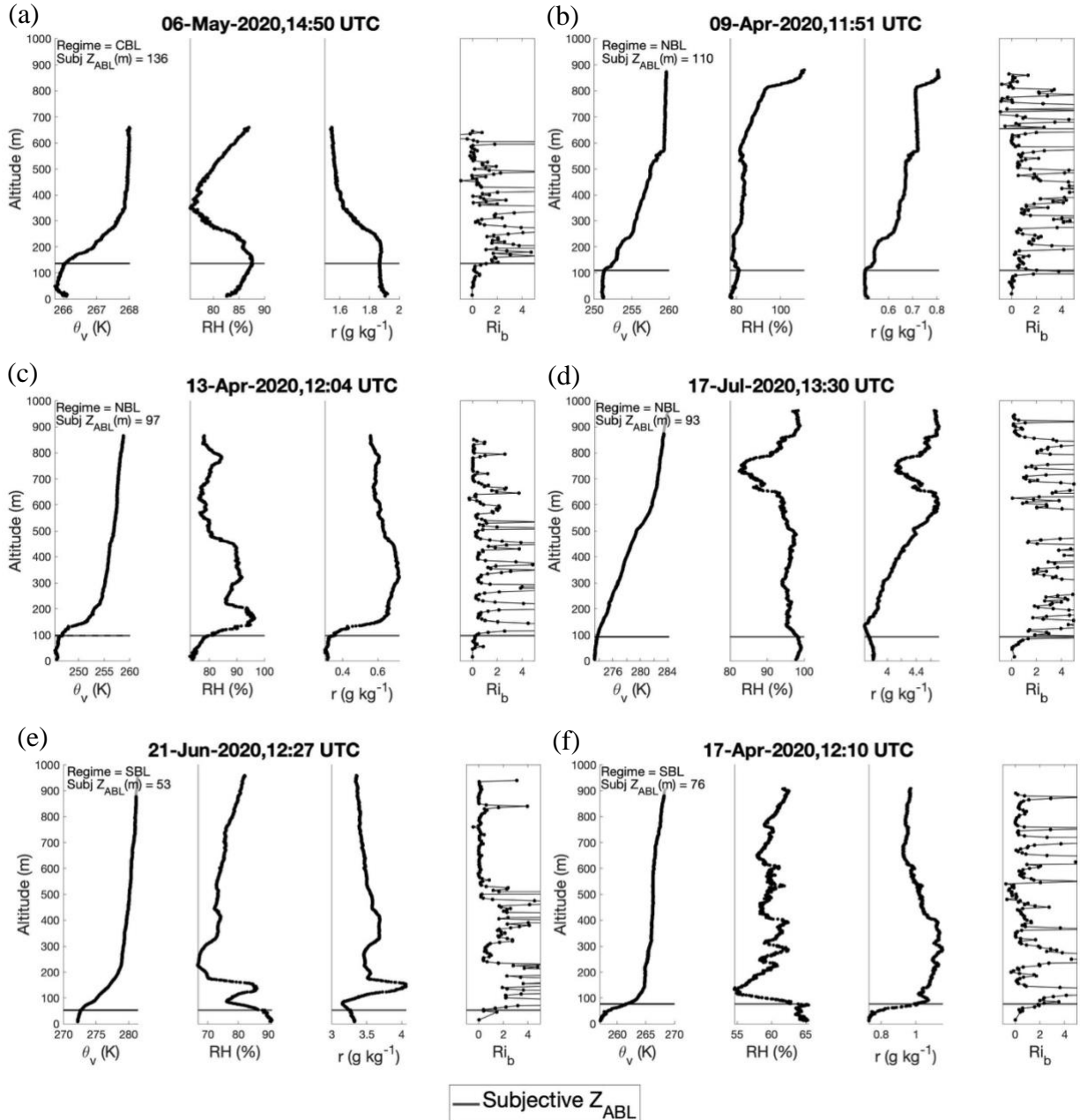


Figure 2: For each flight shown in the figure, the θ_v profile is plotted in the left panel, the RH and mixing ratio profiles are plotted in the middle two panels, and the Ri_b profile is plotted on the right panel. Subjective Z_{ABL} is marked with a horizontal black line on each panel, and is written, along with stability regime, on the left panel. **(a)** Example of a CBL case. **(b-d)** Examples of NBL cases. **(e-f)** Examples of SBL cases.

295

300

When applying the above criteria, Z_{ABL} for the majority of cases (about 90%) was clearly identifiable (i.e., relevant θ_v and humidity kinks were at the same altitude). For the other cases, Z_{ABL} was more ambiguous, meaning there were multiple features that one could argue marked the ABL top (i.e., the θ_v and humidity kinks which could both be interpreted as Z_{ABL} were at different altitudes). In these instances, depending on which feature is chosen, Z_{ABL} could differ by about 10-30 m. Therefore, in such cases, we determine the uncertainty in the subjective Z_{ABL} to be less than

30 m. Uncertainty in the height of a kink in an individual profile is only subject to the vertical averaging procedure and sensor response time, and thus is on the order of only ~1 m.

305 **2.4 Objective identification of atmospheric boundary layer height**

The strength of the subjective methods described above is the knowledge of the expert, which cannot be automated. However, such expert knowledge and the time necessary to individually assess profiles is not always available. Thus, an automated method may often be preferred. Four such methods for objectively determining Z_{ABL} are applied and evaluated. Each of these methods relies on profiles of either $d\theta_v/dz$ or Ri_b , some in combination with the θ_v and/or
310 wind speed profiles. Because the $d\theta_v/dz$ and Ri_b profiles are calculated over an altitude range of 30 m with 5 m resolution, objective Z_{ABL} detection methods which ultimately rely on these profiles can be determined with a resolution of 5 m. If they ultimately rely on the θ_v or wind speed profiles, Z_{ABL} can be determined with 1 m resolution. Figure 3 at the end of Sect. 2.4 shows the application of all objective methods for an SBL and NBL case. A CBL case is not shown, as there were only two CBLs identified in the DH2 profiles, and they are rare in the central Arctic.

315 **2.4.1 Liu-Liang method**

The application of the Liu-Liang method depends on whether the profile includes a CBL, SBL, or NBL, which is determined using Eq. (2)-(4). To implement the Liu-Liang method for a CBL profile, we first find the lowest altitude at which θ_v exceeds its the lowest DH2 value by 0.1 K. Then, Z_{ABL} is identified at the next lowest altitude in which
320 $d\theta_v/dz$ exceeds $0.05 \text{ K } 100 \text{ m}^{-1}$ (Liu and Liang, 2010). For an NBL, Z_{ABL} is identified as the altitude at which $d\theta_v/dz$ first exceeds $2.5 \text{ K } 100 \text{ m}^{-1}$, which is adapted from a threshold of $0.05 \text{ K } 100 \text{ m}^{-1}$ used in Liu and Liang (2010), as this threshold was found to be inappropriate for the current dataset (Z_{ABL} found with the original threshold was always far too low; this may be due to differences in the vertical resolution or smoothing of our data versus that used by Liu and Liang (2010)). The basis of this method is to identify the entrainment zone at the top of the ABL through an increased value of $d\theta_v/dz$.

325 For an SBL, the Liu-Liang method searches for a potential Z_{ABL} associated with either minimal turbulence due to the lack of buoyancy within the ABL, or greater turbulence in the ABL due to the presence of wind shear (Liu and Liang, 2010), both scenarios which may dictate Z_{ABL} for an SBL (Stull, 1988). Thus, SBL height is defined as either the top of the bulk stable (θ_v inversion) layer starting from the ground, or the height of the LLJ maximum if present, whichever is lower (Liu and Liang, 2010). The top of the bulk stable layer is identified where the surface-based θ_v inversion has
330 consistently diminished, and LLJ presence is identified by searching for wind speeds reaching a maximum that is at least 2 m s^{-1} stronger than the local minima above and below (Stull, 1988; Liu and Liang, 2010). For greater detail on these methods, and the guiding equations, see Liu and Liang (2010). Supplementary Figure S1 shows an example of the Liu-Liang method applied to a case for each stability regime.

2.4.2 Heffter method

335 The Heffter method uses θ_v difference across a θ_v inversion ($d\theta_v$) as an indication of Z_{ABL} (Sivaraman et al., 2013), by identifying the lowest θ_v inversion layer where $d\theta_v/dz$ is greater than $0.5 \text{ K } 100 \text{ m}^{-1}$ throughout the θ_v inversion, and $d\theta_v$ is at least 2 K (Heffter, 1980; Pesenson, 2003; Sivaraman et al., 2013). Within this θ_v inversion, the altitude at which θ_v first becomes more than 2 K greater than θ_v at the bottom of the θ_v inversion is labelled as Z_{ABL} (Marsik et al., 1995; Delle Monache et al., 2004; Snyder and Strawbridge, 2004; Sivaraman et al., 2013).

340 For a CBL or NBL, this method is meant to determine the altitude of the elevated θ_v inversion marking the entrainment zone between the well-mixed ABL and free atmosphere (Pesenson, 2003). For an SBL, this method determines where the change in strength of the surface θ_v inversion marks the transition from the ABL to residual layer (if one exists) or free atmosphere above (Stull, 1988). For greater detail on this method, and the guiding equations, see Heffter (1980) or Sivaraman et al. (2013). Supplementary Figure S2 shows an example of the Heffter method applied to a case for
345 each stability regime.

