



1 **Look-up tables resolved by complex refractive index to correct**
2 **particle sizes measured by common research-grade optical particle**
3 **counters**

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19 **Abstract**

20 Optical particle counters (OPCs) are widely used to measure aerosol particle number size
21 distribution at atmospheric ambient conditions and over a large size range. Their measurement
22 principle is based on the dependence of light scattering on particle size. However, this
23 dependence is complex since it is not a monotonic function of size and because light scattering
24 depends on the particle composition (i.e., the complex refractive index, *CRl*) and morphology.
25 Therefore, the conversion of the measured scattered intensity to the desired particle size
26 depends on the microphysical properties of the sampled aerosol population and whose
27 relationship is not necessarily unique at all sizes. While these complexities have been
28 addressed before, corrections are typically applied as needed and are not standardised. This
29 issue is addressed here by providing a consistent database of precomputed correction factors
30 for a wide range of complex refractive index values representing the variability in the



31 composition of atmospheric aerosols. These correction factors are calculated for five different
32 commercial OPCs (USHAS, PCASP, FSSP, GRIMM and its airborne version Sky-GRIMM,
33 CDP) by assuming Mie theory for homogeneous spherical particles, and for the real part of the
34 *CRI* between 1.33 and 1.75 in steps of 0.01 and the imaginary part between 0.0 and 0.4 in
35 steps of 0.001. Correction factors for mineral dust are provided at a *CRI* of $1.53 - 0.003i$ and
36 account for the asphericity of these particles. The datasets described in this paper are
37 distributed through an open access repository: <https://doi.org/10.25326/234> (license CC BY,
38 Formenti et al., 2021) maintained by the French national center for Atmospheric data and
39 services AERIS for data users and geophysicists who determine number size distribution
40 measurements from OPCs in their atmospheric aerosol research. Application and caveats of
41 the *CRI*-corrections factors are presented and discussed.

42 The dataset presented in this paper is not only useful for correcting the size distribution from
43 OPCs when particle refractive indices are known, but also when this parameter can only be
44 estimated. As well, this dataset can be used to calculate uncertainties or sensitivities
45 associated with aerosol volume, mass, or extinction obtained from OPCs given little or no
46 knowledge of the refractive index.

47 **1. Introduction**

48 Aerosol particles are some of the more elusive, but highly climate-relevant components of the
49 atmosphere (Boucher et al., 2013). While airborne, aerosols interact with atmospheric radiation
50 at wavelengths from the ultraviolet to the infrared, and serve as condensation nuclei for liquid
51 and ice clouds (Seinfeld and Pandis, 2006). Upon deposition, they can change the productivity
52 of marine and land ecosystems (Kanakidou et al., 2018). As well as directly influencing
53 atmospheric composition by their emission, they also act as a sink for some reactive gases
54 (e.g., Seinfeld and Pandis, 2006; Kanakidou et al., 2018). Through these processes aerosol
55 particles affect the Earth's climate but they also can have adverse effects on the environment.
56 In particular, aerosol particles can degrade air quality to the detriment of human health
57 (Shiraiwa et al., 2017).

58 These varied effects of aerosols are largely due to by their broad size spectrum, which is.
59 characterised by sizes ranging from a few nanometres to tens of micrometres depending on
60 the sources and the mechanisms of emission, and on the transformations that they undergo
61 whilst airborne (Seinfeld and Pandis, 2006). The typical particle size distribution in the
62 atmosphere is a continuum of four lognormal modes (nucleation, comprising particles with
63 diameters up to 10 nm; Aitken, comprising particles of diameters between 10 nm and 100 nm;
64 accumulation, made up of particles with diameters from 100 nm to approximately 2.5 μm ;



65 coarse, comprising particles with diameters larger than $2.5 \mu\text{m}$), with different amplitudes,
66 mode diameters, relative proportions, and chemical compositions (Seinfeld and Pandis, 2006).
67 The nucleation, Aitken and most of the accumulation modes are due to condensation of
68 precursors gases in the atmosphere, whereas particles in the coarse mode are mostly
69 generated by wind friction on the sea and bare soil surfaces, and from biogenic or volcanic
70 primary emissions. These modes are also characterised by different lifetimes in the
71 atmosphere, ranging from a few minutes (nucleation mode) to days and weeks (Aitken,
72 accumulation, and coarse modes).

73 Whilst the particle size distribution is a critical parameter to assess the effects of aerosols on
74 radiation, clouds, chemistry, ocean and terrestrial productivity, and human health, its
75 measurement is challenging. Instrumental techniques do not cover the entire particle size
76 range, but only portions of it. Furthermore, these different instrumental techniques measure
77 particle size using various operating principles, including light-scattering, aerodynamic and
78 electric mobility properties of aerosol (Baron and Willeke, 2001; Hinds, 1999). Therefore, the
79 particle size measured experimentally is an operational definition based on particle density,
80 real and complex refractive index (*CRI*, the relation between spectral optical properties and
81 chemical composition), and morphology (Baron and Willeke, 2001; Hinds, 1999).

82 Optical particle counters (OPCs) provide fast (better than 1 Hz) measurements over a large
83 dynamic range, both in concentration and size, including sub- and super-micron particles
84 (Baron and Willeke, 2001; Hinds, 1999; Wendisch and Brenguier, 2013). By adapting the
85 geometry of the sensing volume and the wavelength of the light source, the design of the OPC
86 can be customised to different applications (i.e., sampling mostly fine or coarse particles, more
87 or less absorbing aerosols, minimise the effects of asphericity, etc.), therefore making them a
88 versatile tool for atmospheric aerosol research. Research-grade OPCs are used worldwide in
89 laboratory and field studies, including on research aircraft (Collins et al., 2000; Haywood et al.,
90 2003a; 2003b; Reid et al., 2003; Osborne et al., 2008; Ryder et al., 2013; Di Biagio et al., 2015;
91 Denjean et al., 2016; Petzold et al., 2009; Weinzerl et al., 2017; De Perim de Faria et al., 2017;
92 Schafer et al., 2019; Brock et al., 2019; Wu et al., 2020).

93 The operating principle of the OPC is based on the fact that the intensity of monochromatic or
94 white light scattered by an airborne particle (single or ensembles) in a given scattering direction
95 depends on its size (Baron and Willeke, 2001; Hinds, 1999; Wendisch and Brenguier, 2013);
96 and as a consequence, the intensity of light-scattering measured in a known sensing volume
97 and at a known wavelength can be converted into particle size.

98 In principle, the scattering cross-section $C_{sc\alpha}$ measured by an OPC can be converted into an
99 optical equivalent diameter (D_{EO}) based on calibration with non-absorbing spherical particle



100 latex spheres (PSL) or equivalent scattering material of known CRI at the wavelength of light
101 used in the instrument. However, this is complicated by the fact that atmospheric aerosols
102 have different compositions (CRI) than the calibration material, and that the intensity of
103 scattered light also depends on particle morphology (Dubovik et al., 2006; Huang et al., 2021).
104 The differences in CRI and morphology between the calibration spheres and natural aerosols
105 cause the OPC-determined size to be different from the real particle size. The error can
106 propagate to size-relevant datasets, ultimately generating biases in the estimates of aerosol
107 impacts on weather, climate, and human health (Huang et al., 2021). Henceforth, representing
108 the number size distribution of the real atmospheric aerosol requires being able to convert the
109 value of D_{EO} into an equivalent spherical particle geometrical (i.e., volume equivalent) diameter
110 (D_{geo}) corresponding to the CRI of the sampled aerosols at the OPC operating wavelength.

