



1

2 **The S/Z Relationship of Rimed Snow Particles**

3

4 Shelby Fuller ¹, Sam Marlow ¹, Samuel Haimov ¹, Matthew Burkhardt ¹, Kevin Shaffer ¹, Austin
5 Morgan ¹, and Jefferson R. Snider ^{1,2}

6

7 ¹ Department of Atmospheric Science, University of Wyoming, Laramie, WY

8

9

10

11 ² Corresponding Author, jsnider@uwyo.edu

12

13



14 **Abstract**

15 Values of liquid-equivalent snowfall rate (S) at a ground site, and microwave reflectivity
16 (Z) retrieved above the ground site using an airborne W-band radar, were acquired during
17 overflights. Temperature at the ground site was between -6 and -15 °C. At flight level, within
18 clouds containing ice and supercooled liquid water, the temperature was approximately 7 °C
19 colder. Additionally, airborne measurements of snow particle imagery were acquired. The
20 images demonstrate that most of the snow particles were rimed. The S/Z pairs are generally
21 consistent with a published S/Z relationship. The latter was developed with airborne
22 measurements of snow particle imagery, which were used to calculate S , and coincident airborne
23 W-band radar measurements, for Z . Both the previous work and this contribution indicate that
24 most S/Z relationships developed for W-band radars underestimate S in situations with rimed
25 snow particles and with $Z < 1 \text{ mm}^6 \text{ m}^{-3}$.



26 **1 - Introduction**

27 Improvement of methods used to measure snowfall and rainfall are an ongoing focus of
28 meteorological research. The various methods are ground-based instruments that evaluate the
29 mass of precipitation that falls into or onto a collector (precipitation gauges) (Brock and
30 Richardson 2001), ground-based radars (Wilson and Brandes 1979), and airborne and space-
31 borne radars (Matrosov 2007; Kulie and Bennartz 2009; Geerts et al. 2010; Skofronick-Jackson
32 et al. 2017). An objective of these approaches, whether used to make observations independent
33 of other methods (e.g., Kulie and Bennartz 2009), or as a component of multiple observations
34 (e.g., Cocks et al. 2016), is estimation of precipitation rate and accumulation.

35 Many studies have investigated using radar for evaluating rainfall (for a review see
36 Wilson and Brandes 1979). There are two approaches. The first is research, both observational
37 and computational, that probes the relationship between rainfall rate (R) and radar-measured
38 values of backscattered microwave power. The latter is commonly reported as an equivalent
39 radar reflectivity factor (Z_e). The second is operational in the sense that precipitation gauges are
40 used to calibrate measurements acquired using weather surveillance radars. Complications
41 associated with converting Z_e to R , or converting a radar reflectivity factor¹ (Z) to R , can be
42 grouped in four categories: 1) Inaccuracy in quantification of Z , 2) variation of the R/Z
43 relationship stemming from precipitation processes (e.g., evaporation, coalescence, and break
44 up), 3) difference between the volume of a radar range gate versus the much smaller volume of
45 atmosphere sampled as precipitation falls to a gauge, and 4) vertical displacement between a
46 radar range gate and a calibrating gauge, especially at far ranges.

¹ Radars are calibrated to report Z_e (Smith 1984). Herein, radar reflectivities are reported as $Z = Z_e$ and as $\text{dBZ} = 10\log_{10}(Z_e)$.



47 For situations with snowfall, methods employing either gauge or radar are associated
48 with complications beyond that incurred in rainfall (Matrosov 2007; Martinaitis et al. 2015;
49 Cocks et al. 2016). Problems associated with gauge measurements are wind-induced snow
50 particle undercatch, gauge capping, delayed registration, and blowing snow aliasing as snowfall.
51 Moreover, in a situation with snow particles most abundant within a radar range gate, compared
52 to rain drops, and where a measurement of Z is used to infer R via a R/Z relationship, the
53 resultant precipitation rate will likely be inaccurate. This is because hydrometeor shape, density,
54 and dielectric properties are all variable for snow particles while relatively invariant for rain
55 drops. Additionally, a snow particle's terminal fall speed varies with size (as is the case for
56 drops) and with particle shape and particle density. Going forward, we refer to the latter two
57 properties as shape and density.

58 In calculations of Z and liquid-equivalent snowfall rate (S), obtained for the operating
59 wavelength of the nadir-looking radar carried on the CloudSat satellite (wavelength $\lambda = 3.2$ mm),
60 density is an important parameter. In these calculations, density is commonly estimated using
61 empirical data (Matrosov 2007; Kulie and Bennartz 2009; Pokharel and Vali 2011, [PV11]). For
62 graupel, a snow particle that grows via collection of supercooled cloud droplets in a process
63 commonly referred to as riming, paired observations of particle mass and particle size have been
64 used to estimate density. There is considerable uncertainty in this approach. Based on data
65 collected at two northwestern US surface sites (Zikmunda and Vali 1972; Locatelli and Hobbs
66 1974), density values differ by at least a factor of two at particle sizes smaller than $2000 \mu\text{m}$
67 (PV11; their Figure 4). Given that the density of rime ice varies with droplet impact speed,
68 droplet size, and temperature (Macklin 1962), it is not surprising that the density-versus-size
69 relationships analyzed by PV11 are so varied.



70 Our work analyzes values of Z acquired using an airborne radar that operates in the W-
71 band ($\lambda = 3.2$ mm). In satellite, airborne, and ground installations, W-band radars are used to
72 retrieve Z and the latter is converted to S using a S/Z relationship. W-band S/Z relationships are
73 developed in Matrosov (2007), Kulie and Bennartz (2009), Geerts et al. (2010), and PV11. This
74 contribution attempts to refine estimates of S based on W-band radar observations and
75 particularly where the dominant particle type is either rimed crystals or graupel.

76 The following introductory paragraphs overview W-band S/Z relationships being applied
77 in instances of snowfall where mass is acquired by vapor deposition (crystal), by collection of
78 crystals (aggregate), and by riming (rimed crystal and graupel). Henceforth, the latter two snow
79 particle types are collectively referred to as rimed snow particles.

80 In a computational study, Matrosov (2007) reported an upper-limit and a lower-limit S/Z
81 relationship for both the crystal and aggregate particle types. Both were modeled as oblate
82 spheroids. The upper-limit and a lower-limit relationships are $S = 0.11 \cdot Z^{1.25}$ and $S = 0.041 \cdot Z^{1.25}$.
83 At any value of Z these differ by a factor of 2.7. This variance stems from changes in density,
84 shape, and fall speed as these changes are propagated through cloud-microphysical and
85 microwave-scattering computations. Similar analyses were conducted by Kulie and Bennartz
86 (2009). Both Matrosov (2007) and Kulie and Bennartz (2009) state that the S/Z relationships
87 they recommend should be applied cautiously in settings where rimed snow particles dominate.

88 In a hybrid approach (computational and an analysis of airborne observations), PV11
89 concluded that most of the snow particles they imaged were rimed snow particles. Their
90 calculations of S and Z , conducted using two density-size relationships (indicated with ρ_1 and
91 ρ_3), were also presented. They compared their calculated reflectivities to measurements of Z
92 from a W-band radar. That led to their conclusion that “...the lower density assumption...yielded



93 closer correspondence to observed reflectivities.” Their recommendation for S as a function of Z
94 - hereafter the $S(\rho_1)/Z$ relationship - is $S = 0.39 \cdot Z^{0.58}$. In addition to variance in their values of S ,
95 coming from a dependence on density, PV11 state that a value of S derived via their $S(\rho_1)/Z$
96 relationship is uncertain by a factor-of-ten. This uncertainty is evident in the variance seen about
97 the $S(\rho_1)/Z$ relationship in Fig. 11 of PV11. Those investigators, and Geerts et al. (2010),
98 attributed the variance to use of two-dimensional snow particle imagery in calculations of S and
99 to actual variations of density and shape not accounted for in the calculations. Error associated
100 with the radar-retrieved reflectivities reported in PV11 contributed only marginally to the factor-
101 of-ten uncertainty.

102 Our focus is the W-band S/Z relationship for rimed snow particles. Section 2 describes
103 the setting of our study, the instruments we deployed, and recordings we obtained using two data
104 acquisition systems. One of the data systems was operated at a ground site and the other on an
105 aircraft. Section 3 is an analysis of the recordings; this section also considers recordings from
106 two additional, but ancillary, ground sites. Our findings are discussed in Sect. 4 and summarized
107 in Sect. 5. An Appendix (Sect. 6) explains how we averaged recordings of near-surface W-band
108 reflectivities and surface-based recordings of snowfall.



109 **2 - Site, Aircraft, and Instruments**

110 **2.1 - Site**

111 We analyzed aircraft and ground data from 14/15 December 2016, when the analyzed
112 snowfall event spanned a UTC date change, and from 3 January 2017. The ground data were
113 acquired in a forest/prairie ecotone on the eastern slope of the Medicine Bow Mountains in
114 southeastern Wyoming (Figs. 1a-b). No ground-based observers were deployed during the two
115 snowfall events we analyzed.

116 At one of three ground sites (HP in Figs. 1a-b) we deployed a hotplate precipitation
117 gauge (Rasmussen et al. 2011; Zelasko et al. 2018), a GPS receiver, and a data acquisition
118 system. Once per second, the data system ingested a hotplate-generated data string, combined
119 that with time-of-day from the GPS receiver (Coordinated Universal Time (UTC)), and recorded
120 the merged hotplate/UTC data string. The absolute accuracy of the GPS time stamp is no worse
121 than 2 s.

122 Overflights of the hotplate were done by the University of Wyoming King Air (WKA) on
123 14/15 December 2016 and on 3 January 2017. Data acquisition on the WKA was also
124 synchronized with UTC, but with much better accuracy than at the hotplate. Measurements of
125 wind (speed and direction), temperature, and relative humidity from the US-GLE AmeriFlux
126 tower (AF in Figs. 1a-b) are also components of our analysis. The AmeriFlux data were provided
127 to us as 30-minute averages (AmeriFlux 2021; Marlow et al. 2022).