2.4.3 Virtual potential temperature gradient maximum (TGRDM) method

The final $d\theta_v/dz$ -based method used to find Z_{ABL} is the virtual potential temperature gradient maximum (TGRDM) method (Dai et al., 2014). Since the ABL is typically capped by a well-defined θ_v inversion layer (Stull, 1988), even in a weakly stable case, we expect to see a local maximum in the $d\theta_v/dz$ profile at this point. By finding the maximum
350 in the $d\theta_v/dz$ profile, the altitude at which the θ_v inversion is at its strongest and weakens above is identified. To apply this method, local maxima in the $d\theta_v/dz$ profile where $d\theta_v/dz$ is at least $1.75 \text{ K } 100 \text{ m}^{-1}$ greater than the local minimum $d\theta_v/dz$ above are identified. Z_{ABL} is set to the altitude of this lowest peak. Supplementary Figure S3 shows an example of the TGRDM method applied to a case for each stability regime.

2.4.4 Bulk Richardson number method

355 Finally, a bulk Richardson number method for finding the ABL top is applied by determining the altitude at which Ri_b exceeds a critical value, which indicates where turbulence was likely no longer able to form in a laminar atmosphere. Previous literature suggests a wide range of critical values with 0.25 (Stull, 1988) being the most widely accepted value, though a value of 0.5 is also often used (Sivaraman et al., 2013; Zhang et al., 2014). To determine a viable threshold value for the identifying Z_{ABL} in the DH2 data, a comparison between Z_{ABL} determined from a range of
360 threshold values (we used 0.25, 0.5, 0.75, 1.0, 1.25, and 1.5) and the subjective Z_{ABL} was conducted. In identifying Z_{ABL} from these different threshold values, the level above which Ri_b was consistently greater than the threshold value was found. For this dataset, four consecutive datapoints (20 m) were required to be above the threshold value. The altitude of the lowest of these four consecutive points is identified as Z_{ABL} . The threshold values deemed to identify Z_{ABL} closest to that identified by the subjective methods was 0.5 followed by 0.75. Therefore, further Z_{ABL} presented
365 using the Ri_b method is calculated with critical values of 0.5 (hereafter called $Ri_b(0.5)$) and 0.75 (hereafter called $Ri_b(0.75)$). Supplementary Figure S4 shows an example of the Ri_b method applied to a case for each stability regime.

2.5 Applying the objective methods to radiosonde profiles

As discussed above, some of the objective methods used in this study were modified from their original descriptions to better work with the Arctic UAS data. Primarily, this includes changing the altitude range for determining stability regime, adjusting the threshold for calculating Liu-Liang NBL height, adding the 1.75 K 100 m⁻¹ criterion to the TGRDM method, and choosing the best threshold values as well as specifying the necessary vertical distance for the Ri_b method. These adaptations are necessary in part because previous implementations involved analysis of radiosonde profiles, which have a lower resolution than the DH2 profiles, and in mid-latitude locations, where the ABL structure is often quite different than that observed in the Arctic (due to the lack of daytime convection or a diurnal cycle in the Arctic most of the time). Thus, profiles of θ_v , humidity, and wind speed from the balloon-borne radiosondes that were launched at least four times per day from the deck of the *Polarstern* (Maturilli et al., 2021) during MOSAiC are leveraged to determine if the objective methods used to identify Z_{ABL} from the UAS data are robust across platforms, despite differences in sampling methods.

To do this, radiosonde profiles with launch times closest to the DH2 flight times (within at most ~3 hours) are used, repeating the same processes for subjective and objective Z_{ABL} identification and comparison. In eight instances, there were two DH2 flights in closest time proximity to the same radiosonde launch, so we use data from a total of 57 different radiosonde profiles. The specs for the Vaisala RS41-SGP sensor, which recorded the radiosonde variables, are the same as those listed in Table 2 for the DH2's RSS421 sensor, with the addition of a wind uncertainty and resolution of 0.15 m s⁻¹ and 0.1 m s⁻¹ respectively for velocity, and of 2 ° and 0.1 ° respectively for direction. Before proceeding with analysis, profiles of temperature, wind, and humidity from the radiosondes were visually compared to those from the corresponding DH2 flight to confirm that the measurements were similar to each other.

Prior to applying the objective methods, data below 23 m altitude were removed, as the lowest part of the radiosonde profiles were found to show inaccurately warm temperatures for several cases (Maturilli et al., 2021), due to the *Polarstern* acting as a “heat island.” Additionally, in some cases, the radiosonde data showed anomalously warm measurements some distance above 23 m, which is assumed to be the result of the balloon passing through the *Polarstern*'s exhaust plume. These measurements were adjusted by interpolating the temperature between the closest good measurements above and below where the radiosonde was presumably in the ship's plume. Applying these adjustment means that radiosonde data near the surface are not available for determination of stability regime. Therefore, we adapt the methods applied to the DH2 data in Eq. (2)-(4) and instead calculate $d\theta_v$ between the lowest radiosonde measurement and 30 m above, or the subjective Z_{ABL} if lower. We then compare this $d\theta_v$ to the appropriate threshold value, δ_s , that is equal to $(0.2 \text{ K}/40 \text{ m} = 0.005 \text{ K m}^{-1}) \times$ the altitude range used. For example, if the 30 m altitude range is used, the value of δ_s is 0.15 K.

Figure 3 shows two examples (one SBL and one NBL) of all of the objective methods applied to both a DH2 flight and its corresponding radiosonde. These examples show that the subjective Z_{ABL} identified using the DH2 and radiosonde data are similar (differ by only 2 m for the SBL and 12 m for the NBL), and that the objective methods

reveal a similar outcome when applied to the radiosonde data as they do for the DH2 data for both cases. Similar figures for all DH2 and radiosonde profiles used in this study can be found in Supplementary Figures S5-S69.

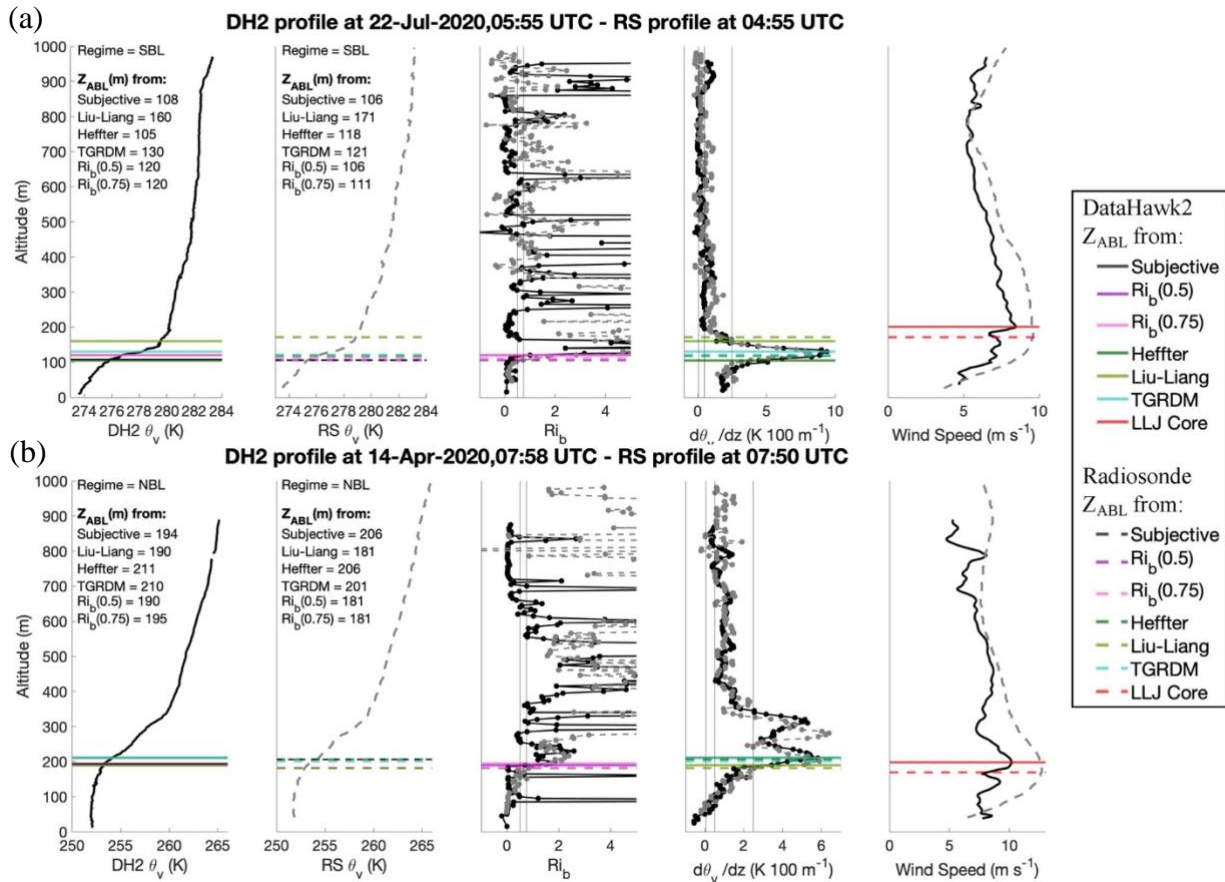
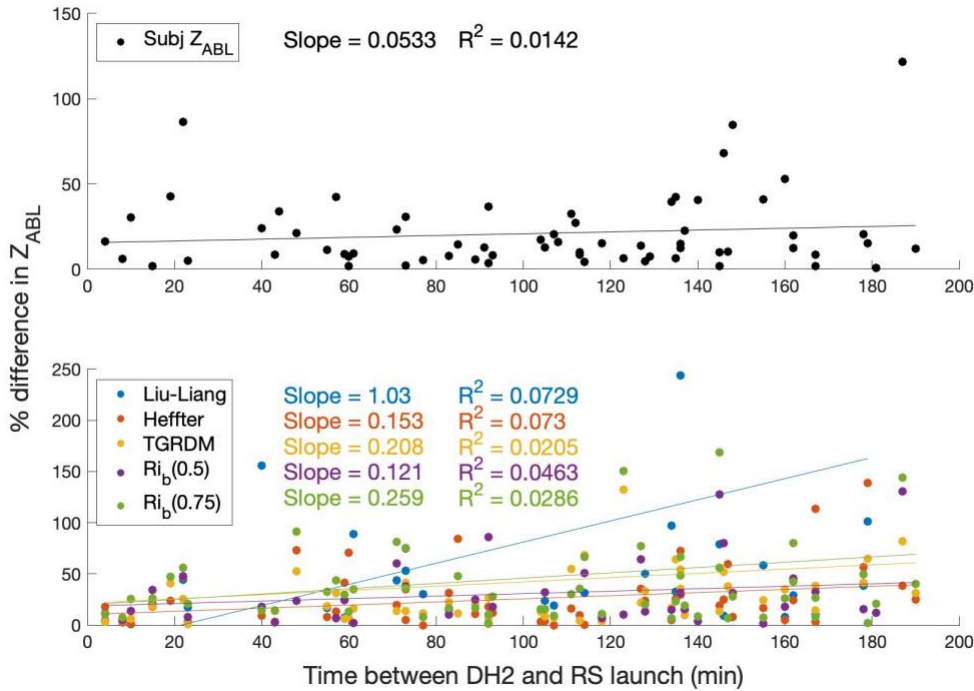


