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3	Detection of small drizzle droplets in a large cloud chamber using
4	ultra high-resolution radar
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Abstract

A large convection cloud chamber has the potential to produce drizzle-sized droplets, thus offering a new opportunity to investigate aerosol-cloud-drizzle interactions at a fundamental level under controlled environmental conditions. One key measurement requirement is the development of methods to detect the low concentration drizzle drops in such a large cloud chamber. In particular, remote sensing methods may overcome some limitations of in situ methods.

Here, the potential of an ultra-high-resolution radar to detect the radar return signal of a small drizzle droplet against the cloud droplet background signal is investigated. It is found that using a small sampling volume is critical to drizzle detection in a cloud chamber to allow a drizzle drop in the radar sampling volume to dominate over the background cloud droplets signal. For instance, a radar volume of 1 cubic centimeter (cm^3) would enable the detection of drizzle embryos with diameter larger than 40 μm . However, the probability of drizzle sampling also decreases as the sample volume reduces, leading to a longer observation time. Thus, the selection of radar volume should consider both of the signal power and the drizzle occurrence probability. Finally, observations from the Pi Convection-Cloud Chamber are used to demonstrate the single drizzle particle detection concept using small radar volume. The results presented in this study also suggest new applications of ultra-high-resolution cloud radar for atmospheric sensing.

1. Introduction

Drizzle formation is one of the most important microphysical processes in warm clouds. Yet the processes controlling drizzle formation remain poorly understood (Wood, 2012). The most challenging aspect is the initial formation of drizzle embryos with diameter around $30 \, \mu m \sim 50 \, \mu m$. The formation of small drizzle particles in this range can neither be adequately explained by the traditionally-defined condensation growth process nor by the traditionally-defined collision-coalescence (C-C) process owing to their low efficiency (Yau and Rogers, 1996;Pruppacher and Klett, 2012;Falkovich et al., 2006;Beard and Ochs III, 1993). Several mechanisms have been hypothesized to explain the efficiency of these processes including i) fine-scale turbulence in cloud (Pinsky and Khain, 1997;Shaw, 2003); ii) giant cloud condensation nuclei (GCCN) (Johnson, 1982;Feingold et al., 1999); and iii) longwave cooling (Roach, 1976;Harrington et al., 2000). Nevertheless, it remains unclear to which extent these proposed mechanisms can adequately explain the origin of drizzle embryos.

One main barrier that hinder our ability to investigate the drizzle initiation process is the lack of observations with sufficient sensitivity and spatiotemporal resolution to detect the early growth of drizzle particles. As such an instrumented large convection cloud chamber with well-controlled initial and boundary conditions might help to improve our understanding of the drizzle initiation mechanism (Shaw et al., 2020). Unlike other types of chambers, a convection-cloud chamber can generate a steady state cloud system for hours in a turbulent environment by maintaining a warm saturated bottom surface, a cold saturated top surface, and a constant aerosol injection rate (Chang et al., 2016). The Michigan Tech Pi convection chamber with a dimension of 2 m x 2 m x 1 m (width x depth x height) has been used to explore aerosol-cloud-turbulence interactions, however, the Pi Chamber is too small to initiate drizzle embryos mainly due to the relatively short lifetime of cloud droplets therein. Results from large eddy simulations indicate that drizzle can be initiated in a large convection-cloud chamber with a height on the order of 10 m (Thomas et al., 2023). However, the drizzle drops are sparse in a large chamber, so the detection of single drizzle embryos in a large cloud chamber is challenging for in-situ probes that generally have a sampling volume of only a few cubic centimeters. On the other hand, active remote sensors have the ability to rapidly sample large volumes and thus offer an attractive option for the detection of small drizzle droplets in a cloud chamber.

Here, we will demonstrate that the detection of an individual drizzle droplet in the presence of numerous cloud droplets is possible with a radar that can achieve a very small sampling volume. The detection of individual drizzle droplets is possible because the radar signal-to-noise ratio (SNR) of a point target (drizzle droplet) is not affected by the radar observational volume, while the SNR of a distributed target (cloud droplets) scales with the radar volume. In the following sections, the detection limits of an individual drizzle particle are investigated using idealized particle size distributions and real particle size distributions from the Michigan Tech Pi Chamber. In the end, the potential of THz radars offering unprecedented sub-centimeter range resolution will be discussed for developing the single drizzle detection radar (Cooper and Chattopadhyay, 2014).

2. Drizzle detection using radar

The detection of early drizzle particles in clouds has been the topic of extensive research. First, the radar needs to have sufficient sensitivity to detect cloud and drizzle droplets. This is typically accomplished using millimeter-wavelength radar (Kollias et al., 2007). Early methodologies for the detection of drizzle drop in clouds employ the use of reflectivity thresholds, ranging from –15 to –20 dBZ, to identify drizzle existence (Frisch et al., 1995;Liu et al., 2008;Comstock et al., 2004). Kollias et al. (2011) introduced the use of the radar Doppler spectra skewness as a more sensitive method for detecting the presence of small drizzle droplets (Acquistapace et al., 2017;Zhu et al., 2022). The radar Doppler spectra technique improved the detection of drizzle droplets that can produce as low as -30 dBZ (Zhu et al., 2022).

However, the use of the radar Doppler spectra technique in a cloud chamber is challenging. First, this will require that the radar point vertically to take advantage of the differential velocity between cloud and drizzle droplets. If we assume a monodisperse droplet size distribution (DSD) and Rayleigh scattering conditions, a drizzle detection limit of -30 dBZ is equivalent to a concentration of 10^{-3} cm⁻³ of drizzle droplets with diameter equal to $100 \mu m$ or a concentration of 6.4×10^{-2} cm⁻³ of drizzle droplets with 50 μm diameter. In the former case, the drizzle particle size is quite

large and not quite an early drizzle droplet detection. In the latter case, the concentration of drizzle droplets is much higher than the concentration observed in nature ($\sim 10^{-4}~cm^{-3}$) (Zhu et al., 2022). Furthermore, the conventional cloud radar has range resolution of tens of meters, which is not applicable in a chamber facility which may be on the order of several meters (approaching the collision mean free path).

