

A novel infrared imager for studies of hydroxyl and oxygen nightglow emissions in the mesopause above northern Scandinavia

5 Peter Dalin¹, Urban Brändström¹, Johan Kero¹, Peter Voelger¹, Takanori Nishiyama²,
Trond Trondsen³, Devin Wyatt³, Craig Unick³, Vladimir Perminov⁴, Nikolay Pertsev⁴,
Jonas Hedin⁵

¹Swedish Institute of Space Physics (IRF), Kiruna, Sweden
10 ²National Institute of Polar Research, Tokyo, Japan
Keo Scientific Ltd., Calgary, Canada
⁴A.M. Obukhov Institute of Atmospheric Physics, Moscow, Russia
⁵Stockholm University, Department of Meteorology, Stockholm, Sweden

15 *Correspondence to:* Peter Dalin (pdalin@irf.se)

Abstract. The paper describes technical characteristics and presents first scientific results of a novel infrared imaging system (imager) for studies of nightglow emissions coming from hydroxyl (OH) and molecular oxygen (O₂) layers in the mesopause region (80-100 km) above northern Scandinavia. The
20 OH imager was put into operation in November 2022 at the Swedish Institute of Space Physics located in Kiruna (67.86°N, 20.42°E, 400 m altitude). The OH imager records selected emission lines in the OH (3-1) band near 1500 nm to obtain intensity and temperature maps at around 87 km altitude. Also, the OH imager registers infrared emissions coming from the O₂ IR A-band airglow at 1268.7 nm in order to obtain O₂ intensity maps at a slightly higher altitude around 94 km. This technique allows
25 tracing wave disturbances both in horizontal and vertical domains in the mesopause region. Validation and comparison of the OH (3-1) rotational temperature with collocated lidar and Aura/MLS satellite temperatures are performed. First scientific results obtained with the OH imager for the first winter season (2022-2023) are discussed.

30 1 Introduction

Besides a well-known scattering of sunlight on air molecules (Rayleigh scattering) and particle scattering, the Earth's atmosphere emits its own radiation 24 hours a day, year-round and on global scales. So-called airglow emissions occur in the mesosphere (50-90 km) and thermosphere (above 90 km). The most prominent airglow emissions are red and green lines of atomic oxygen as well as
35 ultraviolet and infrared bands of molecular oxygen, hydroxyl (OH) emissions in the infrared part of the spectrum, yellow emission line of sodium (Na), ultraviolet and infrared bands of nitric oxide (NO) (e.g. Chamberlain, 1961; Khomich et al., 2008; Savigny, 2017).

Studies of airglow emissions provide important sources of information on physical and chemical states of the middle and upper atmosphere. In particular, temperature of emission layers in the middle
40 and upper atmosphere can be determined by analyzing spectral bands and lines of airglow emissions. By analyzing spatial-temporal variations in the emission intensities and/or temperature one can investigate the short- and long-term dynamical state of the middle and upper atmosphere, which in turn is a hot topic of current atmospheric research (e.g., Fritts and Alexander, 2003; Plougonven and Zhang, 2014; Reid et al., 2014). The key role of the atmospheric dynamical state belongs to waves, being a

45 fundamental property of the atmosphere as they transport energy and momentum across great distances
(Gossard and Hook, 1975).

Observations of airglow emissions are used for inferring parameters of atmospheric waves since
temperature and emission intensity fields in the mesosphere and thermosphere are directly modulated
by gravity and planetary waves, solar thermal and lunar gravitational tides (e.g., Garcia et al., 1997;
50 Taylor et al., 1997; Pautet et al., 2005; Suzuki et al., 2008; Offermann et al., 2009; Gao et al., 2011;
Perminov et al., 2018; Pertsev et al., 2021).

Another very important aspect in atmospheric physics is a long-term monitoring of atmospheric
parameters. The scientific community is wondering to what extent long-term changes in H₂O, CO₂,
N₂O, CH₄, O₃ (greenhouse gases of anthropogenic and natural origin) are contributing to long-term
55 variability in the atmospheric temperature. While CO₂ is thought to endow to global warming on Earth,
it also cools the middle atmosphere (Roble and Dickinson, 1989; Mlynzcak et al., 2022). The increased
H₂O concentration leads to warmer temperatures due to a strong positive water vapor feedback (Dessler
et al., 2008). Methane oxidizes in the stratosphere and mesosphere producing an additional amount of
water vapor in the middle atmosphere (Thomas and Olivero, 2001). The global monthly mean
60 atmospheric methane abundance demonstrates a constant increase since 1987
(https://gml.noaa.gov/ccgg/trends_ch4/), which may have resulted in the observed long-term trend in
water vapor in the mesopause. Lübken et al. (2018) have obtained a summer mesopause cooling of ~1.2
K/decade for the period of 1960-2008 based on LIMA and MIMAS model simulations which included
long-term evolutions of minor atmospheric species CO₂, CH₄, H₂O and O₃. Thus, the long-term
65 temperature trend in the middle atmosphere is a complex function of several atmospheric components.
Estimates of temperatures of the middle atmosphere are only possible with help of remote sensing
instruments. That is why observations of airglow emissions are used for long-term monitoring of the
temperature in the mesopause region (Semenov and Shefov, 1999; Espy and Stegman, 2002;
Offermann et al., 2010; Ammosov et al., 2014; Kalicinsky et al., 2016; Perminov et al., 2018; Dalin et
70 al., 2020; French et al., 2020).

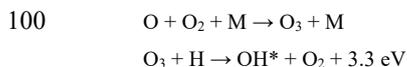
In the present paper, we describe a novel infrared imaging instrument (in the following called the
OH imager) that measures infrared emissions of selected lines in the OH (3-1) band at wavelengths
near 1.5 μm to produce intensity and temperature maps in the mesopause region at around 87 km
altitude. Additionally, the OH imager records nightglow emissions from the O₂ IR A-band at 1268.7
75 nm to obtain O₂ emission intensity maps at a slightly higher altitude of about 94 km. Combining data
from both layers allows tracing wave disturbances in 3-D space, i.e., horizontally and vertically in the
mesopause region. The OH imager was installed in November 2022 at the Swedish Institute of Space
Physics (IRF) in Kiruna (67.86°N, 20.42°E, 400 m altitude) on the eastern slope of the Scandinavian
mountain range. Mountain gravity waves occur frequently at the location, thus the OH imager can be
80 utilized to study the influence of such waves on the thermal and dynamical regime of the mesopause
region over Kiruna. The OH registrations are part of the long-term research infrastructure commitment
of the Kiruna Atmospheric and Geophysical Observatory (KAGO), part of the IRF, whose objective is
to provide present and future scientists with long, continuous archived time series of data of highest
possible scientific quality from its set of instruments.

85 In section 2, we describe technical characteristics of the OH imager as well as the principle of OH
(3-1) rotational temperature measurements. In section 3, the temperature estimation and its uncertainty,
calibration and validation technique are presented. Validation of OH (3-1) temperatures with the
collocated Esrange lidar temperature measurements as well as a comparison of OH temperatures with
simultaneous Aura/MLS satellite temperatures are shown in section 4. In section 5, first scientific
90 results obtained with the OH imager are discussed. Finally, conclusions are given in section 6.

2 Instrument description

2.1 Theory and basic knowledge:

The hydroxyl (OH) airglow emission layer exists in the Earth's atmosphere between 75 and 100
95 km, with the highest concentration at about 87 km and the full width at half maximum of about 8-10
km (Baker and Stair, 1988; Savigny, 2017), although both of these parameters of the OH-layer change
depending on geolocation and season (Perminov et al., 1999, 2018; Melo et al., 2000; Grygalashvily et
al., 2014). Its rotational-vibrational bands, covering the spectral region from 0.5 to 4.5 μm , result
mainly from chemiluminescent processes (Bates and Nicolet, 1950):



where M is a molecule of nitrogen or oxygen, OH* is the OH vibrationally-excited molecule.
Returning to their ground state, the OH* molecules emit visible and infrared radiation in the spectral
105 range between 500 and 4500 nm.

The OH (8-3) band at 732-742 nm and the OH (6-2) band at 830-845 nm are commonly used in OH
airglow observations, which are detectable by commercial CCD cameras. However, at high latitudes,
auroral emissions (molecular and atomic oxygen lines between 730 and 845 nm) have stronger
intensities than those of the OH airglow bands (Christensen et al., 1978; Sivjee et al., 1979). Therefore,
110 measurements of OH nightglow emissions in the Short Wave Infrared Range (SWIR) of the spectrum
(900-1700 nm) are preferable in order to minimize auroral contaminations at high latitudes (Pautet et
al., 2014; Nishiyama et al., 2021). For this purpose, the OH (3-1) band (1470-1550 nm) is used by the
OH imager to measure nightglow emissions at high latitudes. It is important to note that the OH (3-1)
intensity is by two orders of magnitude greater than intensities of the OH (8-3) and OH (6-2) bands.