111 In practice, the equivalence between D_{EO} at the reference CRI and D_{geo} at the CRI of the
112 sampled aerosols is obtained by calculating, at each CRI value, the value of D_{geo} corresponding
113 to the same scattering cross-section C_{sca} as that obtained when using D_{EO} . However, the
114 scattering cross-section may not depend linearly or even monotonically on particle diameter.
115 In the approximation of spherical particles, this effect is due to the Mie resonance and ripple
116 oscillations in the light scattering functions. As a consequence, for a given CRI and D_{EO} , the
117 scattering cross-section C_{sca} could correspond to a number of values of D_{EO} . This well-known
118 and documented problem in the expert community (e.g., Garvey and Pinnick, 1983; Liu et al.,
119 1992; Jaenicke and Hanusch, 1993; Pinnick et al., 2000; Collins et al., 2000; Reid et al., 2003;
120 Osborne et al., 2008; Petzold et al., 2009; Rosenberg et al., 2012; Wendisch and Brenguier,
121 2013; Brock et al., 2016; 2019; Walser et al., 2017; Moore et al., 2021) and is illustrated in
122 Figure 1 for C_{sca} characteristic functions of the CDP instrument (see Table 1). For selected CRI
123 values corresponding to typical atmospheric aerosol types (see section 4), Figure 1 shows how
124 a single value of C_{sca} can correspond to numerous particle diameters, with the exception of the
125 case of highly-absorbing aerosols which causes the oscillations to be smoothed out. Additional
126 uncertainties in the derived size distribution arise from the calibration of the instruments, as
127 well as on their operation (counting statistics, flow rate, and sampling losses; e.g., Brock et al.,
128 2019).

129 Various instrument and data users have reported solutions to these problems previously,
130 corresponding to their interests and expertise:

131 1/ instrument developers and engineers have proposed processing methodologies taking into
132 account the entire chain of instrument operations, in particular the size- and scattering-
133 dependent calibrations (e.g., Rosenberg et al., 2012; Walser et al., 2017);



134 2/ data users and geophysicists who rely on external expertise for calibration and instrument
135 characterization and have proposed methods for adapting the measured size distribution to
136 the ambient refractive indices, which have sometimes been evaluated using concurrent
137 measurements of aerosol composition (e.g., Di Biagio et al. 2015; 2017; Denjean et al., 2016).

138 In this paper, we target the second category of scientists by providing standardized corrections
139 for particle sizing calculations that account for the dependence of angular scattering on particle
140 composition as represented by the *CRI*. Building on reported expert investigations, we address
141 this technical note those who are regular or occasional users of data from some of the most
142 common research-grade OPCs. Such data are available through open access datasets, which
143 are becoming more and more popular through large-scale ground-based and airborne
144 environmental research infrastructures, notably in several European organizations (e.g.,
145 Aerosol, Clouds and Trace Gases, ACTRIS, <https://www.actris.fr/actris-eu/>; In-service Aircraft
146 for a Global Observing System, IAGOS, www.iagos.org; and European Facility for Airborne
147 Research, EUFAR <https://www.eufar.net/>) and within integrative science projects (e.g., the
148 Global Aerosol Synthesis and Science Project, GASSP; Reddington et al., 2017). In using the
149 data for their research and publications, these users are not necessarily data instrument
150 operators, nor do they necessarily have the knowledge, expertise or time to perform and
151 evaluate the *CRI*-adapted corrections.

152 This paper describes lookup tables of pre-computed scattering functions and size correction
153 factors that are provided as downloadable ASCII files for a range of values of *CRI* relevant to
154 atmospheric aerosols. Section 2 (Methods) presents the characteristics of the selected OPCs
155 and the methodology used for the optical calculations and data presentation. Section 3
156 (Description of the datasets) presents the output files of the calculations that are available for
157 interested users. Section 4 (Results) discusses the difficulties of the correction schemes and
158 the associated caveats. Noted are the consequences of the lack of uniqueness in the
159 relationship between particle diameter and scattering intensity and the influence of particle
160 shape. Indeed, the calculations are performed under the assumption of spherical particles,
161 which is a good approximation for a wide range of aerosol types and conditions. Additionally,
162 a shape-dependent formulation is also applied to optimize the correction for aspherical mineral
163 dust (e.g. Chou et al., 2008; Kandler et al., 2009; Huang et al., 2020), which is one of the most
164 abundant aerosol types dominating the aerosol optical depth over large areas of the globe, in
165 particular downwind of their source regions (Kok et al., 2021). Section 4 also presents the
166 consequences on the representation of the size distribution and the calculation of optical
167 properties for different methods commonly used to deal with sizes when the correction is not
168 unique. Finally, section 5 provides the results and recommendations.



169 2. Methods

170 2.1. Instruments

171 Instruments considered in this paper are research-grade OPCs used for ambient atmospheric
172 measurements and laboratory studies. Their nominal technical specifications are summarized
173 in Table 1, which includes nominal size measurement range, light source wavelength, and
174 range of scattering angles. The instruments are:

- 175 1. The Passive Cavity Aerosol Spectrometer Probe (PCASP, Model 100, Droplet
176 Measurement Technologies, Boulder, CO) operates at 632.8 nm and measures light
177 scattering between 35° and 145° collecting light from the direct and the reflected light beam
178 (angular range 35°–120° and 60°–145°, respectively), so that light scattered between 60°
179 and 120° will be counted twice. The particle number size distribution is determined over 31
180 channels for diameters between 0.1 and 3.0 μm (e.g., Liu et al., 1992; Reid and Hobbs,
181 1998; Rosenberg et al., 2012).
- 182 2. The Ultra High Sensitivity Aerosol Spectrometer (UHSAS, Droplet Measurement
183 Technologies, Boulder, CO) has a ground-based version but the airborne version is most
184 commonly used (e.g., Cai et al., 2008; Petzold et al., 2013; Brock et al., 2016; Kupc et al.,
185 2018). It operates at 1054 nm and provides the number size distribution of particles with
186 optical equivalent diameters ranging from 0.04 to 1 μm in 99 nominal size classes. The light
187 scattering sensing angle range (22°–158°) reported by Cai et al. (2008) was subsequently
188 corrected by Petzold et al. (2013) and Brock et al. (2016), who reported that the optically
189 active range is circularly symmetric from 33° to 148°, with a blind region between 75.2 and
190 104.8°.
- 191 3. The Forward Scattering Spectrometer Probe (FSSP, Model 300, Droplet Measurement
192 Technologies, Boulder, CO) is a widely used aircraft probe that measures light scattering at
193 632.8 nm in an optically active volume extending from 3 to 15° to retrieve the number size
194 distribution in a nominal size range from 0.28 to 20.5 μm over 30 size classes (Baumgardner
195 et al., 1992; Petzold et al., 2013).
- 196 4. The Cloud Droplet Probe (CDP, Model 300, Droplet Measurement Technologies, Boulder,
197 CO) measures light scattering at 658 nm in an optically-active volume extending from 4° to
198 12° to retrieve the number size distribution in a nominal size range from 2 to 50 μm over 30
199 size classes (Baumgardner et al., 1992; Petzold et al., 2013).
- 200 5. The ground-based GRIMM and airborne Sky-OPC (Grimm Aerosol Technik, models 1.109
201 and 1.129, Ainring, Germany; hereafter the Sky-OPC will be named Sky-GRIMM) retrieve
202 the particle number distribution over 31 size classes distributed for nominal diameters



203 between 0.25 and 32 μm . These particle counters operate at 655 nm, and measure light
204 scattered by the direct beam from 30° to 150° and by the reflected beam between 81° and
205 98° due to two face-to-face parabolic mirrors (opening angles of 120° and 18°, respectively)
206 that collect light around a mean scattering angle of 90° (Heim et al., 2008). Like the PCASP,
207 the light scattered between 81° and 98° has twice the weight relative to the intensity within
208 30°-81° and 98°-150°.