As a result, we consider alternative methods to increase the probability of early drizzle droplet detection against the cloud droplet signal. As the number concentration of drizzle particle is low, by applying a small radar sampling volume V_{Radar} , it is possible that only one drizzle droplet is presence in V_{Radar} . In this case, the drizzle particle can be considered as a point target with backscattering cross section $\sigma(m^2)$ and the received radar echo power P_r (mW) is commonly expressed as (Battan, 1973):

$$P_{r,drizzle} = P_t \frac{G^2 \lambda^2}{(4\pi)^3 r^4} \sigma(D_d)$$
 (1)

where P_t is the transmit peak power (mW), G is the antenna gain, r (m) is the range of the target relative to the radar receiver and λ (m) is the radar wavelength. It is noteworthy that P_r for a point target does not depend on the radar sampling volume V_{Radar} . For distributed targets such as a cloud droplet population described by a droplet size distribution (DSD) that represents the number concentration of cloud droplets as a function of diameter, the received radar echo power is given by:

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$$P_{r,cloud} = P_t \frac{G^2 \lambda^2}{(4\pi)^3 r^4} \cdot V_{Radar} \cdot \sum_{i=0}^n N_c(D_i) \sigma(D_i) \Delta D_i$$
 (2)

Where n is the number of cloud droplets in the radar volume and $N_c(D)$ is the DSD with units of m^{-4} . In this case, the received radar echo power depends on the radar sampling volume, which is given by the following expression:

$$V_{Radar} = \pi \left(\frac{r\theta_{3dB}}{2}\right)^2 \cdot \Delta R \tag{3}$$

where θ_{3dB} is antenna radiation pattern 3-dB beamwidth in radians and ΔR is the range resolution.

Assuming Rayleigh scattering, the backscatter cross section of the drizzle and cloud droplets is

proportional to the sixth power of the particle diameter and inversely proportional to the fourth

power of the wavelength $(\sigma(D) \sim D^6/\lambda^4)$. Combing Eq. 1 and Eq. 2, the ratio of radar received

echo power from drizzle and cloud is given by the following expression:

is equal to that of the background echo contributed from cloud droplets.

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$$\frac{Signal}{Background} = \frac{P_{r,drizzle}}{P_{r,cloud}} = \frac{1}{V_{Radar}} \cdot \frac{D_d^6}{\sum_{i=0}^n N_c(D_i) D_i^6 \Delta D_i}.$$
 (4)

Eq. 4 indicates that the probability of detecting a single drizzle droplet in the radar sampling volume increases inversely to the radar sampling volume (point vs distributed target).

3. Detection requirement

Here, we will evaluate how small the radar sampling volume needs to be to detect drizzle drops with different diameters against three background (cloud) conditions: i) monodisperse cloud DSD, ii) cloud DSD from a theoretical model and iii) observed cloud DSD from the Michigan Tech Pi Chamber. For simplicity, we will assume that a drizzle drop is detectable if its radar return power

3.1. Monodisperse cloud DSD

We first construct an idealized scenario by considering two categories of droplets, i.e. cloud droplet with a diameter of D_c and a single drizzle drop with a diameter of D_d , the number concentration of cloud droplets in the radar sampling volume is N_c (m⁻³).

In this case, Eq. 4 is simplified as:

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$$\frac{Signal}{Background} = \frac{1}{V_{Radar}} \cdot \frac{D_d^6}{N_c \cdot D_c^6}$$
 (5)

When the signal power equals to the background, the radar sampling volume enabling single drizzle particle detection is estimated as a function of the size ratio $x = \frac{D_d}{D_c}$ shown in Fig. 1. The results are shown for various cloud droplet concentrations. It is noted that the required radar volume for detection depends on the drizzle drop size and the cloud number concentration. Larger radar volume would be required for drizzle detection as the particle size ratio increase; for a given particle size ratio, decreasing cloud number concentration can enhance the required radar volume. For example, if the cloud number concentration is 50 cm^{-3} and the mean cloud diameter (D_c) is $20 \mu m$, then the detection of a drizzle particle with diameter of $40 \mu m$ (x = 2) requires radar volume around 1 cm^3 . Such sampling volume are not achievable with traditional radar systems that employ sampling volumes of the order of 1000 m^3 or more (Kollias et al., 2016).

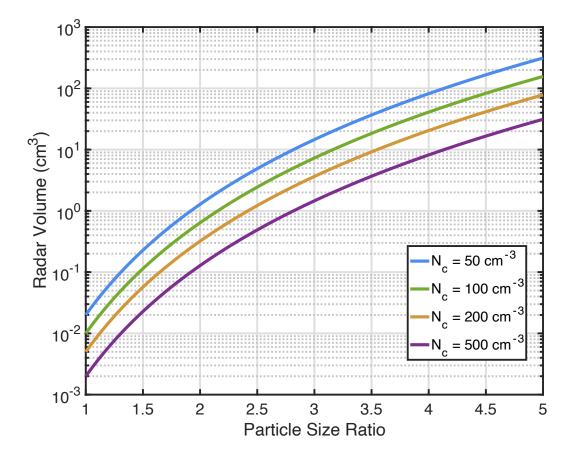


Figure 1: Radar observational volume for single-drizzle-drop detection as a function of particle size ratio $x = \frac{D_d}{D_c}$. Lines of different color represent clouds number concentration (N_c) : 50 cm⁻³ (blue), 100 cm⁻³ (green), 200 cm⁻³ (yellow) and 500 cm⁻³ (purple).