115 Also, the OH (3-1) emission is less affected by water vapor absorption. All these factors make a strong
preference in choosing the OH (3-1) nightglow observation in Kiruna, being located underneath the
auroral oval where auroral emissions occur frequently.

2.2 Instrument description

120 The main component of the OH imager is an infrared camera (Teledyne Princeton Instruments
NIRvana 640) equipped with an Indium-Gallium-Arsenide (InGaAs) sensor for low-light scientific
SWIR imaging applications. The NIRvana 640 camera employs a 640 x 512 InGaAs array, charged-
coupled device (CCD), with a spectral response from 900 to 1700 nm. Each pixel has a size of 20 x 20
 μm . The detector is Peltier-cooled to -80°C to minimize thermally generated noise and, thus improve

125 signal-to-noise ratio. The NIRvana 640 camera has a 16-bit digitization, with low readout noise. Other
components of the instrument are a SWIR wide-angle primary lens ($\sim 120^\circ$ field of view), an $f/1.0$
SWIR-optimized telecentric optical system, a thermally stabilized 8-position filter wheel with four
130 narrow-band (three-cavity) interference filters, brushless DC servo motor drive for the filter wheel,
filter wheel control unit, power supply and, instrument control and data acquisition computer. As the
center wavelengths of interference filters are temperature dependent, the filter wheel's closed-loop
temperature controller heats the filter wheel cavity to $+23 \pm 0.1^\circ\text{C}$. Along with four narrow-band
interference filters, a blocking dark filter is used to collect data for subtraction of dark-current noise
from nighttime measurements. The blocking dark filter's emissivity is matched to that of the narrow-
band interference filters thus ensuring the dark-current noise collected with it is representative of the
135 dark-current noise present in nighttime measurements. The primary wide-angle lens, telecentric optics
and re-imaging lens were custom designed and produced for measurements in the SWIR. The optical
design of the OH imager inscribes the 120° field of view image circle within the short (vertical)
dimension of the 640×512 pixel sensor, resulting in an image circle with a 512 pixel diameter. The
image is therefore cropped to 512×512 pixels for image processing purposes. The OH imager was
140 designed and built by Keo Scientific Ltd (<http://keoscientific.com>). The main instrument characteristics
are summarized in Table 1.

Nightglow observations are restricted to nighttime and twilight (solar elevation angle less than -6.5°)
for optimum signal quality. At the arctic location of Kiruna, we are thus able to conduct
nightglow measurements from September until middle of April. The exposure times for the OH,
145 background and dark filters are set to 30 s, whereas the exposure time for the O_2 IR A-band filter is set
to 40 s, yielding a complete filter cycle in 2.75 min, implying that the instrument can readily resolve
the highest frequency range in the gravity wave spectrum. Figure 1 shows the OH imager installed at
the Knutstorp observatory (67.86°N , 20.42°E), 2 km north of IRF, on 29 November 2022.

150 **3 Temperature estimation and its uncertainty, error analysis, calibration technique**

3.1 Temperature estimation

The temperature of the air in the mesopause region is determined using the brightness ratio of the
two vibrational-rotational lines $P_1(2)$ and $P_1(4)$ in the OH (3-1) band. This method is based on the
assumption that the rotational level populations are in local thermodynamic equilibrium (LTE).
155 Although the OH vibrational states are excited due to non-thermal processes, the rotational level
population is still in LTE due to the long radiative lifetimes of vibrational states (e.g., Savigny, 2017).
This method has been proved to be valid and has been applied for several decades in temperature
measurements of the mesopause region (Sivjee and Hamwey, 1987; Makhlof et al., 1995; Espy and
Stegman, 2002; Pertsev and Perminov, 2008; Suzuki et al., 2008; Offermann et al., 2010; Grygalashvyly
160 et al., 2014; Pautet et al., 2014; Kalicinsky et al., 2016; Perminov et al., 2018).

For the OH rotational temperature estimation, we use Eq. (4) from Pautet et al. (2014). At the same
time, we have corrected this equation for newly estimated Einstein A-coefficients for the OH (3-1) band
at the $P_1(2)$ and $P_1(4)$ lines as calculated by Brooke et al. (2016). New A-coefficients for the $P_1(2)$ and

165 P₁(4) lines are 9.895802 and 13.15222 s⁻¹, respectively. After the correction for the A-coefficients, the following equation for the OH rotational temperature estimation T_r is obtained:

$$T_r = \frac{259.58}{\ln(2.658 \cdot R)} \quad (1)$$

170 where R is the brightness ratio $B(P_1(2))/B(P_1(4))$ inferred from measurements of the P₁(2) and P₁(4) lines. Equation (1) differs from Eq. (4) by Pautet et al. (2014) in terms of the multiplier at the R parameter: 2.658 versus 2.644 used in Eq. (4) by Pautet et al. (2014). This results in 0.7-1.2 K temperature difference between temperature estimations in the present paper and Pautet et al. (2014) for a typical range of winter temperatures (180-240 K) in the mesopause region.

3.2 Absolute calibration

175 The purpose of the absolute calibration of the OH imager is to establish a relationship between the arbitrary, instrument-dependent raw pixel values (counts) and the column emission rate in units of the Rayleigh (R). The imager absolute calibration is a multi-stage process that was performed by Keo Scientific Ltd. First, filter bandpass curves for the three OH(3-1) filters (atmospheric background, P₁(2) and P₁(4) lines) were measured by using a tunable SWIR laser. Second, on-axis broadband sensitivity was determined by imaging the output of a calibrated (National Institute of Standards and Technology, 180 NIST-traceable) integrating sphere through the imager, including the filters, and onto the InGaAs photodiode array sensor. Additionally, an on-axis ‘at-wavelength’ sensitivity was measured for each filter using a collimated laser beam calibrated with a (NIST-traceable) optical power meter through the imager and onto the InGaAs photodiode array sensor. This resulted in the final absolute calibration coefficients, valid on-axis. Finally, a flat-field correction procedure was performed to compensate for 185 off-axis vignetting, resulting in uniformity coefficients, one for each OH filter, allowing one to apply the absolute calibration coefficients to any pixel in the image. A calibration certificate with supporting files was issued by Keo Scientific.

3.3 Geometrical calibration

190 The geometrical calibration of the OH imager is a necessary procedure to measure absolute horizontal coordinates of each pixel of the sensor at the instrument location. The geometrical calibration was performed by analyzing images of a clear night sky with reference stars by using a wide-band filter, also present in the filter wheel. The camera optical model was chosen to be described by the 3rd order polynomial. By comparing theoretical horizontal coordinates of reference stars (more than 50 reference stars have been identified in reference images) with their measured coordinates, ten 195 free coefficients of the 3rd order polynomial were calculated. These ten coefficients describe all possible optical distortions in the whole camera optical path. This procedure finally resulted in calculating absolute horizontal coordinates (elevation and azimuth angles) of each pixel followed by a georeference procedure to project each pixel onto the Earth’s surface. For this, the mean altitudes of the 200 OH and O₂ layers were chosen as 87 km and 94 km, respectively. The spatial horizontal resolution in the center of the image is about 0.3 km, reaching ~1.2 km at the edges of the image. A more detailed

description of the geometrical calibration, georeference procedure and error analysis can be found in Dalin et al. (2015).

205 3.4 Error analysis

In order to estimate the total error of temperature estimation, we need to rewrite the basic Eq. (1). The R brightness ratio in Eq. (1) depends on four variables represented by their raw measurement values: intensity of the $P_1(2)$ line I_{P12} , intensity of the $P_1(4)$ line I_{P14} , atmospheric background intensity I_{bg} and instrumental dark-current noise n_{dc} . Then Eq. (1) can be rewritten in terms of these four variables as follows:

$$T_r = \frac{259.58}{\ln\left(2.658 \cdot \frac{k_1 \cdot (k_2 \cdot I_{P12} - k_3 \cdot I_{bg} + k_4 \cdot n_{dc})}{k_5 \cdot (k_6 \cdot I_{P14} - k_3 \cdot I_{bg} + k_7 \cdot n_{dc})}\right)} \quad (2)$$

where $k_1, k_2, k_3, k_4, k_5, k_6, k_7$ are the coefficients determined in a laboratory during the absolute calibration procedure.