209 Table 1 also reports the reference material used for the calibration for each of the OPCs: NIST-
210 certified polystyrene latex spheres (PSL) for the UHSAS, PCASP-100, FSSP-300, and
211 GRIMM 1.109/Sky-OPC 1.129. The CDP is calibrated with non-absorbing glass bead
212 reference particles. All these (spherical) reference materials have a refractive index of 1.59 –
213 0i at the light-source wavelength.

214 2.2. Optical calculations

215 Following Rosenberg et al. (2012), the scattering cross-section C_{sca} measured by an OPC is
216 calculated as follows:

217

$$218 \quad C_{sca} = \frac{\pi}{k^2} \int_0^{2\pi} \int_{\theta_{min}}^{\theta_{max}} (|S1(\theta, k D_p, CRI)|^2 + |S2(\theta, k D_p, CRI)|^2) \sin(\theta) w_{optics}(\theta, \varphi) d\theta d\varphi \quad (1)$$

219

220 where

- 221 – k is the wavenumber of the light
- 222 – D_p is the particle diameter,
- 223 – CRI is the particle complex refractive index,
- 224 – $S1$ is the light scattering intensity polarized in the parallel plane and $S2$ in the perpendicular
225 plane. Their squared absolute values, represented by bars in Equation 1, integrated over
226 the scattering angle range characteristic of the OPCs is the total light intensity seen by the
227 instrument;
- 228 – θ is the scattering angle between the incident laser beam and the viewing direction,
- 229 – φ is the azimuthal angle for a given scattering direction,
- 230 – $w_{optics}(\theta, \varphi)$ is a weighting function defined by the optical geometry of the OPC. As defined
231 by Rosenberg et al. (2012), $w_{optics}(\theta, \varphi)$ takes into account the fact that, at certain angles,
232 the PCASP and GRIMM/Sky-OPC measure scattered light both directly and after reflection
233 by a mirror. In the case of rotational symmetry around the laser beam, $w_{optics}(\theta, \varphi)$ is a
234 function of the scattering angle θ only.



235 Calculations are done according to two approximations:

- 236 • by assuming that particles are homogeneous spheres, the calculations use Mie theory
237 according to Bohren and Huffman (2007). The particle diameter in Equation 1 is varied
238 between 0.02 and 200 μm in logarithmically–equal steps of 0.004 (1001 values). The real
239 part of the particle CRI (n) is varied between 1.33 and 1.75 (in steps of $\Delta n = 0.01$) and the
240 imaginary part (k) from 0.0 to 0.4 (steps of $\Delta k = 0.001$) encompassing the range of values
241 expected for atmospheric aerosols (e.g., Shettle and Fenn, 1992) at the working
242 wavelengths of the OPCs;
- 243 • that the particles are ellipsoidal shaped to account for the asphericity of mineral dust
244 particles according to Huang et al. (2021), who obtained single-scattering properties of
245 ensembles of ellipsoidal dust particles by combining a shape-resolved single–particle
246 scattering database (i.e., using the method of Meng et al. (2010) who combined four
247 computational methods including Lorenz–Mie theory, T–matrix method, discrete dipole
248 approximation, and an improved geometric optics method) with the globally representative
249 shape distributions of dust aerosols (Huang et al., 2020). In this case, the CRI is set to 1.53
250 $- 0.003i$, at the upper limit of the values measured in the near-infrared for mineral dust
251 aerosols (e.g., Fig. 8 in Di Biagio et al., 2019 which is a synthesis of the then current body
252 of observations).

253 2.3. Calculation of equivalent particle diameter

254 Figure 1 shows that, due to non-monotonic behavior and oscillations, a given C_{sca} (from the
255 nominal bin value) can be associated with several particle diameters (i.e., the value is not
256 unique). The equivalence between D_{EO} at the reference CRI and D_{geo} at the aerosol CRI is
257 obtained by calculating the C_{sca} from Mie theory for the wide range of particle sizes at the
258 aerosol CRI and then, for each value of the bin size boundary (D_{bin}), using the particle size
259 whose C_{sca} is closest to that calculated from D_{EO} , in both C_{sca} intensity and size. This
260 equivalence is ultimately determined by the precision of the floating–point numbers (in our
261 calculations they are 64–bit, double–precision).

262 For each size bin, the midpoint diameter D_{mid} is the geometric mean defined as follows:

263

$$264 \quad D_{mid} = \sqrt{D_{bin,lower} * D_{bin,upper}} \quad (2)$$

265

266 where $D_{bin,lower}$ and $D_{bin,upper}$ are the lower and the upper bin diameter, respectively.



267 The bin width ($d\log D$) is calculated as follows:

268

$$269 \quad d\log D = \log \left(\frac{D_{bin,upper}}{D_{bin,lower}} \right) \quad (3)$$

270

271 3. Description of the dataset

272 Three types of files in ASCII format are provided from this study:

273 1/ The values of scattering cross section C_{sca} for particle diameters between 0.02 and 200 μm
274 in logarithmically equal steps of 0.004 and as a function of CRI ($n=1.33-1.75$; $\Delta n=0.01$, $k=0.0-$
275 0.4 ; $\Delta k=0.001$) are provided in text files whose generic name is `OPC_intensity_real_imag.txt`,
276 where *OPC* is the acronym of the particle counter, *real* is the value of the real part of the *CRI*,
277 and *imag* is the imaginary part of the *CRI*. Each file contains four columns, namely

- 278 - the real part of the *CRI*
- 279 - the imaginary part of the *CRI*
- 280 - the particle diameter used in Equation 1 (in μm)
- 281 - the corresponding value of C_{sca} (in μm^2)

282 2/ the values of the corrected bin diameter as a function of CRI ($n=1.33-1.75$; $\Delta n=0.01$, $k=0.0-$
283 0.4 ; $\Delta k=0.001$) are provided in text files whose generic name is `OPC_Diameter_real_imag.out`,
284 where *OPC* is the acronym of the particle counter, *real* is the value of the real part of the *CRI*,
285 and *imag* is the imaginary part of the *CRI*. Each file contains four columns, namely

- 286 - the nominal bin diameter (in μm) of the calibration *CRI*
- 287 - the bin diameter (in μm) calculated for the atmospheric *CRI*
- 288 - the bin midpoint diameter (D_{mid} , in μm) calculated for the atmospheric *CRI*
- 289 - the bin width ($d\log D$) calculated for the atmospheric *CRI*.

290 3/ The values of the corrected bin diameter in aspherical approximation for mineral dust are
291 provided in text files whose generic name is
292 `OPC_Diameter_non_spherical_1.53000_0.003000000i`, where *OPC* is the acronym of the
293 particle counter. The format is equivalent to that of files `OPC_Diameter_real_imag.out`.

- 294 - the nominal bin diameter (in μm) of the calibration *CRI*
- 295 - the bin diameter (in μm) calculated for the atmospheric *CRI*
- 296 - the bin midpoint diameter (D_{mid} , in μm) calculated for the atmospheric *CRI*
- 297 - the bin width ($d\log D$) calculated for the atmospheric *CRI*.