128



129
130
131
132
133
134
135
136
137
138
139
140
141
142
143
144
145
146
147
148
149
150

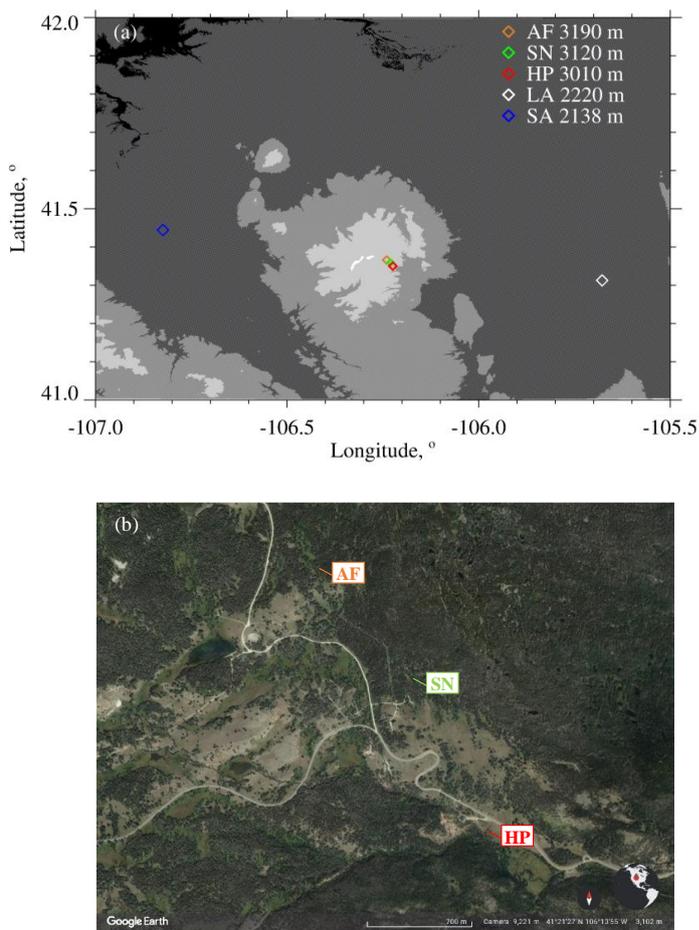


Figure 1 – (a) Southeast Wyoming, airports near the communities of Saratoga, WY (SA) and Laramie, WY (LA), and the ground sites: AF = US-GLE AmeriFlux tower, SN = Brooklyn Lake SNOTEL, and HP = hotplate. Altitudes of the airports and ground sites are in the legend. Altitude thresholds for the digital elevation map are 1500, 2000, 2500, 3000, and 3500 meters. (b) Close up of the AF, SN, and HP ground sites (from © Google Earth).



151
152 **2.2 - University of Wyoming King Air (WKA)**

153 We analyzed the following WKA measurements: aircraft position, ambient temperature,
154 snow particle imagery, and three moments of the cloud droplet size distribution. A
155 Cloud Droplet Probe (CDP; Faber et al. 2018) was the basis for the droplet size distribution
156 measurements and the derived moments. The latter are droplet concentration (N), cloud liquid
157 water content (LWC), and mean droplet diameter ($\langle D \rangle$). Snow particle imagery was obtained
158 using a precipitation particle imaging probe (2DP; Korolev et al. 2011) and a cloud particle
159 imaging probe (2DS; Lawson et al. 2006). These acquired two-dimensional images of particles
160 between 200 to 6400 μm (2DP) and between 10 to 1280 μm (2DS).

161 **2.3 - Wyoming Cloud Radar (WCR)**

162 Measurements from the Wyoming Cloud Radar (WCR), operated on the WKA, were also
163 analyzed. We analyzed values of the vertical-component snow particle Doppler velocity
164 retrieved from below the WKA using the WCR's down-looking antenna. Our starting point for
165 that analysis is the Level 2 WCR data which has snow particle Doppler velocities corrected for
166 aircraft motion (Haimov and Rodi 2013). We use V_D to symbolize the corrected vertical-
167 component Doppler velocity and adopt the convention that $V_D > 0$ indicates upward snow
168 particle motion. Level 2 values of Z retrieved using the up-looking and down-looking WCR
169 antennas were also analyzed.

170 The Level 2 WCR sampling was different on the two flight days and this difference is
171 indicated in Table 1. The flights were conducted in preparation for the SNOWIE field project
172 (Tessendorf et al. 2019) and were flown from the Laramie, WY Airport ("LA" in Fig. 1a).

173



174 Table 1 – Level 2 WCR sampling and the WKA overflight time

Date	Level 2 WCR Vertical Sampling, m	Level 2 WCR Along-track Sampling, s	Overflight Time, UTC
14/15 December 2016	23	0.23	00:00:38 (15 December 2016)
3 January 2017	30	0.36	20:32:03

175
176



177 **2.4 - Hotplate Gauge**

178 Algorithms used to process hotplate measurements are described in Rasmussen et al.
179 (2011), Boudala et al. (2014), and Zelasko et al. (2018). Henceforth, these are referred to as R11,
180 B14, and Z18, respectively. In this section, we describe how we analyzed hotplate measurements
181 acquired at the HP site.

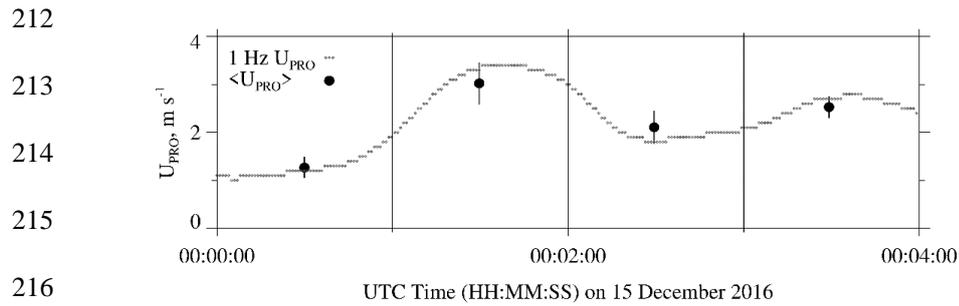
182 Four measurements fundamental to the steady state energy budget of the hotplate's
183 temperature-controlled up-viewing plate are output by the hotplate microprocessor as one-minute
184 running averages (Z18). These running averages were merged with the GPS time and recorded at
185 1 Hz by the data acquisition system (Sect. 2.1). The four measurements are electrical power
186 supplied to the plate, ambient temperature, wind speed, and solar irradiance. With these
187 measurements, calibration data (Marlow et al. 2022), and the algorithm described in Z18, we
188 calculated the liquid-equivalent snowfall rate. The latter is not corrected for the snow particle
189 undercatch.

190 Marlow et al. (2022; their Figure 4b) report the relationship between snow particle catch
191 efficiency and wind speed that we applied in calculations of the undercatch-corrected liquid-
192 equivalent snowfall rate (S , mm h^{-1}). There are three bases for this relationship. The first is the
193 catch efficiencies R11 derived from measurements they obtained using a weighing gauge
194 operated within a double fence intercomparison reference shield and a hotplate gauge. R11
195 plotted their hotplate catch efficiencies versus wind speeds measured at 10 m AGL (their Figure
196 8). The second is Marlow et al.'s adjustment of R11's 10 m AGL wind speeds to 2 m AGL. The
197 basis for that adjustment is surface boundary layer parameters derived for R11's site
198 (Kochendorfer et al. 2018) and Panofsky and Dutton (1984; their Eq. 6.7). The adjustment was
199 made because the hotplate-reported wind speeds, both here and in Marlow et al. (2022), were



200 acquired at approximately 2 m above the snowpack surface. The third is a validation of the
201 Marlow et al. (2022) undercatch relationship. That was done by comparing values of S from the
202 hotplate gauge and a SNOTEL pillow gauge (Serreze et al. 1999). In that validation (Marlow et
203 al. 2022; their Figure 10a), the SNOTEL pillow gauge was at the SN site and the hotplate was at
204 the HP site. The SN and HP sites are in Figs. 1a-b.

205 We also analyzed values of wind speed output by the hotplate microprocessor (U_{PRO}).
206 The basis for this hotplate-derived wind speed is a steady state energy budget of the hotplate's
207 temperature-controlled down-viewing plate and a proprietary algorithm (R11 and Z18). The
208 U_{PRO} were reported by the hotplate as one-minute running averages and we recorded these at 1
209 Hz. Examples are the gray dots in Fig. 2. Additionally, we calculated and analyzed one-minute-
210 averaged values of U_{PRO} and the corresponding standard deviations. Examples of these are the
211 black circles and the short vertical line segments, respectively, in Fig. 2.



218 Figure 2 – Hotplate wind speed measurements (U_{PRO}) 00:00:00 to 00:04:00 on 15 December
219 2016. Gray dots are the one-minute running-average U_{PRO} recorded at 1 Hz. Black circles are the
220 one-minute-averaged U_{PRO} (± 1 standard deviation).
221



222 **3 - Analysis**

223 **3.1 - WKA Overflight Time**

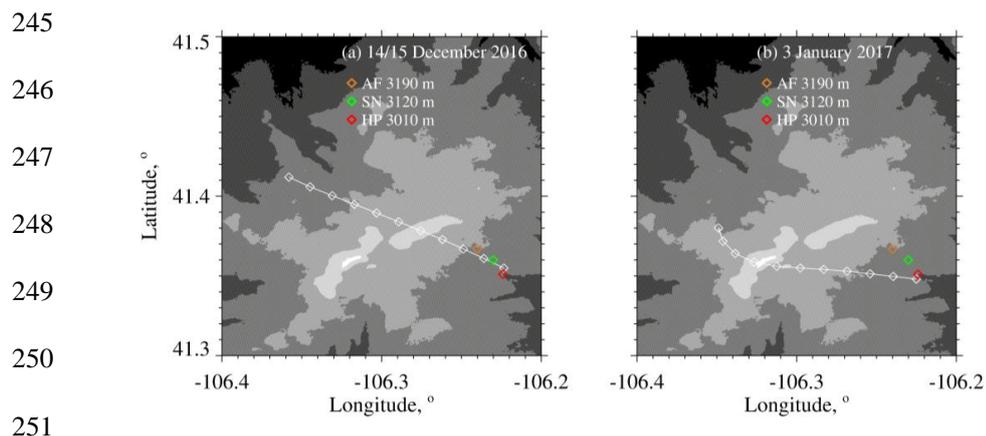
224 The focus of our analysis is the two WKA flight segments shown in Figs. 3a-b. The maps
225 shown in the figures have the three ground sites (AF, SN, and HP) and the WKA flight tracks
226 (white line). The beginning-to-end time interval for the flight tracks is 100 s and these are
227 divided into ten 10-second intervals. The 10 s intervals are indicated with white diamonds. With
228 the exception of the turn evident in Fig. 3b, the flight tracks are straight and level and the track
229 direction is approximately upwind to downwind.

230 Times that the WKA was closest to the HP site were evaluated by finding the point on the
231 flight track where the horizontal position of the WKA was closest to the hotplate's coordinates.
232 These times are symbolized t_o and are referred to as overflight times. In Figs. 3a-b the
233 downwind end of the flight tracks end at the overflight time. The latitude/longitude position of
234 the aircraft was within 390 m of the hotplate at the overflight times. Table 1 has the overflight
235 times on the two flight days.

236 **3.2 - Correction of Doppler Velocity**

237 We accounted for bias in V_D (Sect. 2.3) due to deviation of the down-looking WCR
238 antenna from vertical. This was done by applying the correction described in Zaremba et al.
239 (2022) (their Eq. A4). The west-to-east and south-to-north particle velocities used in the
240 correction were assumed equal to component wind velocities. The latter were expressed as linear
241 functions of altitude using the information in the penultimate and last columns of Table 2. The
242 component velocities as functions of altitude and the linear equations relating velocity and
243 altitude are provided in the Appendix.

244



252

253 Figure 3 – (a) WKA flight track on 14/15 December 2016 for time interval = overflight time -
254 100 s to the overflight time. (b) WKA flight track on 3 January 2017 for time interval =
255 overflight time - 100 s to the overflight time. The white diamonds on the tracks are separated, in
256 time, by 10 s. Altitude thresholds for the digital elevation maps are 2600, 2800, 3000, 3200,
257 3400, and 3600 meters. Altitudes of the ground sites are in the legend.

