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**Detection of small drizzle droplets in a large cloud chamber using
ultra high-resolution radar**

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23 **Abstract**

24

25 A large convection cloud chamber has the potential to produce drizzle-sized droplets, thus offering
26 a new opportunity to investigate aerosol-cloud-drizzle interactions at a fundamental level under
27 controlled environmental conditions. One key measurement requirement is the development of
28 methods to detect the low concentration drizzle drops in such a large cloud chamber. In particular,
29 remote sensing methods may overcome some limitations of in situ methods.

30

31 Here, the potential of an ultra-high-resolution radar to detect the radar return signal of a small
32 drizzle droplet against the cloud droplet background signal is investigated. It is found that using a
33 small sampling volume is critical to drizzle detection in a cloud chamber to allow a drizzle drop in
34 the radar sampling volume to dominate over the background cloud droplets signal. For instance, a
35 radar volume of 1 cubic centimeter (cm^3) would enable the detection of drizzle embryos with
36 diameter larger than $40 \mu m$. However, the probability of drizzle sampling also decreases as the
37 sample volume reduces, leading to a longer observation time. Thus, the selection of radar volume
38 should consider both of the signal power and the drizzle occurrence probability. Finally,
39 observations from the Pi Convection-Cloud Chamber are used to demonstrate the single drizzle
40 particle detection concept using small radar volume. The results presented in this study also
41 suggest new applications of ultra-high-resolution cloud radar for atmospheric sensing.

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52 **1. Introduction**

53 Drizzle formation is one of the most important microphysical processes in warm clouds. Yet the
54 processes controlling drizzle formation remain poorly understood (Wood, 2012). The most
55 challenging aspect is the initial formation of drizzle embryos with diameter around $30 \mu m \sim 50$
56 μm . The formation of small drizzle particles in this range can neither be adequately explained by
57 the traditionally-defined condensation growth process nor by the traditionally-defined collision-
58 coalescence (C-C) process owing to their low efficiency (Yau and Rogers, 1996; Pruppacher and
59 Klett, 2012; Falkovich et al., 2006; Beard and Ochs III, 1993). Several mechanisms have been
60 hypothesized to explain the efficiency of these processes including i) fine-scale turbulence in cloud
61 (Pinsky and Khain, 1997; Shaw, 2003); ii) giant cloud condensation nuclei (GCCN) (Johnson,
62 1982; Feingold et al., 1999); and iii) longwave cooling (Roach, 1976; Harrington et al., 2000).
63 Nevertheless, it remains unclear to which extent these proposed mechanisms can adequately
64 explain the origin of drizzle embryos.

65 One main barrier that hinder our ability to investigate the drizzle initiation process is the lack of
66 observations with sufficient sensitivity and spatiotemporal resolution to detect the early growth of
67 drizzle particles. As such an instrumented large convection cloud chamber with well-controlled
68 initial and boundary conditions might help to improve our understanding of the drizzle initiation
69 mechanism (Shaw et al., 2020). Unlike other types of chambers, a convection-cloud chamber can
70 generate a steady state cloud system for hours in a turbulent environment by maintaining a warm
71 saturated bottom surface, a cold saturated top surface, and a constant aerosol injection rate (Chang
72 et al., 2016). The Michigan Tech Pi convection chamber with a dimension of 2 m x 2 m x 1 m
73 (width x depth x height) has been used to explore aerosol-cloud-turbulence interactions, however,
74 the Pi Chamber is too small to initiate drizzle embryos mainly due to the relatively short lifetime
75 of cloud droplets therein. Results from large eddy simulations indicate that drizzle can be initiated
76 in a large convection-cloud chamber with a height on the order of 10 m (Thomas et al., 2023).
77 However, the drizzle drops are sparse in a large chamber, so the detection of single drizzle embryos
78 in a large cloud chamber is challenging for in-situ probes that generally have a sampling volume
79 of only a few cubic centimeters. On the other hand, active remote sensors have the ability to rapidly
80 sample large volumes and thus offer an attractive option for the detection of small drizzle droplets
81 in a cloud chamber.

82

83 Here, we will demonstrate that the detection of an individual drizzle droplet in the presence of
84 numerous cloud droplets is possible with a radar that can achieve a very small sampling volume.

85 The detection of individual drizzle droplets is possible because the radar signal-to-noise ratio (SNR)
86 of a point target (drizzle droplet) is not affected by the radar observational volume, while the SNR
87 of a distributed target (cloud droplets) scales with the radar volume. In the following sections, the
88 detection limits of an individual drizzle particle are investigated using idealized particle size
89 distributions and real particle size distributions from the Michigan Tech Pi Chamber. In the end,
90 the potential of THz radars offering unprecedented sub-centimeter range resolution will be
91 discussed for developing the single drizzle detection radar (Cooper and Chattopadhyay, 2014).

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95 **2. Drizzle detection using radar**

96

97 The detection of early drizzle particles in clouds has been the topic of extensive research. First, the
98 radar needs to have sufficient sensitivity to detect cloud and drizzle droplets. This is typically
99 accomplished using millimeter-wavelength radar (Kollias et al., 2007). Early methodologies for
100 the detection of drizzle drop in clouds employ the use of reflectivity thresholds, ranging
101 from -15 to -20 dBZ, to identify drizzle existence (Frisch et al., 1995; Liu et al., 2008; Comstock
102 et al., 2004). Kollias et al. (2011) introduced the use of the radar Doppler spectra skewness as a
103 more sensitive method for detecting the presence of small drizzle droplets (Acquistapace et al.,
104 2017; Zhu et al., 2022). The radar Doppler spectra technique improved the detection of drizzle
105 droplets that can produce as low as -30 dBZ (Zhu et al., 2022).

106

107 However, the use of the radar Doppler spectra technique in a cloud chamber is challenging. First,
108 this will require that the radar point vertically to take advantage of the differential velocity between
109 cloud and drizzle droplets. If we assume a monodisperse droplet size distribution (DSD) and
110 Rayleigh scattering conditions, a drizzle detection limit of -30 dBZ is equivalent to a concentration
111 of 10^{-3} cm^{-3} of drizzle droplets with diameter equal to $100 \mu\text{m}$ or a concentration of 6.4×10^{-2}
112 cm^{-3} of drizzle droplets with $50 \mu\text{m}$ diameter. In the former case, the drizzle particle size is quite

113 large and not quite an early drizzle droplet detection. In the latter case, the concentration of drizzle
 114 droplets is much higher than the concentration observed in nature ($\sim 10^{-4} \text{ cm}^{-3}$) (Zhu et al., 2022).
 115 Furthermore, the conventional cloud radar has range resolution of tens of meters, which is not
 116 applicable in a chamber facility which may be on the order of several meters (approaching the
 117 collision mean free path).

118
 119 As a result, we consider alternative methods to increase the probability of early drizzle droplet
 120 detection against the cloud droplet signal. As the number concentration of drizzle particle is low,
 121 by applying a small radar sampling volume V_{Radar} , it is possible that only one drizzle droplet is
 122 present in V_{Radar} . In this case, the drizzle particle can be considered as a point target with
 123 backscattering cross section $\sigma(m^2)$ and the received radar echo power P_r (mW) is commonly
 124 expressed as (Battan, 1973):

$$125 \quad P_{r,drizzle} = P_t \frac{G^2 \lambda^2}{(4\pi)^3 r^4} \sigma(D_d) \quad (1)$$

126
 127 where P_t is the transmit peak power (mW), G is the antenna gain, r (m) is the range of the target
 128 relative to the radar receiver and λ (m) is the radar wavelength. It is noteworthy that P_r for a point
 129 target does not depend on the radar sampling volume V_{Radar} . For distributed targets such as a cloud
 130 droplet population described by a droplet size distribution (DSD) that represents the number
 131 concentration of cloud droplets as a function of diameter, the received radar echo power is given
 132 by:

$$133 \quad P_{r,cloud} = P_t \frac{G^2 \lambda^2}{(4\pi)^3 r^4} \cdot V_{Radar} \cdot \sum_{i=0}^n N_c(D_i) \sigma(D_i) \Delta D_i \quad (2)$$

134
 135
 136 Where n is the number of cloud droplets in the radar volume and $N_c(D)$ is the DSD with units of
 137 m^{-4} . In this case, the received radar echo power depends on the radar sampling volume, which is
 138 given by the following expression:

$$139 \quad V_{Radar} = \pi \left(\frac{r \theta_{3dB}}{2} \right)^2 \cdot \Delta R \quad (3)$$

141

142 where θ_{3dB} is antenna radiation pattern 3-dB beamwidth in radians and ΔR is the range resolution.

143 Assuming Rayleigh scattering, the backscatter cross section of the drizzle and cloud droplets is

144 proportional to the sixth power of the particle diameter and inversely proportional to the fourth

145 power of the wavelength ($\sigma(D) \sim D^6/\lambda^4$). Combing Eq. 1 and Eq. 2, the ratio of radar received

146 echo power from drizzle and cloud is given by the following expression:

147

$$148 \quad \frac{Signal}{Background} = \frac{P_{r,drizzle}}{P_{r,cloud}} = \frac{1}{V_{Radar}} \cdot \frac{D_d^6}{\sum_{i=0}^n N_c(D_i) D_i^6 \Delta D_i}. \quad (4)$$

149

150 Eq. 4 indicates that the probability of detecting a single drizzle droplet in the radar sampling

151 volume increases inversely to the radar sampling volume (point vs distributed target).