298 4. Example results

299 This section describes some of the elements of this analysis to help the reader understand the
300 size correction factors available in the datasets. To do so, we use four example values of CRI ,
301 corresponding to non-absorbing material used for calibration (polystyrene latex spheres or
302 equivalent light-scattering material; $CRI = 1.59 - 0i$), mineral dust ($CRI = 1.53 - 0.003i$; e.g.,
303 Fig. 8 in Di Biagio et al., 2019), urban aerosols ($CRI = 1.56 - 0.087i$; e.g., Radney and
304 Zangmeister, 2018), and marine aerosols ($CRI = 1.38 - 0.001i$; e.g., Zieger et al., 2017).

305 4.1. Scattering function

306 The dependence of C_{sca} values on size for the example CRI values (**Figure 2**) allows
307 exploration of a question: for a given value of CRI , is it possible to infer the particle size on the
308 entire nominal size range of each OPC? The answer is not entirely affirmative. In general
309 terms, the C_{sca} functions for wide-angle instruments (PCASP, UHSAS, and
310 GRIMM/Sky-GRIMM,) are less affected by Mie oscillations than the forward-scattering
311 instruments (FSSP-300 and CDP), and their behavior tends to be monotonic with size.
312 However, the PCASP and the GRIMM/Sky-OPC curves tend to become more flat around 1
313 μm as the imaginary part of the CRI increases. In practical terms, this precludes the possibility
314 of sizing absorbing particles in the range between 0.6 and 2 μm . For FSSP-300 and CDP, the
315 determination of particle size is problematic in the range between 1 and 5 μm . On the contrary,
316 for the UHSAS C_{sca} is monotonic in particle size for nearly the entire size range regardless of
317 particle CRI , as discussed in Moore et al. (2021). As a metric of those considerations, **Figure**
318 **2** also shows the curves of the absolute values of $d\log C_{sca}/d\log D$ representing, for each value
319 of the CRI , the rate of change of the scattering cross section across the size range. The value
320 of $d\log C_{sca}/d\log D$ equal to 1 indicates that, for a given OPC, the dependence of the light-
321 scattering cross section on particle size is linear, while the asymptotic value of 0 represents
322 the situation when C_{sca} is independent of size and the OPC cannot be used to classify the
323 particles. Oscillations and noisy behavior correspond to small scale oscillations, resulting in
324 non-monotonic changes of C_{sca} with size. **Figure S1** extends this approach and shows, for
325 each OPC considered in this paper, the contour plots of $d\log C_{sca}/d\log D$ as function of the bin
326 midpoint diameters and the real and imaginary parts of CRI .

327 **Figure 3** illustrates the effect of asphericity on the calculation of the scattering cross sections
328 C_{sca} for moderately-absorbing mineral dust particles ($CRI = 1.53 - 0.003i$). As pointed out by
329 Huang et al. (2021), calculating C_{sca} for homogeneous aspherical ellipsoid particles results in
330 higher scattering for particles larger than approximately 1 μm . Because of the enhancement in
331 particle surface area, accounting for dust asphericity in an ensemble of dust aerosols has the



332 result of reducing the oscillations that affect the scattering functions in the spherical
333 approximation.

334 **4.2. Size correction factors**

335 **Figure 5** shows the scatterplot of the bin size boundary (D_{bin}) corresponding to the atmospheric
336 CRI of the example aerosol types (mineral dust, urban and marine aerosols) compared to D_{bin}
337 values obtained for the calibration CRI . The figure clearly illustrates the features that can
338 appear, in particular for light absorption, which are evident around $1\ \mu\text{m}$ and most prominent
339 for OPCs that measure side scattering (i.e. PCASP and GRIMM/Sky-OPC). For UHSAS, the
340 increase of the imaginary part of the CRI results in oscillations in the retrieved diameter in the
341 upper part of the size domain, which depends strongly on the interpolation step used. For
342 mineral dust, this is shown for both homogeneous spheres and aspherical ellipsoid particles.
343 As expected from Figure 3, a more realistic representation of the shape distribution of mineral
344 dust reduces the ambiguities of the OPC correction factors and significantly improves the
345 resulting particle sizing. For a given D_{EO} , aspherical dust has a smaller D_{geo} than spherical dust
346 for all the OPC models, because aspherical dust has a larger C_{sca} than volume-equivalent
347 spherical dust (Fig. 2).

348 We therefore recommend care when using the values at sizes where the C_{sca} behavior is not
349 monotonic. These instances can easily be identified by the values of $d\log D$ provided in the
350 OPC_Diameter_real_imag.out data files that are negative, corresponding to the corrected size
351 of the upper bin (at atmospheric CRI) being smaller than the corrected size of the lower bin
352 (Equation 3). As an illustration, Figures 5 and 6 help understand and evaluate the possible
353 consequences on the representation of the number size distribution of mineral dust when the
354 CDP is used in with corrections calculated using the spherical approximation. For the sake of
355 generality, we use a model aerosol distribution reported by Seinfeld and Pandis (2006) that is
356 expressed as the sum of three lognormal modes (see the Appendix for the formulae and modal
357 parameters from the original Table 8.3 of the Seinfeld and Pandis (2006) publication). Figure
358 5 has three panels that show, for a CRI of $1.53-0.003i$, the relationship between the corrected
359 and the calibration bin diameters (left), the behavior of the corrected $d\log D$ with calibration bin
360 diameter (middle), and the representation of the size distribution using the corrected CDP
361 evaluation of particle size (right). In the region between 4 and $10\ \mu\text{m}$, the increase of the
362 calibration D_{bin} does not necessarily correspond to an increasing corrected D_{bin} , thereby
363 resulting in a negative $d\log D$ and in an evident discontinuity in the representation of the size
364 distribution. In the literature, there are two families of data treatment methods that have been
365 proposed to address this kind of occurrence:



- 366 1. grouping or widening of the bins of the OPCs (e.g., Johnson and Osborne, 2011), or
367 excluding specific size ranges (e.g., Denjean et al., 2016) in order to eliminate instances
368 of the scattering cross section flattening or decreasing with increasing size.
- 369 2. smoothing or fitting of the theoretical curves in order to eliminate oscillations (e.g., Liu et
370 al., 1974; Hand and Kreidenweis, 2002; Covert et al., 1990; Johnson et al., 2008; Lance
371 et al., 2010).

372 Figure 6 illustrates the application of these methods to a specific practical case, whose details
373 described in Annex 1 of the supplementary material. For method 1, we grouped instances
374 where $d\log D$ was negative to create a wider bin and for which we assigned the sum of the
375 particle counts of the original bins. For method 2, we fitted the D_{bin} curves with a
376 second-degree polynomial function for each of the three size ranges reported in Annex 1 (a_0 ,
377 a_1 and a_2 parameters in Table S2). Figure 6 shows that both methods improve the
378 representation of the size distribution around $10\ \mu\text{m}$ where the Mie oscillations are the most
379 prominent. However, both choices have potential implications for the calculation of particle
380 optical properties and mass concentrations. Table 2 shows an evaluation of these two methods
381 by comparing the calculations of the optical properties of scattering, absorption and extinction
382 using the two methods compared to those obtained with the base case in which the values of
383 $dN/d\log D$ corresponding to negative values of $d\log D$ are simply discarded. Calculations are
384 done at wavelengths relevant to modelling and remote sensing products (440, 550 and 870
385 nm, and $10\ \mu\text{m}$). It is clear that the data treatment influences their representation in a significant
386 way, in particular by radically changing the spectral dependence of extinction and absorption
387 at the shorter wavelengths (440 and 870 nm) (Table 3). Method 2 also results in increased
388 scattering and extinction at $10\ \mu\text{m}$, due to the artificial increase of the particle number
389 concentration at the highest diameters (not shown).