258



259 Table 2 – Aircraft and atmospheric state averages

260

Date	WKA ^a Track Altitude, m	WKA ^a T, °C	AF ^b T, °C	AF ^b RH, %	WKA ^{a,c} Track Vector	WKA ^{a,c} Wind Vector	AF ^{b,c} Wind Vector
14/15 December 2016	4546	-13.9	-6.3	86	310 / 130	274 / 32	250 / 8.5
3 January 2017	4196	-21.7	-14.6	77	280 / 120	265 / 27	260 / 5.4

261

262

263 ^a Altitude, temperature, track vector, and horizontal wind vector data obtained by averaging 1 Hz WKA
264 measurements. The averaging interval is 60 s and the interval starts at the overflight time, minus 60 s, and
265 ends at the overflight time.

266

267 ^b Temperature, relative humidity, and horizontal wind vector data from sensors on the US-GLE
268 AmeriFlux tower (Sect. 2.1). The wind sensor was deployed at 26 m AGL (3223 m MSL) and the T/RH
269 sensor was deployed at 23 m AGL (3220 m MSL). The AF measurements correspond to 30-minute
270 averages closest to the overpass time. In the AF data set time stamps on the relevant AF recordings are
271 00:00 UTC (15 December 2016) and 20:30 UTC (3 January 2017).

272

273 ^c Vectors are presented in the following format: Direction of motion (degree relative to true north) / speed
274 (m s⁻¹).

275



276
277 **3.3 - Hotplate Measurement of Wind Speed**

278 Here we compare the hotplate-derived wind speed - symbolized U_{PRO} (Z18) - to wind
279 speed derived using an R.M.Young rotating anemometer (R.M.Young 2001). The second of
280 these two speeds we symbolize U_{RMY} and we note that the basis for the first (U_{PRO}) is a
281 proprietary algorithm (Sect. 2.4). We are doing this comparison because B14 showed that U_{PRO}
282 can be high-biased, relative to a conventional anemometer, and because U_{PRO} is the primary
283 determinant of the rate that the up-viewing plate dissipates sensible heat energy. Diagnosis of
284 that heat transfer rate is our basis for calculating the liquid-equivalent snowfall rate (Z18). The
285 U_{PRO} also determines the snow particle catch efficiency. The latter is our basis for calculating the
286 undercatch-corrected liquid-equivalent snowfall rate (Sect. 2.4).

287 Three years before the wind speed comparison presented here, we attempted to compare
288 the U_{PRO} reported by our hotplate² and wind speed reported by a WXT520 Vaisala weather
289 transmitter equipped with an ultrasonic anemometer (Vaisala 2012). These instruments were
290 operated at the HP site in Fig. 1b. However, that data set was difficult to interpret because we did
291 not correctly record the desired 1 Hz wind speed measurements from the WXT520. The
292 comparison reported here was done at the Laramie, WY Airport in December 2019 and January
293 2020. Compared to the HP site, the Laramie Airport site (indicated LA in Fig. 1) is free of
294 obstruction, out to 120 m, and experiences larger wind speeds. By mounting the hotplate and the
295 R.M.Young anemometer on rigid metal pipes, the hotplate's heated horizontal surfaces (the up-
296 and down-viewing plates seen in Figure 1 of Z18) and the anemometer's spinning axis (oriented
297 horizontally) were both positioned at 2 m AGL. The pipes were separated horizontally by 5 m.

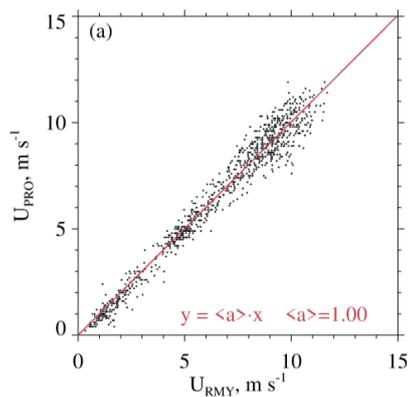
² The hotplate used here is the device described in Wolfe and Snider (2012), in Z18, and in Marlow et al. (2022).



298 There was no precipitation on the days selected for the wind speed comparisons. The values of
299 U_{PRO} and U_{RMY} we analyzed were recorded with a data system that time stamped the 1 Hz U_{PRO}
300 and 1 Hz U_{RMY} with a relative timing accuracy no worse than 1 s.

301 A wind speed comparison - from 13 December 2019 - is shown in Fig. 4a. U_{PRO} was
302 brought into the comparison by sampling it once per minute from files containing 1 Hz
303 recordings of the one-minute running-average U_{PRO} (Sect. 2.4). U_{RMY} was brought into the
304 comparison by starting with files containing 1 Hz recordings and converting these to one-minute
305 averages. Fig. 4a shows no evidence of bias and Fig. 4b demonstrates that the average absolute
306 departure between the U_{PRO} and U_{RMY} (both one-minute averages) is no larger than 0.5 m s^{-1} .
307 Table 3 has eight more precipitation-free comparisons. Included in the table are temperature and
308 wind speed averaged over the comparison intervals (4 to 20 UTC), the slope of the linear-least-
309 squares fit line (forced through the origin), and the lower and upper quartiles of the slope. We
310 calculated the quartiles using the method of Wolfe and Snider (2012). In contrast to Figs. 4a-b,
311 Figs. 4c-d make the comparison using 1 Hz values of U_{PRO} and U_{RMY} . The larger scatter and
312 larger average absolute departure seen in these panels is a consequence of the hotplate's limited
313 time response, compared to the R.M.Young. We quantify the hotplate's response time in terms of
314 a thermal response time. During wintertime at the Laramie Airport, and with wind speed at 5 m
315 s^{-1} , the down-viewing plate's thermal response time is 60 s (results not shown). Because the
316 temperature of down-viewing plate is actively controlled, this does not translate to a 60 s lag
317 between changes in wind speed and the hotplate response. The $U_{\text{PRO}}/U_{\text{RMY}}$ departure is most
318 evident at $U_{\text{PRO}} > 5 \text{ m s}^{-1}$ (Fig. 4d) but this is not a concern for U_{PRO} on 14/15 December 2016 or
319 on 3 January 2017. As we show below, the U_{PRO} was less than 5 m s^{-1} at the hotplate during the
320 two WKA overflights.

321



322

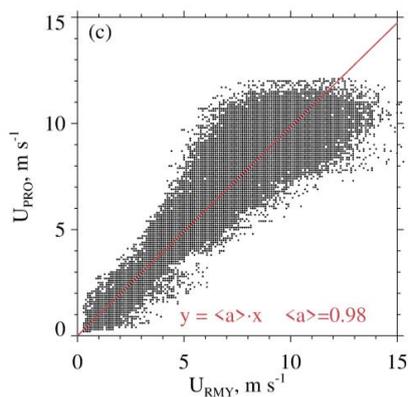
323

324

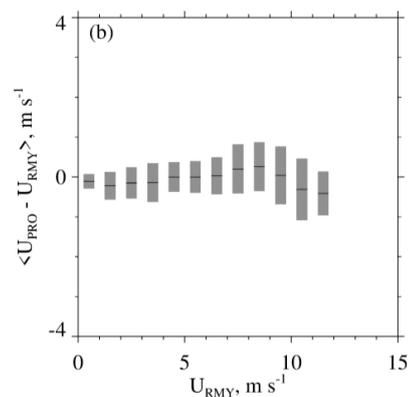
325

326

327



328



329

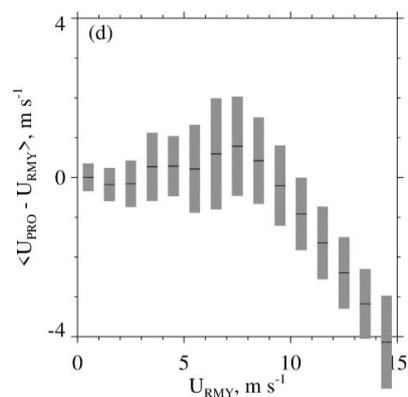
330

331

332

333

334



335

336

337 Figure 4 – (a) Scatterplot of one-minute-averaged U_{PRO} and one-minute-averaged U_{RMY} .

338 Measurements were acquired at the Laramie, WY Airport 13 December 2019. The red line is a

339 linear-least-squares fit line (forced through the origin). (b) Average departure between one-

340 minute-averaged U_{PRO} and one-minute-averaged U_{RMY} . Average departures are computed for

341 discrete U_{RMY} intervals and the averages are indicated with short black horizontal lines. Gray

342 bars indicate ± 1 standard deviation. (c) Same as in (a) except 1 Hz values of U_{PRO} and U_{RMY} . (d)

343 Same as in (b) except for 1 Hz values of U_{PRO} and U_{RMY} .



344
345 Table 3 - U_{PRO} versus U_{RMY} correlations

Date, UTC ¹	$\langle T \rangle^2$, °C	$\langle U \rangle^2$, m s ⁻¹	$\langle a \rangle^3$	a ⁴ First Quartile	a ⁴ Third Quartile
7 December 2019	-0.40	5.40	1.00	0.90	1.04
8 December 2019	2.70	4.10	0.99	0.90	1.04
10 December 2019	-5.20	3.80	0.99	0.83	1.04
13 December 2019	-1.50	6.60	1.00	0.93	1.06
18 December 2019	-6.20	3.60	0.99	0.92	1.04
19 December 2019	-6.90	2.70	0.95	0.84	0.99
6 January 2020	-6.40	8.80	1.01	0.96	1.06
8 January 2020	0.30	4.20	1.00	0.87	1.05
11 January 2020	-7.20	7.00	1.02	0.97	1.08

346
347

348 ¹ Statistics presented are based on one-minute-averaged U_{PRO} and one-minute-averaged U_{RMY}
349 measurements made between 04:00 to 20:00 UTC.