In order to estimate these four raw measurement values and their errors, we follow the method proposed by Pautet et al. (2014). Measurements should be performed during a time when atmospheric and geomagnetic perturbations are at a minimum (i.e., small variations in raw OH emission lines and in the atmospheric background, no contamination by tropospheric clouds, no auroral emissions). A period of one hour on 6 April was chosen to evaluate the four raw measurement variables and their errors. Figure 2 illustrates the raw intensities for the $P_1(2)$, $P_1(4)$ lines as well as for the atmospheric background (BG) measured at the zenith over a single-pixel area ($0.003 \times 0.003^\circ$). To reduce measurement perturbations as much as possible, the raw data were smoothed by a 3-point moving average. Then the residuals were determined by subtracting the smoothed curves from the raw intensity data and their standard deviations were estimated. There have been obtained the following smoothed average values and their standard deviations: $57,711 \pm 290$ counts for the $P_1(2)$ line, $48,552 \pm 110$ counts for the $P_1(4)$ line, and $31,301 \pm 126$ counts for the BG. The measured random temperature error (δT_m) is calculated by using a general equation of error propagation in the case of a function having several variables (Taylor, 1997):

$$\delta T_m = \sqrt{\left(\frac{\partial T}{\partial I_{P12}} \delta I_{P12}\right)^2 + \left(\frac{\partial T}{\partial I_{P14}} \delta I_{P14}\right)^2 + \left(\frac{\partial T}{\partial I_{bg}} \delta I_{bg}\right)^2} \quad (3)$$

where $\partial T/\partial I_{P12}$, $\partial T/\partial I_{P14}$ and $\partial T/\partial I_{bg}$ are the partial derivatives of T_r (see Eq. 2) with respect to raw intensity values $P_1(2)$, $P_1(4)$ and atmospheric background, respectively, and δI_{P12} , δI_{P14} and δI_{bg} are their standard deviations. When inserting the parameters that were estimated above into Eq. (3), a temperature measurement random error of 1.80 K is derived.

The temperature instrumental error is mainly governed by the dark-current noise of the CCD sensor. The dark-current noise was continuously monitored by using the dark blocking filter as part of the

measurement cycle (every 2.75 min), and was subsequently subtracted from the recorded raw intensity values of the OH and O₂ emission lines. Example of the dark-current noise is shown in Fig. 3, demonstrating a linear stability of the dark-current noise during the night of 6-7 April 2023, with the small statistically insignificant regression coefficient equal to 186±355 counts/day. The mean value of the dark-current noise n_{dc} and its standard deviation δn_{dc} is 30,900±183 counts for this night.

The total (instrumental + measurement) error δT_{tot} is calculated by adding a term for the instrumental error to Eq. (3):

$$\delta T_{tot} = \sqrt{\left(\frac{\partial T}{\partial n_{dc}} \delta n_{dc}\right)^2 + \left(\frac{\partial T}{\partial I_{P12}} \delta I_{P12}\right)^2 + \left(\frac{\partial T}{\partial I_{P14}} \delta I_{P14}\right)^2 + \left(\frac{\partial T}{\partial I_{bg}} \delta I_{bg}\right)^2} \quad (4)$$

where the first term $\partial T/\partial n_{dc}$ is the partial derivative of the temperature with respect to dark-current noise n_{dc} and δn_{dc} is its standard deviation. Based on the above estimated mean dark-current noise and its standard deviation, the mean temperature instrumental error is equal to 0.15 K on the night 6-7 April 2023. Finally, the resulting rotational temperature and its total error is equal to 193.90±1.81 K for this particular case shown in Fig. 2.

At the same time, we should note that actual variability of atmospheric conditions due to complex wave dynamics (turbulence, gravity and planetary waves, solar and lunar tides, seasonal variations) is greater resulting in an average standard deviation of the mesopause temperature above Kiruna of ~10 K for the 2023 (see section 5). The most important is that the OH imager enables resolving small temperature variations of 4-6 K due to small-scale gravity waves as will be demonstrated in section 5.

4 Temperature validation and comparison

4.1 OH imager and lidar measurements

A Rayleigh/Mie/Raman backscatter lidar located at Esrangle (~40 km eastward of Kiruna) is used to validate the OH (3-1) temperatures as measured by the OH imager. The Esrangle lidar was developed by the Bonn University to monitor aerosols in the troposphere, stratosphere and mesosphere (Blum and Fricke, 2005). The signal from the aerosol-free part of the atmosphere can be used to determine the temperature assuming atmospheric hydrostatic equilibrium. The vertical and time resolutions of the obtained measurements are 150 m and 4.2 min, respectively. We used lidar backscattered signal from the 532-nm wavelength channel to calculate an atmospheric temperature profile above Esrangle. Note that a temperature profile measured by the Esrangle lidar is located in the field of view of the OH imager, i.e., common volume simultaneous temperature measurements have been conducted to validate the OH (3-1) temperatures.

The Esrangle lidar was operated when weather conditions were favorable during 7 nights in the period from January to March 2023. An example of a temperature comparison, made on the night 2/3 February 2023, between the OH imager and lidar is shown in Fig. 4. The OH temperature measurement has been selected from the temperature map having the closest position to Esrangle with 0.3 km uncertainty. At the upper end of the profile (at around 90 km) a seed temperature equal to 188 K was taken from Aura/MLS satellite temperature measurements. Mean values of the OH (3-1) and lidar

temperatures averaged for the time interval between 18:43 and 05:58 UT on this night are presented. In addition, the lidar temperature profile has been smoothed by using a 1.5 km running average in height in order to reduce large fluctuations in the upper part of the profile. The Aura/MLS temperature profile, having the closest position to Kiruna (about 420 km away) for this case, is shown in Fig. 4 by the red line as well. Since the OH (3-1) temperature is measured across the whole OH layer between about 82 and 92 km, with the mean height of ~ 87 km, for validation purposes the lidar temperature profile has been processed using a Gaussian function with the maximum at 87 km and having the Full Width at Half Maximum (FWHM) of 9.3 km height-weighted between 82 and 90 km; the same technique as was used by Pautet et al. (2014). The average height-weighted lidar temperature is equal to 191.8 ± 12.3 K and the average OH (3-1) temperature is 193.6 ± 8.3 K. The temperature difference between these estimations is 1.8 K for this particular night, i.e., within the error margin. Following the same technique, we have made comparisons between the OH (3-1) and lidar temperatures for all available 7 nights of simultaneous common-volume measurements, resulting in the following statistics: average temperature difference, including sign, between the OH imager and Erange lidar is -0.2 ± 1.6 K. It means that, from statistical point of view, the difference between the OH (3-1) and Erange lidar temperatures is about zero. The maximum temperature difference between these instruments was -2.9 K on the night 5-6 March 2023. Note that we estimate the Gaussian height-weighted lidar temperature between 82 and 90 km with the mean height of 87 km. However, the actual height distribution of the OH layer above Erange is unknown since it varies, in general, in time and space (i.e., Pautet et al., 2014). Thus, one can conclude that there is a good agreement between the OH (3-1) and Erange lidar temperature measurements on individual nights (within 3 K) as well as on an average basis (close to zero), after taking into account dynamical processes in the mesopause region and instrument errors.

4.2 OH imager and Aura/MLS measurements

The Microwave Limb Sounder (MLS) radiometer onboard the Aura satellite provides temperature profiles of the middle atmosphere with near-global coverage. Aura/MLS temperature data (ver.5.0 and level 2 data quality) can be obtained from the NASA public web-site: https://acdisc.gesdisc.eosdis.nasa.gov/data/Aura_MLS_Level2/. The description of the MLS temperature retrieval and its validation can be found in Froidevaux et al. (2006) and Schwartz et al. (2008). We used Aura/MLS temperature profiles as an additional way to compare temperatures that were retrieved with the OH imager.

Vertical resolution of Aura/MLS temperature measurements is about 11-12 km in the mesopause region and temperature precision of individual profiles is about 3.5 K at mesopause heights (see the Aura/MLS data quality and description document available at the NASA web-site). Aura/MLS temperatures at the pressure level of $4.6 \cdot 10^{-3}$ hPa (about 86 km altitude) have been chosen to be closest to the OH imager location (less than 300 km away from Kiruna) and measured simultaneously with OH (3-1) temperatures within the same hour. The comparison between OH imager and Aura/MLS temperature measurements for the period January-April 2023 is illustrated in Fig. 5. In total, 22 individual temperature measurements have been collected. One can see that there are several cases in which the temperatures agree well within their uncertainties (days of year 14, 58, 61, 65, 83, 85, 86, 92,

93, 99, 101). For other days the agreement is worse. Note that even if the Aura/MLS instrument measures temperature in the closest proximity to the OH imager, the Aura/MLS scans horizontally a rather big volume of the mesopause (about 250-280 km) along its trajectory which is not exactly the same mesopause volume as the OH imager measures; besides, the vertical resolution of Aura/MLS temperature data in the mesopause is not equal to the width of the OH layer. All these factors complicate a precise comparison. Natural variabilities of the atmosphere due to wave dynamics and chemical processes result in different temperatures in different atmospheric volumes. At the same time, two important conclusions can be drawn from Fig. 5: a) there is no systematic bias (neither negative nor positive) of the OH (3-1) temperature relative to Aura/MLS temperature since there are OH (3-1) temperatures higher and lower Aura/MLS temperatures; b) the average temperature difference, including sign, between the OH imager and Aura/MLS is 2.8 ± 7.8 K. This is a small difference taking into account the above-mentioned limitations.