390 5. Summary and recommendations

391 In this paper, we describe a set of standardized corrections of particle sizing by OPC
392 instruments in order to account for the dependence of angular scattering on particle
393 composition, as represented by the particle complex refractive index CRI . This dataset of
394 corrections is based on the simple assumption of homogeneous spherical particles and the
395 use of Mie theory, and considers nominal OPC characteristics in terms of scattering angles of
396 the sensing volume and wavelengths of the light sources. The approach covers the range of
397 refractive indices expected for atmospheric aerosols, while, with the exception of mineral dust,
398 it neglects the additional complexity due to the particle asphericity, which could be important
399 for combustion aerosols.



400 In general terms, the analyses described confirm that research-grade OPC probes perform
401 very well for the size ranges and for the particle types for which they were designed, as a result
402 of careful design by experts in the field (see references in Table 1, and the overview of
403 Wendisch and Brenguier (2003)). This contrasts with the relatively poor documentation of low-
404 cost sensors that are indeed useful for complementary applications such as distributed
405 monitoring of air quality (Hagan and Kroll, 2020). The behavior of light scattering intensity with
406 size shown in Figure 2 indicates that the UHSAS performs very well for submicron particles
407 with diameters less than 800 nm, regardless of their refractive index, and represents a very
408 significant improvement compared to the PCASP instrument that operates on a similar, albeit
409 more reduced size range. The FSSP, CDP, and GRIMM/Sky-OPC should be used to size
410 particles larger than approximately 1 μm . However, GRIMM/Sky-GRIMM can be problematic
411 in the range 1-2 μm . The FSSP and CDP can be problematic below 10 μm .

412 We recommend the users consider very carefully instances when light scattering is not
413 monotonic with size, and use care when selecting the method for eliminating them. We
414 recommend users to combine as much as possible, the retrieval of the particle size distribution
415 from OPCs with concurrent, complementary measurements (e.g. particles sizers based on
416 electrical mobility or aerodynamics, lidar measures of the backscattering vertical profile,
417 gravimetric or composition measurements providing the mass concentration and composition)
418 for optical and/or mass closure in order to ensure the robustness and consistency of the
419 dataset, and improved knowledge of the *CRI*. Finally, and in order to make the best use of the
420 possibilities offered by open data policies, we also recommend that users, whenever possible,
421 to make contact with instrument operators to verify the specifics of the OPCs, their calibration,
422 and their performance during field operations.

423 Finally, the significant sensitivity of the light scattering intensity C_{sca} to the particle *CRI* suggests
424 the following considerations:

- 425 - diameters corrected for the *CRI* should be used rather than the calibration diameters even
426 if particle *CRI* is not precisely known. Even with an approximate assumption of the particle
427 origin (i.e. wind direction, time of day, season of year, air mass trajectory), assuming an
428 aerosol type and/or *CRI* based on these other environmental conditions and using the
429 corresponding *CRI*-corrected diameter is likely to be more accurate than using the
430 calibration diameters;
- 431 - The dataset presented in this paper can also be used without any knowledge of the particle
432 refractive index, or to deduce the refractive index of the aerosol if other instruments are
433 available for closure or if more than one OPC are used. This approach can also provide
434 uncertainty or sensitivity of size distribution estimates.



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455 **Data availability**

456 The datasets described in this paper are distributed at open-access repository:
457 <https://doi.org/10.25326/234> (license CC BY, Formenti et al., 2021) maintained by the French
458 national center for Atmospheric data and services, AERIS. Optical calculations with Mie theory
459 for homogeneous spherical particles have been performed with the IDL mie_single.pro routine
460 available at http://www.atm.ox.ac.uk/code/mie/mie_single.html.

461 **Competing interests.**

462 The authors declare no conflict of interest.

463 **Special issue statement**

464 This article is not part of a special issue. It is not associated with a conference.

465 **Author contribution**



466 PF and CDB performed the calculations in the spherical mode and the full data analysis. YH
467 and JFK performed the optical calculations in the aspherical mode. DB provided curation and
468 distribution of the data. MDM and MC provided insight on data analysis and performed
469 literature searches. PF wrote the manuscript with contributions from all co-authors.

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703 **Table 1.** Technical characteristics of the OPCs considered in this paper.

Instrument	Size measurement range D_{Eo} (μm)	Wavelength (nm)	Angular Range of scattered light	Calibration	Primary reference
PCASP-100	0.1–3.0 (31 channels)	632.8	35°–120° (direct beam) + 60°–145° (reflected beam)	PSL (1.59–0 <i>i</i>)	Liu et al. (1992)
UHSAS	0.04 – 1 (99 channels)	1054	33° –75.2° + 104.8°–148°	PSL (1.59–0 <i>i</i>)	Cai et al. (2008)
FSSP-300	0.275–20.5 (30 channels)	632.8	3°–15°	PSL (1.59–0 <i>i</i>)	Baumgardner et al. (1992)
CDP-300	2–50 (30 channels)	658	4°–12°	Glass beads (1.59–0 <i>i</i>)	Lance et al. (2010)
GRIMM model 1.109	0.25 – 32 (31 channels)	655 [§]	30°–150° (direct beam) + 81°–98° (reflected beam)	PSL (1.59–0 <i>i</i>)	Heim et al. (2008)
Sky-GRIMM model 1.129[#]	0.25 – 32 (31 channels)	655	30°–150° (direct beam) + 81°–98° (reflected beam)	PSL (1.59–0 <i>i</i>)	Grimm and Eatough (2009)

704 [§] Company specifications. Heim et al. (2008) reported that the working wavelength of the GRIMM
 705 1.109 is 683 nm.

706 [#] The technical characteristics of the GRIMM and Sky-OPC instruments relevant to this paper are
 707 identical and the calculations performed hold for both.
 708



709 **Table 2.** Calculated optical scattering, absorption and extinction coefficients (σ_{sca} , σ_{abs} , σ_{ext}
 710 expressed in Mm^{-1} ; $1 Mm^{-1} = 10^{-6} m^{-1}$), single scattering albedo (SSA, unitless) , mass
 711 extinction and absorption coefficients (MEC and MAC, respectively; units $m^2 g^{-1}$) obtained from
 712 Mie calculations for homogeneous spherical particles ($CRI = 1.53 - 0.003i$) assuming the size
 713 distribution of mineral dust measured by the CDP (base case, see Fig.5) and two methods for
 714 correcting instances when the corrected values of $dlogD$ become negative. Calculations are
 715 performed at wavelengths relevant to modelling and remote sensing products (440, 550 and
 716 870 nm, and $10 \mu m$).

	Base case				Method 1 Grouping/widening				Method 2 Smoothing or fitting			
	440	550	870	10	440	550	870	10	440	550	870	10
Scattering coeff. σ_{sca}	41.0	34.3	43.4	3.6	41.2	34.4	43.6	3.9	36.1	39.4	36.9	6.5
Absorption coeff. σ_{abs}	4.7	3.6	3.4	0.1	4.8	3.7	3.4	0.1	5.1	4.7	3.1	0.2
Extinction coeff. σ_{ext}	45.8	37.9	46.8	3.7	46.0	38.1	47.0	4.0	41.1	44.1	40.0	6.7
SSA	0.90	0.91	0.93	0.96	0.90	0.90	0.93	0.97	0.88	0.90	0.92	0.97
MEC	0.50	0.41	0.51	0.04	0.50	0.41	0.51	0.04	0.40	0.43	0.39	0.07
MAC	0.052	0.039	0.037	0.002	0.051	0.039	0.037	0.002	0.049	0.045	0.029	0.002

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719 **Table 3.** Same as Table 2 for the Angstrom exponents for extinction and scattering (AEE and
720 AEA, respectively; unitless) calculated between 440 and 870 nm, and mass concentration (M_c ,
721 expressed in $\mu\text{g m}^{-3}$).