350

351 ² Averaged temperature and interval-averaged wind speed

352

353 ³ Slope of the one-minute-averaged U_{PRO} versus one-minute-averaged U_{RMY} linear-least-squares fit line,
354 forced through the origin

355

356 ⁴ Quartiles of the slope (see text)

357

358



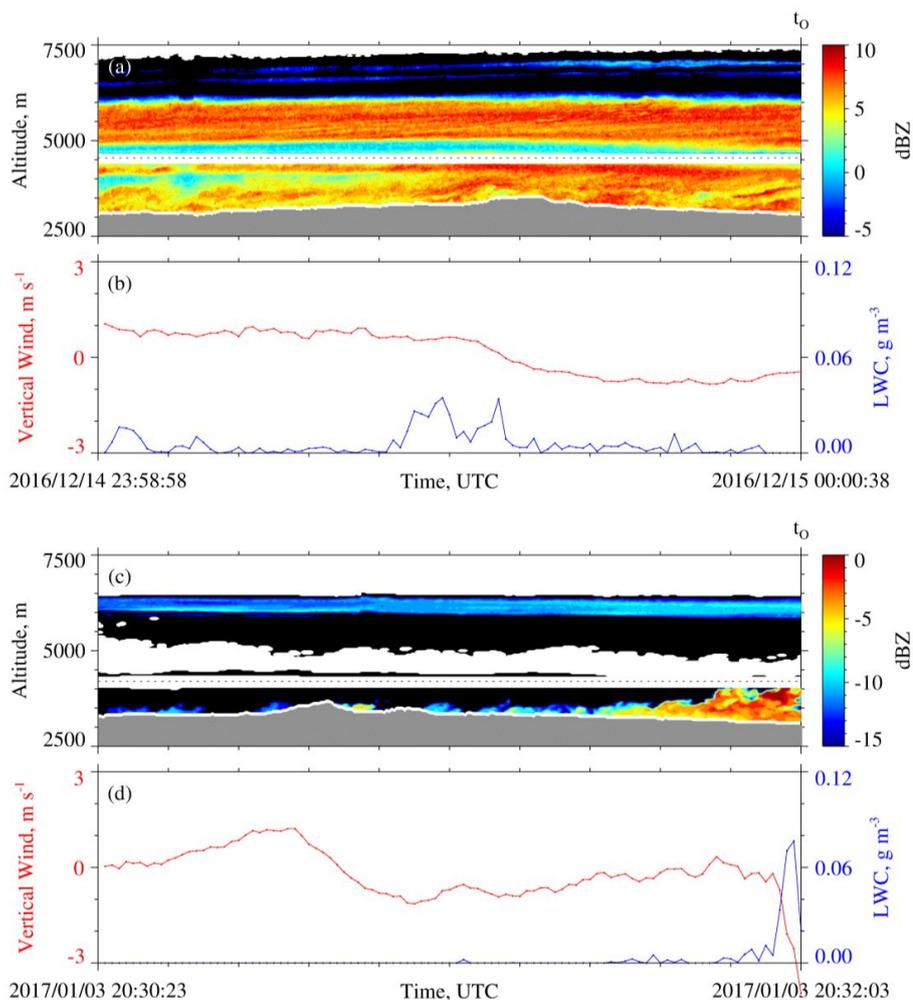
359

360 3.4 - WCR Measurements

361 Figure 5 has WCR and WKA measurements starting 100 s prior to t_o and completing at
362 t_o . The sequences in Figs. 5a and 5c are reflectivities from both the up- and down-looking
363 antennas. In Fig. 5a the flight track (black dashed horizontal line) is at 4550 m and in Fig. 5c the
364 flight track is at 4200 m. At the t_o in Fig. 5a, below the WKA, the maximum radar echo is +6
365 dBZ ($Z = 4 \text{ mm}^6 \text{ m}^{-3}$) while in Fig. 5c that maximum is -3 dBZ ($Z = 0.5 \text{ mm}^6 \text{ m}^{-3}$). Supercooled
366 liquid water was detected as the aircraft approached the ridgeline (Fig. 5b) and during the last 3
367 seconds of the time sequence in Fig. 5d. During these encounters with supercooled liquid, the
368 maximum LWC values were 0.03×10^{-3} and $0.08 \times 10^{-3} \text{ kg m}^{-3}$ on 14 December 2016 and 3
369 January 2017, respectively. Values of N (Sect. 2.2) at times of maximal LWC were 3×10^6 and
370 $100 \times 10^6 \text{ m}^{-3}$ on 14 December 2016 and 3 January 2017, respectively. Even on 3 January 2017,
371 the $\langle D \rangle$ (Sect. 2.2) associated with maximum LWC was sufficient for hexagonal plate crystals
372 with diameter larger than $100 \mu\text{m}$ to collide with the observed droplets with efficiencies > 0.1
373 (Wang and Ji 2000).

374 We spatially averaged the values of Z we compared with time-averaged values of S.
375 There are two reasons for this: 1) As discussed in Sect. 3.1, the WCR did not sample Z exactly
376 over the hotplate, and furthermore, the width of radar beam at 1500 m range - roughly the
377 distance between the aircraft and the ground at the overflight times - is 30 m and thus
378 considerably smaller than the minimum horizontal distance between the aircraft and the HP. 2)
379 Compared to the WCR, the hotplate is a relatively slow-response measurement system whose
380 output is commonly averaged over one-minute intervals (Z18).

381



401 Figure 5 – (a) 100 s of WCR reflectivity from 14/15 December 2016 ending at t_0 . (b) 100 s of
402 LWC and gust probe vertical wind velocity from 14/15 December 2016 ending at t_0 . (c) 100 s
403 of WCR reflectivity from 3 January 2017 ending at t_0 . (d) 100 s of LWC and gust probe
404 vertical wind velocity from 3 January 2017 ending at t_0 . In (a) and (c), above and below the
405 flight track is the roughly 200-m-deep WCR blind zone, reflectivity above (below) the flight
406 track is from the up-looking (down-looking) WCR antenna, black indicates reflectivity [dBZ]
407 smaller than minimum indicated in the color bar, white immediately above the terrain indicates
408 echo that was discarded because of ground clutter, and white above the ground clutter and
409 outside of the blind zone indicate dBZ < minimum detectable signal.



410
411 Our averaging scheme starts with the HP averaging intervals begun at t_o and at $t_o + 60$
412 s. The duration of these intervals is one minute. Figure 6a is a schematic of an HP averaging
413 interval started at t_o , Fig. 6b is a schematic of the corresponding WCR averaging domain, and
414 Figs. 6c-d are schematics of an adjacent averaging interval/domain. Figures 6a and 6c also show
415 that the indexes $i = 0$ and $i = 1$ are used to indicate HP averaging intervals begun at t_o and $t_o +$
416 60 s, respectively. Figures 7b and 8b show hotplate snowfall measurements from 14/15
417 December 2016 and 3 January 2017 and how we label the HP averaging intervals begun at t_o .
418 In these and subsequent figures, colored circles surround the indexes, blue is used to color-code
419 15 December 2016, and red is used to color-code 3 January 2017. The Appendix explains the
420 averaging in greater detail. Two aspects not discussed here, but are discussed in the Appendix,
421 are how the “i” indexes were used to calculate the WCR averaging start and end times and how
422 the lines defining the top of the WCR averaging domains, seen in Fig. 6b and 6d, were
423 calculated.

424 Figure 9a and Fig. 10a have enlarged views of the reflectivity structures recorded on the
425 two flight days. Different from Fig. 5a and Fig. 5c, these measurements are only from the
426 WCR’s down-looking antenna. Additional differences are the following: 1) The plots are set up
427 so that Z and V_D structures downwind of the hotplate can be seen. These structures are discussed
428 in the following section. 2) The WCR measurements are shown for 50 s of flight. With the WKA
429 ground speed approximately 125 m s^{-1} (Table 2), the distance along the abscissa is 6250 m. 3)
430 Colored circles that surround the $i = 0$ index are placed below the WCR averaging domains
431 (reflectivity and Doppler velocity). The domains are drawn with solid black lines and these are



432 seen to overlay both the Z data (Fig. 9a and Fig. 10a) and the V_D data (Fig. 9b and Fig. 10b).
433 Consistent with Fig. 6b, and the Appendix, one of these black lines is vertical and the other is
434 negatively sloped.
435

436

437

438

439

440

441

442

443

444

445

446

447

448

449

450

451

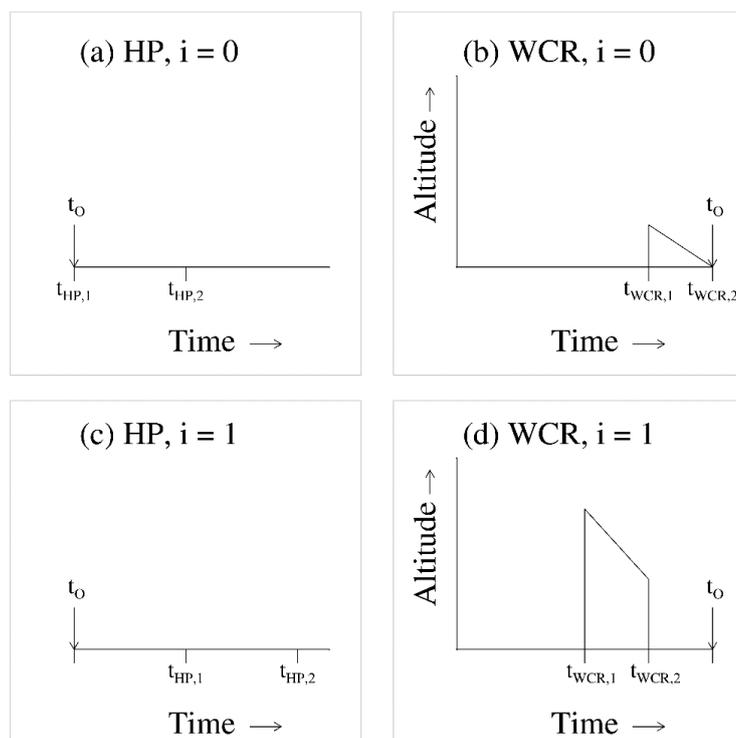
452

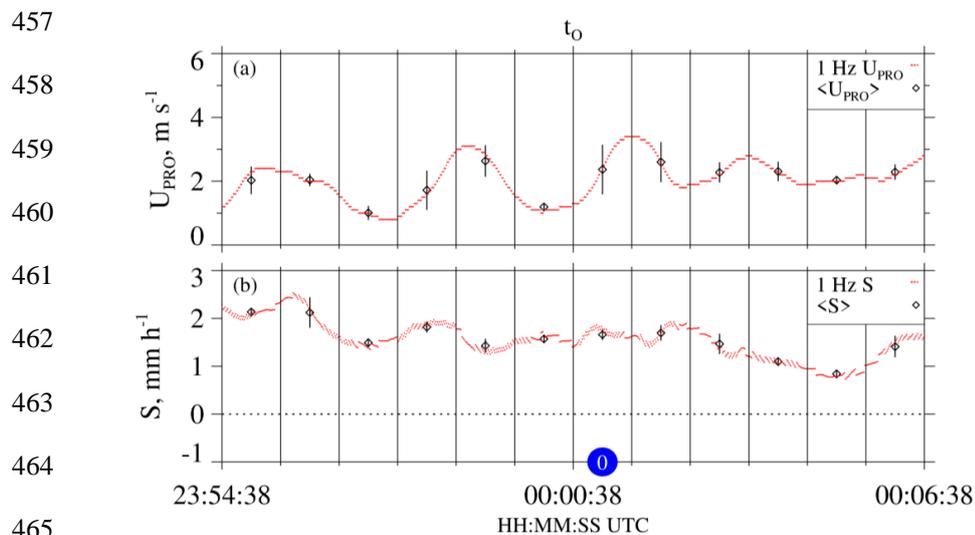
453 Figure 6 – (a and c) Schematic diagrams of the $i=0$ and $i=1$ one-minute HP averaging

454 intervals. (b and d) Schematic diagrams of the $i=0$ and $i=1$ WCR averaging domains with the

455 lowest-retrievable weather target at the low end of the ordinate. The t_o is shown in all panels.