3.2 Drizzle detection against an idealized cloud droplet background

In a realistic cloud chamber environment, we expect a population of cloud droplets with various sizes that can be represented by a DSD. Particularly, when condensation and fallout are the main sources and sinks for the evolution equation for the DSD, the DSD in the cloud chamber can be approximately described by theoretically derived distributions (Saito et al., 2019; Chandrakar et al., 2020; Krueger, 2020). Here we adapt the theoretical DSD formula derived by Krueger (2020) to investigate the ability of a radar to detect a drizzle embryo present in a small sample volume under different chamber environment conditions. To better represent the cloud DSD under different environments, the analytical DSD is rearranged to be expressed as a function of liquid water content (LWC_c; g m⁻³) and number concentration (N_c ; m⁻³) as:

$$N(D_c) = \frac{2N_c D_c}{\pi^{1/2}} \left(\frac{4\Gamma\left(\frac{5}{4}\right) \pi^{\frac{1}{2}} \rho_l N_c}{3LW C_c} \right)^{2/3} \exp\left(-\left(\frac{4\Gamma\left(\frac{5}{4}\right) \pi^{\frac{1}{2}} \rho_l N_c}{3LW C_c} \right)^{\frac{4}{3}} \left(\frac{D_c}{2} \right)^4 \right)$$
(6)

where ρ_l is liquid water density $(g \ m^{-3})$, and D_c is cloud droplet diameter (m). $N(D_c)$ represents the number concentration of cloud droplet for the given diameter $(cm^{-3} \ \mu m^{-1})$.

Here we define the minimal drizzle drop $(D_{d,min})$ as the size of particle with radar return power equal to the total return power from cloud droplets in a given radar volume (V). Given the cloud DSD described by Eq. 6, $D_{d,min}$ can be estimated as:

$$D_{d,min}^6 = \int VN(D_c)D_c^6 dD \tag{7}$$

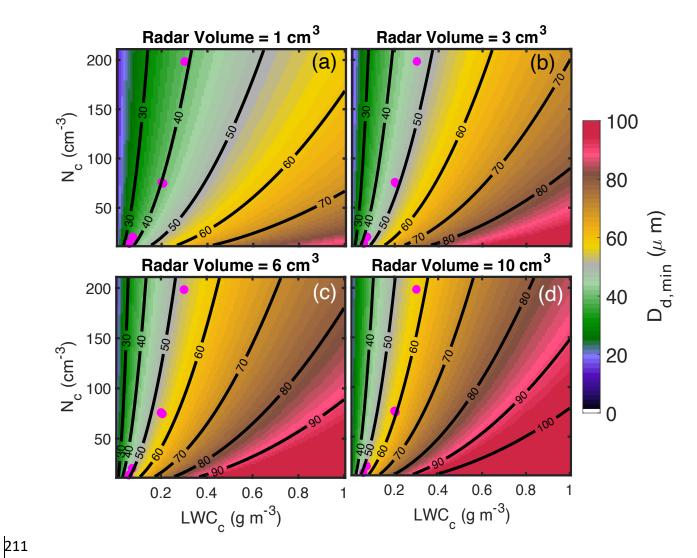


Figure 2: The minimal detectable drizzle particle ($D_{d,min}$) under different LWC_c and N_c conditions with radar sampling volume of (a) 1 cm^3 , (b) 3 cm^3 , (c) 6 cm^3 and (d) 10 cm^3 . The black lines are the $D_{d,min}$ contour of 30 μm , 40 μm , 50 μm , 60 μm , 70 μm . The magenta dots indicate the LWC_c and N_c observed in the Pi-cloud chamber.

Fig. 2 illustrates $D_{d,min}$ under different LWC_c and N_c combinations for various radar volumes. For a given steady-state cloud in a convection chamber (i.e., fixed LWC_c and N_c), $D_{d,min}$ generally increases as the radar volume increases. This is because larger radar volumes contain more cloud droplets that produce stronger background power, thus only a larger drizzle particle with a higher backscattering signal would be detectable. On the other hand, for a given radar observational volume, $D_{d,min}$ is jointly determined by LWC_c and N_c which are inversely proportional. As such,

 $D_{d,min}$ increases rapidly with increasing LWC_c but slightly decreases with increasing N_c . This contrasting relationship is caused by a larger sensitivity of radar reflectivity to particle size than to number concentration. Thus, increasing LWC_c can increase mean cloud particle size and greatly enhance the background power, leading to a larger detectable $D_{d,min}$. On the other side, when LWC_c is fixed, increasing cloud total number concentration tends to decrease particle size. The reduced cloud particle size would reduce the backscattering power and more than compensate for the power enhancement contributed from the increased number concentration.

It should be noted that LWC_c and N_c in a convection cloud chamber have a stronger correlation compared with those in atmospheric clouds (Shaw et al., 2023). Instead, the LWC_c and N_c often exhibit a positive covariance relationship. To understand the typical value of these two quantities in the chamber environment, we refer to typical measurement data from the Pi Chamber (magenta dots in Fig. 2). The data are from experiments conducted by Chandrakar et al. (2018). We can notice that for this specific experiment set up, drizzle embryos with diameter ranging from 40 μm to 60 μm can be detected using radar observational volume from 1 cm^3 to 10 cm^3 .