152

153 3. Detection requirement

154

155 Here, we will evaluate how small the radar sampling volume needs to be to detect drizzle drops

156 with different diameters against three background (cloud) conditions: i) monodisperse cloud DSD,

157 ii) cloud DSD from a theoretical model and iii) observed cloud DSD from the Michigan Tech Pi

158 Chamber. For simplicity, we will assume that a drizzle drop is detectable if its radar return power

159 is equal to that of the background echo contributed from cloud droplets.

160

161 3.1. Monodisperse cloud DSD

162

163 We first construct an idealized scenario by considering two categories of droplets, i.e. cloud droplet

164 with a diameter of D_c and a single drizzle drop with a diameter of D_d , the number concentration

165 of cloud droplets in the radar sampling volume is N_c (m^{-3}).

166

167 In this case, Eq. 4 is simplified as:

168

169

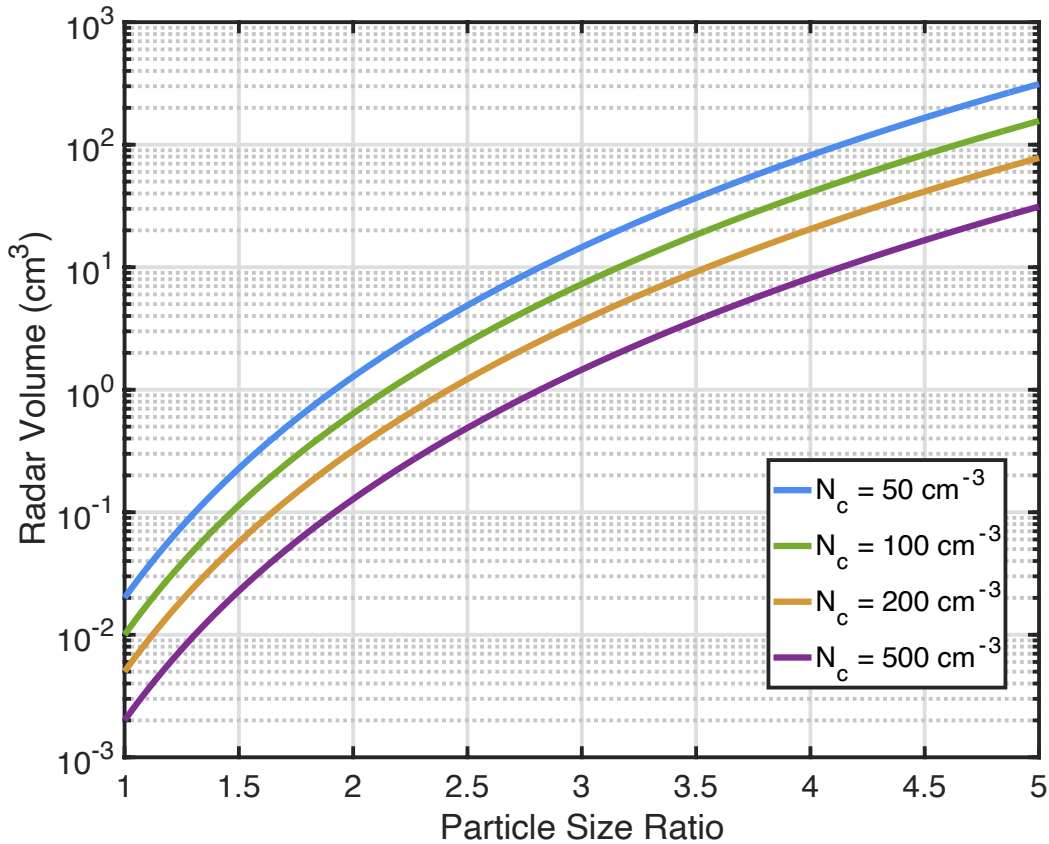
170

$$\frac{Signal}{Background} = \frac{1}{V_{Radar}} \cdot \frac{D_d^6}{N_c \cdot D_c^6} \quad (5)$$

171

172 When the signal power equals to the background, the radar sampling volume enabling single
 173 drizzle particle detection is estimated as a function of the size ratio $x = \frac{D_d}{D_c}$ shown in Fig. 1. The
 174 results are shown for various cloud droplet concentrations. It is noted that the required radar
 175 volume for detection depends on the drizzle drop size and the cloud number concentration. Larger
 176 radar volume would be required for drizzle detection as the particle size ratio increase; for a given
 177 particle size ratio, decreasing cloud number concentration can enhance the required radar volume.
 178 For example, if the cloud number concentration is 50 cm^{-3} and the mean cloud diameter (D_c) is
 179 $20 \mu\text{m}$, then the detection of a drizzle particle with diameter of $40 \mu\text{m}$ ($x = 2$) requires radar
 180 volume around 1 cm^3 . Such sampling volume are not achievable with traditional radar systems
 181 that employ sampling volumes of the order of 1000 m^3 or more (Kollias et al., 2016).

182



183

184 Figure 1: Radar observational volume for single-drizzle-drop detection as a function of particle
 185 size ratio $x = \frac{D_d}{D_c}$. Lines of different color represent clouds number concentration (N_c): 50 cm^{-3}
 186 (blue), 100 cm^{-3} (green), 200 cm^{-3} (yellow) and 500 cm^{-3} (purple).

187

188

189 3.2 Drizzle detection against an idealized cloud droplet background

190

191 In a realistic cloud chamber environment, we expect a population of cloud droplets with various
 192 sizes that can be represented by a DSD. Particularly, when condensation and fallout are the main
 193 sources and sinks for the evolution equation for the DSD, the DSD in the cloud chamber can be
 194 approximately described by theoretically derived distributions (Saito et al., 2019; Chandrakar et al.,
 195 2020; Krueger, 2020). Here we adapt the theoretical DSD formula derived by Krueger (2020) to
 196 investigate the ability of a radar to detect a drizzle embryo present in a small sample volume under
 197 different chamber environment conditions. To better represent the cloud DSD under different
 198 environments, the analytical DSD is rearranged to be expressed as a function of liquid water
 199 content (LWC_c ; $g \text{ m}^{-3}$) and number concentration (N_c ; m^{-3}) as:

200

$$201 \quad N(D_c) = \frac{2N_c D_c}{\pi^{1/2}} \left(\frac{4\Gamma\left(\frac{5}{4}\right) \pi^{1/2} \rho_l N_c}{3LWC_c} \right)^{2/3} \exp \left(- \left(\frac{4\Gamma\left(\frac{5}{4}\right) \pi^{1/2} \rho_l N_c}{3LWC_c} \right)^{\frac{4}{3}} \left(\frac{D_c}{2} \right)^4 \right) \quad (6)$$

202

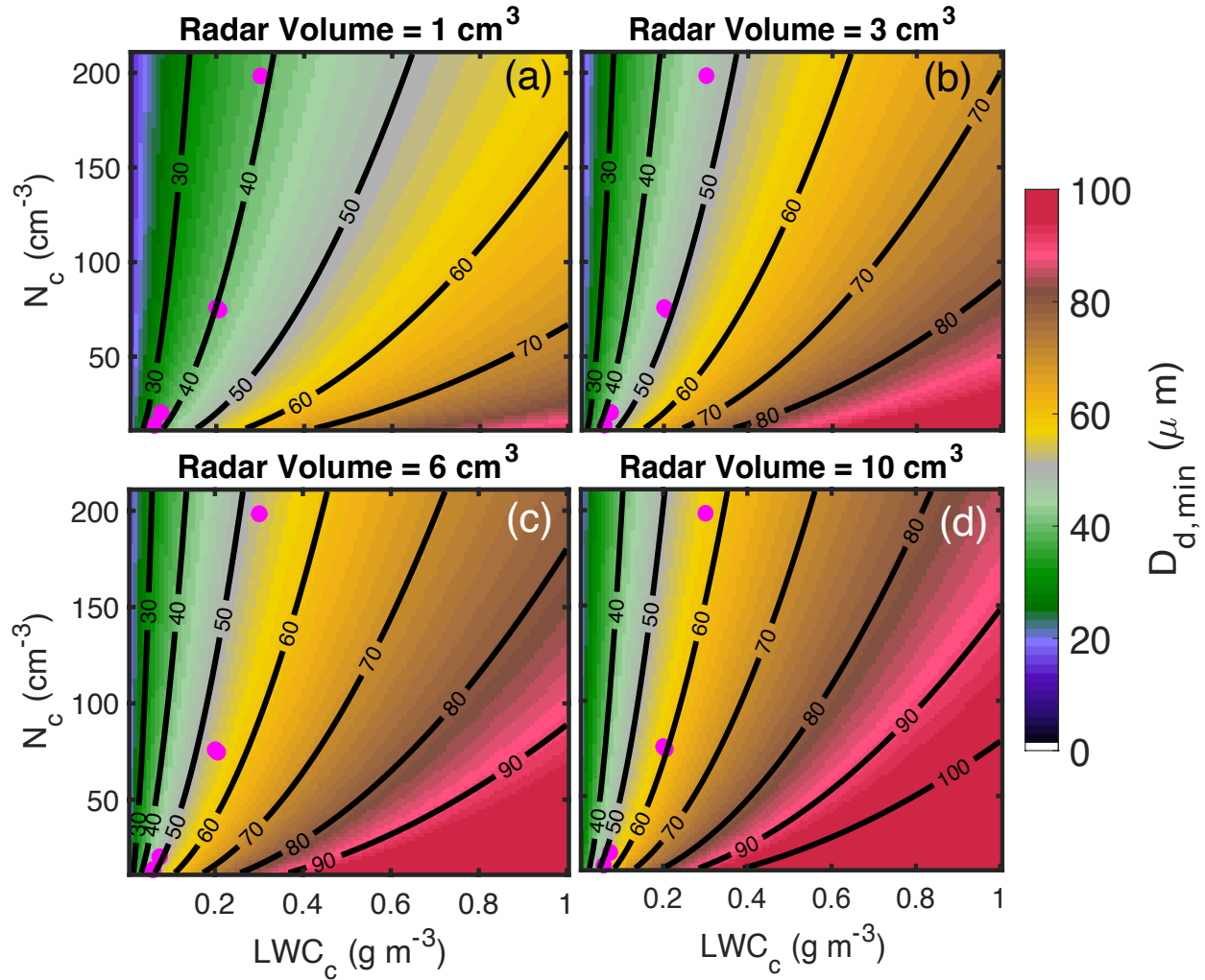
203 where ρ_l is liquid water density ($g \text{ m}^{-3}$), and D_c is cloud droplet diameter (m). $N(D_c)$ represents
 204 the number concentration of cloud droplet for the given diameter ($cm^{-3} \mu m^{-1}$).