Figure 3: Demonstration of Z_{ABL} identification using all objective methods on both the DH2 (represented by solid lines) and corresponding radiosonde (represented by dashed lines) for an (a) SBL and (b) NBL case. Panel 1: θ_v profile from the DH2. Panel 2: θ_v profile from the radiosonde. Panel 3: Ri_b profiles from the DH2 (solid black) and the radiosonde (dashed grey). Panel 4: $d\theta_v/dz$ profiles from the DH2 (solid black) and the radiosonde (dashed grey). Panel 5: wind speed profiles from the DH2 (solid black) and radiosonde (dashed grey). The legend on the right indicates the Z_{ABL} detection method associated with each horizontal line in the figure. LLJ core is not in itself a Z_{ABL} detection method, but plays into the Liu-Liang method, so it is included. Each Z_{ABL} is written on the corresponding platform's θ_v profile.

While the radiosonde and DH2 profiles generally exhibit a similar structure due to the close time and space proximity (the radiosondes were launched <600 m from the DH2 flights), the subjective Z_{ABL} identified in those profiles differ by 1-101 m. In general, the deviation between Z_{ABL} from the DH2 and the radiosonde increases with increasing time proximity. Figure 4 shows the percent difference between DH2 and radiosonde subjective Z_{ABL} (top panel), as well as the percent difference between the DH2 and radiosonde objective Z_{ABL} for each method (bottom panel) as a function of time difference in minutes between the DH2 and radiosonde launch. The best fit linear regression for each method shows that as time between the DH2 and radiosonde launch increases, the differences in Z_{ABL} increase as well. However, the increase in percent difference between subjective Z_{ABL} from the DH2 and radiosonde as time between the launches increases is not significant at the 5% significance level (probability value of 0.34). Therefore, we are

confident that Z_{ABL} does not significantly change for DH2 and radiosonde launches up to 3.16 hours apart, which justifies the use of the radiosonde closest in time to each DH2 to test if there is similar efficacy of the different objective methods.



425 **Figure 4:** Relative difference between subjective Z_{ABL} from the DH2 and subjective Z_{ABL} from the radiosonde
 closest in time to the DH2 launch (black dots, top panel) and relative difference between objective Z_{ABL} from the
 DH2 and objective Z_{ABL} from the radiosonde closest in time to the DH2 launch (colored dots, bottom panel) versus
 absolute time difference in minutes between the DH2 and radiosonde launches. A few outlier points are not shown,
 430 as they lie outside the y-axis range. Lines of best fit are included for the subjective Z_{ABL} and for each objective
 method, and the slope and R^2 values of each line is written next to the legend.

3. Results and discussion

3.1 Efficacy of objective Z_{ABL} identification methods

Whereas the objective methods all rely on information from one variable (or two, in the case of the Liu-Liang method
 for an SBL), the subjective methods use a combination of methods which can only be weighted properly by visual
 435 analysis. This is why the subjective methods result in more accurate Z_{ABL} identification and provide a good basis for
 comparison with Z_{ABL} identified by the objective methods.

To determine how well the different objective methods worked, Z_{ABL} identified by each objective method is compared
 to the subjective Z_{ABL} . Figure 5 shows scatter plots comparing the objective to the subjective Z_{ABL} in each case, along
 with the associated best fit linear regression, coefficient of determination (R^2), slope, and probability value (p-value)
 440 resulting from a paired two sample T-test. For instances in which there were two DH2 flights in closest time proximity
 to the same radiosonde launch, the results from that radiosonde profile are plotted only once.

The R^2 value demonstrates how much of the variation in objective Z_{ABL} can be explained by the difference in subjective Z_{ABL} . Slope values are also included to help evaluate the level of correspondence between the subjective and objective Z_{ABL} . Additionally, looking at the intercept combined with the slope value tells us whether the objective method tends to over- or underestimate Z_{ABL} compared to the subjective method. Lastly, the p-value tells us whether the relationship between subjective and objective Z_{ABL} can be considered statistically significant at the 5% significance level (a p-value less than 0.05 indicates that there is a 95% chance the relationship is due to true correlation).

Based on the DH2 data in these scatter plots, the method that gives the greatest R^2 value is the $Ri_b(0.5)$ method (R^2 of 0.653, Fig. 5d), followed by the $Ri_b(0.75)$ method (R^2 of 0.537, Fig. 5e). These are followed closely by the Heffter method (R^2 of 0.485, Fig. 5b). The TGRDM method has the fourth highest R^2 value (R^2 of 0.316, Fig. 5c). The only objective method with a very low R^2 value is the Liu-Liang method (R^2 of 0.0907, Fig. 5a). The slope values for all methods are within 0.3 of 1, the closest to 1 being the $Ri_b(0.75)$ method (slope of 1.02), followed by the TGRDM method (slope of 1.1) and Heffter method (slope of 1.18). These slope values greater than 1 and positive intercept indicate that these methods generally overestimate Z_{ABL} when applied to the DH2 data, compared to the subjective Z_{ABL} . The results of the $Ri_b(0.5)$ method and the Liu-Liang method, however, are more complex, as the slope values are both less than 1 (0.721 and 0.708 respectively), but the intercepts are both positive. This indicates that these methods overestimate Z_{ABL} for a shallow ABL, but underestimate it for a deep ABL when applied to the DH2 data. Comparing the p-values for all relationships to the 5% significance level, the relationship between subjective and objective Z_{ABL} is significant for every method (p-value is less than 0.05). These p-values follow the same order as the R^2 values, with the lowest p-value found for the $Ri_b(0.5)$ (indicating the highest significance) and the highest p-value for the Liu-Liang method (indicating the lowest significance).

The radiosonde data gives a slightly different conclusion. Here, the method that gives the greatest R^2 value is the Heffter method (R^2 of 0.558, Fig. 5b), followed by the $Ri_b(0.5)$ method (R^2 of 0.42, Fig. 5d). The $Ri_b(0.75)$ method and the TGRDM method have lower R^2 values (0.207 and 0.225 in Fig. 5e and 5c, respectively). As was the case for the DH2 data, the only objective method with a very low R^2 is the Liu-Liang method (R^2 of 0.00597, Fig. 5a), which is also echoed by a slope value far from 1, of 0.171. The slope values for the rest of the methods are not as close to 1 as they are for the DH2 data, but they are all within 0.5 of 1. The TGRDM has a slope value of 1, and the method with the next closest value to 1 is the Heffter method at 1.13. Both of these methods have a positive intercept, which indicates that these method tends to overestimate Z_{ABL} when applied to the radiosonde data used in the current study. The rest of the methods have a slope of less than 1 and positive intercept, indicating that they tend to overestimate Z_{ABL} for a shallow ABL, but underestimate it for a deep ABL when applied to the radiosonde data used in the current study. Lastly, the p-values follow the same order as the R^2 values, with the lowest p-value found for the Heffter method (indicating the highest significance) and the highest p-value for the Liu-Liang method (indicating the lowest significance). Unlike the DH2 results, for the radiosonde, the p-values for all relationships compared to the 5% significance level show that the relationship between subjective and objective Z_{ABL} is significant for every method except the Liu-Liang method, in which the p-value is greater than 0.05.

Lastly, Fig. 5f compares subjective Z_{ABL} from the radiosondes to subjective Z_{ABL} from the DH2. The high R^2 value (0.752) indicates a strong correlation between subjective Z_{ABL} from both platforms, which demonstrates that Z_{ABL} usually did not change much between the DH2 and radiosonde launches in each case. Interestingly, there is enhanced deviation from the line of best fit for a shallower ABL, and better agreement for a deeper ABL. However, this might simply be due to the greater number of samples with Z_{ABL} below ~ 200 m. The very low p-value of $2.62e-18$ demonstrates the high significance in the relationship between Z_{ABL} from the DH2 and radiosondes.