The aforementioned estimation is conducted under the assumption that signal (drizzle) power is equal to the background (cloud) power. In practice, to reduce the detection false alarms, the drizzle signal should be larger than the backgrounds. Here we define the signal to noise ratio (SNR) to investigate the drizzle detectability in the chamber environment:

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$$SNR = 10 \log 10 \left(\frac{D_d^6}{\int VP(D_c) D_c^6 dD} \right).$$
 (8)

Fig. 3 shows the estimated SNR for four drizzle particles under varying LWC_c and N_c conditions with a radar volume of 1 cm^3 . Generally, a smaller LWC_c and a larger N_c correspond to a large SNR, which is preferable for drizzle detection. If we arbitrarily choose SNR > 3 as the detection threshold, to detect a drizzle drop with diameter of 50 μm in a radar volume of 1 cm^3 (Fig. 3c), LWC_c in the cloud chamber should be lower than 0.3 g m^{-3} and N_c should be high than 90 cm^{-3} . The required LWC_c and N_c would be different for different drizzle particle targets: to detect drops with diameter of 60 μm , LWC_c should be lower than 0.5 g m^{-3} and N_c should be higher than 90

 cm^{-3} . It should be noted that although a drizzle drop is more likely to be detected by the radar at a lower LWC_c, drizzle initiation is generally more likely to occur at a higher LWC_c because the collision-coalescence rate is thought to be proportional to the square of LWC (Kostinski and Shaw, 2005). This suggests appropriate LWC_c and N_c combinations should be achieved such that drizzle can form by the C-C process in a convection cloud chamber and it can also be detected by radar in a small sampling volume. It is also noted that the results shown in Fig.3 are based on a radar volume of 1 cm^3 , and the estimated SNR would change if a different radar volume size is applied. For instance, increasing the radar volume will enhance the background power thus decreasing the SNR for the given cloud chamber environment.

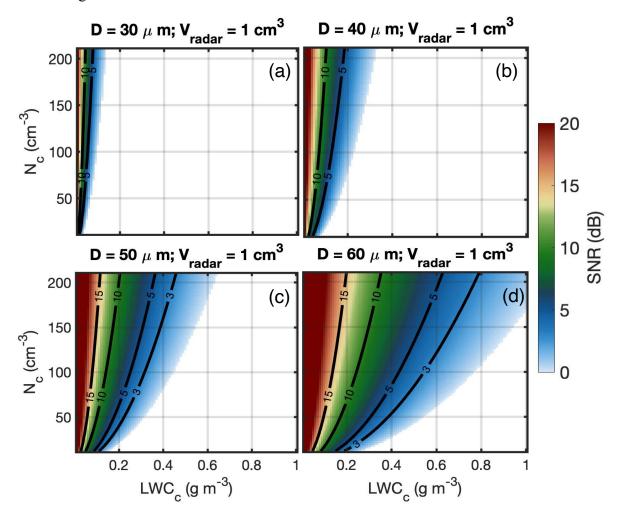


Figure 3: SNR of the drizzle signal under different LWC and N conditions in a 1 cm^3 radar sample volume for drizzle diameters of 30 μm , 40 μm , 50 μm , 60 μm . The black lines are SNR contours of 3 dB, 5 dB, 10 dB and 15 dB. SNR lower than 0 is indicated as the blank region.

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270 In the previous section, it was demonstrated that a radar with very small sampling volume ($\sim cm^3$) 271 can plausibly achieve the detection of single drizzle droplets against a cloud background signal.

On the other hand, the smaller the radar volume, the lower the probability of a drizzle particle 272 273 encountering the volume. To illustrate this trade-off scenario, we define the probability of drizzle

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 $p(D_d) = \begin{cases} 1, & VN(D_d)\Delta D \ge 1\\ VN(D_d), & VN(D_d)\Delta D < 1 \end{cases}$ (9)

Specifically, the product of V and $N(D_d)$ represents the expected number of drizzle drops in the radar volume. If the product is smaller than 1, it indicates the probability of the occurrence of

4. Probability of detection due to drizzle concentration

occurrence in the radar volume $(p(D_d))$ as:

drizzle particle in a given volume; while if the product is larger than 1, it means, statistically, at

least one drizzle drop with a diameter of D_d exists in the radar volume, and thus we set $p(D_d)=1$.

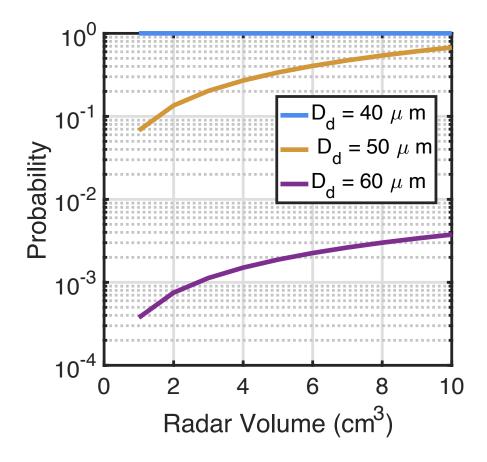


Figure 4: The probability of drizzle occurrence as a function of radar observational volume. The blue, yellow, and purple lines indicate the drizzle particle with diameters of 40 μm , 50 μm , and 60 μm .

The probability of occurrence for three selected drizzle particles as a function of radar volume is shown in Fig. 4. The N(D) in Eq. 9 is adapted from the size distribution described by Eq.6, with LWC_c and N_c set as $0.5 \ g \ m^{-3}$ and $50 \ cm^{-3}$, respectively. For these conditions, drizzle droplets with a diameter of $40 \ \mu m$ have sufficiently high concentration to be on average always present in volumes larger than $1 \ cm^3$. For drizzle droplets with a diameter of $50 \ \mu m$ or $60 \ \mu m$, their concentration is low enough that their probability of being found in a $10 \ cm^3$ volume is on average below 1. It is also noticed that the occurrence probability is strongly sensitive to the particle size: the probability of drizzle with $60 \ \mu m$ diameter occurring in the volume is almost two magnitudes smaller than that for a particle with $50 \ \mu m$ diameter. A smaller drizzle occurrence in the volume

indicates that a larger number of radar samples would be required to find one particle, leading to a longer observational time.