205

206 Here we define the minimal drizzle drop ($D_{d,min}$) as the size of particle with radar return power
 207 equal to the total return power from cloud droplets in a given radar volume (V). Given the cloud
 208 DSD described by Eq. 6, $D_{d,min}$ can be estimated as:

209

$$210 \quad D_{d,min}^6 = \int V N(D_c) D_c^6 dD \quad (7)$$



211

212 Figure2: The minimal detectable drizzle particle ($D_{d,min}$) under different LWC_c and N_c conditions
 213 with radar sampling volume of (a) 1 cm^3 , (b) 3 cm^3 , (c) 6 cm^3 and (d) 10 cm^3 . The black lines are
 214 the $D_{d,min}$ contour of $30 \mu\text{m}$, $40 \mu\text{m}$, $50 \mu\text{m}$, $60 \mu\text{m}$, $70 \mu\text{m}$. The magenta dots indicate the LWC_c
 215 and N_c observed in the Pi-cloud chamber.

216

217 Fig. 2 illustrates $D_{d,min}$ under different LWC_c and N_c combinations for various radar volumes. For
 218 a given steady-state cloud in a convection chamber (i.e., fixed LWC_c and N_c), $D_{d,min}$ generally
 219 increases as the radar volume increases. This is because larger radar volumes contain more cloud
 220 droplets that produce stronger background power, thus only a larger drizzle particle with a higher
 221 backscattering signal would be detectable. On the other hand, for a given radar observational
 222 volume, $D_{d,min}$ is jointly determined by LWC_c and N_c which are inversely proportional. As such,

223 $D_{d,min}$ increases rapidly with increasing LWC_c but slightly decreases with increasing N_c . This
 224 contrasting relationship is caused by a larger sensitivity of radar reflectivity to particle size than to
 225 number concentration. Thus, increasing LWC_c can increase mean cloud particle size and greatly
 226 enhance the background power, leading to a larger detectable $D_{d,min}$. On the other side, when
 227 LWC_c is fixed, increasing cloud total number concentration tends to decrease particle size. The
 228 reduced cloud particle size would reduce the backscattering power and more than compensate for
 229 the power enhancement contributed from the increased number concentration.

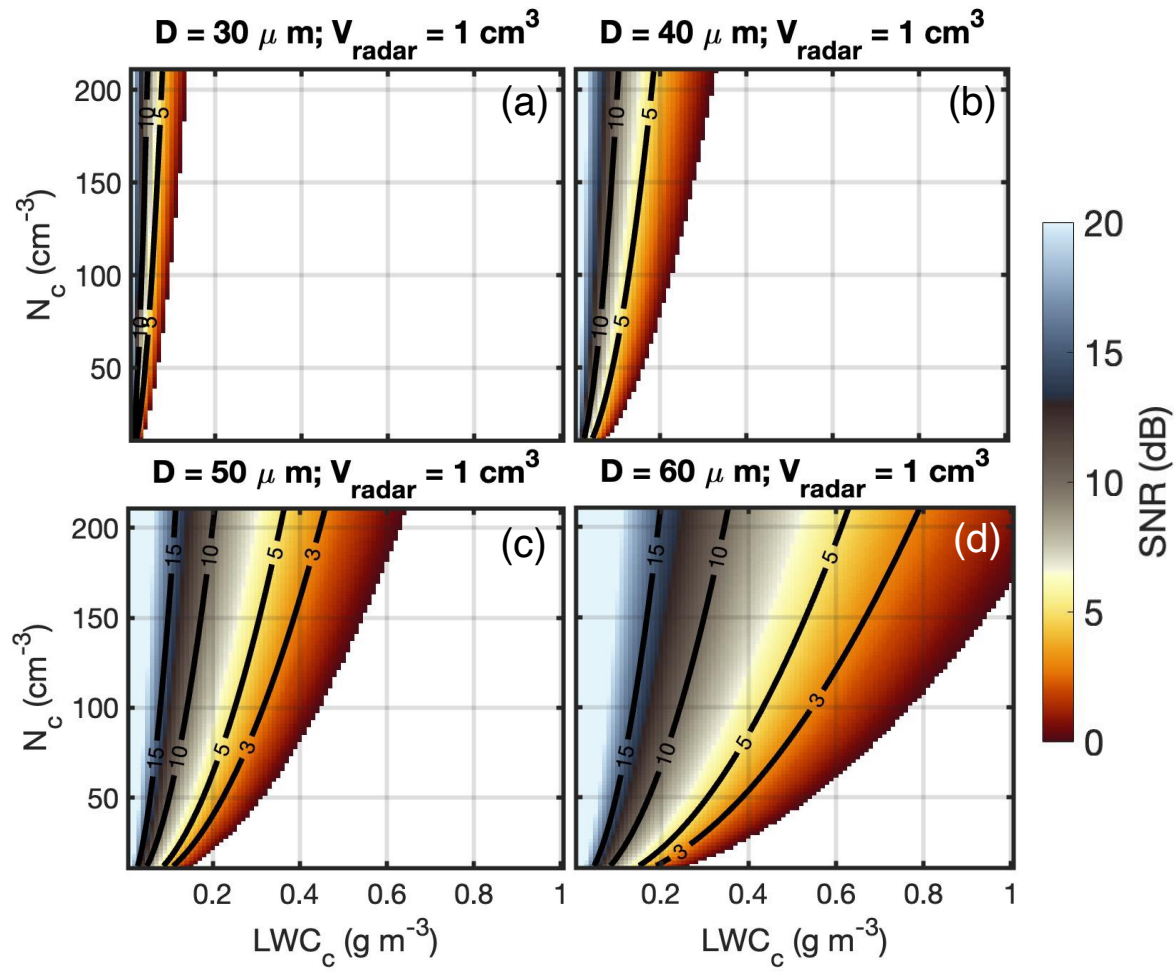
230
 231 It should be noted that LWC_c and N_c in a convection cloud chamber have a stronger correlation
 232 compared with those in atmospheric clouds (Shaw et al., 2023). Instead, the LWC_c and N_c often
 233 exhibit a positive covariance relationship. To understand the typical value of these two quantities
 234 in the chamber environment, we refer to typical measurement data from the Pi Chamber (magenta
 235 dots in Fig. 2). The data are from experiments conducted by Chandrakar et al. (2018). We can
 236 notice that for this specific experiment set up, drizzle embryos with diameter ranging from $40 \mu m$
 237 to $60 \mu m$ can be detected using radar observational volume from $1 cm^3$ to $10 cm^3$.

238
 239 The aforementioned estimation is conducted under the assumption that signal (drizzle) power is
 240 equal to the background (cloud) power. In practice, to reduce the detection false alarms, the drizzle
 241 signal should be larger than the backgrounds. Here we define the signal to noise ratio (SNR) to
 242 investigate the drizzle detectability in the chamber environment:

243
 244
$$SNR = 10 \log_{10} \left(\frac{D_d^6}{\int VP(D_c) D_c^6 dD} \right). \quad (8)$$

245
 246 Fig. 3 shows the estimated SNR for four drizzle particles under varying LWC_c and N_c conditions
 247 with a radar volume of $1 cm^3$. Generally, a smaller LWC_c and a larger N_c correspond to a large
 248 SNR, which is preferable for drizzle detection. If we arbitrarily choose $SNR > 3$ as the detection
 249 threshold, to detect a drizzle drop with diameter of $50 \mu m$ in a radar volume of $1 cm^3$ (Fig. 3c),
 250 LWC_c in the cloud chamber should be lower than $0.3 g m^{-3}$ and N_c should be high than $90 cm^{-3}$.
 251 The required LWC_c and N_c would be different for different drizzle particle targets: to detect drops
 252 with diameter of $60 \mu m$, LWC_c should be lower than $0.5 g m^{-3}$ and N_c should be higher than 90

253 cm^{-3} . It should be noted that although a drizzle drop is more likely to be detected by the radar at a
 254 lower LWC_c , drizzle initiation is generally more likely to occur at a higher LWC_c because the
 255 collision-coalescence rate is thought to be proportional to the square of LWC (Kostinski and Shaw,
 256 2005). This suggests appropriate LWC_c and N_c combinations should be achieved such that drizzle
 257 can form by the C-C process in a convection cloud chamber and it can also be detected by radar in
 258 a small sampling volume. It is also noted that the results shown in Fig.3 are based on a radar
 259 volume of 1 cm^3 , and the estimated SNR would change if a different radar volume size is applied.
 260 For instance, increasing the radar volume will enhance the background power thus decreasing the
 261 SNR for the given cloud chamber environment.