456





457
458
459
460
461
462
463
464
465
466
467 Figure 7 – Twelve minutes of hotplate data from 14/15 December 2016. (a) Wind speed. (b)
468 Snowfall rate. In (a), red dots are the one-minute running-average U_{PRO} , recorded at 1 Hz, and in
469 (b), red dots are values of S computed using hotplate output recorded at 1 Hz. In both panels the
470 black diamonds are the one-minute-averaged values (± 1 standard deviation). The t_0 is shown
471 above the top panel and the large blue circle indicates the $i=0$ HP averaging interval.
472

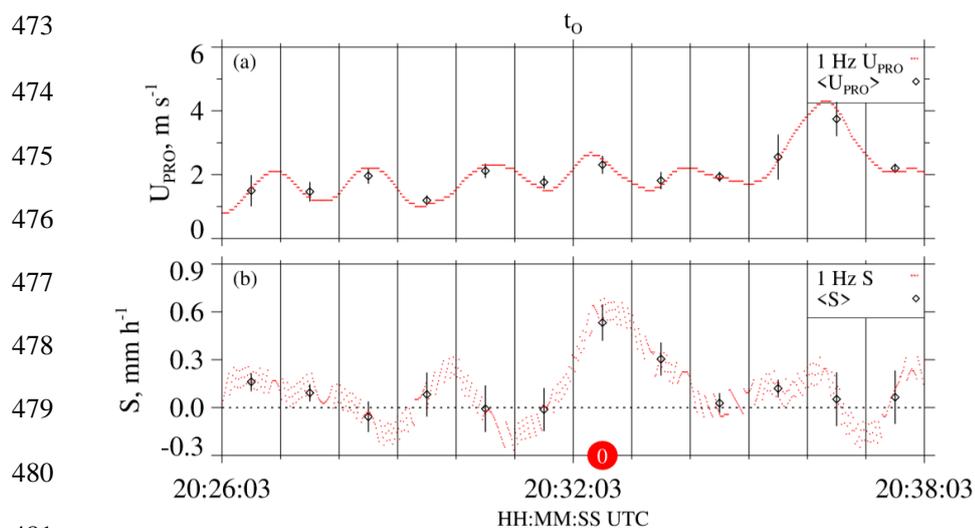


Figure 8 – Twelve minutes of hotplate data from 3 January 2017. (a) Wind speed. (b) Snowfall rate. In (a), red dots are the one-minute running-average U_{PRO} , recorded at 1 Hz, and in (b) red dots are values of S computed using hotplate output recorded at 1 Hz. In both panels the black diamonds are the one-minute-averaged values (± 1 standard deviation). The t_0 is shown above the top panel and the large red circle indicates the $i=0$ HP averaging interval.



490

491

492

493

494

495

496

497

498

499

500

501

502

503

504

505

506

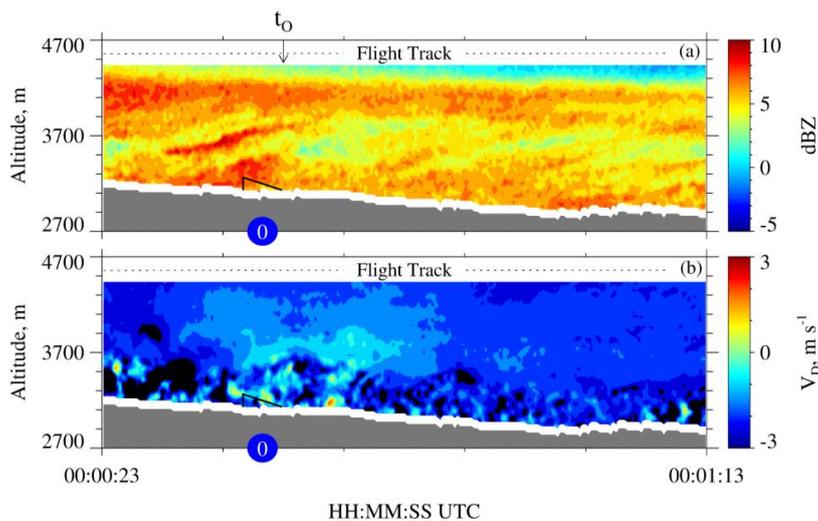


Figure 9 – 50 s of measurements from the down-looking WCR antenna on 15 December 2016.

(a) Crosssection of reflectivity $t_0 - 15$ s to $t_0 + 35$ s. (b) Crosssection of Doppler velocity $t_0 - 15$

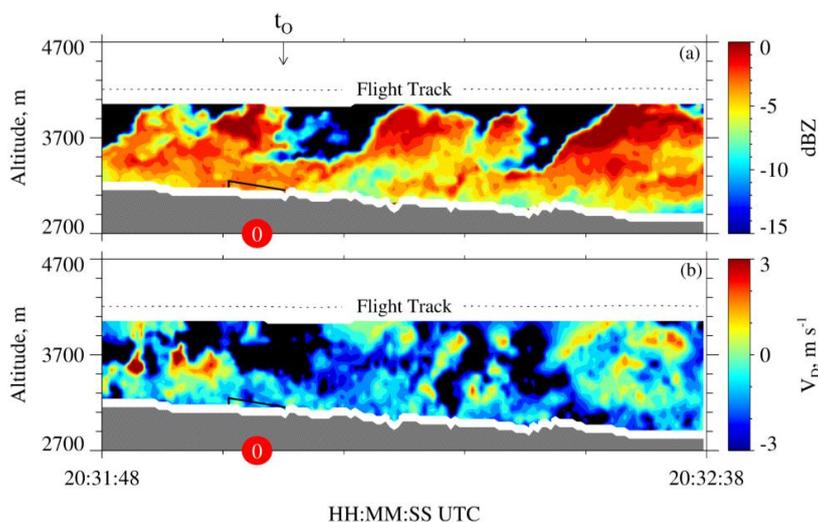
s to $t_0 + 35$ s. The t_0 is shown above the top panel. In both panels, the solid black lines (vertical

and sloped) encompass the $i = 0$ WCR averaging domain and the blue circles have the $i = 0$

index.



507
508
509
510
511
512
513
514
515
516
517



518 Figure 10 – 50 s of measurements from the down-looking WCR antenna on 3 January 2017. (a)
519 Crosssection of reflectivity $t_0 - 15$ s to $t_0 + 35$ s. (b) Crosssection of Doppler velocity $t_0 - 15$ s to
520 $t_0 + 35$ s. The t_0 is shown above the top panel. In both panels, the solid black lines (vertical and
521 sloped) encompass the $i=0$ WCR averaging domain and the red circles have the $i=0$ index.
522



523 The $i = 0$ averages of S and Z are presented in Table 4 and the corresponding averaging
524 intervals/domains are viewable in Fig. 7b and Fig. 9a (15 December 2016) and in Fig. 8b and
525 Fig. 10a (3 January 2017). The $i = 1$ averages are also presented in Table 4. According to the
526 averaging scheme (Fig. 6), the $i = 1$ HP averaging interval is time-shifted positively compared to
527 the $i = 0$ HP averaging interval and the $i = 1$ WCR averaging domain is time-shifted negatively
528 compared of the $i = 0$ WCR averaging domain. This arrangement of the averaging
529 intervals/domains is one way to average while also accounting for wind advection of the snow
530 particles.
531



532 Table 4 – Averaged wind, hotplate, and WCR measurements

Date	v_w^a , m s ⁻¹	i index	$\langle S \rangle \pm \sigma_S^b$, mm h ⁻¹	WCR Samples ^c	$\langle V_D \rangle^d$, m s ⁻¹	$\sigma_{V_D}^e$, m s ⁻¹	v_p^f , m s ⁻¹	$\langle Z \rangle \pm \sigma_Z^g$, mm ⁶ m ⁻³
15 December 2016	7.4	0	1.7±0.1	42	-1.3	0.9	2.2	4.9±2.1
15 December 2016	7.4	1	1.7±0.2	149	-1.8	1.2	3.0	5.6±1.1
3 January 2017	8.9	0	0.5±0.1	20	-0.8	0.8	1.6	0.49±0.05
3 January 2017	8.9	1	0.3±0.1	35	-0.8	0.4	1.2	0.50±0.10

533

534 ^a Horizontal wind advection speed (Eq. A5) calculated using values from the penultimate and last
 535 columns of Table 2.

536

537 ^b One-minute average of the undercatch-corrected liquid-equivalent snowfall rate (± 1 standard deviation).
 538 Example averaging intervals are the $i = 0$ intervals in Fig. 7b and Fig. 8b.

539

540 ^c Number of samples used to calculate WCR statistics in the penultimate four columns. The averaging
 541 domains (e.g., the $i = 0$ domains in Figs. 9a-b and 10a-b) encompass the averaged WCR samples.

542

543 ^d Average of Doppler velocity within the averaging domains.

544

545 ^e Standard deviation of Doppler velocity within the averaging domains.

546

547 ^f Maximum likely snow particle speed toward the surface (Eq. A6).

548

549 ^g Average reflectivity (± 1 standard deviation)

550



551 3.5 - Snow Particle Imagery

552 In Fig. 9a and Fig. 10a, the time for a snow particle to move the abscissa and ordinate
553 distances is different. The ratio of these two times is 2.6. This follows from our choice of
554 abscissa and ordinate ranges, from values of particle fall speed (1 m s^{-1}) and horizontal wind
555 advection speed (8 m s^{-1}), which we assumed, and from the WKA ground speed ($g_s \sim 125 \text{ m s}^{-1}$;
556 Table 2). The assumed values are approximately consistent with values of $\langle V_D \rangle$ and v_w , in
557 Table 4, and with the V_D sign convention (Sect. 2.3). We used $g_s = 125 \text{ m s}^{-1}$ to scale (virtually)
558 the time axes in Fig. 9a and Fig. 10a to a horizontal distance. Within the scaled coordinate
559 frames, we assumed that all snow particle trajectories have negative slope ($\Delta z / \Delta x = -1 \text{ m s}^{-1} / 8$
560 $\text{m s}^{-1} = -0.12$) and that all trajectories are stationary. However, both assumptions seem
561 inconsistent with the reflectivity structures in Fig. 5a, where positively-sloped particle fall
562 streaks are evident at $\sim 5500 \text{ m}$, inconsistent with Fig. 9a where positively-sloped fall streaks are
563 at $\sim 3500 \text{ m}$, and inconsistent with the positively-sloped fall streaks in Fig. 10a. On both flight
564 days, the fall streaks evince particle sources that move horizontally and with a horizontal speed
565 that is larger than the $v_w = 8 \text{ m s}^{-1}$ we applied in our estimate of the trajectory slope and in our
566 evaluation of the WCR averaging domains (Sect. 3.4). It may be that the source's horizontal
567 speed is comparable to the flight-level WKA-derived horizontal wind (27 to 32 m s^{-1} ; Table 2)
568 but we do not have data needed to verify that assertion. Based on our qualitative interpretation of
569 the fall streaks, and the assumption that snow particles followed the fall streaks while both were
570 advecting horizontally, we looked *downwind* of the hotplate - at a time later than t_O in Fig. 9a
571 and Fig. 10a - for particles that became those that produced snowfall at the hotplate.