	Base case	Method 1 Grouping/widening	Method 2 Smoothing or fitting
Angstrom exponent, extinction (AEE)	-0.031	-0.031	0.043
Angstrom exponent, absorption (AEA)	0.497	0.498	0.749
Mass concentration (M_c)	91.4	93.0	103.2

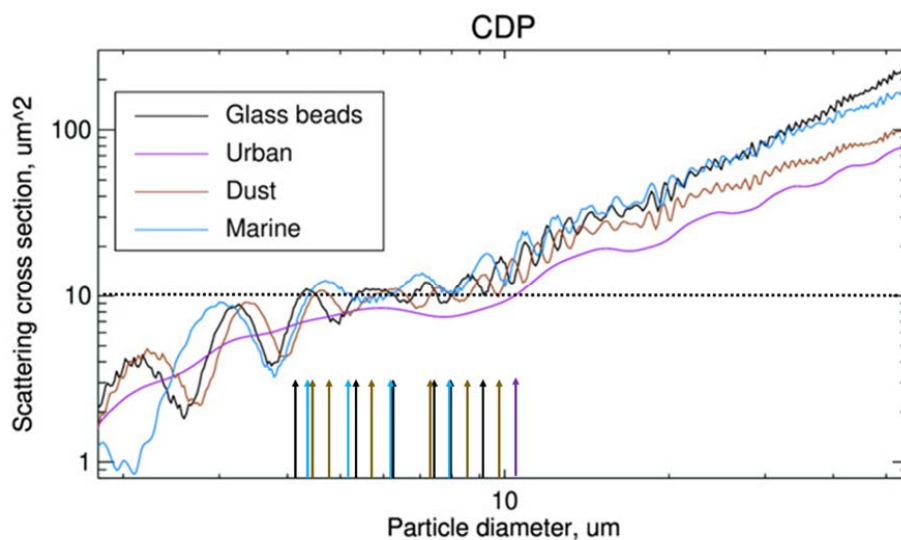
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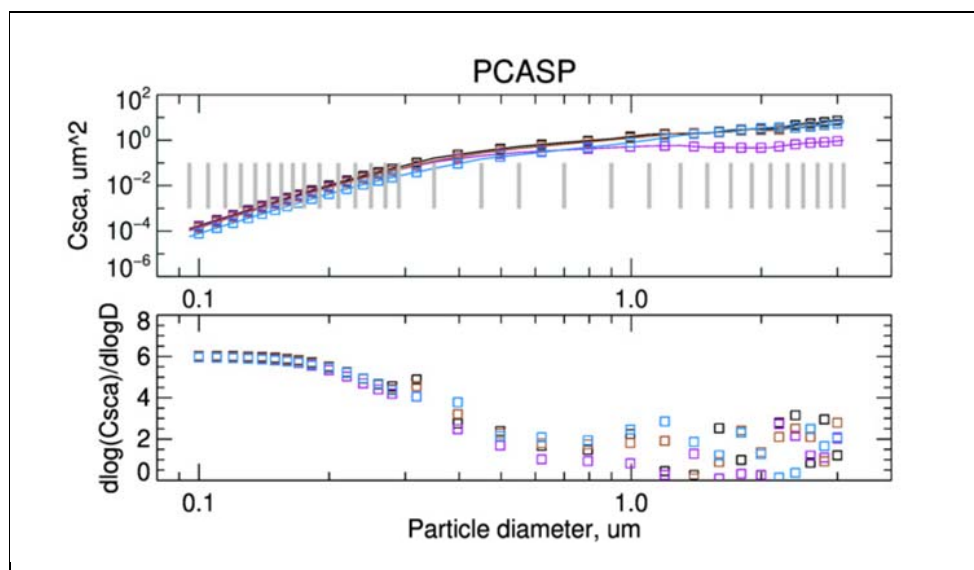
725 **Figure 1.** Scattering cross section values, C_{sca} , as a function of particle diameter calculated
726 using Mie theory for the CDP OPC considered in this paper. The black lines represent C_{sca}
727 for the calibration particles (PSL and equivalent light-scattering material). The purple lines
728 represent the C_{sca} function for absorbing urban aerosols ($CRl = 1.56 - 0.087i$), while the brown
729 lines represent C_{sca} for moderately-absorbing mineral dust ($CRl = 1.53 - 0.003i$), and the light
730 blue lines represent C_{sca} for marine aerosols ($CRl = 1.38 - 0.001i$). The dotted line represents
731 a single C_{sca} value which corresponds many values of particle diameter, represented by the
732 coloured arrows.