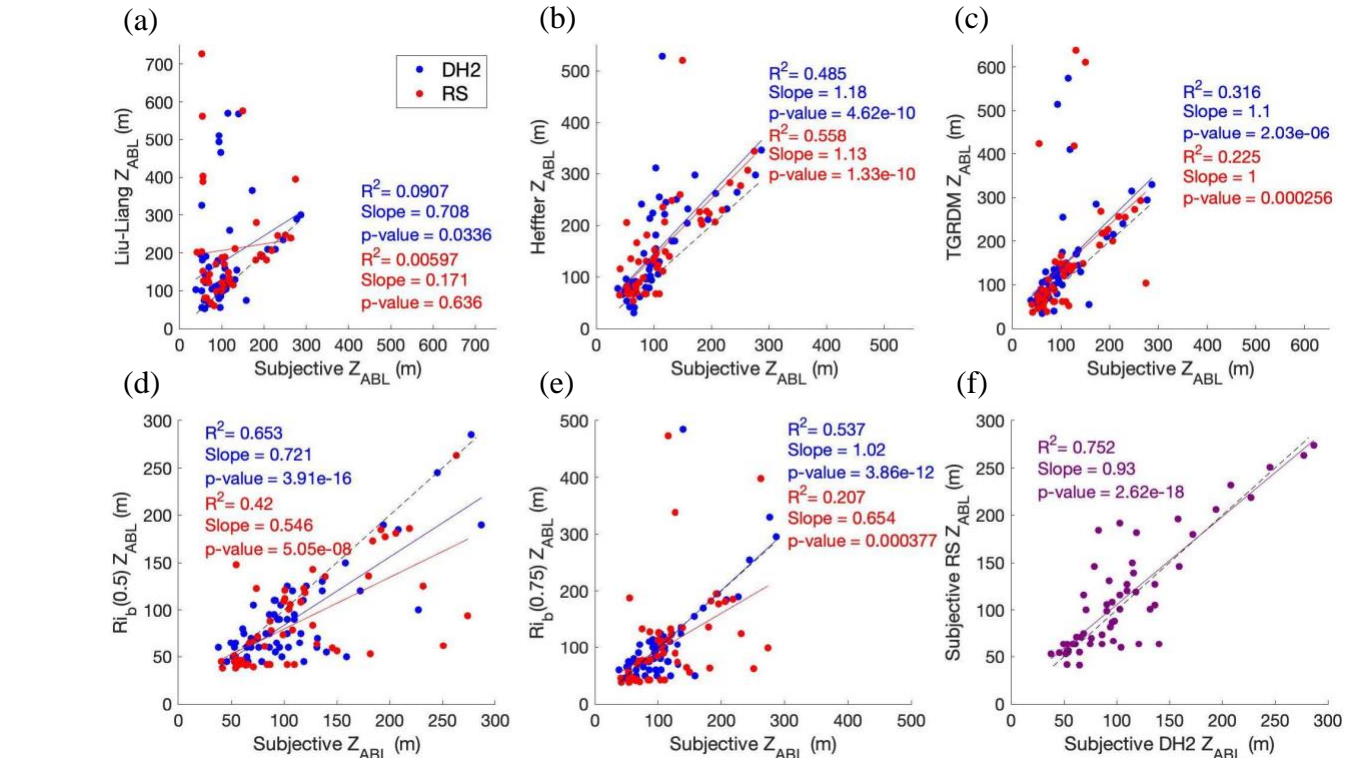


Figure 5: Relationships between subjective Z_{ABL} and objective Z_{ABL} from the (a) Liu-Liang method (50 DH2 samples and 40 RS samples), (b) Heffter method (61 DH2 samples and 53 RS samples), (c) TGRDM method (62 DH2 samples and 55 RS samples), and (d, e) Ri_b method (65 DH2 samples and 57 RS samples). Blue dots represent DH2 data and red dots represent radiosonde data. The solid blue line (solid red line) on each panel is the line of best fit for the DH2 (radiosonde) data. (f) Relationship between subjective Z_{ABL} from the radiosonde and subjective Z_{ABL} from the DH2 with line of best fit in purple (57 samples). Each panel is overlaid by the corresponding R^2 , slope value, and p-value. The dashed black line on each panel is a line with slope of 1 and y-intercept of 0, for reference.

Supplementary Figures S70 and S71 show the results presented in Fig. 5, but separated by stability regime, where S70 shows results for only SBLs, and S71 shows results for only NBLs. The primary takeaways from separating the results into stability regime is that, for SBLs, the Heffter and TGRDM methods have the highest R^2 values, and for NBLs, the $Ri_b(0.5)$ and $Ri_b(0.75)$ methods have the highest R^2 values for the DH2 data. For the radiosonde data, the Heffter and $Ri_b(0.5)$ methods have the highest R^2 values for SBLs; for NBLs, the same methods, plus the Liu-Liang method, have the highest R^2 values.

Additional analysis was completed to assess the cumulative frequency distribution for the difference in objective Z_{ABL} relative to the subjective Z_{ABL} . To do this, relative difference between the objective and subjective Z_{ABL} in each case

500 and for each method was determined. These results are included in Fig. 6a for the DH2 profiles, and in Fig. 6b for the radiosonde profiles. For example, about 26% of the time, the Liu-Liang Z_{ABL} was within 10% of the subjective Z_{ABL} for the DH2 data.

505 Figure 6a shows that, for the DH2 profiles, the $Ri_b(0.75)$ method results in the highest percent of cases to be within 10% of the subjective Z_{ABL} , followed by the $Ri_b(0.5)$ method. Interestingly, the Liu-Liang method results in the third highest percent of cases to be within 10% of the subjective Z_{ABL} . However, the Liu-Liang method falls behind other methods as the percent difference range is increased above 20%. Additionally, the Liu-Liang method has the highest percent of cases in which no Z_{ABL} is found at all for the DH2 profiles, as well as about 20% of cases that have greater than 100% difference from the subjective Z_{ABL} . This trend indicates that, while the Liu-Liang method sometimes works to find a Z_{ABL} close to the subjective Z_{ABL} , it also fails to find a Z_{ABL} close to the subjective Z_{ABL} , or to find any Z_{ABL} , in many cases. Another important finding is that the Ri_b method using either critical value never fails to find a Z_{ABL} , and the number of cases within each relative difference range is greater for the Ri_b method than that for all other methods.

510 The information presented in the bar graph for the radiosonde profiles (Fig. 6b) leads to a similar conclusion. As for the DH2 profiles, the Ri_b method results in the highest percent of cases to be within 10% of the subjective Z_{ABL} (but for this platform, the $Ri_b(0.5)$ method does best). Here, the Liu-Liang method results in the fourth highest percent of cases to be within 10% of the subjective Z_{ABL} , and performs more poorly as the percent difference range is increased. The Liu-Liang method also has the highest percent of cases in which no Z_{ABL} is found at all, followed by the Heffter and TGRDM methods, which was also true for the DH2 data. As for the DH2, there are no radiosonde cases in which the Ri_b method with either critical value finds no Z_{ABL} . The main difference between Fig. 6b of the radiosonde data and Fig. 6a of the DH2 data is that, while the $Ri_b(0.75)$ method applied to the DH2 data was always more successful than the $Ri_b(0.5)$ method for percent difference ranges below 70%, for the radiosonde data, the $Ri_b(0.5)$ method proves to always be more successful than the $Ri_b(0.75)$ method. We suspect that this results from the radiosonde data being more smoothed, which produces less sporadic Ri_b values as the atmosphere transitions from the ABL to the free atmosphere, compared to the less smoothed DH2 data. This smoothing of the radiosonde data is applied by the Vaisala software to remove any effect of the chaotic pendulum swing directly after launch, while the wire unwinds. Thus, a lower critical Ri_b threshold value may be better applicable when more smoothing or filtering procedures are applied to a dataset.

515
520
525

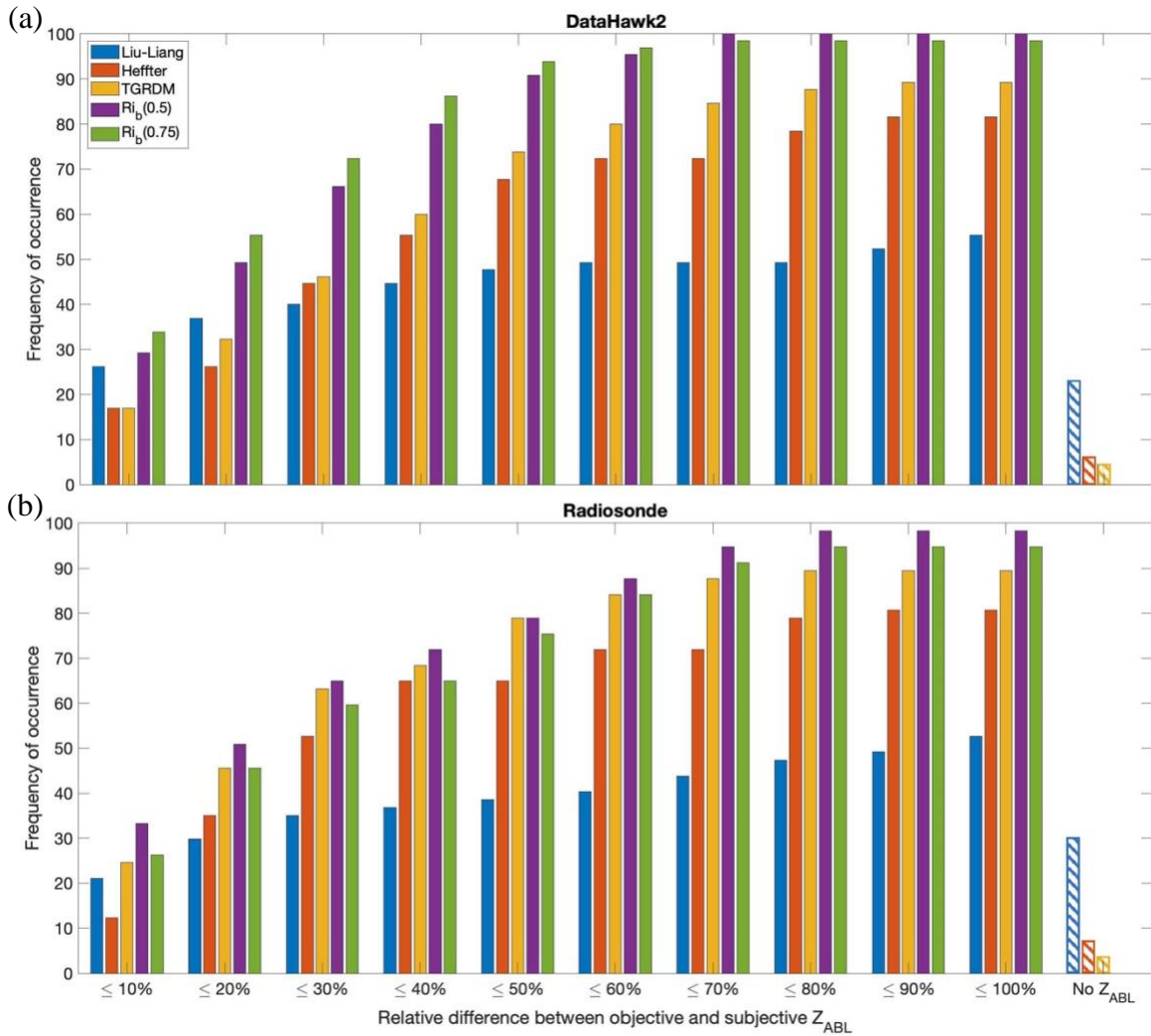


Figure 6: Bar plot showing what percent of (a) DH2 cases and (b) radiosonde cases give an objective Z_{ABL} within different percent difference ranges from the subjective Z_{ABL} using the different objective methods. Plot also shows the percent of cases for each method where no Z_{ABL} is found (labelled as “No Z_{ABL} ”).