The probability of a drizzle drop to be in the radar sampling volume or passing through the radar volume within a finite time period should be an important consideration for a practical measurement system. Fig. 5 shows the probability of the occurrence of drizzle particle under different chamber environments same as Fig. 3. The blank region in Fig. 5 indicates the corresponding SNR shown in Fig. 3 is lower than 0 (i.e., cannot be detected by the radar even they exist in the sampling volume). Generally, it is noticed that the probability of occurrence differs in various chamber environment for different droplet size: large droplets have lower occurrence probability under small LWC and high *N* conditions.

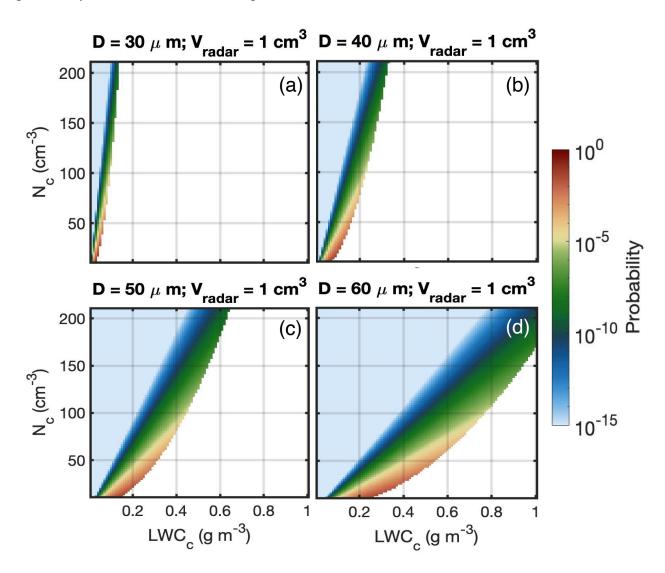


Figure 5: Drizzle occurrence probability under different LWC and N conditions for a 1 cm^3 radar volume with particle diameter of (a) 30 μm , (b) 40 μm , (c) 50 μm , and (d) 60 μm . The blank region indicates that the associated SNR is smaller than 0 (Fig. 3).

Comparison of Fig. 3 and Fig. 5 reveals that conditions that favor high radar SNR (i.e., larger drops or smaller radar sampling volume) are associated with a lower probability of occurrence of the drizzle droplet in the radar volume and subsequently increase the radar sampling time. For example, to detect a drizzle particle of 50 μm diameter under the condition of LWC_c and N_c of 0.3 g m^{-3} , 90 cm⁻³, the particle occurrence probability is on the magnitude of 10⁻⁸ (Fig. 5c) for SNR equals to 3 (Fig. 3c). A 1 dB enhancement of SNR threshold would decrease the occurrence probability to 10⁻¹ 11. This implies that on average, a volume of air equal to 10¹¹ times the size of the radar sampling volume needs to be sampled before a drizzle droplet will be detected. Assuming an air mean flow within the cloud chamber of 1 m s⁻¹, this implies that a radar sampling volume with a typical dimension of 1 cm will be updated (through advection) 100 times per second. If the radar is sampling along 1000 range gates (i.e. assuming a chamber with height of 10 m), this suggests that the radar can sample a volume equal to 10^5 its radar sampling volume each second. To reach the average required sampling volume (1011), it will take 106 seconds or 11.5 days. This is an unrealistically long observational time. For practical application, we want to work with sampling configurations that will not require to sample more than 10^9 times the radar sampling volume (~10s of minutes).

Another factor to consider in estimating the probability of drizzle occurrence with a certain diameter in a specific volume is the realism of using Eq. 6 for describing the N(D) in a cloud chamber. Eq. 6 describes the cloud droplet distribution controlled by the condensation process alone, thus the results may underestimate the actual drizzle occurrence as condensation is inefficient to produce large drizzle particle. In nature or in a large convection cloud chamber, the C-C mechanism is expected to be a more efficient process to increase the size and concentration of drizzle droplets.

Here we apply the ClusColl model to demonstrate that Fig. 4 and Fig. 5 may underestimate the drizzle occurrence probability with the collision-coalescence process being activate. ClusColl is a

simulation method for describing droplet motions and collisions in turbulent flows (Krueger and Kerstein, 2018). ClusColl simulates the movement of individual droplets in a vertical column due to turbulence and gravitational sedimentation. The unique capability of the ClusColl model is its capability to efficiently simulate the droplet collisions and coalescence process. Fig. 6 shows the simulated DSD with and without the collision-coalescence process for a 10-m-height cloud chamber and with cloud number concentration of 100 cm⁻³. The temperature difference between top and bottom walls is 40 °C. Noticeable differences can be identified at the right tail of the distribution, particularly for droplet diameter larger than 40 µm: more larger droplets are generated if the collision-coalescence process is active. The higher concentration of large drops results in a significantly shorter waiting time for detection compared to what was calculated for the condensation-only examples given in the earlier part. For instance, for the generated particle with diameter of 50 μm , the C-C process can generate number concentration more than 100 higher than the one without C-C process included. Reviewing the earlier estimation, to detect drizzle particle with diameter of 50 µm with SNR higher than 4, the required 10⁶ s becomes 10⁴ s which is approximately 3 hours, which is much more achievable for laboratory experiment. Thus, the estimation based on the condensation-only distribution (Eq. 6) is the most-conservative scenario. The actual radar measurement time would likely to be much shorter when the C-C process is activated.