262 Figure 3: SNR of the drizzle signal under different LWC and N conditions in a 1 cm^3 radar sample
 263 volume for drizzle diameters of $30\ \mu\text{m}$, $40\ \mu\text{m}$, $50\ \mu\text{m}$, $60\ \mu\text{m}$. The black lines are SNR contours
 264 of 3 dB, 5 dB, 10 dB and 15 dB. SNR lower than 0 is indicated as the blank region.
 265

266

267 4. Probability of detection due to drizzle concentration

268

269 In the previous section, it was demonstrated that a radar with very small sampling volume ($\sim cm^3$)
270 can plausibly achieve the detection of single drizzle droplets against a cloud background signal.

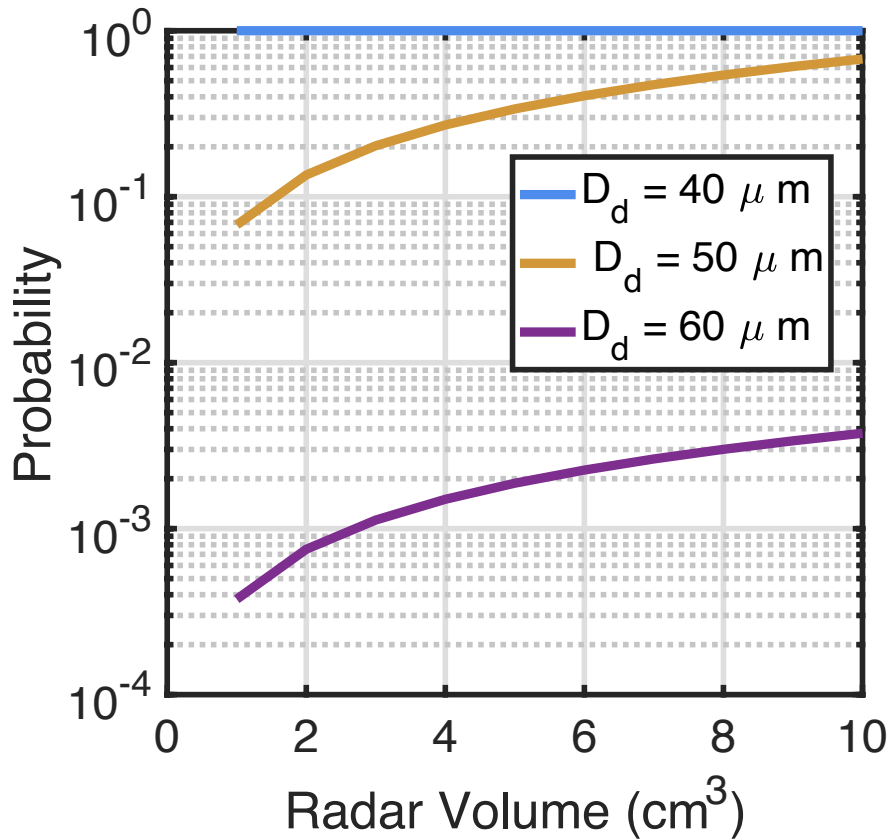
271 On the other hand, the smaller the radar volume, the lower the probability of a drizzle particle
272 encountering the volume. To illustrate this trade-off scenario, we define the probability of drizzle
273 occurrence in the radar volume ($p(D_d)$) as:

274

$$275 \quad p(D_d) = \begin{cases} 1, & VN(D_d)\Delta D \geq 1 \\ VN(D_d), & VN(D_d)\Delta D < 1 \end{cases} \quad (9)$$

276

277 Specifically, the product of V and $N(D_d)$ represents the expected number of drizzle drops in the
278 radar volume. If the product is smaller than 1, it indicates the probability of the occurrence of
279 drizzle particle in a given volume; while if the product is larger than 1, it means, statistically, at
280 least one drizzle drop with a diameter of D_d exists in the radar volume, and thus we set $p(D_d)=1$.



282

283 Figure 4: The probability of drizzle occurrence as a function of radar observational volume. The
 284 blue, yellow, and purple lines indicate the drizzle particle with diameters of $40 \mu m$, $50 \mu m$, and 60
 285 μm .

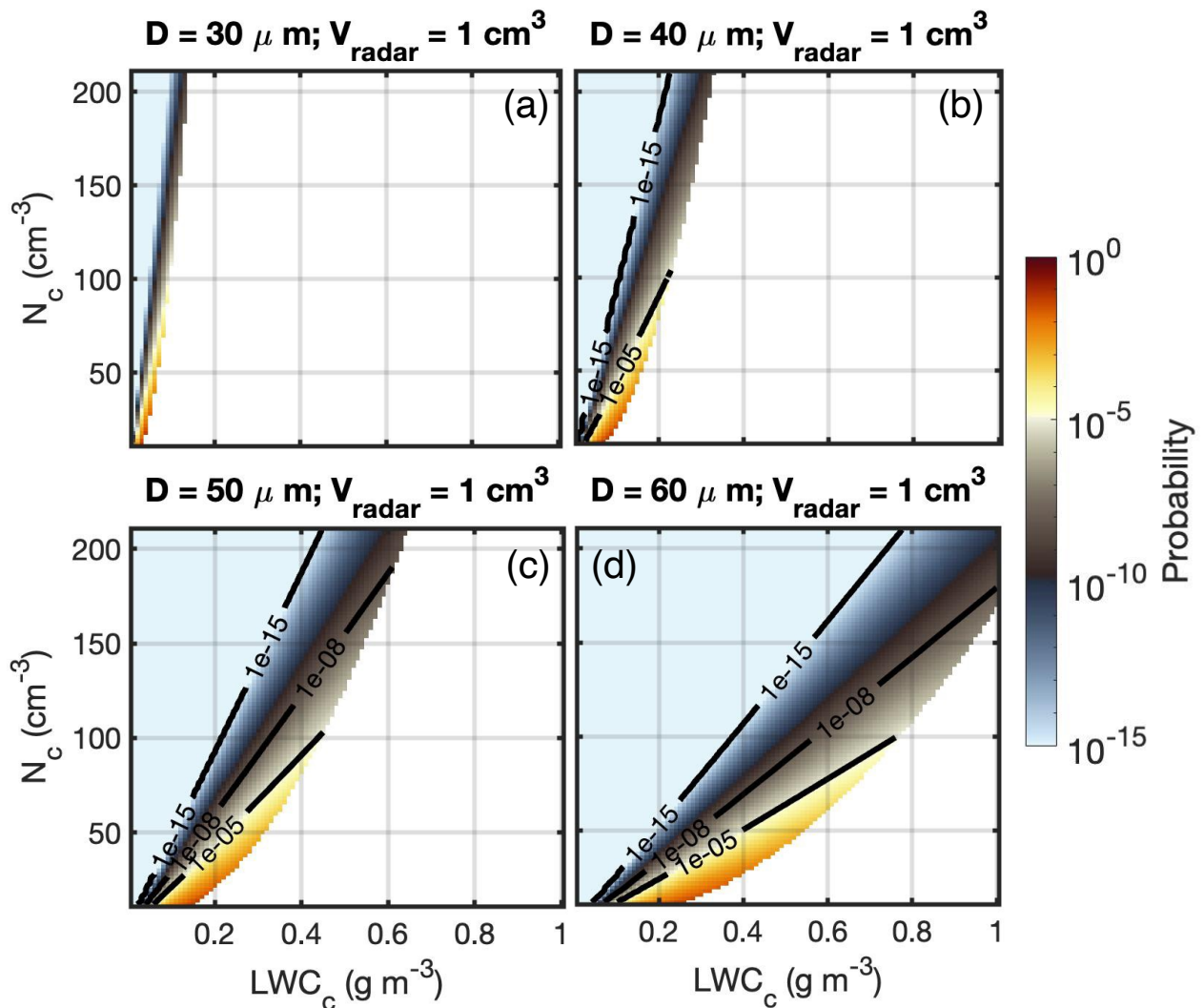
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287 The probability of occurrence for three selected drizzle particles as a function of radar volume is
 288 shown in Fig. 4. The $N(D)$ in Eq. 9 is adapted from the size distribution described by Eq.6, with
 289 LWC_c and N_c set as $0.5 g m^{-3}$ and $50 cm^{-3}$, respectively. For these conditions, drizzle droplets with
 290 a diameter of $40 \mu m$ have sufficiently high concentration to be on average always present in
 291 volumes larger than $1 cm^3$. For drizzle droplets with a diameter of $50 \mu m$ or $60 \mu m$, their
 292 concentration is low enough that their probability of being found in a $10 cm^3$ volume is on average
 293 below 1. It is also noticed that the occurrence probability is strongly sensitive to the particle size:
 294 the probability of drizzle with $60 \mu m$ diameter occurring in the volume is almost two magnitudes
 295 smaller than that for a particle with $50 \mu m$ diameter. A smaller drizzle occurrence in the volume

296 indicates that a larger number of radar samples would be required to find one particle, leading to
 297 a longer observational time.

298

299 The probability of a drizzle drop to be in the radar sampling volume or passing through the radar
 300 volume within a finite time period should be an important consideration for a practical
 301 measurement system. Fig. 5 shows the probability of the occurrence of drizzle particle under
 302 different chamber environments same as Fig. 3. The blank region in Fig. 5 indicates the
 303 corresponding SNR shown in Fig. 3 is lower than 0 (i.e., cannot be detected by the radar even they
 304 exist in the sampling volume). Generally, it is noticed that the probability of occurrence differs in
 305 various chamber environment for different droplet size: large droplets have lower occurrence
 306 probability under small LWC and high N conditions.