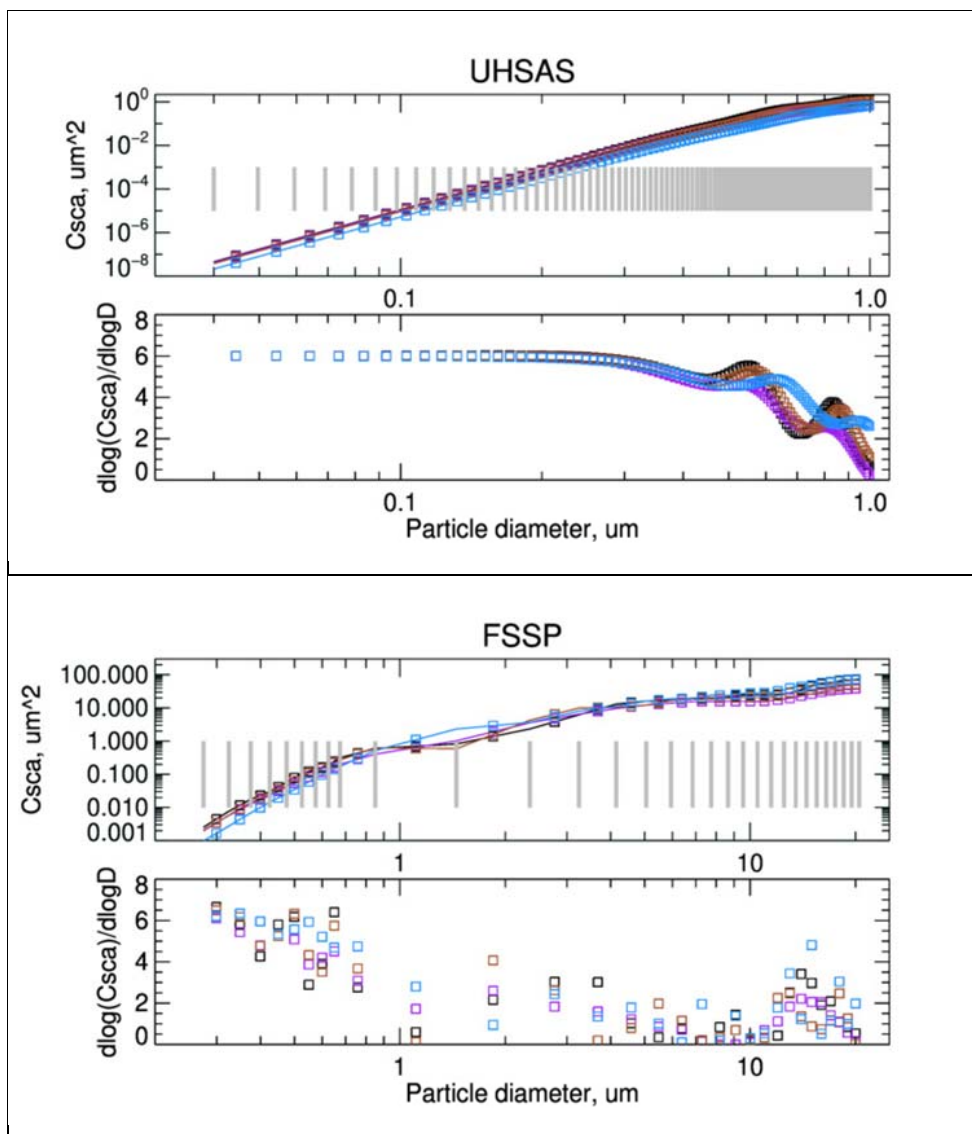


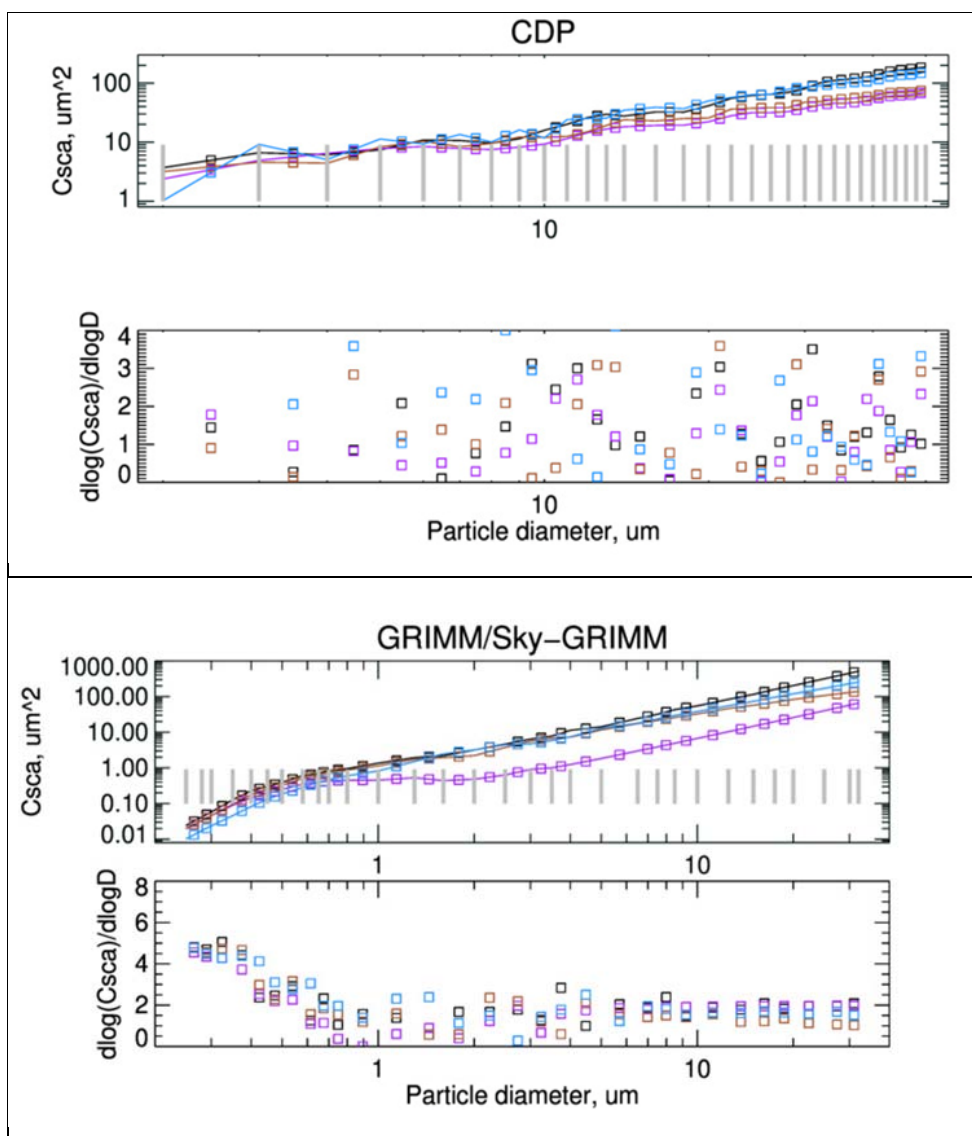
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736 **Figure 2.** (Top in each group) Scattering cross sections, C_{sca} , as a function of particle diameter
737 for the OPCs considered in this paper. The black lines represent C_{sca} for the calibration
738 particles (PSL and equivalent light-scattering material). The purple lines represent the C_{sca}
739 function for absorbing urban aerosols ($CRI = 1.56 - 0.087i$), while the brown lines represent
740 C_{sca} for moderately-absorbing mineral dust ($CRI = 1.53 - 0.003i$), and the light blue lines
741 represent C_{sca} or marine aerosols ($CRI = 1.38 - 0.001i$). The symbols represent the average
742 C_{sca} over the nominal size bins of each OPC, whose boundaries are illustrated by the vertical
743 thick grey lines.
744 (Bottom in each group) Absolute values of the first derivative in log-space of C_{sca} versus the
745 bin width $d\log D$, which is a measure of the sensitivity of the instrument to particle size.