Supplementary Figures S72 and S73 show the results presented in Fig. 6, but separated by stability regime, where S72 shows results for only SBLs, and S73 shows results for only NBLs. The primary takeaways from separating the results into stability regime is that, for both the DH2 and radiosonde, the Ri_b method has the most cases and the Liu-Liang method has the least cases with objective Z_{ABL} within 10% of the subjective Z_{ABL} for SBLs, though the Heffter and TGRDM methods also do well. For NBLs, the Liu-Liang method actually has the most cases with objective Z_{ABL} within 10% of the subjective Z_{ABL} , followed by the Ri_b method, for both platforms.

After comparing Z_{ABL} from the different objective methods to the subjective Z_{ABL} for both the DH2 and the radiosondes (Fig. 5-6), it is found that, with the exception of the Liu-Liang method, all other methods generally provide a reasonable estimate of Z_{ABL} for both datasets, with the Ri_b method being most favorable. This is in agreement with Siebert et al. (2000), Dai et al. (2014) and Zhang et al. (2014) which found an Ri_b -based method to be preferred when mechanically-produced turbulence dominates, as is true in the central Arctic (Brooks et al., 2017). Additionally, the

efficacy of each method is similar for the DH2 and the radiosonde data, as is indicated by similar patterns in the scatter plots (Fig. 5) and bar plots (Fig. 6), despite occasional differences in radiosonde versus DH2-based Z_{ABL} estimates, which likely result from the differences in sampling methods between the two platforms. Most specifically, the DH2 samples directly from the surface in most cases, so it observes important ABL features that support accurate stability and Z_{ABL} identification, whereas the radiosonde, which only samples down to 23 m at the lowest, may miss these features. Additionally, the DH2 samples with higher resolution, again contributing to its ability to record complex fine scale features which the radiosonde might miss. However, the similarity in efficacy of the objective methods between both platforms supports the fact that the objective Z_{ABL} identification methods that were improved using the high resolution DH2 data are indeed robust across platforms with different sampling methods.

This is further explored by re-running the analysis with DH2 profiles averaged over 5 m, 10 m, and 20 m bins instead of 1 m bins, to determine how sensitive the efficacy of the methods is to the vertical resolution of the data. When comparing objective Z_{ABL} found using the coarser data to the original subjective Z_{ABL} for each method, the F-test reveals that generally the R^2 values do not differ significantly from those found using 1 m binned data at the 5% significance level. The only exceptions are the Liu-Liang method at all larger bin sizes, and the Heffter method when using a 10 or 20 m bin size, which all manifest in lower R^2 value than those found using 1 m binned data. This reveals that the Liu-Liang method performs even more poorly at lower vertical resolution, and the Heffter method starts to perform more poorly at a vertical resolution of 10 m. On the other hand, the Ri_b and TGRDM methods remain just as successful when vertical resolution is reduced, and the preferred Ri_b critical value does not appear to depend on vertical resolution. For vertical resolution of 30 m or coarser, the altitude range over which Ri_b is calculated would have to be increased, and at this point a higher critical Ri_b value may be more applicable.

3.2 When the objective methods fail

While the Liu-Liang method sometimes works well, it is not reliable across a wide range of different profile structures. Table 4 lists the most common features which cause the Liu-Liang method to fail, along with the corresponding failure (either over- or underestimation, or no Z_{ABL} found) and an example of such a situation shown in the Supplementary Figures. Option 1a causes failure because the $d\theta_v/dz$ criteria are not met anywhere in the profile, meaning that the method reverts to using the LLJ core height as Z_{ABL} . However, the LLJ core was observed to usually be above the subjective Z_{ABL} (supported by Stull, 1988; Jakobson et al., 2013; and Mahrt et al., 2014). This cause for failure agrees with Dai et al. (2011) which found that using LLJ core height to define SBL top produces results inconsistent with those from other methods.

Any of the other objective methods would be a good choice for objectively determining Z_{ABL} for a dataset similar to the DH2 and radiosonde datasets (high resolution profiles in the central Arctic environment). However, each method still struggles in some situations. As for the Liu-Liang method, this information is shared for all other objective methods in Table 4. The primary downfall of the Heffter method is that it identifies Z_{ABL} as the point where θ_v is 2 K warmer than θ_v at the bottom of the θ_v inversion. Failures noted in options 1-3 in Table 4 all occur when this criterion does not accurately identify the ABL top. The primary downfall of the TGRDM method, as noted in options 1-2 in

Table 4, is that the strongest point of the θ_v inversion is not always at the ABL top. The TGRDM method also fails to find any Z_{ABL} if there is no θ_v inversion strong enough to exceed the threshold necessary for Z_{ABL} identification as laid out in Sect. 2.4.3. Lastly, the failure of the Ri_b method occurs due to the difficulty of defining an accurate critical value which correctly captures the likelihood of turbulence for all cases.

The last column in Table 4 lists the cases in which the objective Z_{ABL} differs by more than 50% from the subjective Z_{ABL} for the DH2 data, or no Z_{ABL} was found, which can be referenced in the Supplementary Figures for all examples of the profile structures that are not as conducive to the success of the different objective methods.

Table 4: Summary of the features which lead to failure by each objective method, along with examples of DH2 cases that exemplify each failure, which can be found in the Supplementary Figures. The last column indicates the Supplementary Figures associated with cases in which the objective Z_{ABL} was greater than 50% different than the subjective Z_{ABL} , or no objective Z_{ABL} was found.

Objective method	Features which lead to failure	Resulting failure	Examples	Cases with >50% difference in Z_{ABL}
Liu-Liang	<p>1. A weak θ_v inversion persists throughout the whole profile</p> <p>a. LLJ core altitude is well above the ABL top</p> <p>b. No LLJ</p> <p>2. NBL capped by weak θ_v inversion</p>	<p>1a. Overestimation of Z_{ABL}</p> <p>1b. No Z_{ABL} found</p> <p>2. Overestimation of Z_{ABL}</p>	<p>1a. S6 on 24 March at 12:09 UTC</p> <p>1b. S33 on 30 April at 14:07 UTC</p> <p>2. S54 on 17 July at 13:30 UTC</p>	S6, S9, S10, S11, S13, S14, S17, S18, S19, S24, S29, S30, S31, S32, S33, S34, S35, S39, S41, S46, S48, S49, S52, S54, S55, S57, S58, S59, S60, S62, S64, S65, S66, S68
Heffter	<p>1. SBL height is not the altitude at which θ_v is 2 K warmer than θ_v at the surface</p> <p>a. SBL extends higher</p> <p>b. SBL does not extend as high</p> <p>2. NBL capped by weak θ_v inversion</p> <p>3. Only shallow, weak θ_v inversion(s)</p>	<p>1a. Underestimation of Z_{ABL}</p> <p>1b. Overestimation of Z_{ABL}</p> <p>2. Overestimation of Z_{ABL}</p> <p>3. No Z_{ABL} found</p>	<p>1a. S5 on 23 March at 13:52 UTC</p> <p>1b. S42 on 21 June at 13:13 UTC</p> <p>2. S52 on 18 July at 13:10 UTC</p> <p>3. S40 on 6 May at 14:50 UTC</p>	S4, S15, S16, S17, S25, S29, S32, S33, S34, S40, S41, S45, S47, S51, S52, S54, S55, S56, S58, S59, S66
TGRDM	<p>1. θ_v inversion is strongest at the surface</p> <p>2. θ_v inversion is strongest within the entrainment zone</p> <p>3. Only shallow, weak θ_v inversion(s)</p>	<p>1. Underestimation of Z_{ABL}</p> <p>2. Overestimation of Z_{ABL}</p> <p>3. No Z_{ABL} found</p>	<p>1. S10 on 7 April (radiosonde profile)</p> <p>2. S64 on 22 July at 7:37 UTC</p> <p>3. S57 on 20 July at 11:28 UTC</p>	S12, S13, S14, S24, S25, S29, S32, S45, S46, S52, S54, S57, S58, S59, S60, S64, S66
Ri_b	<p>1. Ri_b is not capturing transition from turbulent to laminar atmosphere</p> <p>2. Critical value is not accurate</p>	1/2. Over- or underestimation of Z_{ABL}	1/2. S8 on 29 March at 12:24 UTC and S45 on 30 June at 8:39 UTC	<p>$Ri_b(0.5)$: S8, S17, S18, S52, S57, S66</p> <p>$Ri_b(0.75)$: S17, S52, S57, S66</p>

When applying these objective methods to a large dataset to automatically identify Z_{ABL} , it is recommended that some level of pre-screening is applied to flag cases that contain the features or structural patterns summarized in Table 4 that would make certain objective methods have difficulty identifying Z_{ABL} , and choosing which objective method to use based on that.