307 Figure 5: Drizzle occurrence probability under different LWC and N conditions for a 1 cm^3 radar
308 volume with particle diameter of (a) $30\ \mu\text{m}$, (b) $40\ \mu\text{m}$, (c) $50\ \mu\text{m}$, and (d) $60\ \mu\text{m}$. The black lines
309 are probability contours of 10^{-15} , 10^{-8} , 10^{-5} . The blank region indicates that the associated SNR is
310 smaller than 0 (Fig. 3).

311

312 Comparison of Fig. 3 and Fig. 5 reveals that conditions that favor high radar SNR (i.e., larger drops
313 or smaller radar sampling volume) are associated with a lower probability of occurrence of the
314 drizzle droplet in the radar volume and subsequently increase the radar sampling time. For example,
315 to detect a drizzle particle of $50\ \mu\text{m}$ diameter under the condition of LWC_c and N_c of $0.3\ \text{g m}^{-3}$, 90
316 cm^{-3} , the particle occurrence probability is on the magnitude of 10^{-8} (Fig. 5c) for SNR equals to 3
317 (Fig. 3c). A 1 dB enhancement of SNR threshold would decrease the occurrence probability to 10^{-
318 11 . This implies that on average, a volume of air equal to 10^{11} times the size of the radar sampling
319 volume needs to be sampled before a drizzle droplet will be detected. Assuming an air mean flow
320 within the cloud chamber of $1\ \text{m s}^{-1}$, this implies that a radar sampling volume with a typical
321 dimension of $1\ \text{cm}$ will be updated (through advection) 100 times per second. If the radar is
322 sampling along 1000 range gates (i.e. assuming a chamber with height of 10 m), this suggests that
323 the radar can sample a volume equal to 10^5 its radar sampling volume each second. To reach the
324 average required sampling volume (10^{11}), it will take 10^6 seconds or 11.5 days. This is an
325 unrealistically long observational time. For practical application, we want to work with sampling
326 configurations that will not require to sample more than 10^9 times the radar sampling volume ($\sim 10\text{s}$
327 of minutes).

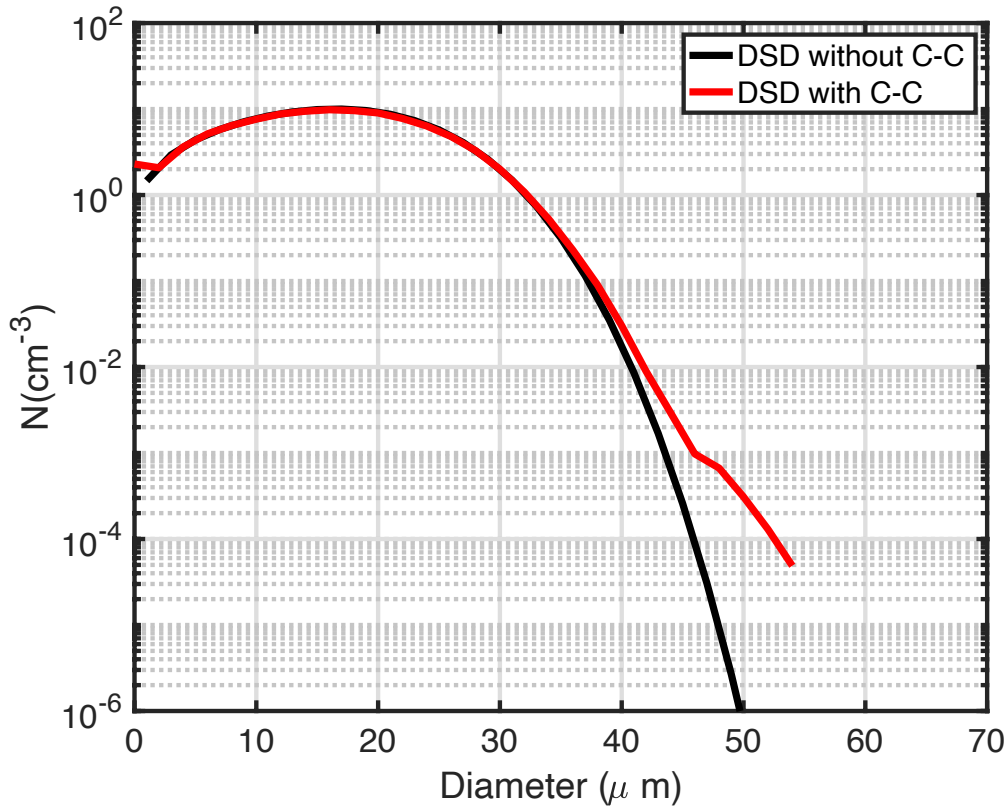
328

329 Another factor to consider in estimating the probability of drizzle occurrence with a certain
330 diameter in a specific volume is the realism of using Eq. 6 for describing the $N(D)$ in a cloud
331 chamber. Eq. 6 describes the cloud droplet distribution controlled by the condensation process
332 alone, thus the results may underestimate the actual drizzle occurrence as condensation is
333 inefficient to produce large drizzle particle. In nature or in a large convection cloud chamber, the
334 C-C mechanism is expected to be a more efficient process to increase the size and concentration
335 of drizzle droplets.

336

337 Here we apply the ClusColl model to demonstrate that Fig. 4 and Fig. 5 may underestimate the
338 drizzle occurrence probability with the collision-coalescence process being activate. ClusColl is a
339 simulation method for describing droplet motions and collisions in turbulent flows (Krueger and
340 Kerstein, 2018). ClusColl simulates the movement of individual droplets in a vertical column due
341 to turbulence and gravitational sedimentation. The unique capability of the ClusColl model is its
342 capability to efficiently simulate the droplet collisions and coalescence process. Fig. 6 shows the
343 simulated DSD with and without the collision-coalescence process for a 10-*m*-height cloud
344 chamber and with cloud number concentration of 100 cm^{-3} . The temperature difference between
345 top and bottom walls is $40 \text{ }^\circ\text{C}$. Noticeable differences can be identified at the right tail of the
346 distribution, particularly for droplet diameter larger than $40 \text{ }\mu\text{m}$: more larger droplets are generated
347 if the collision-coalescence process is active. The higher concentration of large drops results in a
348 significantly shorter waiting time for detection compared to what was calculated for the
349 condensation-only examples given in the earlier part. For instance, for the generated particle with
350 diameter of $50 \text{ }\mu\text{m}$, the C-C process can generate number concentration more than 100 higher than
351 the one without C-C process included. Reviewing the earlier estimation, to detect drizzle particle
352 with diameter of $50 \text{ }\mu\text{m}$ with SNR higher than 4, the required 10^6 s becomes 10^4 s which is
353 approximately 3 hours, which is much more achievable for laboratory experiment. Thus, the
354 estimation based on the condensation-only distribution (Eq. 6) is the most-conservative scenario.
355 The actual radar measurement time would likely to be much shorter when the C-C process is
356 activated.

357



358

359 Figure 6: DSD simulated from the ClusColl model with (red line) and without (black line) droplet
 360 growth by collision-coalescence. In both cases, growth by condensation in a uniform
 361 supersaturation field, and removal by size-dependent droplet sedimentation are calculated.
 362 Therefore, the black line is described by the distribution given by Equation 6.

363

364

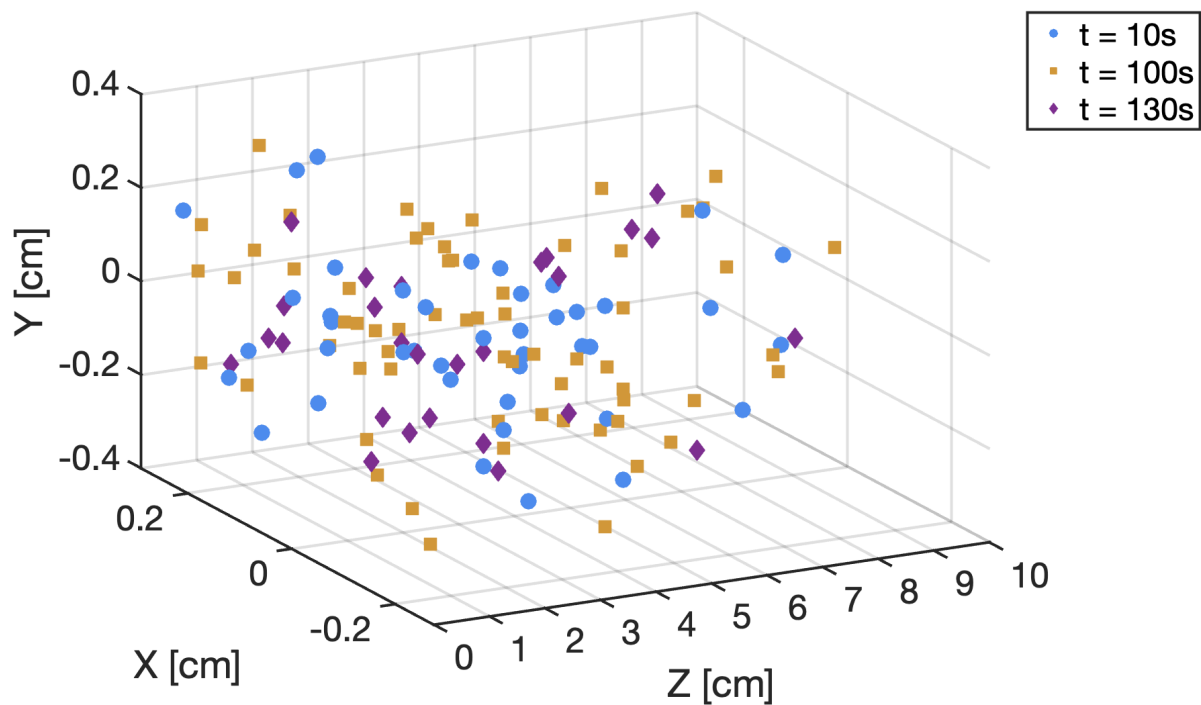
365

366 5. Evaluation from cloud chamber observations

367

368 In a cloud chamber and in the real atmosphere, the DSD in the radar sampling volume is expected
 369 to be time dependent due to turbulent fluctuations. To better quantify the particle backscattering
 370 power and its fluctuation in a small volume, observations made in the Pi Chamber using a
 371 holographic system (Holo-Pi) are used. Holo-Pi uses the principle of in-line digital holography to
 372 measure the spatial distribution and sizes of cloud particles (Fugal and Shaw, 2009; Beals et al.,
 373 2015), and is specifically designed for the Pi-chamber environment (Desai et al., 2018). In contrast