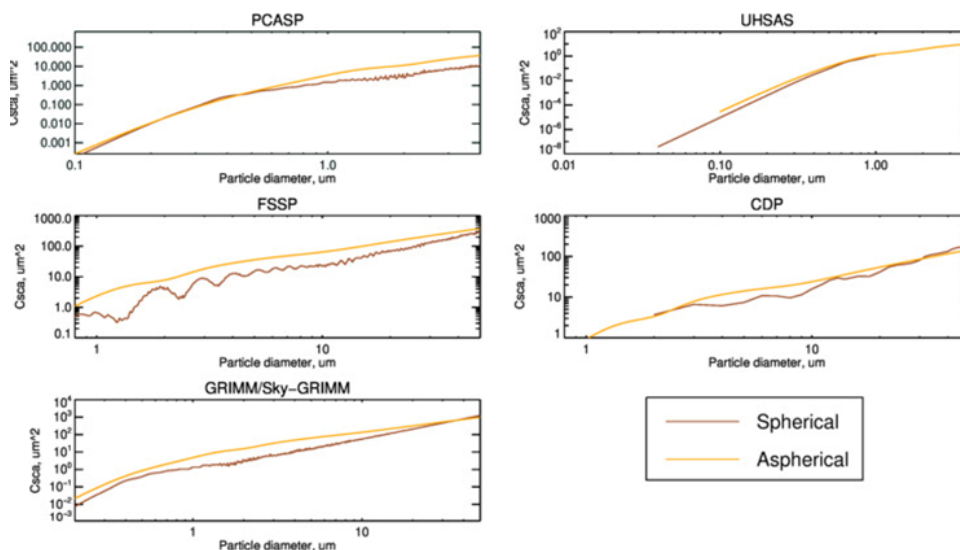




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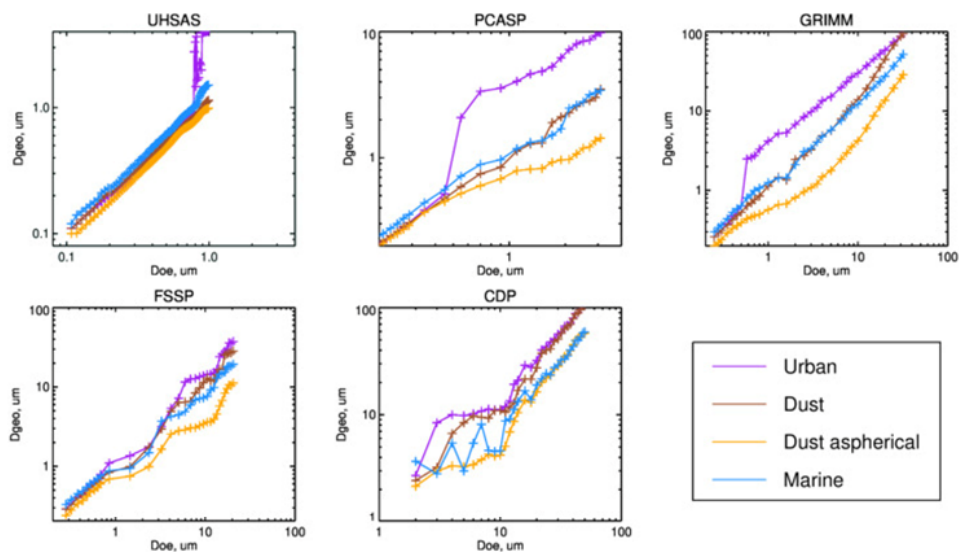
747 **Figure 3.** Scattering cross sections, C_{sca} , as a function of particle diameter for the OPCs
748 considered in this paper for moderately-absorbing mineral dust ($CRI = 1.53 - 0.003i$). The
749 brown line represents C_{sca} calculated by Mie theory assuming homogeneous spherical
750 particles, while the orange line represents C_{sca} calculated according to Huang et al. (2021)
751 assuming homogeneous aspherical ellipsoid particles.



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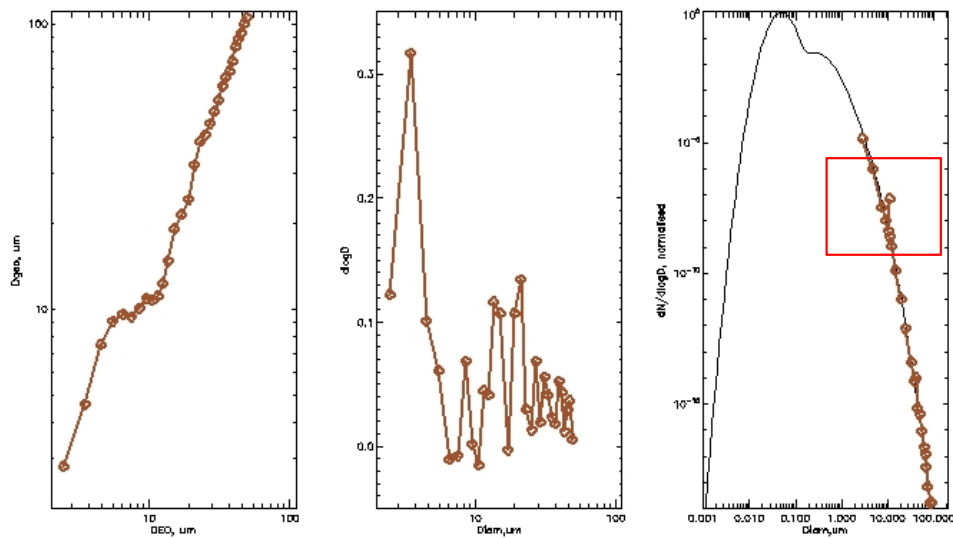
756 **Figure 4.** Scatterplot of geometric–equivalent bin size boundary (D_{geo}) corresponding to the
757 CRI of the example aerosol types (mineral dust, urban and marine) compared to the
758 optical–equivalent bin size boundary (D_{EO}) obtained if the calibration CRI is used. For mineral
759 dust, values are shown for both the aspherical and spherical approximation (brown and yellow
760 lines, respectively).



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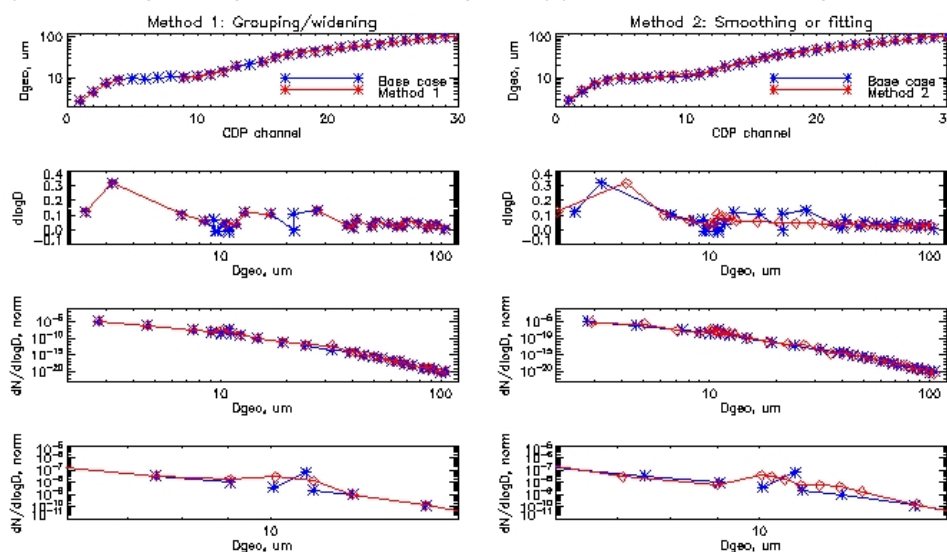
763 **Figure 5.** Representation of a model log-normal size distribution of mineral dust (see Seinfeld
764 and Pandis, 2006) for the CDP, taking into account the corrected bin diameter. The first panel
765 (left) shows the relationship between the corrected and the calibration bin diameters for a CRF
766 of $1.53-0.003i$. The middle panel shows the behaviour of the corrected $d\log D$ with calibration
767 bin diameter, while the right panel shows the representation of the size distribution using the
768 corrected CDP evaluation of particle size. The red box indicates the discontinuity of the size
769 distribution around a diameter of $10\ \mu\text{m}$.



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772 **Figure 6.** Correction methods in the CDP measurements for the caes of CRI of mineral dust
773 ($1.53 - 0.003i$). The base case is based on simply discarding the values of $dN/d\log D$
774 corresponding to negative values of $d\log D$. The left panels (from top to bottom) correspond to
775 the application of Method 1, that is, grouping instances where $d\log D$ was negative to create a
776 wider bin to which the sum of the particle counts of the original bins was attributed. The right
777 panels correspond to the application of Method 2, that is, fitting of the D_{bin} curves with a
778 second-degree polynomial function adapted to three size ranges (below $2.43 \mu\text{m}$, in between
779 2.43 and $16.85 \mu\text{m}$ and above $16.85 \mu\text{m}$, see Annex 1 for fit parameters). Panels show, from
780 top to bottom: (1) the size range of application for all the CDP channels; (2) the resulting $d\log D$;
781 (3) the resulting $dN/d\log D$ for the full size range and (4) for a restricted size range.



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