On the simplest level, one could choose which objective Z_{ABL} detection method to use based on stability regime. Given the results in Supplementary Figures S70-73, the best choice to use for SBLs might be the Heffter method (highest R^2 and higher frequency of cases within 10% of the subjective Z_{ABL} when compared to NBL cases, from both the DH2 and radiosonde data) and the best choice to use for NBLs might be the Ri_b method with either critical value (highest R^2 's from the DH2 data and higher frequency of cases within 10% of the subjective Z_{ABL} when compared to SBL cases, from both the DH2 and radiosonde data). However, when separating out the efficacy of the objective methods depending on stability regime, the Ri_b method has a combination of a high R^2 values and a high percentage of cases with objective Z_{ABL} within 10% of the subjective Z_{ABL} for both stability regimes, so this would be the best choice to apply to all profiles if one wanted to choose a single method, preferably with critical value of 0.5.

Overall, the objective methods are more likely to agree with each other as well as with the subjective Z_{ABL} for cases with more simplistic structures, such as those with strong θ_v inversions with a base at or just below the top of the ABL, those with LLJ core altitude at or just above the top of the ABL, and those with consistently and somewhat gradually increasing θ_v with altitude above the entrainment zone.

605 4. Summary and conclusions

By comparing subjective Z_{ABL} identified visually in θ_v , humidity (both RH and mixing ratio), and Ri_b profiles to objectively determined Z_{ABL} , the performance of several published methods (i.e., Liu-Liang, Heffter, TGRDM, and Ri_b) are evaluated across 65 DH2 UAS profiles. When comparing objective to subjective Z_{ABL} for each DH2 case, the method that is most successful (combination of high R^2 value, low p-value, and slope close to 1) is the Ri_b method with either critical value of 0.5 or 0.75 (Fig. 5). When calculating the percent of DH2 cases in which the objective Z_{ABL} is within certain percent difference ranges from the subjective Z_{ABL} , the Ri_b method is also most successful (Fig. 6). The Heffter and TGRDM methods also produce reasonable results according to Fig. 5 and 6. The only objective method that largely fails at accurately identifying Z_{ABL} is the Liu-Liang method.

In the process of applying these different objective methods to the DH2 data, some threshold and qualifying values were modified to be better applicable to the UAS dataset. While these adjustments were made to best suit the 65 DH2 profiles analyzed in this study which occurred between March and July of 2020, these adjustments should yield better results for identifying Z_{ABL} over sea ice during any season and location in the central Arctic. We hypothesize this because the ABL structures sampled by the DH2 in the current study were diverse and encompass the variety of ABL structures commonly observed in the central Arctic (which are typically shallow and either stable or neutral) throughout the entire year. Additionally, since the locations of the DH2 flights in this study range from deep in the sea

ice pack to near the sea ice edge, we are confident that the adjustments made will be applicable for identifying Z_{ABL} in either environment.

625 Testing these adjustments outside of the 65 DH2 flights, the modified techniques were also applied to the radiosonde profiles closest in time to each DH2 flight, to determine if the methods work similarly on data from another sensing platform with different sampling methods. Radiosonde profiles closest in time proximity to the DH2 flights were used under the assumption that the ABL structure would change minimally between the launch of the two platforms (supported by Fig. 4), and thus applying the methods of subjective and objective Z_{ABL} detection would lead to a similar conclusion. For the radiosonde data, the Heffter and Ri_b methods prove most successful in terms of having a high R^2 value, low p-value, and slope closest to 1 when compared to the other objective methods (Fig. 5). Additionally, the 630 Ri_b method also proves most successful when looking at the percent of cases in which the objective Z_{ABL} was within different percent difference ranges for the radiosondes, as it did for the DH2 (Fig. 6). Once again, the only method that consistently provided unfavorable results is the Liu-Liang method. These similar conclusions demonstrate that the adapted objective methods are indeed robust across platforms despite differences in sampling method, which suggest that one can take the methods and apply them to UAS, radiosonde, or other profile data alike, without having 635 to tweak them.

These findings show that no single method works well 100% of the time. Given this, the best way to accurately identify Z_{ABL} across a variety of conditions in the Arctic atmosphere is to visually analyze the θ_v , humidity, and Ri_b profiles for each case individually. However, as subjective identification is time consuming and requires expert knowledge of the physical processes that dictate ABL structure, then in the case of large datasets that require automated processing 640 techniques, the current study reveals that the Ri_b , Heffter, or TGRDM methods are most suitable for such a task, with the preferred method being the Ri_b method with critical value of 0.5. For data with vertical resolution of 10 m or coarser, the Heffter method is no longer recommended. The Liu-Liang method does not provide consistent results in accurately identifying Arctic Z_{ABL} in many cases, especially for SBLs (Fig. S72). The most common occurrence of failure of the objective methods exists for NBLs capped by a weak θ_v inversion, so that a clear θ_v slope change between the ABL and entrainment zone is difficult for automated methods to find. In such cases, the Ri_b method was found to 645 be most reliable for identifying Z_{ABL} . A full list of features which cause each objective method to fail is provided in Table 4 above. The objective methods may also fail if the near-surface atmosphere is not well sampled, for example in the case of the radiosonde data; if ABL stability is defined by what is happening near the surface (e.g., a shallow convective layer), then this is missed by radiosonde profiles which only begin 23 m or higher, and stability regime 650 could be incorrectly diagnosed. This highlights the value of platforms which can sample the near-surface atmosphere, such as the DH2. To accommodate the above problems, a semi-automatic approach may be beneficial in which one would apply all the recommended objective methods, and visually inspect only the profiles for which the resulting Z_{ABL} diverges greatly.

The methods and results of this study for stability regime and Z_{ABL} identification are currently being applied to the 655 entire year of radiosonde data collected during the MOSAiC expedition (October 2019 – September 2020) to create a

data product containing year-long statistics on ABL characteristics in the central Arctic. Additional metrics, such as LLJ height and speed and temperature inversion layer depth and strength will be included in this product for eventual publication. Value from the DH2 data and methods used in the current study comes from the uniqueness of the location and timing of the profiles collected. Therefore, these data provide a unique opportunity to evaluate any additional Z_{ABL} detection schemes that were not addressed in this study, or that have yet to be developed, as well as can be used to learn about the intricacies of additional structural components of the Arctic atmosphere such as the entrainment zone. Lastly, we are working to derive turbulence parameters from the DH2 fine wire measurements which will enhance the value of the DH2 data in ABL studies.

660

Data availability

665

All DataHawk2 data used in this study are openly available from the National Science Foundation Arctic Data Center at <https://doi.org/10.18739/A2Z60C34R> (Jozef et al., 2021) as described in de Boer et al. (submitted). The radiosonde data are available at the PANGAEA Data Publisher at <https://doi.org/10.1594/PANGAEA.928656> (Maturilli et al., 2021). These data are subject to the MOSAiC Data Policy (Immerz et al., 2019) and will be openly available after 1 January 2023.

670

Author contributions

GdB and JC planned the DH2 data collection and acquired funding; GJ and JC conducted DH2 flights; SD provided the radiosonde data; GJ, JC, and GdB conceptualized the analysis presented in this paper; GJ analyzed the data; GJ wrote the manuscript; JC, GdB, and SD reviewed and edited the manuscript.

Competing interests

675

The authors declare that they have no conflict of interest.

Acknowledgments

680

Data used in this paper were produced as part of RV *Polarstern* cruise AWI_PS122 and of the international Multidisciplinary drifting Observatory for the Study of the Arctic Climate (MOSAiC) with the tag MOSAiC20192020. We thank all those who contributed to MOSAiC and made this endeavor possible (Nixdorf et al., 2021). Radiosonde data were obtained through a partnership between the leading Alfred Wegener Institute (AWI), the Atmospheric Radiation Measurement (ARM) User Facility, a US Department of Energy facility managed by the Biological and Environmental Research Program, and the German Weather Service (DWD). Non-author contributors to DH2 design, data collection, and data processing include Dale Lawrence¹, Jonathan Hamilton^{1,2,4}, Radiance Calmer^{2,3}, Brian Argrow^{1,5}, Steven Borenstein^{1,5}, Abhiram Doddi¹, Julia Schmale⁶, and Andreas Preußer⁷.

685

¹Dept. of Aerospace Engineering Sciences, University of Colorado Boulder

²Cooperative Institute for Research in Environmental Sciences, University of Colorado Boulder

³National Snow and Ice Data Center, University of Colorado Boulder

⁴NOAA Physical Sciences Laboratory

⁵Integrated Remote and In-Situ Sensing, University of Colorado Boulder

690 ⁶Swiss Federal Institute of Technology Lausanne

⁷University of Trier

Financial support

Collection and analysis of atmospheric boundary layer data with the DataHawk2 was funded by the National Science Foundation (award OPP 1805569, de Boer, PI). Additional funding and support were provided by the Cooperative
695 Institute for Research in Environmental Sciences, the National Oceanic and Atmospheric Administration Physical Sciences Laboratory, and the Alfred Wegener Institute.