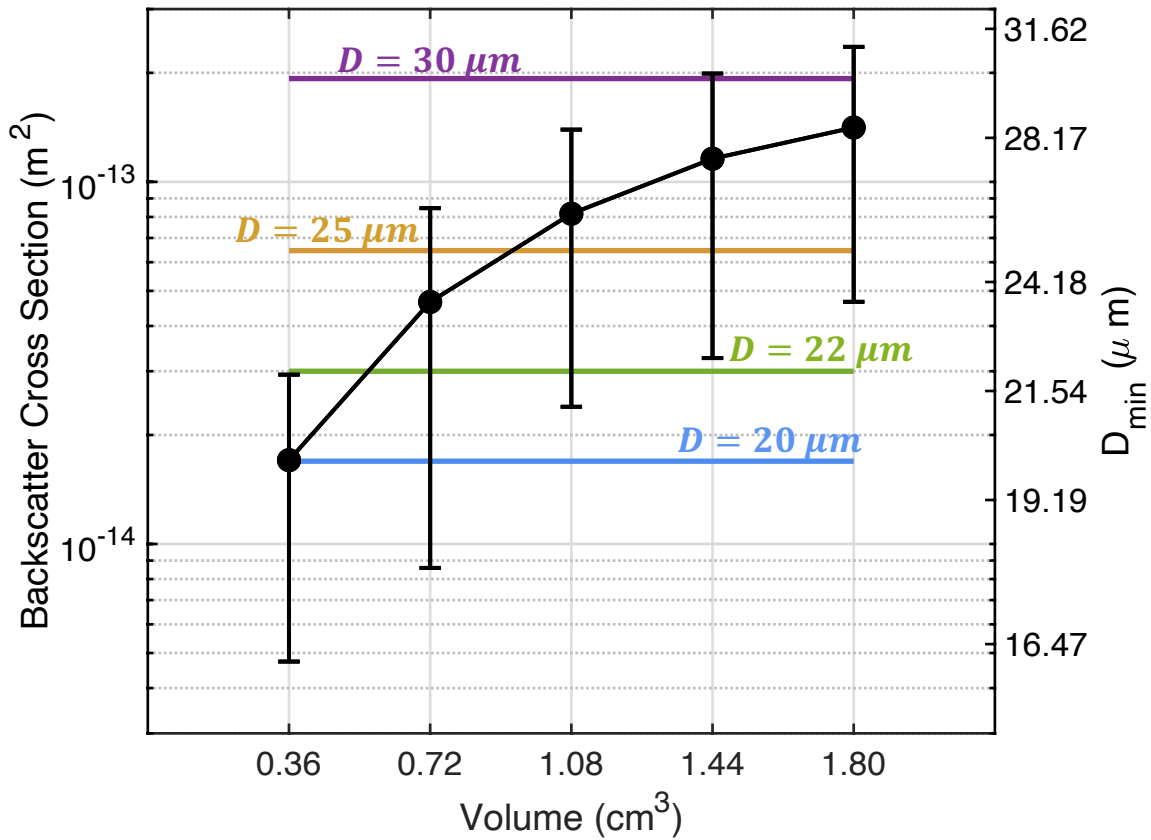
374 to the typical measurement strategy of single particle detections requiring time averaging, Holo-
 375 Pi captures instantaneous snapshots of all cloud droplets in the sample volume of 3.6 cm^3 (0.6 cm
 376 $\times 0.6 \text{ cm} \times 10 \text{ cm}$) and is well suited to measure the temporal variations of cloud droplet
 377 populations within a sample volume similar to plausible radar sample volumes. The inability to
 378 resolve the smallest cloud droplets in the size distribution is not expected to be a significant
 379 limitation as the backscattering radar power is more sensitive to larger particle diameters. For the
 380 results presented here, cloud droplets are formed in the Pi Chamber by activation of size-selected
 381 sodium chloride aerosol particles (dry particle diameter $\approx 130 \text{ nm}$) injected into a supersaturated
 382 turbulent flow sustained by an unstable temperature difference of 20 K. An illustration of the 3D
 383 view of the cloud droplets measured by Holo-Pi at different time instants in the Pi-chamber is
 384 shown in Fig. 7. The sample volume used for our calculations is limited to a vertical extent of 5
 385 cm as particle detectability falls off beyond this point; this results in a total sample volume of 1.8
 386 cm^3 . The Holo-Pi system is set up to capture a hologram every 10 seconds during a 720-s period.
 387 For the optical configuration used here, the Holo-Pi has a lower size resolution of $12 \mu\text{m}$
 388 throughout its sample volume.



389
 390 Figure 7: A 3D view of the particle locations observed by Holo-Pi in the Pi Chamber. Different
 391 colors/symbols represent observations taken at different timesteps.

392

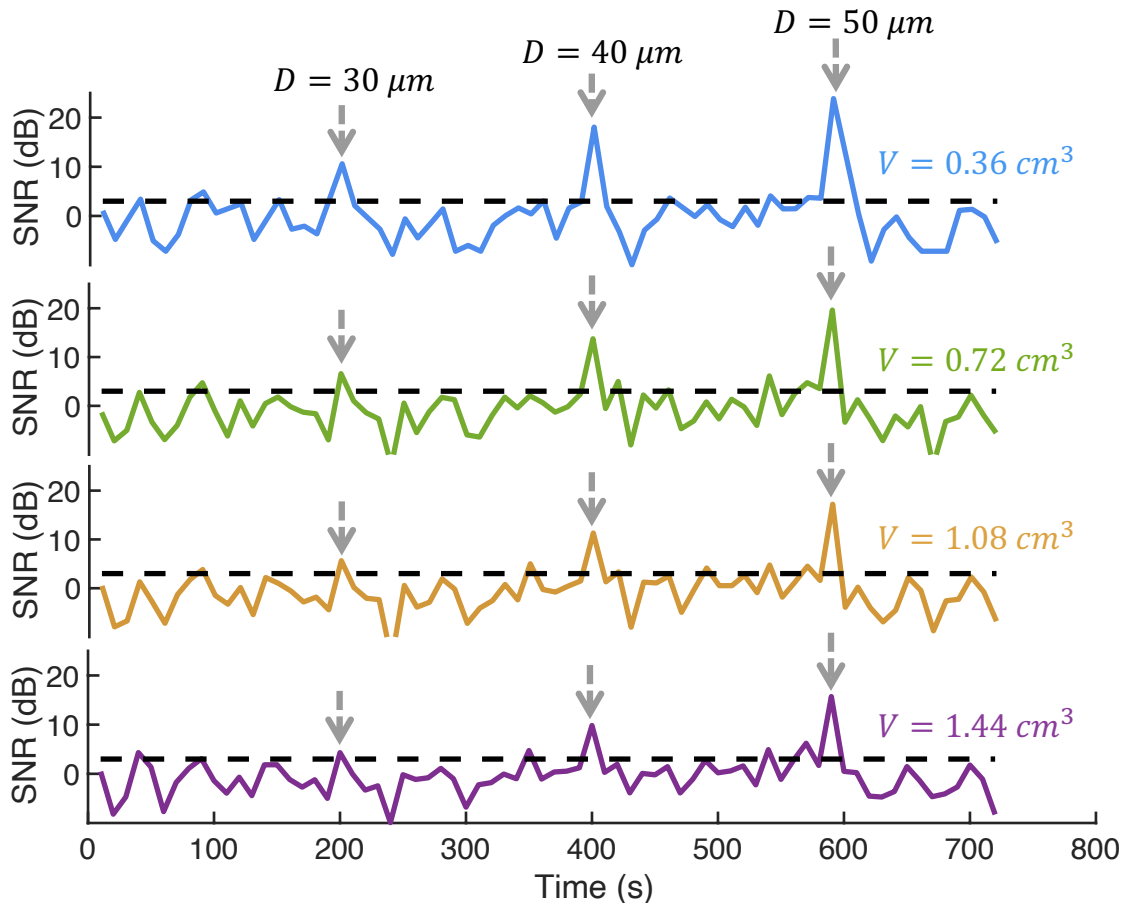
393 The Holo-Pi observational volume is divided into five sub volumes with the cross section of 0.36
394 cm^2 and the depth increasing from 1 to 5 cm with an increment of 1 cm , thus, corresponding to a
395 volume of 0.36, 0.72, 1.08, 1.44 and 1.8 cm^3 . Within each sub-volume at each time step, the total
396 backscattering cross section for the detected droplets is estimated using a THz radar with
397 wavelength of 0.44 mm . The calculated radar backscattering cross section as a function of volume
398 size is shown in Fig. 8. Similar to the previous estimation, we see that the background power
399 increases with volume size due to the increment of cloud droplets. Importantly, the uncertainty
400 bars shown in Fig. 8 represent the standard deviation of the backscattering cross section during the
401 observational time, which indicates the background power fluctuations. We notice that the cloud
402 distribution in a small radar sampling volume is highly heterogeneous in time, and the magnitude
403 of the background fluctuation varies by approximately a factor of 10. In order to detect drizzle
404 drops, the backscattering power of the drizzle drop should be larger than the range of background
405 fluctuations. For example, a radar volume smaller than 0.36 cm^3 should be utilized to detect a
406 droplet with a diameter larger than 22 μm , and a radar sampling volume smaller than 1 cm^3 is
407 needed to detect droplet with diameter larger than 30 μm for this particular Pi Chamber experiment
408 set up.