5 Results and discussion

Maps of OH (3-1) intensities in the Rayleigh and O₂ IR A-band intensity in relative units as well as the OH (3-1) temperature map on 16 February 2023 are illustrated in Fig. 6. One can see there is a fine modulation of the temperature (Fig. 6d) due to small-scale gravity waves (ripples) with wavelengths of 4-8 km, overlying larger gravity wave crests; the temperature amplitudes due to these small-scale waves are in range of 4-6 K. It means that this case clearly demonstrates the ability of the OH imager to resolve small-amplitude small-scale dynamical temperature variations that is of importance for studying a high frequency range of the gravity wave spectrum. Another important feature of this case there is a medium-scale modulation, oriented in the direction from north-west to south-east, having horizontal wavelengths of 20-40 km and temperature amplitudes of 7-10 K. These medium-scale gravity wave crests are clearly seen on all the four maps, meaning that the same gravity wave package was propagating through both the OH layer (maximum is at ~87 km) and the O₂ layer (maximum is at ~94 km). An analysis of a sequence of temperature maps have demonstrated that these medium-scale gravity waves were moving from south-west to north-east, with the observed horizontal phase speed of about 30-35 m/s. This type of medium-scale gravity waves is common in the mesopause region (e.g., Pautet et al., 2011; Demissie et al., 2014). Thus, the OH imager allows retrieving information on gravity waves propagating both in horizontal and vertical domains in the mesopause region. A detailed analysis of gravity wave characteristics and of the gravity wave spectrum is beyond the scope of this paper and will be addressed in future studies.

The OH (3-1) daily mean rotational temperatures measured above Kiruna for the 2023 winter season (13 January-16 April) are illustrated in Fig. 7 by the black circles. The OH (3-1) temperature values are within the range of 178-224 K. The average daily winter temperature for this particular period is equal to 203 ± 10 K. For comparison, Aura/MLS temperatures, close to Kiruna location (less than 675 km away), are indicated by the red circles. There is a good agreement between these time series, with a small average temperature difference of 1.9 ± 6.9 K. The temperature behavior agrees well with a general seasonal temperature behavior in the winter mesopause having a maximum in the period of December-February followed by a rapid temperature decrease in March-April (e.g., Perminov et al.,

2018). Kim et al. (2017) have analyzed mesopause temperatures from OH airglow measurements with a Fourier Transform Spectrometer (FTS) for winter periods of 2003-2014 at Kiruna. The authors have found variations in the winter mesopause temperature above Kiruna being in the range of 180-230 K. This temperature range agrees well with the OH (3-1) temperature values presented in Fig. 7. Cho et al. (2011) have studied OH and O₂ airglow temperatures for winter periods of 2001-2007 measured with a Spectral Airglow Temperature Imager (SATI) instrument and a Michelson interferometer at high northern latitudes at Resolute Bay and Esrange, respectively. Cho et al. (2011) have demonstrated winter mesopause temperatures to vary between 180 and 245 K that again agrees well with the temperature range obtained by the OH imager. Sigernes et al. (2003) have studied daily mesospheric winter temperatures for more than 20 years of ground-based spectral measurements of the OH layer over Svalbard from 1980 to 2001. The authors have found the average daily winter temperature to be equal to 208±15 K, with the maximum and minimum temperatures of 257 K and 168 K, respectively, emphasizing that the mesospheric temperature variations over Svalbard in winter are extremely high. The average daily winter OH (3-1) temperature that was estimated in the present study is close to the value found by Sigernes et al. (2003), taking into account large temperature deviations of the average daily values found in both papers.

It is interesting to discuss the following feature in the OH (3-1) temperature behavior. There have been observed two rather strong temperature enhancements on day 52 (21 February 2023) and on day 66 (7 March 2023). The 2023 Sudden Stratospheric Warming (SSW) event started around 1st February 2023 (<https://www.climate.gov/news-features/event-tracker/disrupted-polar-vortex-brings-sudden-stratospheric-warming-february>). It is known that a SSW is accompanied by both cooling and warming of the mesopause region: cooling is observed at the time of a SSW onset or a few days later while a warming is observed with a delay of about 20-25 days (Cho et al., 2004; Shepherd et al., 2014). Sometimes several mesopause warmings (wavy oscillations) can be generated after a SSW onset (see Fig. 1b,f in Shepherd et al., 2014). This is exactly what was observed in the OH (3-1) temperature behavior: first, there was the OH (3-1) temperature decrease between days 32-36 (close to the 2023 SSW onset), then the two strong temperature maxima occurred on days 52 and 66, i.e., about 20 and 34 days after the 2023 SSW onset. A detailed link of the OH (3-1) temperature above Kiruna to SSW events will be investigated in future papers.

385

6 Conclusions

A novel infrared imager (OH imager) began operating in November 2022 in Kiruna (northern Sweden), aiming at studies of nightglow emissions coming from hydroxyl and molecular oxygen layers in the mesopause region. The OH imager registers selected emissions lines (1523 and 1542 nm) in the OH (3-1) band to obtain emission intensity and temperature maps at around 87 km altitude. Also, the imager records infrared emissions coming from the O₂ IR A-band airglow at 1269 nm in order to obtain O₂ emission intensity maps at a slightly higher altitude at ~94 km. This technique allows tracing gravity wave disturbances both in horizontal and vertical domains in the mesopause region. The lee location of the OH imager favors studies of mountain waves, their probable influence on the mesopause temperature and dynamics.

395

Temperature maps demonstrate the ability of the OH imager to resolve small-scale (4-8 km) small-amplitude (4-6 K) gravity waves as well as medium-scale gravity waves having horizontal wavelengths of 20-40 km and temperature amplitudes of 7-10 K. Such medium-scale waves have propagated through both the OH and O₂ layers.

400 The comparison between the OH imager and collocated Esrange lidar shows a good agreement between OH (3-1) and Esrange lidar temperature measurements on individual nights (within 3 K) as well as on an average basis (-0.2 ± 1.6 K). The comparison between the OH imager and Aura/MLS radiometer demonstrates no systematic temperature bias between these instruments as well as small average temperature difference of 2.8 ± 7.8 K, which is likely to be caused by the variable state of the mesopause temperature over spatial scales of their different observational volumes.

405 OH (3-1) daily mean rotational temperatures measured above Kiruna for the 2023 winter season (January-April) demonstrate typical seasonal variations with a maximum in February and a rapid temperature decrease in April in the winter polar mesopause. The OH (3-1) temperature values are within the range of 178-224 K that agrees well with other high latitude temperature measurements in 410 the winter mesopause. The average daily winter temperature for the considered period is 203 ± 10 K.

Code availability

Matlab codes used in this study for the OH image processing are available upon request from the corresponding author (pdalin@irf.se).

415

Data availability

Processed OH image data published in the paper's figures are available via the Swedish National Data Center (<https://snd.gu.se/en/catalogue/dataset/preview/d5776ec9-d346-4e6a-a8c3-1445e75201b6/1>).

420 Underlying research data are available upon request to Peter Dalin (pdalin@irf.se). Once the commissioning period of the OH imager instrument is completed, registrations will be made freely available for non-commercial scientific usage in accordance with other data obtained from the Kiruna Atmospheric and Geophysical Observatory (KAGO) at the Swedish Institute of Space Physics (IRF). Aura/MLS temperature data (ver.5.0 and level 2 data quality) can be obtained from the NASA public web-site: https://acdsc.gesdisc.eosdis.nasa.gov/data/Aura_MLS_Level2/. Esrange lidar data can be 425 made available upon request to Peter Dalin (pdalin@irf.se) or Jonas Hedin (jhedi@misu.su.se).

Author contributions

Author contributions. PD: conceptualization, methodology, installation of the OH imager, data processing and validation, matlab and python codes, statistical analysis, data visualization and archiving, optical calibration, writing the original manuscript. UB: conceptualization, writing a 430 proposal to the Swedish Research Council (Vetenskapsrådet), installation of the OH imager, providing data storage. JK: conceptualization, writing a proposal to the Swedish Research Council (Vetenskapsrådet). PV: installation of the OH imager, software development, lidar operation. TN: conceptualization, installation of the OH imager, writing python codes. TT, DW and CU: methodology, design, manufacturing and absolute calibration of the OH imager. VP and NP: conceptualization,

435 methodology, assisted with the interpretation of scientific results. JH: Esrange lidar preparation and operation. All co-authors contributed to writing and reviewing the paper.

Competing interests

The authors declare that they have no conflict of interest.

Acknowledgements

440 We are thankful to Daria Mikhaylova and Martin Rönnfalk (IRF-Kiruna) for their help with the installation of the OH imager.

Financial support

The OH imager installation was financed and performed as part of the research infrastructure grant 2021-00360 from the Swedish Research Council (Vetenskapsrådet) to the Kiruna Atmospheric and
445 Geophysical Observatory (KAGO) at the Swedish Institute of Space Physics (IRF).