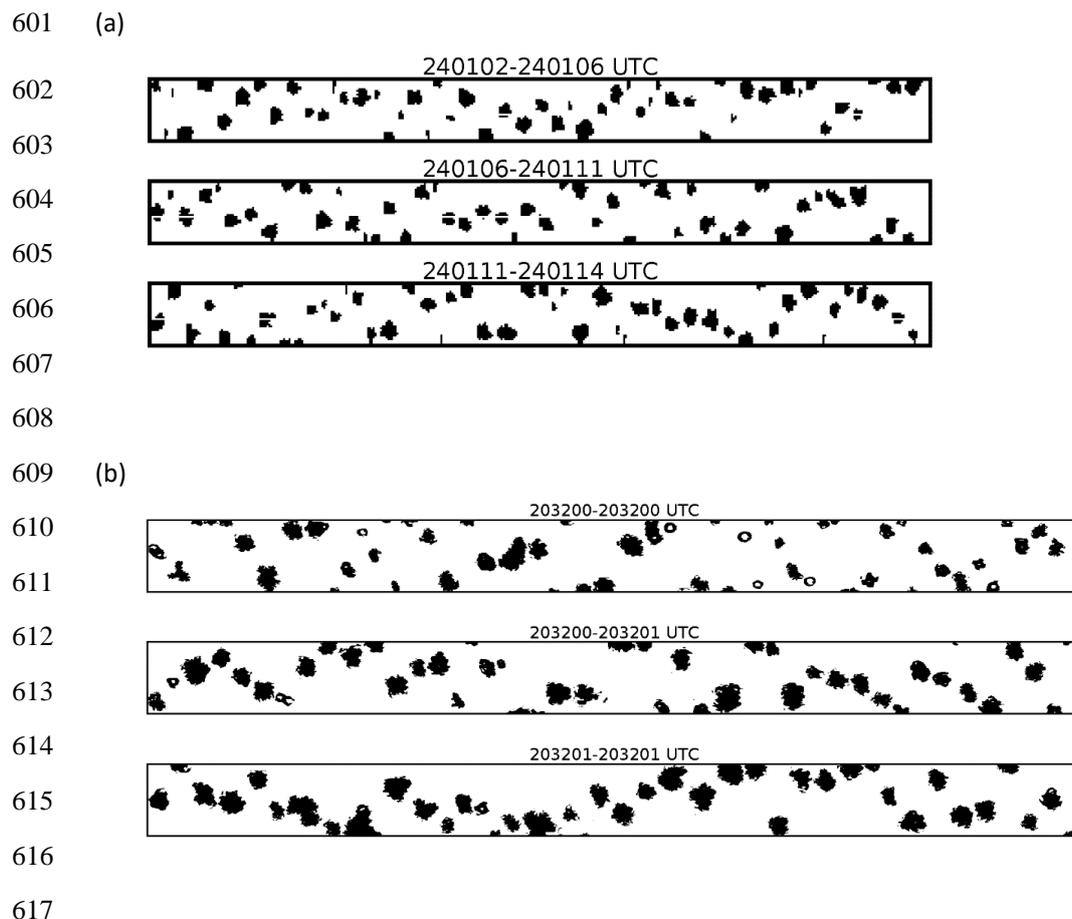


572 Particle images from 15 December 2016 were analyzed using the 2DP. With this
573 instrument the maximum all-in particle size, in the horizontal direction perpendicular to flight, is
574 6400 μm and the particle size resolution is 200 μm (Sect. 2.2). Within the time interval we
575 picked for this analysis (discussed below), particles sizing in the smaller of the two spectral
576 modes, with mode size $\sim 400 \mu\text{m}$, were more numerous (results not shown). Because the 400 μm
577 particles are poorly resolved by the 2DP, and the same can be said for somewhat larger particles,
578 those smaller than 1000 μm were excluded from the following analysis. Figure 11a shows
579 imagery from 12 s of measurements acquired near the end of the sequence in Fig. 9a (00:01:02 to
580 00:01:14). This time interval was selected by tracing backwards from t_O , along the slope of the
581 fall streaks in Fig. 9a, to the flight level. Many of the particles are rounded (indicating riming)
582 and a few have arms likely due to incomplete conversion of branched crystals to rimed snow
583 particles. The mode size corresponding to these images is 1600 μm . No liquid water was
584 detected with these particles ($\text{LWC} < 0.01 \times 10^{-3} \text{ kg m}^{-3}$; Fuller 2020; her Figure 8), but liquid was
585 detected, at $\sim 00:00:00$, as the aircraft approached the ridgeline (Sect. 3.4).

586 Turning to imagery from 3 January 2017, the most appropriate location for analysis
587 would be through the second billow structure evident in Fig. 10a. This billow sourced a fall
588 streak that terminated at the hotplate (i.e., at the time t_O indicated in the figure). However, the
589 aircraft only clipped the top of this billow, and it was only when sampling the billow seen ~ 13 s
590 earlier that larger ice particle concentrations ($\sim 20,000 \text{ m}^{-3}$) (Fuller 2020; her Figure 10) and
591 larger LWC ($> 0.06 \times 10^{-3} \text{ kg m}^{-3}$; Fig. 5d) were detected. Maximum reflectivities were the same
592 in all three billows ($Z \sim 1 \text{ mm}^6 \text{ m}^{-3}$; 0 dBZ), so we assumed that imagery collected in the first
593 billow (20:32:00 to 20:32:02) was representative of what was falling toward the hotplate. The
594 2DS was used to image these particles (Fig. 11b); with this instrument the maximum all-in



595 particle size, in the horizontal direction perpendicular to flight, is 1280 μm and the size
596 resolution is 10 μm (Sect. 2.2). Most of the objects in Fig. 11b appear to be rimed and their mode
597 size is $\sim 400 \mu\text{m}$. It is also noted that we eliminated particles smaller than 100 μm from these
598 images, however, compared to the $\sim 400 \mu\text{m}$ particles those smaller than 100 μm were
599 significantly less abundant (results not shown).
600



618 Figure 11 – (a) 2DP particle imagery from 15 December 2016. The height of the strips is 6400
619 μm . These particles are estimated to be representative of those that fell from flight level toward
620 the hotplate. (b) 2DS particle imagery from 3 January 2017. The height of the strips is 1280 μm .
621 These particles are estimated to be representative of those that fell from flight level toward the
622 hotplate.

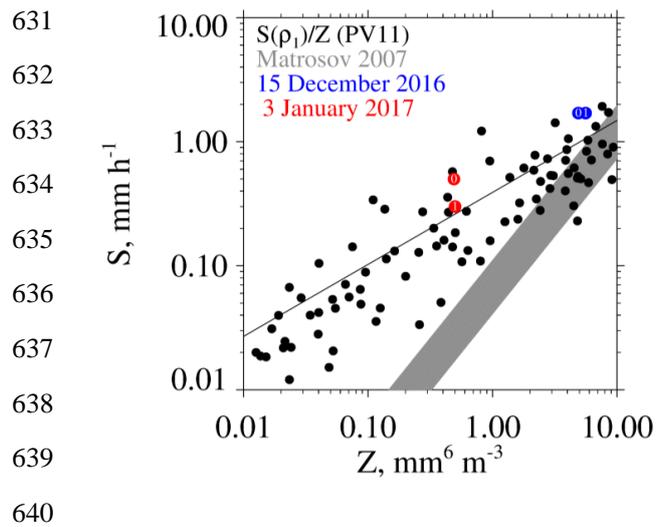
623



624 **3.6 – S/Z Relationships**

625 The S/Z pairs presented by PV11 in their Figure 11 vary by a factor-of-ten about their
626 best-fit relationship ($S(\rho_1)/Z$). Those results are shown in Fig. 12 with black circles and a black
627 line. Our S/Z pairs are presented in Table 4 and are plotted in Fig. 12 where we used the indexes
628 ($i = 0$ and $i = 1$) to indicate the averaging intervals/domains. Our data pairs plot above the
629 $S(\rho_1)/Z$ relationship but within the variability.

630



641 Figure 12 – Snowfall rate (S) versus radar reflectivity (Z). The $S(\rho_1)/Z$ data points are from
642 PV11 (their Figure 11). The $S(\rho_1)/Z$ relationship and the upper- and lower-limit Matrosov
643 (2007) S/Z relationships are presented as math functions in Sect. 1. Results for $i = 0$ and $i = 1$
644 averaging intervals/domains are shown with colored circles.



646
647 There are two potential biases in the values of S we tabulate (Table 4) and plot (Fig. 12).
648 First, the two snowfall events had flight-level vertical wind velocities that were positive
649 (upward) upwind of the summit, and vice versa downwind of the summit. Except for the
650 strongest downdraft on 3 January 2017, the magnitude of this variance is $\sim 1 \text{ m s}^{-1}$ (Figs. 5b and
651 5d). Assuming 1 m s^{-1} was the downward wind immediately over the hotplate, the snow particles
652 would have approached the gauge faster than their fall speed, and especially so on 3 January
653 2017. Our basis for stating this is fall speeds for the mode sizes discussed in Sect. 3.6 (1600 and
654 $400 \mu\text{m}$) and our assumption that the particles were graupel (Table 5). However, the conjectured
655 downdraft speed is likely an overestimate - because of divergence occurring as the draft
656 approached the surface - and because the sizes in Table 5 likely underestimate what fell to the
657 hotplate. Relevant to the last of these assertions, we used the T/RH/altitude measurements (Table
658 2) to calculate the vertical distance available for growth via riming, and thus for a fall speed
659 increase, between the flight level and the lifted condensation level. Assuming an adiabatically-
660 stratified supercooled cloud and unit collection efficiency (overestimates growth by riming), and
661 no change of particle crosssection (underestimates growth by riming), our calculations indicate
662 that relative increases of size and fall speed were 40 and 20 %, respectively, on 3 January 2017,
663 and that these relative increases were a factor-of-two larger on 15 December 2016.
664



665 Table 5 – Estimates of snow particle fall speed

Date	Mode Size, μm	Assumed Particle Type	Fall Speed, m s^{-1}	Reference
15 December 2015	1600	graupel	1.4	PV11; assuming ρ_1 in their Figure 5
3 January 2016	400	graupel	0.7	PV11; assuming ρ_1 in their Figure 5

666
667



668 Second, there is concern that values of S from 3 January 2017 are underestimated.
669 Although, values of S must be > 0 , we presented 1 Hz values (red points, Fig. 8b) as small as -
670 0.3 mm h^{-1} . Negative values resulted because we did not impose a threshold of 0 mm h^{-1} on the
671 uncorrected snowfall rates (this thresholding is discussed in Z18) and because negative snowfall
672 rate values (uncorrected for catch inefficiency) are amplified by the gauge-catch correction (Sect.
673 2.4). The implication is that 0.2 mm h^{-1} could be added to the one-minute averaged values of S in
674 Table 4 and in Fig. 12. Here, the assumption is that an averaged S of -0.2 mm h^{-1} , in Fig. 8b, is
675 indicating no snowfall and no surface deposition of blowing snow; however, because the hotplate
676 was operated autonomously (Sect. 2.1) we have no way to verify the assumption.



677 **4 – Results**

678 Our paired values of surface-measured precipitation and aircraft-measured radar
679 reflectivity provide evidence - in addition to PV11 - that most prior determinations of the S/Z
680 relationship for W-band radars lead to underestimation of S in situations with rimed snow
681 particles and particularly so in situations with Z smaller than $1 \text{ mm}^6 \text{ m}^{-3}$. We assert that the
682 underestimate stems from the smaller density implicit in most computationally-based S/Z
683 relationships and especially those which assume that snow particles consist of vapor-grown
684 crystals or aggregates of vapor-grown crystals. Values of density are quite different for these two
685 particle types versus that for rimed snow particles. For example, in Matrosov (2007), assuming a
686 2 mm aggregate, the density is $\sim 30 \text{ kg m}^{-3}$, whereas in PV11, assuming a 2 mm graupel particle,
687 the density is $\sim 200 \text{ kg m}^{-3}$.