410

411 Figure 8: Dots and uncertainty bars indicate the mean and standard deviation of the total
 412 backscattering cross-section (with units of m^2) of droplets measured in different volumes by Holo-
 413 Pi during the observational period. The right axis and the horizontal lines represent the diameter
 414 of a single drizzle drop with backscattering power equivalent to the background.

415 To further demonstrate the single drizzle detection concept using a small radar volume, a pseudo-
416 radar observation experiment is conducted based on the Holo-Pi measurements. The Holo-Pi
417 observational volume is divided into 4 sub-volumes indicated as different lines shown in Fig. 9. In
418 each volume, we consider the mean radar backscattering power from all cloud particles sampled
419 during the observational period as the background noise, and the power estimated at each time step
420 as the signal, such that the SNR as a function of observation time is estimated. To simulate the
421 drizzle occurring events, artificial drizzle droplets with diameter of $30 \mu m$, $40 \mu m$, $50 \mu m$ are
422 added to the volume at 200 s, 400 s, and 600 s, respectively, and the associated SNR is estimated.
423 Fig. 9 shows a clear SNR enhancement when the drizzle droplets are added. The signal
424 enhancement is more significant when using a small sampling volume and for larger drizzle drop
425 diameter, which is consistent with the theoretical estimation in Sec. 3. For instance, a drizzle drop
426 with a diameter of $50 \mu m$ can have a SNR of 23 dB with a volume of 0.36 cm^3 while it has a SNR
427 of 15 dB with a volume of 1.44 cm^3 . For a drizzle drop of $30 \mu m$, the SNR with a volume of 0.36
428 cm^3 can reach to 10 dB, which is an adequate SNR value for radar detection, while with a volume



429 of 1.44 cm^3 , the drizzle drop SNR is overwhelmed by background fluctuation and it is unable to
430 be detected.

431

432 Figure 9: Simulated SNR of radar measurements during the Holo-Pi observational period using
433 four sampling volumes: 0.36 cm^3 (blue line), 0.72 cm^3 (green line), 1.08 cm^3 (yellow line) and
434 1.44 cm^3 (purple line). The grey arrows indicate an artificial drizzle particle is added at the
435 indicated time step. The black dashed line indicates a SNR of 3, which is used as a threshold to
436 distinguish the signal (drizzle) from the background (clouds) in Fig. 3.

437

438 6. Summary

439

440 Recent simulation results suggest that drizzle initiation could occur in a large convection-cloud
441 chamber. Such a facility would provide measurements in a controlled environment that can
442 advance our understanding of warm rain formation in clouds. One of the critical measurements in
443 a large convection-cloud chamber is the detection of low-concentration drizzle droplets in the
444 presence of numerous cloud droplets. Early in the drizzle initiation, those drizzle drops are rare
445 and in-homogeneously distributed in the chamber, presenting a significant detection challenge for
446 conventional in-situ probes. Here, the potential of a radar with ultra-fine sampling volume for
447 drizzle detection is investigated. It was demonstrated that if the radar sampling volume becomes
448 orders of magnitude smaller (e.g., several cm^3), compared to those typically available in research
449 radars ($\sim 10^3$ - 10^6 m^3), isolated drizzle particles can be detected against the cloud background signal.
450 This concept is based on the notion that the SNR of point targets (i.e., drizzle droplet) is
451 independent of the radar sampling volume while the SNR of background (i.e., high concentration
452 cloud droplets) scales with the sampling volume.

453

454 A theoretical DSD was adapted to represent the distribution of cloud droplets in a convection cloud
455 chamber and to estimate properties of a detectable drizzle particle. It was shown that the minimum
456 size of an isolated drizzle droplet that can be detected with such a radar depended on the radar
457 sampling volume and the strength of the background signal (i.e., cloud droplets radar return), that
458 in turn, depends on LWC and N_c . To minimize the false alarm drizzle detection, we require that
459 the backscattering power from a drizzle particle should be larger than the backscattered power

460 contributed from the cloud particles ($\text{SNR} > 1$). It is demonstrated that the application of a small
461 radar volume can significantly enhance SNR under a given chamber environment. On the other
462 hand, the smaller the radar sampling volume the lower the probability of an isolated drizzle droplet
463 to be sampled. Thus, the determination of the radar volume for drizzle detection should account
464 for the size of drizzle particle of interest, the environment conditions that favor drizzle initiation
465 and the required observational time.

466
467 In addition to analytical estimates, real observations from the MTU Pi convection-cloud chamber
468 are used to demonstrate the single drizzle particle detection framework. The Holo-Pi system(Desai
469 et al., 2018) is applied to provide detailed 3D imaging of the cloud particles in the cloud chamber,
470 from which the fluctuations of the backscattering power in a small volume can be well estimated.
471 Generally, the observational results are consistent with the theoretical estimation showing that the
472 background power is decreased and the ability to detect drizzle particles is enhanced as radar
473 sample volume is decreased. It is also noticed that the magnitude of the background fluctuation is
474 comparable to the mean power, which indicate that the distribution of cloud droplets is highly
475 inhomogeneous in the small volume. Thus, the power from a drizzle particle should at least
476 dominate the background power fluctuation in order to be detected. With the cloud chamber
477 environment from the experiment, drizzle particles with diameter larger than $30 \mu\text{m}$ can be
478 confidently detected using a radar sampling volume of 1 cm^3 or lower.

479
480 The key remaining question is the technological feasibility of building a radar that can operate
481 within a box (large convection cloud chamber) and achieve the required ultra-fine range resolution.
482 In fact, the effort of using “small” radar volumes for single particle detection has already been
483 achieved in previous studies. For example, Schmidt et al. (2012) utilized a C-band radar with a 14-
484 m^3 observational volume and successfully detected the trajectories of rain droplet with diameter
485 down to 0.5 mm . In our case, the required radar sampling volume for drizzle detection is much
486 smaller (with several cm^3). Such ultra-fine range resolution can be achieved using a THz radar
487 operating at 340 or 680 GHz that can support wide bandwidth waveforms and thus enable sub-
488 centimeter range resolution (Cooper and Chattopadhyay, 2014). If the radar operates at a very high
489 carrier frequency (THz) it can afford a very wide bandwidth for pulse modulation. In this case, the
490 range resolution is not dictated by the pulse length but from the radar bandwidth (Cooper and

491 Chattopadhyay, 2014). The ultra-fine range resolution along with a reflector that minimizes the
492 angular spread of the radar beam can result in radar sampling volumes of a few cm^3 . Such radar
493 imaging capabilities have been extensively used for security screening at airports, for example. In
494 our context, additional complexity is introduced by the fact that this radar needs to operate in a
495 chamber with typical dimension of ~ 10 m. These technical design issues will be the focus of a
496 follow-up paper study that will include real observations of drizzle droplets from a THz radar
497 system.

498

499 To conclude, we outline three issues that will need to be properly addressed before a radar can be
500 applied to the drizzle-detection problem in a cloud chamber:

501

502 1) Does the radar have enough sensitivity to detect a single drizzle particle? With the
503 development of the THz technology, radar with centimeter resolution is achievable,
504 however, the currently developed THz radars are mainly used to detect relatively hard
505 targets that do not require ultra-high sensitivity. For the purpose of drizzle detection,
506 however, the backscattering cross-section is on the order of $10^{-13} m^2$; such lower receiving
507 power would require the radar to have a much higher transmitting power or a larger antenna
508 size. Fortunately, an advantage for the drizzle detection in a cloud chamber is that the radar
509 detection range is only several meters depending on the size of the chamber. According to
510 Eq. 1, radar receiving power is inversely proportional to the fourth power of the target
511 distance. Thus, the small detection range may greatly relieve the demand for high
512 sensitivity in the radar design. In addition, recent advancements in THz transmitters allow
513 us to utilize higher power output transmitter (~ 50 to $200 mW$) at THz frequencies such as
514 340 GHz.

515

516 2) What are the appropriate radar sampling strategies for drizzle detection in a cloud chamber
517 facility? Most of the cloud radars applied in the atmosphere are vertically pointing and can
518 provide continuous observation at a given location along the radar beam. However as
519 discussed in the paper, drizzle occurrence in the chamber is extremely rare and
520 inhomogeneous in space and time. If the radar is vertically pointing, with the radar beam
521 of several cm width, it may wait significant time for the radar to detect one drizzle drop.

522 Adding a scanning capability to the radar may provide a more efficient way to observe and
523 detect drizzle in the cloud chamber.

524

525 3) How to eliminate or reduce the degradation effect of the chamber environment on a radar
526 signal? In particular, the cloud chamber is a humid environment with liquid particles
527 continually falling towards the bottom. Accumulation of water on the radar antenna can
528 also severally attenuate the transmitting power and degrade the radar detectability.
529 Furthermore, the chamber walls and the in-situ instruments mounted inside would produce
530 strong backscattering signals and pollute the backscattering signal from hydrometeors.
531 Thus, the design of the radar should also account for radar instrument design and sampling
532 strategies that minimize these noise sources so that the best possible detection capability
533 can be achieved.