References

- Ammosov, P., Gavriyeva, G., Ammosova, A., and Koltovskoi, I.: Response of the mesopause temperatures to solar activity over Yakutia in 1999-2013, *Adv. Space Res.*, 54(12), 2518-2524,
450 <http://dx.doi.org/10.1016/j.asr.2014.06.007>, 2014.
- Baker, D. J., and Stair, A. T.: Rocket measurements of the altitude distributions of the hydroxyl airglow, *Phys. Scr.*, 37, 611-622, <http://dx.doi.org/10.1088/0031-8949/37/4/021>, 1988.
- Bates, D. R., and Nicolet, M.: The photochemistry of atmospheric water vapour, *J. Geophys. Res.*, 55, 301-327, <http://dx.doi.org/10.1029/JZ055i003p00301>, 1950.
- 455 Blum, U., and Fricke, K. H.: The Bonn University lidar at the Esrange: technical description and capabilities for atmospheric research, *Ann. Geophys.*, 23, 1645-1658, <https://doi.org/10.5194/angeo-23-1645-2005>, 2005.
- Brooke, J. S. A., Bernath, P. F., Western, C. M., Sneden, C., Afşar, M., Li, G., and Gordon, I. E.: Line strengths of rovibrational and rotational transitions in the X²Π ground state of OH, *J. Quant. Spectrosc. Radiat. Transf.*, 168, 142-157, <https://doi.org/10.1016/j.jqsrt.2015.07.021>, 2016.
- 460 Chamberlain, J. W.: *Physics of the aurora and airglow*, 1st Edition, International Geophysics Series, 2, 722, Academic Press., ISBN: 9781483222530, 1961.
- Cho, Y.-M., Shepherd, G. G., Won, Y.-I., Sargoytchev, S., Brown, S., and Solheim, B.: MLT cooling during stratospheric warming events, *Geophys. Res. Lett.*, 31, L10104,
465 <https://doi.org/10.1029/2004GL019552>, 2004.
- Cho, Y.-M., Shepherd, G. G., Shepherd, M. G., Hocking, W. K., Mitchell, N. J., and Won, Y.-I.: A study of temperature and meridional wind relationships at high northern latitudes, *J. Atmos. Solar-Terr. Phys.*, 73, 936-943, <https://doi.org/10.1016/j.jastp.2010.08.011>, 2011.
- Christensen, A. B., Rees, M. H., Romick, G. J., and Sivjee, G. G.: OI (7774 Å) and OI (8446 Å)
470 emissions in aurora, *J. Geophys. Res. Space Phys.*, 83(A4), 1421-1425,
<https://doi.org/10.1029/JA083iA04p01421>, 1978.

- 475 Dalin, P., Pogoreltsev, A., Pertsev, N., Perminov, V., Shevchuk, N., Dubietis, A., Zalcik, M.,
Kulikov, S., Zadorozhny, A., Kudabayeva, D., Solodovnik, A., Salakhutdinov, G., and Grigoryeva, I.:
Evidence of the formation of noctilucent clouds due to propagation of an isolated gravity wave caused
by a tropospheric occluded front, *Geophys. Res. Lett.*, 42, 2037-2046,
<https://doi.org/10.1029/JA083iA04p0142110.1002/2014GL062776>, 2015.
- Dalin, P., Perminov, V., Pertsev, N., and Romejko, V.: Updated long-term trends in mesopause
temperature, airglow emissions, and noctilucent clouds, *J. Geophys. Res. Atmos.*, 125,
e2019JD030814, <https://doi.org/10.1029/2019JD030814>, 2020.
- 480 Demissie, T. D., Espy, P. J., Kleinknecht, N. H., Hatlen, M., Kaifler, N., and Baumgarten, G.:
Characteristics and sources of gravity waves observed in noctilucent cloud over Norway, *Atmos.*
Chem. Phys., 14, 12133-12142, <https://doi.org/10.5194/acp-14-12133-2014>, 2014.
- Dessler, A. E., Zhang, Z., and Yang, P.: Water-vapor climate feedback inferred from climate
fluctuations, 2003-2008, *Geophys. Res. Lett.*, 35, L20704, <https://doi.org/10.1029/2008GL035333>,
485 2008.
- Espy, P. J., and Stegman, J.: Trends and variability of mesospheric temperature at high latitudes,
Phys. Chem. Earth, 27, 543-553, [https://doi.org/10.1016/S1474-7065\(02\)00036-0](https://doi.org/10.1016/S1474-7065(02)00036-0), 2002.
- French, W. J. R., Mulligan, F. J., and Klekociuk, A. R.: Analysis of 24 years of mesopause region
OH rotational temperature observations at Davis, Antarctica - Part I: long-term trends, *Atmos. Chem.*
490 *Phys.*, 20, 6379-6394, <https://doi.org/10.5194/acp-20-6379-2020>, 2020.
- Fritts, D. C., and Alexander, M. J.: Gravity wave dynamics and effects in the middle atmosphere,
Rev. Geophys., 41(1), 1003, <https://doi.org/10.5194/acp-14-12133-201410.1029/2001RG000106>,
2003.
- Froidevaux, L., Livesey, N. J., Read, W. G., Jiang, Y. B., Jiménez, C. C., Filipiak, M. J., Schwartz,
495 M. J., et al.: Early validation analyses of atmospheric profiles from EOS MLS on the Aura satellite,
IEEE Transactions on Geoscience and Remote Sensing, 44(5), 1106-1121, 2006.
- Gao, H., Xu, J. Y., Chen, G. M., Yuan, W., and Beletsky, A. B.: Global distributions of OH and O₂
(1.27 μm) nightglow emissions observed by TIMED satellite, *Sci. China Technol. Sci.*, 54, 447-456,
<https://doi.org/10.1007/s11431-010-4236-5>, 2011.
- 500 Garcia, F. J., Taylor, M. J., and Kelley, M. C.: Two-dimensional spectral analysis of mesospheric
airglow image data, *Appl. Opt.*, 36, 7374-7385, <https://doi.org/10.1364/AO.36.007374>, 1997.
- Gossard, E. E., and Hook, W. H.: *Waves in the atmosphere*, Elsevier Scientific Publishing
Company, Amsterdam, 532 p., ISBN 978-0444411969, 1975.
- Grygalashvyly, M., Sonnemann, G. R., Lübken, F.-J., Hartogh, P., and Berger, U.: Hydroxyl layer:
505 mean state and trends at midlatitudes, *J. Geophys. Res. Atmos.*, 119, 12391-12419,
<http://dx.doi.org/10.1002/2014JD022094>, 2014.
- Kalicinsky, C., Knieling, P., Koppmann, R., Offermann, D., Steinbrech, W., and Wintel, J.: Long-
term dynamics of OH* temperatures over central Europe: trends and solar correlations, *Atmos. Chem.*
Phys., 16, 15033-15047, <http://dx.doi.org/10.5194/acp-16-15033-2016>, 2016.

- 510 Khomich, V. Yu., Semenov, A. I., and Shefov, N. N.: Airglow as an indicator of upper atmospheric structure and dynamics, Springer-Verlag, Berlin, Heidelberg. <http://dx.doi.org/10.1007/978-3-540-75833-4>, 2008.
- Kim, G., Kim, J.-H., Kim, Y. H., and Lee, Y. S.: Long-term trend of mesospheric temperatures over Kiruna (68°N, 21°E) during 2003-2014, *J. Atmos. Sol.-Terr. Phys.*, 161, 83-87, <https://doi.org/10.1016/j.jastp.2017.06.018>, 2017.
- 515 Lübken, F.-J., Berger, U., and Baumgarten, G.: On the anthropogenic impact on long-term evolution of noctilucent clouds, *Geophys. Res. Lett.*, 45, 6681-6689, <https://doi.org/10.1029/2018GL077719>, 2018.
- Makhlouf, U. B., Picard, R. H., and Winick, J. R.: Photochemical-dynamical modeling of the measured response of airglow to gravity waves: 1. Basic model for OH airglow, *J. Geophys. Res.*, 100, 11,289-11,311, <https://doi.org/10.1029/94JD03327>, 1995.
- 520 Mlynczak, M. G., Hunt, L. A., Garcia, R. R., Harvey, V. L., Marshall, B. T., Yue, J., Mertens, C. J., and Russell, J. M.: Cooling and contraction of the mesosphere and lower thermosphere from 2002 to 2021, *J. Geophys. Res. Atmos.*, 127, e2022JD036767, <https://doi.org/10.1029/2022JD036767>, 2022.
- 525 Melo, S. M. L., Lowe, R. P., and Russell, J. P.: Double-peaked hydroxyl airglow profiles observed from WINDII/UARS, *J. Geophys. Res.*, 105(D10), 12,397-12,403, <https://doi.org/10.1029/1999JD901169>, 2000.
- Nishiyama, T., Taguchi, M., Suzuki, H., Dalin, P., Ogawa, Y., Brändström, U., and Sakanoi, T.: Temporal evolutions of N₂⁺ Meinel (1,2) band near 1.5 μm associated with aurora breakup and their effects on mesopause temperature estimations from OH Meinel (3,1) band, *Earth, Planets and Space*, 73(30), <https://doi.org/10.1186/s40623-021-01360-0>, 2021.
- 530 Offermann, D., Gusev, O., Donner, M., Forbes, J. M., Hagan, M., Mlynczak, M. G., Oberheide, J., Preusse, P., Schmidt, H., and Russell III, J. M.: Relative intensities of middle atmosphere waves, *J. Geophys. Res.*, 114, D06110, <https://doi.org/10.1029/2008JD010662>, 2009.
- 535 Offermann, D., Hoffmann, P., Knieling, P., Koppmann, R., Oberheide, J., Steinbrecht, W.: Long-term trend and solar cycle variations of mesospheric temperature and dynamics, *J. Geophys. Res.*, 115, D18127, <https://doi.org/10.1029/2009JD013363>, 2010.
- Pautet, P.-D., Taylor, M. J., Liu, A. Z., and Swenson, G. R.: Climatology of short-period gravity waves observed over northern Australia during the Darwin area wave experiment (DAWEX) and their dominant source regions, *J. Geophys. Res.*, 110, D03S90, <https://doi.org/10.1029/2004JD004954>, 2005.
- 540 Pautet, P.-D., Stegman, J., Wrasse, C. M., Nielsen, K., Takahashi, H., Taylor, M. J., et al.: Analysis of gravity waves structures visible in noctilucent cloud images, *J. Atmos. Sol. Terr. Phys.*, 73(14-15), 2082-2090, <https://doi.org/10.1016/j.jastp.2010.06.001>, 2011.
- 545 Pautet, P.-D., Taylor, M. J., Pendleton, W. R., Zhao, Jr., Y., Yuan, T., Esplin, R., and McLain, D.: Advanced mesospheric temperature mapper for high-latitude airglow studies, *Applied Optics*, 53(26), <http://dx.doi.org/10.1364/AO.53.005934>, 2014.
- Perminov, V. I., Lowe, R. P., and Pertsev, N. N.: Longitudinal variations in the hydroxyl nightglow, *Adv. Space Res.*, 24(11), 1609-1612, [https://doi.org/10.1016/S0273-1177\(99\)00887-X](https://doi.org/10.1016/S0273-1177(99)00887-X), 1999.