688 In the previous paragraph, a cutoff at $Z = 1 \text{ mm}^6 \text{ m}^{-3}$ was specified because that is where
689 the separation between the Matrosov (2007) and both our and PV11's WCR observations
690 become evident (Fig. 12). The cutoff was also picked because Kulie et al. (2016) apply it in an
691 analysis of snowfall retrieved using the W-band radar on CloudSat. They concluded that 74% of
692 shallow cumuliform cloud structures, and 37% of nimbostratus cloud structures, have near-
693 surface reflectivities $< 1 \text{ mm}^6 \text{ m}^{-3}$. Depending on which snowfall process dominates in these
694 structures (vapor growth, aggregation, or riming) an alteration of S for $Z < 1 \text{ mm}^6 \text{ m}^{-3}$ (e.g., Fig.
695 12) could have a significant effect on W-band retrievals. For example, the analysis of Kulie et al.
696 2016 (their Figure 6) suggests that the Greenland, Norwegian, and Barents Seas regions may be
697 susceptible to this alteration.

698 Some computationally based S/Z relationships (Surussavadee and Staelin (2006) and
699 Kulie and Bennartz (2009)) do plot between PV11's $S(\rho_1)/Z$ relationship - the black line in Fig.



700 12 - and Matrosov's upper-limit S/Z relationship (the top of the gray area in Fig. 12). Of these
701 the Surussavadee and Staelin relationship assumes that the snow particles are spheres. This
702 seems reasonable for rimed snow particles but not for the crystal and aggregate types modeled by
703 Matrosov (2007) where the particles are approximated as low-density oblate spheroids with their
704 major axis (on average) oriented horizontal. Because of this, proposed space-based platforms
705 may carry instrumentation that can guide selection of a scene-appropriate S/Z relationship. Both
706 lidar and radiometers can sense supercooled liquid water from space, and if combined with
707 Doppler radar, can diagnose precipitation attributable to rimed snow particles. These approaches
708 are being tested in ground-based field studies (Moisseev et al. 2017; Mason et al. 2018).
709



710 **5 - Conclusions**

711 This study is significant because it brings together direct measurements of snowfall rate,
712 measured at the ground, and measurements of reflectivity from an airborne W-band radar.
713 Compared to PV11's $S(\rho_1)/Z$ relationship, shown in Fig. 12, our observations do not depart
714 significantly; however, they do plot somewhat larger. This excess could be consistent with
715 downslope flow that occurs in lee of the Medicine Bow Mountains (Sect. 3.6) or with
716 calculations which indicate that larger density is associated with larger S , in the S -versus- Z
717 coordinate system (PV11), combined with the intrinsic variability of the density of rime ice
718 (Macklin 1962).

719 If the downslope flow hypothesis is correct, we expect it to manifest as negatively-biased
720 retrievals of S , in settings leeward of a ridgeline, where snowfall is produced by riming, and
721 PV11's $S(\rho_1)/Z$ relationship is applied in the retrieval. This follows because PV11, and all other
722 S/Z relationship developers, do not account for the effect of vertical air motion on S values
723 incorporated into their S/Z relationships. Furthermore, the sign of the hypothesized bias will vary
724 from positive (radar-retrieved S larger than a surface-measured S) to negative (radar-retrieved S
725 smaller than a surface-measured S) in the downwind direction across a ridgeline. Finally, we
726 expect the relative magnitude of the hypothesized biases will be enhanced in a situation where Z
727 is measured, snowfall is produced via the diffusion growth of crystals, and the scene-appropriate
728 S/Z relationship is applied.

729



730 **6 - Appendix**

731 This appendix explains how the HP (hotplate) and WCR (Wyoming Cloud Radar)
732 averages were evaluated.

733 With the overflight time symbolized t_O and i an index equal to either 0 or 1 the start
734 and stop times for a one-minute HP average are

$$735 \quad t_{HP,1} = t_O + i \cdot 60 \quad (A1)$$

$$736 \quad t_{HP,2} = t_O + (i + 1) \cdot 60 \quad (A2)$$

737 Examples of $t_{HP,1}$ and $t_{HP,2}$ are at the left and right edges of the $i = 0$ one-minute HP averaging
738 intervals in Fig. 7b and Fig. 8b.

739 The stop time for WCR averaging was calculated as

$$740 \quad t_{WCR,2} = t_O - i \cdot 60 \cdot v_w / g_s \quad (A3)$$

741 Here v_w is a wind advection speed (discussed below) and the second term on the rhs is a wind
742 advection distance divided by the WKA (Wyoming King Air) ground speed (g_s). The start time
743 for WCR averaging was calculated as

$$744 \quad t_{WCR,1} = t_{WCR,2} - (1 + i) \cdot 60 \cdot v_w / g_s \quad (A4)$$

745 The wind advection speed (v_w) in Eqs. A3-A4 was calculated using an altitude-
746 dependent west-to-east wind velocity (u) and an altitude-dependent south-to-north wind
747 velocity (v). These altitude-dependent component velocities were calculated using the horizontal
748 wind vectors in the penultimate and last columns of Table 2. Plots of the component velocities
749 versus altitude and the linear functions used to relate component velocities to altitude are
750 presented in Figs. A1a-b.



751 We assumed an altitude ($z' = 3400$ m) for evaluating the horizontal wind advection
752 vector. This is the altitude of the ridges west and northwest of the HP site (Figs. 3a-b). Picking
753 the altitude to be either $z' = 3200$ m or $z' = 3600$ m does not substantially alter our conclusions.

754 The WKA track vector (Table 2) defines the vertical plane of the WCR measurements.
755 We assumed that wind advection of snow particles occurred parallel to this vector. With the
756 assumption stated in the previous paragraph, the horizontal wind advection speed (v_w) was
757 calculated as the projection of the horizontal wind vector onto the track vector.

$$758 \quad v_w = \frac{u(z') \cdot g s_x + v(z') \cdot g s_y}{(g s_x^2 + g s_y^2)^{1/2}}. \quad (\text{A5})$$

759 In Eq. A5 the west-to-east and south-to-north components of the track vector are symbolized
760 $g s_x$ and $g s_y$. Vector representations of the track vector are in Table 2. On 14/15 December 2016
761 and 3 January 2017 the values of v_w are 7.4 and 8.9 m s^{-1} , respectively.

762 In addition to the properties $g s$ and v_w used to calculate the averaging times (Eqs. A3-
763 A4), the WCR averages were derived using a snow particle downward speed (Eq. A6).

$$764 \quad v_p = |\langle V_D \rangle| + \sigma_{V_D} \quad (\text{A6})$$

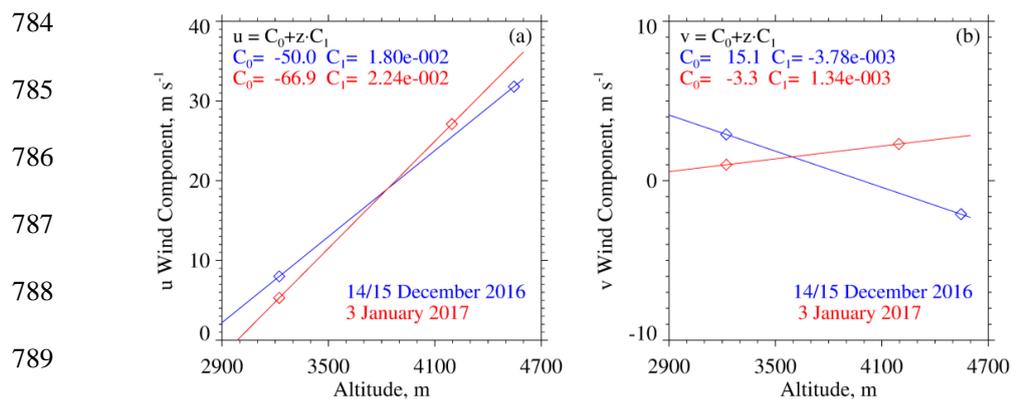
765 Here, v_p is a snow particle downward speed (discussed below), $\langle V_D \rangle$ is the average of
766 Doppler velocities within an averaging domain, $|\langle V_D \rangle|$ is the absolute value of the average,
767 and σ_{V_D} is the standard deviation of the average. On both the lhs and rhs of Eq. A6, all properties
768 are greater than zero.

769 We interpret v_p as the maximum likely snow particle speed toward the surface. There are
770 three reasons for this: 1) For the WCR averaging domains we analyzed, values of $\langle V_D \rangle$ are



771 consistently less than zero. This indicates that snow particles (on average) were moving toward
772 the surface. 2) Again for the WCR averaging domains we analyzed, σ_{V_D} is comparable to
773 $|<V_D>|$ indicating that turbulent eddies transported snow particles upward and downward at
774 speeds comparable to the fall speed of the snow particles in quiescent air. 3) The V_D are
775 reflectivity weighted (Haimov and Rodi 2013) and are thus indicative of the motion of the largest
776 particles within the WCR viewing volume.

777 We now focus on the top of the WCR averaging domains shown schematically in Fig. 6.
778 The slope defining this upper boundary was calculated as $-v_p \cdot g_s / v_w$. That is, particles below
779 this boundary were moving downward sufficiently fast and horizontally sufficiently slow to
780 advect reasonably close to the hotplate. Starting with diagnosed values of g_s and v_w , values of
781 v_p and thus values of the slope, were derived iteratively. The precision of the derived v_p is ± 0.1
782 m s^{-1} .
783



791 Figure A1 – (a) West-to-east (u) wind velocity derived using measurements from the WKA and
792 the AmeriFlux (AF) tower. Also shown is the altitude-dependent linear function used to relate u
793 to altitude. (b) South-to-north (v) wind velocity derived using measurements from the WKA and
794 AF. Also shown is the altitude-dependent linear function used to relate v to altitude. WKA and
795 AF velocities are presented as vectors in the penultimate and last columns of Table 2.

796



797 Data Availability. The WKA and WCR measurements can be obtained from the SNOWIE data
798 archive of NCAR/EOL, which is sponsored by the National Science Foundation. Hotplate gauge
799 measurements are at <https://doi.org/10.15786/20103146>. The US-GLE AmeriFlux measurements
800 are at <https://ameriflux.lbl.gov/>. The Brooklyn Lake SNOTEL gauge measurements are at
801 <https://www.wcc.nrcs.usda.gov/snow/>.

802



803 Author contributions. JRS and MB wrote a successful proposal that funded this research. Field
804 measurements were performed by SF, SM, SH, MB, and JRS. SF wrote her MS dissertation, and
805 this was adapted for this paper by JRS. KS processed the snow particle imagery. AM maintained
806 the measurement sites. All authors contributed to the editing of this paper.
807



808 **Acknowledgements –**

809 We acknowledge the expert technical assistance provided by David Plummer, Larry
810 Oolman, Zane Little, Brent Glover, Edward Sigel, Thomas Drew, and Brett Wadsworth. We
811 thank SNOWIE project PI Jeffery French, who provided the flight data, Gabor Vali who
812 provided the S/Z data points in Figure 12, and John Frank and John Korfmacher who acquire and
813 maintain the GLE-US AmeriFlux data set. This work was supported by the United States
814 National Science Foundation (Award Number 1850809) and the John P. Ellbogen Foundation.