534

535 In short, this paper demonstrates the conceptual feasibility of THz radars for rare drizzle detection
536 in a laboratory context. Undoubtedly, the development of a high-resolution radar for drizzle
537 detection in a cloud chamber needs close collaborations between cloud physics scientists and radar
538 engineers moving forward.

539

540 **Competing interests.**

541 **P. K.** is an associate editor of AMT and the peer-review process was handled by an independent
542 editor. The authors have no other competing interests to declare.

543

544 **Code/Data availability**

545 The codes and observations used to conduct all the analyses in this study are available upon
546 request.

547

548

549 **Author contributions**

550 Zeen Zhu conceptualized and implemented the method, performed the analysis, produced the
551 figures, and wrote the initial draft of the manuscript. Pavlos Kollias and Fan Yang contributed to
552 the conceptualization and provide guidance on the analysis. Raymond A. Shaw and Alex B.
553 Kostinski contributed to the simplified model described in Section 3.1. Steve Krueger provided
554 the ClusColl model. Nithin Allwayin provided the Holo-Pi measurements from the Pi Chamber.
555 All authors read the manuscript draft and contributed comments and manuscript editing.

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581 **Reference**

582

583 Acquistapace, C., Kneifel, S., Löhnert, U., Kollias, P., Maahn, M., and Bauer-Pfundstein, M.:
584 Optimizing observations of drizzle onset with millimeter-wavelength radars, *Atmospheric*
585 *Measurement Techniques*, 10, 1783-1802, 2017.

586 Battan, L. J.: Radar observation of the atmosphere, 1973.

587 Beals, M. A., Fugal, J. P., Shaw, R. A., Lu, J., Spuler, S. M., and Stith, J. L.: Holographic
588 measurements of inhomogeneous cloud mixing at the centimeter scale, *Science*, 350, 87-90,
589 2015.

590 Beard, K. V., and Ochs III, H. T.: Warm-rain initiation: An overview of microphysical mechanisms,
591 *Journal of Applied Meteorology and Climatology*, 32, 608-625, 1993.

592 Chandrakar, K., Cantrell, W., Kostinski, A., and Shaw, R.: Data supporting the paper "Dispersion
593 aerosol indirect effect in turbulent clouds: laboratory measurements of effective radius", 2018.

594 Chandrakar, K. K., Saito, I., Yang, F., Cantrell, W., Gotoh, T., and Shaw, R. A.: Droplet size
595 distributions in turbulent clouds: Experimental evaluation of theoretical distributions, *Quarterly*
596 *Journal of the Royal Meteorological Society*, 146, 483-504, 2020.

597 Chang, K., Bench, J., Brege, M., Cantrell, W., Chandrakar, K., Ciochetto, D., Mazzoleni, C.,
598 Mazzoleni, L., Niedermeier, D., and Shaw, R.: A laboratory facility to study gas-aerosol-cloud
599 interactions in a turbulent environment: The Π chamber, *Bulletin of the American Meteorological*
600 *Society*, 97, 2343-2358, 2016.

601 Comstock, K. K., Wood, R., Yuter, S. E., and Bretherton, C. S.: Reflectivity and rain rate in and
602 below drizzling stratocumulus, *Quarterly Journal of the Royal Meteorological Society: A journal*
603 *of the atmospheric sciences, applied meteorology and physical oceanography*, 130, 2891-2918,
604 2004.

605 Cooper, K. B., and Chattopadhyay, G.: Submillimeter-wave radar: Solid-state system design and
606 applications, *IEEE microwave magazine*, 15, 51-67, 2014.

607 Desai, N., Chandrakar, K. K., Chang, K., Cantrell, W., and Shaw, R.: Influence of microphysical
608 variability on stochastic condensation in a turbulent laboratory cloud, *Journal of the Atmospheric*
609 *Sciences*, 75, 189-201, 2018.

610 Falkovich, G., Stepanov, M. G., and Vucelja, M.: Rain initiation time in turbulent warm clouds,
611 *Journal of applied meteorology and climatology*, 45, 591-599, 2006.

612 Feingold, G., Cotton, W. R., Kreidenweis, S. M., and Davis, J. T.: The impact of giant cloud
613 condensation nuclei on drizzle formation in stratocumulus: Implications for cloud radiative
614 properties, *Journal of the atmospheric sciences*, 56, 4100-4117, 1999.

615 Frisch, A., Fairall, C., and Snider, J.: Measurement of stratus cloud and drizzle parameters in ASTEX
616 with a Ka-band Doppler radar and a microwave radiometer, *Journal of the Atmospheric Sciences*,
617 52, 2788-2799, 1995.

618 Fugal, J., and Shaw, R.: Cloud particle size distributions measured with an airborne digital in-line
619 holographic instrument, *Atmospheric Measurement Techniques*, 2, 259-271, 2009.

620 Johnson, D. B.: The role of giant and ultragiant aerosol particles in warm rain initiation, *Journal*
621 *of Atmospheric Sciences*, 39, 448-460, 1982.

622 Kollias, Clothiaux, E., Miller, M., Albrecht, B., Stephens, G., and Ackerman, T.: Millimeter-
623 wavelength radars: New frontier in atmospheric cloud and precipitation research, *Bulletin of the*
624 *American Meteorological Society*, 88, 1608-1624, 2007.

625 Kollias, Remillard, J., Luke, E., and Szyrmer, W.: Cloud radar Doppler spectra in drizzling stratiform
626 clouds: 1. Forward modeling and remote sensing applications, *Journal of Geophysical Research-*
627 *Atmospheres*, 116, 10.1029/2010jd015237, 2011.

628 Kollias, Clothiaux, E. E., Ackerman, T. P., Albrecht, B. A., Widener, K. B., Moran, K. P., Luke, E. P.,
629 Johnson, K. L., Bharadwaj, N., and Mead, J. B.: Development and applications of ARM millimeter-
630 wavelength cloud radars, *Meteorological Monographs*, 57, 17.11-17.19, 2016.

631 Kostinski, A. B., and Shaw, R. A.: Fluctuations and luck in droplet growth by coalescence, *Bulletin*
632 *of the American Meteorological Society*, 86, 235-244, 2005.

633 Krueger, S. K., and Kerstein, A. R.: An economical model for simulating turbulence enhancement
634 of droplet collisions and coalescence, *Journal of Advances in Modeling Earth Systems*, 10, 1858-
635 1881, 2018.

636 Krueger, S. K.: Equilibrium droplet size distributions in a turbulent cloud chamber with uniform
637 supersaturation, 2020.

638 Liu, Y., Geerts, B., Miller, M., Daum, P., and McGraw, R.: Threshold radar reflectivity for drizzling
639 clouds, *Geophysical research letters*, 35, 2008.

640 Pinsky, and Khain, A.: Turbulence effects on droplet growth and size distribution in clouds—A
641 review, *Journal of aerosol science*, 28, 1177-1214, 1997.

642 Pruppacher, H. R., and Klett, J. D.: *Microphysics of Clouds and Precipitation: Reprinted 1980*,
643 Springer Science & Business Media, 2012.

644 Saito, I., Gotoh, T., and Watanabe, T.: Broadening of cloud droplet size distributions by
645 condensation in turbulence, *Journal of the Meteorological Society of Japan. Ser. II*, 2019.

646 Schmidt, J. M., Flatau, P. J., Harasti, P. R., Yates, R. D., Littleton, R., Pritchard, M. S., Fischer, J. M.,
647 Fischer, E. J., Kohri, W. J., and Vetter, J. R.: Radar observations of individual rain drops in the free
648 atmosphere, *Proceedings of the National Academy of Sciences*, 109, 9293-9298, 2012.

649 Shaw, R. A.: Particle-turbulence interactions in atmospheric clouds, *Annual Review of Fluid*
650 *Mechanics*, 35, 183-227, 2003.

651 Shaw, R. A., Cantrell, W., Chen, S., Chuang, P., Donahue, N., Feingold, G., Kollias, P., Korolev, A.,
652 Kreidenweis, S., and Krueger, S.: Cloud–aerosol–turbulence interactions: Science priorities and
653 concepts for a large-scale laboratory facility, *Bulletin of the American Meteorological Society*,
654 101, E1026-E1035, 2020.

655 Shaw, R. A., Thomas, S., Prabhakaran, P., Cantrell, W., Ovchinnikov, M., and Yang, F.: Fast and
656 slow microphysics regimes in a minimalist model of cloudy Rayleigh-Bénard convection,
657 *Physical Review Research*, 5, 043018, 10.1103/PhysRevResearch.5.043018, 2023.

658 Thomas, S., Yang, F., Ovchinnikov, M., Cantrell, W., and Shaw, R. A.: Scaling of turbulence and
659 microphysics in a convection–cloud chamber of varying height, *Journal of Advances in Modeling*
660 *Earth Systems*, 15, e2022MS003304, 2023.

661 Wood: Stratocumulus Clouds, *Monthly Weather Review*, 140, 2373-2423, 10.1175/mwr-d-11-
662 00121.1, 2012.

663 Yau, M. K., and Rogers, R. R.: *A short course in cloud physics*, Elsevier, 1996.

664 Zhu, Z., Kollias, P., Luke, E., and Yang, F.: New insights on the prevalence of drizzle in marine
665 stratocumulus clouds based on a machine learning algorithm applied to radar Doppler spectra,
666 Atmospheric Chemistry and Physics, 22, 7405-7416, 2022.
667