- 550 Perminov, V. I., Semenov, A. I., Pertsev, N. N., Medvedeva, I. V., Dalin, P. A., and Sukhodoev, V. A.: Multi-year behaviour of the midnight OH* temperature according to observations at Zvenigorod over 2000-2016, *Adv. Space Res.*, 61(7), 1901-1908, <http://dx.doi.org/10.1016/j.asr.2017.07.020>, 2018.
- Pertsev, N., and Perminov, V.: Response of the mesopause airglow to solar activity inferred from measurements at Zvenigorod, Russia. *Ann. Geophys.*, 26(5), 1049-1056,
555 <http://dx.doi.org/10.5194/angeo-26-1049-2008>, 2008.
- Pertsev, N. N., Dalin, P. A., and Perminov, V. I.: Lunar tides in the mesopause region obtained from summer temperature of the hydroxyl emission layer, *Geomagn. Aeron.*, 61(2), 259-265,
<https://doi.org/10.1134/S0016793221020109>, 2021.
- Plougouven, R., and Zhang, F.: Internal gravity waves from atmospheric jets and fronts, *Rev.*
560 *Geophys.*, 52, 33-76, <https://doi.org/10.1002/2012RG000419>, 2014.
- Reid, I. M., Spargo, A. J., and Woithe, J. M.: Seasonal variations of the nighttime O(1S) and OH(8-3) airglow intensity at Adelaide, Australia, *J. Geophys. Res. Atmos.*, 119, 6991-7013,
<https://doi.org/10.1002/2013JD020906>, 2014.
- Roble, R. G., and Dickinson, R. E.: How will changes in carbon dioxide and methane modify the
565 mean structure of the mesosphere and thermosphere?, *Geophys. Res. Lett.*, 16, 1441-1444,
<https://doi.org/10.1029/GL016i012p01441>, 1989.
- Savigny, C.: Airglow in the Earth atmosphere: basic characteristics and excitation mechanisms, *ChemTexts*, 3(14), <https://doi.org/10.1007/s40828-017-0051-y>, 2017.
- Schwartz, M. J., Lambert, A., Manney, G. L., Read, W. G., Livesey, N. J., Froidevaux, L., Ao, C.
570 O., Bernath, P. F., Boone, C. D., Cofield, R. E., et al.: Validation of the Aura Microwave Limb
Sounder temperature and geopotential height measurements, *J. Geophys. Res.*, 113, D15S11,
<https://doi.org/10.1029/2007JD008783>, 2008.
- Semenov, A. I., and Shefov, N. N.: Empirical model of hydroxyl emission variations, *Int. J.*
Geomagn. Aeron., 1, 229-242, <https://doi.org/10.1134/S0016793207010161>, 1999.
- 575 Shepherd, M. G., Beagley, S. R., and Fomichev, V. I.: Stratospheric warming influence on the
mesosphere/lower thermosphere as seen by the extended CMAM, *Ann. Geophys.*, 32, 589-608,
<https://doi.org/10.5194/angeo-32-589-2014>, 2014.
- Sigernes, F., Shumilov, N., Deehr, C. S., Nielsen, K. P., Svenøe, T., and Havnes, O.: Hydroxyl
rotational temperature record from the auroral station in Adventdalen, Svalbard (78°N, 15°E), *J.*
580 *Geophys. Res.*, 108(A9), 1342, <https://doi.org/10.1029/2001JA009023>, 2003.
- Sivjee, G. G., Romick, G. J., and Rees, M. H.: Intensity ratio and center wavelengths of [O II]
(7320-7330 Å) line emissions, *The Astrophysical Journal*, 229, 432-438,
<https://doi.org/10.1086/156971>, 1979.
- Sivjee, G. G., and Hamwey, R. M.: Temperature and chemistry of the polar mesopause OH, *J.*
585 *Geophys. Res.*, 92, 4663-4672, <https://doi.org/10.1029/JA092iA05p04663>, 1987.
- Suzuki, H., Shiokawa, K., Tsutsumi, M., Nakamura, T., and Taguchi, M.: Atmospheric gravity
waves identified by ground-based observations of the intensity and rotational temperature of OH
airglow, *Polar Science*, 2, 1-8, <https://doi.org/10.1016/j.polar.2007.12.002>, 2008.



- Taylor, J. R.: An introduction to error analysis: the study of uncertainties in physical measurements,
590 Second Edition, University Science Books, California, 327 p., ISBN 9780935702422, 1997.
- Taylor, M. J., Pendleton Jr., W. R., Clark, S., Takahashi, H., Gobbi, D., and Goldberg, R. A.: Image
measurements of short-period gravity waves at equatorial latitudes, *J. Geophys. Res.*, 102(D22),
26,283-26,299, <https://doi.org/10.1029/96JD03515>, 1997.
- Thomas, G. E., and Olivero, J.: Noctilucent clouds as possible indicators of global changes in the
595 mesosphere, *Adv. Space Res.*, 28(7), 937-946, [https://doi.org/10.1016/S0273-1177\(01\)80021-1](https://doi.org/10.1016/S0273-1177(01)80021-1), 2001.



Table 1. The main characteristics of the OH imager in Kiruna.

Parameter	Value
CCD image sensor	InGaAs focal plane array
CCD format/ used for image processing	640 x 512 / 512 x 512 pixels
CCD image area	12.80 x 10.24 mm
Pixel size	20 x 20 μm
Pixel full well high capacity at 2 MHz	677,800 e ⁻
Quantum efficiency	~85%
Dark current at 2 MHz at -80°C	21.1 e ⁻ /pix/sec
Temperature of detector using a Peltier cooler	-80°C at ambient temperature 21°C
Digitization	16 bit
Field of view	120°
Filters diameter	76.2 mm
Narrow-band three-cavity filters centered at:	
OH(3-1) background ($\Delta\lambda$ FWHM)	1521.0 nm ($\Delta\lambda=2.2$ nm)
OH(3-1) P1(2) line ($\Delta\lambda$ FWHM)	1523.7 nm ($\Delta\lambda=2.2$ nm)
OH(3-1) P1(4) line ($\Delta\lambda$ FWHM)	1542.8 nm ($\Delta\lambda=2.1$ nm)
O ₂ IR A-band line	1268.7 nm ($\Delta\lambda=2.9$ nm)
Exposure times for OH and O ₂ filters	30 and 40 s
Time resolution of the whole cycle	165 s
Spatial resolution of OH measurements at 87 km	0.3 km in the image center
Horizontal coverage at 87 km	300 x 300 km