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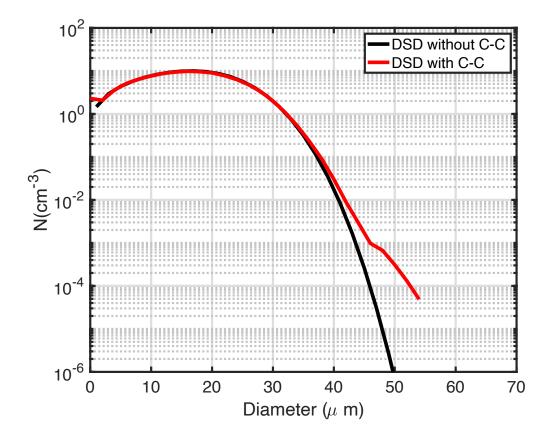


Figure 6: DSD simulated from the ClusColl model with (red line) and without (black line) droplet growth by collision-coalescence. In both cases, growth by condensation in a uniform supersaturation field, and removal by size-dependent droplet sedimentation are calculated. Therefore, the black line is described by the distribution given by Equation 6.

5. Evaluation from cloud chamber observations

In a cloud chamber and in the real atmosphere, the DSD in the radar sampling volume is expected to be time dependent due to turbulent fluctuations. To better quantify the particle backscattering power and its fluctuation in a small volume, observations made in the Pi Chamber using a holographic system (Holo-Pi) are used. Holo-Pi uses the principle of in-line digital holography to

measure the spatial distribution and sizes of cloud particles (Fugal and Shaw, 2009; Beals et al., 2015), and is specifically designed for the Pi-chamber environment (Desai et al., 2018). In contrast to the typical measurement strategy of single particle detections requiring time averaging, Holo-Pi captures instantaneous snapshots of all cloud droplets in the sample volume of 3.6 cm³ (0.6 cm \times 0.6 cm \times 10 cm) and is well suited to measure the temporal variations of cloud droplet populations within a sample volume similar to plausible radar sample volumes. The inability to resolve the smallest cloud droplets in the size distribution is not expected to be a significant limitation as the backscattering radar power is more sensitive to larger particle diameters. For the results presented here, cloud droplets are formed in the Pi Chamber by activation of size-selected sodium chloride aerosol particles (dry particle diameter $\approx 130 \text{ nm}$) injected into a supersaturated turbulent flow sustained by an unstable temperature difference of 20 K. An illustration of the 3D view of the cloud droplets measured by Holo-Pi at different time instants in the Pi-chamber is shown in Fig. 7. The sample volume used for our calculations is limited to a vertical extent of 5 cm as particle detectability falls off beyond this point; this results in a total sample volume of 1.8 cm³. The Holo-Pi system is set up to capture a hologram every 10 seconds during a 720-s period. For the optical configuration used here, the Holo-Pi has a lower size resolution of 12 µm throughout its sample volume.

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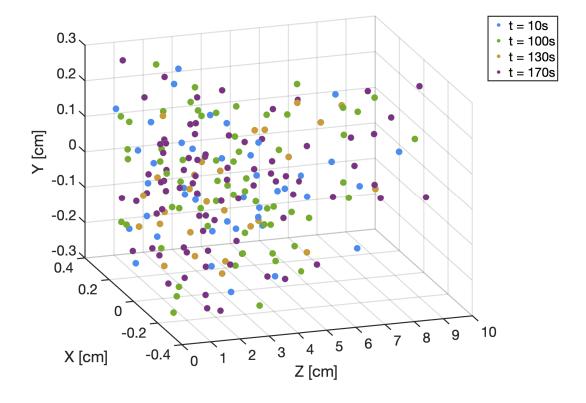


Figure 7: A 3D view of the particle locations observed by Holo-Pi in the Pi Chamber. Different colors represent observations taken at different timesteps.

The Holo-Pi observational volume is divided into five sub volumes with the cross section of 0.36 cm^2 and the depth increasing from 1 to 5 cm with an increment of 1 cm, thus, corresponding to a volume of 0.36, 0.72, 1.08, 1.44 and 1.8 cm^3 . Within each sub-volume at each time step, the total backscattering cross section for the detected droplets is estimated using a THz radar with wavelength of 0.44 mm. The calculated radar backscattering cross section as a function of volume size is shown in Fig. 8. Similar to the previous estimation, we see that the background power increases with volume size due to the increment of cloud droplets. Importantly, the uncertainty bars shown in Fig. 8 represent the standard deviation of the backscattering cross section during the observational time, which indicates the background power fluctuations. We notice that the cloud distribution in a small radar sampling volume is highly heterogeneous in time, and the magnitude of the background fluctuation varies by approximately a factor of 10. In order to detect drizzle drops, the backscattering power of the drizzle drop should be larger than the range of background

fluctuations. For example, a radar volume smaller than $0.36~cm^3$ should be utilized to detect a droplet with a diameter larger than $22~\mu m$, and a radar sampling volume smaller than $1~cm^3$ is needed to detect droplet with diameter larger than $30~\mu m$ for this particular Pi Chamber experiment set up.