References

- Alfred-Wegener-Institut Helmholtz-Zentrum für Polar- und Meeresforschung: Polar Research and Supply Vessel
POLARSTERN operated by the Alfred-Wegener-Institute, *Journal of Large-Scale Research Facilities*, 3,
700 A119, <https://doi.org/10.17815/jlsrf-3-163>, 2017.
- AMS Glossary of Meteorology: Bulk Richardson number: https://glossary.ametsoc.org/wiki/Bulk_richardson_number, last access: 4 March 2022, 2012.
- Andreas, E. L., Horst, T. W., Grachev, A. A., Persson, P. O. G., Fairall, C. W., Guest, P. S., and Jordan, R. E.:
705 Parametrizing turbulent exchange over summer sea ice and the marginal ice zone, *Q. J. Roy. Meteor. Soc.*,
136, 927–943, <https://doi.org/10.1002/qj.618>, 2010.
- Banta, R. M., Pechugina, Y. L., and Newsom, R. K.: Relationship between Low-Level Jet Properties and Turbulence
Kinetic Energy in the Nocturnal Stable Boundary Layer, *J. Atm. Sci.*, 60, 2549-2555,
[https://doi.org/10.1175/1520-0469\(2003\)060<2549:RBLJPA>2.0.CO;2](https://doi.org/10.1175/1520-0469(2003)060<2549:RBLJPA>2.0.CO;2), 2003.
- Brooks, I. M., Tjernström, M., Persson, P. O. G., Shupe, M. D., Atkinson, R. A., Canut, G., Birch, C. E., Mauritsen,
710 T., Sedlar, J., and Brooks, B. J.: The Turbulent Structure of the Arctic Summer Boundary Layer During The
Arctic Summer Cloud-Ocean Study, *J. Geophys. Res.-Atmos.*, 122, 9685–9704,
<https://doi.org/10.1002/2017JD027234>, 2017.
- Collaud Coen, M., Praz, C., Haefele, A., Ruffieux, D., Kaufmann, P., and Calpini, B.: Determination and
climatology of the planetary boundary layer height above the Swiss plateau by in situ and remote sensing
715 measurements as well as by the COSMO-2 model, *Atmos. Chem. Phys.*, 14, 13205–13221,
<https://doi.org/10.5194/acp-14-13205-2014>, 2014.
- Dai, C., Gao, Z., Wang, Q., Cheng, G.: Analysis of Atmospheric Boundary Layer Height Characteristics over the
Arctic Ocean Using the Aircraft and GPS Soundings, *Atmos. Ocean. Sci. Lett.*, 4, 124-130,
<https://doi.org/10.1080/16742834.2011.11446916>, 2011.
- 720 Dai, C., Wang, Q., Kalogiros, J. A., Lenschow, D. H., Gao, Z., and Zhou, M.: Determining Boundary-Layer Height
from Aircraft Measurements, *Bound.-Lay. Meteorol.*, 152, 277–302, <https://doi.org/10.1007/s10546-014-9929-z>, 2014.

- de Boer, G., Chapman, W., Kay, J. E., Medeiros, B., Shupe, M. D., Vavrus, S., and Walsh, J.: A Characterization of the Present-day Arctic Atmosphere in CCSM4, *J. Climate*, 25, 2676–2695, <https://doi.org/10.1175/JCLI-D-11-00228.1>, 2012.
- 725 de Boer, G., Calmer, R., Jozef, G., Cassano, J., Hamilton, J., Lawrence, D., Borenstein, S., Doddi, A., Cox, C., Schmale, J., Preußner, A. and Argrow, B.: Observing the Central Arctic Atmosphere and Surface with University of Colorado Uncrewed Aircraft Systems, *Nat. Sci. Data*, submitted 2022.
- Delle Monache, L., Perry, K. D., Cederwall, R. T., and Ogren, J. A.: In situ aerosol profiles over the Southern Great Plains cloud and radiation test bed site: 2. Effects of mixing height on aerosol properties, *J. Geophys. Res.-Atmos.*, 109, D6, <https://doi.org/10.1029/2003jd004024>, 2004.
- 730 Esau, I., and Sorokina, S.: Climatology of the Arctic Planetary Boundary Layer, in: *Atmospheric Turbulence, Meteorological Modeling and Aerodynamics*, edited by: Lang, P. R. and Lombargo, F. S., Nova Science Publishers, Inc., New York, 3-58, ISBN:978-1-60741-091-1, 2010.
- 735 Garratt, J. R.: Review: the atmospheric boundary layer, *Earth-Sci. Rev.*, 37, 89-134, [https://doi.org/10.1016/0012-8252\(94\)90026-4](https://doi.org/10.1016/0012-8252(94)90026-4), 1994.
- Grachev, A. A., Fairall, C. W., Persson, P. O. G., Andreas, E. L., and Guest, P. S.: Stable Boundary-layer Scaling Regimes: The SHEBA Data, *Bound.-Lay. Meteorol.*, 116, 201–235, <https://doi.org/10.1007/s10546-004-2729-0>, 2005.
- 740 Hamilton, J., de Boer, G., Lawrence, D.: The DataHawk2 Uncrewed Aircraft System for Atmospheric Research, *Atmos. Meas. Tech.*, submitted 2021.
- Heffter, J. L.: Transport layer depth calculations, in: *Proceedings of the 2nd Joint Conference on Applications of Air Pollution Modelling*, American Meteorological Society, New Orleans, LA, American Meteorological Society, 787–791, 1980.
- 745 Immerz, A., Frickenhaus, S., von der Gathen, P., Shupe, M., Morris, S., Nicolaus, M., Schneebeil, M., Regnery, J., Fong, A., Snoeijs-Leijonmalm, P., Geibert, W., Rabe, B., Herber, A., Krumpfen, T., Singha, S., Jaiser, R., Ransby, D., Schumacher, S., Driemel, A., Gerchow, P., Schäfer, A., Schewe, I., Ajjan, M., Glöckner, F. O., Schäfer-Neth, C., Jones, C., Goldstein, J., Jones, M., Prakash, G., Rex, M.: MOSAiC Data Policy, Zenodo, <https://doi.org/10.5281/zenodo.4537178>, 2019.
- 750 Jakobson, L., Vihma, T., Jakobson, E., Palo, T., Männik, A., and Jaagus, J.: Low-level jet characteristics over the Arctic Ocean in spring and summer, *Atmos. Chem. Phys.*, 13, 11089-11099, <https://doi.org/10.5194/acp-13-11089-2013>, 2013.
- Jenkins, A. D., Paskyabi, M. B., Fer, I., Gupta, A., and Adakudlu, M.: Modelling the Effect of Ocean Waves on the Atmospheric and Ocean Boundary Layers, *Energy Proced.*, 24, 166–175, <https://doi.org/10.1016/j.egypro.2012.06.098>, 2012.
- 755 Jozef, G., de Boer, G., Cassano, J., Calmer, R., Hamilton, J., Lawrence, D., Borenstein, S., Doddi, A., Schmale, J., Preußner, A., and Argrow, B.: DataHawk2 Uncrewed Aircraft System data from the Multidisciplinary drifting Observatory for the Study of Arctic Climate (MOSAiC) campaign, B1 level, Arctic Data Center [data set], <https://doi.org/10.18739/A2Z60C34R>, 2021.

- 760 Karlsson, J., and Svensson, G.: Consequences of poor representation of Arctic sea-ice albedo and cloud-radiation interactions in the CMIP5 model ensemble, *Geophys. Res. Lett.*, 40, 4374–4379, <https://doi.org/10.1002/grl.50768>, 2013.
- Kruppen, T., von Albedyll, L., Goessling, H. F., Hendricks, S., Juhls, B., Spreen, G., Willmes, S., Jakob Belter, H., Dethloff, K., Hass, C., Kaleschke, L., Katlein, C., Tian-Kunze, X., Ricker, R., Rostosky, P., Rückert, J., Singha, S., and Sokolova, J.: MOSAiC drift expedition from October 2019 to July 2020: sea ice conditions from space and comparison with previous years, *The Cryosphere*, 15, 3897-3920, <https://doi.org/10.5194/tc-15-3897-2021>, 2021.
- 765 Lesins, G., Duck, T. J., and Drummond, J. R.: Surface Energy Balance Framework for Arctic Amplification of Climate Chang, *J. Climate*, 25, 8277–8288, <https://doi.org/10.1175/JCLI-D-11-00711.1>, 2012.
- 770 Liu, S., and Liang, X. Z.: Observed Diurnal Cycle Climatology of Planetary Boundary Layer Height, *J. Climate*, 23, 5790–5809, <https://doi.org/10.1175/2010JCLI3552.1>, 2010.
- Lüpkes, C., Vihma, T., Birnbaum, G., and Wacker, U.: Influence of leads in sea ice on the temperature of the atmospheric boundary layer during polar night, *Geophys. Res. Lett.*, 35, L03805, <https://doi.org/10.1029/2007GL032461>, 2008.
- 775 Mahrt, L., Vickers, D., and Andreas, E. L.: Low-Level Wind Maxima and Structure of the Stable Stratified Boundary Layer in the Coastal Zone, *J. Appl. Meteorol. Clim.*, 53, 363-376, <https://doi.org/10.1175/JAMC-D-13-0170.1>, 2014.
- Marsik, F. J., Fischer, K. W., McDonald, T. D., and Samson, P. J.: Comparison of Methods for Estimating Mixing Height Used during the 1992 Atlanta Filed Intensive, *J. Appl. Meteorol. Clim.*, 34, 1802–1814, [https://doi.org/10.1175/1520-0450\(1995\)034<1802:COMFEM>2.0.CO;2](https://doi.org/10.1175/1520-0450(1995)034<1802:COMFEM>2.0.CO;2), 1995.
- 780 Maturilli, M., Holdridge, D. J., Dahlke, S., Graeser, J., Sommerfeld, A., Jaiser, R., Deckelmann, H., and Schulz, A.: Initial radiosonde data from 2019-10 to 2020-09 during project MOSAiC, Alfred Wegener Institute, Helmholtz Centre for Polar and Marine Research, Bremerhaven, PANGAEA [dataset], <https://doi.org/10.1594/PANGAEA.928656>, 2021.
- 785 Mayer, S., Jonassen, M. O., Sandvik, A., and Reuder, J.: Profiling the Arctic Stable Boundary Layer in Advent Valley, Svalbard: Measurements and Simulations, *Bound.-Lay. Meteorol.*, 143, 507–526, <https://doi.org/10.1007/s10546-012-9709-6>, 2012.
- Nicolaus M., Perovich D., Spreen G., Granskog M., Albedyll L., Angelopoulos M., Anhaus P., Arndt S., Belter H., Bessonov V., Birnbaum G., Brauchle J., Calmer R., Cardellach E., Cheng B., Clemens-Sewall D., Dadic R., Damm E., Boer G., Demir O., Dethloff K., Divine D., Fong A., Fons S., Frey M., Fuchs N., Gabarró C., Gerland S., Goessling H., Gradinger R., Haapala J., Haas C., Hamilton J., Hannula H.-R., Hendricks S., Herber A., Heuzé C., Hoppmann M., Høyland K., Huntemann M., Hutchings J., Hwang B., Itkin P., Jacobi H.-W., Jaggi M., Jutila A., Kaleschke L., Katlein C., Kolabutin N., Krampe D., Kristensen S., Kruppen T., Kurtz N., Lampert A., Lange B., Lei R., Light B., Linhardt F., Liston G., Loose B., Macfarlane A., Mahmud M., Matero I., Maus S., Morgenstern A., Naderpour R., Nandan V., Niubom A., Oggier M., Oppelt N., Pätzold F., Perron C., Petrovsky T., Pirazzini R., Polashenski C., Rabe B., Raphael I., Regnery J., Rex M., Ricker R., Riemann-
- 795