600 **Figures:**

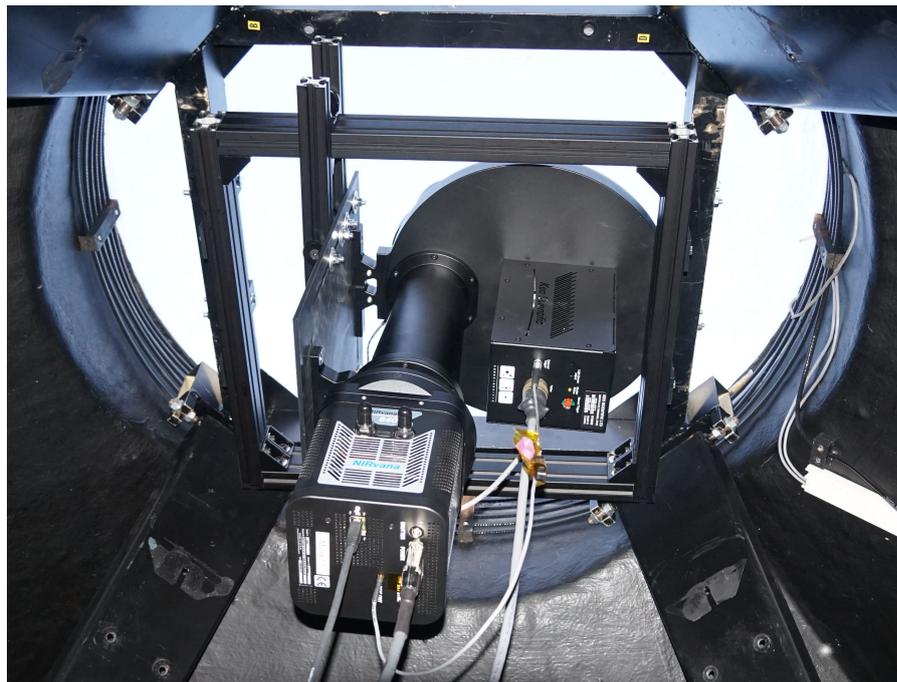
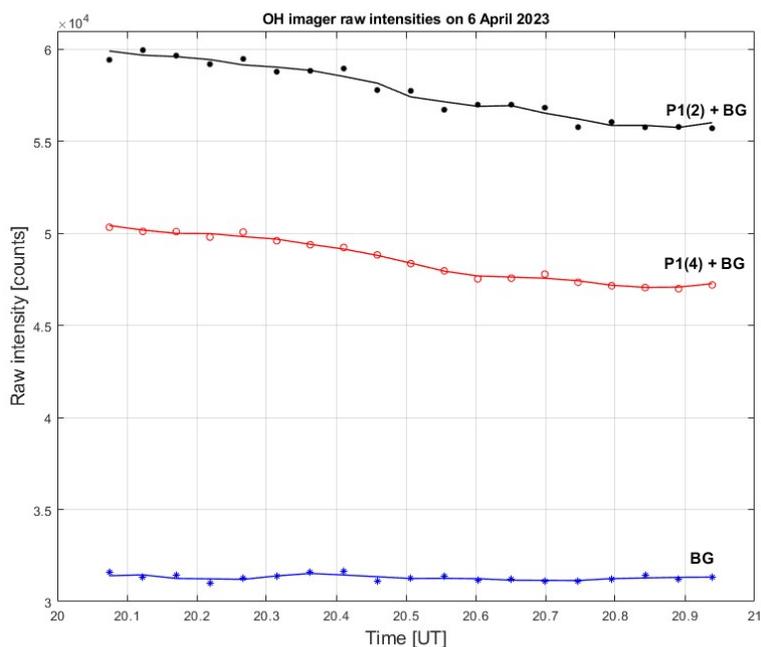


Figure 1: The Keo Scientific OH imager installed in November 2022 at the Swedish Institute of Space Physics (IRF), located in Kiruna, Sweden (67.86°N, 20.42°E). The image was taken by Peter Dalin.

605



610 **Figure 2:** Raw intensities of emission lines at the zenith as measured by the OH imager in Kiruna on 6 April 2023. The black dots, red circles and blue asterisks are raw intensities of the OH P₁(2), P₁(4) lines and atmospheric background (BG), respectively. The black, red and blue lines are 3-point moving averages.

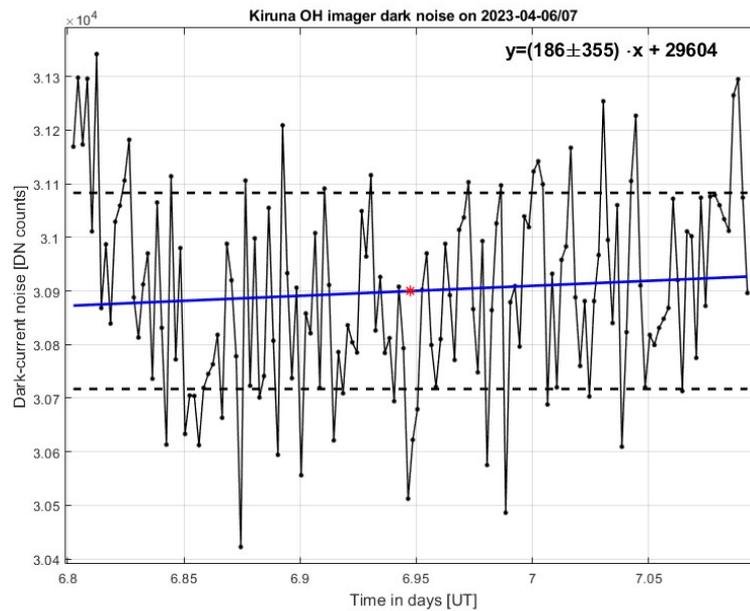
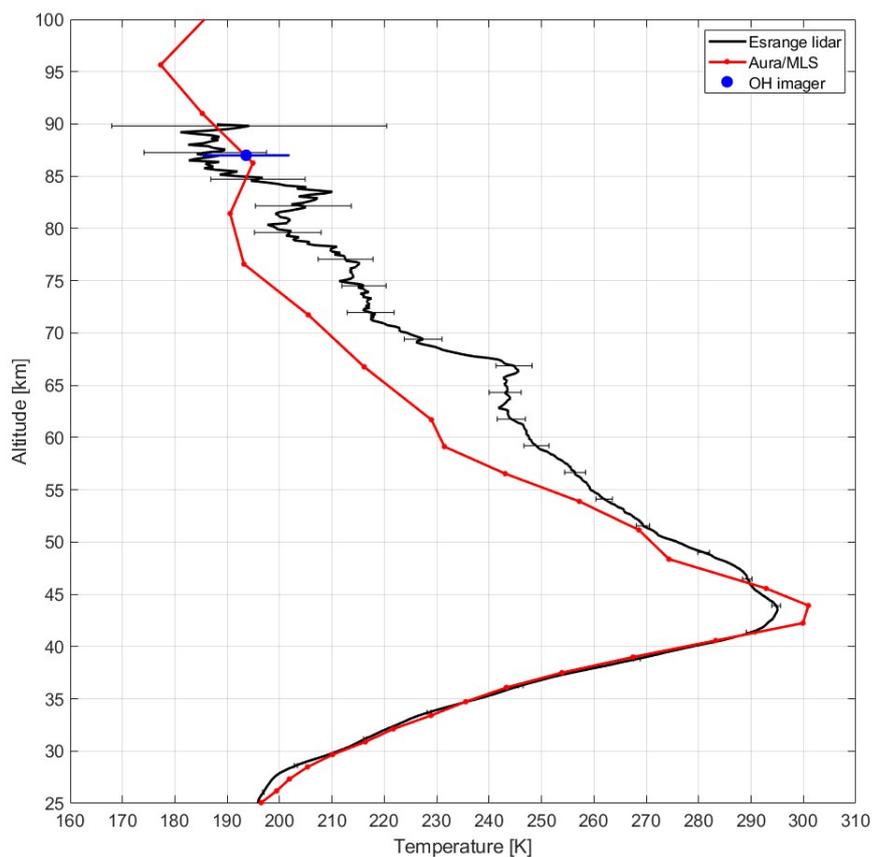
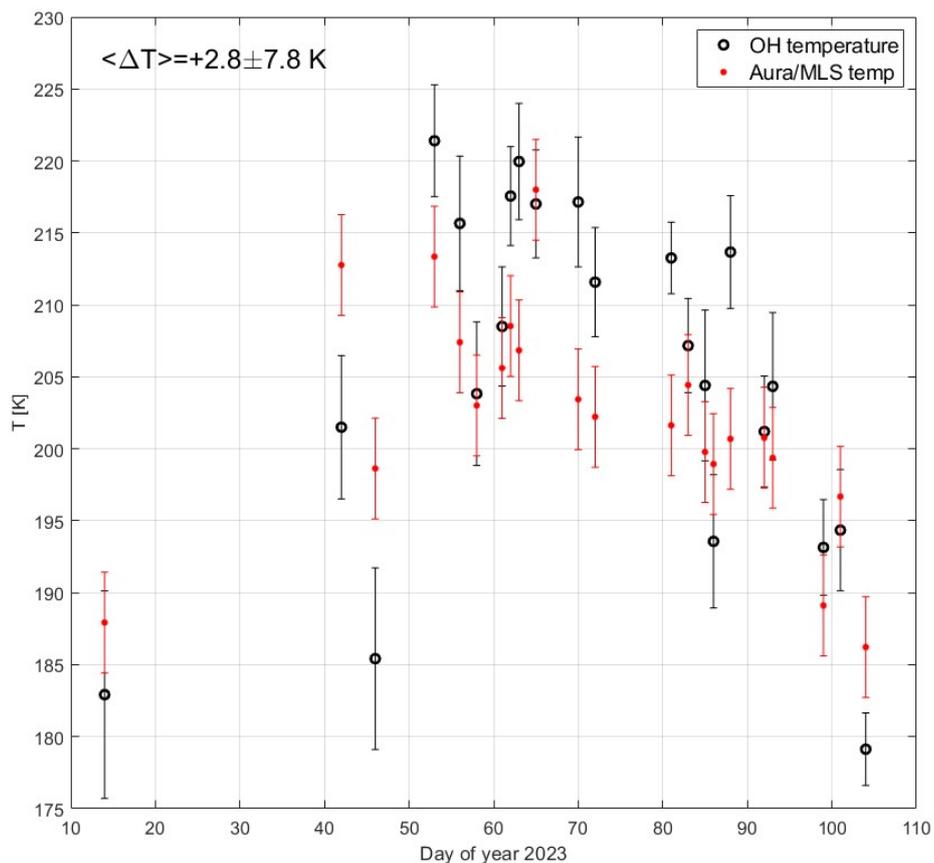


Figure 3: Example of the dark-current noise of the OH imager in Kiruna on the night of 6-7 April 2023. The black dots are dark-current noise counts in the detector center, the blue line is the linear regression, the red asterisk is the mean dark-current noise and the dashed lines are one standard deviation around the mean noise.