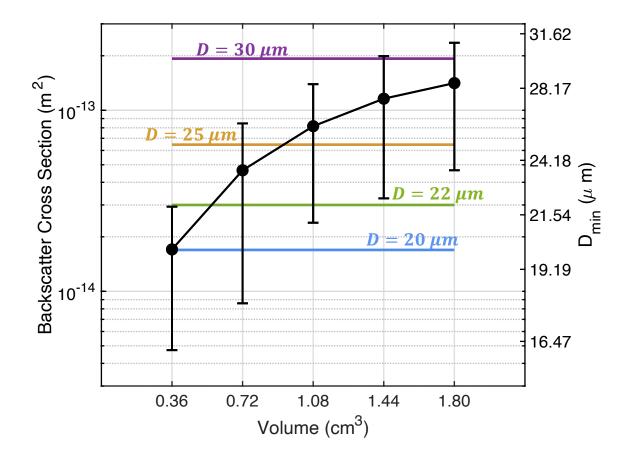
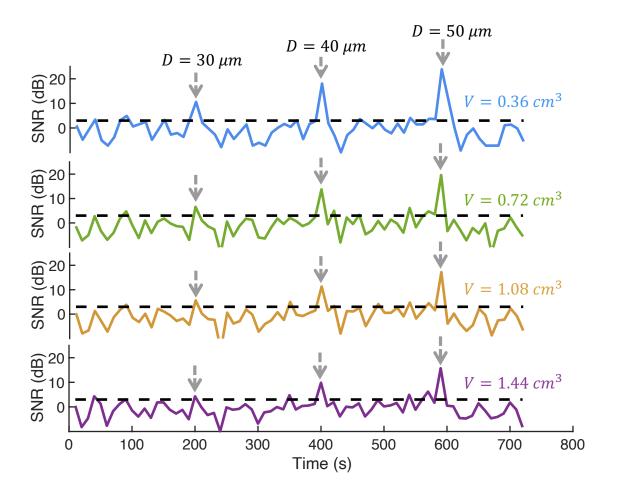


Figure 8: Dots and uncertainty bars indicate the mean and standard deviation of the total backscattering cross-section (with units of m^2) of droplets measured in different volumes by Holo-Pi during the observational period. The right axis and the horizontal lines represent the diameter of a single drizzle drop with backscattering power equivalent to the background.

To further demonstrate the single drizzle detection concept using a small radar volume, a pseudoradar observation experiment is conducted based on the Holo-Pi measurements. The Holo-Pi observational volume is divided into 4 sub-volumes indicated as different lines shown in Fig. 9. In each volume, we consider the mean radar backscattering power from all cloud particles sampled during the observational period as the background noise, and the power estimated at each time step as the signal, such that the SNR as a function of observation time is estimated. To simulate the drizzle occurring events, artificial drizzle droplets with diameter of $30 \,\mu m$, $40 \,\mu m$, $50 \,\mu m$ are added to the volume at $200 \, \text{s}$, $400 \, \text{s}$, and $600 \, \text{s}$, respectively, and the associated SNR is estimated. Fig. 9 shows a clear SNR enhancement when the drizzle droplets are added. The signal enhancement is more significant when using a small sampling volume and for larger drizzle drop diameter, which is consistent with the theoretical estimation in Sec. 3. For instance, a drizzle drop with a diameter of $50 \,\mu m$ can have a SNR of 23 dB with a volume of $0.36 \, cm^3$ while it has a SNR of 15 dB with a volume of $1.44 \, cm^3$. For a drizzle drop of $30 \,\mu m$, the SNR with a volume of $0.36 \, cm^3$ can reach to $10 \, \text{dB}$, which is an adequate SNR value for radar detection, while with a volume



of 1.44 cm³, the drizzle drop SNR is overwhelmed by background fluctuation and it is unable to be detected.

Figure 9: Simulated SNR of radar measurements during the Holo-Pi observational period using four sampling volumes: $0.36 \ cm^3$ (blue line), $0.72 \ cm^3$ (green line), $1.08 \ cm^3$ (yellow line) and $1.44 \ cm^3$ (purple line). The grey arrows indicate an artificial drizzle particle is added at the indicated time step. The black dashed line indicates a SNR of 3, which is used as a threshold to distinguish the signal (drizzle) from the background (clouds) in Fig. 3.

6. Summary

Recent simulation results suggest that drizzle initiation could occur in a large convection-cloud chamber. Such a facility would provide measurements in a controlled environment that can advance our understanding of warm rain formation in clouds. One of the critical measurements in a large convection-cloud chamber is the detection of low-concentration drizzle droplets in the presence of numerous cloud droplets. Early in the drizzle initiation, those drizzle drops are rare and in-homogeneously distributed in the chamber, presenting a significant detection challenge for conventional in-situ probes. Here, the potential of a radar with ultra-fine sampling volume for drizzle detection is investigated. It was demonstrated that if the radar sampling volume becomes orders of magnitude smaller (e.g., several cm^3), compared to those typically available in research radars ($\sim 10^3 - 10^6 \,\mathrm{m}^3$), isolated drizzle particles can be detected against the cloud background signal. This concept is based on the notion that the SNR of point targets (i.e., drizzle droplet) is independent of the radar sampling volume while the SNR of background (i.e., high concentration cloud droplets) scales with the sampling volume.

A theoretical DSD was adapted to represent the distribution of cloud droplets in a convection cloud chamber and to estimate properties of a detectable drizzle particle. It was shown that the minimum size of an isolated drizzle droplet that can be detected with such a radar depended on the radar sampling volume and the strength of the background signal (i.e., cloud droplets radar return), that in turn, depends on LWC and N_c. To minimize the false alarm drizzle detection, we require that the backscattering power from a drizzle particle should be larger than the backscattered power

contributed from the cloud particles (SNR>1). It is demonstrated that the application of a small radar volume can significantly enhance SNR under a given chamber environment. On the other hand, the smaller the radar sampling volume the lower the probability of an isolated drizzle droplet to be sampled. Thus, the determination of the radar volume for drizzle detection should account for the size of drizzle particle of interest, the environment conditions that favor drizzle initiation and the required observational time.