815 **References**

- 816 AmeriFlux, <https://ameriflux.lbl.gov/>, 2021
- 817 Boudala, F.S., R. Rasmussen, G.A. Isaac, and B. Scott, Performance of Hot Plate for Measuring
818 Solid Precipitation in Complex Terrain during the 2010 Vancouver Winter Olympics, J.
819 Atmos. Oceanic Technol., 31, 437–446, <https://doi.org/10.1175/JTECH-D-12-00247.1>,
820 2014
- 821 Brock, F. V., and Richardson, S. J., Meteorological Measurement Systems, Oxford University Press,
822 New York, 304 pp., 2001
- 823 Cocks, S.B., S.M. Martinaitis, B. Kaney, J. Zhang, and K. Howard, MRMS QPE Performance during
824 the 2013/14 Cool Season, J. Hydrometeor., 17, 791–810, <https://doi.org/10.1175/JHM-D-15-0095.1>, 2016
825
- 826 Faber, S., French, J. R., and Jackson, R., Laboratory and in-flight evaluation of measurement
827 uncertainties from a commercial Cloud Droplet Probe (CDP), Atmos. Meas. Tech., 11,
828 3645–3659, <https://doi.org/10.5194/amt-11-3645-2018>, 2018
- 829 Fuller, S.E., Improvement of the Snowfall / Reflectivity Relationship for W-band Radars, MS
830 Thesis, Department of Atmospheric Science, University of Wyoming, 2020
- 831 Geerts, B., Q. Miao, Y. Yang, R. Rasmussen, and D. Breed, An Airborne Profiling Radar Study of
832 the Impact of Glaciogenic Cloud Seeding on Snowfall from Winter Orographic Clouds, J.
833 Atmos. Sci., 67, 3286–3302, <https://doi.org/10.1175/2010JAS3496.1>, 2010
- 834 Haimov, S., and Rodi, A., Fixed-Antenna Pointing-Angle Calibration of Airborne Doppler Cloud
835 Radar, Journal of Atmospheric and Oceanic Technology, 30, 2320-2335,
836 <https://doi.org/10.1175/JTECH-D-12-00262.1>, 2013
- 837 Kochendorfer, J., Nitu, R., Wolff, M., Mekis, E., Rasmussen, R., Baker, B., and Jachcik, A, Testing
838 and development of transfer functions for weighing precipitation gauges in WMO-SPICE,
839 Hydrology and Earth System Sciences, 2, 1437-1452, <https://doi.org/10.5194/hess-22-1437-2018>, 2018
840
- 841 Korolev, A. V., E. F. Emery, J. W. Strapp, S. G. Cober, G. A. Isaac, M. Wasey, and D. Marcotte, Small
842 ice particles in tropospheric clouds: Fact or artifact? Airborne Icing Instrumentation
843 Evaluation Experiment, Bull. Amer. Meteor. Soc., 92, 967–973,
844 <https://doi.org/10.1175/2010BAMS3141.1>, 2011
- 845 Kulie, M. S., and R. Bennartz, Utilizing Spaceborne Radars to Retrieve Dry Snowfall, J. Appl.
846 Meteor. Climatol., 48, 2564–2580, <https://doi.org/10.1175/2009JAMC2193.1>, 2009
- 847 Kulie, M. S., Milani, L., Wood, N. B., Tushaus, S. A., Bennartz, R., and L’Ecuyer, T. S., A Shallow
848 Cumuliform Snowfall Census Using Spaceborne Radar, Journal of Hydrometeorology, 4,
849 1261-1279, <https://doi.org/10.1175/JHM-D-15-0123.1>, 2016



- 850 Lawson, R. P., O'Connor, D., Zmarzly, P., Weaver, K., Baker, B., Mo, Q., and Jonsson, H., The 2D-S
851 (Stereo) Probe: Design and Preliminary Tests of a New Airborne, High-Speed, High-
852 Resolution Particle Imaging Probe, *J. Atmos. Ocean. Tech.*, 23, 1462–1477,
853 <https://doi.org/10.1175/JTECH1927.1>, 2006
- 854 Locatelli, J.D. and Hobbs, P.V., Fall speed and masses of solid precipitation particles, *J. Geophys.*
855 *Res.*, 79, 2185–2197, <https://doi.org/10.1029/JC079i015p02185>, 1974
- 856 Macklin, W.C., The density and structure of ice formed by accretion, *Q.J.R.Meteorol.Soc.*, 88: 30-
857 50. doi:10.1002/qj.49708837504, <https://doi.org/10.1002/qj.49708837504>, 1962
- 858 Marlow, S.A, J.M. Frank, M. Burkhart, B. Borkhuu, S.E. Fuller, and J.R. Snider, Snowfall
859 measurements in mountainous terrain, in revision for the *Journal of Applied Meteorology*
860 and *Climatology*, http://www-das.uwyo.edu/~jsnider/JAMC-D-22-0093_6.pdf, 2022
- 861 Martinaitis, S.M., S.B. Cocks, Y. Qi, B.T. Kaney, J. Zhang, and K. Howard, Understanding winter
862 precipitation impacts on automated gauge observations within a real-rime system, *J.*
863 *Hydrometeor.*, 16, 2345-2363, <https://doi.org/10.1175/JHM-D-15-0020.1>, 2015
- 864 Mason, S. L., Chiu, C. J., Hogan, R. J., Moisseev, D., and Kneifel, S., Retrievals of riming and snow
865 density from vertically pointing Doppler radars, *Journal of Geophysical Research:*
866 *Atmospheres*, 123, 13,807– 13,834, <https://doi.org/10.1029/2018JD028603>, 2018
- 867 Matrosov, S.Y., Modeling Backscatter Properties of Snowfall at Millimeter Wavelengths, *J. Atmos.*
868 *Sci.*, 64, 1727-1736, <https://doi.org/10.1175/JAS3904.1>, 2007
- 869 Moisseev, D., von Lerber, A., and Tiira, J., Quantifying the effect of riming on snowfall using
870 ground-based observations, *J. Geophys. Res. Atmos.*, 122, 4019–4037,
871 doi:10.1002/2016JD026272, 2017
- 872 Panofsky, H.A. and Dutton, J.A., *Atmospheric Turbulence*, Wiley-Interscience, New York, 397 pp.,
873 1984
- 874 Pokharel, B. and G. Vali, Evaluation of Collocated Measurements of Radar Reflectivity and Particle
875 Sizes in Ice Clouds, *J. Appl. Meteor. Climatol.*, 50, 2104–2119,
876 <https://doi.org/10.1175/JAMC-D-10-05010.1>, 2011
- 877 Rasmussen, R.M., J. Hallett, R. Purcell, S.D. Landolt, and J. Cole, The Hotplate precipitation gauge,
878 *J. Atmos. Oceanic Technol.*, 28, 148-164, <https://doi.org/10.1175/2010JTECHA1375.1>,
879 2011
- 880 R.M. Young Company, Model 05103 Wind Monitor, 2001
- 881 Serreze, M. C., M. P Clark, and R. L. Armstrong, D. A. MacGinnis, and R. S. Pulwarty, Characteristics
882 of the western United States snowpack from snowpack telemetry (SNOTEL) data, *Water*
883 *Resources Research*, 35, 2145-2160, <https://doi.org/10.1029/1999WR900090>, 1999
- 884 Skofronick-Jackson, G., and Coauthors, The Global Precipitation Measurement (GPM)
885 Mission for science and society, *Bull. Amer. Meteor. Soc.*, 98, 1679–1695,
886 <https://doi.org/10.1175/BAMS-D-15-00306.1>, 2017



- 887
888 Smith, P.L., Equivalent radar reflectivity factors for snow and ice particles, *J. Climatol. Appl.*
889 *Meteor.*, 23, 1258–1260, <https://doi.org/10.1175/1520->
890 0450(1984)023<1258:ERRFFS>2.0.CO;2, 1984
- 891 Surussavadee, C. and D.H. Staelin, Comparison of AMSU millimeter-wave satellite observations,
892 MM5/TBSCAT predicted radiances, and electromagnetic models for hydrometeors, *IEEE*
893 *Trans. Geosci. Remote Sens.*, 44, 2667–2678, 10.1109/TGRS.2006.873275, 2006
- 894 Tessendorf, S. A., and Coauthors, A transformational approach to winter orographic weather
895 modification research: The SNOWIE Project, *Bulletin of the American Meteorological*
896 *Society*, 100, 71–92, <https://doi.org/10.1175/BAMS-D-17-0152.1>, 2019
- 897 Vaisala, User's Guide, Vaisala Weather Transmitter, WXT520, 2012
- 898 Wang, P.K., and W. Ji, Collision Efficiencies of Ice Crystals at Low–Intermediate Reynolds Numbers
899 Colliding with Supercooled Cloud Droplets: A Numerical Study, *Journal of the Atmospheric*
900 *Sciences*, 57, 1001–1009, <https://doi.org/10.1175/1520->
901 0469(2000)057<1001:CEOICA>2.0.CO;2, 2000
- 902 Wilson, J., and E. Brandes, Radar measurement of rainfall—A summary, *Bull. Amer. Meteor. Soc.*,
903 60, 1048–1058, [https://doi.org/10.1175/1520-0477\(1979\)060<1048:RMORS>2.0.CO;2](https://doi.org/10.1175/1520-0477(1979)060<1048:RMORS>2.0.CO;2),
904 1979
- 905 Wolfe, J.P., and J.R. Snider, A relationship between reflectivity and snow rate for a high-altitude
906 S-band radar, *J. Appl. Meteor. Climatol.*, 51, 1111–1128, <https://doi.org/10.1175/JAMC->
907 D-11-0112.1, 2012
- 908 Zaremba, T.J., and Coauthors, Vertical motions in orographic cloud systems over the Payette
909 River Basin. Part 1: Recovery of vertical motions and their uncertainty from airborne
910 Doppler radial Velocity Measurements, in press at the *Journal of Applied Meteorology*
911 *and Climatology*, <https://doi.org/10.1175/JAMC-D-21-0228.1>, 2022
- 912 Zelasko, N., Wettlaufer, A., Borkhuu, B., Burkhart, M., Campbell, L. S., Steenburgh, W. J., and
913 Snider, J.R., Hotplate precipitation gauge calibrations and field measurements, *Atmos.*
914 *Meas. Tech.*, 11, 441–458, <https://doi.org/10.5194/amt-11-441-2018>, 2018
- 915 Zikmunda, J. and Vali, G., Fall patterns and fall velocities of rimed ice crystals, *J. Atmos. Sci.*, 29,
916 1334–1347, [https://doi.org/10.1175/1520-0469\(1972\)029<1334:FPAFVO>2.0.CO;2](https://doi.org/10.1175/1520-0469(1972)029<1334:FPAFVO>2.0.CO;2),
917 1972