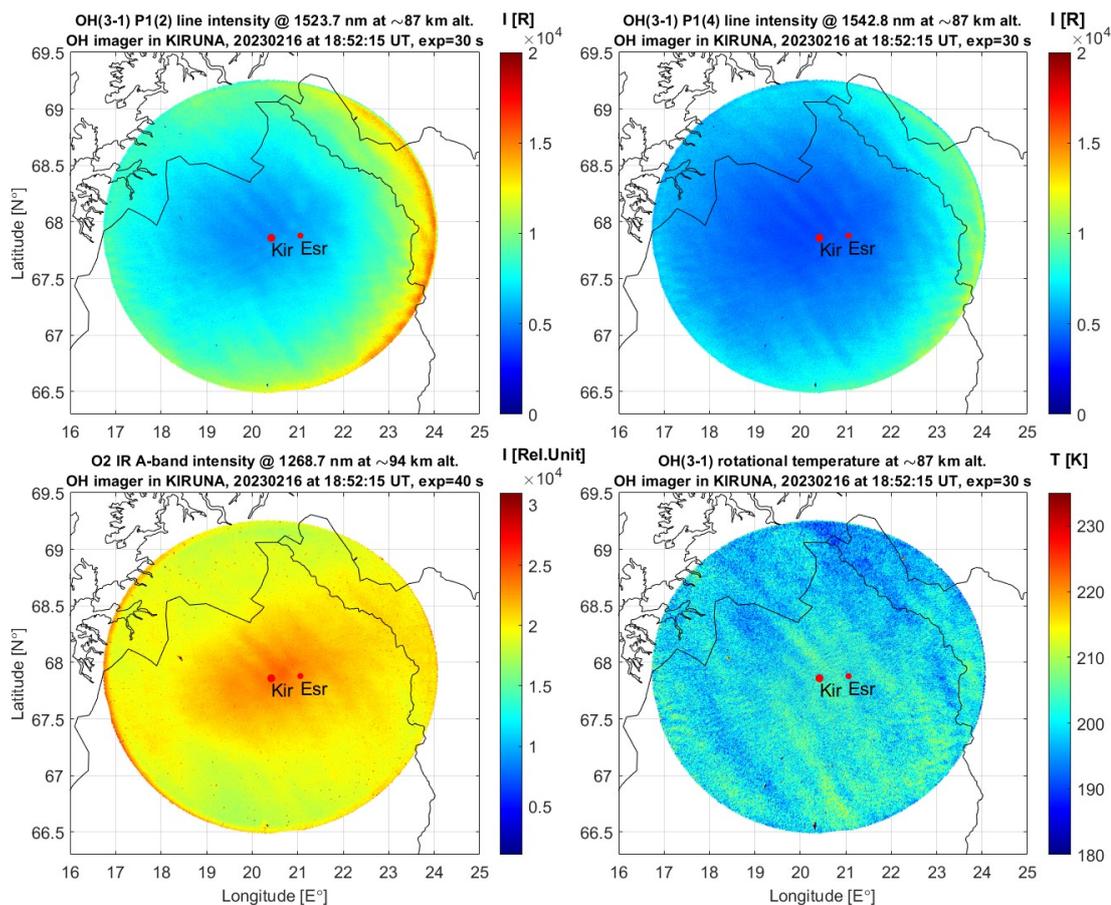


620 **Figure 4: Temperature measurements made between 18:43 and 05:58 UT on the night 2/3 February 2023. The blue dot is the mean OH (3-1) temperature at 87 km averaged over the same time period. The black line is the Esrange lidar mean temperature profile. Error bars for the OH imager and lidar measurements are one standard deviation of the temperature variations for the night. The red line is the Aura/MLS temperature profile which was chosen at 70.3°N 12.2°E, closest to the position of Kiruna (about 420 km away from Kiruna), measured at 02:32 UT on 3 February 2023.**

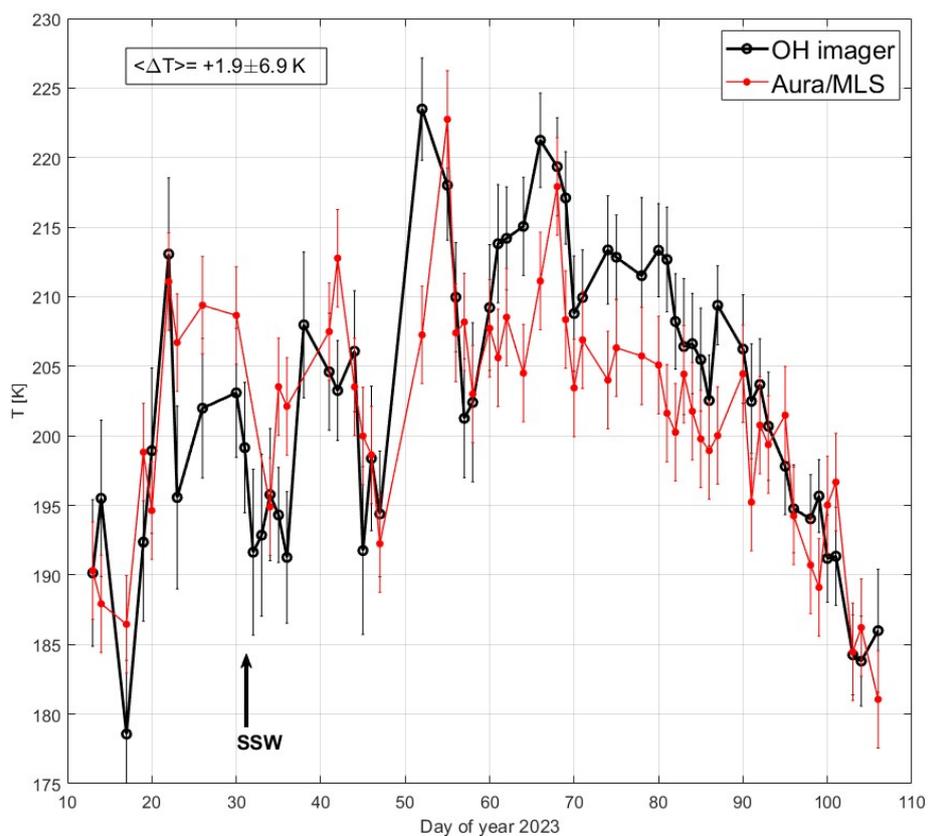


625 **Figure 5: A comparison between simultaneous measurements of the mesopause temperature performed by the OH imager (black circles) above Kiruna at ~87 km and Aura/MLS (red points) at 0.0046 hPa, ~86 km, close to the Kiruna location (distance difference is less than 300 km). Error bars of the OH (3-1) temperature are one standard deviations for a particular hour. Error bars of the Aura/MLS temperature are equal to 3.5 K (see the text).**

630



635 **Figure 6: Maps of airglow intensities and temperature obtained by the OH imager on 16 February 2023. (A) Intensity in the Rayleigh of the OH (3-1) P1(2) line at 1523.7 nm. (B) Intensity in the Rayleigh of the OH (3-1) P1(4) line at 1542.8 nm. (C) Intensity in relative units of the O₂ IR A-band at 1268.7 nm. (D) OH (3-1) rotational temperature estimated by the brightness ratio of the two OH (3-1) emission lines (see the text).**



640 **Figure 7: Daily mean OH (3-1) rotational temperature above Kiruna, at zenith, at ~87 km for the 2023 winter season (black circles). Aura/MLS temperature instantaneous measurements close to the Kiruna location at ~86 km (red points). Error bars of the OH (3-1) temperature are one standard deviation for a particular night. Error bars of the Aura/MLS temperature are equal to 3.5 K (see the text). The onset of the sudden stratospheric warming is marked by the black arrow. The considered period is from 13 January until 16 April 2023.**