In addition to analytical estimates, real observations from the MTU Pi convection-cloud chamber are used to demonstrate the single drizzle particle detection framework. The Holo-Pi system(Desai et al., 2018) is applied to provide detailed 3D imaging of the cloud particles in the cloud chamber, from which the fluctuations of the backscattering power in a small volume can be well estimated. Generally, the observational results are consistent with the theoretical estimation showing that the background power is decreased and the ability to detect drizzle particles is enhanced as radar sample volume is decreased. It is also noticed that the magnitude of the background fluctuation is comparable to the mean power, which indicate that the distribution of cloud droplets is highly inhomogeneous in the small volume. Thus, the power from a drizzle particle should at least dominate the background power fluctuation in order to be detected. With the cloud chamber environment from the experiment, drizzle particles with diameter larger than 30 μm can be confidently detected using a radar sampling volume of 1 cm^3 or lower.

The key remaining question is the technological feasibility of building a radar that can operate within a box (large convection cloud chamber) and achieve the required ultra-fine range resolution. In fact, the effort of using "small" radar volumes for single particle detection has already been achieved in previous studies. For example, Schmidt et al. (2012) utilized a C-band radar with a 14- m^3 observational volume and successfully detected the trajectories of rain droplet with diameter down to 0.5 mm. In our case, the required radar sampling volume for drizzle detection is much smaller (with several cm^3). Such ultra-fine range resolution can be achieved using a THz radar operating at 340 or 680 GHz that can support wide bandwidth waveforms and thus enable subcentimeter range resolution (Cooper and Chattopadhyay, 2014). If the radar operates at a very high carrier frequency (THz) it can afford a very wide bandwidth for pulse modulation. In this case, the range resolution in not dictated by the pulse length but from the radar bandwidth (Cooper and

Chattopadhyay, 2014). The ultra-fine range resolution along with a reflector that minimizes the angular spread of the radar beam can result in radar sampling volumes of a few cm^3 . Such radar imaging capabilities have been extensively used for security screening at airports, for example. In our context, additional complexity is introduced by the fact that this radar needs to operate in a chamber with typical dimension of ~ 10 m. These technical design issues will be the focus of a follow-up paper study that will include real observations of drizzle droplets from a THz radar system.

To conclude, we outline three issues that will need to be properly addressed before a radar can be applied to the drizzle-detection problem in a cloud chamber:

1) Does the radar have enough sensitivity to detect a single drizzle particle? With the development of the THz technology, radar with centimeter resolution is achievable, however, the currently developed THz radars are mainly used to detect relatively hard targets that do not require ultra-high sensitivity. For the purpose of drizzle detection, however, the backscattering cross-section is on the order of $10^{-13} \, m^2$; such lower receiving power would require the radar to have a much higher transmitting power or a larger antenna size. Fortunately, an advantage for the drizzle detection in a cloud chamber is that the radar detection range is only several meters depending on the size of the chamber. According to Eq. 1, radar receiving power is inversely proportional to the fourth power of the target distance. Thus, the small detection range may greatly relieve the demand for high sensitivity in the radar design. In addition, recent advancements in THz transmitters allow us to utilize higher power output transmitter (~50 to 200 mW) at THz frequencies such as 340 GHz.

2) What are the appropriate radar sampling strategies for drizzle detection in a cloud chamber facility? Most of the cloud radars applied in the atmosphere are vertically pointing and can provide continuous observation at a given location along the radar beam. However as discussed in the paper, drizzle occurrence in the chamber is extremely rare and inhomogeneous in space and time. If the radar is vertically pointing, with the radar beam of several *cm* width, it may wait significant time for the radar to detect one drizzle drop.

Adding a scanning capability to the radar may provide a more efficient way to observe and detect drizzle in the cloud chamber.

3) How to eliminate or reduce the degradation effect of the chamber environment on a radar signal? In particular, the cloud chamber is a humid environment with liquid particles continually falling towards the bottom. Accumulation of water on the radar antenna can also severally attenuate the transmitting power and degrade the radar detectability. Furthermore, the chamber walls and the in-situ instruments mounted inside would produce strong backscattering signals and pollute the backscattering signal from hydrometeors. Thus, the design of the radar should also account for radar instrument design and sampling strategies that minimize these noise sources so that the best possible detection capability can be achieved.

In short, this paper demonstrates the conceptual feasibility of THz radars for rare drizzle detection in a laboratory context. Undoubtedly, the development of a high-resolution radar for drizzle detection in a cloud chamber needs close collaborations between cloud physics scientists and radar engineers moving forward.

Competing interests.

P. K. is an associate editor of AMT and the peer-review process was handled by an independent editor. The authors have no other competing interests to declare.

Code/Data availability

The codes and observations used to conduct all the analyses in this study are available upon request.

Author contributions Zeen Zhu conceptualized and implemented the method, performed the analysis, produced the figures, and wrote the initial draft of the manuscript. Pavlos Kollias and Fan Yang contributed to the conceptualization and provide guidance on the analysis. Raymond A. Shaw and Alex B. Kostinski contributed to the simplified model described in Section 3.1. Steve Krueger provided the ClusColl model. Nithin Allwayin provided the Holo-Pi measurements from the Pi Chamber. All authors read the manuscript draft and contributed comments and manuscript editing. Acknowledgements Authors from Brookhaven National Laboratory were supported by the Office of Biological and Environmental Research in the Department of Energy, Office of Science, through the United States Department of Energy Contract No. DE-SC0012704. Authors from Michigan Technological University and from Stony Brook University (subaward 2105003Z8) acknowledge support from National Science Foundation award AGS-2133229. A. Kostinski, S. Krueger, R. Shaw, and F. Yang accomplished some of this work during a visit at the Kavli Institute for Theoretical Physics as part of the Multiphase Flows in Geophysics and the Environment program. KITP is supported in part by National Science Foundation under grant number PHY-1748958. We also want to thank Dr. Ken Cooper for the constructive discussion and for providing feedback to the manuscript.

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