- Campe K., Rinke A., Rohde J., Salganik E., Scharien R., Schiller M., Schneebeli M., Semmling M., Shimanchuk E., Shupe M., Smith M., Smolyanitsky V., Sokolov V., Stanton T., Stroeve J., Thielke L., Timofeeva A., Tonboe R., Tavri A., Tsamados M., Wagner D., Watkins D., Webster M., Wendisch M.:
800 Overview of the MOSAiC expedition – Snow and sea ice, *Elementa Science of the Anthropocene*, 9, <https://doi.org/10.1525/elementa.2021.000046>, 2022.
- Nixdorf, U., Dethloff, K., Rex, M., Shupe, M., Sommerfeld, A., Perovich, D., Nicolaus, M., Heuzé, C., Rabe, B., Loose, B., Damm, E., Gradinger, R., Fong, A., Maslowski, W., Rinke, A., Kwok, R., Spreen, G., Wendisch, M., Herber, A., Hirsekorn, M., Mohaupt, V., Frickenhaus, S., Immerz, A., Weiss-Tuider, K., König, B.,
805 Mengedoht, D., Regnery, J., Gerchow, P., Ransby, D., Krumpfen, T., Morgenstern, A., Haas, C., Kanzow, T., Rack, F. R., Saitzev, V., Sokolov, V., Makarov, A., Schwarze, S., Wunderlich, T., Wurr, K., and Boetius, A.: MOSAiC Extended Acknowledgement, Zenodo, <https://doi.org/10.5281/zenodo.5179738>, 2021.
- Pesenson, I.: Implementation and evaluation of the Heffter Method to calculate the height of the planetary boundary layer above the ARM Southern Great Plains site, Lawrence Berkeley National Laboratory, <https://escholarship.org/uc/item/6pp1d93m>, 2003.
810
- Seibert, P., Beyrich, F., Gryning, S.-E., Joffre, S., Rasmussen, A., and Tercier, P.: Review and intercomparison of operational methods for the determination of the mixing height, *Atmos. Environ.*, 34, 1001–1027, [https://doi.org/10.1016/S1352-2310\(99\)00349-0](https://doi.org/10.1016/S1352-2310(99)00349-0), 2000.
- Serreze, M. and Barry, R.: Processes and impacts of Arctic amplification: A research synthesis, *Global and Planetary Change*, 77, 85-96, <https://doi.org/10.1016/j.gloplacha.2011.03.004>, 2011.
815
- Shupe, M. D., Rex, M., Dethloff, K., Damm, E., Fong, A. A., Gradinger, R., Heuzé, C., Loose, C., B., Makarov, A., Maslowski, W., Nicolaus, M., Perovich, D., Rabe, B., Rinke, A., Sokolov, V., and Sommerfeld, A.: The MOSAiC Expedition : A Year Drifting with the Arctic Sea Ice, NOAA Arctic Report Card, 1–8, <https://doi.org/10.25923/9g3v-xh92>, 2020.
- 820 Shupe, M. D., Rex, M., Blomquist, B., Persson, P. O. G., Schmale, J., Uttal, T., Althausen, D., Angot, H., Archer, S., Bariteau, L., Beck, I., Bilberry, J., Bucci, S., Buck, C., Boyer, M., Brasseur, Z., Brooks, I. M., Calmer, R., Cassano, J., Castro, V., Chu, D., Costa, D., Cox, C. J., Creamean, J., Crewell, S., Dahlke, S., Damm, E., de Boer, G., Deckelmann, H., Dethloff, K., Dütsch, M., Ebell, K., Ehrlich, A., Ellis, J., Engelmann, R., Fong, A. A., Frey, M. M., Gallagher, M. R., Ganzeveld, L., Gradinger, R., Graeser, J., Greenamyre, V., Griesche, H.,
825 Griffiths, S., Hamilton, J., Heinemann, G., Helmig, D., Herber, A., Heuzé, C., Hofer, J., Houchens, T., Howard, D., Inoue, J., Jacobi, H.-W., Jaiser, R., Jokinen, T., Jourdan, O., Jozef, G., King, W., Kirchgaessner, A., Klingebiel, M., Krassovski, M., Krumpfen, T., Lampert, A., Landing, W., Laurila, T., Lawrence, D., Lonardi, M., Loose, B., Lüpkes, C., Maahn, M., Macke, A., Maslowski, W., Marsay, C., Maturilli, M., Mech, M., Morris, S., Moser, M., Nicolaus, M., Ortega, P., Osborn, J., Pätzold, F., Perovich, D. K., Petäjä, T., Pilz, C., Pirazzini, R., Posman, K., Powers, H., Pratt, K. A., Preußner, A., Quéléver, L., Radenz, M., Rabe, B., Rinke, A., Sachs, T., Schulz, A., Siebert, H., Silva, T., Solomon, A., Sommerfeld, A., Spreen, G., Stephens, M.,
830 Stohl, A., Svensson, G., Uin, J., Viegas, J., Voigt, C., von der Gathen, P., Wehner, B., Welker, J. M., Wendisch, M., Werner, M., Xie, Z. Q., Yue, F.: Overview of the MOSAiC expedition: Atmosphere, *Elementa*:

- Science of the Anthropocene, 10, <https://doi.org/10.1525/elementa.2021.00060>, 2022.
- 835 Sivaraman, C., Mcfarlane, S., Chapman, E., Jensen, M., Toto, T., Liu, S., and Fischer, M.: Planetary Boundary Layer (PBL) Height Value Added Product (VAP) : Radiosonde Retrievals. ARM User Facility, DOE/SC-ARM/TR-132, 2013.
- Snyder, B. J., and Strawbridge, K. B.: Meteorological analysis of the Pacific 2001 air quality field study, *Atmos. Environ.*, 38, 5733–5743, <https://doi.org/10.1016/j.atmosenv.2004.02.068>, 2004.
- 840 Steeneveld, G. J., van de Wiel, B. J. H., and Holtslag, A. A. M.: Diagnostic Equations for the Stable Boundary Layer height: Evaluation and Dimensional Analysis, *J. Appl. Meteorol. Clim.*, 46, 212–225, <https://doi.org/10.1175/JAM2454.1>, 2007.
- Stroeve, J., Holland, M. M., Meier, W., Scambos, T., and Serreze, M.: Arctic sea ice decline: Faster than forecast, *Geophys. Res. Lett.*, 34, L09501, <https://doi.org/10.1029/2007GL029703>, 2007.
- 845 Stroeve, J. C., Kattsov, V., Barrett, A., Serreze, M., Pavlova, T., Holland, M., and Meier, W. N.: Trends in Arctic sea ice extent from CMIP5, CMIP3 and observations, *Geophys. Res. Lett.*, 39, L16502, <https://doi.org/10.1029/2012GL052676>, 2012.
- Stull, R. B.: *An Introduction to Boundary Layer Meteorology*, Kluwer Academic Publishers, The Netherlands, 670 pp., 1988.
- 850 Tjernström, M., Leck, C., Persson, P. O. G., Jensen, M. L., Oncley, S. P., and Targino, A.: The Summertime Arctic Atmosphere: Meteorological Measurements during the Arctic Ocean Experiment 2001, *B. Am. Meteorol. Soc.*, 85, 1305–1321, <https://doi.org/10.1175/BAMS-85-9-1305>, 2004.
- Tjernström, M., Birch, C. E., Brooks, I. M., Shupe, M. D., Persson, P. O. G., Sedlar, J., Mauritsen, T., Leck, C., Paatero, J., Szczodrak, M., and Wheeler, C. R.: Meteorological conditions in the central Arctic summer during the Arctic Summer Cloud Ocean Study (ASCOS), *Atmos. Chem. Phys.*, 12, 6863–6889, <https://doi.org/10.5194/acp-12-6863-2012>, 2012.
- 855 Vihma, T., Lüpkes, C., Hartmann, J., and Savijärvi, H.: Observations and Modelling of Cold-air Advection over Arctic Sea Ice, *Bound.-Lay. Meteorol.*, 117, 275–300, <https://doi.org/10.1007/s10546-004-6005-0>, 2005.
- Zhang, Y., Gao, Z., Li, D., Li, Y., Zhang, N., Zhao, X., and Chen, J.: On the computation of planetary boundary-layer height using the bulk Richardson number method, *Geosci. Model Dev.*, 7, 2599–2611, <https://doi.org/10.5194/gmd-7-2599-2014>, 2014.
- 860 Zilitinkevich, S., and Baklanov, A.: Calculation of the Height of the Stable Boundary Layer in Practical Applications, *Bound.-Lay. Meteorol.*, 105, 389–409, <https://doi.org/10.1023/A:1020376832738>, 2002.